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**TEORÍA CUÁNTICA DE CAMPOS
RELATIVISTAS. FORMALIZACIÓN
TEÓRICA**

**QUANTUM THEORY OF RELATIVISTIC FIELDS.
THEORETICAL FORMALIZATION**

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Teoría Cuántica de Campos Relativistas. Formalización Teórica

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RESUMEN

El propósito del presente manuscrito, es sentar las bases de la teoría cuántica de campos relativistas, deducida referencialmente en artículos iniciales y condensada en trabajos más recientes. Llámese también teoría cuántica en espacios curvos. Por tanto, para todos los efectos constantes en este artículo y en los trabajos relacionados que le preceden a éste, entiéndase por teoría cuántica en espacios o campos curvos, a la teoría cuántica de campos o espacios relativistas y viceversa, y lo propio a sus equivalentes. Esta teoría, en sentido estricto, surge como iniciativa de conciliación entre la relatividad general y la mecánica cuántica, a escala subatómica. El desarrollo matemático vinculante a la teoría en proposición, ha sido ampliamente desarrollado en trabajos previos, sin embargo, en este manuscrito, me ocuparé única y exclusivamente de los fundamentos teóricos, sin descuidar el lenguaje matemático.

Palabras clave: Campos de gauge, supersimetrías, agujeros negros cuánticos, gravedad cuántica, superpembranas

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Quantum Theory of Relativistic Fields. Theoretical Formalization

ABSTRACT

The purpose of this manuscript is to lay the foundations of the quantum theory of relativistic fields, deduced referentially in initial articles and condensed in more recent works. Also called quantum theory in curved spaces. Therefore, for all the constant effects in this article and in the related works that precede it, quantum theory in curved spaces or fields is understood to mean the quantum theory of relativistic fields or spaces and vice versa, and the same to their equivalents. This theory, in the strict sense, arises as an initiative to reconcile general relativity and quantum mechanics, at the subatomic scale. The mathematical development binding to the theory in proposition has been extensively developed in previous works, however, in this manuscript, I will deal solely and exclusively with the theoretical foundations, without neglecting mathematical language.

Keywords: Gauge fields, supersymmetries, quantum black holes, quantum gravity, supermembranes

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INTRODUCCIÓN

Es bien sabido y hasta la saciedad, que la relatividad general y la mecánica cuántica, son incompatibles, esencialmente, a propósito de lo que describen y miden. La gravedad esencialmente, al igual que los objetos interestelares, se comportan de manera distinta en relación a las interacciones que se suscitan en el espacio cuántico de que se trate, esto es, entre partículas subatómicas. Mientras que, a nivel cosmológico, la métrica es puramente determinista, pues, a escala microscópica o cuántica, poco o nada se tiene de certeza, verbigracia, el principio de indeterminación de Heisenberg, que postula que las variables dinámicas de una partícula, esto es, momento angular, momento lineal, posición, etc, son perturbativas, por lo que, no obedecen a la mecánica clásica.

Sin embargo, luego de extensos trabajos puramente matemáticos aunque marginalmente teóricos, y a través de un paralelismo parcial con la teoría de Witten, he deducido una teoría cuántica de campos, en la que, extrapolamos las ecuaciones einstenianas de campo a espacios cuánticos indiscriminados, diseñando supersimetrías que describen el comportamiento de las partículas subatómicas, especialmente en tratándose de aquellas, cuya masa es superlativa o superior, según sea el caso, o en su defecto, de aquellas que se aproximan, igual o superan la velocidad de la luz, éstas últimas, tengan o no masa. El núcleo medular de esta teoría, radica principalmente en la gravedad, entendida, no como fuerza fundamental, sino como una distorsión del espacio – tiempo cuántico provocada por una partícula cosmológica o en su defecto, por la interacción de las propias partículas subatómicas en un espacio cuántico indiscriminado. En el apartado de resultados y discusión, se explicarán cada uno de los postulados que contiene la teoría cuántica de campos relativistas, entendida como una alternativa a la teoría del todo.

RESULTADOS Y DISCUSIÓN.

En este punto, paso a explicar el modelo:

En términos generales, los campos cuánticos relativistas o curvos, comportan la distorsión tensorial del espacio – tiempo a escala subatómica. En este punto, corresponde explicar, las posibles reglas por las que, se despliega este fenómeno:



1. Por la existencia de una partícula cosmológica. Por partícula cosmológica, entendemos en sentido estricto, al gravitón, esto es, una partícula fundamental, que al contrario de lo teorizado en 1930, tratase de una partícula de naturaleza bosónica o posiblemente tratase de una partícula compañera en relación a otra partícula – gauge de origen fermiónico, con una masa superior a los $1,6 \times 10^{-69}$ kg, cuyo campo permea un espacio cuántico indiscriminado, al igual que el mecanismo bosónico de Higgs, cuya diferencia, radica en que esta partícula, a diferencia del bosón de Higgs, que transfiere o dota de masa a las demás partículas subatómicas a través del contacto con el campo cuántico de Higgs, pues ésta, es decir, el gravitón, transfiere o dota de gravedad a las partículas que gozan de masa y energía cinética o potencial suficientes para entrar en contacto con el campo cuántico del gravitón (el cual es curvo), lo que provoca que éstas, llamadas también superpartículas, deformen el espacio – tiempo en el que se interrelacionan, afectando la interacción cuántica con las demás partículas fermiónicas o bosónicas, según sea el caso, que ineludiblemente carezcan de masa o tengan menos masa que la superpartícula perturbativa. Las superpartículas, cabe aclarar, pueden ser de origen bosónico o fermiónico, según sea el caso.

Por tanto, la conexión entre la geometría del espacio – tiempo y la materia, está dada por la métrica $\int_{\mathcal{T}_{\mu\nu}} \eta_{\mu\nu} + \hbar_{\mu\nu} \equiv g_{\mu\nu}$, en la que, $\hbar_{\mu\nu}$ es la materia, es decir, la superpartícula, $\eta_{\mu\nu}$, el campo cuántico del gravitón, $g_{\mu\nu}$, la gravedad soportada y transferida por el gravitón a través de su campo cuántico, el cual permea el espacio y finalmente, la constante $\mathcal{T}_{\mu\nu}$ que es el tensor stress–energy. Al encontrarnos así, en un campo cuántico curvo o relativista, es decir, perturbativo – armónico, la regla de oro de Fermi es aplicable dada la transición de estados por la transferencia de permeo, lo que, en hamiltoniano (estado inicial y estados finales), se dirá:

$$\mathcal{H}_0 u_k(x) = E_k u_k(x) = \hbar \omega_k u_k(x)$$

Aquí los valores de energía se miden a través de la frecuencia $\omega_k = E_k/\hbar$. Siendo así, tenemos:

$$\psi(x, t) = \sum_k c_k(0) e^{-i\omega_k t} u_k(x)$$

Hasta lograr el estado de equilibrio:

$$\begin{aligned} \psi(x, 0) &= u_i(x) \\ \mathcal{H} &= \mathcal{H}_0 + \hat{V}(x) \end{aligned}$$

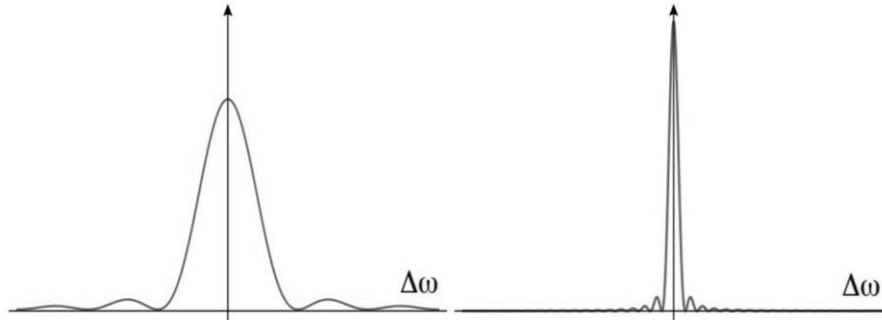


$$\begin{aligned}
u_i(x) &= \sum_h d_h(0)v_h \rightarrow \psi'(x, t) = \sum_h d_h(0)e^{-iE_h^0 t/\hbar}v_h(x). \\
\psi'(x, t) &= \sum_k c_k(t)e^{-i\omega_k t}u_k(x) \\
i\hbar \frac{\partial \psi'}{\partial t} &= \mathcal{H}_0 \psi' + \hat{V} \psi' \\
i\hbar \sum_k [\dot{c}_k(t)e^{-i\omega_k t}u_k(x) - i\omega_k c_k(t)e^{-i\omega_k t}u_k(x)] &= \sum_k c_k(t)e^{-i\omega_k t}(\mathcal{H}_0 u_k(x) + \hat{V}[u_k(x)]) \\
\sum_k [i\hbar \dot{c}_k(t)e^{-i\omega_k t}u_k(x) + \hbar\omega_k c_k(t)e^{-i\omega_k t}u_k(x)] &= \sum_k [c_k(t)e^{-i\omega_k t}\hbar\omega_k u_k(x) + c_k(t)e^{-i\omega_k t}\hat{V}[u_k(x)]] \\
\sum_k i\hbar \dot{c}_k(t)e^{-i\omega_k t}u_k(x) &= \sum_k c_k(t)e^{-i\omega_k t}\hat{V}[u_k(x)] \\
\sum_k i\hbar \dot{c}_k(t)e^{-i\omega_k t} \int_{-\infty}^{\infty} u_h^*(x)u_k(x)dx &= \sum_k c_k(t)e^{-i\omega_k t} \int_{-\infty}^{\infty} u_h^*(x)\hat{V}[u_k(x)]dx \\
\sum_k i\hbar \dot{c}_k(t)e^{-i\omega_k t} \int_{-\infty}^{\infty} u_h^*(x)u_k(x)dx &= i\hbar \dot{c}_h(t)e^{-i\omega_h t} \\
V_{hk} &= \int_{-\infty}^{\infty} u_h^*(x)\hat{V}[u_k(x)]dx \\
\dot{c}_h(t) &= -\frac{i}{\hbar} \sum_k c_k(t)e^{i(\omega_h - \omega_k)t}V_{hk} \\
c_h(t) &= -\frac{i}{\hbar} \sum_k \int_0^t c_k(t')e^{i(\omega_h - \omega_k)t'}V_{hk}dt' + c_h(0) \\
c_h(t) &= -\frac{i}{\hbar} \sum_k c_k(0) \int_0^t e^{i(\omega_h - \omega_k)t'}V_{hk}dt' + c_h(0) \\
c_h(t) &= -\frac{i}{\hbar} \int_0^t e^{i(\omega_h - \omega_i)t'}V_{hi}dt' \\
c_h(t) &= -\frac{i}{\hbar} V_{hi} \int_0^t e^{i\Delta\omega_h t'} dt' = -\frac{V_{hi}}{\hbar\Delta\omega_h} (1 - e^{i\Delta\omega_h t}) \\
P(i \rightarrow h) &= \frac{4|V_{hi}|^2}{\hbar^2 \Delta\omega_h^2} \sin^2\left(\frac{\Delta\omega_h t}{2}\right) \\
P(i \rightarrow h) &= \frac{2\pi|V_{hi}|^2 t}{\hbar^2} \delta(\Delta\omega_h)
\end{aligned}$$

Y la tasa de transición por dispersión energética, se mide:

$$\begin{aligned}
W_{ih} &= \frac{dP(i \rightarrow h)}{dt} \\
W_{ih} &= \frac{2\pi}{\hbar^2} |V_{hi}|^2 \delta(\Delta\omega_h) \\
W_{ih} &= \frac{2\pi}{\hbar} |V_{hi}|^2 \rho(E_h) \Big|_{E_h=E_i}
\end{aligned}$$





En tanto que, en relación al tiempo deformado, la tasa de transición está dada por:

$$\begin{aligned}
 V(t) &= V\Theta(t - t_0) = \begin{cases} 0 & t < t_0 \\ V & t \geq t_0 \end{cases} \\
 b_k &= -\frac{i}{\hbar} \int_{t_0}^t d\tau e^{i\omega_{kl}(\tau-t_0)} V_{k\ell}(\tau) \\
 b_k &= -\frac{i}{\hbar} V_{k\ell} \int_0^t d\tau e^{i\omega_{k\ell}\tau} \\
 &= -\frac{V_{k\ell}}{E_k - E_\ell} [\exp(i\omega_{kl}t) - 1] \\
 &= -\frac{2iV_{k\ell}e^{i\omega_{kl}t/2}}{E_k - E_\ell} \sin(\omega_{kl}t/2) \\
 e^{i\theta} - 1 &= 2ie^{i\theta/2} \sin(\theta/2). \\
 P_k &= |b_k|^2 \\
 &= \frac{4|V_{k\ell}|^2}{|E_k - E_\ell|^2} \sin^2\left(\frac{\omega_{kl}t}{2}\right) \\
 \Delta &= (E_k - E_\ell)/2 \\
 P_k &= \frac{V^2}{\Delta^2} \sin^2(\Delta t/\hbar) \\
 P_k &= \frac{V^2}{V^2 + \Delta^2} \sin^2\left(\sqrt{\Delta^2 + V^2}t/\hbar\right) \\
 P_k &= \frac{V^2 t^2}{\hbar^2} \text{sinc}^2(\Delta t/2\hbar) \\
 \text{sinc}(x) &= \frac{\sin(x)}{x} \\
 \lim_{x \rightarrow 0} \text{sinc}(x) &= 1, \\
 \lim_{\Delta \rightarrow 0} P_k &= \frac{V^2 t^2}{\hbar^2} \\
 \Delta E \cdot \Delta t &\geq 2\pi\hbar \\
 \Delta p \cdot \Delta x &\geq 2\pi\hbar. \\
 \lim_{t \rightarrow \infty} \frac{\sin^2(ax/2)}{ax^2} &= \frac{\pi}{2} \delta(x) \\
 \lim_{t \rightarrow \infty} P_k(t) &= \frac{2\pi}{\hbar} |V_{k\ell}|^2 \delta(E_k - E_\ell) t \\
 w_k(t) &= \frac{\partial P_k(t)}{\partial t} = \frac{2\pi |V_{k\ell}|^2}{\hbar} \delta(E_k - E_\ell) \\
 \frac{dP}{dt} &= -wP. \\
 V(t) &= V \cos \omega t \Theta(t)
 \end{aligned}$$

$$\begin{aligned}
V_{k\ell}(t) &= V_{k\ell} \cos \omega t \\
&= \frac{V_{k\ell}}{2} [e^{-i\omega t} + e^{i\omega t}] \\
b_k &= \frac{-i}{\hbar} \int_{t_0}^t d\tau V_{k\ell}(\tau) e^{i\omega_{kl}\tau} \\
&= \frac{-iV_{k\ell}}{2\hbar} \int_0^t d\tau [e^{i(\omega_{kl}-\omega)\tau} + e^{i(\omega_{kl}+\omega)\tau}] \\
&= \frac{-iV_{k\ell}}{2\hbar} \left[\frac{e^{i(\omega_{kl}-\omega)t} - 1}{\omega_{kl} - \omega} + \frac{e^{i(\omega_{kl}+\omega)t} - 1}{\omega_{kl} + \omega} \right] \\
&\quad e^{\theta} - 1 = 2ie^{i\theta/2} \sin(\theta/2) \\
b_k &= \frac{V_{k\ell}}{\hbar} \left[\underbrace{\frac{e^{i(\omega_{kl}-\omega)t/2} \sin[(\omega_{kl} - \omega)t/2]}{\omega_{kl} - \omega}}_{\text{estado inicial}} + \underbrace{\frac{e^{i(\omega_{kl}+\omega)t/2} \sin[(\omega_{kl} + \omega)t/2]}{\omega_{kl} + \omega}}_{\text{estado final}} \right] \\
P_{k\ell} = |b_k|^2 &= \frac{|V_{kl}|^2}{\hbar^2(\omega_{kt} - \omega)^2} \sin^2 \left[\frac{1}{2}(\omega_{kl} - \omega)t \right] \\
P_{k\ell} = |b_k|^2 &= \frac{|V_{kl}|^2}{\hbar^2(\omega_{kt} - \omega)^2 + |V_{kl}|^2} \sin^2 \left[\frac{1}{2\hbar} \sqrt{|V_{kl}|^2 + (\omega_{kl} - \omega)^2} t \right] \\
w_{k\ell} &= \frac{\pi}{2\hbar^2} |V_{k\ell}|^2 [\delta(\omega_{kl} - \omega) + \delta(\omega_{kl} + \omega)] \\
\lim_{\Delta\omega \rightarrow 0} P_k(t) &= \frac{|V_{k\ell}|^2}{4\hbar^2} t^2 \\
t &\ll \frac{2\hbar}{V_{k\ell}} \\
t &> \frac{1}{\omega} \approx \frac{1}{\omega_{k\ell}} \\
V_{k\ell} &\ll \hbar\omega_{k\ell}
\end{aligned}$$

Calculada la transición de estados, a propósito de la interacción entre el campo gravitónico y las partículas elementales susceptibles o repercutibles, que entran en contacto con dicho campo y por ende, engendran gravedad deformando la geometría espacio – tiempo en la que interactúan (superpartículas), procedemos en este punto, a proponer la ecuación maestra de un campo cuántico gravitónico o del gravitón:

En un campo bosónico, el lagrangiano será:



$$\begin{aligned}
\mathcal{L}_{SM\text{gravitón-campo bosónico}} = & -\frac{1}{2}\partial_\nu g_\mu^\alpha \partial_\nu g_\mu^\alpha - g_s f^{abc} \partial_\mu g_\nu^a g_\mu^b g_\nu^c - \frac{1}{4} g_s^2 f^{abc} f^{ade} g_\mu^b g_\nu^c g_\mu^d g_\nu^e + \frac{1}{2} i g_s^2 (\bar{q}_i^g \gamma^\mu q_j^g) g_\mu^a + \\
& G^a \partial^2 G^a + g_s f^{abc} \partial_\mu G^a G^b g_\mu^c - \partial_\nu W_\mu^+ \partial_\nu W_\mu^- - M^2 W_\mu^+ W_\mu^- - \frac{1}{2} \partial_\nu Z_\mu^0 \partial_\nu Z_\mu^0 - \frac{1}{2c_w^2} M^2 Z_\mu^0 Z_\mu^0 - \\
& \frac{1}{2} \partial_\mu A_\nu \partial_\mu A_\nu - \frac{1}{2} \partial_\mu H \partial_\mu H - \frac{1}{2} m_h^2 H^2 - \partial_\mu \phi^+ \partial_\mu \phi^- - M^2 \phi^+ \phi^- - \frac{1}{2} \partial_\mu \phi^0 \partial_\mu \phi^0 - \\
& \frac{1}{2c^2} M \phi^0 \phi^0 - \beta_h \left[\frac{2M^2}{g^2} + \frac{2M}{g} H + \frac{1}{2} (H^2 + \phi^0 \phi^0 + 2\phi^+ \phi^-) \right] + \frac{2M^4}{g^2} \alpha_h - i g c_w [\partial_\nu Z_\mu^0 (W_\mu^+ W_\nu^- - \\
& W_\nu^+ W_\mu^-) - Z_\nu^0 (W_\mu^+ \partial_\nu W_\mu^- - W_\mu^- \partial_\nu W_\mu^+) + Z_\mu^0 (W_\nu^+ \partial_\nu W_\mu^- - W_\nu^- \partial_\nu W_\mu^+)] - i g s_w [\partial_\nu A_\mu (W_\mu^+ W_\nu^- - \\
& W_\nu^+ W_\mu^-) - A_\nu (W_\mu^+ \partial_\nu W_\mu^- - W_\mu^- \partial_\nu W_\mu^+) + A_\mu (W_\nu^+ \partial_\nu W_\mu^- - W_\nu^- \partial_\nu W_\mu^+)] - \frac{1}{2} g^2 W_\mu^+ W_\mu^- W_\nu^+ W_\nu^- + \\
& \frac{1}{2} g^2 W_\mu^+ W_\nu^- W_\mu^+ W_\nu^- + g^2 c_w^2 (Z_\mu^0 W_\mu^+ Z_\nu^0 W_\nu^- - Z_\mu^0 Z_\nu^0 W_\mu^+ W_\nu^-) + g^2 s_w^2 (A_\mu W_\mu^+ A_\nu W_\nu^- - \\
& A_\mu A_\nu W_\mu^+ W_\nu^-) + g^2 s_w c_w [A_\mu Z_\nu^0 (W_\mu^+ W_\nu^- - W_\nu^+ W_\mu^-) - 2A_\mu Z_\mu^0 W_\nu^+ W_\nu^-] - g \alpha [H^3 + \\
& H \phi^0 \phi^0 + 2H \phi^+ \phi^-] - \frac{1}{8} g^2 \alpha_h [H^4 + (\phi^0)^4 + 4(\phi^+ \phi^-)^2 + 4(\phi^0)^2 \phi^+ \phi^- + 4H^2 \phi^+ \phi^- + \\
& 2(\phi^0)^2 H^2] - g M W_\mu^+ W_\mu^- H - \frac{1}{2} g \frac{M}{c_w^2} Z_\mu^0 Z_\mu^0 H - \frac{1}{2} i g [W_\mu^+ (\phi^0 \partial_\mu \phi^- - \phi^- \partial_\mu \phi^0) - W_\mu^- (\phi^0 \partial_\mu \phi^+ - \\
& \phi^+ \partial_\mu \phi^0)] + \frac{1}{2} g [W_\mu^+ (H \partial_\mu \phi^- - \phi^- \partial_\mu H) - W_\mu^- (H \partial_\mu \phi^+ - \phi^+ \partial_\mu H)] + \frac{1}{2} g \frac{1}{c_w} Z_\mu^0 (H \partial_\mu \phi^0 - \\
& \phi^0 \partial_\mu H) - i g \frac{g_e^2}{c_w} M Z_\mu^0 (W_\mu^+ \phi^- - W_\mu^- \phi^+) + i g s_w M A_\mu (W_\mu^+ \phi^- - W_\mu^- \phi^+) - i g \frac{1 - 2c_w^2}{2c_w} Z_\mu^0 (\phi^+ \partial_\mu \phi^- - \\
& \phi^- \partial_\mu \phi^+) + i g s_w A_\mu (\phi^+ \partial_\mu \phi^- - \phi^- \partial_\mu \phi^+) - \frac{1}{4} g^2 W_\mu^+ W_\mu^- [H^2 + (\phi^0)^2 + 2\phi^+ \phi^-] - \\
& \frac{1}{4} g^2 \frac{1}{c_s^2} Z_\mu^0 Z_\mu^0 [H^2 + (\phi^0)^2 + 2(2s_w^2 - 1)^2 \phi^+ \phi^-] - \frac{1}{2} g^2 \frac{s_c^2}{c_w} Z_\mu^0 \phi^0 (W_\mu^+ \phi^- + W_\mu^- \phi^+) - \\
& \frac{1}{2} i g^2 \frac{s_m^2}{c_w} Z_\mu^0 H (W_\mu^+ \phi^- - W_\mu^- \phi^+) + \frac{1}{2} g^2 s_w A_\mu \phi^0 (W_\mu^+ \phi^- + W_\mu^- \phi^+) + \frac{1}{2} i g^2 s_w A_\mu H (W_\mu^+ \phi^- - \\
& W_\mu^- \phi^+) - g^2 \frac{A_\mu}{c_\omega} (2c_w^2 - 1) Z_\mu^0 A_\mu \phi^+ \phi^- - g^1 s_\omega^2 A_\mu A_\mu \phi^+ \phi^- - \bar{e}^\lambda (\gamma \partial + m_e^\lambda) e^\lambda - \\
& \bar{\nu}^\lambda \gamma \partial \nu^\lambda - \bar{u}_j^\lambda (\gamma \partial + m_u^\lambda) u_j^\lambda - \bar{d}_j^\lambda (\gamma \partial + m_d^\lambda) d_j^\lambda + i g s_w A_\mu \left[-(\bar{e}^\lambda \gamma^\mu e^\lambda) + \frac{2}{3} (\bar{u}_j^\lambda \gamma^\mu u_j^\lambda) - \right. \\
& \left. \frac{1}{3} (d_j^\lambda \gamma^\mu d_j^\lambda) \right] + \frac{i g}{4c_w} Z_\mu^0 \left[(\nu^\lambda \gamma^\mu (1 + \gamma^5) \nu^\lambda) + (e^\lambda \gamma^\mu (4s_w^2 - 1 - \gamma^5) e^\lambda) + (\bar{u}_j^\lambda \gamma^\mu \left(\frac{4}{3} s_w^2 - \right. \right. \right. \\
& \left. \left. 1 - \gamma^5 \right) u_j^\lambda) + (d_j^\lambda \gamma^\mu \left(1 - \frac{8}{3} s_w^2 - \gamma^5 \right) d_j^\lambda) \right] + \frac{i g}{2\sqrt{2}} W_\mu^+ [(\nu^\lambda \gamma^\mu (1 + \gamma^5) e^\lambda) + (\bar{u}_j^\lambda \gamma^\mu (1 + \\
& \gamma^5) C_{\lambda k} d_j^k)] + \frac{i g}{2\sqrt{2}} W_\mu^- [(\bar{e}^\lambda \gamma^\mu (1 + \gamma^5) \nu^\lambda) + (d_j^k C_{\lambda k}^\dagger \gamma^\mu (1 + \gamma^5) u_j^\lambda)] + \frac{i g}{2\sqrt{2}} \frac{m^\lambda}{M} [-\phi^+ (\bar{\nu}^\lambda (1 - \\
& \gamma^5) e^\lambda) + \phi^- (\bar{e}^\lambda (1 + \gamma^5) \nu^\lambda)] - \frac{g m_\lambda^\lambda}{2M} [H (\bar{e}^\lambda e^\lambda) + i \phi^0 (\bar{e}^\lambda \gamma^5 e^\lambda)] + \frac{i g}{2M\sqrt{2}} \phi^+ [-m_d^\kappa (\bar{u}_j^\lambda C_{\lambda k} (1 - \\
& \gamma^5) d_j^k) + m_u^\lambda (\bar{u}_j^\lambda C_{\lambda k} (1 + \gamma^5) d_j^k) + \frac{i g}{2M\sqrt{2}} \phi^- [m_d^\lambda (\bar{d}_j^\lambda C_{\lambda k}^\dagger (1 + \gamma^5) u_j^k) - m_u^\kappa (\bar{d}_j^\lambda C_{\lambda k}^\dagger (1 - \\
& \gamma^5) u_j^k) - \frac{g m_n^\lambda}{2M} H (\bar{u}_j^\lambda u_j^\lambda) - \frac{g m_\lambda^\lambda}{2M} H (d_j^\lambda d_j^\lambda) + \frac{i g m_j^\lambda}{2M} \phi^0 (\bar{u}_j^\lambda \gamma^5 u_j^\lambda) - \frac{i g m_\lambda^\lambda}{2M} \phi^0 (d_j^\lambda \gamma^5 d_j^\lambda) + \\
& \bar{X}^+ (\partial^2 - M^2) X^+ + \bar{X}^- (\partial^2 - M^2) X^- + \bar{X}^0 \left(\partial^2 - \frac{M^2}{c_0^2} \right) X^0 + \bar{Y} \partial^2 Y + i g c_w W_\mu^+ (\partial_\mu \bar{X}^0 X^- - \\
& \partial_\mu \bar{X}^+ X^0) + i g s_w W_\mu^+ (\partial_\mu \bar{Y} X^- - \partial_\mu \bar{X}^+ Y) + i g c_w W_\mu^- (\partial_\mu \bar{X}^- X^0 - \partial_\mu \bar{X}^0 X^+) + \\
& i g s_w W_\mu^- (\partial_\mu \bar{X}^- Y - \partial_\mu \bar{Y} X^+) + i g c_w Z_\mu^0 (\partial_\mu \bar{X} + X^+ - \partial_\mu \bar{X} - X^-) + i g s_w A_\mu (\partial_\mu \bar{X}^+ X^+ - \\
& \partial_\mu \bar{X}^- X^-) - \frac{1}{2} g M \left[\bar{X}^+ X^+ H + \bar{X}^- X^- H + \frac{1}{\epsilon_0^2} \bar{X}^0 X^0 H \right] + \frac{1 - 2c_m^2}{2\omega_w} i g M [\bar{X}^+ X^0 \phi^+ - \\
& \bar{X}^- X^0 \phi^-] + \frac{1}{2c_z} i g M [\bar{X}^0 X^+ \phi^+ - \bar{X}^0 X^+ \phi^-] + i g M s_w [\bar{X}^0 X^- \phi^+ - \bar{X}^0 X^+ \phi^-] + \\
& \frac{1}{2} i g M [\bar{X}^+ X^+ \phi^0 - \bar{X}^- X^- \phi^0]
\end{aligned}$$

En un campo fermiónico, el lagrangiano será:



$$\begin{aligned}
\mathcal{L}_{SM}^{\text{gravitón-campo fermiónico}} = & -\frac{1}{2}\partial_\nu g_\mu^\alpha \partial_\nu g_\mu^\alpha - g_s f^{abc} \partial_\mu g_\nu^a g_\mu^b g_\nu^c - \frac{1}{4}g_s^2 f^{abc} f^{ade} g_\mu^b g_\nu^c g_\mu^d g_\nu^e + \\
& \frac{1}{2}ig_s^2(\bar{q}_i^\sigma \gamma^\mu q_j^\sigma)g_\mu^\alpha + \bar{G}^a \partial^2 G^a + g_s f^{abc} \partial_\mu \bar{G}^a G^b g_\mu^c - \partial_\nu W_\mu^+ \partial_\nu W_\mu^- - \\
& M^2 W_\mu^+ W_\mu^- - \frac{1}{2}\partial_\nu Z_\mu^0 \partial_\nu Z_\mu^0 - \frac{1}{2c_w^2}M^2 Z_\mu^0 Z_\mu^0 - \frac{1}{2}\partial_\mu A_\nu \partial_\mu A_\nu - \frac{1}{2}\partial_\mu H \partial_\mu H - \\
& \frac{1}{2}m_h^2 H^2 - \partial_\mu \phi^+ \partial_\mu \phi^- - M^2 \phi^+ \phi^- - \frac{1}{2}\partial_\mu \phi^0 \partial_\mu \phi^0 - \frac{1}{2c_w^2}M\phi^0 \phi^0 - \beta_h \left[\frac{2M^2}{g^2} + \right. \\
& \left. \frac{2M}{g}H + \frac{1}{2}(H^2 + \phi^0 \phi^0 + 2\phi^+ \phi^-) \right] + \frac{2M^4}{g^2} \alpha_h - ig_{c_w}[\partial_\nu Z_\mu^0 (W_\mu^+ W_\nu^- - \\
& W_\nu^+ W_\mu^-) - Z_\nu^0 (W_\mu^+ \partial_\nu W_\mu^- - W_\mu^- \partial_\nu W_\mu^+) + Z_\mu^0 (W_\nu^+ \partial_\nu W_\mu^- - \\
& W_\nu^- \partial_\nu W_\mu^+)] - ig_{s_w}[\partial_\nu A_\mu (W_\mu^+ W_\nu^- - W_\nu^+ W_\mu^-) - A_\nu (W_\mu^+ \partial_\nu W_\mu^- - \\
& W_\mu^- \partial_\nu W_\mu^+) + A_\mu (W_\nu^+ \partial_\nu W_\mu^- - W_\nu^- \partial_\nu W_\mu^+)] - \frac{1}{2}g^2 W_\mu^+ W_\mu^- W_\nu^+ W_\nu^- + \\
& \frac{1}{2}g^2 W_\mu^+ W_\nu^- W_\mu^+ W_\nu^- + g^2 c_w^2 (Z_\mu^0 W_\mu^+ Z_\nu^0 W_\nu^- - Z_\mu^0 Z_\nu^0 W_\mu^+ W_\nu^-) + \\
& g^2 s_w^2 (A_\mu W_\mu^+ A_\nu W_\nu^- - A_\mu A_\nu W_\mu^+ W_\nu^-) + g^2 s_w c_w [A_\mu Z_\nu^0 (W_\mu^+ W_\nu^- - \\
& W_\nu^+ W_\mu^-) - 2A_\mu Z_\mu^0 W_\nu^+ W_\nu^-] - g\alpha [H^3 + H\phi^0 \phi^0 + 2H\phi^+ \phi^-] - \\
& \frac{1}{8}g^2 \alpha_h [H^4 + (\phi^0)^4 + 4(\phi^+ \phi^-)^2 + 4(\phi^0)^2 \phi^+ \phi^- + 4H^2 \phi^+ \phi^- + 2(\phi^0)^2 H^2] - \\
& gMW_\mu^+ W_\mu^- H - \frac{1}{2}g\frac{M}{c_w^2}Z_\mu^0 Z_\mu^0 H - \frac{1}{2}ig[W_\mu^+ (\phi^0 \partial_\mu \phi^- - \phi^- \partial_\mu \phi^0) - \\
& W_\mu^- (\phi^0 \partial_\mu \phi^+ - \phi^+ \partial_\mu \phi^0)] + \frac{1}{2}g[W_\mu^+ (H\partial_\mu \phi^- - \phi^- \partial_\mu H) - W_\mu^- (H\partial_\mu \phi^+ - \\
& \phi^+ \partial_\mu H)] + \frac{1}{2}g\frac{1}{c_w}Z_\mu^0 (H\partial_\mu \phi^0 - \phi^0 \partial_\mu H) - ig\frac{s_w^2}{c_w}MZ_\mu^0 (W_\mu^+ \phi^- - W_\mu^- \phi^+) + \\
& ig_{s_w}MA_\mu (W_\mu^+ \phi^- - W_\mu^- \phi^+) - ig\frac{1-2c_w^2}{2c_w}Z_\mu^0 (\phi^+ \partial_\mu \phi^- - \phi^- \partial_\mu \phi^+) + \\
& ig_{s_w}A_\mu (\phi^+ \partial_\mu \phi^- - \phi^- \partial_\mu \phi^+) - \frac{1}{4}g^2 W_\mu^+ W_\mu^- [H^2 + (\phi^0)^2 + 2\phi^+ \phi^-] - \\
& \frac{1}{4}g^2 \frac{1}{c_w^2}Z_\mu^0 Z_\mu^0 [H^2 + (\phi^0)^2 + 2(2s_w^2 - 1)^2 \phi^+ \phi^-] - \frac{1}{2}g^2 \frac{s_w^2}{c_w}Z_\mu^0 \phi^0 (W_\mu^+ \phi^- + \\
& W_\mu^- \phi^+) - \frac{1}{2}ig^2 \frac{s_w^2}{c_w}Z_\mu^0 H (W_\mu^+ \phi^- - W_\mu^- \phi^+) + \frac{1}{2}g^2 s_w A_\mu \phi^0 (W_\mu^+ \phi^- + \\
& W_\mu^- \phi^+) + \frac{1}{2}ig^2 s_w A_\mu H (W_\mu^+ \phi^- - W_\mu^- \phi^+) - g^2 c_w (2c_w^2 - 1)Z_\mu^0 A_\mu \phi^+ \phi^- - \\
& g^1 s_w^2 A_\mu A_\mu \phi^+ \phi^- - \bar{e}^\lambda (\gamma \partial + m_e^2) e^\lambda - \bar{\nu}^\lambda \gamma \partial \nu^\lambda - \bar{u}_j^\lambda (\gamma \partial + m_u^\lambda) u_j^\lambda - \\
& \bar{d}_j^\lambda (\gamma \partial + m_d^\lambda) d_j^\lambda + ig_{s_w} A_\mu \left[-(\bar{e}^\lambda \gamma^\mu e^\lambda) + \frac{2}{3}(\bar{u}_j^\lambda \gamma^\mu u_j^\lambda) - \frac{1}{3}(\bar{d}_j^\lambda \gamma^\mu d_j^\lambda) \right] + \\
& \frac{ig}{4c_w} Z_\mu^0 \left[(\bar{\nu}^\lambda \gamma^\mu (1 + \gamma^5) \nu^\lambda) + (\bar{e}^\lambda \gamma^\mu (4s_w^2 - 1 - \gamma^5) e^\lambda) + (\bar{u}_j^\lambda \gamma^\mu \left(\frac{4}{3}s_w^2 - \right. \right. \right. \\
& \left. \left. \left. 1 - \gamma^5 \right) u_j^\lambda) + (\bar{d}_j^\lambda \gamma^\mu \left(1 - \frac{8}{3}s_w^2 - \gamma^5 \right) d_j^\lambda) \right] + \frac{ig}{2\sqrt{2}} W_\mu^+ [(\bar{\nu}^\lambda \gamma^\mu (1 + \gamma^5) \bar{s}^5) + \\
& (\bar{u}_j^\lambda \gamma^\mu (1 + \gamma^5) C_{\lambda k} d_j^k)] + \frac{ig}{2\sqrt{2}} W_\mu^- [(\bar{e}^\lambda \gamma^\mu (1 + \gamma^5) \nu^\lambda) + (\bar{d}_j^\lambda C_{\lambda k}^\dagger \gamma^\mu (1 + \\
& \gamma^5) u_j^\lambda)] + \frac{ig}{2\sqrt{2}} \frac{m_\lambda^1}{M} [-\phi^+ (\bar{\nu}^\lambda (1 - \gamma^5) e^\lambda) + \phi^- (\bar{e}^\lambda (1 + \gamma^5) \nu^\lambda)] - \\
& \frac{2m_\lambda^1}{2M} [H(\bar{e}^\lambda e^\lambda) + i\phi^0 (\bar{e}^\lambda \gamma^5 e^\lambda)] + \frac{ig}{2M\sqrt{2}} \phi^+ [-m_\lambda^k (\bar{u}_j^\lambda C_{\lambda k} (1 - \gamma^5) d_j^k) + \\
& m_\lambda^k (\bar{u}_j^\lambda C_{\lambda k} (1 + \gamma^5) d_j^k)] + \frac{ig}{2M\sqrt{2}} \phi^- [m_\lambda^k (\bar{d}_j^\lambda C_{\lambda k}^\dagger (1 + \gamma^5) u_j^k) - m_\lambda^k (\bar{d}_j^\lambda C_{\lambda k}^\dagger (1 - \\
& \gamma^5) u_j^k)] - \frac{2m_\lambda^1}{2M} H(\bar{u}_j^\lambda u_j^\lambda) - \frac{2m_\lambda^1}{2M} H(\bar{d}_j^\lambda d_j^\lambda) + \frac{ig}{2M} \phi^0 (\bar{u}_j^\lambda \gamma^5 u_j^\lambda) - \\
& \frac{ig}{2M} \phi^0 (\bar{d}_j^\lambda \gamma^5 d_j^\lambda) + X^+ (\partial^2 - M^2) X^+ + X^- (\partial^2 - M^2) X^- + X^0 (\partial^2 - \\
& \frac{M^2}{c_w^2}) X^0 + \bar{Y} \partial^2 Y + ig_{c_w} W_\mu^+ (\partial_\mu \bar{X}^0 X^- - \partial_\mu \bar{X}^+ X^0) + ig_{s_w} W_\mu^+ (\partial_\mu \bar{Y} X^- - \\
& \partial_\mu \bar{X} + Y) + ig_{c_w} W_\mu^- (\partial_\mu \bar{X} - X^0 - \partial_\mu \bar{X}^0 X^+) + ig_{s_w} W_\mu^- (\partial_\mu \bar{X} - Y - \\
& \partial_\mu \bar{Y} X^+) + ig_{c_w} Z_\mu^0 (\partial_\mu \bar{X}^+ X^+ - \partial_\mu \bar{X}^- X^-) + ig_{s_w} A_\mu (\partial_\mu \bar{X}^+ X^+ - \\
& \partial_\mu \bar{X}^- X^-) - \frac{1}{2}gM \left[\bar{X}^+ X^+ H + \bar{X}^- X^- H + \frac{1}{c_\xi^2} \bar{X}^0 X^0 H \right] + \\
& \frac{1-2c_w^2}{2c_w} igM [\bar{X}^+ X^0 \phi^+ - \bar{X}^- X^0 \phi^-] + \frac{1}{2c_w} igM [\bar{X}^0 X^- \phi^+ - \bar{X}^0 X^+ \phi^-] + \\
& igM s_w [\bar{X}^0 X^- \phi^+ - \bar{X}^0 X^+ \phi^-] + \frac{1}{2} igM [\bar{X} + X^+ \phi^0 - \bar{X}^- X^- \phi^0]
\end{aligned}$$



Cuya solución a las ecuaciones maestras es:

$$\begin{aligned}
 S &= \int d^4x \mathcal{L} \\
 \delta S &= \delta \int d^4x \mathcal{L}(\phi_a(x), \partial_\mu \phi_a(x)) = 0 \\
 \frac{\partial \mathcal{L}}{\partial \phi_a} - \partial_\mu \frac{\partial \mathcal{L}}{\partial (\partial_\mu \phi_a)} &= 0, (a = 1, 2, 3, \dots, n) \\
 \pi_a(x) &= \frac{\partial \mathcal{L}}{\partial (\partial_0 \phi_a)} \\
 [\phi_a(\mathbf{x}, t), \phi_b(\mathbf{y}, t)] &= [\pi_a(\mathbf{x}, t), \pi_b(\mathbf{y}, t)] = 0 \\
 [\phi_a(\mathbf{x}, t), \pi_b(\mathbf{y}, t)] &= i \delta_{ab} \delta^3(\mathbf{x} - \mathbf{y}) \\
 \mathcal{L} &\rightarrow \mathcal{L} + \alpha \partial_\mu J^\mu \\
 \alpha \Delta \mathcal{L} &= \frac{\partial \mathcal{L}}{\partial \phi_a} (\alpha \Delta \phi_a) + \left(\frac{\partial \mathcal{L}}{\partial (\partial_\mu \phi_a)} \right) \partial_\mu (\alpha \Delta \phi_a). \\
 \alpha \partial_\mu \left(\frac{\partial \mathcal{L}}{\partial (\partial_\mu \phi_a)} \Delta \phi_a \right) &- \alpha \partial_\mu \left(\frac{\partial \mathcal{L}}{\partial (\partial_\mu \phi_a)} \right) \Delta \phi_a \\
 &= \alpha \partial_\mu \left(\frac{\partial \mathcal{L}}{\partial (\partial_\mu \phi_a)} \right) \Delta \phi_a \\
 &+ \alpha \frac{\partial \mathcal{L}}{\partial (\partial_\mu \phi_a)} \partial_\mu (\Delta \phi_a) - \alpha \partial_\mu \left(\frac{\partial \mathcal{L}}{\partial (\partial_\mu \phi_a)} \right) \Delta \phi_a \\
 &= \left(\frac{\partial \mathcal{L}}{\partial (\partial_\mu \phi_a)} \right) \partial_\mu (\alpha \Delta \phi_a) \\
 \alpha \Delta \mathcal{L} &= \alpha \left[\frac{\partial \mathcal{L}}{\partial \phi_a} - \partial_\mu \left(\frac{\partial \mathcal{L}}{\partial (\partial_\mu \phi_a)} \right) \right] \Delta \phi_a + \alpha \partial_\mu \left(\frac{\partial \mathcal{L}}{\partial (\partial_\mu \phi_a)} \Delta \phi_a \right) \\
 \alpha \Delta \mathcal{L} &= \alpha \partial_\mu \left(\frac{\partial \mathcal{L}}{\partial (\partial_\mu \phi_a)} \Delta \phi_a \right) \\
 \partial_\mu j^\mu &= 0, \text{ for } j^\mu = \frac{\partial \mathcal{L}}{\partial (\partial_\mu \phi_a)} \Delta \phi_a - J^\mu \\
 \frac{\partial}{\partial t} j^0 + \vec{\nabla} \cdot \vec{j} &= 0 \\
 \int d^3x \left(\frac{\partial}{\partial t} j^0 + \vec{\nabla} \cdot \vec{j} \right) &= 0 \\
 Q(t) &= \int d^3x j^0 \\
 \frac{dQ(t)}{dt} &= \frac{d}{dt} \int d^3x j^0 \\
 \frac{dQ(t)}{dt} &= 0 \\
 Q(t) &= \int d^3x j^0(\mathbf{x}, t) \\
 \psi'(x) &= U_\theta \psi(x) = e^{-i\theta} \psi(x) \\
 \psi'(x) &= e^{-i\theta(x)} \psi(x)
 \end{aligned}$$



$$\begin{aligned}
\mathcal{L} \rightarrow \mathcal{L}' &= \bar{\psi}'(x)(i\gamma^\mu \partial_\mu - m)\psi'(x) \\
&= e^{i\theta(x)}\bar{\psi}(x)(i\gamma^\mu \partial_\mu - m)e^{-i\theta(x)}\psi(x) \\
&= e^{i\theta(x)}\bar{\psi}(x)\gamma^\mu e^{-i\theta(x)}\psi(x)\partial_\mu\theta(x) \\
&\quad + e^{i\theta(x)}\bar{\psi}(x)e^{-i\theta(x)}(i\gamma^\mu \partial_\mu - m)\psi'(x) \\
&= \bar{\psi}(x)(i\gamma^\mu \partial_\mu - m)\psi'(x) + \bar{\psi}(x)\gamma^\mu\psi(x)\partial_\mu\theta(x) \\
&= \mathcal{L} + j^\mu(x)\partial_\mu\theta(x)
\end{aligned}$$

$$\begin{aligned}
\partial_\mu &\rightarrow D_\mu \equiv \partial_\mu - iqA_\mu(x), \\
\mathcal{L} &\rightarrow \mathcal{L} + qj^\mu A_\mu.
\end{aligned}$$

$$\begin{aligned}
\mathcal{L}(\psi, \bar{\psi}, A_\mu) &= \bar{\psi}(i\gamma^\mu D_\mu - m)\psi \\
&= \bar{\psi}(i\gamma^\mu \partial_\mu - m)\psi + q\bar{\psi}\gamma_\mu\psi A_\mu
\end{aligned}$$

$$A_\mu \rightarrow A'_\mu = A_\mu - \frac{1}{q}\partial_\mu\theta(x)$$

$$\mathcal{L} = \bar{\psi}(i\gamma^\mu D_\mu - m)\psi$$

$$(D_\mu\psi(x))' = e^{-i\theta(x)}D_\mu\psi(x).$$

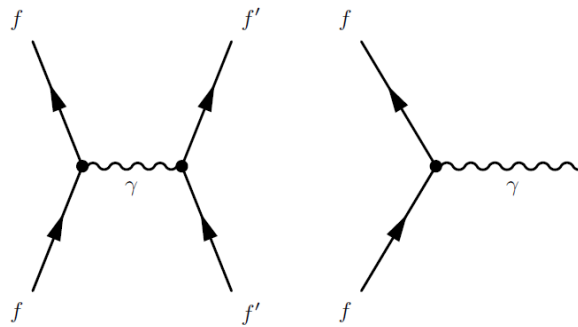
$$\mathcal{L}_\gamma = \frac{m^2}{2}A_\mu A^\mu$$

$$A^\mu A_\mu \rightarrow (A^\mu - \partial^\mu\theta)(A_\mu - \partial_\mu\theta) \neq A^\mu A_\mu$$

$$\mathcal{L} = -\frac{1}{4}(\partial_\nu A_\mu - \partial_\mu A_\nu)(\partial^\nu A^\mu - \partial^\mu A^\nu)$$

$$\mathcal{L} = -\frac{1}{4}F_{\mu\nu}F^{\mu\nu}$$

$$\mathcal{L}_{QED} = \bar{\psi}(x)(i\gamma^\mu D_\mu - m)\psi(x) - \frac{1}{4}F_{\mu\nu}F^{\mu\nu}$$



$$\mathcal{L} = \mathcal{L}_1 + \mathcal{L}_2$$

$$= \bar{\psi}_1(x)(i\gamma^\mu \partial_\mu - m_1)\psi_1(x)$$

$$+ \bar{\psi}_2(x)(i\gamma^\mu \partial_\mu - m_2)\psi_2(x).$$

$$\psi = \begin{pmatrix} \psi_1 \\ \psi_2 \end{pmatrix}$$

$$\bar{\psi} = (\bar{\psi}_1 \quad \bar{\psi}_2)$$

$$M = \begin{pmatrix} m_1 & 0 \\ 0 & m_2 \end{pmatrix}$$

$$\mathcal{L} = \bar{\psi}(i\gamma^\mu \partial_\mu - M)\psi$$

$$\psi \rightarrow \psi' = U\psi$$

$$\bar{\psi} \rightarrow \psi' = \bar{\psi}U^\dagger$$

$$H = a_0I + a_1\tau_1 + a_2\tau_2 + a_3\tau_3 = a_0I + \mathbf{a} \cdot \boldsymbol{\tau}$$

$$U = e^{ia_0}e^{i\mathbf{a} \cdot \boldsymbol{\tau}}$$

$$\psi' = e^{i\mathbf{a} \cdot \boldsymbol{\tau}}\psi$$

$$\psi' = e^{-ig\frac{\boldsymbol{\tau}}{2}\theta(x)}\psi$$

$$\bar{\psi}' = e^{ig\frac{\boldsymbol{\tau}}{2}\theta(x)}\bar{\psi}$$

$$\begin{aligned}
\partial_\mu \psi' &= (\partial_\mu U)\psi + U(\partial_\mu \psi) \\
\mathcal{L}' &= \bar{\psi}(i\gamma^\mu \partial_\mu - M)\psi \\
&= e^{ig\frac{\tau}{2}\theta(x)} \bar{\psi}(i\gamma^\mu \partial_\mu - M)e^{-ig\frac{\tau}{2}\theta(x)} \psi \\
&= \bar{\psi}(i\gamma^\mu \partial_\mu - M)\psi + (\partial_\mu \theta(x)) \bar{\psi} \left(\gamma^\mu g \frac{\tau}{2} \right) \psi \\
&= \mathcal{L} + (\partial_\mu \theta(x)) \bar{\psi} \left(\gamma^\mu g \frac{\tau}{2} \right) \psi.
\end{aligned}$$

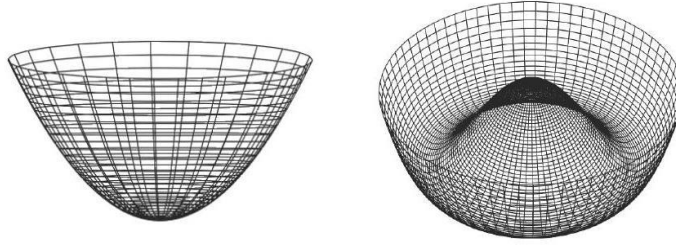
$$\begin{aligned}
\partial_\mu &\rightarrow D_\mu = \partial_\mu - ig\vec{A}_\mu \\
(D_\mu \psi)' &= D'_\mu \psi' = U(D_\mu \psi) \\
D'_\mu \psi' &= (\partial_\mu - ig\vec{A}'_\mu)\psi' \\
&= (\partial_\mu U)\psi + U(\partial_\mu \psi) - igA'_\mu(U\psi), \\
U(D_\mu \psi) &= U(\partial_\mu - ig\vec{A}_\mu)\psi \\
&= U(\partial_\mu \psi) - igU(A_\mu \psi) \\
igA'_\mu(U\psi) &= (\partial_\mu U)\psi + igU(A_\mu \psi) \\
igA'_\mu(UU^{-1}U\psi) &= (\partial_\mu U)U^{-1}U\psi + igU(A_\mu U^{-1}U\psi) \\
igA'_\mu \psi' &= (\partial_\mu U)U^{-1}\psi' + igU(A_\mu U^{-1}\psi') \\
A'_\mu &= UA_\mu U^{-1} - \frac{i}{g}(\partial_\mu U)U^{-1} \\
F_{\mu\nu} &\rightarrow UF_{\mu\nu}U^{-1}
\end{aligned}$$

$$\begin{aligned}
F_{\mu\nu} &= \partial_\mu A_\nu - \partial_\nu A_\mu \\
&\rightarrow \partial_\mu \left[UA_\nu U^{-1} - \frac{i}{g}(\partial_\nu U)U^{-1} \right] - \partial_\nu \left[UA_\mu U^{-1} - \frac{i}{g}(\partial_\mu U)U^{-1} \right] \\
&= \partial_\mu(UA_\nu U^{-1}) - \frac{i}{g}(\partial_\mu \partial_\nu U)U^{-1} - \frac{i}{g}(\partial_\nu U)(\partial_\mu U^{-1}) - \partial_\nu(UA_\mu U^{-1}) + \frac{i}{g}(\partial_\nu \partial_\mu U)U^{-1} + \frac{i}{g}(\partial_\mu U)(\partial_\nu U^{-1}) \\
&= \frac{i}{g}(\partial_\mu U \partial_\nu U^{-1} - \partial_\nu U \partial_\mu U^{-1}) \partial_\mu UA_\nu U^{-1} + U \partial_\mu A_\nu U^{-1} + UA_\nu \partial_\mu U^{-1} - \partial_\nu UA_\mu U^{-1} - U \partial_\nu A_\mu U^{-1} - UA_\mu \partial_\nu U^{-1} \\
&= \frac{i}{g}(\partial_\mu U \partial_\nu U^{-1} - \partial_\nu U \partial_\mu U^{-1}) + \partial_\mu UA_\nu U^{-1} + UA_\nu \partial_\mu U^{-1} - \partial_\nu UA_\mu U^{-1} - UA_\mu \partial_\nu U^{-1} + U(\partial_\mu A_\nu - \partial_\nu A_\mu)U^{-1} \\
&ig(A_\mu A_\nu - A_\nu A_\mu) \\
&\rightarrow ig \left[\left(UA_\mu U^{-1} - \frac{i}{g}(\partial_\mu U)U^{-1} \right) \left(UA_\nu U^{-1} - \frac{i}{g}(\partial_\nu U)U^{-1} \right) - \left(UA_\nu U^{-1} - \frac{i}{g}(\partial_\nu U)U^{-1} \right) \left(UA_\mu U^{-1} - \frac{i}{g}(\partial_\mu U)U^{-1} \right) \right] \\
&= ig \left[UA_\mu A_\nu U^{-1} - \frac{i}{g}(\partial_\mu UA_\nu U^{-1} - UA_\mu \partial_\nu U^{-1}) + \frac{1}{g^2} \partial_\mu U \partial_\nu U^{-1} - UA_\nu A_\mu U^{-1} + \frac{i}{g}(\partial_\nu UA_\mu U^{-1} - UA_\nu \partial_\mu U^{-1}) \right. \\
&\quad \left. - \frac{1}{g^2} \partial_\nu U \partial_\mu U^{-1} \right]
\end{aligned}$$

$$\begin{aligned}
F_{\mu\nu} &= \partial_\mu A_\nu - \partial_\nu A_\mu - ig[A_\mu, A_\nu] \\
\mathcal{L}_{\text{gauge}} &= -\frac{1}{2} \text{Tr}(\vec{F}_{\mu\nu} \cdot \vec{F}^{\mu\nu}) \\
&= -\frac{1}{2} \sum_{i,j=1}^3 \text{Tr} \left(\frac{\tau^i}{2} F_{\mu\nu}^i \frac{\tau^j}{2} F^{j\mu\nu} \right) = -\frac{1}{4} F_{\mu\nu}^i F^{i\mu\nu}
\end{aligned}$$

$$\begin{aligned}
\mathcal{L}_{YM} &= \mathcal{L}_F + \mathcal{L}_{\text{gauge}} \\
&= \bar{\psi}(i\gamma^\mu D_\mu - M)\psi - \frac{1}{4} F_{\mu\nu}^i F^{i\mu\nu} \\
\mathcal{L} &= \mathcal{L}_{\text{sym}} + \mathcal{L}_{\text{breaking}} \\
\mathcal{L}_o &= \partial_\mu \phi^\dagger \partial^\mu \phi - m^2 \phi^\dagger \phi \\
\phi(x) &\leftrightarrow \phi^\dagger(x), \\
\mathcal{L}_I &= -\frac{\lambda}{4} (\phi^\dagger \phi)^2.
\end{aligned}$$

$$\mathcal{L} = \frac{1}{2} \partial_\mu \varphi_1 \partial^\mu \varphi_1 + \frac{1}{2} \partial_\mu \varphi_2 \partial^\mu \varphi_2 - V(\varphi_1^2 + \varphi_2^2),$$



$$\begin{pmatrix} \varphi_1 \\ \varphi_2 \end{pmatrix} \rightarrow \begin{pmatrix} \varphi'_1 \\ \varphi'_2 \end{pmatrix} = \begin{pmatrix} \cos \theta & -\sin \theta \\ \sin \theta & \cos \theta \end{pmatrix} \begin{pmatrix} \varphi_1 \\ \varphi_2 \end{pmatrix}.$$

$$V(\varphi_1, \varphi_2) = V(\varphi_{01}, \varphi_{02}) + \sum_{a=1,2} \left(\frac{\partial V}{\partial \varphi_a} \right)_0 (\varphi_0 - \varphi_{0a})$$

$$+ \frac{1}{2} \sum_{a,b=1,2} \left(\frac{\partial^2 V}{\partial \varphi_a \partial \varphi_b} \right)_0 (\varphi_a - \varphi_{0a})(\varphi_b - \varphi_{0b}) + \dots$$

$$m_{ab}^2 = \begin{pmatrix} m^2 & 0 \\ 0 & m^2 \end{pmatrix},$$

$$V(\varphi_1^2 + \varphi_2^2) = -\frac{\mu^2}{2}(\varphi_1^2 + \varphi_2^2) + \frac{\lambda}{4}(\varphi_1^2 + \varphi_2^2)^2.$$

$$\left(\frac{\partial V}{\partial \varphi_1} \right)_0 = -\mu^2 \varphi_{01} + \lambda \varphi_{01}(\varphi_1^2 + \varphi_2^2) = 0$$

$$\left(\frac{\partial V}{\partial \varphi_2} \right)_0 = -\mu^2 \varphi_{02} + \lambda \varphi_{02}(\varphi_1^2 + \varphi_2^2) = 0$$

$$\varphi_{01}^2 + \varphi_{02}^2 = \frac{\mu^2}{\lambda} = v^2$$

$$(\phi^\dagger \phi)_0 = |\phi_0|^2 = \frac{\mu^2}{2\lambda} = \frac{v^2}{2}$$

$$\frac{\partial^2 V}{\partial \varphi_1^2} = -\mu^2 + \lambda(\varphi_1^2 + \varphi_2^2) + 2\lambda\varphi_1^2$$

$$\frac{\partial^2 V}{\partial \varphi_2^2} = -\mu^2 + \lambda(\varphi_1^2 + \varphi_2^2) + 2\lambda\varphi_2^2$$

$$\frac{\partial^2 V}{\partial \varphi_1 \partial \varphi_2} = 2\lambda\varphi_1\varphi_2$$

$$m_{ab}^2 = \begin{pmatrix} 2\lambda v^2 & 0 \\ 0 & 0 \end{pmatrix}$$

$$\eta(x) = \varphi_1(x) - \sqrt{\frac{\mu^2}{\lambda}}, \xi = \varphi_2$$

$$\mathcal{L} = \left[\frac{1}{2} (\partial_\mu \eta)(\partial^\mu \eta) - \mu^2 \eta^2 \right] + \left[\frac{1}{2} (\partial_\mu \xi)(\partial^\mu \xi) \right]$$

$$- \lambda \sqrt{\frac{\mu^2}{\lambda}} \eta(\eta^2 + \xi^2) - \frac{\lambda}{4}(\eta^2 + \xi^2)^2 + \frac{\mu^4}{4\lambda}$$

$$\mathcal{L} = |D^\mu \phi|^2 - \mu^2 |\phi|^2 - \lambda (\phi^\dagger \phi)^2 - \frac{1}{4} F_{\mu\nu} F^{\mu\nu}$$

$$\phi = \frac{\varphi_1 + i\varphi_2}{\sqrt{2}}$$

$$D_\mu = \partial_\mu - iqA_\mu$$

$$F_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu$$

$$\phi(x) \rightarrow \phi'(x) = e^{-i\alpha} \phi(x)$$

$$\begin{aligned}
\phi(x) &\rightarrow \phi'(x) = e^{-i\alpha(x)}\phi(x) \\
A_\mu &\rightarrow A'_\mu = A_\mu - \frac{1}{q}\partial_\mu\alpha(x) \\
|\phi_0|^2 &= \frac{\mu^2}{2|\lambda|} = \frac{v^2}{2} \\
\langle\phi\rangle_0 &= \frac{v}{\sqrt{2}} \\
\phi' &= \phi - \langle\phi\rangle_0 \\
\phi(x) &= \frac{1}{\sqrt{2}}(v + \eta)e^{i\xi/v} \\
&\approx \frac{v + \eta + i\xi}{\sqrt{2}} \\
\mathcal{L}_{so} &= \frac{1}{2}[(\partial_\mu\eta)(\partial^\mu\eta) + 2\mu^2\eta^2] + \frac{1}{2}(\partial_\mu\xi)(\partial^\mu\xi) \\
&\quad - \frac{1}{4}F_{\mu\nu}F^{\mu\nu} + qvA_\mu(\partial^\mu\xi) + \frac{q^2v^2}{2}A_\mu A^\mu + \dots \\
\phi(x) &\rightarrow \phi'(x) = e^{i\xi(x)/v}\phi(x) = \frac{1}{\sqrt{2}}(v + \eta) \\
A_\mu(x) &\rightarrow B_\mu(x) = A_\mu(x) - \frac{1}{qv}\partial_\mu\xi(x) \\
D_\mu\phi(x) &\rightarrow D'_\mu\phi'(x) = (\partial_\mu - iqB_\mu)\frac{1}{\sqrt{2}}(v + \eta) \\
F_{\mu\nu}(A) &\rightarrow F_{\mu\nu}(B) = \partial_\mu B_\nu - \partial_\nu B_\mu \\
\mathcal{L}_{so} &= \frac{1}{2}|\partial_\mu\eta - iqB_\mu(v + \eta)|^2 - \frac{\mu^2}{2}(v + \eta)^2 - \frac{\lambda}{4}(v + \eta)^4 \\
&\quad - \frac{1}{4}F_{\mu\nu}(B)F^{\mu\nu}(B) \\
&= \frac{1}{2}[(\partial_\mu\eta)(\partial^\mu\eta) - 2\mu^2\eta^2] - \frac{1}{4}F_{\mu\nu}(B)F^{\mu\nu}(B) \\
&\quad + \frac{q^2v^2}{2}B_\mu B^\mu + \frac{1}{2}q^2B_\mu B^\mu\eta(\eta + 2v) - \lambda v\eta^3 - \frac{\lambda}{4}\eta^4 \\
\phi &= \begin{pmatrix} \phi_1 \\ \phi_2 \end{pmatrix} \\
\mathcal{L} &= (D_\mu\phi)^\dagger(D^\mu\phi) - \frac{1}{4}F_{\mu\nu}^i F^{i\mu\nu} - V(\phi^\dagger\phi) \\
F_{\mu\nu}^i &= \partial_\mu A_\nu^i - \partial_\nu A_\mu^i + g\varepsilon_{ijk}A_\mu^j A_\nu^k, \\
D_\mu\phi &= \left(\partial_\mu - ig\frac{\tau^i}{2}A_\mu^i\right)\phi, (i = 1, 2, 3) \\
V(\phi^\dagger\phi) &= -\mu^2\phi^\dagger\phi + \lambda(\phi^\dagger\phi)^2. (\mu^2 > 0) \\
\phi(x) &= \frac{1}{\sqrt{2}}e^{i\tau^i\xi^i(x)/2v} \begin{pmatrix} 0 \\ v + H(x) \end{pmatrix} \\
\phi(x) &\rightarrow \phi'(x) = U(x)\phi(x) = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ v + H(x) \end{pmatrix} \\
\vec{A}_\mu &\rightarrow \vec{B}_\mu = U\vec{A}_\mu U^{-1} - \frac{i}{g}(\partial_\mu U)U^{-1} \\
U(x) &= \frac{1}{\sqrt{2}}e^{i\tau^i\xi^i(x)/2v} \\
D_\mu\phi &\rightarrow (D_\mu\phi)' = \left(\partial_\mu - ig\frac{\tau^i}{2}B_\mu^i\right)\frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ v + H(x) \end{pmatrix} \\
F_{\mu\nu}^i(A)F^{i\mu\nu}(A) &\rightarrow F_{\mu\nu}^i(B)F^{i\mu\nu}(B) = F_{\mu\nu}^i(A)F^{i\mu\nu}(A)
\end{aligned}$$

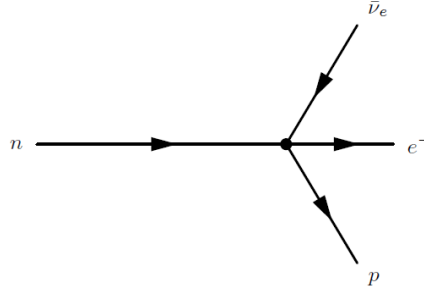
$$F_{\mu\nu}^i(B) = \partial_\mu B_\nu^i - \partial_\nu B_\mu^i + g\varepsilon_{ijk}B_\mu^j B_\nu^k$$

$$\mathcal{L} = (D_\mu\phi)'^\dagger (D^\mu\phi)' - \frac{1}{4}F_{\mu\nu}^i F^{i\mu\nu} + \mu^2\phi'^\dagger\phi' - \lambda(\phi'^\dagger\phi')^2$$

$$[(D_\mu\phi)']^{\dagger a} (D^\mu\phi)'_a = \frac{1}{2}\partial_\mu H\partial^\mu H + g^2 B_\mu^i B^{j\mu} \left(\frac{\tau^i}{2}\right)_b^a \left(\frac{\tau^j}{2}\right)_c^a b'^b b'_c = \frac{1}{2}\partial_\mu H\partial^\mu H + g^2 B_\mu^i B^{j\mu} (v+H)^2$$

$$\mathcal{L} = \frac{1}{2}(\partial_\mu H\partial^\mu H - 2\mu^2 H^2) - \frac{1}{4}F_{\mu\nu}^i(B)F^{i\mu\nu}(B) + \frac{g^2 v^2}{8}B_\mu^i B^{i\mu} + \frac{g^2}{8}B_\mu^i B^{i\mu}H(2v+H) - \lambda v H^3 - \frac{\lambda}{4}H^4 - \frac{v^4}{4}$$

$$\mathcal{H}_I = \frac{G_F}{\sqrt{2}}J_\mu^\dagger J^\mu$$



$$SU(2)_L \times U(1)_Y \xrightarrow{\text{SSB}} U(1)_{EM}$$

$$L = \begin{pmatrix} \nu_e \\ e \end{pmatrix}_L, R = e_R$$

$$SU(2)_L: L \rightarrow L' = e^{-i\alpha^i(x)\frac{\tau^i}{2}}L, R \rightarrow R' = R$$

$$U(1)_Y: L \rightarrow L' = e^{i\beta(x)}L, R \rightarrow R' = e^{i\beta(x)}R$$

$$\mathcal{L}_F = \bar{L}i\gamma^\mu \left(\partial_\mu - ig\frac{\vec{\tau}}{2} \cdot \vec{A}_\mu + \frac{i}{2}g'B_\mu \right) L$$

$$+ \bar{R}i\gamma^\mu (\partial_\mu + ig'B_\mu)R$$

$$D_\mu = \partial_\mu - ig\frac{\vec{\tau}}{2} \cdot \vec{A}_\mu - ig'\frac{Y}{2}B_\mu$$

$$\mathcal{L}_G = -\frac{1}{4}F_{\mu\nu}^i F^{i\mu\nu} - \frac{1}{4}B_{\mu\nu}B^{\mu\nu}$$

$$F_{\mu\nu}^i = \partial_\mu A_\nu^i - \partial_\nu A_\mu^i + g\varepsilon_{ijk}A_\mu^j A_\nu^k$$

$$B_{\mu\nu} = \partial_\mu B_\nu - \partial_\nu B_\mu$$

$$\phi = \begin{pmatrix} \phi^+ \\ \phi^0 \end{pmatrix}$$

$$\mathcal{L}_S = (D_\mu\phi)^\dagger (D^\mu\phi) - V(\phi^\dagger\phi)$$

$$D_\mu\phi = \left(\partial_\mu - ig\frac{\vec{\tau}}{2} \cdot \vec{A}_\mu - \frac{i}{2}g'B_\mu \right) \phi$$

$$V(\phi^\dagger\phi) = m^2\phi^\dagger\phi + \lambda(\phi^\dagger\phi)^2$$

$$\mathcal{L}_Y = -G_e(\bar{L}\phi R + \bar{R}\phi^\dagger L) + h.c.$$

$$\mathcal{L} = \mathcal{L}_F + \mathcal{L}_G + \mathcal{L}_S + \mathcal{L}_Y$$

$$\phi_0 = \langle 0|\phi|0\rangle = \begin{pmatrix} 0 \\ v/\sqrt{2} \end{pmatrix}$$

$$\phi = \begin{pmatrix} \phi^+ \\ \phi^0 \end{pmatrix} = e^{i\vec{\tau} \cdot \vec{\xi}/2v} \begin{pmatrix} 0 \\ (v+H)/\sqrt{2} \end{pmatrix}$$

$$\langle 0|\xi_i|0\rangle = \langle 0|H|0\rangle = 0$$

$$U(\xi) = e^{-i\vec{\tau} \cdot \vec{\xi}/2v}$$

$$\begin{aligned}
\phi' &= U(\xi)\phi = \begin{pmatrix} 0 \\ (v+H)/\sqrt{2} \end{pmatrix} = \frac{1}{\sqrt{2}}(v+H)\chi \\
L' &= U(\xi)L \\
\vec{A}'_\mu &= U(\xi)\vec{A}_\mu U(\xi)^{-1} - \frac{i}{g}(\partial_\mu U(\xi))U^\dagger(\xi) \\
R' &= R \\
B'_\mu &= B_\mu \\
\mathcal{L}_F &= \bar{L}'i\gamma^\mu \left(\partial_\mu - ig\frac{\vec{\tau}}{2} \cdot \vec{A}'_\mu + \frac{i}{2}g'B'_\mu \right) L' \\
&\quad + \bar{R}'i\gamma^\mu (\partial_\mu + ig'B'_\mu) R' \\
\mathcal{L}_G &= -\frac{1}{4}F_{\mu\nu}^i F'^i{}_{\mu\nu} - \frac{1}{4}B'_{\mu\nu} B'^{\mu\nu} \\
\mathcal{L}_S &= (D_\mu\phi)'(D^\mu\phi)' - V(\phi'^\dagger\phi') \\
\mathcal{L}_Y &= -G_e(\bar{L}'\phi'R' + \bar{R}'\phi'^\dagger L') + h.c. \\
\mathcal{L}_S &= (D_\mu\phi)'(D^\mu\phi)' - V(\phi'^\dagger\phi') \\
(D_\mu\phi)' &= \left(\partial_\mu - ig\frac{\vec{\tau}}{2} \cdot \vec{A}'_\mu - \frac{i}{2}g'B'_\mu \right) \phi' \\
&= \left(\partial_\mu - ig\frac{\vec{\tau}}{2} \cdot \vec{A}'_\mu - \frac{i}{2}g'B'_\mu \right) \frac{1}{\sqrt{2}}(v+H)\chi \\
\mathcal{L}_{\text{mass}} &= \frac{v^2}{2}\chi^\dagger \left(g\frac{\vec{\tau}}{2} \cdot \vec{A}'_\mu + \frac{g'}{2}B'_\mu \right) \left(g\frac{\vec{\tau}}{2} \cdot \vec{A}'^\mu + \frac{g'}{2}B'^\mu \right) \chi \\
&= \frac{v^2}{8} (g^2 \vec{A}'_\mu \cdot \vec{A}'^\mu + g'^2 B'_\mu B'^\mu - 2gg' B'_\mu A'^{3\mu}) \\
&= \frac{v^2}{8} (g^2 A'^{1\mu} A'^{1\mu} + g^2 A'^{2\mu} A'^{2\mu} + (gA'^{3\mu} - g'B'_\mu)^2), \\
W_\mu^\pm &= \frac{A'^{1\mu} \mp iA'^{2\mu}}{\sqrt{2}} \\
M_W &= \frac{1}{2}gv \\
\frac{v^2}{8} \begin{pmatrix} A'^3_\mu & B'_\mu \end{pmatrix} \begin{pmatrix} g^2 & -gg' \\ -gg' & g'^2 \end{pmatrix} \begin{pmatrix} A'^{3\mu} \\ B'^\mu \end{pmatrix} \\
\frac{v^2}{8} \begin{pmatrix} Z_\mu & A_\mu \end{pmatrix} \begin{pmatrix} g^2 + g'^2 & 0 \\ 0 & 0 \end{pmatrix} \begin{pmatrix} Z^\mu \\ A_\mu \end{pmatrix} &= \frac{v^2}{8} (g^2 + g'^2) Z_\mu Z^\mu \\
\begin{pmatrix} Z_\mu \\ A_\mu \end{pmatrix} &= \begin{pmatrix} \cos\theta_W & -\sin\theta_W \\ \sin\theta_W & \cos\theta_W \end{pmatrix} \begin{pmatrix} A'^3_\mu \\ B'_\mu \end{pmatrix} \\
\tan\theta_W &= \frac{g'}{g} \\
\sin\theta_W &= \frac{g'}{\sqrt{g^2 + g'^2}}, \quad \cos\theta_W = \frac{g}{\sqrt{g^2 + g'^2}} \\
M_Z &= \frac{1}{2}v\sqrt{g^2 + g'^2} \\
M_Z &= \frac{M_W}{\cos\theta_W} \\
M_W &\approx \frac{37.22}{\sin\theta_W} \text{ GeV} > 37 \text{ GeV} \\
M_Z &= \frac{M_W}{\cos\theta_W} \approx \frac{74.44}{\sin 2\theta_W} \text{ GeV} > 74 \text{ GeV}. \\
M_W &= 79.8 \pm 0.8 \text{ GeV}, \quad M_Z = 90.8 \pm 0.6 \text{ GeV}
\end{aligned}$$

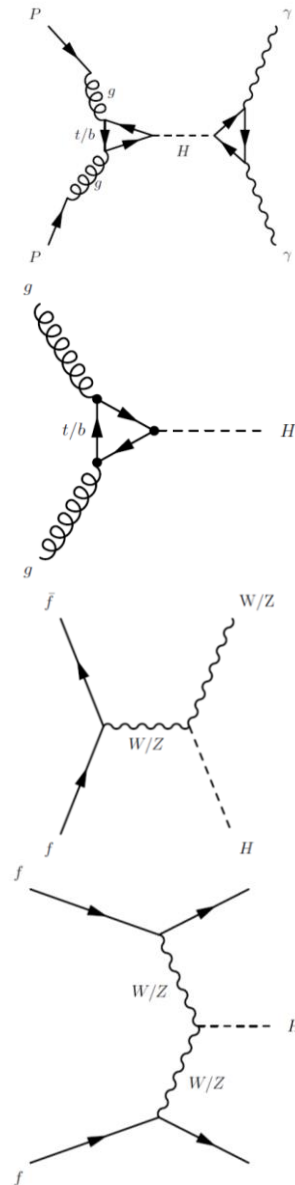


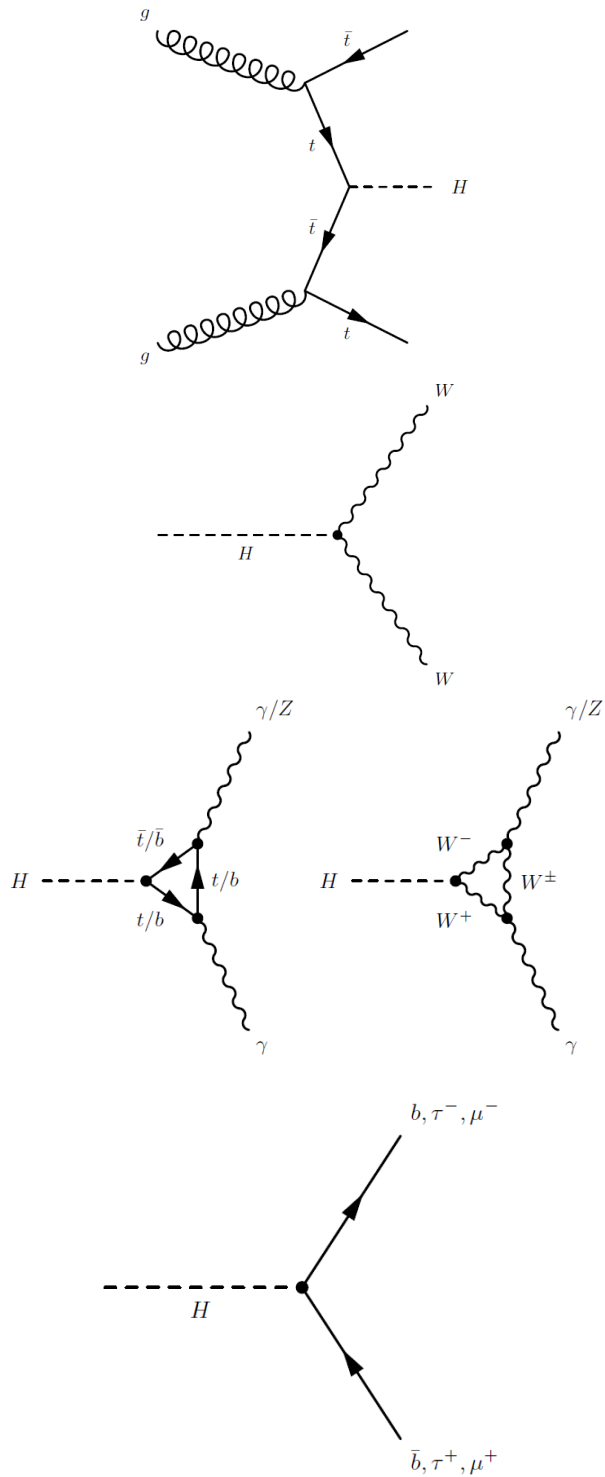
$$M_W = 80.4335 \pm 0.094 \text{ GeV}, M_Z = 91.1876 \pm 0.021 \text{ GeV}$$

$$\begin{aligned} V(\phi'^{\dagger}\phi') &= -\frac{\mu^2}{2}(v+H)^2\chi^{\dagger}\chi + \frac{\lambda}{4}(v+H)^4(\chi^{\dagger}\chi)^2 \\ &= \frac{-\mu^2 v^2}{4} + \frac{1}{2}(2\mu^2)H^2 + \lambda v H^3 + \frac{\lambda}{4}H^4. \\ M_H &= \sqrt{2\mu^2} \end{aligned}$$

$$\begin{aligned} \mathcal{L}_S &= (D_{\mu}\phi)'\dagger(D^{\mu}\phi)' - V(\phi'^{\dagger}\phi') \\ &= \frac{1}{2}\partial_{\mu}H\partial^{\mu}H - \frac{1}{2}M_H^2H^2 - \lambda v H^3 - \frac{\lambda}{4}H^4 + \frac{g^2}{8}(H^2 + 2Hv)\left[\frac{1}{\cos^2\theta_W}Z_{\mu}Z^{\mu} + 2W_{\mu}^{+}W^{-\mu}\right] + M_W^2W_{\mu}^{+}W^{-\mu} + \frac{1}{2}M_Z^2Z_{\mu}Z^{\mu} \\ \mathcal{L}_Y &= -G_e(\bar{L}'\phi'R' + \bar{R}'\phi'\dagger L') + h.c. = -G_e\left(\bar{e}'_L\frac{1}{\sqrt{2}}(v+H)e'_R + \bar{e}'_R\frac{1}{\sqrt{2}}(v+H)e'_L\right) + h.c. \\ &= -\frac{G_e v}{\sqrt{2}}\bar{e}'e' - \frac{G_e}{\sqrt{2}}H\bar{e}'e'. \end{aligned}$$

$$\begin{aligned} m_e &= \frac{G_e}{\sqrt{2}}v \\ \frac{G_e}{\sqrt{2}} &= \frac{m_e}{v}, \end{aligned}$$





$$i \frac{\partial \psi}{\partial t} + \frac{1}{2m} \nabla^2 \psi - V(\mathbf{r}, t) \psi = 0$$

$$E = \sqrt{p^2 + m^2},$$

$$\mathbf{p}^2 + m^2 = \mathbf{E}^2$$

$$\left(\frac{\partial^2}{\partial t^2} - \nabla^2 + m^2 \right) \psi = 0$$

$$-\eta^{\mu\nu} \partial_\mu \partial_\nu \equiv \sum_{\mu=0}^3 \sum_{\nu=0}^3 -\eta^{\mu\nu} \partial_\mu \partial_\nu = \frac{\partial^2}{\partial t^2} - \nabla^2$$

$$(\square + m^2) \psi = 0,$$



$$\begin{aligned}
& (\square + m^2)\psi^* = 0. \\
& \psi^* \square \psi - \psi \square \psi^* = 0 \\
& \psi^* \left(\frac{\partial^2}{\partial t^2} - \nabla^2 \right) \psi - \psi \left(\frac{\partial^2}{\partial t^2} - \nabla^2 \right) \psi^* = 0 \\
& \psi^* \frac{\partial^2 \psi}{\partial t^2} - \psi^* \nabla^2 \psi - \psi \frac{\partial^2 \psi^*}{\partial t^2} + \psi \nabla^2 \psi^* = 0 \\
& \frac{\partial}{\partial t} \left(\psi^* \frac{\partial \psi}{\partial t} - \psi \frac{\partial \psi^*}{\partial t} \right) - \nabla \cdot (\psi^* \nabla \psi - \psi \nabla \psi^*) = 0 \\
& \frac{\partial \rho}{\partial t} + \nabla \cdot \mathbf{J} = 0 \\
& \mathbf{J} = \frac{1}{2im} (\psi^* \nabla \psi - \psi \nabla \psi^*) \\
& \rho = \frac{i}{2m} \left(\psi^* \frac{\partial \psi}{\partial t} - \psi \frac{\partial \psi^*}{\partial t} \right) \\
& i \frac{\partial \psi}{\partial t} = \hat{H} \psi \\
& i \frac{\partial \psi}{\partial t} = \left[-i \left(\hat{\alpha}_1 \frac{\partial}{\partial x^1} + \hat{\alpha}_2 \frac{\partial}{\partial x^2} + \hat{\alpha}_3 \frac{\partial}{\partial x^3} \right) + \hat{\beta} m \right] \psi \\
& = \left[-i \sum_{k=1}^N \hat{\alpha}_k \frac{\partial}{\partial x^k} + \hat{\beta} m \right] \psi = \hat{H}_D \psi \\
& \psi = \begin{pmatrix} \psi_1(\mathbf{x}, t) \\ \psi_2(\mathbf{x}, t) \\ \vdots \\ \psi_N(\mathbf{x}, t) \end{pmatrix} \\
& \rho = \psi^\dagger \psi = \sum_{i=1}^N \psi_i^\dagger \psi_i \\
& -\frac{\partial^2 \psi}{\partial t^2} = (-\nabla^2 + m^2)\psi = 0 \\
& -\frac{\partial^2 \psi}{\partial t^2} = -\sum_{i,j=1}^3 \frac{\hat{\alpha}_i \hat{\alpha}_j + \hat{\alpha}_j \hat{\alpha}_i}{2} \frac{\partial \psi}{\partial x^i \partial x^j} \\
& -im \sum_{i=1}^3 (\hat{\alpha}_i \hat{\beta} + \hat{\beta} \hat{\alpha}_i) \frac{\partial \psi}{\partial x^i \partial x^j} + \hat{\beta}^2 m^2 \psi \\
& \hat{\alpha}_i \hat{\alpha}_j + \hat{\alpha}_j \hat{\alpha}_i = \{\hat{\alpha}_i, \hat{\alpha}_j\} = 2\delta_{ij} I_N \\
& \hat{\alpha}_i \hat{\beta} + \hat{\beta} \hat{\alpha}_i = 0 \\
& \hat{\alpha}^2 = \hat{\beta}^2 = I_N \\
& \hat{\alpha}_i = -\hat{\beta} \hat{\alpha}_i \hat{\beta} \\
& \text{Tr} \hat{\alpha}_i = -\text{Tr} \hat{\beta} \hat{\alpha}_i \hat{\beta} = -\text{Tr} \hat{\beta}^2 \hat{\alpha}_i = -\text{Tr} \hat{\alpha}_i \\
& \hat{\sigma}_1 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \hat{\sigma}_2 = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}, \hat{\sigma}_3 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \\
& \hat{\alpha}_i = \begin{pmatrix} 0 & \hat{\sigma}_i \\ \hat{\sigma}_i & 0 \end{pmatrix}, \hat{\beta} = \begin{pmatrix} I_2 & 0 \\ 0 & I_2 \end{pmatrix} \\
& i\psi^\dagger \frac{\partial \psi}{\partial t} = -i \sum_{k=1}^N \psi^\dagger \hat{\alpha}_k \frac{\partial \psi}{\partial x^k} + \psi^\dagger \hat{\beta} m \psi. \\
& -i\psi \frac{\partial \psi^\dagger}{\partial t} = i \sum_{k=1}^N \psi \hat{\alpha}_k \frac{\partial \psi^\dagger}{\partial x^k} + \psi \hat{\beta} m \psi^\dagger
\end{aligned}$$

$$\begin{aligned}
i\psi^\dagger \frac{\partial \psi}{\partial t} + i\psi \frac{\partial \psi^\dagger}{\partial t} &= -i \sum_{k=1}^N \psi^\dagger \hat{\alpha}_k \frac{\partial \psi}{\partial x^k} + \psi^\dagger \hat{\beta} m \psi \\
&\quad -i \sum_{k=1}^N \psi \hat{\alpha}_k \frac{\partial \psi^\dagger}{\partial x^k} - \psi \hat{\beta} m \psi^\dagger \\
\frac{\partial}{\partial t} (\psi^\dagger \psi) + \sum_{k=1}^N \frac{\partial}{\partial x^k} (\psi^\dagger \hat{\alpha}_k \psi) &= 0 \\
\frac{\partial \rho}{\partial t} + \nabla \cdot J &= 0 \\
\rho = \psi^\dagger \psi = \sum_{i=1}^4 \psi_i^* \psi & \\
\frac{\partial}{\partial t} \int d^3x \psi^\dagger \psi &= 0 \\
\left(i \frac{\partial}{\partial t} + i \sum_{k=1}^3 \hat{\alpha}_k \frac{\partial}{\partial x^k} - \hat{\beta} m \right) \psi &= 0 \\
\hat{\beta} \left(i \frac{\partial}{\partial t} + i \sum_{k=1}^3 \hat{\beta} \hat{\alpha}_k \frac{\partial}{\partial x^k} - m \right) \psi &= 0 \\
i \left(\gamma^0 \frac{\partial}{\partial x^0} + \gamma^1 \frac{\partial}{\partial x^1} + \gamma^2 \frac{\partial}{\partial x^2} + \gamma^3 \frac{\partial}{\partial x^3} \right) \psi - m \psi &= 0 \\
\gamma^\mu \gamma^\nu + \gamma^\nu \gamma^\mu &= 2g^{\mu\nu} I_4 \\
\gamma^\mu = \begin{pmatrix} 0 & \hat{\sigma}^\mu \\ -\hat{\sigma}^\mu & 0 \end{pmatrix}, \gamma^0 = \begin{pmatrix} I_2 & 0 \\ 0 & -I_2 \end{pmatrix} & \\
(i\gamma^\mu \partial_\mu - m)\psi &= 0 \\
\mathcal{L}(x) = \bar{\psi}(x)(i\gamma^\mu \partial_\mu - m)\psi(x) & \\
\frac{\partial \mathcal{L}}{\partial \psi(x)} = -\bar{\psi}(x)m & \\
\frac{\partial \mathcal{L}}{\partial (\partial_\mu \psi(x))} = i\gamma^\mu \bar{\psi}(x) & \\
\frac{\partial \mathcal{L}}{\partial \bar{\psi}} - \partial_\mu \frac{\partial \mathcal{L}}{\partial (\partial_\mu \bar{\psi})} = -\bar{\psi}(x)m - i\gamma^\mu \partial_\mu \bar{\psi}(x) &= 0 \\
\bar{\psi}(x)(i\gamma^\mu \partial_\mu + m) &= 0 \\
\psi^\dagger(x)\gamma^0(i\gamma^\mu \partial_\mu + m) &= 0 \\
\gamma^0(-i(\gamma^0 \gamma^\mu \gamma^0) \partial_\mu + m)\psi &= 0 \\
(-i(\gamma^\mu \gamma^0) \partial_\mu + m\gamma^0)\psi &= 0 \\
(i\gamma^\mu \partial_\mu - m)\psi &= 0
\end{aligned}$$

Los operadores perturbativos del gravitón y por ende, del campo gravitónico, están dados por:

$$\hat{H}_{\text{int},l} = \frac{L}{\pi^2} \sqrt{\frac{M\hbar}{\omega_l}} \frac{(-1)^{\frac{l-1}{2}}}{l^2} (\hat{b}_l + \hat{b}_l^\dagger) \ddot{h}$$

Satisfaciendo la relación de dispersión, así:

$$\hat{h} = \sum_{\mathbf{k}} h_{q,\mathbf{k}} (\hat{a}_{\mathbf{k}} + \hat{a}_{\mathbf{k}}^\dagger), \text{ donde } h_{q,\mathbf{k}} = \frac{1}{c} \sqrt{\frac{8\pi G \hbar}{V \nu_{\mathbf{k}}}} \text{ es igual a } \hat{a}_{\mathbf{k}}(t) (\hat{a}_{\mathbf{k}}^\dagger(t))$$



En tanto que, el efecto cuántico gravitacional estimulado por el campo gravitónico, viene dado por la escala de Planck, la misma que se expresa así:

$$\begin{aligned}\Gamma_{\text{efecto cuántico gravitacional estimulado}} &= \frac{ML^2\omega_l^2 h^2}{4l^4\pi^5\hbar} = \frac{v_s^2}{4l^2\pi^3\hbar} Mh^2. \\ \Gamma_{\text{efecto cuántico gravitacional estimulado}} &= \frac{2\pi}{\hbar^2} |\langle\alpha|\langle 1|\hat{H}_{\text{int}}|\alpha\rangle|0\rangle|^2 D(\omega) \\ \Gamma_{\text{efecto cuántico gravitacional estimulado}} &= \frac{|\alpha|^2 8GML^2\omega_l^4}{l^4 \pi^4 c^5} \\ \Gamma_{\text{efecto cuántico gravitacional estimulado}} &= \frac{1}{l^4} \frac{ML^2\omega_l^2 h_0^2}{4\pi^5\hbar} \\ P(\omega, t) &\approx D(\omega) \frac{(h_0^2\omega^3 ML^2)t}{2\hbar\pi^3} = \Gamma_{\text{efecto cuántico gravitacional estimulado}} t \\ \Gamma_{\text{efecto cuántico gravitacional estimulado}} &= D(\omega) \frac{(h_0^2\omega^3 ML^2)}{2\hbar\pi^3} = \frac{Vh_0^2\omega^5 ML^2}{4\hbar\pi^5 c^3} = \frac{h_0^2 M v_s^2}{4\hbar\pi^3}.\end{aligned}$$

Más, el efecto cuántico gravitacional espontáneo, es decir, cuando no interviene el campo gravitónico, viene dado por la misma escala, pero así:

$$\Gamma_{\text{efecto cuántico gravitacional espontáneo}} = \frac{8GML^2\omega_l^4}{l^4\pi^4 c^5} = \frac{8\pi G\rho v_s^4 R^2}{Lc^5}$$

Calculándose la densidad y distribución de la masa de la superpartícula, sea por permeabilidad del campo gravitónico o en su defecto, por gravedad cuántica intrínseca, en cualquiera de los casos anteriores, se usará la siguiente métrica:

$$\begin{aligned}M &= \frac{\pi^2 \hbar \omega^3}{v_s^2 \chi(h, \omega, t)^2} \\ m\ddot{\xi}_n + m\omega_D^2(2\xi_n - \xi_{n-2} - \xi_{n+2}) &= 0. \\ \xi_n(t) &= e^{-i\omega t} (Ae^{ikna/2} + Be^{-ikna/2}) + \text{H.c.} \\ \mu_l &= \int_{-L/2}^{L/2} \rho(x) dx \chi_l(x)^2 = \frac{M}{L} \int_{-L/2}^{L/2} dx \chi_l(x)^2 = M/2\end{aligned}$$

Para una onda monocromática $h(t) = h_0 \sin(\nu t)$ tenemos que $\chi \sim \frac{1}{2} h_0 \nu^2 t \text{sinc}\left(\frac{t\Delta}{2}\right)$ en la que $\Delta = |\nu - \omega|$, cuya aproximación es $\Delta \ll \omega, \nu$ y cuyo límite es $t\Delta \rightarrow \infty$.

En cuanto a la expansión de Taylor es fase $\Phi(s)$ en relación a $h(s)e^{i\omega s}$, desde un punto estacionario $s = s^*$ donde $\Phi'(s)|_{s=s^*} = 0$, evaluamos la integral de Gauss para una solución aproximada de $\chi(h, \omega, t)$.

Calculando lo anterior, tenemos:

$$\begin{aligned}v(t) &= \left(\frac{1}{v_0^{8/3}} - \frac{8}{3} kt \right)^{-3/8} \\ k &= \frac{k_f}{(2\pi)^{8/3}} = \frac{96}{5} \pi \left(\frac{\pi G M_c}{c^3} \right)^{5/3} \frac{1}{(2\pi)^{8/3}} = \frac{48}{5} \left(\frac{G M_c}{2c^3} \right)^{5/3}\end{aligned}$$

En la que la masa $M_c = (m_1 m_2)^{3/5} / (m_1 + m_2)^{1/5}$ es efectiva para un sistema binario con masas m_1 and m_2 , lo que comporta un campo gravitónico.

Estimamos el tiempo $\tau [\omega - \Delta\omega, \omega + \Delta\omega]$, lo que da como resultado:

$$\tau = t(v = \omega + \Delta\omega) - t(v = \omega - \Delta\omega)$$



$$\tau = \frac{2\Delta\omega}{k\omega^{11/3}}$$

$$\chi = \left| \int_0^t ds e^{i\omega s} h_0 v^2 \sin(v s) \right| \approx h_0 v^2 \frac{t}{2} \operatorname{sinc}\left(\frac{\delta t}{2}\right)$$

Aquí asumimos que $\omega + v \gg \omega - v = \delta$, por lo que:

$$2\Delta\omega = \frac{8}{T}$$

Según la escala del tiempo, tenemos:

$$\tau = 2 \sqrt{\frac{2}{k}} \omega^{-11/6}$$

$$\chi \approx h_0 \omega^2 \left| \int_0^\tau ds e^{i\omega s} \sin(\omega s) \right|$$

$$\chi(\tau) \approx h_0 \frac{\omega^2 \tau}{2}$$

$$\chi \approx h_0 \sqrt{\frac{2}{k}} \omega^{1/6} = h_0 \sqrt{\frac{5}{24} \left(\frac{2c^3}{GM_c}\right)^{5/6}} \omega^{1/6}$$

$$M = \frac{\pi^2 \hbar \omega^3}{v_s^2 \chi^2} \approx \frac{\pi^2 \hbar k}{2v_s^2 h_0^2} \omega^{8/3} = \frac{24\pi^2}{5} \frac{\hbar}{h_0^2 v_s^2} \left(\frac{GM_c}{2c^3}\right)^{5/3} \omega^{8/3}$$

En tanto que la densidad de la energía de la superpartícula, se calculará así:

$$\left. \frac{d\xi_n}{dn} \right|_{n=\pm N} = 0$$

$$\xi_n(t) = \sum_{l=0,2,..}^{N-1} \chi_l(t) \cos[l\pi n/(2N)]$$

$$+ \sum_{l=1,3,..}^N \chi_l(t) \sin[l\pi n/(2N)].$$

$$\sum_{n=-N}^N \cos\left[\frac{l\pi n}{2(N+1)}\right] \cos\left[\frac{l'\pi n}{2(N+1)}\right] = \frac{N+1}{2} \delta_{ll'}$$

$$\sum_{n=-N}^N \sin\left[\frac{l\pi n}{2(N+1)}\right] \sin\left[\frac{l'\pi n}{2(N+1)}\right] = \frac{N+1}{2} \delta_{ll'}$$

$$\sum_{n=-N}^N \cos\left[\frac{l\pi n}{2(N+1)}\right] \sin\left[\frac{l'\pi n}{2(N+1)}\right] = 0.$$

$$\xi_n(t) = \sum_{l=0,2,..}^{N-1} \chi_l(t) \cos\left[\frac{l\pi n}{2(N+1)}\right]$$

$$+ \sum_{l=1,3,..}^N \chi_l(t) \sin\left[\frac{l\pi n}{2(N+1)}\right].$$

$$\ddot{\chi}_l + \omega_l^2 \chi_l = 0$$



$$\begin{aligned}
E &= \frac{m}{2} \sum_{n=-N}^N \xi_n^2 + \frac{m\omega_D^2}{2} \sum_{n=-N}^{N-2} (\xi_{n+2} - \xi_n)^2 \\
&= \frac{M}{4} \sum_{l=0}^N \chi_l^2 + \frac{M}{4} \sum_{l=0}^N \omega_l^2 \chi_l^2. \\
&\int_{-L/2}^{L/2} dx \chi_l(x)^2 = L/2 \\
\hat{H}_{\omega_l} &= \hbar \omega_l \left(\hat{b}_l^\dagger \hat{b}_l + \frac{1}{2} \right) \\
|\psi_M\rangle &= (2\pi t_m)^{-\frac{1}{4}} \int dx e^{-\frac{x^2}{4t_m}} |x\rangle \\
\hat{H}_{\text{int}}^M dt &= \sqrt{dt} \hat{p} \hat{N} \\
\hat{M}_{\hat{N}}(y) &= \langle y | e^{-i\hat{H}_{\text{int}}^M dt} | \psi_M \rangle \\
&= (2\pi t_m)^{-\frac{1}{4}} \exp \left\{ \left[-\frac{(y - \hat{N}\sqrt{dt})^2}{4t_m} \right] \right\}. \\
\hat{M}_{\hat{N}}(r) &= (2\pi t_m/dt)^{-\frac{1}{4}} \exp \left\{ \left[-\frac{dt(r - \hat{N})^2}{4t_m} \right] \right\} \\
r(t) &\approx \langle \hat{N}(t) \rangle + \sqrt{t_m} \zeta(t), \zeta(t), \langle \zeta(t) \zeta(t') \rangle = \delta(t - t') \otimes \rho(t + dt)
\end{aligned}$$

En este punto, cabe precisar, que las ondas gravitacionales causadas por efecto de la distorsión de la geometría del espacio – tiempo cuántico, a propósito de las interacciones de las superpartículas engendradas por la permeabilidad del campo gravitónico con la materia y en consecuencia, el potencial gravitacional del gravitón, en relación a la densidad de los estados inicial y final, se calculan así:

$$\begin{aligned}
f(x) &= -m\nabla\phi = -m\nabla \left(\frac{1}{2} \frac{\partial^2 \phi}{\partial x^2} x^2 + \dots \right) \\
&= m \frac{\ddot{h}_{xx}}{4} \nabla(x^2) = m \frac{\ddot{h}_{xx}}{2} x \\
f_n &= m \frac{\ddot{h}_{xx}}{2} (x_n + \xi_n) \\
H_I &= -m \frac{\ddot{h}_{xx}}{2} \sum_{n=-N}^N (x_n \xi_n + \xi_n^2/2) \\
-m \frac{\ddot{h}_{xx}}{2} \sum_{n=-N}^N x_n \xi_n &\approx -\frac{ML\ddot{h}_{xx}}{\pi^2} \sum_{l=1,3,\dots}^N (-1)^{\frac{l-1}{2}} \frac{1}{l^2} \chi_l(t) \\
-m \frac{\ddot{h}_{xx}}{4} \sum_{n=-N}^N \xi_n^2 &= -\frac{M\ddot{h}_{xx}}{8} \sum_{l=0}^N \chi_l^2 \\
\hat{H}_I &= \sum_{l=0}^N \hat{H}_l^I = -\frac{ML\ddot{h}_{xx}}{\pi^2} \sum_{l=1,3,\dots}^N (-1)^{\frac{l-1}{2}} \frac{1}{l^2} \hat{\chi}_l \\
&\quad -\frac{M\ddot{h}_{xx}}{8} \sum_{l=0}^N \hat{\chi}_l^2. \\
\hat{H}_l^I &= -\frac{ML\ddot{h}_{xx}}{\pi^2} (-1)^{\frac{l-1}{2}} \frac{1}{l^2} \sqrt{\frac{\hbar}{M\omega_l}} (\hat{b}_l + \hat{b}_l^\dagger) \\
&\quad -\frac{\ddot{h}_{xx}}{8} \frac{\hbar}{\omega_l} (\hat{b}_l + \hat{b}_l^\dagger)^2
\end{aligned}$$

$$\hat{H}_l^l = -\frac{\hbar_{xx}}{8} \frac{\hbar}{\omega_l} (\hat{b}_l + \hat{b}_l^\dagger)^2$$

$$\frac{\hat{H}_{int}}{\hbar} = \sqrt{(-1)^{l-1} \frac{8\pi GMv^3 L^2}{\omega_l c^2 V \pi^4 l^4}} (\hat{b}_l + \hat{b}_l^\dagger) (\hat{a} e^{-i\omega t} + \hat{a}^\dagger e^{i\omega t}),$$

$$N = \frac{h_0^2 c^5}{32\pi G \hbar v^2}$$

$$P_{0 \Rightarrow 1} \approx \frac{L^2}{4\pi^4} \frac{M}{\omega \hbar} h_0^2 v^4 t^2 \text{sinc}^2\left(\frac{\omega - \nu}{2} t\right)$$

$$P(t) = \frac{L^2}{\pi^4} \frac{M}{\omega \hbar} |\chi(t)|^2 = \frac{h_0^2 \omega t^2 M v_s^2}{4\pi^2 \hbar}$$

$$\Gamma_{mc} = \frac{dP(t)}{dt} = \frac{h_0^2 \omega t M v_s^2}{2\pi^2 \hbar} = \frac{h_0^2 N_c M v_s^2}{\pi \hbar}$$

$$h_0 = \sqrt{\frac{\pi k_B T}{M v_s^2 Q N_c}}$$

$$P(\nu, \omega, t) \approx |\beta(\nu, \omega, t)|^2 \approx \frac{(h_0^2 \omega^3 M L^2) \sin^2\left[\frac{1}{2} t(\nu - \omega)\right]}{\hbar(\nu - \omega)^2 \pi^4}$$

$$P(\omega, t) \approx \sum_{\nu} |\beta(\nu, \omega, t)|^2$$

$$= \int_{\omega - \delta/2}^{\omega + \delta/2} d\nu D(\nu) |\beta(\nu, \omega, t)|^2$$

$$= D(\omega) \frac{(h_0^2 \omega^3 M L^2)}{\hbar \pi^4} \int_{\omega - \delta/2}^{\omega + \delta/2} d\nu \frac{\sin^2\left(\frac{1}{2} t(\nu - \omega)\right)}{(\nu - \omega)^2}$$

$$= D(\omega) \frac{(h_0^2 \omega^3 M L^2)}{\hbar \pi^4} \Xi(t)$$

$$\sum_{\nu} = \int d\nu D(\nu) = \int d\nu \frac{V \nu^2}{2\pi^2 c^3}$$

$$\lim_{t \delta \gg 1} \Xi(t) \rightarrow \frac{1}{2} \pi t$$

$$h_c \equiv 2\pi \sqrt{\frac{\pi k_B T}{M v_s^2 Q}}$$

En el que, el operador – oscilador y los conmutadores del gravitón en relación a la partícula permeada por el campo gravitónico y viceversa, se calculan de la siguiente manera:

$$\hat{H} = \hbar \omega \hat{b}^\dagger \hat{b} + \frac{1}{\pi^2} L \sqrt{\frac{M \hbar}{\omega}} \ddot{h}(t) (\hat{b} + \hat{b}^\dagger)$$

$$\hat{U}_{int} = \hat{T} e^{-i \int_0^t ds (g(s) \hat{b}(s) + g^*(s) \hat{b}^\dagger(s))}$$

$$g(t) = \frac{1}{\pi^2} L \sqrt{\frac{M}{\hbar \omega}} \ddot{h}(t)$$

$$\hat{U} = e^{\Omega(t)}$$

$$\Omega(t) = \int_0^t dt_1 \hat{A}(t_1) + \frac{1}{2} \int_0^t dt_1 \int_0^{t_1} dt_2 [\hat{A}(t_1), \hat{A}(t_2)] + \dots$$

$$\begin{aligned}
\hat{A}(t) &= -i(g(t)\hat{b}(t) + g^*(t)\hat{b}^\dagger(t)), [\hat{A}(t_1), \hat{A}(t_2)] = -(g(t_1)g^*(t_2)e^{-i\omega(t_1-t_2)} - \\
&\quad g(t_2)g^*(t_1)e^{+i\omega(t_1-t_2)}) \\
\hat{U}_{int} &= e^{-i\int_0^t ds(g(s)\hat{b}(s) + g^*(s)\hat{b}^\dagger(s))} e^{-i\varphi} \\
\varphi &= \int_0^t dt_1 \int_0^{t_1} \text{Im}[g(t_1)g^*(t_2)e^{-i\omega(t_1-t_2)}] \\
\hat{U} &= e^{-i\varphi} e^{-i\omega t \hat{b}^\dagger \hat{b}} \hat{D}(\beta) \\
\hat{D}(\beta) &= e^{\beta \hat{b}^\dagger - \beta^* \hat{b}} \\
\beta &= -i \int_0^t ds g^*(s) e^{i\omega s} \\
&\quad e^{-i\varphi} |\beta e^{-i\omega t}\rangle \\
|\beta| &= \frac{L}{\pi^2} \sqrt{\frac{M}{\omega \hbar}} \chi(h, \omega, t) \\
\chi(h, \omega, t) &= \left| \int_0^t ds \ddot{h}(s) e^{i\omega s} \right|
\end{aligned}$$

2. Efecto cuántico gravitacional implícito, intrínseco o endógeno. Este fenómeno cuántico, se materializa cuando una partícula, a propósito de su masa hamiltoniana y densidad lagrangiana, se tiene como supermasiva, por lo que, aunque no se vea permeada por el campo de gauge gravitónico y en consecuencia, no interactúe con el gravitón, es capaz, por sí misma, de distorsionar el espacio – tiempo e incluso, colapsar, provocando un agujero negro de mecánica cuántica. Por ejemplo, la partícula subatómica más pesada que se conoce es el quark top, cuya masa es de aproximadamente 173 GeV (gigaelectronvoltios). El bosón de Higgs, que también es bastante pesado, tiene una masa de 125 GeV.

Paso a calcular, extrapolando las ecuaciones de campo de Einstein:

$$\begin{aligned}
L(\phi_a, \partial\phi_a) &= \int dV_x \mathcal{L}(\phi_a, \partial\phi_a) \\
S &= \int dv_x \mathcal{L}(\phi_a, \partial\phi_a) \\
\partial_\mu \left(\frac{\partial \mathcal{L}}{\partial(\partial_\mu \phi_a)} \right) - \frac{\partial \mathcal{L}}{\partial \phi_a} &= 0 \\
\Pi^a &= \frac{\partial \mathcal{L}}{\partial(\partial_0 \phi_a)} \\
[\phi_a(\vec{x}, t), \phi_b(\vec{x}', t)] &= [\Pi^a(\vec{x}, t), \Pi^b(\vec{x}', t)] = 0 \\
[\phi_a(\vec{x}, t), \Pi^b(\vec{x}', t)] &= i\delta_a^b \delta(\vec{x} - \vec{x}') \\
\phi_a(\vec{x})| \phi'_a \rangle &= \phi'_a | \phi'_a \rangle \\
i \frac{d}{dt} |\psi(t)\rangle &= H(\phi, \Pi) |\psi(t)\rangle \\
H &= \int dV_x \mathcal{H} = \int dV_x (\Pi^a \partial_0 \phi_a - \mathcal{L}) \\
\Theta^\mu{}_\nu &= \frac{\partial \mathcal{L}}{\partial(\partial_\mu \phi_a)} \partial_\nu \phi_a - \delta^\mu{}_\nu \mathcal{L} \\
\partial_\mu \Theta^\mu{}_\nu &= 0 \\
\int dv_x \partial_\mu \Theta^\mu{}_\nu &= 0 \\
\frac{d}{dt} \int dV_x \Theta^0{}_\nu &= 0 \\
P_\nu &= \int dV_x \Theta^0{}_\nu \\
x^\mu \rightarrow x'^\mu &= x^\mu + \delta x^\mu \\
\phi_a \rightarrow \phi'_a &= \phi_a + \delta_0 \phi_a
\end{aligned}$$



$$\begin{aligned}
\delta S &= G(t_2) - G(t_1), \\
G(t) &= \int dV_x (\Pi^a \delta \phi_a - \Theta_v^0 \delta x^v) \\
\mathcal{L}(\phi'_a, \partial^t \phi'_a) &= \mathcal{L}(\phi_a, \partial \phi_a) \\
J^\mu &= \frac{\partial \mathcal{L}}{\partial(\partial_\mu \phi_a)} \delta \phi_a - \Theta_v^\mu \delta x^v \\
\partial_\mu J^\mu &= 0 \\
\mathcal{L} &= \frac{1}{2} (\eta^{\mu\nu} \partial_\mu \phi \partial_\nu \phi - m^2 \phi^2) \\
(\partial^\mu \partial_\mu \phi + m^2) \phi &= 0 \\
g_{\vec{k}}(\vec{x}, t) &= \frac{1}{\sqrt{(2\pi)^{n-1} 2\omega_k}} e^{i(\vec{k}\vec{x} - \omega_k t)} \\
\phi &= \int d^{n-1}k (a_{\vec{k}}^\dagger g_{\vec{k}}^* + a_{\vec{k}} g_{\vec{k}}) \\
[a_{\vec{k}}, a_{\vec{k}'}] &= [a_{\vec{k}}^\dagger, a_{\vec{k}'}^\dagger] = 0 \\
[a_{\vec{k}}, a_{\vec{k}'}^\dagger] &= \delta(\vec{k} - \vec{k}') \\
H &= \frac{1}{(2\pi)^{n-1}} \int d^{n-1}k \omega_k a_{\vec{k}}^\dagger a_{\vec{k}} \\
[H, a_{\vec{k}}^\dagger] &= \omega_k a_{\vec{k}}^\dagger, \\
[H, a_{\vec{k}}] &= -\omega_k a_{\vec{k}} \\
|\vec{k}\rangle &= a_{\vec{k}}^\dagger |0\rangle. \\
a_{\vec{k}_1}^\dagger a_{\vec{k}_2}^\dagger a_{\vec{k}_3}^\dagger |0\rangle &= |\vec{k}_1, \vec{k}_2, \vec{k}_3\rangle \\
N_{\vec{k}} &= a_{\vec{k}}^\dagger a_{\vec{k}}, \\
G_{\mu\nu} &= 8\pi T_{\mu\nu} \\
ds^2 &= g^{\mu\nu} dx_\mu dx_\nu \\
S(\phi'_a, \nabla' \phi'_a, g'_{\mu\nu}(x')) &= S(\phi_a, \nabla \phi_a, g_{\mu\nu}(x)) \\
\eta_{\mu\nu} &\rightarrow g_{\mu\nu} \\
\partial_\mu &\rightarrow \nabla_\mu \\
d^n x &\rightarrow |g|^{1/2} d^n x = dv_x \\
S &= \int d^n x \mathcal{L}(\phi, \nabla \phi, g_{\mu\nu}) \\
\partial_\mu \left(\frac{\partial \mathcal{L}}{\partial(\partial_\mu \phi_a)} \right) - \frac{\partial \mathcal{L}}{\partial \phi_a} &= 0 \\
\delta S &= \int dv_x \left(\frac{\partial \mathcal{L}}{\partial \phi_a} \delta \phi_a + \frac{\partial \mathcal{L}}{\partial(\nabla_\mu \phi_a)} \nabla_\mu \delta \phi_a \right). \\
\frac{\partial \mathcal{L}}{\partial(\nabla_\mu \phi_a)} \nabla_\mu \phi_a &= \nabla_\mu \left(\frac{\partial \mathcal{L}}{\partial(\nabla_\mu \phi_a)} \delta \phi_a \right) - \nabla_\mu \left(\frac{\partial \mathcal{L}}{\partial(\nabla_\mu \phi_a)} \right) \delta \phi_a \\
\delta S &= \int dv_x \left(\frac{\partial \mathcal{L}}{\partial \phi_a} - \nabla_\mu \left(\frac{\partial \mathcal{L}}{\partial(\nabla_\mu \phi_a)} \right) \right) \delta \phi_a \\
\nabla_\mu \left(\frac{\partial \mathcal{L}}{\partial(\nabla_\mu \phi_a)} \right) - \frac{\partial \mathcal{L}}{\partial \phi_a} &= 0 \\
x^\mu &\rightarrow x'^\mu = x^\mu - \varepsilon^\mu(x) \\
g_{\mu\nu}(x) &\rightarrow g'_{\mu\nu}(x') = \frac{\partial x^\sigma}{\partial x'^\mu} \frac{\partial x^\lambda}{\partial x'^\nu} g_{\sigma\lambda}(x) \\
g'_{\mu\nu}(x') &= g'_{\mu\nu}(x - \varepsilon) = g'_{\mu\nu}(x) - \varepsilon^\rho \partial_\rho g'_{\mu\nu}(x)
\end{aligned}$$

$$\begin{aligned}
g'_{\mu\nu}(x') &= (\delta_\mu^\sigma - \varepsilon_{,\mu}^\sigma)(\delta_\nu^\lambda - \varepsilon_{,\nu}^\lambda)g_{\mu\nu}(x) \\
g'_{\mu\nu}(x) - \varepsilon^\rho g'_{\mu\nu,\rho}(x) &= g_{\mu\nu}(x) + g_{\mu\lambda}(x)\varepsilon_{,\nu}^\lambda + g_{\nu\sigma}(x)\varepsilon_{,\mu}^\sigma \\
\delta_0 g_{\mu\nu}(x) &\equiv g'_{\mu\nu}(x) - g_{\mu\nu}(x) \\
\delta_0 g_{\mu\nu}(x) &= \varepsilon_{\mu;\nu} + \varepsilon_{\nu;\mu} = \mathcal{L}_\varepsilon g_{\mu\nu} \\
S(g_{\mu\nu} + \delta_0 g_{\mu\nu}) &= S(g_{\mu\nu}) + \int d^n x \frac{\delta S}{\delta g_{\mu\nu}} \delta_0 g_{\mu\nu} \\
\delta S = S(g_{\mu\nu} + \delta_0 g_{\mu\nu}) - S(g_{\mu\nu}) &= \int d^n x \frac{\delta S}{\delta g_{\mu\nu}} \delta_0 g_{\mu\nu} = 0 \\
T^{\mu\nu} &\equiv -2|g|^{-1/2} \frac{\delta S}{\delta g_{\mu\nu}}, \\
- \int dv_x T^{\mu\nu} \varepsilon_{\nu;\mu} &= 0 \\
\nabla_\mu (T^{\mu\nu} \varepsilon_\nu) &= T_{;\mu}^{\mu\nu} \varepsilon_\nu + T^{\mu\nu} \varepsilon_{\nu;\mu} \\
\delta S = - \int dv_x \nabla_\mu (T^{\mu\nu} \varepsilon_\nu) + \int dv_x (\nabla_\mu T^{\mu\nu}) \varepsilon_\nu \\
\int dv_x (\nabla_\mu T^{\mu\nu}) \varepsilon_\nu &= 0, \\
\nabla_\mu T^{\mu\nu} &= 0. \\
\delta(g^{\mu\sigma} g_{\sigma\nu}) = \delta(\delta_\nu^\mu) &= 0, \\
-\delta(g_{\lambda\nu}) = \delta(g^{\mu\sigma}) g_{\sigma\nu} g_{\lambda\mu}, \\
\frac{\delta}{\delta g_{\mu\nu}} &= -g^{\mu\lambda} g^{\nu\sigma} \frac{\delta}{\delta g^{\lambda\sigma}}. \\
T_{\mu\nu} = T^{\alpha\beta} g_{\alpha\mu} g_{\beta\nu} &= -\frac{2}{|g|^{1/2}} \frac{\delta S}{\delta g_{\alpha\beta}} g_{\alpha\nu} g_{\beta\nu} = 2|g|^{-1/2} \frac{\delta S}{\delta g^{\mu\nu}}. \\
x^\mu \rightarrow x'^\mu = x^\mu + \delta x^\mu \\
\phi_a(x) \rightarrow \phi'_a(x) &= \phi(x) + \delta_0 \phi_a(x) \\
\delta S = S' - S &= \int dv'_x L(\phi'(x')_a, \nabla' \phi'_a(x), g_{\mu\nu}) - \int dv_x L(\phi(x)_a, \nabla \phi_a(x), g_{\mu\nu}) \\
S' &= \int_{V'} dv_x L(\phi_a(x) + \delta_0 \phi(x), \nabla \phi_a(x) + \nabla \delta_0 \phi_a(x), g_{\mu\nu}) = \\
&\int_{V'} dv_x L(\phi_a(x), \nabla \phi_a(x), g_{\mu\nu}) + \int_V dv_x \left(\frac{\partial L}{\partial \phi_a} \delta_0 \phi_a + \frac{\partial L}{\partial (\nabla_\mu \phi_a)} \nabla_\mu \delta_0 \phi_a \right) \\
&\int_{\partial V} d\sigma_\mu \delta x^\mu L(\phi_a, \nabla \phi_a) = \int_{V'} dv_x L(\phi_a, \nabla \phi_a) - \int_V dv_x L(\phi_a, \nabla \phi_a) \\
\delta S = S' - S &= \int_{\partial V} d\sigma_\mu \delta x^\mu L(\phi_a, \nabla \phi_a) + \int_V dv_x \left(\frac{\partial L}{\partial \phi_a} \delta_0 \phi_a + \frac{\partial L}{\partial (\nabla_\mu \phi_a)} \nabla_\mu \delta_0 \phi_a \right) \\
\frac{\partial L}{\partial (\nabla_\mu \phi_a)} \nabla_\mu (\delta_0 \phi_a) &= \nabla_\mu \left(\frac{\partial L}{\partial (\nabla_\mu \phi_a)} \delta_0 \phi_a \right) - \nabla_\mu \left(\frac{\partial L}{\partial (\nabla_\mu \phi_a)} \right) \delta_0 \phi_a \\
\delta S &= \int_{\partial V} d\sigma_\mu \delta x^\mu L(\phi_a, \nabla \phi_a) + \int_V dv_x \nabla_\mu \left(\frac{\partial L}{\partial (\nabla_\mu \phi_a)} \delta_0 \phi_a \right) \\
\delta S &= \int_V dv_x \nabla_\mu \left(\delta x^\mu L + \frac{\partial L}{\partial (\nabla_\mu \phi_a)} \delta_0 \phi_a \right) \\
\phi'_a(x) &= \phi'_a(x - \delta x) = \phi'(x') - \nabla_\mu \phi_a(x) \delta x^\mu \\
\delta S &= \int d^{n-1} x \left(\frac{\partial \mathcal{L}}{\partial (\partial_0 \phi_a)} \delta_0 \phi_a - \Theta_\nu^0 \delta x^\nu \right) \Big|_{t_1}^{t_2} \\
G(t) &= \int d^{n-1} x (\Pi^a \delta \phi_a - \Theta_\nu^0 \delta x^\nu) \\
\delta S &= G(t_2) - G(t_1) \\
i\delta_0 F &= [F, G]
\end{aligned}$$



$$\begin{aligned}
[\phi_a(\vec{x}, t), \phi_b(\vec{x}', t)] &= [\Pi_a(\vec{x}, t), \Pi_b(\vec{x}', t)] = 0 \\
[\phi_a(\vec{x}, t), \Pi_b(\vec{x}', t)] &= i\delta_{a,b}\delta^{(n-1)}(\vec{x}' - \vec{x}) \\
\delta\mathcal{L} &= \partial_\mu \left(\frac{\partial\mathcal{L}}{\partial(\partial_\mu\phi_a)} \delta\phi_a \right) = 0 \\
J^\mu &= \frac{\partial\mathcal{L}}{\partial(\partial_\mu\phi_a)} \delta\phi_a \\
\frac{d}{dt}G(t) &= \int d^{n-1}x \frac{\partial}{\partial t} (\Pi_a \delta\phi_a) = \int d^{n-1}x \partial_i \left(\frac{\partial\mathcal{L}}{\partial(\partial_i\phi_a)} \delta\phi_a \right) = 0 \\
p_\lambda &= \int d^{n-1}x \Theta_\nu^0 \\
x^\mu \rightarrow x'^\mu &= x^\mu - \varepsilon\xi^\mu(x), \\
\nabla_\mu(T^{\mu\nu}\xi_\nu) &= \frac{1}{|g|^{1/2}} \partial_\mu (|g|^{1/2} T^{\mu\nu} \xi_\nu) = 0 \\
P_\xi &= \int dV_x T^{0\nu} \xi_\nu \\
P_\lambda &= \int dV_x T^0_\nu \\
S(\phi_a, g_{\mu\nu}) &= S(\tilde{\phi}_a, \tilde{g}_{\mu\nu}) + \int_{\partial V} dV_x \dots, \\
\tilde{g}_{\mu\nu}(x) &= \Omega^2(x) g_{\mu\nu}(x), \\
\tilde{\phi}_a(x) &= \Omega^{2p}(x) \phi_a(x), \\
\Omega^2(x) &= 1 + \lambda(x), \\
\delta_0 g_{\mu\nu} &= \tilde{g}_{\mu\nu} - g_{\mu\nu} = \lambda g_{\mu\nu} \\
\delta_0 \phi_a &= \tilde{\phi}_a - \phi_a = p\lambda(x) \phi_a(x) \\
\delta S &= \int d^n x \left(\frac{\delta S}{\delta\phi_a} \delta_0 \phi_a + \frac{\delta S}{\delta g_{\mu\nu}} \delta_0 g_{\mu\nu} \right) = \int d^n x \left(\frac{\delta S}{\delta\phi_a} p\lambda\phi_a + \frac{\delta S}{\delta g_{\mu\nu}} \lambda g_{\mu\nu} \right) \\
\delta S &= \int d^n x \left(\frac{\delta S}{\delta g_{\mu\nu}} \lambda g_{\mu\nu} \right) \\
0 &= \frac{\delta S}{\delta g_{\mu\nu}} g_{\mu\nu} = -\frac{1}{2} |g|^{1/2} T^{\mu\nu} g_{\mu\nu} \\
S &= \int d^n x \frac{1}{2} |g|^{1/2} (g^{\mu\nu} \partial_\mu \phi \partial_\nu \phi - m^2 \phi^2) \\
S &= \int d^n x \frac{1}{2} |g|^{1/2} (g^{\mu\nu} \partial_\mu \phi \partial_\nu \phi - m^2 \phi^2 - \xi R \phi^2) \\
g_{\mu\nu}(x) \rightarrow \tilde{g}_{\mu\nu}(x) &= (1 + \lambda(x)) g_{\mu\nu}(x) \\
\phi(x) \rightarrow \tilde{\phi}(x) &= \left(1 - \frac{1}{2} \lambda(x)\right) \phi(x) \\
|g|^{1/2} \rightarrow |\tilde{g}|^{1/2} &= (1 + 2\lambda(x)) |g|^{1/2} \\
g^{\mu\nu} \rightarrow \tilde{g}^{\mu\nu} &= (1 - \lambda(x)) g^{\mu\nu} \\
\Gamma_{\beta\gamma}^\alpha \rightarrow \tilde{\Gamma}_{\beta\gamma}^\alpha &= \Gamma_{\beta\gamma}^\alpha + \frac{1}{2} (\delta_\gamma^\alpha \lambda_{,\beta} + \delta_\beta^\alpha \lambda_{,\gamma} - g_{\beta\gamma} \lambda^{,\alpha}), \\
R \rightarrow \tilde{R} &= (1 - \lambda)R + 3 \square \lambda(x) \\
\tilde{\mathcal{L}} &= \frac{1}{2} |\tilde{g}|^{1/2} \left(\tilde{g}^{\mu\nu} \partial_\mu \tilde{\phi} \partial_\nu \tilde{\phi} - \frac{1}{6} \tilde{R} \tilde{\phi}^2 \right) \\
&= \frac{1}{2} (1 + 2\lambda) |g|^{1/2} \left[(1 - \lambda) g^{\mu\nu} \partial_\mu ((1 - 1/2\lambda)\phi) \partial_\nu ((1 - 1/2\lambda)\phi) \right. \\
&\quad \left. - \frac{1}{6} (1 + 2\lambda) (1 - \lambda)^2 R \phi^2 - \frac{1}{2} (1 + 2\lambda) \square \lambda \phi^2 \right] \\
\tilde{\mathcal{L}} &= \frac{1}{2} |g|^{1/2} \left[g^{\mu\nu} (\partial_\mu \phi \partial_\nu \phi - \phi \partial_\mu \phi \partial_\nu \lambda) - \frac{1}{6} R \phi^2 - \frac{1}{2} \square \lambda \phi^2 \right]
\end{aligned}$$

$$\begin{aligned}
\tilde{\mathcal{L}} &= \mathcal{L} - \frac{1}{2}|g|^{1/2} \left[g^{\mu\nu} \phi \partial_\mu \phi \partial_\nu \lambda + \frac{1}{2} \square \lambda \phi^2 \right] \\
\tilde{\mathcal{L}} &= \mathcal{L} - \partial_\mu (|g|^{1/2} g^{\mu\nu} \phi^2 \partial_\nu \lambda) \\
g_{\mu\nu} &\rightarrow \tilde{g}_{\mu\nu} = \Omega^2 g_{\mu\nu} \\
\phi &\rightarrow \tilde{\phi} = \Omega^{-1} \phi \\
\tilde{\mathcal{L}} &= \mathcal{L} - \partial_\mu \left(\frac{1}{2} |g|^{1/2} g^{\mu\nu} \phi^2 \partial_\nu \log \Omega \right) \\
(\square + m^2 + \xi R) \phi &= 0. \\
T^{\mu\nu} &= \nabla^\mu \nabla^\nu \phi - \frac{1}{2} g^{\mu\nu} \nabla^\rho \phi \nabla_\rho \phi + \frac{1}{2} g^{\mu\nu} m^2 \phi^2 - \xi \left(R^{\mu\nu} - \frac{1}{2} g^{\mu\nu} R \right) \phi^2 + \\
&\quad \xi [g^{\mu\nu} \square (\phi^2) - \nabla^\mu \nabla^\nu (\phi^2)] \\
\delta g^{\mu\nu} &= -g^{\mu\rho} g^{\nu\sigma} \delta g_{\rho\sigma} \\
\delta |g|^{1/2} &= \frac{1}{2} |g|^{1/2} g^{\mu\nu} \delta g_{\mu\nu} \\
\delta R &= \delta (R_{\mu\nu} g^{\mu\nu}) = \delta R_{\mu\nu} g^{\mu\nu} + R_{\mu\nu} \delta g^{\mu\nu} = -R^{\mu\nu} \delta g_{\mu\nu} + g^{\mu\nu} \delta R_{\mu\nu} \\
\delta R_{\mu\nu} &= \delta \Gamma_{\mu\lambda;\nu}^\lambda - \delta \Gamma_{\mu\nu;\lambda}^\lambda \\
\delta \Gamma_{\mu\nu;\lambda}^\lambda &= g^{\rho\sigma} \delta g_{\rho\mu;\sigma\nu} \\
\delta R &= -R^{\mu\nu} \delta g_{\mu\nu} + g^{\rho\sigma} g^{\mu\nu} (\delta g_{\rho\sigma;\mu\nu} - \delta g_{\rho\mu;\sigma\nu}) \\
\delta S &= \frac{1}{2} \int d^n x |g|^{1/2} \left[\frac{1}{2} g^{\mu\nu} \delta g_{\mu\nu} (g^{\rho\sigma} \partial_\rho \phi \partial_\sigma \phi - m^2 \phi^2 - \xi R \phi^2) - \delta g_{\mu\nu} \nabla^\mu \phi \nabla^\nu \phi \right. \\
&\quad \left. - \xi \left(-R^{\mu\nu} \delta g_{\mu\nu} + g^{\rho\sigma} g^{\mu\nu} (\delta g_{\rho\sigma;\mu\nu} - \delta g_{\rho\mu;\sigma\nu}) \right) \phi^2 \right]. \\
&\int d^n x |g|^{1/2} g^{\rho\sigma} g^{\mu\nu} \delta g_{\rho\sigma;\mu\nu} \phi^2 = \int d^n x |g|^{1/2} g^{\rho\sigma} \delta g_{\rho\sigma} \square (\phi^2) \\
&\int d^n x |g|^{1/2} g^{\rho\sigma} g^{\mu\nu} \delta g_{\rho\mu;\sigma\nu} \phi^2 = \int d^n x |g|^{1/2} g^{\sigma\mu} g^{\lambda\nu} \delta g_{\mu\nu} \nabla_\sigma \nabla_\lambda (\phi^2) \\
\delta S &= -\frac{1}{2} \int d^n x |g|^{1/2} T^{\mu\nu} \delta g_{\mu\nu} \\
\partial^\mu \partial_\mu \phi + m^2 \phi &= 0 \\
(f_1, f_2) &= i \int dV_x (f_1^*(\vec{x}, t) \partial_0 f_2(\vec{x}, t) - \partial_0 f_1^*(\vec{x}, t) f_2(\vec{x}, t)) = i \int dV_x (f_1^* \vec{\partial}_0 f_2) \\
(f_1, f_2)^* &= (f_2, f_1) = -(f_1^*, f_2^*) \\
(f_1, f_2) &= i \int d^{n-1} x |g|^{1/2} g^{0\nu} f_1^* \vec{\partial}_\nu f_2 \\
\frac{d}{dt} (f_1, f_2) &= i \int d^{n-1} \partial_0 (|g|^{1/2} g^{0\nu} f_1^* \vec{\partial}_\nu f_2) \\
&= i \int d^{n-1} |g|^{1/2} \nabla_\mu (g^{\mu\nu} f_1^* \vec{\partial}_\nu f_2) - i \int d^{n-1} \partial_i (|g|^{1/2} g^{i\nu} f_1^* \vec{\partial}_\nu f_2) \\
\nabla_\mu (g^{\mu\nu} f_1^* \vec{\partial}_\nu f_2) &= g^{\mu\nu} \nabla_\mu (f_1^* \partial_\nu f_2 - \partial_\nu f_1^* f_2) = \\
&= g^{\mu\nu} (\partial_\mu f_1^* \partial_\nu f_2 + f_1^* \nabla_\mu \partial_\nu f_2 - \nabla_\mu \partial_\nu f_1^* f_2 - \partial_\nu f_1^* \nabla_\mu f_2) = \\
f_1^* \square f_2 - \square f_1^* f_2 &= f_1^* (-m^2 - \xi R) f_2 - f_2 (-m^2 - \xi R) f_1^* = 0 \\
\frac{d}{dt} (f_1, f_2) &= 0 \\
(f_1, f_2) &= i \int d\sigma |g|^{1/2} n^\mu f_1^* \vec{\partial}_\mu f_2 \\
(f_1, f_2)_{\sigma'} - (f_1, f_2)_\sigma &= i \int d\sigma' |g|^{1/2} n'^\mu f_1^* \vec{\partial}_\mu f_2 - i \int d\sigma |g|^{1/2} n^\mu f_1^* \vec{\partial}_\mu f_2 = \\
&\quad i \int dV_x \nabla^\mu (f_1^* \vec{\partial}_\mu f_2) = 0 \\
ds^2 &= dt^2 - a(t)^2 (dx^2 + dy^2 + dz^2) \\
a(t) &= \begin{cases} a_1, & t \rightarrow -\infty, \\ a_2, & t \rightarrow +\infty. \end{cases}
\end{aligned}$$



$$\begin{aligned}
\Box \phi &= \frac{1}{|g|^{1/2}} \partial_\mu (|g|^{1/2} \partial^\mu \phi) = 0 \\
\frac{1}{a(t)^3} \partial_t (a(t)^3 \partial_t \phi) - \frac{1}{a(t)^2} \sum_{i=1}^3 \partial_i \phi &= 0 \\
\phi &= \sum_{\vec{k}} \left(A_{\vec{k}} f_{\vec{k}}(x) + A_{\vec{k}}^\dagger f_{\vec{k}}^*(x) \right), \\
f_{\vec{k}} &= V^{-1/2} e^{i\vec{k}\vec{x}} \psi_k(\tau), \\
\tau &= \int_{t_0}^t a(t')^{-3} dt' \\
\frac{d^2 \psi_k(\tau)}{d\tau^2} + k^2 a^4 \psi_k &= 0 \\
\psi_k &\sim e^{-i\frac{k}{a_1} t} \\
f_{\vec{k}} &\sim \frac{1}{\sqrt{2V a_1^3 \omega_{1k}}} e^{i(\vec{k}\vec{x} - \omega_{1k} t)} \\
x^i &\rightarrow x^i = a_1 x^i \\
k'^i &= k^i / a_1, |\vec{k}'| = k / a_1 = \omega_{1k} \\
(f_{\vec{k}}, f_{\vec{k}'}) &= i \int d^3 x |g|^{1/2} g^{0\nu} f_{\vec{k}} \partial_\nu f_{\vec{k}'} \\
&= i \int d^3 x \frac{1}{2V (\omega_{1k} \omega_{1k'})^{1/2}} (-i) (\omega_{1k'} + \omega_{1k}) e^{i(\omega_{1k} - \omega_{1k'}) t} e^{i(\vec{k}' - \vec{k})\vec{x}} \\
&= \delta_{\vec{k}, \vec{k}'} \\
(f_{\vec{k}}, f_{\vec{k}'}^*) &= i \int d^3 x |g|^{1/2} g^{0\nu} f_{\vec{k}} \partial_\nu f_{\vec{k}'}^* \\
&= i \int d^3 x \frac{1}{2V (\omega_{1k} \omega_{1k'})^{1/2}} i (\omega_{1k'} - \omega_{1k}) e^{i(\omega_{1k} + \omega_{1k'}) t} e^{-i(\vec{k}' + \vec{k})\vec{x}} \\
&= 0 \\
[A_{\vec{k}}, A_{\vec{k}'}] &= [A_{\vec{k}}^\dagger, A_{\vec{k}'}^\dagger] = 0 \\
[A_{\vec{k}}, A_{\vec{k}'}^\dagger] &= \delta_{\vec{k}, \vec{k}'} \\
[\phi(\vec{x}, t), \phi(\vec{x}', t)] &= [\Pi(\vec{x}, t), \Pi(\vec{x}', t)] = 0 \\
i\delta^{(3)}(\vec{x} - \vec{x}') &= [\phi(\vec{x}, t), \Pi(\vec{x}', t)]. \\
[\phi(\vec{x}, t), \phi(\vec{x}', t)] &= \sum_{\vec{k}, \vec{k}'} \left((A_{\vec{k}} f_{\vec{k}}(x) + A_{\vec{k}}^\dagger f_{\vec{k}}^*(x)) (A_{\vec{k}'} f_{\vec{k}'}(x') + A_{\vec{k}'}^\dagger f_{\vec{k}'}^*(x')) - \right. \\
&\quad \left. (A_{\vec{k}'} f_{\vec{k}'}(x') + A_{\vec{k}'}^\dagger f_{\vec{k}'}^*(x')) (A_{\vec{k}} f_{\vec{k}}(x) + A_{\vec{k}}^\dagger f_{\vec{k}}^*(x)) \right) \\
&= \sum_{\vec{k}, \vec{k}'} \left([A_{\vec{k}}, A_{\vec{k}'}] f_{\vec{k}}(x) f_{\vec{k}'}(x') + [A_{\vec{k}}^\dagger, A_{\vec{k}'}^\dagger] f_{\vec{k}'}^*(x') f_{\vec{k}}^*(x) + [A_{\vec{k}}, A_{\vec{k}'}^\dagger] f_{\vec{k}}(x) f_{\vec{k}'}^*(x') \right. \\
&\quad \left. + [A_{\vec{k}}^\dagger, A_{\vec{k}'}] f_{\vec{k}'}^*(x') f_{\vec{k}}^*(x) \right) = 0 \\
\Pi &= \frac{\partial \mathcal{L}}{\partial(\partial_0 \phi)} = a^3 \partial_t \phi = \partial_t \phi \\
[\phi(\vec{x}, t), \Pi(\vec{x}', t)] &= a_1^3 \sum_{\vec{k}, \vec{k}'} \left((A_{\vec{k}} f_{\vec{k}}(x) + A_{\vec{k}}^\dagger f_{\vec{k}}^*(x)) (A_{\vec{k}'} \partial_t f_{\vec{k}'}(x') + A_{\vec{k}'}^\dagger \partial_t f_{\vec{k}'}^*(x')) - \right. \\
&\quad \left. (A_{\vec{k}'} \partial_t f_{\vec{k}'}(x') + A_{\vec{k}'}^\dagger \partial_t f_{\vec{k}'}^*(x')) (A_{\vec{k}} f_{\vec{k}}(x) + A_{\vec{k}}^\dagger f_{\vec{k}}^*(x)) \right)
\end{aligned}$$



$$\begin{aligned}
&= a_1^3 \sum_{\vec{k}, \vec{k}'} \left(\delta_{\vec{k}, \vec{k}'} f_{\vec{k}}(x) \partial_t f_{\vec{k}'}^*(x') - \delta_{\vec{k}, \vec{k}'} f_{\vec{k}}^*(x) \partial_t f_{\vec{k}'}(x') \right) \\
&= a_1^3 \sum_{\vec{k}} \left(f_{\vec{k}}(x) \partial_t f_{\vec{k}}^*(x) - f_{\vec{k}}^*(x) \partial_t f_{\vec{k}}(x) \right) \\
&= i \frac{1}{2V} \sum_{\vec{k}} \cos(\vec{k}(\vec{x} - \vec{x}')) = i \delta^{(3)}(\vec{x} - \vec{x}')
\end{aligned}$$

$$\begin{aligned}
A_{\vec{k}}|0\rangle &= 0, \forall \vec{k} \\
\psi_k^{(\pm)}(\tau) &\sim \frac{1}{(2a_2^3 \omega_{2k})^{-1/2}} e^{\mp i \omega_{2k} a_2^3 \tau} \\
\psi_k(\tau) &= \alpha_k \psi_k^{(+)}(\tau) + \beta_k \psi_k^{(-)}(\tau) \\
\psi_k(\tau) &\sim \frac{1}{(2a_2^3 \omega_{2k})^{-1/2}} (\alpha_k e^{-i \omega_{2k} a_2^3 \tau} + \beta_k e^{i \omega_{2k} a_2^3 \tau}) \\
\psi_k \partial_\tau \psi_k^* - \psi_k^* \partial_\tau \psi_k &= i \\
1 &= |\alpha_k|^2 - |\beta_k|^2 \\
f_{\vec{k}} &\sim \frac{1}{(2V a_2^3 \omega_{2k})^{1/2}} e^{i \vec{k} \vec{x}} (\alpha_k e^{-i \omega_{2k} t} + \beta_k e^{i \omega_{2k} t}) \\
\phi &= \sum_{\vec{k}} \left(a_{\vec{k}} g_{\vec{k}}(x) + a_{\vec{k}}^\dagger g_{\vec{k}}^*(x) \right), \\
g_{\vec{k}}(x) &\sim \frac{1}{\sqrt{2V a_1^3 \omega_{2k}}} e^{i(\vec{k} \vec{x} - \omega_{2k} t)}
\end{aligned}$$

$$\begin{aligned}
\phi &= \sum_{\vec{k}} \left(A_{\vec{k}} f_{\vec{k}}(x) + A_{\vec{k}}^\dagger f_{\vec{k}}^*(x) \right) \\
&= \sum_{\vec{k}} \frac{1}{(2V a_2^3 \omega_{2k})^{1/2}} \left(A_{\vec{k}} \alpha_k e^{i(\vec{k} \vec{x} - \omega_{2k} t)} + A_{\vec{k}} \beta_k e^{i(\vec{k} \vec{x} + \omega_{2k} t)} + A_{\vec{k}}^\dagger \alpha_k^* e^{-i(\vec{k} \vec{x} - \omega_{2k} t)} + \right. \\
&= A_{\vec{k}}^\dagger \beta_k^* e^{-i(\vec{k} \vec{x} + \omega_{2k} t)} \left. \right) \\
&= \sum_{\vec{k}} \left((\alpha_k A_{\vec{k}} + \beta_k^* A_{-\vec{k}}^\dagger) g_{\vec{k}}(x) + (\alpha_k^* A_{\vec{k}}^\dagger + \beta_k A_{-\vec{k}}) g_{\vec{k}}^*(x) \right)
\end{aligned}$$

$$\begin{aligned}
a_{\vec{k}} &= \alpha_k A_{\vec{k}} + \beta_k^* A_{-\vec{k}}^\dagger \\
a_{\vec{k}}^\dagger &= \alpha_k^* A_{\vec{k}}^\dagger + \beta_k A_{-\vec{k}} \\
[a_{\vec{k}}, a_{\vec{k}}^\dagger] &= (\alpha_k A_{\vec{k}} + \beta_k^* A_{-\vec{k}}^\dagger) (\alpha_k^* A_{\vec{k}}^\dagger + \beta_k A_{-\vec{k}}) - (\alpha_k^* A_{\vec{k}}^\dagger + \beta_k A_{-\vec{k}}) (\alpha_k A_{\vec{k}} + \beta_k^* A_{-\vec{k}}^\dagger) \\
&= \delta_{\vec{k}, \vec{k}'} (|\alpha_k|^2 - |\beta_k|^2) = \delta_{\vec{k}, \vec{k}'} \\
\langle N_{\vec{k}} \rangle_{t \rightarrow -\infty} &= \langle 0 | A_{\vec{k}}^\dagger A_{\vec{k}} | 0 \rangle = 0 \\
\langle N_{\vec{k}} \rangle_{t \rightarrow \infty} &= \langle 0 | a_{\vec{k}}^\dagger a_{\vec{k}} | 0 \rangle \\
&= \langle 0 | (\alpha_k^* A_{\vec{k}}^\dagger + \beta_k A_{-\vec{k}}) (\alpha_k A_{\vec{k}} + \beta_k^* A_{-\vec{k}}^\dagger) | 0 \rangle = |\beta_k|^2 \\
\Box \phi + (m^2 + \xi R) \phi &= 0 \\
\phi &= \sum_{\vec{k}} \left(A_{\vec{k}} f_{\vec{k}}(x) + A_{\vec{k}}^\dagger f_{\vec{k}}^*(x) \right) \\
f_{\vec{k}} &\sim \frac{1}{\sqrt{2V a_1^3 \omega_{1k}}} e^{i(\vec{k} \vec{x} - \omega_{1k} t)}
\end{aligned}$$



$$\phi = \sum_{\vec{k}} \left(a_{\vec{k}} g_{\vec{k}}(x) + a_{\vec{k}}^\dagger g_{\vec{k}}^*(x) \right)$$

$$g_{\vec{k}} \sim \frac{1}{\sqrt{2V a_2^3 \omega_k}} e^{i(\vec{k}\vec{x} - \omega_k t)}$$

$$a_{\vec{k}} = \alpha_k A_{\vec{k}} + \beta_k^* A_{-\vec{k}}^\dagger$$

$$a_{\vec{k}}^\dagger = \alpha_k^* A_{\vec{k}}^\dagger + \beta_k A_{-\vec{k}}$$

$$1 = |\alpha_k|^2 - |\beta_k|^2$$

$$\{A_{\vec{k}}^\dagger, A_{\vec{k}'}^\dagger\}_\pm = \{A_{\vec{k}}, A_{\vec{k}'}\}_\pm = 0$$

$$\{A_{\vec{k}}, A_{\vec{k}'}^\dagger\}_\pm = \delta_{\vec{k}, \vec{k}'}$$

$$\{a_{\vec{k}}, a_{\vec{k}'}^\dagger\}_\pm = (\alpha_k A_{\vec{k}} + \beta_k^* A_{-\vec{k}}^\dagger) (\alpha_{k'} A_{\vec{k}'}^\dagger + \beta_{k'}^* A_{-\vec{k}'}^\dagger) \pm$$

$$(\alpha_{k'} A_{\vec{k}'}^\dagger + \beta_{k'}^* A_{-\vec{k}'}^\dagger) (\alpha_k A_{\vec{k}} + \beta_k^* A_{-\vec{k}}^\dagger)$$

$$= (\alpha_k \beta_{k'}^* \pm \beta_k^* \alpha_{k'}) \delta_{\vec{k}, -\vec{k}'}$$

$$\{a_{\vec{k}}, a_{\vec{k}'}^\dagger\}_\pm = (\alpha_k A_{\vec{k}} + \beta_k^* A_{-\vec{k}}^\dagger) (\alpha_{k'}^* A_{\vec{k}'}^\dagger + \beta_{k'} A_{-\vec{k}'}) \pm$$

$$(\alpha_{k'}^* A_{\vec{k}'}^\dagger + \beta_{k'} A_{-\vec{k}'}) (\alpha_k A_{\vec{k}} + \beta_k^* A_{-\vec{k}}^\dagger)$$

$$= (|\alpha_k|^2 \pm |\beta_k|^2) \delta_{\vec{k}, \vec{k}'}$$

$$\mathcal{L} = \frac{1}{2} |g|^{1/2} \left(g^{\mu\nu} \partial_\mu \phi \partial_\nu \phi - \frac{1}{6} R \phi^2 \right)$$

$$g_{\mu\nu}(x) \rightarrow \tilde{g}_{\mu\nu}(x) = \Omega^2(x) g_{\mu\nu}(x)$$

$$\phi(x) \rightarrow \tilde{\phi}(x) = \Omega^{-1}(x) \phi(x)$$

$$\frac{\delta S}{\delta \phi} = \frac{\delta \tilde{S}}{\delta \tilde{\phi}} = \frac{\delta \tilde{S}}{\delta \tilde{\phi}} \Omega^{-1}$$

$$\left(\square + \frac{1}{6} R \right) \phi = \Omega^3 \left(\tilde{\square} + \frac{1}{6} \tilde{R} \right) \tilde{\phi}$$

$$ds^2 = dt^2 - a(t)^2 (dx^2 + dy^2 + dz^2)$$

$$\eta = \int_{-\infty}^t a^{-1}(t') dt'$$

$$ds^2 = a(t)^2 (d\eta^2 - dx^2 - dy^2 - dz^2)$$

$$\tilde{f}_{\vec{k}}(x) = \frac{1}{(2V k)^{1/2}} e^{i(\vec{k}\vec{x} - k\eta)}$$

$$f_{\vec{k}}(x) = a^{-1}(t) \tilde{f}_{\vec{k}}(x) = \frac{1}{(2V a^3(t) \omega_k(t))^{1/2}} e^{i(\vec{k}\vec{x} - \int_{-\infty}^t \omega_k(t') dt')}$$

$$\phi = \sum_{\vec{k}} \left(A_{\vec{k}} f_{\vec{k}} + A_{\vec{k}}^\dagger f_{\vec{k}}^* \right),$$

$$A_{\vec{k}} |0\rangle = 0$$

$$ds^2 = \left(1 - \frac{2M}{r} \right) dt^2 - \left(1 - \frac{2M}{r} \right)^{-1} dr^2 - r^2 d\theta^2 - r^2 \sin^2 \theta d\varphi^2$$

$$\frac{D}{D\lambda} \left(\frac{dx^\mu}{d\lambda} \right) = 0$$

$$\int_a^b \mathcal{L} d\lambda, \mathcal{L} = \frac{1}{2} g^{\mu\nu} \frac{dx^\mu}{d\lambda} \frac{dx^\nu}{d\lambda}$$

$$p_\mu = g_{\mu\nu} \frac{dx^\nu}{d\lambda} = \frac{\partial \mathcal{L}}{\partial (dx^\mu/d\lambda)}$$

$$\begin{aligned}
E &= p_t = \left(1 - \frac{2M}{r}\right) \frac{dt}{d\lambda} \\
L &= r^2 \frac{d\varphi}{d\lambda} \\
\left(1 - \frac{2M}{r}\right) \left(\frac{dt}{d\lambda}\right)^2 - \left(1 - \frac{2M}{r}\right)^{-1} \left(\frac{dr}{d\lambda}\right)^2 - r^2 \left(\frac{d\varphi}{d\lambda}\right)^2 &= 0 \\
E^2 - \left(\frac{dr}{d\lambda}\right)^2 - \frac{L^2}{r^2} \left(1 - \frac{2M}{r}\right) &= 0 \\
\frac{dr}{d\lambda} &= \pm E \\
\frac{dt}{d\lambda} \mp \left(1 - \frac{2M}{r}\right)^{-1} \frac{dr}{d\lambda} &= 0 \\
\frac{dr^*}{dr} &= \left(1 - \frac{2M}{r}\right)^{-1} \\
r^* &= r + 2M \ln(r - 2M) \\
\frac{d}{d\lambda} (t \mp r^*) &= 0 \\
\frac{du}{d\lambda} &= \frac{dt}{d\lambda} - \frac{dr^*}{d\lambda} \\
\frac{dr^*}{d\lambda} &= \frac{dr^*}{dr} \frac{dr}{d\lambda} = - \left(1 - \frac{2M}{r}\right)^{-1} E \\
\frac{dt}{d\lambda} &= \left(1 - \frac{2M}{r}\right)^{-1} E \\
\frac{du}{d\lambda} &= \frac{2}{\left(1 - \frac{2M}{r}\right)} E \\
r - 2M &= E \\
\frac{du}{d\lambda} &= 2E - \frac{4M}{\lambda} \\
u(\lambda) &= 2E\lambda - 4M \ln(\lambda/K_1)
\end{aligned}$$

$$\begin{aligned}
u(v) &= -4M \ln(\lambda/K_1) \\
v_0 - v &= K_2 \lambda \\
u(v) &= -4M \ln\left(\frac{v_0 - v}{K_1 K_2}\right) \\
\Box \phi &= \frac{1}{|g|^{1/2}} \partial_\mu (|g|^{1/2} g^{\mu\nu} \partial_\nu \phi) = 0 \\
\phi &= \int_0^\infty d\omega (a_\omega f_\omega + a_\omega^\dagger f_\omega^*) \\
(f_{\omega_1}, f_{\omega_2}) &= \delta(\omega_1 - \omega_2) \\
[\phi(\vec{x}, t), \phi(\vec{x}', t)] &= [\pi(\vec{x}, t), \pi(\vec{x}', t)] = 0 \\
[\phi(\vec{x}, t), \pi(\vec{x}', t)] &= i\delta^{(3)}(\vec{x} - \vec{x}') \\
[a_{\omega_1}, a_{\omega_2}] &= [a_{\omega_1}^\dagger, a_{\omega_2}^\dagger] = 0 \\
[a_{\omega_1}, a_{\omega_2}^\dagger] &= \delta(\omega_1 - \omega_2) \\
f_\omega &\sim \frac{1}{\sqrt{\omega r}} e^{-i\omega v} S(\theta, \varphi) \\
\Box f_\omega &= \frac{1}{1 - \frac{2M}{r}} \partial_t^2 f_\omega - \left(1 - \frac{2M}{r}\right) \frac{2}{r} \partial_r f_\omega - \left(1 - \frac{2M}{r}\right) \partial_r^2 f_\omega \\
&= -\frac{2i\omega M}{r^2 - 2Mr} = O(r^{-2}) \\
(p_{\omega_1}, p_{\omega_2}) &= \delta(\omega_1 - \omega_2)
\end{aligned}$$



$$\begin{aligned}
p_\omega &\sim \frac{1}{\sqrt{\omega r}} e^{-i\omega u} S(\theta, \varphi) \\
\Box p_\omega &= O(r^{-2}) \\
(q_{\omega_1}, q_{\omega_2}) &= \delta(\omega_1, \omega_2) \\
(q_{\omega_1}, p_{\omega_2}) &= (q_{\omega_1}, p_{\omega_2}^*) = 0 \\
\phi &= \int_0^\infty d\omega (b_\omega p_\omega + c_\omega q_\omega + b_\omega^\dagger p_\omega^* + c_\omega^\dagger q_\omega^*) \\
[b_{\omega_1}, b_{\omega_2}^\dagger] &= [c_{\omega_1}, c_{\omega_2}^\dagger] = \delta(\omega_1 - \omega_2), \\
a_\omega |0\rangle &= 0, \forall \omega \\
p_\omega &= \int_0^\infty d\omega' (\alpha_{\omega\omega'} f_{\omega'} + \beta_{\omega\omega'} f_{\omega'}^*) \\
(p_\omega, \phi) &= \left(p_\omega, \int_0^\infty d\omega' (b_{\omega'} p_{\omega'} + c_{\omega'} q_{\omega'} + b_{\omega'}^\dagger p_{\omega'}^* + c_{\omega'}^\dagger q_{\omega'}^*) \right) \\
&= \int_0^\infty d\omega' b_{\omega'} \delta(\omega - \omega') = b_\omega \\
(p_\omega, \phi) &= \left(\left(\int_0^\infty d\omega' (\alpha_{\omega\omega'} f_{\omega'} + \beta_{\omega\omega'} f_{\omega'}^*) \right), \int_0^\infty d\omega'' (a_{\omega''} f_{\omega''} + a_{\omega''}^\dagger f_{\omega''}^*) \right) \\
&= \int_0^\infty d\omega' \int_0^\infty d\omega'' (\alpha_{\omega\omega'} a_{\omega''} \delta(\omega' - \omega'') - \beta_{\omega\omega'} a_{\omega''}^\dagger \delta(\omega' - \omega'')) \\
&= \int_0^\infty d\omega' (\alpha_{\omega\omega'} a_{\omega'} - \beta_{\omega\omega'} a_{\omega'}^\dagger) \\
(p_{\omega_1}, p_{\omega_2}) &= \left(\int_0^\infty d\omega' (\alpha_{\omega_1\omega'} f_{\omega'} + \beta_{\omega_1\omega'} f_{\omega'}^*), \int_0^\infty d\omega'' (\alpha_{\omega_2\omega''} f_{\omega''} + \beta_{\omega_2\omega''} f_{\omega''}^*) \right) \\
&= \int_0^\infty (\alpha_{\omega_1\omega'}^* \alpha_{\omega_2\omega'} - \beta_{\omega_1\omega'}^* \beta_{\omega_2\omega'}) \\
(f_{\omega'}, p_\omega) &= \int_0^\infty d\omega'' (\alpha_{\omega\omega''} (f_{\omega'}, f_{\omega''}) + \beta_{\omega\omega''} (f_{\omega'}, f_{\omega''}^*)) \\
&= \int_0^\infty d\omega'' (\alpha_{\omega\omega''} \delta(\omega' - \omega'')) = \alpha_{\omega\omega'} \\
-(f_{\omega'}^*, p_\omega) &= - \int_0^\infty d\omega'' (\alpha_{\omega\omega''} (f_{\omega'}^*, f_{\omega''}) + \beta_{\omega\omega''} (f_{\omega'}^*, f_{\omega''}^*)) \\
&= - \int_0^\infty d\omega'' (\beta_{\omega\omega''} (-\delta(\omega' - \omega''))) = \beta_{\omega\omega'} \\
u(v) &= -4M \ln \left(\frac{v - v_0}{K} \right), K = K_1 K_2 > 0 \\
p_\omega &\sim \frac{1}{\sqrt{\omega}} \frac{e^{-i\omega u(v)}}{r} S(\theta, \varphi) \\
f_{\omega'} &\sim \frac{1}{\sqrt{\omega'}} \frac{e^{-i\omega' v}}{r} S(\theta, \varphi)
\end{aligned}$$



$$\alpha_{\omega\omega'} = (f_{\omega'}, p_{\omega}) = i \int dV_x \frac{1}{1 - \frac{2M}{r}} (f_{\omega'}^* \partial_t p_{\omega} - \partial_t f_{\omega'}^* p_{\omega})$$

$$= C \int_{2M}^{\infty} \frac{r^2}{1 - \frac{2M}{r}} \frac{e^{i\omega'v} e^{-i\omega u(v)}}{r^2} \left(\sqrt{\frac{\omega'}{\omega}} + \sqrt{\frac{\omega}{\omega'}} \right) dr$$

$$= -C \int_{-\infty}^{v_0} dv \sqrt{\frac{\omega'}{\omega}} e^{i\omega'v} e^{-i\omega u(v)},$$

$$\beta_{\omega\omega'} = -(f_{\omega'}^*, p_{\omega}) = C \int_{-\infty}^{v_0} dv \sqrt{\frac{\omega'}{\omega}} e^{-i\omega'v} e^{-i\omega u(v)}$$

$$\alpha_{\omega\omega'} = -C \int_{\infty}^0 ds \sqrt{\frac{\omega'}{\omega}} e^{-i\omega's} e^{i\omega'v_0} e^{i\omega 4M \ln(s/K)}$$

$$\beta_{\omega\omega'} = C \int_{-\infty}^0 ds \sqrt{\frac{\omega'}{\omega}} e^{-i\omega's} e^{-i\omega'v_0} e^{i\omega 4M \ln(-s/K)}$$

$$\alpha_{\omega\omega'} = iC \int_{-\infty}^0 ds' \sqrt{\frac{\omega'}{\omega}} e^{\omega's'} e^{i\omega'v_0} e^{i\omega 4M \log(is'/K)}$$

$$\beta_{\omega\omega'} = -iC \int_{-\infty}^0 ds' \sqrt{\frac{\omega'}{\omega}} e^{\omega's'} e^{-i\omega'v_0} e^{i\omega 4M \log(-is'/K)}$$

$$\log\left(\frac{is'}{K}\right) = \ln\left(\frac{|s'|}{K}\right) - i\frac{\pi}{2}$$

$$\log\left(\frac{-is'}{K}\right) = \ln\left(\frac{|s'|}{K}\right) + i\frac{\pi}{2}$$

$$\alpha_{\omega\omega'} = iC e^{i\omega'v_0} e^{2\omega M\pi} \int_{-\infty}^0 ds' \sqrt{\frac{\omega'}{\omega}} e^{\omega's'} e^{i\omega 4M \ln(|s'|/K)}$$

$$\beta_{\omega\omega'} = -iC e^{-i\omega'v_0} e^{-2\omega M\pi} \int_{-\infty}^0 ds' \sqrt{\frac{\omega'}{\omega}} e^{\omega's'} e^{i\omega 4M \ln(|s'|/K)}$$

$$|\alpha_{\omega\omega'}|^2 = e^{8\pi\omega M} |\beta_{\omega\omega'}|^2$$

$$(p_{\omega_1}, p_{\omega_2}) = \Gamma(\omega_1) \delta(\omega_1 - \omega_2)$$

$$(p_{\omega_1}, p_{\omega_2}) = (p_{\omega_1}^{(1)}, p_{\omega_2}^{(1)}) + (p_{\omega_1}^{(2)}, p_{\omega_2}^{(2)})$$

$$(p_{\omega_1}^{(1)}, p_{\omega_2}^{(1)}) = (1 - \Gamma(\omega_1)) \delta(\omega_1 - \omega_2)$$

$$\Gamma(\omega_1) \delta(\omega_1 - \omega_2) = \int_0^{\infty} d\omega' (\alpha_{\omega_1\omega'}^* \alpha_{\omega_2\omega'} - \beta_{\omega_1\omega'}^* \beta_{\omega_2\omega'})$$

$$b_{\omega} = (p_{\omega}^{(2)}, \phi) = \int_0^{\infty} d\omega' (\alpha_{\omega\omega'} a_{\omega'} + \beta_{\omega\omega'} a_{\omega'}^{\dagger})$$

$$\langle N \rangle = \langle 0 | b_{\omega}^{\dagger} b_{\omega} | 0 \rangle = \int_0^{\infty} d\omega' \beta_{\omega\omega'} \langle \omega' | \int_0^{\infty} d\omega'' \beta_{\omega\omega''}^* | \omega'' \rangle = \int_0^{\infty} d\omega' |\beta_{\omega\omega'}|^2$$

$$\Gamma(\omega) \delta(0) = \int_0^{\infty} d\omega' (|\alpha_{\omega\omega'}|^2 - |\beta_{\omega\omega'}|^2) = (e^{8\pi M\omega} - 1) \int_0^{\infty} d\omega' |\beta_{\omega\omega'}|^2$$

$$\delta(\omega_1 - \omega_2) = \lim_{T \rightarrow \infty} \frac{1}{2\pi} \int_{-T/2}^{T/2} dt e^{i(\omega_1 - \omega_2)t}$$



$$\Gamma(\omega) \lim_{T \rightarrow \infty} \frac{T}{2\pi} = (e^{8\pi M\omega} - 1) \int_0^\infty d\omega' |\beta_{\omega\omega'}|^2$$

$$\langle N \rangle = \lim_{T \rightarrow \infty} \frac{T}{2\pi} \frac{\Gamma(\omega)}{\Gamma(\omega)} \frac{1}{e^{8\pi M\omega} - 1}$$

$$T = \frac{1}{8\pi k_B M} = \frac{\kappa}{2\pi}$$

$$p = \int d\omega A(\omega) e^{i\gamma(\omega)} p_\omega$$

$$T = \frac{\hbar c^3}{8\pi G M k_B} \approx 10^{-7} \left(\frac{M_\odot}{M} \right) K$$

$$\frac{dE}{dt} = 4\pi r_s^2 \sigma T^4$$

$$\frac{dM}{dt} = -\beta \frac{m_p^3}{t_p} \frac{1}{M^2},$$

$$M(t) = \left(M_0^3 - 3\beta \frac{m_p^3}{t_p} t \right)^{1/3}$$

$$\Delta t = \frac{t_p}{3} \beta \left(\frac{M_0}{m_p} \right)^3$$

3. Supersimetrías, agujeros negros de mecánica cuántica y supermembranas – Modelo de Supergravedad Cuántica. Debido a la permeabilidad del campo gravitónico en espacios infinitos e indiscriminados, las superpartículas comportan la relación entre bosones y fermiones, a propósito de la interacción con el gravitón, el cual, como se ha dicho también, puede ser una partícula supersimétrica o compañera de otra partícula – gauge de naturaleza fermiónica, todo esto, sin perjuicio del efecto cuántico gravitacional implícito respecto de aquellas partículas subatómicas supermasivas, en las que, por la distorsión del espacio – tiempo cuántico, crean supersimetrías. En este punto, es importante precisar que los campos cuánticos, no son irreducibles, es decir, que coexisten y se afectan entre sí, pues, bosones y fermiones son en sí y a propósito de la supersimetría, partículas compañeras que se afectan unas a otras, según el espectro gravitacional involucrado.

Ahora bien, cuando una superpartícula, deforma el espacio – tiempo cuántico, sea por efecto cuántico gravitacional extrínseco (campo gravitónico – gravitón) o intrínseco (en relación a las partículas supermasivas), según sea el caso, produce diversos fenómenos físicos, desde microagujeros o agujeros negros cuánticos, que por su naturaleza son pequeños y que resultan de la colisión de partículas, producto de la distorsión geométrica del espacio – tiempo cuántico, así como supermembranas, es decir, distintas capas dimensionales que se desdoblán desde la singularidad de la curvatura y que dan lugar a diferentes estados transicionales de la materia, lo que equivale a la supergravedad a escala cuántica.

El modelo matemático, comporta:

La métrica de un gravitón está dada por G_{MN} con $M, N = 0, 1, \dots, D - 1$. En un espacio de D dimensiones, resulta en $\frac{1}{2}(D - 1)(D - 2) - 1$. Para lograr la supersimetría, el gravitón tiene un compañero, esto es, el gravitino $\Psi_{M,\alpha}$, la cual es una partícula – gauge, las cuales interactúan en superespacios. Los cálculos de los espinores, calibre, escalares, tensores y vectores en un superespacio – tiempo infinito, de gauge y curvado bajo parámetros de transformación supersimétricos, por acoplamiento del gravitón y el gravitino, en un campo de gauge gravitónico, comprenden:

$$(\Gamma^M \Psi_M)_\alpha = 0$$

Dimensión espacio-tiempo D	Espinor dimensión \tilde{D}	Gravitón $\frac{1}{2}(D - 1)(D - 2) - 1$	Gravitino $\frac{1}{2}(D - 3)\tilde{D}$	Diferencia
------------------------------	-------------------------------	--	---	------------



4	4	2	2	0
5	8	5	8	3
6	8	9	12	3
7	16	14	32	18
8	16	20	40	20
9	16	27	48	21
10	16	35	56	21
11	32	44	128	84
12	64	54	288	234

Ahora bien, esta relación híbrida, que da como resultado, la dinámica del campo gravitónico, matemáticamente comporta lo que sigue, siguiendo el sistema de gauge de Yang – Mills – Einstein – Newton – Dirac – Transformaciones de Supergravedad y reparación de gauge.

$$S_{11d} = \frac{1}{16\pi\mathcal{G}_{(11)}} \left[\int d^{11}x \sqrt{\|G\|} \left(R - \frac{1}{2}|G|^2 + \frac{1}{6} \int C \wedge G \wedge G - \frac{i}{2} \bar{\psi}_M \Gamma^{MNP} D_N \left(\frac{\omega + \hat{\omega}}{2} \right) \psi_P \right. \right. \\ \left. \left. - \frac{i}{384} (\bar{\psi}_M \Gamma^{MNABCD} \psi_N + 12 \bar{\psi}^A \Gamma^{BC} \psi^D) (G_{ABCD} + \hat{G}_{ABCD}) \right) \right]$$

$$16\pi\mathcal{G}_{(11)} = \frac{(2\pi\ell_p)^9}{2\pi}$$

$$D_M(\omega)\psi_N \equiv \partial_M \psi_N - \frac{1}{4} \omega_{MAB} \Gamma^{AB} \psi_N$$

$$\omega_{MAB} = \omega_{MAB}^{(0)}(E) + \frac{i}{16} [\bar{\psi}_N \Gamma_{MAB}{}^{NP} \psi_P - 2(\bar{\psi}_M \Gamma_B \psi_A - \bar{\psi}_M \Gamma_A \psi_B + \bar{\psi}_B \Gamma_M \psi_A)]$$

$$\hat{\omega}_{MAB} = \omega_{MAB} - \frac{i}{16} \bar{\psi}_N \Gamma_{MAB}{}^{NP} \psi_P.$$

$$G_{LMNP} \equiv 4\partial_{[L} C_{MNP]}$$

$$\hat{G}_{LMNP} \equiv G_{LMNP} + \frac{3}{2} i \bar{\psi}_L \Gamma_{MN} \psi_P.$$

$$|G|^2 = |dC|^2 \equiv \frac{1}{4!} G_{LMNP} G^{LMNP}$$

$$\delta E_M^A = \frac{i}{2} \bar{\epsilon} \Gamma^A \Psi_M$$

$$\delta C_{MNP} = -\frac{3}{2} i \bar{\epsilon} \Gamma_{[MN} \Psi_{P]}$$

$$\delta C = d\Lambda$$

$$G_{\mu\nu}^{(11)} = G_{\mu\nu}^{(10)} + e^{2\gamma} A_\mu A_\nu, G_{\mu 10}^{(11)} = e^{2\gamma} A_\mu, G_{1010}^{(11)} = e^{2\gamma}$$

$$C_{\mu\nu\rho}^{(11)} = C_{\mu\nu\rho}^{(10)}, C_{\mu\nu 10}^{(11)} = B_{\mu\nu}.$$

$$\left\| \begin{matrix} A & B \\ C & D \end{matrix} \right\| = \|A - B^T D^{-1} C\| \|D\|$$

$$\sqrt{\|G^{(11)}\|} = \sqrt{\|G^{(10)}\|} e^\gamma$$

$$\frac{2\pi}{(2\pi\ell_p)^9} \int d^{11}x \sqrt{\|G^{(11)}\|} R \rightarrow \frac{(2\pi)^2 R_{10}}{(2\pi\ell_p)^9} \int d^{10}x \sqrt{\|G^{(10)}\|} \left(e^\gamma \left(R - \frac{1}{2} |d\gamma|^2 \right) - \frac{1}{2} e^{3\gamma} |dA|^2 \right)$$



$$\begin{aligned}
& -\frac{2\pi}{(2\pi\ell_p)^9} \frac{1}{2} \int d^{11}x \sqrt{\|G^{(11)}\| \|G\|^2} \rightarrow -\frac{(2\pi)^2 R_{10}}{(2\pi\ell_p)^9} \frac{1}{2} \int d^{10}x \sqrt{\|G^{(10)}\|} (e^\gamma |dC^{(10)}|^2 + e^{-\gamma} |dB|^2) + \dots \\
& -\frac{2\pi}{(2\pi\ell_p)^9} \frac{1}{6} \int C \wedge G \wedge G \rightarrow -\frac{(2\pi)^2 R_{10}}{(2\pi\ell_p)^9} \frac{1}{2} \int B \wedge dC^{(10)} \wedge dC^{(10)} \\
S_{IIA} = & \frac{2\pi}{(2\pi\ell_s)^8} \int d^{10}x \sqrt{\|G^{(10)}\|} \left[e^{-2\Phi} \left(R + |d\Phi|^2 - \frac{1}{2} |dB|^2 \right) - \left(\frac{1}{2} |dA|^2 + \frac{1}{2} |dC|^2 \right) \right] \\
& -\frac{2\pi}{(2\pi\ell_s)^8} \frac{1}{2} \int B \wedge dC \wedge dC + \dots,
\end{aligned}$$

$$\begin{aligned}
G_{\mu\nu} & \rightarrow e^{-\gamma} G_{\mu\nu}, \quad \Phi = \frac{3}{2} \gamma, \\
G_{\mu\nu}^{(11)} & = e^{-\gamma} G_{\mu\nu}^{(10)} = e^{-\frac{2}{3}\Phi} G_{\mu\nu}^{(10)}. \\
\frac{L}{\ell_p} & = e^{-\frac{1}{2}\gamma} \frac{L}{\ell_s} = e^{-\frac{1}{3}\Phi} \frac{L}{\ell_s}.
\end{aligned}$$

$$\begin{aligned}
e^{(\Phi)} & = g_s \\
\ell_p & = g_s^{\frac{1}{3}} \ell_s \\
\frac{(2\pi)^2 R_{10}}{(2\pi\ell_p)^9} & = \frac{2\pi}{(2\pi\ell_s)^8} \frac{1}{g_s^2},
\end{aligned}$$

$$R_{10} = g_s \ell_s \int A_M dX^M$$

$$\int A_M(X(\tau)) \frac{dX^M}{d\tau} d\tau = \int A_\tau d\tau$$

$$A_\tau \equiv A_M \frac{dX^M}{d\tau}$$

$$\begin{aligned}
\int B_{MN} dX^M \wedge dX^N & = \int B_{\mu\nu} d\xi^\mu \wedge d\xi^\nu \\
B_{\mu\nu} & \equiv B_{MN} \frac{dX^M}{d\xi^\mu} \frac{dX^N}{d\xi^\nu}
\end{aligned}$$

$$\int C_{MNP} dX^M \wedge dX^N \wedge dX^P = \int C_{\mu\nu\lambda} d\xi^\mu \wedge d\xi^\nu \wedge d\xi^\lambda = \frac{1}{6} \int d^3\xi \epsilon^{\mu\nu\lambda} C_{\mu\nu\lambda}$$

$$C_{\mu\nu\lambda} \equiv C_{MNP} \frac{dX^M}{d\xi^\mu} \frac{dX^N}{d\xi^\nu} \frac{dX^P}{d\xi^\lambda}$$

$$S_{M2}^{\text{bosonic}} = \int d^3\xi \left(\frac{1}{2} \sqrt{|g|} g^{\mu\nu} G_{MN} \partial_\mu X^M \partial_\nu X^N - \frac{1}{2} \sqrt{|g|} + \frac{1}{6} \epsilon^{\mu\nu\lambda} C_{MNP} \partial_\mu X^M \partial_\nu X^N \partial_\lambda X^P \right)$$

$$\begin{aligned}
g_{\mu\nu} & = G_{MN} \partial_\mu X^M \partial_\nu X^N \\
\Pi_\mu^M & \equiv \partial_\mu X^M - i \bar{\theta} \Gamma^M \partial_\mu \theta
\end{aligned}$$

$$\begin{aligned}
S_{M2}^{\text{susy}} & = \int d^3\xi \left(\frac{1}{2} \sqrt{|g|} g^{\mu\nu} \Pi_\mu^M \Pi_\nu^N - \frac{1}{2} \sqrt{|g|} \right. \\
& \quad \left. + \frac{i}{2} \epsilon^{\mu\nu\lambda} \bar{\theta} \Gamma_{MN} \partial_\mu \theta \left[\Pi_\nu^M \Pi_\lambda^N + i \Pi_\nu^M \bar{\theta} \Gamma^N \partial_\lambda \theta - \frac{1}{3} \bar{\theta} \Gamma^M \partial_\nu \theta \bar{\theta} \Gamma^N \partial_\lambda \theta \right] \right).
\end{aligned}$$

$$\delta X^M = -i \bar{\theta} \Gamma^M \epsilon$$

$$\delta \theta = \epsilon$$

$$\delta g_{\mu\nu} = 0$$

$$\Pi_\mu \equiv \Pi_\mu^M \Gamma_M$$

$$\tau^\mu \equiv \frac{1}{2\sqrt{|g|}} \epsilon^{\mu\nu\lambda} \Pi_\nu \Pi_\lambda$$

$$\Gamma \equiv \frac{1}{6\sqrt{|g|}} \epsilon^{\mu\nu\lambda} \Pi_\mu \Pi_\nu \Pi_\lambda$$



$$\begin{aligned}\delta X^M &= i\bar{\theta}\Gamma^M(1+\Gamma)\kappa \\ \delta\theta &= (1+\Gamma)\kappa \\ \delta(\sqrt{|g|}g^{\mu\nu}) &= ig^{\sigma(\mu}\epsilon^{\nu)\lambda\rho}\bar{\kappa}(1+\Gamma)\partial_\sigma\theta\Pi_\lambda\Pi_\rho \\ &\quad + \frac{2i}{3\sqrt{|g|}}\epsilon^{\sigma\tau(\mu}\epsilon^{\nu)\lambda\rho}\bar{\kappa}\Pi^\alpha\partial_\alpha\theta(\Pi_\sigma^M\Pi_{\lambda M}\Pi_\tau^N\Pi_{\rho N} + \Pi_\sigma^M\Pi_{\lambda M}g_{\tau\rho} + g_{\sigma\lambda}g_{\tau\rho})\end{aligned}$$

$$\begin{aligned}g_{\mu\nu} &= \Pi_{\mu M}\Pi_\nu^M \\ \Gamma^2 &= 1 \\ \tau^\mu &= g^{\mu\nu}\Pi_\nu\Gamma = g^{\mu\nu}\Gamma\Pi_\nu \\ \{\tau^\mu, \tau^\nu\} &= 2g^{\mu\nu}.\end{aligned}$$

$$A^M \equiv \partial_\mu \left\{ \sqrt{|g|}g^{\mu\nu}\Pi_\nu^M - i\epsilon^{\mu\nu\lambda}(\bar{\theta}\Gamma^{MN}\partial_\nu\theta) \left(\Pi_{\lambda N} + \frac{i}{2}\bar{\theta}\Gamma_N\partial_\lambda\theta \right) \right\} = 0$$

$$(1-\Gamma)g^{\mu\nu}\Pi_\mu^M\Gamma_M\partial_\nu\theta = 0$$

$$B = (1-\Gamma)\tau^\mu\partial_\mu\theta = 0$$

$$\Pi_\mu^MA_M = -2i\sqrt{|g|}\partial_\mu\bar{\theta}B$$

$$X^M \rightarrow (X^\mu, X^I), \mu = 0, 1, 2; I = 3, 4, \dots, 10,$$

$$X^\mu = \xi^\mu, \mu = 0, 1, 2.$$

$$\delta X^M = \eta^\nu(\xi)\partial_\nu X^M + \Lambda_N^M X^N$$

$$\delta X^\mu = \eta^\mu + \Lambda_\nu^\mu \xi^\nu + \Lambda_I^\mu X^I$$

$$\eta^\mu = -\Lambda_\nu^\mu \xi^\nu - \Lambda_I^\mu X^I$$

$$\delta X^I = -(\Lambda_\nu^\mu \xi^\nu + \Lambda_I^\mu X^I)\partial_\mu X^I + \Lambda_J^I X^J$$

$$\xi^\mu \rightarrow \xi^\mu + \ell_\nu^\mu \xi^\nu$$

$$\delta\phi = \ell_\nu^\mu \xi^\nu \partial_\mu\phi$$

$$\delta\theta = -\Lambda^\mu{}_\nu \xi^\nu \partial_\mu\theta + \frac{1}{4}\Lambda_{\mu\nu}\Gamma^{\mu\nu}\theta + \frac{1}{4}\Lambda_{IJ}\Gamma^{IJ}\theta,$$

$$(1+\Gamma^*)\theta = 0$$

$$\Gamma^* \equiv \Gamma^1\Gamma^2 \dots \Gamma^8$$

$$S_{M2, \text{ bosonic}}^{\text{static gauge}} = -T_{M2} \int d^3\xi \sqrt{-\det\left(\eta_{\mu\nu} + \frac{1}{T_{M2}}\partial_\mu X^I\partial_\nu X^I\right)} \sim -\frac{1}{2} \int \partial_\mu X^I \partial^\mu X^I + \frac{1}{T_{M2}} O(\partial X)^4 + \dots,$$

$$\delta\Psi_M \equiv D_M(\omega)\epsilon - \frac{i}{288}(\Gamma_M^{PQRS} + 8\Gamma^{QRS}\delta_M^P)G_{PQRS}\epsilon = 0$$

$$G_7 = \star G_4 - \frac{1}{2}C_3 \wedge G_4$$

$$ds^2 = f_{(1)}(r)dy^\mu dy_\mu + f_{(2)}(r)dx^I dx^I$$

$$G_{012r} = f_{(3)}(r),$$

$$r = \sqrt{(x^1)^2 + (x^2)^2 + \dots + (x^8)^2}$$

$$H_{M2}(r) = 1 + \frac{(r_{M2})^6}{r^6}$$

$$f_{(1)}(r) = H_{M2}(r)^{-\frac{2}{3}}$$

$$f_{(2)}(r) = H_{M2}(r)^{\frac{1}{3}}$$

$$f_{(3)}(r) = -\frac{\partial}{\partial r}(H_{M2}(r)^{-1})$$

$$G_{J_1 J_2 \dots J_7} = 6(r_{M2})^6 \epsilon_{J_1 J_2 \dots J_7} \frac{x^I}{r^8}$$

$$G_{\theta_1 \theta_2 \dots \theta_7} = 6(r_{M2})^6 \epsilon_{\theta_1 \theta_2 \dots \theta_7}$$

$$Q^{(e)} = 6(r_{M2})^6 \Omega_7 = 2\pi^4 (r_{M2})^6,$$

$$R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R = 8\pi\mathcal{G}_{(D)}T_{\mu\nu}$$

$$R_{00} = \frac{D-3}{D-2} 8\pi\mathcal{G}_{(D)}\rho$$



$$\begin{aligned}
\nabla^2 g_{00} &= -2 \frac{D-3}{D-2} 8\pi \mathcal{G}_{(D)} \rho \\
\nabla^2 \phi &= 4\pi \mathcal{G}_{(D)} \rho \\
g_{00} &= -\left(1 + 4 \frac{D-3}{D-2} \phi\right) \\
\phi(x) &= -\frac{4\pi \mathcal{G}_{(D)} M}{(D-3)\Omega_{D-2} r^{D-3}} \\
g_{00} &= -\left(1 - \frac{16\pi \mathcal{G}_{(D)} M}{(D-2)\Omega_{D-2} r^{D-3}}\right) \\
R_{00} &= \frac{D-p-3}{D-2} 8\pi \mathcal{G}_{(D)} \rho \\
g_{00} &= -\left(1 + 4 \frac{D-p-3}{D-2} \phi\right) \\
\phi(x) &= -\frac{4\pi \mathcal{G}_{(D)} T_p}{(D-p-3)\Omega_{D-p-2} r^{D-p-3}} \\
g_{00} &= -\left(1 - \frac{16\pi \mathcal{G}_{(D)} T_p}{(D-2)\Omega_{D-p-2} r^{D-p-3}}\right) \\
(r_{M2})^6 &= \frac{8\pi \mathcal{G}_{(11)}}{3 \Omega_7} n_2 T_{M2} \\
(r_{M2})^6 &= 128\pi^4 n_2 \ell_p^9 T_{M2} \\
ds^2 &= g_{(1)}(r) dy^\mu dy_\mu + g_{(2)}(r) dx^I dx^I \\
r &= \sqrt{(x^1)^2 + (x^2)^2 + \dots + (x^5)^2} \\
H_{M5}(r) &= 1 + \frac{(r_{M5})^3}{r^3} \\
g_{(1)}(r) &= H_{M5}(r)^{-\frac{1}{3}} \\
g_{(2)}(r) &= H_{M5}(r)^{\frac{2}{3}} \\
G_{\theta_1 \theta_2 \theta_3 \theta_4} &= 3(r_{M5})^3 \epsilon_{\theta_1 \theta_2 \theta_3 \theta_4} \\
Q^{(m)} &= \int_{S^4} G = 3(r_{M5})^3 \Omega_4 = 8\pi^2 (r_{M5})^3 \\
(r_{M5})^3 &= 32\pi^6 n_5 \ell_p^9 T_{M5} \\
T_{M2} &\sim \frac{1}{\ell_p^3}, T_{M5} \sim \frac{1}{\ell_p^6} \\
\frac{1}{16\pi \mathcal{G}_{(11)}} Q^{(e)} Q^{(m)} &= 2\pi n, \\
\frac{2\pi}{(2\pi \ell_p)^9} Q^{(e)} Q^{(m)} &= (2\pi)^8 \ell_p^9 T_{M2} T_{M5} \\
T_{M2} T_{M5} &= \frac{(2\pi)^2}{(2\pi \ell_p)^9} \\
T_{M2} &= \frac{2\pi}{(2\pi \ell_p)^3}, T_{M5} = \frac{2\pi}{(2\pi \ell_p)^6} \\
T_p &= \frac{1}{g_s} \frac{2\pi}{(2\pi \ell_s)^{p+1}}. \\
T_{F1} &= \frac{2\pi}{(2\pi \ell_s)^2}, T_{NS5} = \frac{1}{g_s^2} \frac{2\pi}{(2\pi \ell_s)^6}. \\
T_{M2} &= \frac{1}{4\pi^2 \ell_p^3}
\end{aligned}$$



$$\begin{aligned}
T_{M2}^{\text{wrapped}} &= T_{M2} \times 2\pi R_{10} \\
&= \frac{1}{4\pi^2 g_s \ell_s^3} \times 2\pi g_s \ell_s \\
&= \frac{1}{2\pi \ell_s^2}, \\
T_{M2} &= \frac{1}{4\pi^2 \ell_p^3} \\
&= \frac{1}{4\pi^2 \left(g_s^{\frac{1}{3}} \ell_s\right)^3} \\
&= \frac{1}{g_s} \frac{1}{4\pi^2 \ell_s^3} \\
&= T_{D2}.
\end{aligned}$$

$$\begin{aligned}
T_{M5} &= \frac{1}{32\pi^5 \ell_p^6} \\
T_{M5}^{\text{wrapped}} &= T_{M5} \times 2\pi R_{10} \\
&= \frac{g_s \ell_s}{16\pi^4 g_s^2 \ell_s^6} \\
&= \frac{1}{g_s} \frac{2\pi}{(2\pi \ell_s)^5} \\
&= T_{D4}, \\
T_{M5} &= \frac{1}{32\pi^5 \ell_p^6} \\
&= \frac{1}{g_s^2} \frac{1}{32\pi^5 \ell_s^6} \\
&= \frac{1}{g_s^2} \frac{2\pi}{(2\pi \ell_s)^6} \\
&= T_{NS5},
\end{aligned}$$

$$\begin{aligned}
T_0 &= \frac{1}{g_s \ell_s} = \frac{1}{R_{10}}. \\
p &= \frac{2\pi n}{L}.
\end{aligned}$$

$$E^2 = p_1^2 + \dots + p_9^2 + p_{10}^2$$

$$ds_{\text{Taub-NUT}}^2 = U(\vec{x}) d\vec{x} \cdot d\vec{x} + \frac{1}{U(\vec{x})} (dy + \vec{A} \cdot d\vec{x})^2$$

$$\vec{B} = \vec{\nabla} \times \vec{A}$$

$$\vec{\nabla} U = -\vec{B}.$$

$$U(\vec{x}) = 1 + \frac{R}{2r}$$

$$\vec{B} = \frac{R}{2} \frac{\vec{x}}{r^3}$$

$$\frac{R}{2r} dr^2 + \frac{2r}{R} dy^2$$

$$\tilde{r} = \sqrt{2rR}$$

$$d\tilde{r}^2 + \frac{\tilde{r}^2}{R^2} dy^2$$



$$\begin{aligned}
T_{KK6} &= \frac{2\pi}{(2\pi\ell_p)^9} \times 2\pi R_{10} \int d^3x \vec{\nabla}^2 U \\
&= \frac{2\pi}{(2\pi\ell_p)^9} \times (2\pi R_{10})^2 \\
&= \frac{1}{g_s} \frac{2\pi}{(2\pi\ell_s)^7} = T_{D6} \\
&\quad g_s \rightarrow \frac{1}{g_s}, \ell_s \rightarrow \sqrt{g_s} \ell_s \\
T_{p,q} &= \frac{1}{2\pi\ell_s^2} \sqrt{p^2 + \frac{q^2}{g_s^2}} \\
g_s(\text{IIB}) &= \frac{R_{10}}{R_9}, \ell_s(\text{IIB}) = \sqrt{\frac{\ell_p^3}{R_{10}}} \\
T_{M2}^{\text{wrapped}} &= T_{M2} \sqrt{p(2\pi R_{10})^2 + q(2\pi R_9)^2} \\
&= \frac{1}{2\pi\ell_s^2} \sqrt{p^2 + \frac{q^2}{g_s^2}} = T_{p,q} \\
S_{M2}^{\text{bosonic}} &= -\frac{1}{(2\pi)^2 \ell_p^3} \int d^3\xi \sqrt{-\det(\eta_{\mu\nu} + (2\pi)^2 \ell_p^3 \partial_\mu X^I \partial_\nu X^I)} \\
S_{D2}^{\text{bosonic}} &= -\frac{1}{(2\pi\alpha')^2 g_{YM}^2} \int d^3\xi \sqrt{-\det(\eta_{\mu\nu} + (2\pi\alpha')^2 \partial_\mu X^i \partial_\nu X^i + 2\pi\alpha' F_{\mu\nu})} \\
&\quad g_{YM}^2 = \frac{g_s}{\alpha'} \\
&\quad \frac{1}{g_{YM}^2} \left(-\frac{1}{2} \partial_\mu X^i \partial^\mu X^i - \frac{1}{4} F_{\mu\nu} F^{\mu\nu} \right) \\
\mathcal{L} &= \frac{1}{2} \varepsilon^{\mu\nu\lambda} B_\mu F_{\nu\lambda} - \frac{1}{(2\pi\alpha')^2 g_{YM}^2} \sqrt{-\det(\eta_{\mu\nu} + (2\pi\alpha')^2 \partial_\mu X^i \partial_\nu X^i + (2\pi\alpha')^2 g_{YM}^4 B_\mu B_\nu)} \\
&\quad B_\mu \rightarrow -\frac{1}{g_{YM}} \partial_\mu X^8 \\
\mathcal{L} &= -\frac{1}{(2\pi)^2 \ell_p^3} \sqrt{-\det(\eta_{\mu\nu} + (2\pi)^2 \ell_p^3 \partial_\mu X^I \partial_\nu X^I)} \\
&\quad X^8 \sim X^8 + 2\pi g_{YM} \\
S_{M2}^{\text{susy}} \stackrel{\ell_p \rightarrow 0}{=} &\int d^3\xi \left(-\frac{1}{2} \partial_\mu X^I \partial^\mu X^I + \frac{i}{2} \bar{\psi}^A \gamma^\mu \partial_\mu \psi^A \right) \\
S_{D2}^{\text{susy}} \stackrel{\ell_s \rightarrow 0}{=} &\frac{1}{g_{YM}^2} \int d^3\xi \left(-\frac{1}{2} \partial_\mu X^i \partial^\mu X^i - \frac{1}{4} F_{\mu\nu} F^{\mu\nu} + \frac{i}{2} \bar{\psi}^A \gamma^\mu \partial_\mu \psi^A \right) \\
S_{nD2} \stackrel{\ell_s \rightarrow 0}{=} &\frac{1}{g_{YM}^2} \int d^3\xi \text{Tr} \left(-\frac{1}{2} D_\mu \mathbf{X}^i D^\mu \mathbf{X}^i + \frac{1}{4} [\mathbf{X}^i, \mathbf{X}^j]^2 - \frac{1}{4} \mathbf{F}_{\mu\nu} \mathbf{F}^{\mu\nu} + \frac{i}{2} \boldsymbol{\psi}^A \gamma^\mu D_\mu \boldsymbol{\psi}^A - \boldsymbol{\psi}^A \Gamma_{AB}^i [\mathbf{X}^i, \boldsymbol{\psi}^B] \right) \\
&\quad D_\mu \mathbf{X}^i = \partial_\mu \mathbf{X}^i - i[\mathbf{A}_\mu, \mathbf{X}^i] \\
&\quad \mathbf{F}_{\mu\nu} = \partial_\mu \mathbf{A}_\nu - \partial_\nu \mathbf{A}_\mu - i[\mathbf{A}_\mu, \mathbf{A}_\nu] \\
&\quad X^9 = \frac{N}{2r}, F_{\theta\phi} = -r^2 \partial_r X^9, \\
&\quad \frac{1}{2\pi} \int F = N \\
\frac{\partial \mathbf{X}^i}{\partial x^9} &= \pm \frac{i}{2} \epsilon^{ijk} [\mathbf{X}^j, \mathbf{X}^k], i, j, k \in 1, 2, 3 \\
&\quad [\boldsymbol{\alpha}^i, \boldsymbol{\alpha}^j] = 2i \epsilon^{ijk} \boldsymbol{\alpha}^k
\end{aligned}$$

$$\begin{aligned}
X^i &= \pm \frac{1}{2x^9} \alpha^i, i = 1, 2, 3 \\
R^2 &= \frac{(2\pi\alpha')^2}{N} \text{Tr} \sum_{i=1,2,3} (X^i)^2 \\
\sum_{i=1,2,3} (\alpha^i)^2 &= N^2 - 1, \\
R &= \frac{\pi\alpha'\sqrt{N^2 - 1}}{x^9} \\
X^{10} &\sim \frac{N}{r^2} \\
\mathcal{R}^+ &= \left(\frac{n+1}{4}, \frac{n-1}{4}\right), \mathcal{R}^- = \left(\frac{n-1}{4}, \frac{n+1}{4}\right). \\
G_5 &\equiv \mathcal{P}_{\mathcal{R}^+} - \mathcal{P}_{\mathcal{R}^-} \\
[A_1, A_2, A_3, A_4] &\equiv \sum_{\text{permutations } \sigma} \text{sign}(\sigma) A_{\sigma(1)} A_{\sigma(2)} A_{\sigma(3)} A_{\sigma(4)} \\
\frac{\partial X^i}{\partial x^{10}} &= \frac{1}{4!} \frac{b}{8\pi\ell_p^3} \epsilon^{ijkl} [G_5, X^j, X^k, X^l] \\
\rho_s(\Gamma^i) &\equiv \mathbb{1} \otimes \dots \otimes \Gamma^i \otimes \dots \otimes \mathbb{1}, s = 1, 2, \dots, n \\
\sum_{s=1}^n \rho_s(\Gamma^i) &= (\Gamma^i \otimes \mathbb{1} \otimes \dots \otimes \mathbb{1}) + (\mathbb{1} \otimes \Gamma^i \otimes \dots \otimes \mathbb{1}) + \dots + (\mathbb{1} \otimes \dots \otimes \mathbb{1} \otimes \Gamma^i)
\end{aligned}$$

$$\mathbf{G}^i = \mathcal{P}_{\mathcal{R}^+} \sum_{s=1}^n \rho_s(\Gamma^i P_-) \mathcal{P}_{\mathcal{R}^-} + \mathcal{P}_{\mathcal{R}^-} \sum_{s=1}^n \rho_s(\Gamma^i P_+) \mathcal{P}_{\mathcal{R}^+}$$

$$\begin{aligned}
X^i(x^{10}) &= i\hat{R}(x^{10})\mathbf{G}^i. \\
\epsilon^{ijkl} \mathbf{G}_5 \mathbf{G}^i \mathbf{G}^j \mathbf{G}^k &= -2(n+2)\mathbf{G}^i,
\end{aligned}$$

$$\hat{R}(x^{10}) = \sqrt{\frac{2\pi\ell_p^3}{(n+2)bx^{10}}}.$$

$$R^2 = \frac{1}{N} \left| \text{Tr} \sum_{i=1}^4 (X^i)^2 \right|.$$

$$R = \sqrt{N} |\hat{R}|.$$

$$x^{10} = \frac{2\pi\ell_p^3 N}{(n+2)bR^2}.$$

$$E = T_{M2} \int d^2\sigma \text{Tr} \left[\left(\frac{dX^i}{dx^{10}} + \frac{b}{8\pi\ell_p^3} \epsilon^{ijkl} \mathbf{G}_5 X^j X^k X^l \right)^2 + \left(1 - \frac{b}{16\pi\ell_p^3} \epsilon^{ijkl} \left\{ \frac{dX^i}{dx^{10}}, \mathbf{G}_5 X^j X^k X^l \right\} \right)^2 \right]^{\frac{1}{2}}$$

$$E|_{BH} = T_{M2} \int d^2\sigma \left(1 - \frac{b}{16\pi\ell_p^3} \epsilon^{ijkl} \left\{ \frac{dX^i}{dx^{10}}, \mathbf{G}_5 X^j X^k X^l \right\} \right)$$

$$E = NT_{M2}L \int dx^{10} + T_{M5}L \int 2\pi^2 R^3 dR$$

$$\mathbf{H}^{KLM} \equiv [X^K, X^L, X^M] \equiv \{[X^K, X^L], X^M\} + \{[X^L, X^M], X^K\} + \{[X^M, X^K], X^L\}$$

$$S = -T_{M2} \int d^3\sigma \text{Tr} \left[1 + (\partial_a X^M)^2 - \frac{b^2}{12} (\mathbf{H}^{KLM})^2 + \frac{b^2}{48} [\partial_a X^{[K}, \mathbf{H}^{LMN}]^2 \right]^{\frac{1}{2}}$$

$$\mathcal{L}_{CS} = \frac{k}{4\pi} \text{Tr} \left(A \wedge dA - \frac{2i}{3} A \wedge A \wedge A \right)$$

$$\mathcal{L}_{\mathcal{N}=1} = \frac{k}{4\pi} \text{Tr} \left(A \wedge dA - \frac{2i}{3} A \wedge A \wedge A - i\bar{\chi}\chi \right)$$



$$\begin{aligned}
\delta A_\mu &= i\bar{\epsilon}\gamma_\mu\chi \\
\delta\chi &= -\frac{1}{2}\gamma^{\mu\nu}F_{\mu\nu}\epsilon. \\
\mathcal{L}_{\mathcal{N}=1}^{\text{matter}} &= -\frac{1}{2}\partial_\mu\phi^a\partial^\mu\phi^a + \frac{i}{2}\bar{\psi}^a\gamma^\mu\partial_\mu\psi^a + \frac{1}{2}C^aC^a + t_{abcd}\phi^a\phi^b\left(\frac{1}{3}\phi^cC^d - \frac{1}{2}\bar{\psi}^c\psi^d\right) \\
\delta\phi^a &= i\bar{\epsilon}\psi^a \\
\delta\psi^a &= -(\gamma^\mu\partial_\mu\phi^a - C^a)\epsilon \\
\delta C^a &= -i\bar{\epsilon}\gamma^\mu\partial_\mu\psi^a \\
\partial_\mu\phi^a &\rightarrow \partial_\mu\phi^a - iA_\mu^I(T^I)_{ab}\phi^b \\
&\quad \phi^a\bar{\chi}^IT_{ab}^I\psi^b \\
\mathcal{L}_{\mathcal{N}=2} &= \frac{k}{4\pi}\text{Tr}\left(\mathbf{A}\wedge d\mathbf{A} - \frac{2i}{3}\mathbf{A}\wedge\mathbf{A}\wedge\mathbf{A} - i\bar{\chi}\chi + 2\mathbf{D}\sigma\right) \\
\delta A_\mu &= \frac{i}{2}(\bar{\epsilon}\gamma_\mu\chi - \bar{\chi}\gamma_\mu\epsilon) \\
\delta\sigma &= -\frac{1}{2}(\bar{\epsilon}\chi - \bar{\chi}\epsilon) \\
\delta\mathbf{D} &= \frac{1}{2}(\bar{\epsilon}\gamma^\mu D_\mu\chi + D_\mu\bar{\chi}\gamma^\mu\epsilon) - \frac{1}{2}(\bar{\epsilon}\chi, \sigma] + [\bar{\chi}, \sigma]\epsilon \\
\delta\chi &= \left(-\frac{1}{2}\gamma^{\mu\nu}F_{\mu\nu} + i\gamma^\mu D_\mu\sigma - i\mathbf{D}\right)\epsilon. \\
-\frac{1}{2}\partial_\mu\phi_A\partial^\mu\phi^A + i\bar{\psi}_A\gamma^\mu\partial_\mu\psi^A + F_A F^A + (F^A W_{,A} + c.c.) \\
\delta\phi^A &= i\bar{\epsilon}\psi^A \\
\delta\psi^A &= -\gamma^\mu\partial_\mu\phi^A + F^A\epsilon^* \\
\delta F^A &= -i\bar{\epsilon}^*\gamma^\mu\partial_\mu\psi^A \\
-\sigma^I\sigma^J(\phi_A T^I T^J \phi^A) + D^I(\phi_A T^I \phi^A) - i\sigma^I(\bar{\psi}_A T^I \psi^A) - \phi_A\bar{\chi}^I T^I \psi^A + \phi^A T^I \bar{\psi}_A \chi^I \\
\mathcal{L}_{\mathcal{N}=2, \text{gauged}} &= \mathcal{L}_{CS} + \mathcal{L}_{\text{kinetic}} + \mathcal{L}_{\text{scalar-fermion}} - V(\phi), \\
\mathcal{L}_{\text{scalar-fermion}} &= -\frac{4\pi i}{k}(\phi_A T^I \phi^A)(\bar{\psi}_B T^I \psi^B) - \frac{8\pi i}{k}(\bar{\psi}_A T^I \phi^A)(\phi_B T^I \psi^B) \\
V(\phi) &= \frac{16\pi^2}{k^2}(\phi_A T^I \phi^A)(\phi_B T^J \phi^B)(\phi_C T^I T^J \phi^C) \\
W(Q, \tilde{Q}) &= \alpha(\tilde{Q}T^I Q)(\tilde{Q}T^I Q) \\
\Gamma_{012}\epsilon &= \epsilon \\
X^I &= X_a^I T^a \\
\Psi &= \Psi_a T^a. \\
\delta X_d^I &= i\bar{\epsilon}\Gamma^I\Psi_d \\
\delta\Psi_d &= \partial_\mu X_d^I \Gamma^\mu \Gamma^I \epsilon - \frac{1}{3!}X_a^I X_b^J X_c^K f_d^{abc} \Gamma^{IJK} \epsilon \\
[T^a, T^b, T^c] &= f^{abc}{}_d T^d \\
[\delta_1, \delta_2]X_d^I &= -2i\bar{\epsilon}_2\Gamma^\mu\epsilon_1\partial_\mu X_d^I - (2i\bar{\epsilon}_2\Gamma^{JK}\epsilon_1 X_a^J X_b^K f_d^{abc})X_d^I \\
\delta X_d^I &= \tilde{\Lambda}_d^c X_c^I, \tilde{\Lambda}_d^c = -2i\bar{\epsilon}_2\Gamma^{JK}\epsilon_1 X_a^J X_b^K f_d^{abc} \\
\delta X^I &= \alpha_{JK}[X^I, X^J, X^K] \\
\delta X &= [X, A, B] \\
\delta[X, Y, Z] &= [\delta X, Y, Z] + [X, \delta Y, Z] + [X, Y, \delta Z] \\
[A, B, [X, Y, Z]] &= [[A, B, X], Y, Z] + [X, [A, B, Y], Z] + [X, Y, [A, B, Z]] \\
f^{abc}{}_g f^{efg}{}_d &= f^{efa}{}_g f^{gbc}{}_d + f^{agc}{}_d f^{efb}{}_g + f^{abg}{}_d f^{efc}{}_g \\
D_\mu X_d^I &= \partial_\mu X_d^I - \tilde{A}_\mu^c X_c^I \\
\delta\tilde{A}_\mu^c{}_d &= \partial_\mu\tilde{\Lambda}^c{}_d + \tilde{A}_\mu^c{}_e \tilde{\Lambda}^e{}_d - \tilde{\Lambda}^c{}_e \tilde{A}_\mu^e{}_d
\end{aligned}$$



$$\begin{aligned}
\tilde{F}_{\mu\nu}{}^a{}_b &= \partial_\nu \tilde{A}_\mu{}^a{}_b - \partial_\mu \tilde{A}_\nu{}^a{}_b - \tilde{A}_\mu{}^a{}_c \tilde{A}_\nu{}^c{}_b + \tilde{A}_\nu{}^a{}_c \tilde{A}_\mu{}^c{}_b \\
\delta X_a^I &= i\bar{\epsilon} \Gamma^I \Psi_a \\
\delta \Psi_a &= D_\mu X_a^I \Gamma^\mu \Gamma^I \epsilon - \frac{1}{6} X_b^I X_c^J X_d^K f^{abcd} \Gamma^{IJK} \epsilon \\
\delta \tilde{A}_\mu{}^b{}_a &= i\bar{\epsilon} \Gamma_\mu \Gamma_I X_c^I \Psi_a f^{cdb}{}_a \\
[\delta_1, \delta_2] X_a^I &= v^\lambda \partial_\lambda X_a^I + (\tilde{\Lambda}^b{}_a - v^\lambda \tilde{A}_\lambda{}^b{}_a) X_b^I \\
[\delta_1, \delta_2] \Psi_a &= v^\lambda \partial_\lambda \Psi_a + (\tilde{\Lambda}^b{}_a - v^\lambda \tilde{A}_\lambda{}^b{}_a) \Psi_b \\
[\delta_1, \delta_2] \tilde{A}_\mu{}^b{}_a &= v^\lambda \partial_\lambda \tilde{A}_\mu{}^b{}_a + \tilde{D}_\mu (\tilde{\Lambda}^b{}_a - v^\lambda \tilde{A}_\lambda{}^b{}_a), \\
&\Gamma^\mu D_\mu \Psi_a + \frac{1}{2} \Gamma_{IJ} X_c^I X_d^J \Psi_b f^{cdb}{}_a = 0 \\
\tilde{F}_{\mu\nu}{}^b{}_a + \varepsilon_{\mu\nu\lambda} \left(X_c^J D^\lambda X_d^I + \frac{i}{2} \tilde{\Psi}_c \Gamma^\lambda \Psi_d \right) f^{cdb}{}_a &= 0. \\
D^2 X_a^I - \frac{i}{2} \tilde{\Psi}_c \Gamma^I X_d^J \Psi_b f^{cdb}{}_a + \frac{1}{2} X_b^J X_c^K X_e^I X_f^J X_g^K f^{abcd}{}_a f^{efg}{}_d &= 0 \\
\langle X, Y \rangle &= h^{ab} X_a Y_b \\
f^{abcd} &= f^{[abcd]} \\
\mathcal{L} &= -\frac{1}{2} D_\mu X^{aI} D^\mu X_a^I + \frac{i}{2} \tilde{\Psi}^a \Gamma^\mu D_\mu \Psi_a + \frac{i}{4} \tilde{\Psi}_b \Gamma_{IJ} X_c^I X_d^J \Psi_a f^{abcd} - V + \mathcal{L}_{CS} \\
V &= \frac{1}{12} X_a^I X_b^J X_c^K X_e^I X_f^J X_g^K f^{abcd} f^{efg}{}_d \\
\mathcal{L}_{CS} &= \frac{1}{2} \varepsilon^{\mu\nu\lambda} \left(f^{abcd} A_{\mu ab} \partial_\nu A_{\lambda cd} + \frac{2}{3} f^{cda}{}_g f^{efgb} A_{\mu ab} A_{\nu cd} A_{\lambda ef} \right), \\
\tilde{A}_\mu{}^b{}_a &= A_{\mu cd} f^{cdb}{}_a \\
f^{abcd} &= \frac{2\pi}{k} \varepsilon^{abcd}, h^{ab} = \delta^{ab} \\
\tilde{A}_\mu{}^a{}_b &= \tilde{A}_\mu^{+a}{}_b + \tilde{A}_\mu^{-a}{}_b \\
\mathcal{L}_{CS} &= \frac{k}{8\pi} \varepsilon^{\mu\nu\lambda} \left(\tilde{A}_\mu^{+a}{}_b \partial_\nu \tilde{A}_\lambda^{+b}{}_a + \frac{2}{3} \tilde{A}_\mu^{+a}{}_b \tilde{A}_\nu^{+b}{}_c \tilde{A}_\lambda^{+c}{}_a \right) - \frac{k}{8\pi} \varepsilon^{\mu\nu\lambda} \left(\tilde{A}_\mu^{-a}{}_b \partial_\nu \tilde{A}_\lambda^{-b}{}_a + \frac{2}{3} \tilde{A}_\mu^{-a}{}_b \tilde{A}_\nu^{-b}{}_c \tilde{A}_\lambda^{-c}{}_a \right) \\
A_\mu^L &= \frac{1}{2} \varepsilon_i{}^{jk} \sigma_k \tilde{A}_\mu^{+i}{}_j, A_\mu^R = \frac{1}{2} \varepsilon_i{}^{jk} \sigma_k \tilde{A}_\mu^{-i}{}_j \\
\frac{k}{4\pi} \varepsilon^{\mu\nu\lambda} \text{Tr} \left[\left(A_\mu^L \partial_\nu A_\lambda^L - \frac{2i}{3} A_\mu^L A_\nu^L A_\lambda^L \right) - \left(A_\mu^R \partial_\nu A_\lambda^R - \frac{2i}{3} A_\mu^R A_\nu^R A_\lambda^R \right) \right] & \\
X_{\alpha\beta}^I &= \left(\frac{1}{2} X_4^I \mathbb{1} + \frac{i}{2} X_i^I \sigma^i \right)_{\alpha\beta}, i = 1, 2, 3 \\
X^{\dagger I \alpha \beta} &= \varepsilon^{\alpha\dot{\alpha}} \varepsilon^{\beta\dot{\beta}} X_{\dot{\beta}\dot{\alpha}}^I \\
D_\mu X^I &= \partial_\mu X^I - i A_\mu^L X^I + i X^I A_\mu^R \\
V(X) &= \frac{8}{3} \text{Tr} (X^{\dagger I} X^{\dagger J} X^{\dagger K} X^{\dagger L} X^{\dagger M} X^{\dagger N}) \\
\mathcal{L}_{su(2)} &= \frac{k}{4\pi} \varepsilon^{\mu\nu\lambda} \text{Tr} \left(A_\mu \partial_\nu A_\lambda - \frac{2i}{3} A_\mu A_\nu A_\lambda \right) \\
\int d^3x \mathcal{L}_{su(n)} &\rightarrow \int d^3x \mathcal{L}_{su(n)} + 2\pi k w \\
\delta Z_d^A &= i\bar{\epsilon}^{AB} \psi_{Bd} \\
\delta \psi_{Bd} &= \gamma^\mu D_\mu Z_d^A \epsilon_{AB} + f^{ab}{}_{cd} Z_a^C Z_b^A \bar{Z}_C^c \epsilon_{AB} + f^{ab}{}_{cd} Z_a^C Z_b^D \bar{Z}_B^c \epsilon_{CD} \\
\delta \tilde{A}_\mu{}^c{}_d &= -i\bar{\epsilon}_{AB} \gamma_\mu Z_a^A \psi^{Bb} f^{ca}{}_{bd} + i\bar{\epsilon}^{AB} \gamma_\mu \bar{Z}_a^A \psi_{Bb} f^{ca}{}_{bd}, \\
[\delta_1, \delta_2] Z_d^A &= v^\lambda D_\lambda Z_d^A + \tilde{\Lambda}_d^a Z_a^A \\
v^\lambda &= \frac{i}{2} \bar{\epsilon}_2^{CD} \gamma^\lambda \epsilon_{1CD}, \\
\tilde{\Lambda}_d^a &= \Lambda^c{}_b f^{ab}{}_{cd}, \Lambda^c{}_b = i(\bar{\epsilon}_2^{DE} \epsilon_{1CE} - \bar{\epsilon}_1^{DE} \epsilon_{2CE}) \bar{Z}_D^c Z_b^c, \\
\gamma^\mu D_\mu \psi_{cd} &= f^{ab}{}_{cd} \psi_{ca} Z_b^D \bar{Z}_D^c + 2f^{ab}{}_{cd} \psi_{Da} Z_b^D \bar{Z}_C^c - \varepsilon_{CDEF} f^{ab}{}_{cd} \psi^{Dc} Z_a^E Z_b^F \\
\tilde{F}_{\mu\nu}{}^c{}_d &= -\varepsilon_{\mu\nu\lambda} (D^\lambda Z_a^A \bar{Z}_A^b - Z_a^A D^\lambda \bar{Z}_A^b - i\tilde{\Psi}^{Ab} \gamma^\lambda \psi_{Aa}) f^{ca}{}_{bd}
\end{aligned}$$



$$\begin{aligned}
& [T^a, T^b; \bar{T}_c] = f^{ab}{}_{cd} T^d \\
& f^{ef}{}_{gb} f^{cb}{}_{ad} + f^{fe}{}_{ab} f^{cb}{}_{gd} + f^{ga}{}_{fb} f^{ce}{}_{bd} + f^{ag}{}_{eb} f^{cf}{}_{bd} = 0. \\
& f_{ab}^{*cd} = (f^{ab}{}_{cd})^* \\
& \Lambda_a^{*b} = (\Lambda^a{}_b)^* \\
& \langle \bar{X}, Y \rangle = \bar{X}^a Y_a \\
& \delta_\Lambda \bar{Z}_A^d = \tilde{\Lambda}^{*d}{}_{Aa} \bar{Z}_A^a \\
& \tilde{\Lambda}_a^{*b} = (\tilde{\Lambda}^a{}_b)^* = \Lambda^*{}^c{}_d f_{ac}{}^{*db}. \\
& \tilde{\Lambda}_b^{*a} = -\tilde{\Lambda}^a{}_b. \\
& f_{cd}^{ab} = f_{cd}^{*ab} \\
& f^{ge}{}_{fd} f^{ab}{}_{cg} = f^{ae}{}_{fg} f^{gb}{}_{cd} + f^{be}{}_{fg} f^{ag}{}_{cd} - f_{cf}^{*eg} f^{ab}{}_{gd} \\
\mathcal{L} = & -D^\mu \bar{Z}_A^a D_\mu Z_A^a - i\bar{\psi}^{Aa} \gamma^\mu D_\mu \psi_{Aa} - V + \mathcal{L}_{CS} \\
& -if^{ab}{}_{cd} \bar{\psi}^{Ad} \psi_{Aa} Z_b^B \bar{Z}_B^c + 2if^{ab}{}_{cd} \bar{\psi}^{Ad} \psi_{Ba} Z_b^B \bar{Z}_B^c \\
& + \frac{i}{2} \epsilon_{ABCD} f^{ab}{}_{cd} \bar{\psi}^{Ad} \psi^{Bc} Z_a^C Z_b^D - \frac{i}{2} \epsilon^{ABCD} f^{cd}{}_{ab} \bar{\psi}_{Ac} \psi_{Bd} \bar{Z}_C^a \bar{Z}_D^b. \\
V = & \frac{2}{3} Y_{Bd}^{CD} \bar{Y}_{CD}^{Bd} \\
Y_{Bd}^{CD} = & f^{ab}{}_{cd} Z_a^C Z_b^D \bar{Z}_B^c - \frac{1}{2} \delta_B^C f^{ab}{}_{cd} Z_a^E Z_b^D \bar{Z}_E^c + \frac{1}{2} \delta_B^D f^{ab}{}_{cd} Z_a^E Z_b^C \bar{Z}_E^c \\
\mathcal{L}_{CS} = & \frac{1}{2} \epsilon^{\mu\nu\lambda} \left(f^{ab}{}_{cd} A_\mu{}^c{}_b \partial_\nu A_\lambda{}^d{}_a + \frac{2}{3} f^{ac}{}_{dg} f^{ge}{}_{fb} A_\mu{}^b{}_a A_\nu{}^d{}_c A_\lambda{}^f{}_e \right) \\
& \frac{\delta \mathcal{L}_{CS}}{\delta A_\lambda{}^b{}_a} f_{ab}^{ac} = \frac{1}{2} \epsilon^{\lambda\mu\nu} \tilde{F}_{\mu\nu}{}^c{}_d \\
[X, Y; Z] = & -\frac{2\pi}{k} (XZ^\dagger Y - YZ^\dagger X) \\
\delta X_{dl} = [X, Y; Z]_{dl} = & f^{ajib}{}_{ckdl} \Lambda^{ck}{}_{bj} X_{ai} \\
= & -\frac{2\pi}{k} (X_{ak} Z^{+kb} Y_{bl} - Y_{dk} Z^{+kb} X_{bl}). \\
f^{ajib}{}_{ckdl} = & -\frac{2\pi}{k} (\delta^a{}_d \delta^b{}_c \delta^i{}_k \delta^j{}_l - \delta^a{}_c \delta^b{}_d \delta^i{}_l \delta^j{}_k) \\
\delta X_{dl} = \tilde{\Lambda}_{dl}^{ai} X_{ai} = & -\frac{2\pi}{k} (\delta_l^i \Lambda_{dj}^{aj} - \delta_d^a \Lambda_{bl}^{bi}) X_{ai} \\
f^{ab}{}_{cd} = & -\frac{2\pi}{k} (J^{ab} J_{cd} + (\delta^a{}_c \delta^b{}_d - \delta^a{}_d \delta^b{}_c)) \\
\delta X_d = \tilde{\Lambda}_d^a X_a = & -\frac{2\pi}{k} [(\Lambda_d^a + \Lambda_d^a) - \delta_d^a \Lambda_b^b] X_a. \\
Z_d^A = & Z_{\alpha\dot{\alpha}}^A \bar{\sigma}_d^{\dot{\alpha}\alpha} \\
(\bar{\sigma}^a \sigma^b \bar{\sigma}^c - \bar{\sigma}^c \sigma^b \bar{\sigma}^a)^{\dot{\alpha}\alpha} = & -2\epsilon^{abcd} \bar{\sigma}_d^{\dot{\alpha}\alpha} \\
\delta Z_d^A = & i\bar{\epsilon}^{AD} \Psi_{Dd} + i\bar{\eta} \Psi_d^A \\
\delta \Psi_D^d = & \gamma^\mu \epsilon_{AD} D_\mu Z^{Ad} + \gamma^\mu \eta D_\mu \bar{Z}_D^d \\
& + \epsilon^{abcd} Z_a^A Z_b^B \bar{Z}_{Dc} \epsilon_{AB} - \epsilon^{abcd} Z_a^A Z_b^B \bar{Z}_{Bc} \epsilon_{AD} \\
& - \epsilon^{abcd} Z_a^A \bar{Z}_{Ab} \bar{Z}_{Dc} \eta - \frac{1}{3} \epsilon_{ABCD} \epsilon^{abcd} \eta^* Z_a^A Z_b^B Z_c^C \\
[\delta_1, \delta_2] \Psi_{Dd} = & v^\mu D_\mu \Psi_{Dd} + \tilde{\Lambda}_d^a \Psi_{Da} \\
& + \frac{i}{2} \epsilon_{[2}^{CB} \epsilon_{1]CD} E_{Bd} - \frac{i}{4} \epsilon_2^{BE} \gamma^\mu \epsilon_{1BE} \gamma_\mu E_{Dd} \\
& + i\bar{\eta}_{[2} \epsilon_{1]CD} E_d^C - \frac{i}{2} (\bar{\eta}_{[2} \eta_{1]}^* + \bar{\eta}_{[2} \gamma^\mu \eta_{1]} \gamma_\mu) E_{Dd} \\
\delta \tilde{A}_\mu^{ad} = & -i\epsilon^{abcd} \bar{\epsilon}_{BC} \gamma_\mu \Psi_b^B Z_c^C - i\epsilon^{abcd} \bar{\epsilon}^{BC} \gamma_\mu \Psi_{Bb} \bar{Z}_{Cc} \\
& + i\epsilon^{abcd} \bar{\eta}^* \gamma_\mu \Psi_{Bb} Z_c^B + i\epsilon^{abcd} \bar{\eta} \gamma_\mu \Psi_b^B \bar{Z}_{Bc}
\end{aligned}$$



$$\begin{aligned}
\delta Z_a^A &= g^{AB} \bar{Z}_{Ba} \\
\delta \Psi_{Ba} &= -\frac{1}{2} \epsilon_{BCDE} g^{DE} \Psi_a^C \\
\delta \epsilon^{AB} &= g^{AB} \eta^* + \frac{1}{2} \epsilon^{ABCD} g_{CD}^* \eta \\
\delta \eta &= -\frac{1}{2} g^{AB} \epsilon_{AB}, \\
D_\mu Z_a^A &= \partial_\mu Z_a^A - \tilde{A}_\mu^b{}_a Z_b^A - i B_\mu \delta_a^b Z_b^A \\
&\quad B_\mu \rightarrow B_\mu + \partial_\mu \theta \\
&\quad \delta B_\mu = 0 \\
\delta \mathcal{L}_{\text{SU}(n) \times \text{SU}(n)}^{\text{gauged}} &= -\frac{1}{2} G_{\mu\nu} \bar{\epsilon}_{AB} \gamma^{\mu\nu} \psi^{Aa} Z_a^B + \frac{1}{2} G_{\mu\nu} \bar{\epsilon}^{AB} \gamma^{\mu\nu} \psi_{Aa} \bar{Z}_B^a \\
&= -\frac{1}{2} \epsilon^{\mu\nu\lambda} G_{\mu\nu} \bar{\epsilon}_{AB} \gamma_\lambda \psi^{Aa} Z_a^B + \frac{1}{2} \epsilon^{\mu\nu\lambda} \bar{\epsilon}^{AB} G_{\mu\nu} \bar{\epsilon} \gamma_\lambda \psi_{Aa} \bar{Z}_B^a, \\
\mathcal{L}_{\text{U}(n) \times \text{U}(n)} &= \mathcal{L}_{\text{SU}(n) \times \text{SU}(n)}^{\text{gauged}} + \frac{k'}{8\pi} \epsilon^{\mu\nu\lambda} G_{\mu\nu} Q_\lambda \\
\delta Q_\lambda &= \frac{4\pi}{k'} \bar{\epsilon}_{AB} \gamma_\lambda \psi^{Aa} Z_a^B - \frac{4\pi}{k'} \bar{\epsilon}^{AB} \gamma_\lambda \psi_{Aa} \bar{Z}_B^a \\
[\delta_1, \delta_2] B_\mu &= v^\nu G_{\nu\mu} v^\nu = \frac{i}{2} (\bar{\epsilon}_2^{CD} \gamma^\nu \epsilon_{CD}^1) \\
&\quad H_{\mu\nu} = \partial_\mu Q_\nu - \partial_\nu Q_\mu \\
[\delta_1, \delta_2] Q_\mu &= \frac{k'}{4\pi} v^\nu \epsilon_{\mu\nu\lambda} (i Z_a^A D^\lambda \bar{Z}_A^a - i D^\lambda Z_a^A \bar{Z}_A^a - \bar{\psi}_a^A \gamma^\lambda \psi_a^A) + D_\mu \Lambda \\
H_{\mu\nu} &= -\frac{k'}{4\pi} \epsilon_{\mu\nu\lambda} (i Z_a^A D^\lambda \bar{Z}_A^a - i D^\lambda Z_a^A \bar{Z}_A^a - \bar{\psi}_a^A \gamma^\lambda \psi_a^A) \\
[\delta_1, \delta_2] Q_\mu &= v^\nu H_{\nu\mu} + D_\mu \Lambda \\
\mathcal{L}_{\text{U}(1) \times \text{U}(1) \text{CS}} &= \frac{k'}{4\pi} \epsilon^{\mu\nu\lambda} A_\mu^L \partial_\nu A_\lambda^L - \frac{k'}{4\pi} \epsilon^{\mu\nu\lambda} A_\mu^R \partial_\nu A_\lambda^R \\
&\quad k' = nk, \\
\delta A_\lambda^R &= \delta A_\lambda^L = \frac{2\pi}{nk} \bar{\epsilon}_{AB} \gamma_\lambda \psi^{Aa} Z_a^B - \frac{2\pi}{nk} \bar{\epsilon}^{AB} \gamma_\lambda \psi_{Aa} \bar{Z}_B^a \\
[\cdot, \cdot, \cdot] &: \mathcal{V} \otimes \mathcal{V} \otimes \mathcal{V} \rightarrow \mathcal{V} \\
[\cdot, \cdot, \cdot] &: \mathcal{V} \otimes \mathcal{V} \otimes \bar{\mathcal{V}} \rightarrow \mathcal{V} \\
[\cdot, \cdot, \cdot] &: \bar{\mathcal{V}} \otimes \bar{\mathcal{V}} \otimes \mathcal{V} \rightarrow \bar{\mathcal{V}} \\
\varphi_{U, \bar{V}}(X) &= [X, U; \bar{V}] \quad \varphi_{U, \bar{V}}(\bar{X}) = -[\bar{X}, \bar{V}; U] \\
\varphi_{U, \bar{V}}([X, Y; \bar{Z}]) &= [\varphi_{U, \bar{V}}(X), Y; \bar{Z}] + [X, \varphi_{U, \bar{V}}(Y); \bar{Z}] + [X, Y; \varphi_{U, \bar{V}}(\bar{Z})] \\
[[X, Y; \bar{Z}], U; \bar{V}] &= [[X, U; \bar{V}], Y; \bar{Z}] + [X, [Y, U; \bar{V}]; \bar{Z}] - [X, Y; [\bar{Z}, \bar{V}, U]] \\
\langle [X, \bar{V}; U], Y \rangle &= \langle X, [Y, U; \bar{V}] \rangle \\
[\varphi_{U, \bar{V}}, \varphi_{Y, \bar{Z}}](X) &= \varphi_{\varphi_{U, \bar{V}}(Y), Z}(X) - \varphi_{Y, \varphi_{U, \bar{V}}(\bar{Z})}(X) \\
(\varphi_{Y, \bar{Z}}, \varphi_{U, \bar{V}}) &= \langle \bar{Z}, [Y, U; \bar{V}] \rangle \\
\varphi &: \mathcal{V} \times \bar{\mathcal{V}} \rightarrow \mathcal{G}, \\
\varphi^*(g)_{U, \bar{V}} &= \langle \bar{V}, g(U) \rangle, \\
\varphi^*(g)_{U, \bar{V}} &= (\varphi_{U, \bar{V}}, g) \\
[W, U; \bar{V}] &= \varphi_{U, \bar{V}}(W). \\
[g, \varphi_{U, \bar{V}}] &= \varphi_{g(U), \bar{V}} + \varphi_{U, g(\bar{V})} \\
\varphi_{U, \bar{V}}(g)_r &= h_{rs} (T^s)^a{}_b U_a \bar{V}_e g^{be} \\
f^{abc}{}_d &= (T^r)^b{}_e (T^s)^a{}_d h_{rs} g^{ce}. \\
f^{abc}{}_d &= \sum_R c_R (T^r)^b{}_e (T^s)^a{}_d k_{rs}^R g^{ce} \\
[X, Y, Z] &= [[X, Y], Z] \\
\int d^4\theta Q^\dagger e^V Q, \int d^4\theta \tilde{Q}^\dagger e^{-V} \tilde{Q}
\end{aligned}$$

$$\begin{aligned}
ds^2 &= U dx^a dx_a + U^{-1} (d\phi + \omega^a dx_a)^2 \\
U &= U_\infty + \frac{q}{2|\vec{x}|} \\
U^{ij} \partial_{(i)}^a \partial_{(j)a} U_{kl} &= 0 \\
\partial_{(j)}^a \omega_{ki}^b - \partial_{(k)}^b \omega_{ji}^a &= \epsilon^{abc} \partial_{(j)c} U_{ki} \\
\{x_a^{(1)}\} &= (x^3, x^4, x^5), \{x_a^{(2)}\} = (x^7, x^8, x^9) \\
\tau &= -\frac{U_{12}}{U_{11}} + i \frac{\sqrt{\det U}}{U_{11}} \\
ds^2 &= U_{ij} dx_a^{(i)} dx_a^{(j)} + U^{ij} (d\phi_i + \omega_{ik}^a dx_a^{(k)}) (d\phi_j + \omega_{jl}^b dx_b^{(l)}) \\
U_1 &= U_\infty + \begin{pmatrix} h_1 & 0 \\ 0 & 0 \end{pmatrix}, h_1 = \frac{1}{2|\vec{x}^{(1)}|}, U_\infty = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \\
U_2' &= U_\infty' + \begin{pmatrix} h_2 & kh_2 \\ kh_2 & k^2 h_2 \end{pmatrix}, h_2 = \frac{1}{2|\vec{x}^{(1)} + k\vec{x}^{(2)}|}, U_\infty' = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \\
G_j^i &= \begin{pmatrix} 1 & 0 \\ -k^{-1} & k^{-1} \end{pmatrix} \\
U_2 &= U_\infty + \begin{pmatrix} 0 & 0 \\ 0 & h_2 \end{pmatrix}, h_2 = \frac{1}{2|\vec{x}^{(2)}|}, U_\infty = k^{-2} \begin{pmatrix} k^2 + 1 & -1 \\ -1 & 1 \end{pmatrix} \\
ds^2 &= \sum_{i=1,2} \left(\frac{1}{2|\vec{x}^{(i)}|} d\vec{x}^{(i)} \cdot d\vec{x}^{(i)} + 2|\vec{x}^{(i)}| \left(d\phi_i' + \frac{1}{2} \cos \vartheta_i d\phi_i \right)^2 \right) \\
ds^2 &= \sum_{i=1,2} \left(dr_i^2 + r_i^2 \left(d\phi_i' + \frac{1}{2} \cos \vartheta_i d\phi_i \right)^2 + \frac{r_i^2}{4} (d\vartheta_i^2 + \sin^2 \vartheta_i d\phi_i^2) \right) \\
z^1 &= r_1 \cos \frac{\vartheta_1}{2} e^{-i\phi_1' - \frac{i}{2}\varphi_1} \quad z^2 = r_2 \cos \frac{\vartheta_2}{2} e^{i\phi_2' + \frac{i}{2}\varphi_2} \\
z^3 &= r_1 \sin \frac{\vartheta_1}{2} e^{-i\phi_1' + \frac{i}{2}\varphi_1} \quad z^4 = r_2 \sin \frac{\vartheta_2}{2} e^{i\phi_2' - \frac{i}{2}\varphi_2} \\
ds^2 &= \sum_{A=1}^4 |dz^A|^2 \\
&\quad e^{\pi(\Gamma_{34} + \Gamma_{56} + \Gamma_{78} + \Gamma_{910})/k} \\
&\quad (\Gamma_{34} + \Gamma_{56} + \Gamma_{78} + \Gamma_{910}) \epsilon = 0 \\
\mathcal{L} &= -\text{Tr}(D^\mu \bar{Z}_A, D_\mu Z^A) - i \text{Tr}(\bar{\psi}^A, \gamma^\mu D_\mu \psi_A) - V + \mathcal{L}_{CS} \\
&\quad + \frac{2\pi i}{k} \text{Tr}(\bar{\psi}^A \psi_A \bar{Z}_B Z^B - \bar{\psi}^A Z^B \bar{Z}_B \psi_A) - \frac{4\pi i}{k} \text{Tr}(\bar{\psi}^A \psi_B \bar{Z}_A Z^B - \bar{\psi}^A Z^B \bar{Z}_A \psi_B) \\
&\quad - \frac{2\pi i}{k} \epsilon_{ABCD} \text{Tr}(\bar{\psi}^A Z^C \psi^B Z^D) + \frac{2\pi i}{k} \epsilon^{ABCD} \text{Tr}(\bar{\psi}_A \bar{Z}_C \psi_B \bar{Z}_D), \\
V &= \frac{1}{3} \text{Tr}(4Z^A \bar{Z}_A Z^B \bar{Z}_B Z^C \bar{Z}_C - 4Z^A \bar{Z}_B Z^C \bar{Z}_A Z^B Z_C - Z^A \bar{Z}_A Z^B \bar{Z}_B Z^C \bar{Z}_C - \bar{Z}_A Z^A \bar{Z}_B Z^B \bar{Z}_C Z^C) \\
&\quad D_\mu Z^A = \partial_\mu Z^A - i A_\mu^L Z^A + i Z^A A_\mu^R \\
\mathcal{L}_{CS} &= \frac{k}{4\pi} \epsilon^{\mu\nu\lambda} \left(\text{Tr} \left(A_\mu^L \partial_\nu A_\lambda^L - \frac{2}{3} i A_\mu^L A_\nu^L A_\lambda^L \right) - \text{Tr} \left(A_\mu^R \partial_\nu A_\lambda^R - \frac{2}{3} i A_\mu^R A_\nu^R A_\lambda^R \right) \right), \\
\delta Z^A &= i \bar{\epsilon}^{AB} \Psi_B, \\
\delta \psi_B &= \gamma^\mu \epsilon_{AB} D_\mu Z^A + \frac{2\pi}{k} (Z^C \bar{Z}_B Z^D - Z^D \bar{Z}_B Z^C) \epsilon_{CD} - \frac{2\pi}{k} (Z^A \bar{Z}_C Z^C - Z^C \bar{Z}_C Z^A) \epsilon_{AB} \\
\delta A_\mu^L &= -\frac{2\pi}{k} (\bar{\epsilon}_{AB} \gamma_\mu Z^B \bar{\psi}^A - \bar{\epsilon}^{AB} \gamma_\mu \psi_A \bar{Z}_B) \\
\delta A_\mu^R &= -\frac{2\pi}{k} (\bar{\epsilon}_{AB} \gamma_\mu \bar{\psi}^A Z^B - \bar{\epsilon}^{AB} \gamma_\mu \bar{Z}_B \psi_A). \\
ds^2 &= \frac{R^2}{4} ds_{\text{AdS}_4}^2 + R^2 ds_{S^7/\mathbb{Z}_k}^2
\end{aligned}$$

$$\begin{aligned}
ds_{S^7/\mathbb{Z}_k}^2 &= \frac{1}{k^2} (d\phi + k\omega)^2 + ds_{\mathbb{CP}^3}^2 \\
ds_{\mathbb{dP}^3}^2 &= \frac{1}{|z|^4} (|z|^2 dz_i \bar{d}z_i - \bar{z}_i z_j dz_i d\bar{z}_j) \\
J &\sim id \left(\frac{z_i}{|z|} \right) \wedge d \left(\frac{\bar{z}_i}{|z|} \right) \\
ds_{IIA}^2 &= \frac{R^3}{k} \left(\frac{1}{4} ds_{AdS_4}^2 + ds_{\mathbb{CP}^3}^2 \right) \\
e^{2\Phi} &= \frac{R^3}{k^3} \sim \frac{1}{n^2} \left(\frac{n}{k} \right)^{\frac{5}{2}} \\
\Phi^i &= \text{diag}(x_1^i, x_2^i, \dots, x_n^i) \\
(\mathbb{R}^6)^n / \mathbb{S}_n &\equiv \text{Sym}_n(\mathbb{R}^6). \\
Z^A &= \text{diag}(z_1^A, z_2^A) \\
g_{12} &: \text{diag}(z_1^A, z_2^A) \cong \text{diag}(z_2^A, z_1^A) \\
\mathcal{L} &= -D_\mu z_1^A D^\mu z_{1A} - D_\mu z_2^A D^\mu z_{2A} + \frac{k}{2\pi} \varepsilon^{\mu\nu\lambda} B_\mu \partial_\nu Q_\lambda \\
\mathcal{L} &= -D_\mu z_1^A D^\mu z_{1A} - D_\mu z_2^A D^\mu z_{2A} + \frac{k}{4\pi} \varepsilon^{\mu\nu\lambda} B_\mu H_{\nu\lambda} + \frac{1}{4\pi} \varepsilon^{\mu\nu\lambda} \sigma \partial_\mu H_{\nu\lambda} \\
w_1^A &= e^{i\sigma/k} z_1^A, w_2^A = e^{-i\sigma/k} z_2^A \\
\mathcal{L} &= -\partial_\mu w_1^A \partial^\mu w_{1A} - \partial_\mu w_2^A \partial^\mu w_{2A} \\
\Psi_\gamma &= e^{i\oint_\gamma \square A} \square \Psi_0 \\
\Psi_\gamma &= e^{i \int_D F} \Psi_0 \\
e^{i \int_{D-D'} F} &= 1 \\
\frac{1}{4\pi} \int_\Sigma H &\in \mathbb{Z}, \\
\frac{1}{4\pi} \int \varepsilon^{\mu\nu\lambda} \partial_\mu H_{\nu\lambda} &\in 2\mathbb{Z} \\
\frac{1}{4\pi} \int \varepsilon^{\mu\nu\lambda} \partial_\mu H_{\nu\lambda} &\in \mathbb{Z} \\
g_{\text{SU}(2)} &= \begin{cases} z_1^A \cong e^{\pi i/k} z_1^A & z_2^A \cong e^{-\pi i/k} z_2^A & \text{SU}(2) \times \text{SU}(2) \\ z_1^A \cong e^{2\pi i/k} z_1^A & z_2^A \cong e^{-2\pi i/k} z_2^A & (\text{SU}(2) \times \text{SU}(2))/\mathbb{Z}_2 \\ (\mathbb{C}^4 \times \mathbb{C}^4)/\mathbb{D}_{4k} & \text{for } \text{SU}(2) \times \text{SU}(2) \\ (\mathbb{C}^4 \times \mathbb{C}^4)/\mathbb{D}_{2k} & \text{for } (\text{SU}(2) \times \text{SU}(2))/\mathbb{Z}_2 \end{cases} \\
Z^C Z_B Z^D - Z^D Z_B Z^C &= 0 \\
Z^A &= \text{diag}(z_1^A, \dots, z_n^A) \\
A_\mu^L &= \text{diag}(A_\mu^{L1}, \dots, A_\mu^{Ln}), A_\mu^R = \text{diag}(A_\mu^{R1}, \dots, A_\mu^{Rn}) \\
\mathcal{L} &= -\frac{1}{2} \sum_{i=1}^n D_\mu z_i^A D^\mu z_A^i + \frac{k}{4\pi} \varepsilon^{\mu\nu\lambda} \sum_{i=1}^n B_\mu^i \partial_\nu Q_\lambda^i \\
\mathcal{L} &= -\frac{1}{2} \sum_i D_\mu z_i^A D^\mu z_A^i + \frac{k}{8\pi} \varepsilon^{\mu\nu\lambda} \sum_i B_{i\mu} H_{i\nu\lambda} + \frac{1}{8\pi} \varepsilon^{\mu\nu\lambda} \sum_i \sigma_i \partial_\mu H_{i\nu\lambda} \\
\mathcal{L} &= -\frac{1}{2} \sum_i \partial_\mu w_i^A \partial^\mu w_A^i \\
\frac{1}{8\pi} \varepsilon^{\mu\nu\lambda} \int \partial_\mu H_{i\nu\lambda} &\in \mathbb{Z} \\
w_i^A &\cong e^{2\pi i/k} w_i^A \\
\mathcal{M}_k &= (\mathbb{C}^4/\mathbb{Z}_k)^n / \mathbb{S}_n \equiv \text{Sym}_n(\mathbb{C}^4/\mathbb{Z}_k) \\
L_{CS} &= \frac{k}{4\pi} \text{Tr} \left(A \wedge dA + \frac{2}{3} A \wedge A \wedge A - \tilde{A} \wedge d\tilde{A} - \frac{2}{3} \tilde{A} \wedge \tilde{A} \wedge \tilde{A} \right) \\
\delta\Phi &= -\Lambda\Phi + \Phi\tilde{\Lambda}
\end{aligned}$$



$$\begin{aligned}
D_\mu \Phi &= \partial_\mu \Phi + A_\mu \Phi - \Phi \tilde{A}_\mu \\
&\quad \frac{k}{4\pi} \text{Tr}(D_\mu \Phi^\dagger D^\mu \Phi) \\
&\quad \frac{k}{4\pi} \text{Tr}|A_\mu \Phi - \Phi \tilde{A}_\mu|^2 \\
&\quad \frac{k}{4\pi} v^2 \text{Tr}(A_\mu - \tilde{A}_\mu)^2, \\
B &= \frac{1}{2}(A - \tilde{A}), C = \frac{1}{2}(A + \tilde{A}). \\
\mathcal{L} &= \frac{k}{\pi} \text{Tr} \left(B \wedge F^{(C)} + \frac{1}{3} B \wedge B \wedge B - v^2 B \wedge *B \right) \\
&\quad F^{(C)} + B \wedge B - 2v^2 *B = 0 \\
B &= -\frac{1}{2v^2} F^* F^{(C)} - \frac{1}{2v^2} *(B \wedge B) \\
&= -\frac{1}{2v^2} *F^{(C)} - \frac{1}{8v^6} **F^{(C)} \wedge *F^{(C)} + \dots \\
\mathcal{L} &= \frac{k}{\pi} \left(-\frac{1}{4v^2} F^{(C)} \wedge *F^{(C)} - \frac{1}{24v^6} F^*(C) \wedge *F^{(C)} \wedge *F^{(C)} + \dots \right) \\
\mathcal{L} &= -\frac{1}{2} D_\mu X^{aI} D^\mu X_a^I + \frac{i}{2} \bar{\Psi}^a \Gamma^\mu D_\mu \Psi_a + \frac{i}{4} f_{abcd} \bar{\Psi}^b \Gamma^{IJ} X^{cI} X^{dJ} \Psi^a \\
&\quad -\frac{1}{12} (f_{abcd} X^{aI} X^{bJ} X^{cK}) (f_{efg}{}^d X^{eI} X^{fJ} X^{gK}) \\
&\quad + \frac{1}{2} \epsilon^{\mu\nu\lambda} \left(f_{abcd} A_\mu^{ab} \partial_\nu A_\lambda^{cd} + \frac{2}{3} f_{aef}{}^g f_{bcdg} A_\mu^{ab} A_\nu^{cd} A_\lambda^{ef} \right) \\
D_\mu X^{aI} &= \partial_\mu X^{aI} + f_{bcd}^a A_\mu^{cd} X^{bI} \\
V(X) &= \frac{1}{12} \sum_{I,J,K=1}^8 (\epsilon_{abcd} \epsilon_{efg}{}^d X^{aI} X^{bJ} X^{cK} X^{eI} X^{fJ} X^{gK}) \\
&= \frac{1}{2} \sum_{i<j}^7 (\epsilon_{abcd} \epsilon_{efg}{}^d X^{aI} X^{bJ} X^{c(8)} X^{eI} X^{fJ} X^{g(8)}) \\
&\quad + \frac{1}{2} \sum_{i<j<k}^7 (\epsilon_{abcd} \epsilon_{efg}{}^d X^{aI} X^{bJ} X^{cK} X^{eI} X^{fJ} X^{gk}) \\
&= \frac{1}{2} v^2 \sum_{i<j}^7 (\epsilon_{AB4D} \epsilon_{EF4}{}^D X^{Ai} X^{Bj} X^{Ei} X^{Fj}) + vO(X^5) + O(X^6). \\
&\quad \frac{1}{2} v^2 \sum_{i<j=1}^7 (\epsilon_{ABC} \epsilon_{EF}^C X^{Ai} X^{Bj} X^{Ei} X^{Fj}) \\
\frac{i}{4} \epsilon_{abcd} \bar{\Psi}^b \Gamma^{IJ} X^{cI} X^{dJ} \Psi^a &= \frac{i}{2} v \epsilon_{ABC} \bar{\Psi}^B \Gamma_i X^{Ci} \Psi^A + O(X^2 \Psi^2) \\
A_\mu^{A4} &\equiv A_\mu^A \quad \text{and} \quad \frac{1}{2} \epsilon_{BC}^A A_\mu^{BC} \equiv B_\mu^A \\
\frac{1}{2} \epsilon^{\mu\nu\lambda} \epsilon_{abcd} A_\mu^{ab} \partial_\nu A_\lambda^{cd} &= 2\epsilon^{\mu\nu\lambda} \epsilon_{ABC} A_\mu^{AB} \partial_\nu A_\lambda^C = 4\epsilon^{\mu\nu\lambda} B_\mu^A \partial_\nu A_{\lambda A} \\
\frac{1}{3} \epsilon^{\mu\nu\lambda} \epsilon_{aef}{}^g \epsilon_{bcdg} A_\mu^{ab} A_\nu^{cd} A_\lambda^{ef} &= -4\epsilon^{\mu\nu\lambda} \epsilon_{ABC} B_\mu^A B_\nu^B A_\lambda^C - \frac{4}{3} \epsilon^{\mu\nu\lambda} \epsilon_{ABC} B_\mu^A B_\nu^B B_\lambda^C. \\
D_\mu X^{AI} &= \partial_\mu X^{AI} + \epsilon_{bcd}^A A_\mu^{cd} X^{bI} \\
&= \partial_\mu X^{AI} + 2\epsilon_{BC}^A A_\mu^C X^{BI} + 2B_\mu^A X^{4(I)} \\
D_\mu X^{4I} &= \partial_\mu X^{4I} - 2B_{\mu A} X^{AI}
\end{aligned}$$

$$\begin{aligned}
\mathcal{L}_{\text{kinetic}} &= -2v^2 B_\mu^A B_A^\mu - 2B_\mu^A X^{4I} D'^\mu X_A^I - 2v B_\mu^A D'^\mu X_A^{(8)} \\
&\quad - 2B_{\mu A} X^{AI} B_B^\mu X^{BI} - 2B_\mu^A B_A^\mu X^{4I} X_{4I} + 2B_\mu^A X^{AI} \partial_\mu X^{4I} + \dots, \\
&\quad D'_\mu X^{AI} = \partial_\mu X^{AI} - 2\varepsilon_{BC}^A A_\mu^B X^{CI} \\
\mathcal{L}_{\text{CS}} &= 2\varepsilon^{\mu\nu\lambda} B_\mu^A F'_{\nu\lambda A} - \frac{4}{3} \varepsilon^{\mu\nu\lambda} \varepsilon_{ABC} B_\mu^A B_\nu^B B_\lambda^C + \dots, \\
&\quad F'_{\nu\lambda}{}^A = \partial_\nu A_\lambda^A - \partial_\lambda A_\nu^A - 2\varepsilon_{BC}^A A_\nu^B A_\lambda^C. \\
\mathcal{L} &= -2v^2 B_\mu^A B_A^\mu - 2v B_\mu^A D'^\mu X_A^{(8)} + 2\varepsilon^{\mu\nu\lambda} B_\mu^A F'_{\nu\lambda A} \\
&\quad B_\mu^A = \frac{1}{2v^2} \varepsilon_\mu^{\nu\lambda} F'_{\nu\lambda}{}^A - \frac{1}{2v} D'_\mu X^{A(8)}. \\
&\quad -\frac{1}{v^2} F'_{\mu\nu}{}^A F'^{\mu\nu}{}_A - \frac{1}{2} \partial_\mu X^{4I} \partial^\mu X_4^I - \frac{1}{2} D_\mu X^{Ai} D^\mu X_A^i + O(BX\partial X) + O(B^2 X^2) + O(B^3) \\
&\quad A \rightarrow \frac{1}{2} A \\
&\quad D'_\mu X^{AI} \rightarrow D_\mu X^{AI} \equiv \partial_\mu X^{AI} - \varepsilon_{BC}^A A_\mu^B X^{CI} \\
&\quad F'_{\mu\nu}{}^A \rightarrow \frac{1}{2} F_{\mu\nu}^A \equiv \frac{1}{2} (\partial_\mu A_\nu^A - \partial_\nu A_\mu^A - \varepsilon_{BC}^A A_\mu^B A_\nu^C). \\
&\quad -\frac{1}{4v^2} F_{\mu\nu}^A F^{\mu\nu}{}_A - \frac{1}{2} \partial_\mu X^{4I} \partial^\mu X_4^I - \frac{1}{2} D_\mu X^{Ai} D^\mu X_A^i + \frac{1}{v} O(X\partial X(F/v + DX)) \\
&\quad + \frac{1}{v} O(X^2(F/v + DX)^2) + \frac{1}{v^3} O((F/v + DX)^3) \\
&\quad \frac{i}{2} \bar{\Psi}^a \Gamma^\mu D_\mu \Psi_a \rightarrow \frac{i}{2} \bar{\Psi}^A \Gamma^\mu D_\mu \Psi_A + \frac{i}{2} \bar{\Psi}^4 \Gamma^\mu \partial_\mu \Psi_4 \\
&\quad \mathcal{L} = \mathcal{L}_{\text{SU}(2)} + \mathcal{L}_{\text{U}(1)} \\
&\quad \mathcal{L}_{\text{U}(1)} = -\frac{1}{2} \partial_\mu X^{4I} \partial^\mu X_4^I + \frac{i}{2} \bar{\Psi}^4 \Gamma^\mu \partial_\mu \Psi_4 \\
&\quad \mathcal{L}_{\text{SU}(2)} = \frac{1}{v^2} \mathcal{L}_0 + \frac{1}{v^3} \mathcal{L}_1 + O\left(\frac{1}{v^4}\right) \\
&\quad \mathcal{L}_0 = -\frac{1}{4} F_{\mu\nu A} F^{\mu\nu A} - \frac{1}{2} D_\mu X^{Ai} D^\mu X_A^i + \frac{1}{4} (\varepsilon_{ABC} X^{Ai} X^{Bj}) (\varepsilon_{DE}^C X^{Di} X^{Ej}) \\
&\quad + \frac{i}{2} \bar{\Psi}^A \not{D} \Psi_A + \frac{i}{2} \varepsilon_{ABC} \bar{\Psi}^A \Gamma^i X^{Bi} \Psi^C \\
&\quad F_{\mu\nu}^A = \partial_\mu A_\nu^A - \partial_\nu A_\mu^A - \varepsilon_{BC}^A A_\mu^B A_\nu^C \text{ and } D_\mu^{AB} = \partial_\mu \delta^{AB} + \varepsilon_C^{AB} A_\mu^C \\
\mathcal{L} &= -\text{Tr}(D^\mu \bar{Z}_A, D_\mu Z^A) + \frac{k}{4\pi} \varepsilon^{\mu\nu\lambda} \left(\text{Tr} \left(A_\mu^L \partial_\nu A_\lambda^L - \frac{2}{3} i A_\mu^L A_\nu^L A_\lambda^L \right) - \text{Tr} \left(A_\mu^R \partial_\nu A_\lambda^R - \frac{2}{3} i A_\mu^R A_\nu^R A_\lambda^R \right) \right) \\
&\quad - \frac{1}{3} \text{Tr}(4Z^A \bar{Z}_A Z^B \bar{Z}_B Z^C \bar{Z}_C - 4Z^A \bar{Z}_B Z^C \bar{Z}_A Z^B Z_C - Z^A \bar{Z}_A Z^B \bar{Z}_B Z^C \bar{Z}_C - \bar{Z}_A Z^A \bar{Z}_B Z^B \bar{Z}_C Z^C) \\
&\quad \hat{D}_\mu Z^A = \partial_\mu Z^A - i A_\mu^L Z + i Z^A A_\mu^R \\
&\quad Z^A = v \delta^{A4} + z^A \\
&\quad Z^A = v \delta^{A4} \mathbb{1}_{N \times N} + \frac{1}{\sqrt{2}} X^A + i \frac{1}{\sqrt{2}} X^{A+4} \\
&\quad A_\mu^+ = \frac{1}{2} (A_\mu^L + A_\mu^R), A_\mu^- = \frac{1}{2} (A_\mu^L - A_\mu^R) \\
&\quad \hat{D}_\mu Z^A = D_\mu Z^A - i \{A_\mu^-, Z^A\} \\
&\quad D_\mu Z^A = \partial_\mu Z^A - i [A_\mu^+, Z^A] \\
&\quad F_{\mu\nu}^+ = \partial_\mu A_\nu^+ - i [A_\mu^+, A_\nu^+]. \\
S_{\text{CS}} &= \int d^3x \frac{k}{2\pi} \varepsilon^{\mu\nu\lambda} \text{Tr} \left(A_\mu^- F_{\nu\lambda}^+ - \frac{2i}{3} A_\mu^- A_\nu^- A_\lambda^- \right) \\
&\quad Z^A = Z_0^A T^0 + i Z_a^A T^a, \\
&\quad A_\mu^L = A_\mu^{L0} T^0 + A_\mu^{La} T^a, \\
&\quad A_\mu^R = A_\mu^{R0} T^0 + A_\mu^{Ra} T^a,
\end{aligned}$$

$$\begin{aligned}
Z^A &= \left(\frac{X_0^A}{\sqrt{2}} + v\delta^{A,4} \right) T^0 + i \frac{X_0^{A+4}}{\sqrt{2}} T^0 + i \frac{X_a^A}{\sqrt{2}} T^a - \frac{X_a^{A+4}}{\sqrt{2}} T^a. \\
\hat{D}_\mu Z^A &= \frac{\partial_\mu X_0^A}{\sqrt{2}} T^0 - \frac{D_\mu X_a^{A+4}}{\sqrt{2}} T^a + \frac{i\partial_\mu X_0^{A+4}}{\sqrt{2}} T^0 + \frac{iD_\mu X_a^A}{\sqrt{2}} T^a \\
&\quad - 2ivA_{\mu a}^- T^a \delta^{A4} - i\sqrt{2}A_{\mu a}^- X_0^A T^a + \sqrt{2}A_{\mu a}^- X_0^{A+4} T^a - A_{\mu}^{-a} d_{abc} T^c X_b^{A+4} \\
&\quad + iA_{\mu}^{-a} d_{abc} T^c X_b^A - 2ivA_{\mu 0}^- T^0 \delta^{A4} - i\sqrt{2}A_{\mu 0}^- X_0^A T^0 + \sqrt{2}A_{\mu 0}^- X_0^{A+4} T^0 \\
&\quad + \sqrt{2}A_{\mu 0}^- X_a^A T^a + i\sqrt{2}A_{\mu 0}^- X_a^{A+4} T^a, \\
\text{Tr}|\hat{D}_\mu Z^A|^2 &= N \frac{(\partial_\mu X_0^A)^2}{2} + \left(\frac{(D_\mu X)_c^{A+4}}{\sqrt{2}} \right)^2 + \left(2vA_{\mu c}^- \delta^{A4} - \frac{(D_\mu X)_c^A}{\sqrt{2}} \right)^2 \\
&\quad + N \left(\frac{1}{\sqrt{2}} \partial_\mu X_0^{A+4} - 2vA_{\mu 0}^- \delta^{A4} \right)^2 \\
&\quad - \frac{k}{2\pi} \epsilon^{\mu\nu\lambda} \frac{1}{v} \frac{1}{2\sqrt{2}} (D_\mu X)_a F_{\nu\lambda}^{+a} - \frac{nk}{2\pi} \epsilon^{\mu\nu\lambda} \frac{1}{v} \frac{1}{2\sqrt{2}} (\partial_\mu X_0^8) F_{\nu\lambda}^{+0} \\
S &= \int d^3x \left[\frac{k}{2\pi} e^{\mu\nu\lambda} \text{Tr} \left(A_{\mu}^- F_{\nu\lambda}^+ - \frac{2i}{3} A_{\mu}^- A_{\nu}^- A_{\lambda}^- \right) - \text{Tr}|\hat{D}_\mu Z^A|^2 \right] \\
&= \int d^3x \left[\frac{k}{2\pi} e^{\mu\nu\lambda} \left(A_{\mu a}^- - \frac{1}{2v} \frac{1}{\sqrt{2}} (D_\mu X)_a \right) F_{\nu\lambda}^{+a} + \frac{nk}{2\pi} e^{\mu\nu\lambda} \left(A_{\mu 0}^- - \frac{1}{2v} \frac{1}{\sqrt{2}} (\partial_\mu X)_0^8 \right) F_{\nu\lambda}^{+0} \right. \\
&\quad - \left(2vA_{\mu a}^- - \frac{1}{\sqrt{2}} (D_\mu X)_a \right)^2 - n \left(\frac{1}{\sqrt{2}} \partial_\mu X_0^{A+4} - 2vA_{\mu 0}^- \delta^{A4} \right)^2 \\
&\quad \left. - \frac{1}{2} (D_\mu X)_a^{I'} (D^\mu X)_a^{I'} - \frac{1}{2} n \partial_\mu X_0^A \partial^\mu X_0^A \right] \\
&\quad A_{\mu a}^- \rightarrow A_{\mu a}^- + \frac{1}{2v} \frac{1}{\sqrt{2}} (D_\mu X)_a \quad \text{and} \quad A_{\mu 0}^- \rightarrow A_{\mu 0}^- + \frac{1}{2v} \frac{1}{\sqrt{2}} (\partial_\mu X)_0^8 \\
S &= \int d^3x \left(\frac{k}{2\pi} e^{\mu\nu\lambda} (A_{\mu a}^- F_{\nu\lambda}^{a+} + NA_{\mu 0}^- F_{\nu\lambda}^{+0}) - 4v^2 A_{\mu a}^- A_a^{-\mu} - 4nv^2 A_{\mu 0}^- A_0^{-\mu} \right. \\
&\quad \left. - \frac{1}{2} (D_\mu X)_a^{I'} (D^\mu X)_a^{I'} - \frac{1}{2} n \partial_\mu X_0^{I'} \partial^\mu X_0^{I'} \right), \\
&\quad A_{\mu}^- = \frac{k}{16\pi v^2} \epsilon_{\mu\nu\lambda} F^{+\nu\lambda} \\
S &= \int d^3x \left[-\text{Tr} \left(\frac{k^2}{32\pi^2 v^2} F^{+\mu\nu} F_{\mu\nu}^+ \right) - \frac{1}{2} (D_\mu X)_a^{I'} (D^\mu X)_a^{I'} - \frac{1}{2} n \partial_\mu X_0^{I'} \partial^\mu X_0^{I'} \right] \\
&\quad \frac{k^2}{32\pi^2 v^2} = \frac{1}{4g_{YM}^2} \\
&\quad -V_6 \rightarrow -\frac{g_{YM}^2}{4} \text{Tr}([X^{I'}, X^{J'}][X^{J'}, X^{I'}]) \\
Z^4 &\rightarrow Z^4 e^{2\pi i/k} \simeq Z^4 \left(1 + 2\pi i \frac{1}{k} + \dots \right) \simeq Z^4 + 2\pi i \frac{Z^4}{k} \\
Z^4 T_{M2}^{-\frac{1}{2}} &\rightarrow Z^4 T_{M2}^{-\frac{1}{2}} + 2\pi i R \\
T^a &= \left\{ -\frac{i}{\sqrt{2}} \sigma_1, -\frac{i}{\sqrt{2}} \sigma_2, -\frac{i}{\sqrt{2}} \sigma_3, \frac{1}{\sqrt{2}} \mathbb{1}_{2 \times 2} \right\}, \\
f^{abcd} &= \frac{2\pi}{k} \epsilon^{abcd} \quad \text{and} \quad \text{Tr}(T_b T^a) = \delta_b^a \\
\mathcal{L}_{\text{u}(n) \times \text{u}(n)} &= \mathcal{L}_{\text{su}(n) \times \text{su}(n)}^{\text{gauged}} + \frac{nk}{4\pi} \epsilon^{\mu\nu\lambda} B_\mu \partial_\nu Q_\lambda \\
\mathcal{L}_{\text{u}(n) \times \text{u}(n)} &= \mathcal{L}_{\text{su}(n) \times \text{su}(n)}^{\text{gauged}} + \frac{nk}{8\pi} \epsilon^{\mu\nu\lambda} B_\mu H_{\nu\lambda} + \frac{n}{8\pi} \sigma \epsilon^{\mu\nu\lambda} \partial_\mu H_{\nu\lambda}
\end{aligned}$$



$$\begin{aligned}
\mathcal{L}_{\mathfrak{u}(n)\oplus\mathfrak{u}(n)} &= \mathcal{L}_{\mathfrak{su}(n)\oplus\mathfrak{su}(n)}^{\text{gauged}} + \frac{nk}{8\pi} \varepsilon^{\mu\nu\lambda} B_\mu H_{\nu\lambda} - \frac{n}{8\pi} \varepsilon^{\mu\nu\lambda} \partial_\mu \sigma H_{\nu\lambda} \\
&\quad B_\mu = \frac{1}{k} \partial_\mu \sigma \\
&\quad \sigma \rightarrow \sigma + k\theta \\
\mathcal{L}_{\mathfrak{u}(n)\oplus\mathfrak{u}(n)}(Z^A, \psi_A, \tilde{A}_{\mu b}^a, B_\mu, Q_\mu) &\cong \mathcal{L}_{\mathfrak{su}(n)\oplus\mathfrak{su}(n)}(e^{i\sigma/k} Z^A, e^{i\sigma/k} \psi_A, \tilde{A}_{\mu b}^a) \\
&\quad \int dF^{L/R} \in \frac{2\pi}{n} \mathbb{Z} \\
&\quad \int dH = \int \frac{1}{2} \varepsilon^{\mu\nu\lambda} \partial_\mu H_{\nu\lambda} \in \frac{4\pi}{n} \mathbb{Z} \\
\hat{Z}^A &\cong e^{2\pi i/k} \hat{Z}^A \quad \text{and} \quad \hat{\psi}_A \cong e^{2\pi i/k} \hat{\psi}_A \\
Z^A &= \text{diag}(z_1^A, \dots, z_n^A) \\
\det(g_{\text{U}(1)}^{l_B} g_0) &= \det(g_{\text{U}(1)}^{l_B}) = e^{2\pi i n l_B/k} \\
\det(g_1^{l_1} \dots g_n^{l_n}) &= e^{2\pi i (l_1 + \dots + l_n)/k} \\
l &= n l_B \text{ mod } k \\
l &= p k \text{ mod } n. \\
r_1^A &= z_1^A + z_2^A, r_2^A = i(z_1^A - z_2^A), \\
g_{12}: r_1^A &\cong r_1^A, \quad r_2^A \cong -r_2^A \\
g_{12} g_{\text{SU}(2)}^2: r_1^A &\cong -r_1^A, \quad r_2^A \cong r_2^A \\
g_{\text{SU}(2)} g_{12}: r_1^A &\cong r_2^A, \quad r_2^A \cong r_1^A \\
\mathcal{L} &= -\frac{1}{2} \sum_{A=1}^4 \partial_\mu R^A \partial^\mu R^A - \frac{1}{2} \sum_{A=1}^4 (R^A)^2 (\partial_\mu \theta^A - B_\mu) (\partial^\mu \theta^A - B^\mu) + \frac{k}{4\pi} \varepsilon^{\mu\nu\lambda} B_\mu \partial_\nu Q_\lambda \\
&\quad \Pi_{R^A} = \partial_0 R^A \\
&\quad \Pi_{\theta^A} = (R^A)^2 (\partial_0 \theta^A - B_0) \\
\Pi_{Q_j} &= \frac{k}{4\pi} \varepsilon^{ij} B_i \\
\Pi_{Q_0} &= \Pi_{B_0} = 0 \\
H &= \int d^2x \left\{ \frac{1}{2} \sum_A \Pi_{R^A}^2 + \frac{1}{2} \sum_A (R^A)^{-2} \Pi_{\theta^A}^2 + \frac{1}{2} \sum_A (\partial_i R^A)^2 + \frac{1}{2} \sum_A (R^A)^2 (\partial_i \theta^A - B_i)^2 \right. \\
&\quad \left. - \frac{k}{4\pi} F_{12} Q_0 + \left(\sum_A \Pi_{\theta^A} + \frac{k}{4\pi} H_{12} \right) B_0 \right\}, \\
&\quad F_{12} = 0 \\
H_{12} &= -\frac{4\pi}{k} \sum_A \Pi_{\theta^A} \\
&\quad \int d^2x H_{12} \in 4\pi \mathbb{Z} \\
&\quad \sum_A P_{\theta^A} \in k\mathbb{Z} \\
\mathcal{L} &= -\text{Tr}(D_\mu Z^A D^\mu \bar{Z}_A) - i\text{Tr}(\bar{\psi}^A \gamma^\mu D_\mu \psi) + \mathcal{L}_{\text{Yukawa}} - V \\
&\quad + \frac{k}{4\pi} \varepsilon^{\mu\nu\lambda} \text{Tr} \left(A_\mu^L \partial_\nu A_\lambda^L - \frac{2i}{3} A_\mu^L A_\nu^L A_\lambda^L \right) - \frac{k}{4\pi} \varepsilon^{\mu\nu\lambda} \text{Tr} \left(A_\mu^R \partial_\nu A_\lambda^R - \frac{2i}{3} A_\mu^R A_\nu^R A_\lambda^R \right) \\
H &= \int d^2x \text{Tr}(\Pi_{Z^A} \Pi_{\bar{Z}_A}) + \text{Tr}(D_i Z^A D^i \bar{Z}_A) - \mathcal{L}_{\text{Yukawa}} + V \\
&\quad + \text{Tr} \left(iZ^A \Pi_{Z^A} - i\Pi_{\bar{Z}_A} \bar{Z}_A + i\psi_A \psi^A - \frac{k}{2\pi} F_{12}^L \right) A_0^L + \text{Tr} \left(i\bar{Z}_A \Pi_{\bar{Z}_A} - i\Pi_{Z^A} Z^A - i\psi^A \psi_A + \frac{k}{2\pi} F_{12}^R \right) A_0^R
\end{aligned}$$

$$\frac{k}{2\pi} F_{12}^L = iZ^A \Pi_{Z^A} - i\Pi_{\bar{Z}^A} \bar{Z}^A + i\psi_A \psi^A$$

$$\frac{k}{2\pi} F_{12}^R = i\Pi_{Z^A} Z^A - i\bar{Z}^A \Pi_{\bar{Z}^A} + i\psi^A \psi_A$$

$$Z^A = \begin{pmatrix} z_1^A \\ z_2^A \\ \vdots \\ z_n^A \end{pmatrix}$$

$$F_{12}^L = F_{12}^R = -\frac{2\pi}{k} \begin{pmatrix} \sum_A \Pi_{\theta_1^A} & & & & & & \\ & \sum_A \Pi_{\theta_2^A} & & & & & \\ & & \ddots & & & & \\ & & & \ddots & & & \\ & & & & \sum_A \Pi_{\theta_n^A} & & \end{pmatrix},$$

$$F = \star \frac{Q_M}{2} d \left(\frac{1}{|x - x_0|} \right)$$

$$e^{2\pi i Q_M} = 1$$

$$Q_M = \vec{q} \cdot \vec{H} + q \mathbb{1}$$

$$|\vec{\mu}^1\rangle = |\vec{\lambda}^1\rangle, |\vec{\mu}^2\rangle = |\vec{\lambda}^1 - \vec{\alpha}_1\rangle, |\vec{\mu}^3\rangle = |\vec{\lambda}^1 - \vec{\alpha}_1 - \vec{\alpha}_2\rangle, \dots, |\vec{\mu}^n\rangle = |\vec{\lambda}^1 - \vec{\alpha}_1 - \dots - \vec{\alpha}_{n-1}\rangle$$

$$q_i = \langle \vec{\mu}^i | Q_M | \vec{\mu}^i \rangle = \vec{q} \cdot \vec{\mu}^i + q, i = 1, \dots, n$$

$$w_i = \vec{q} \cdot \vec{\alpha}_i = \vec{q} \cdot (\vec{\mu}^i - \vec{\mu}^{i+1}) = q_i - q_{i+1}, i = 1, \dots, n-1$$

$$\vec{q} = (q_1 - q_2)^L \vec{\lambda}^1 + (q_2 - q_3)^L \vec{\lambda}^2 \dots + (q_{n-1} - q_n)^L \vec{\lambda}^{n-1}$$

$$q = \frac{1}{n} (q_1 + \dots + q_n)$$

$$\vec{q} = q_1 \vec{\alpha}_1 + (q_1 + q_2) \vec{\alpha}_2 + \dots + (q_1 + q_2 + \dots + q_{n-1}) \vec{\alpha}_{n-1}$$

$$\Phi = \text{diag}(v, -v) + \frac{Q_M}{4\pi r} + \dots$$

$$\begin{pmatrix} 2 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}, \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 0 \end{pmatrix}, \begin{pmatrix} 0 & 0 & 0 \\ 0 & 2 & 0 \\ 0 & 0 & 0 \end{pmatrix}, \begin{pmatrix} 1 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 1 \end{pmatrix}, \begin{pmatrix} 0 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix}, \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 2 \end{pmatrix}$$

$$\begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 0 \end{pmatrix},$$

$$\frac{1}{2\pi} \int_{S^2} F^L = \frac{1}{2\pi} \int_{S^2} F^R = Q_M \in \mathfrak{u}(1)^n$$

$$\mathcal{M}_{Q_M}(x) \rightarrow e^{ik/2\pi \text{tr} \int (D\omega_L \wedge F^L - D\omega_R \wedge F^R)} \mathcal{M}_{Q_M}(x)$$

$$= e^{ik \text{tr}((\omega_L(x) - \omega_R(x)) Q_M)} \mathcal{M}_{Q_M}(x).$$

$$\vec{\Lambda} = k\vec{q} \oplus -k\vec{q}$$

$$\mathcal{M}_{\vec{\Lambda}}(x),$$

$$\delta\psi_A = \gamma^\mu D_\mu Z^B \epsilon_{BA}$$

$$D_\mu Z^1 = 0$$

$$\frac{k}{2\pi} \varepsilon_{\mu\nu\lambda} F_L^{\nu\lambda} = iZ^1 D_\mu \bar{Z}_1 = D_\mu (iZ^1 \bar{Z}_1)$$

$$\frac{k}{2\pi} \varepsilon_{\mu\nu\lambda} F_R^{\nu\lambda} = iD_\mu \bar{Z}_1 Z^1 = D_\mu (i\bar{Z}_1 Z^1)$$

$$Z^1 \bar{Z}_1 = -\frac{2\pi i}{k} \frac{Q_M}{|x - x_0|}$$

$$\text{Tr}(Z^A \bar{Z}^B Z^C \dots)$$

$$\text{Tr}(\mathcal{M}_{\vec{\Lambda}} Z^{A_1} Z^{A_2} Z^{A_3} \dots Z^{A_p})$$



$$\begin{aligned}
\vec{\Lambda} &= p\vec{\lambda}^{n-1} \\
Q_B^i &= \frac{k}{4\pi} \int H_{12}^i + \Sigma_A \Pi_{\theta_i^A} \\
Q_Q^i &= \frac{k}{4\pi} \int F_{12}^i. \\
\mathcal{M}_i &= e^{i\sigma_i(x)} \\
Z &= \int [dz_i^A][dA_i^L][dA_i^R] e^{i \int d^3x \mathcal{L}} \\
\frac{1}{8\pi} \int d^3y \varepsilon^{\mu\nu\lambda} \partial_\mu H_{i\nu\lambda}(y) &\rightarrow \frac{1}{8\pi} \int d^3y (\varepsilon^{\mu\nu\lambda} \partial_\mu H_{i\nu\lambda}(y) + 8\pi p \delta(x-y)) \\
e^{i\sigma_i(x)} &= e^{i \int_Y d\sigma_i} = e^{ik \int_Y B_i} \\
D_\mu e^{i\sigma_i} &= (i\partial_\mu \sigma_i - ik B_{\mu i}) e^{i\sigma_i} = 0 \\
J_{\mu B}^A &= \text{Tr}(Z^A D_\mu \bar{Z}_B - D_\mu Z^A \bar{Z}_B + i\psi^A \gamma_\mu \psi_B) \\
J_\mu^{AB} &= \text{Tr}((\mathcal{M}_{2\vec{\lambda}^{n-1}})(Z^A D_\mu Z^B - D_\mu Z^A Z^B + i\varepsilon^{ABCD} \psi_C \gamma_\mu \psi_D)) \\
&\quad \text{Tr}(\mathcal{M}_{2\vec{\lambda}^{n-1}} D_\mu Z^A \psi_A) \\
&\quad \text{Tr}(\mathcal{M}_{\vec{\lambda}^{n-1}} Z^A), \\
S &= -T_{M2} \int d^3\sigma \sqrt{-\det(\partial_\mu x^M \partial_\nu x^N g_{MN})} \\
&\quad + \frac{T_{M2}}{3!} \int d^3\sigma \varepsilon^{\mu\nu\lambda} \partial_\mu x^M \partial_\nu x^N \partial_\lambda x^P C_{MNP} \\
S_C &= \frac{1}{3!} e^{\mu\nu\lambda} \int d^3x (aT_{M2} C_{\mu\nu\lambda} + 3b C_{\mu I J} \text{Tr}(D_\nu X^I, D_\lambda X^J) \\
&\quad + 12c C_{\mu\nu I J K L} \text{Tr}(D_\lambda X^I, [X^J, X^K, X^L]) \\
&\quad + 12d C_{[\mu I J} C_{\nu K L]} \text{Tr}(D_\lambda X^I, [X^J, X^K, X^L]) + \dots), \\
S_C &= \frac{1}{3!} \varepsilon^{\mu\nu\lambda} \int d^3x (nT_{M2} C_{\mu\nu\lambda} \\
&\quad + \frac{3}{2} G_{\mu\nu I J} \text{Tr}(X^I, D_\lambda X^J) - \frac{3}{2} C_{\mu I J} \text{Tr}(X^I, \tilde{F}_{\nu\lambda} X^J) \\
&\quad - c G_{\mu\nu\lambda I J K L} \text{Tr}(X^I, [X^J, X^K, X^L])), \\
S_F &= \frac{1}{4} \varepsilon^{\mu\nu\lambda} \int d^3x \text{Tr}(X^I, \tilde{F}_{\mu\nu} X^J) C_{\lambda I J} \\
S_{CG} &= -\frac{c}{2 \cdot 3!} \varepsilon^{\mu\nu\lambda} \int d^3x \text{Tr}(X^I, [X^J, X^K, X^L]) (C_3 \wedge G_4)_{\mu\nu\lambda I J K L} \\
S_{\text{flux}} &= S_C + S_F + S_{CG} \\
&= \frac{1}{3!} e^{\mu\nu\lambda} \int d^3x \left(nT_{M2} C_{\mu\nu\lambda} + \frac{3}{2} G_{\mu\nu I J} \text{Tr}(X^I, D_\lambda X^J) \right. \\
&\quad \left. - c \left(G_7 + \frac{1}{2} C_3 \wedge G_4 \right)_{\mu\nu\lambda I J K L} \text{Tr}(X^I, [X^J, X^K, X^L]) \right). \\
\mathcal{L}_{\text{flux}} &= c \tilde{G}_{I J K L} \text{Tr}(X^I, [X^J, X^K, X^L]) \\
\tilde{G}_{I J K L} &= -\frac{1}{3!} \varepsilon^{\mu\nu\lambda} \left(G_7 + \frac{1}{2} C_3 \wedge G_4 \right)_{\mu\nu\lambda I J K L} \\
&= \frac{1}{4!} \varepsilon_{I J K L M N P Q} G^{M N P Q} \\
\mathcal{L} &= \mathcal{L}_{\mathcal{N}=8} + \mathcal{L}_{\text{mass}} + \mathcal{L}_{\text{flux}} \\
\mathcal{L}_{\text{mass}} &= -\frac{1}{2} m^2 \delta_{I J} \text{Tr}(X^I, X^J) - \frac{ic}{16} \text{Tr}(\tilde{\Psi} \Gamma^{I J K L}, \Psi) \tilde{G}_{I J K L} \\
m^2 &= \frac{c^2}{32 \cdot 4!} G^2
\end{aligned}$$

$$\begin{aligned}
\delta' X_a^I &= 0 \\
\delta' \tilde{A}_\mu^b{}_a &= 0 \\
\delta' \Psi_a &= \frac{c}{8} \Gamma^{IJKL} \Gamma^M \epsilon X_a^M \tilde{G}_{IJKL} \\
G_{MN[IJ} G_{KL]}^{MN} &= 0 \\
G &= \mu(dx^3 \wedge dx^4 \wedge dx^5 \wedge dx^6 + dx^7 \wedge dx^8 \wedge dx^9 \wedge dx^{10}) \\
S_C &= \frac{1}{3!} e^{\mu\nu\lambda} \int d^3x \left(nT_{M2} C_{\mu\nu\lambda} + \frac{3}{2} C_{\mu}^A{}_B \text{Tr}(D_\nu \bar{Z}_A, D_\lambda Z^B) + \frac{3}{2} C_{\mu A}^B \text{Tr}(D_\nu Z^A, D_\lambda \bar{Z}_B) \right. \\
&\quad \left. + \frac{3c}{2} C_{\mu\nu AB}^{CD} \text{Tr}([D_\lambda \bar{Z}_D, [Z^A, Z^B; \bar{Z}_C]) + \frac{3c}{2} C_{\mu\nu}^{AB}{}^C{}_D \text{Tr}([D_\lambda Z^D, [\bar{Z}_A, \bar{Z}_B; Z^C]]) \right) \\
S_F &= \frac{1}{8} \epsilon^{\mu\nu\lambda} \int d^3x C_{\mu}^A{}_B \text{Tr}(\bar{Z}_A, \tilde{F}_{\nu\lambda} Z^B) + C_{\mu A}^B \text{Tr}(Z^A, \tilde{F}_{\nu\lambda} \bar{Z}_B) \\
S_{CG} &= -\frac{c}{8 \cdot 3!} \epsilon^{\mu\nu\lambda} \int d^3x (C_3 \wedge G_4)_{\mu\nu AB}^{CD} \text{Tr}(\bar{Z}_D, [Z^A, Z^B; \bar{Z}_C]) \\
S_{\text{flux}} &= S_C + S_F + S_{CG} \\
&= \frac{1}{3!} \epsilon^{\mu\nu\lambda} \int d^3x \left(nT_{M2} C_{\mu\nu\lambda} \right. \\
&\quad \left. + \frac{3}{4} G_{\mu\nu}^A{}_B \text{Tr}(\bar{Z}_A, D_\lambda Z^B) + \frac{3}{4} G_{\mu\nu A}^B \text{Tr}(Z^A, D_\lambda \bar{Z}_B) \right. \\
&\quad \left. - \frac{c}{4} \left(G_7 + \frac{1}{2} C_3 \wedge G_4 \right)_{\mu\nu\lambda AB}^{CD} \text{Tr}([\bar{Z}_D, [Z^A, Z^B; \bar{Z}_C]]) \right) \\
\mathcal{L} &= \mathcal{L}_{\mathcal{N}=6} + \mathcal{L}_{\text{mass}} + \mathcal{L}_{\text{flux}} \\
\mathcal{L}_{\text{flux}} &= \frac{c}{4} \text{Tr}([\bar{Z}_D, [Z^A, Z^B; \bar{Z}_C]) \tilde{G}_{AB}^{CD} \\
\tilde{G}_{AB}^{CD} &= -\frac{1}{3!} \epsilon^{\mu\nu\lambda} \left(G_7 + \frac{1}{2} C_3 \wedge G_4 \right)_{\mu\nu\lambda AB}^{CD} \\
&= \frac{1}{4} \epsilon_{ABEF} \epsilon^{CDGH} G^{EF} \\
\mathcal{L}_{\text{mass}} &= -m^2 \text{Tr}(\bar{Z}_A, Z^A) + \frac{ic}{4} \text{Tr}(\bar{\psi}^A, \psi_F) \tilde{G}_{AE}^{EF} \\
\delta' \psi_{Ad} &= \frac{c}{4} \epsilon_{DF} Z_d^F \tilde{G}_{AE}^{ED} \\
\tilde{G}_{AE}^{EB} \tilde{G}_{BF}^{FC} &= \frac{16m^2}{c^2} \delta_A^C \\
\tilde{G}_{AB}^{CD} &= \frac{1}{2} \delta_B^C \tilde{G}_{AE}^{ED} - \frac{1}{2} \delta_A^C \tilde{G}_{BE}^{ED} - \frac{1}{2} \delta_B^D \tilde{G}_{AE}^{EC} + \frac{1}{2} \delta_A^D \tilde{G}_{BE}^{EC} \\
m^2 &= \frac{1}{32 \cdot 4!} c^2 G^2 \\
\tilde{G}_{AB}^{BC} &= \begin{pmatrix} \mu & 0 & 0 & 0 \\ 0 & \mu & 0 & 0 \\ 0 & 0 & -\mu & 0 \\ 0 & 0 & 0 & -\mu \end{pmatrix} \\
R_{mn} - \frac{1}{2} g_{mn} R &= \frac{1}{2 \cdot 3!} G_{mpqr} G_n^{pqr} - \frac{1}{4 \cdot 4!} g_{mn} G^2 \\
d \star G_4 - \frac{1}{2} G_4 \wedge G_4 &= 0. \\
g_{mn} &= \begin{pmatrix} e^{2\omega} \eta_{\mu\nu} & 0 \\ 0 & g_{IJ} \end{pmatrix} \\
\omega = \omega(x^I) &= \omega \left(X^I / T_{M2}^{\frac{1}{2}} \right) g_{IJ} = g_{IJ}(x^I) = g_{IJ} \left(X^I / T_{M2}^{\frac{1}{2}} \right)
\end{aligned}$$

$$\begin{aligned}
S_1 &= -T_{M_2} \int d^3x \sqrt{-\det(e^{2\omega}\eta_{\mu\nu} + \partial_\mu x^I \partial_\nu x^J g_{IJ})} \\
&= -T_{M_2} \int d^3x e^{3\omega} \left(1 + \frac{1}{2} e^{-2\omega} \partial_\mu x^I \partial^\mu x^J g_{IJ} + \dots \right) \\
&= - \int d^3x \left(T_{M_2} e^{3\omega} + \frac{1}{2} e^\omega \partial_\mu X^I \partial^\mu X^J g_{IJ} + \dots \right). \\
e^{2\omega(x)} &= e^{2\omega(X^I/\sqrt{T_{M_2}})} = 1 + \frac{2}{T_{M_2}} \omega_{IJ} X^I X^J + \dots \\
g_{IJ}(x) &= g_{IJ}(X^I/\sqrt{T_{M_2}}) = \delta_{IJ} + \dots \\
S_1 &= - \int d^3x \left(T_{M_2} + 3\omega_{IJ} X^I X^J + \frac{1}{2} \partial_\mu X^I \partial^\mu X^J \delta_{IJ} + \dots \right)
\end{aligned}$$

$$\partial_I \partial^I e^{2\omega} = \frac{1}{3 \cdot 4!} G^2$$

$$e^{2\omega} = 1 + \frac{1}{48 \cdot 4!} G^2 \delta_{IJ} x^I x^J$$

$$S_1 = - \int d^3x \frac{1}{32 \cdot 4!} G^2 X^2$$

$$S_2 = \frac{T_{M_2}}{3!} \int d^3x \epsilon^{\mu\nu\lambda} C_{\mu\nu\lambda}$$

$$\partial_I \partial^I C_0 = \frac{1}{2 \cdot 4!} G^2$$

$$C_0 = \frac{1}{32 \cdot 4!} G^2 \delta_{IJ} x^I x^J$$

$$S_2 = - \int d^3x \frac{1}{32 \cdot 4!} G^2 X^2$$

$$m^2 = \frac{1}{8 \cdot 4!} G^2$$

$$\begin{aligned}
S_{\text{ABJM}} &= \int d^3x \left[\frac{k}{4\pi} e^{\mu\nu\lambda} \text{Tr} \left(A_\mu^L \partial_\nu A_\lambda^L + \frac{2i}{3} A_\mu^L A_\nu^L A_\lambda^L - A_\mu^R \partial_\nu A_\lambda^R - \frac{2i}{3} A_\mu^R A_\nu^R A_\lambda^R \right) - \text{Tr} (D_\mu \bar{Z}_A D^\mu Z^A) \right. \\
&\quad \left. + \frac{4\pi^2}{3k^2} \text{Tr} (Z^A \bar{Z}_A Z^B \bar{Z}_B Z^C \bar{Z}_C + \bar{Z}_A Z^A \bar{Z}_B Z^B \bar{Z}_C Z^C + 4Z^A \bar{Z}_B Z^C \bar{Z}_A Z^B \bar{Z}_C - 6Z^A \bar{Z}_B Z^B \bar{Z}_A Z^C \bar{Z}_C) \right]
\end{aligned}$$

$$V = |M^\alpha|^2 + |N^\alpha|^2$$

$$M^\alpha = \mu Q^\alpha + \frac{2\pi}{k} (2Q^{[\alpha} \bar{Q}_\beta Q^{\beta]} + R^\beta \bar{R}_\beta Q^\alpha - Q^\alpha \bar{R}_\beta R^\beta + 2Q^\beta \bar{R}_\beta R^\alpha)$$

$$N^\alpha = -\mu R^\alpha + \frac{2\pi}{k} (2R^{[\alpha} \bar{R}_\beta R^{\beta]} + Q^\beta \bar{Q}_\beta R^\alpha - R^\alpha \bar{Q}_\beta Q^\beta + 2R^\beta \bar{Q}_\beta Q^\alpha),$$

$$\frac{1}{2} M_B^C \epsilon_{CD} Z^D + ([Z^C, Z^D; Z_B] + [Z^E, Z^C; \bar{Z}_E] \delta_B^D) \epsilon_{CD} = 0$$

$$M_B^C = \begin{pmatrix} 2\mu & 0 & 0 & 0 \\ 0 & 2\mu & 0 & 0 \\ 0 & 0 & -2\mu & 0 \\ 0 & 0 & 0 & -2\mu \end{pmatrix}$$

$$\frac{1}{4} M_B^C Z^D - \frac{1}{4} M_B^D Z^C + [Z^C, Z^D; Z_B] + \frac{1}{2} [Z^E, Z^C; \bar{Z}_E] \delta_B^D - \frac{1}{2} [Z^E, Z^D; \bar{Z}_E] \delta_B^C = 0$$

$$\frac{1}{2} M_B^C Z^B = [Z^D, Z^E; \bar{Z}_E]$$

$$R^\alpha = \frac{2\pi}{\mu k} (R^\alpha \bar{R}_\beta R^\beta - R^\beta \bar{R}_\beta R^\alpha),$$

$$R^\alpha = f G^\alpha$$

$$G^\alpha = G^\alpha \bar{G}_\beta G^\beta - G^\beta \bar{G}_\beta G^\alpha$$



$$\begin{aligned}
(G^1)_{m,l} &= \sqrt{m-1} \delta_{m,l} \\
(G^2)_{m,l} &= \sqrt{(n-m)} \delta_{m+1,l} \\
(\bar{G}_1)_{m,l} &= \sqrt{m-1} \delta_{m,l} \\
(\bar{G}_2)_{m,l} &= \sqrt{(n-l)} \delta_{l+1,m} \\
J_\beta^\alpha &= G^\alpha \bar{G}_\beta \quad \text{and} \quad \bar{J}_\alpha^\beta = \bar{G}_\alpha G^\beta \\
J_i &= (\tilde{\sigma}_i)^\alpha{}_\beta G^\beta \bar{G}_\alpha = (\tilde{\sigma}_i)^\alpha{}_\beta J^\beta{}_\alpha \equiv (\sigma_i)_\beta{}^\alpha J^\beta{}_\alpha \\
\bar{J}_i &= (\tilde{\sigma}_i)^\alpha{}_\beta \bar{G}_\alpha G^\beta = (\tilde{\sigma}_i)^\alpha{}_\beta \bar{J}_\alpha{}^\beta \equiv (\sigma_i)_\beta{}^\alpha \bar{J}_\alpha{}^\beta, \\
[J_i, J_j] &= 2i \epsilon_{ijk} J_k \quad \text{and} \quad [\bar{J}_i, \bar{J}_j] = 2i \epsilon_{ijk} \bar{J}_k \\
J_i G^\alpha - G^\alpha \bar{J}_i &= (\tilde{\sigma}_i)^\alpha{}_\beta G^\beta \\
x_i &\simeq \frac{J_i}{n} \quad \text{and} \quad \bar{x}_i \simeq \frac{\bar{J}_i}{n} \\
g^\alpha &\simeq \frac{G^\alpha}{\sqrt{n}} \quad \text{and} \quad g_\alpha^* \simeq \frac{\bar{G}_\alpha}{\sqrt{n}} \\
x_i &= (\tilde{\sigma}_i)^\alpha{}_\beta g^\beta g_\alpha^* \\
\bar{x}_i &= (\tilde{\sigma}_i)^\alpha{}_\beta g_\alpha^* g^\beta, \\
S &= - \int d^3x \text{Tr}(D_\mu Z^A D^\mu \bar{Z}_A) \\
S_{phys} &= -T_{M2} \int d^3x \text{Tr}(D_\mu Z^A D^{\mu-} \bar{Z}_A) \\
R_{ph}^2 &= \frac{2}{n} \text{Tr}(Z^A \bar{Z}_A) = 8\pi^2 n f^2 \ell_p^3 \\
R_{11} &= \frac{R_{ph}}{k} \\
R_{ph} &= 2n\mu\lambda \\
0 &= \left(\partial_2 X_a^{L'} - \frac{1}{3!} \epsilon^{I'J'K'L'} f^{cd}{}_a X_c^{I'} X_d^{J'} X_b^{K'} \right) \Gamma_2 \Gamma^{L'} \epsilon \\
\frac{dX^{I'}}{dx^2} &= -\frac{1}{3!} \epsilon^{I'J'K'L'} [X^{J'}, X^{K'}, X^{L'}] \\
0 &= \gamma^2 \partial_2 Z^\alpha \epsilon_{\alpha B} + [Z^\gamma, Z^\alpha; \bar{Z}_\gamma] \epsilon_{\alpha B} + [Z^\gamma, Z^\delta; \bar{Z}_B] \epsilon_{\gamma\delta} \\
0 &= \gamma^2 \partial_2 Z^\alpha \epsilon_{\alpha\beta'} + [Z^\gamma, Z^\alpha; \bar{Z}_\gamma] \epsilon_{\alpha\beta'} \\
\frac{dZ^\alpha}{dx^2} &= [Z^\gamma, Z^\alpha; \bar{Z}_\gamma] = \frac{2\pi}{k} (Z^\gamma Z_\gamma^\dagger Z^\alpha - Z^\alpha Z_\gamma^\dagger Z^\gamma) \\
0 &= \gamma^2 \partial_2 Z^\alpha \epsilon_{\alpha\beta} + [Z^\gamma, Z^\alpha; \bar{Z}_\gamma] \epsilon_{\alpha\beta} + [Z^\gamma, Z^\delta; \bar{Z}_\beta] \epsilon_{\gamma\delta} \\
& [Z^\gamma, Z^\delta; \bar{Z}_\beta] \epsilon_{\gamma\delta} = 2[Z^1, Z^2; \bar{Z}_\beta] \epsilon_{12} \\
& = 2\epsilon_{\beta\alpha} [Z^\gamma, Z^\alpha; \bar{Z}_\gamma] \epsilon_{12} \\
& = -2[Z^\gamma, Z^\alpha; \bar{Z}_\gamma] \epsilon_{\alpha\beta}, \\
Z^\alpha &= f(x^2) G^\alpha \\
\frac{df}{dx^2} &= -\frac{2\pi}{k} f^3 \\
f &= \sqrt{\frac{k}{2\pi} \frac{2}{\sqrt{x^2}}} \\
f_d^{abc} &= f_d^{[abc]}, (f_d^{abc})^* = f_d^{abc} \\
f^{ab}{}_{cd} &= -f^{ba}{}_{dc} = -f^{ab}{}_{cd}, (f^{ab}{}_{cd})^* = f^{cd}{}_{ab} \\
\epsilon^{AB} &= -\epsilon^{BA}, \epsilon^{AB} \omega_{AB} = 0 \\
\epsilon_{AB} &= \omega_{AC} \omega_{BD} \epsilon^{CD} \\
(Z_a^A)^* &= \bar{Z}_A^a = -J^{ab} \omega_{AB} Z_b^B \\
(\Psi_{Aa})^* &= \Psi^{Aa} = -J^{ab} \omega^{AB} \Psi_{Bb},
\end{aligned}$$



$$\begin{aligned}
\delta Z_d^A &= i\epsilon^{AD}\Psi_{Dd} \\
\delta\Psi_{Dd} &= \gamma^\mu\epsilon_{AD}D_\mu Z_d^A + h^{abc}{}_d Z_a^A Z_b^B Z_c^C \epsilon_{AB}\omega_{DC} + j^{abc}{}_d Z_a^A Z_b^B Z_c^C \epsilon_{DC}\omega_{AB}, \\
D_\mu Z_d^A &= \partial_\mu Z_d^A - \tilde{A}_\mu{}^a{}_d Z_a^A \\
(h^{abcd})^* &= h_{abcd} = \omega_{ae}\omega_{bf}\omega_{cg}\omega_{dh}h^{efgh} \\
(j^{abcd})^* &= j_{abcd} = \omega_{ae}\omega_{bf}\omega_{cg}\omega_{dh}j^{efgh}, \\
[\delta_1, \delta_2]Z_d^A &= v^\mu D_\mu Z_d^A + \tilde{\Lambda}_d^a Z_a^A \\
\tilde{\Lambda}_d^a &= ih^{abc}{}_d Z_b^B Z_c^C \omega_{DC} \bar{\epsilon}_{[2}^{DF} \epsilon_{1]BF} \\
j^{abc}{}_d &= \frac{1}{2}(h^{bca}{}_d - h^{acb}{}_d) \\
\delta\Psi_{Dd} &= \gamma^\mu\epsilon_{AD}D_\mu Z_d^A + h^{abc}{}_d Z_a^A Z_b^B Z_c^C \epsilon_{AB}\omega_{DC} - h^{acb}{}_d Z_a^A Z_b^B Z_c^C \epsilon_{DC}\omega_{AB} \\
[\delta_1, \delta_2]\Psi_{Dd} &= v^\mu D_\mu \Psi_{Dd} + \tilde{\Lambda}_d^a \Psi_{Da} - \frac{i}{2}\bar{\epsilon}_{[1}^{AC} \epsilon_{2]AD} E_{Cd} + \frac{i}{4}(\bar{\epsilon}_1^{AB} \gamma_\nu \epsilon_{2AB})\gamma^\nu E_{Dd}, \\
E_{Dd} &= \gamma^\mu D_\mu \Psi_{Dd} - h^{abc}{}_d (\Psi_{Dc} Z_a^A Z_b^B + \Psi_{Db} Z_a^A Z_c^C)\omega_{AB} + 2h^{abc}{}_d (\Psi_{Ab} Z_a^A Z_c^C + \Psi_{Ac} Z_a^A Z_b^B)\omega_{DC} = 0. \\
\delta\tilde{A}_\mu{}^a{}_d &= -i(h^{acb}{}_d + h^{abc}{}_d)\omega^{BE}\bar{\epsilon}_{EC}\gamma_\mu\Psi_{Bb}Z_c^C. \\
h^{abc}{}_g (h^{edg}{}_f + h^{egd}{}_f)Z_a^A Z_b^B Z_c^C Z_d^D \omega_{AD}\omega_{BC} &= 0 \\
h^{abc}{}_g (h^{edg}{}_f + h^{egd}{}_f)Z_a^A Z_b^B Z_c^C Z_d^D \bar{\xi}_{AB[1}\gamma^\mu \xi_{2]CD} &= 0. \\
\delta Z_d^A &= i\epsilon^{AD}\Psi_{Dd} \\
\delta\Psi_{Dd} &= \gamma^\mu\epsilon_{AD}D_\mu Z_d^A + f^{ab}{}_{cd} Z_a^A Z_b^B \bar{Z}_c^C \epsilon_{BD} + f^{ab}{}_{cd} Z_a^A Z_b^B \bar{Z}_d^C \epsilon_{AB} \\
\delta\tilde{A}_\mu{}^a{}_d &= -if^{ab}{}_{cd}(\bar{\epsilon}^{BC}\gamma_\mu\Psi_{Bb}\bar{Z}_c^C + \bar{\epsilon}_{BC}\gamma_\mu\Psi^{Cc}Z_b^B) \\
h^{abe}{}_d J_{ce} &= f^{ab}{}_{cd} = -f^{ba}{}_{cd} = -f^{ab}{}_{dc}. \\
h^{abc}{}_d &= g^{acb}{}_d - g^{bca}{}_d, \\
g^{acbd} &= g^{cabd} = g^{bdac}. \\
g^{(acb)d} &= 0 \\
J_{gj}(g^{afbg}g^{jchd} + g^{afgd}g^{hjbc} + g^{afhg}g^{jdbc} + g^{afgc}g^{bjhd}) &= 0. \\
\delta Z_d^A &= i\epsilon^{AD}\Psi_{Dd} \\
\delta\Psi_{Dd} &= \gamma^\mu\epsilon_{AD}D_\mu Z_d^A - g^{abc}{}_d Z_a^A Z_b^B Z_c^C \epsilon_{DB}\omega_{AC} + 2g^{abc}{}_d Z_a^A Z_b^B Z_c^C \epsilon_{AC}\omega_{DB} \\
\delta\tilde{A}_\mu{}^a{}_d &= 3ig^{bca}{}_d \omega^{BE}\bar{\epsilon}_{EC}\gamma_\mu\Psi_{Bb}Z_c^C. \\
\tilde{\Lambda}_d^a &= -\frac{3i}{2}g^{bca}{}_d Z_b^B Z_c^C \omega_{DC} \bar{\epsilon}_{[2}^{DF} \epsilon_{1]BF}. \\
\mathcal{L} &= -D^\mu \bar{Z}_A^a D_\mu Z_a^A - i\bar{\Psi}^{Aa}\gamma^\mu D_\mu \Psi_{Aa} - V + \mathcal{L}_{CS} \\
&\quad - 3ig^{abcd}\omega_{AB}\omega_{CD}(Z_a^A Z_b^B \bar{\Psi}_c^C \Psi_d^D - 2Z_a^A Z_b^B \bar{\Psi}_c^C \Psi_d^D) \\
V &= \frac{12}{5}\bar{\Upsilon}_{ABC}^d \Upsilon_d^{ABC} \\
\Upsilon_d^{ABC} &= g^{abc}{}_d \left(Z_a^A Z_b^B Z_c^C + \frac{1}{4}\omega^{BC} Z_a^A Z_b^D Z_{Dc} \right) \\
\delta Z_d^A &= \tilde{\Lambda}_d^a Z_a^A = g^{bca}{}_d \Lambda_{bc} Z_a^A \\
g_1^{aibjckdl} &= (\delta^{ac}\delta^{bd} - \delta^{ad}\delta^{bc})J^{ijkl} \\
g_2^{aibjckdl} &= (J^{ik}J^{jl} + J^{jk}J^{il})\delta^{ab}\delta^{cd} \\
g_3^{(\pm)aibjckdl} &= (\delta^{ac}\delta^{bd} \pm \delta^{ad}\delta^{bc})(J^{ik}J^{jl} \pm J^{jk}J^{il}) \\
g^{aibjckdl} &= -\frac{2\pi}{k}[g_1^{aibjckdl} - g_2^{aibjckdl}] \\
g^{aibjckdl} &= -\frac{2\pi}{k}[g_3^{(+)aibjckdl} + g_3^{(-)aibjckdl}] \\
g^{aibjckdl} &= -\frac{2\pi}{k}[(\delta^{ac}\delta^{bd} - \delta^{ad}\delta^{bc})J^{ijkl} - \delta^{ab}\delta^{cd}(J^{ik}J^{jl} + J^{jk}J^{il})]
\end{aligned}$$



$$\begin{aligned}
\delta Z^{Adl} &= -\frac{2\pi}{k} [(\delta^{ba}\delta^{cd} - \delta^{bd}\delta^{ca})J^{jk}J^{il} - \delta^{bc}\delta^{ad}(J^{ji}J^{kl} + J^{ki}J^{jl})]\Lambda_{bjck}Z_{ai}^A \\
&= -\frac{2\pi}{k} [(\Lambda_{ajdk} - \Lambda_{djak})J^{jk}J^{il} - \delta^{ad}(J^{ji}J^{kl} + J^{ki}J^{jl})\Lambda_{bjck}]Z_{ai}^A \\
g^{aibjckdl} &= -\frac{2\pi}{k} [J^{ik}J^{jl}\delta^{ac}\delta^{bd} + J^{il}J^{jk}\delta^{ad}\delta^{bc}] \\
g^{aibjckdl} &= -\frac{2\pi}{k} [g_1^{aibjckdl} - g_2^{aibjckdl} + \alpha\varepsilon^{abcd}J^{ij}J^{kl}] \\
g^{aibjckdl} &= -\frac{2\pi}{k} [g_1^{aibjckdl} - g_2^{aibjckdl} + \beta C^{abcd}J^{ij}J^{kl}] \\
C^{abcd} &= \frac{1}{3!}\varepsilon^{abcdefg}C_{efg} \\
g^{aibjckdl} &= -\frac{2\pi}{k} \left[(\delta^{ac}\delta^{bd} - \delta^{ad}\delta^{bc} + \frac{1}{2}C^{abcd})J^{ij}J^{kl} - \delta^{ab}\delta^{cd}(J^{ik}J^{jl} + J^{jk}J^{il}) \right] \\
\delta Z^{Adl} &= g^{bjckaidl}\Lambda_{bjck}Z_{ai}^A \\
&= -\frac{2\pi}{k} \left[J^{il} \left(\delta^{ba}\delta^{cd} - \delta^{bd}\delta^{ca} + \frac{1}{2}C^{bcad} \right) J^{jk}\Lambda_{bjck} - \delta^{ad}(J^{ji}J^{kl} + J^{ki}J^{jl})\Lambda_{bjck} \right] \cdot \\
&\quad \mathcal{P}_{14}^{abcd} = \frac{1}{3} \left(\delta^{ab}\delta^{cd} - \delta^{ac}\delta^{bd} + \frac{1}{2}C^{abcd} \right) \\
&\quad \mathcal{P}_{14}^{abcd}C_{bce} = 0 \\
g^{aibjckdl} &= -\frac{2\pi}{k} \left[\delta^{ab}\delta^{cd}(J^{ik}J^{jl} + J^{jk}J^{il}) - \frac{1}{6}\Gamma_{mn}^{ab}\Gamma_{mn}^{cd}J^{ij}J^{kl} \right], \\
\delta Z^{Adl} &= g^{bjckaidl}\Lambda_{bjck}Z_{ai}^A \\
&= -\frac{2\pi}{k} \left[\delta^{ad}(J^{ji}J^{kl} + J^{ki}J^{jl})\Lambda_{bjck} - \frac{1}{6}J^{jl}\Gamma_{mn}^{ad}\Gamma_{mn}^{bc}J^{jk}\Lambda_{bjck} \right] Z_{ai}^A \\
g^{abcd} &= h_{mn}(\tau^{ma} e^{jbe})(\tau^{nc} f^{jdf}) \\
Z_a^A &= Z_{a1}^A + iZ_{a2}^A \\
\bar{Z}_A^a &= \bar{Z}_A^{a1} - i\bar{Z}_A^{a2} \\
\Xi_{Aa} &= \Psi_{Aa1} + i\Psi_{Aa2} \\
\Xi^{*Aa} &= \Psi^{Aa1} - i\Psi^{Aa2} \\
\bar{Z}_A^{ai} &= -\omega_{AB}J^{ab}\delta^{ij}Z_{bj}^B \\
\Psi^{Aai} &= -\omega^{AB}J^{ab}\delta^{ij}\Psi_{Bbj} \text{ [for Sp}(n) \times \text{SO}(2)] \\
\bar{Z}_A^{ai} &= -\omega_{AB}\delta^{ab}\varepsilon^{ij}Z_{bj}^B \\
\Psi^{Aai} &= -\omega^{AB}\delta^{ab}\varepsilon^{ij}\Psi_{Bbj} \text{ [for SO}(4) \times \text{SU}(2)] \\
\bar{Z}_A^a &= -\omega_{AB}J^{ab}(Z_{b1}^B - iZ_{b2}^B) \\
g^{aibjckdl} &= \frac{4\pi}{3k} [(\delta^{ik}\delta^{jl} - \delta^{il}\delta^{jk})J^{ab}J^{cd} - \delta^{ij}\delta^{kl}(J^{ac}J^{bd} + J^{bc}J^{ad})] \\
f^{ab}{}_{cd} &= -\frac{2\pi}{k} [J^{ab}J_{cd} + (\delta_c^a\delta_d^b - \delta_d^a\delta_c^b)] \\
\delta Z_{dl}^A &= -\omega^{AD}\bar{\eta}\Psi_{Ddl} \\
\delta\Psi_{Ddl} &= -i\gamma^\mu\omega_{AD}\eta D_\mu Z_{dl}^A + if^{ab}{}_{cd}(\omega_{AB}\omega_{CD} - \omega_{AC}\omega_{BD}) \times (\varepsilon_{ik}\varepsilon_{jl} + \varepsilon_{jk}\varepsilon_{il} + i\delta_{ij}\varepsilon_{kl})Z_{ai}^AZ_{bj}^BZ_{ck}^C\eta \\
\delta\tilde{A}_\mu^{aidl} &= if^{abcd}(\bar{\eta}\gamma_\mu\Psi_{Bbj}Z_{ck}^B - \bar{\eta}\gamma_\mu\Psi_{Bck}Z_{bj}^B)(\delta^{jk}\varepsilon^{il} + \varepsilon^{jk}\delta^{il}) \\
&\quad g^{aibjckdl} \rightarrow \alpha\varepsilon^{abcd}\varepsilon^{ij}\varepsilon^{kl} \\
\delta Z_{dl}^D &\rightarrow \alpha\varepsilon^{abcd}\varepsilon^{jk}\varepsilon^{il}\Lambda_{bjck}Z_{ai}^A \\
\bar{Z}_{Aa} &= -i\omega_{AB}\delta_{ab}(Z_{b1}^B - iZ_{b2}^B) \\
g^{aibjckdl} &= \frac{4\pi}{3k}\varepsilon^{abcd}\varepsilon^{ij}\varepsilon^{kl} \\
f^{abcd} &= -\frac{2\pi}{k}\varepsilon^{abcd}
\end{aligned}$$



$$\begin{aligned}
\delta Z_{dl}^A &= -\omega^{AD}\bar{\eta}\Psi_{Ddl} \\
\delta\Psi_{Ddl} &= -i\gamma^\mu\omega_{AD}\eta D_\mu Z_{dl}^A - \frac{4\pi}{k}\varepsilon^{abcd}\omega_{AB}\omega_{CD}\delta_{ik}\delta_{jl}Z_{ai}^AZ_{bj}^BZ_{ck}^C\eta \\
\delta\tilde{A}_\mu{}^{aidl} &= i\frac{4\pi}{k}\varepsilon^{abcd}\varepsilon^{il}\bar{\eta}\gamma_\mu\Psi_{Bbj}Z_{cj}^B \\
\eta_{AA} &= \begin{pmatrix} a & b \\ -b^* & a^* \end{pmatrix}, \eta^{AA} = \varepsilon^{AB}\varepsilon^{\dot{A}\dot{B}}\eta_{B\dot{B}} = \begin{pmatrix} a^* & b^* \\ -b & a \end{pmatrix}. \\
\delta Z_d^A &= i\bar{\eta}^{A\dot{A}}\Psi_{\dot{A}d}. \\
\begin{pmatrix} Z_d^A \\ 0 \end{pmatrix} &\rightarrow Z_d^A, \quad \begin{pmatrix} \bar{Z}_d^{\dot{A}} \\ 0 \end{pmatrix} \rightarrow \bar{Z}_d^{\dot{A}} \\
\begin{pmatrix} 0 \\ \Psi_{\dot{A}d} \end{pmatrix} &\rightarrow \Psi_{Aa}, \quad \begin{pmatrix} 0 \\ \Psi_{\dot{A}a} \end{pmatrix} \rightarrow \Psi^{Aa}, \\
\begin{pmatrix} 0 & \eta^{A\dot{B}} \\ -(\eta^T)^{\dot{A}B} & 0 \end{pmatrix} &\rightarrow \varepsilon^{AB}, \quad \begin{pmatrix} 0 & \eta_{A\dot{B}} \\ -(\eta^T)_{\dot{A}B} & 0 \end{pmatrix} \rightarrow \varepsilon_{AB} \\
\delta Z_d^A &= i\bar{\eta}^{A\dot{D}}\Psi_{\dot{D}d} \\
\delta\Psi_{\dot{D}d} &= \gamma^\mu\eta_{A\dot{D}}D_\mu Z_d^A + f^{ab}{}_{cd}Z_d^AZ_b^B\bar{Z}_A^C\eta_{B\dot{D}} \\
\delta\tilde{A}_\mu{}^a{}_d &= -if^{ab}{}_{cd}(\bar{\eta}^{C\dot{B}}\gamma_\mu\Psi_{\dot{B}b}\bar{Z}_c^C + \bar{\eta}_{B\dot{C}}\gamma_\mu\Psi^{\dot{C}c}Z_b^B) \\
(Z_d^A)^* &= \bar{Z}_A^{\dot{a}} = -J^{ab}\varepsilon_{AB}Z_b^B \\
(\Psi_{\dot{A}a})^* &= \Psi^{\dot{A}a} = -J^{ab}\varepsilon^{\dot{A}\dot{B}}\Psi_{\dot{B}b}, \\
\begin{pmatrix} \varepsilon_{AB} & 0 \\ 0 & \varepsilon_{\dot{A}\dot{B}} \end{pmatrix} &\rightarrow \omega_{AB}, \quad \begin{pmatrix} \varepsilon^{AB} & 0 \\ 0 & \varepsilon^{\dot{A}\dot{B}} \end{pmatrix} \rightarrow \omega^{AB} \\
\delta Z_d^A &= i\bar{\eta}^{A\dot{D}}\Psi_{\dot{D}d} \\
\delta\Psi_{\dot{D}d} &= \gamma^\mu\eta_{A\dot{D}}D_\mu Z_d^A + g^{abc}{}_dZ_d^AZ_b^BZ_c^C\eta_{B\dot{D}}\varepsilon_{AC} \\
\delta\tilde{A}_\mu{}^a{}_d &= -3ig^{bca}{}_d\varepsilon^{\dot{B}\dot{E}}\bar{\eta}_{C\dot{E}}\gamma_\mu\Psi_{\dot{B}b}Z_c^C, \\
\begin{pmatrix} Z_d^A \\ Z_{\dot{a}}^{\dot{A}} \end{pmatrix} &\rightarrow Z_d^A, \quad \begin{pmatrix} \bar{Z}_d^{\dot{A}} \\ Z_{\dot{a}}^{\dot{A}} \end{pmatrix} \rightarrow \bar{Z}_d^{\dot{A}} \\
\begin{pmatrix} \Psi_{\dot{A}a} \\ \Psi_{\dot{A}a} \end{pmatrix} &\rightarrow \Psi_{Aa}, \quad \begin{pmatrix} \Psi_{\dot{A}a} \\ \Psi_{\dot{A}a} \end{pmatrix} \rightarrow \Psi^{Aa}, \\
\delta Z_d^A &= i\bar{\eta}^{A\dot{D}}\Psi_{\dot{D}d} \\
\delta Z_{\dot{a}}^{\dot{A}} &= -i\bar{\eta}^{D\dot{A}}\Psi_{D\dot{a}} \\
\delta\Psi_{\dot{D}d} &= \gamma^\mu\eta_{A\dot{D}}D_\mu Z_d^A + f^{ab}{}_{cd}Z_d^AZ_b^B\bar{Z}_A^C\eta_{B\dot{D}} + f^{\dot{a}b}{}_{c\dot{d}}Z_{\dot{a}}^{\dot{A}}Z_b^B\bar{Z}_A^C\eta_{B\dot{D}} - 2f^{\dot{a}b}{}_{cd}Z_{\dot{a}}^{\dot{A}}Z_b^B\bar{Z}_D^C\eta_{B\dot{A}} \\
\delta\Psi_{D\dot{a}} &= -\gamma^\mu\eta_{D\dot{A}}D_\mu Z_{\dot{a}}^{\dot{A}} - f^{\dot{a}b}{}_{c\dot{d}}Z_{\dot{a}}^{\dot{A}}Z_b^B\bar{Z}_A^C\eta_{D\dot{B}} - f^{a\dot{a}}{}_{cd}Z_d^AZ_b^B\bar{Z}_A^C\eta_{D\dot{B}} + 2f^{a\dot{a}}{}_{cd}Z_d^AZ_b^B\bar{Z}_D^C\eta_{A\dot{B}}, \\
f^{ab}{}_{cd} &= h_{mn}\tau^{ma}{}_c\tau^{nb}{}_d, f^{a\dot{a}}{}_{cd} = f^{\dot{b}a}{}_{dc} = h_{mn}\tau^{ma}{}_c\tau^{nb}{}_d, \\
(Z_d^A)^* &= \bar{Z}_A^{\dot{a}} = -J^{ab}\varepsilon_{AB}Z_b^B \\
(Z_{\dot{a}}^{\dot{A}})^* &= \bar{Z}_{\dot{A}}^{\dot{a}} = -J^{\dot{a}b}\varepsilon_{\dot{A}\dot{B}}Z_b^B, \\
\delta Z_d^A &= i\bar{\eta}^{A\dot{D}}\Psi_{\dot{D}d} \\
\delta Z_{\dot{a}}^{\dot{A}} &= -i\bar{\eta}^{D\dot{A}}\Psi_{D\dot{a}} \\
\delta\Psi_{\dot{D}d} &= \gamma^\mu\eta_{A\dot{D}}D_\mu Z_d^A + g^{abc}{}_dZ_d^AZ_b^BZ_c^C\eta_{B\dot{D}}\varepsilon_{AC} - 3g^{\dot{a}bc}{}_dZ_{\dot{a}}^{\dot{A}}Z_b^BZ_c^C\eta_{C\dot{A}}\omega_{\dot{D}\dot{B}} \\
\delta\Psi_{D\dot{a}} &= -\gamma^\mu\eta_{D\dot{A}}D_\mu Z_{\dot{a}}^{\dot{A}} - g^{\dot{a}bc}{}_dZ_{\dot{a}}^{\dot{A}}Z_b^BZ_c^C\eta_{D\dot{B}}\varepsilon_{\dot{A}\dot{C}} + 3g^{abc}{}_dZ_d^AZ_b^BZ_c^C\eta_{AC}\varepsilon_{D\dot{B}}, \\
g^{abcd} &= h_{mn}J^{eb}J^fd\tau^{ma}{}_e\tau^{ncb}{}_f, g^{\dot{a}bcd} = h_{mn}J^{eb}J^fd\tau^{ma}{}_e\tau^{nc}{}_f, \\
-\frac{1}{4g_{YM}^2}\mathbf{F}^{\mu\nu}\mathbf{F}_{\mu\nu} &\rightarrow \frac{1}{2}\varepsilon^{\mu\nu\lambda}\mathbf{B}_\mu\mathbf{F}_{\nu\lambda} - \frac{1}{2}(D_\mu\boldsymbol{\phi} - g_{YM}\mathbf{B}_\mu)^2 \\
\delta\boldsymbol{\phi} &= g_{YM}\mathbf{M}, \delta\mathbf{B}_\mu = D_\mu\mathbf{M} \\
L &= \text{Tr}\left(\frac{1}{2}\varepsilon^{\mu\nu\lambda}\mathbf{B}_\mu\mathbf{F}_{\nu\lambda} - \frac{1}{2}(D_\mu\boldsymbol{\phi} - g_{YM}\mathbf{B}_\mu)^2 - \frac{1}{2}D_\mu\mathbf{X}^iD^\mu\mathbf{X}^i - \frac{g_{YM}^2}{4}[\mathbf{X}^i,\mathbf{X}^j]^2 + \text{fermions}\right)
\end{aligned}$$



$$\begin{aligned}
& -\frac{1}{2}\hat{D}_\mu X^I \hat{D}^\mu X^I = -\frac{1}{2}(\partial_\mu X^I - [A_\mu, X^I] - g_{YM}^I \mathbf{B}_\mu)^2 \\
\frac{g_{YM}^2}{4} [X^I, X^J]^2 &= \frac{1}{12}(g_{YM}^I [X^J, X^K] + g_{YM}^J [X^K, X^I] + g_{YM}^K [X^I, X^J])^2. \\
& g_{YM}^I \rightarrow X_+^I(x) \\
L_C &= (C_i^\mu - \partial^\mu X_-^I) \partial_\mu X_+^I \\
\delta X_-^I &= \lambda^I, \delta C_\mu^I = \partial_\mu \lambda^I \\
L &= \text{Tr} \left(\frac{1}{2} e^{\mu\nu\lambda} B_\mu F_{\nu\lambda} - \frac{1}{2} \hat{D}_\mu X^I \hat{D}_\mu X^I - \frac{1}{12} (X_+^I [X^J, X^K] + X_+^J [X^K, X^I] + X_+^K [X^I, X^J])^2 \right) \\
& + (C^{\mu I} - \partial^\mu X_-^I) \partial_\mu X_+^I + L_{\text{gauge-fixing}} + \mathcal{L}_{\text{fermions}} \cdot \\
X^{IJK} &\equiv X_+^I [X^J, X^K] + X_+^J [X^K, X^I] + X_+^K [X^I, X^J] \\
f^{+abc} &= f^{abc}, f^{-abc} = f^{+-ab} = f^{abcd} = 0 \\
[T^+, T^{+i}, T_{\bar{m}}^a] &= m^i T_{\bar{m}}^a \\
[T_{\bar{m}}^a, T_{\bar{n}}^b, T_{\bar{p}}^c] &= -i f^{abc} T^- \delta_{\bar{m}+\bar{n}+\bar{p}, \bar{0}} \\
[T^+, T^{+i}, T_{\bar{m}}^a] &= m^i T_{\bar{m}}^a \\
[T^+, T_{\bar{m}}^a, T_{\bar{n}}^b] &= m^i T^{-i} h^{ab} \delta_{\bar{m}, -\bar{n}} + i f^{ab} c T_{\bar{m}+\bar{n}}^c \\
[T_{\bar{m}}^a, T_{\bar{n}}^b, T_{\bar{p}}^c] &= -i f^{abc} T^- \delta_{\bar{m}+\bar{n}+\bar{p}, \bar{0}} \\
\langle T^+, T^- \rangle &= 1 \\
\langle T^{+i}, T^-_j \rangle &= \delta_j^i \\
\langle T_{\bar{m}}^a, T_{\bar{n}}^b \rangle &= h^{ab} \delta_{\bar{m}, -\bar{n}} \\
[X^I, X^{J\dagger}, X^K] &= \frac{1}{3} (X^I X^J X^K - X^I X^K X^J + X^K X^I X^J) \\
\text{STr}(AB^\dagger CD^\dagger) &= \frac{1}{12} \text{Tr} [A(B^\dagger CD^\dagger + B^\dagger DC^\dagger + C^\dagger DB^\dagger + C^\dagger BD^\dagger + D^\dagger BC^\dagger + D^\dagger CB^\dagger) + \text{h.c.}] \\
\text{STr}[X^{IJK} X^{IJL\dagger} X^{MKN} X^{MNL\dagger}] &= 2 \text{STr}[X^{IJM} X^{KLM\dagger} X^{IKN} X^{JLN\dagger}] \\
&= \frac{1}{3} \text{STr}[X^{IJK} X^{IJK\dagger} X^{LMN} X^{LMN\dagger}], \\
X^{IJK} &= X^I X^J X^K \\
(DX)^4: k^2 \text{STr}[\mathbf{a} D^\mu X^I D_\mu X^{J\dagger} D^\nu X^J D_\nu X^{I\dagger} + \mathbf{b} D^\mu X^I D_\mu X^{I\dagger} D^\nu X^J D_\nu X^{J\dagger}] \\
X^{IJK} (DX)^3: k^2 \varepsilon^{\mu\nu\lambda} \text{STr}[\mathbf{c} X^{IJK} D_\mu X^{I\dagger} D_\nu X^J D_\lambda X^{K\dagger}] \\
(X^{IJK})^2 (DX)^2: k^2 \text{STr}[\mathbf{d} X^{IJK} X^{IJK\dagger} D_\mu X^L D^\mu X^{L\dagger} + \mathbf{e} X^{IJK} X^{IJL\dagger} D_\mu X^K D^\mu X^{L\dagger}] \\
(X^{IJK})^4: k^2 \text{STr}[\mathbf{f} X^{IJK} X^{IJK\dagger} X^{LMN} X^{LMN\dagger}], \\
D^\mu X^8 &\rightarrow \frac{1}{v} \mathbf{f}^\mu, \quad D^\mu X^i \rightarrow \frac{1}{v} D^\mu X^i, \quad X^{ij8} \rightarrow -\frac{1}{4v} X^{ij}, \\
D^\mu X^{8\dagger} &\rightarrow -\frac{1}{v} \mathbf{f}^\mu, \quad X^{ijk} \rightarrow O\left(\frac{1}{v^3}\right) \\
D^\mu X^{i\dagger} &\rightarrow -\frac{1}{v} D^\mu X^i, \quad X^{ij8\dagger} \rightarrow \frac{1}{4v} X^{ij}, \\
& X^{ijk\dagger} \rightarrow O\left(\frac{1}{v^3}\right),
\end{aligned}$$

$$S_a^b = \mathbf{a} \left(\frac{k}{v^2}\right)^2 \int d^3x \text{STr} [D^\mu X^i D_\mu X^j D^\nu X^i D_\nu X^j + 2D^\mu X^i D_\nu X^i \mathbf{f}^\mu \mathbf{f}_\nu + \mathbf{f}^\mu \mathbf{f}_\mu \mathbf{f}^\nu \mathbf{f}_\nu]$$

$$S_b^b = \mathbf{b} \left(\frac{k}{v^2}\right)^2 \int d^3x \text{STr} [D^\mu X^i D_\mu X^i D^\nu X^j D_\nu X^j + 2D^\mu X^i D_\mu X^i \mathbf{f}^\nu \mathbf{f}_\nu + \mathbf{f}^\mu \mathbf{f}_\mu \mathbf{f}^\nu \mathbf{f}_\nu]$$

$$S_c^b = \mathbf{c} \left(\frac{k}{v^2}\right)^2 \int d^3x \text{STr} \left[\frac{3}{4} \varepsilon^{\mu\nu\lambda} D_\mu X^i \mathbf{f}_\nu D_\lambda X^j X^{ij} \right]$$

$$S_d^b = \mathbf{d} \left(\frac{k}{v^2}\right)^2 \int d^3x \text{STr} \left[\frac{3}{16} D^\mu X^i D_\mu X^i X^{jk} X^{jk} + \frac{3}{16} \mathbf{f}^\mu \mathbf{f}_\mu X^{ij} X^{ij} \right]$$

$$S_e^b = \mathbf{e} \left(\frac{k}{v^2}\right)^2 \int d^3x \text{STr} \left[\frac{1}{8} D^\mu X^i X^{ij} X^{kj} D_\mu X^k + \frac{1}{16} \mathbf{f}^\mu \mathbf{f}_\mu X^{ij} X^{ij} \right]$$

$$S_f^b = \mathbf{f} \left(\frac{k}{v^2}\right)^2 \int d^3x \text{STr} \left[\frac{9}{256} X^{ij} X^{ji} X^{kl} X^{lk} \right]$$

$$(2\pi)^2 \ell_p^3 \left(\frac{k}{2\pi v^2}\right)^2 = \frac{(2\pi\alpha')^2}{g_{YM}^2}$$

$$S_{\alpha'^2}^b = \frac{(2\pi\alpha')^2}{g_{YM}^2} \int d^3x \text{STr} \left[\frac{1}{4} F_{\mu\nu} F^{\nu\rho} F_{\rho\sigma} F^{\sigma\mu} - \frac{1}{16} F^{\mu\nu} F_{\mu\nu} F^{\rho\sigma} F_{\rho\sigma} - \frac{1}{4} D_\mu X^i D^\mu X^i D_\nu X^j D^\nu X^j \right. \\ \left. + \frac{1}{2} D_\mu X^i D^\nu X^i D_\nu X^j D^\mu X^j + \frac{1}{4} X^{ij} X^{jk} X^{kl} X^{li} - \frac{1}{16} X^{ij} X^{ij} X^{kl} X^{kl} \right. \\ \left. - F_{\mu\nu} F^{\nu\rho} D_\rho X^i D^\mu X^i - \frac{1}{4} F_{\mu\nu} F^{\mu\nu} D_\rho X^i D^\rho X^i - \frac{1}{8} F_{\mu\nu} F^{\mu\nu} X^{kl} X^{kl} \right. \\ \left. - \frac{1}{4} D_\mu X^i D^\mu X^i X^{kl} X^{kl} - X^{ij} X^{jk} D^\mu X^k D_\mu X^i - F_{\mu\nu} D^\nu X^i D^\mu X^j X^{ij} \right]$$

$$\text{STr} [X^{ij} X^{jk} X^{kl} X^{li}] = \frac{1}{2} \text{STr} [X^{ij} X^{ij} X^{kl} X^{kl}]$$

$$\mathbf{a} = \frac{1}{2}, \quad \mathbf{b} = -\frac{1}{4}, \quad \mathbf{c} = \frac{4}{3}$$

$$\mathbf{d} = -\frac{4}{3}, \quad \mathbf{e} = 8, \quad \mathbf{f} = \frac{16}{9}$$

$$\text{STr}(T^a T^b T^c T^d) = mh^{(ab} h^{cd)}$$

$$\text{STr}(T^i T^j T^k T^l) = 2 \text{STr} \left(\frac{\sigma^i \sigma^j \sigma^k \sigma^l}{2 \cdot 2 \cdot 2 \cdot 2} \right) = \frac{1}{4} \delta^{(ij} \delta^{kl)}$$

$$S_{\ell_p^3}^b = (2\pi)^2 \ell_p^3 \int d^3x \text{STr} \left[\frac{1}{4} D^\mu X^I D_\mu X^J D^\nu X^I D_\nu X^J - \frac{1}{8} D^\mu X^I D_\mu X^I D^\nu X^J D_\nu X^J \right.$$

$$\left. + \frac{1}{6} \varepsilon^{\mu\nu\lambda} X^{IJK} D_\mu X^I D_\nu X^J D_\lambda X^K \right.$$

$$\left. + \frac{1}{4} X^{IJK} X^{IJL} D^\mu X^K D_\mu X^L - \frac{1}{24} X^{IJK} X^{IJK} D^\mu X^L D_\mu X^L \right.$$

$$\left. + \frac{1}{288} X^{IJK} X^{IJK} X^{LMN} X^{LMN} \right]$$

$$X^{IJK} = [X^I, X^J, X^K].$$

$$\hat{D}_\mu \hat{X}^I = \partial_\mu \hat{X}^I - [\hat{A}_\mu, \hat{X}^I] - \hat{B}_\mu \hat{X}_+^I$$

$$\hat{X}^{IJK} = \hat{X}_+^I [\hat{X}^J, \hat{X}^K] + \hat{X}_+^J [\hat{X}^K, \hat{X}^I] + \hat{X}_+^K [\hat{X}^I, \hat{X}^J].$$



$$\begin{aligned}
& (\hat{D}\hat{X})^4: \frac{1}{4} \text{STr} \left(\hat{D}^\mu \hat{X}^I \hat{D}_\mu \hat{X}^J \hat{D}^\nu \hat{X}^I \hat{D}_\nu \hat{X}^J - \frac{1}{2} \hat{D}^\mu \hat{X}^I \hat{D}_\mu \hat{X}^I \hat{D}^\nu \hat{X}^J \hat{D}_\nu \hat{X}^J \right) \\
& \hat{X}^{IJK} (\hat{D}\hat{X})^3: \frac{1}{6} \varepsilon^{\mu\nu\lambda} \text{STr} (\hat{X}^{IJK} \hat{D}_\mu \hat{X}^I \hat{D}_\nu \hat{X}^J \hat{D}_\lambda \hat{X}^K) \\
& (\hat{X}^{IJK})^2 (\hat{D}\hat{X})^2: \frac{1}{4} \text{STr} \left(\hat{X}^{IJK} \hat{X}^{IJL} \hat{D}^\mu \hat{X}^K \hat{D}_\mu \hat{X}^L - \frac{1}{6} \hat{X}^{IJK} \hat{X}^{IJK} \hat{D}^\mu \hat{X}^L \hat{D}_\mu \hat{X}^L \right) \\
& (\hat{X}^{IJK})^4: \frac{1}{24} \text{STr} \left(\hat{X}^{IJM} \hat{X}^{KLM} \hat{X}^{IKN} \hat{X}^{JLN} - \frac{1}{12} \hat{X}^{IJK} \hat{X}^{IJK} \hat{X}^{LMN} \hat{X}^{LMN} \right). \\
\text{STr}(\hat{X}^{IJK} \hat{X}^{IJL} \hat{X}^{MKN} \hat{X}^{MNL}) &= \text{STr} \left(\frac{4}{3} \hat{X}^{IJM} \hat{X}^{KLM} \hat{X}^{IKN} \hat{X}^{JLN} + \frac{1}{9} \hat{X}^{IJK} \hat{X}^{IJK} \hat{X}^{LMN} \hat{X}^{LMN} \right) \\
S_{\text{BLG}, \ell_p^3}^b &= \ell_p^3 \int d^3x \text{STr} \left[\frac{1}{4} \left(D^\mu X^I D_\mu X^J D^\nu X^I D_\nu X^J - \frac{1}{2} D^\mu X^I D_\mu X^I D^\nu X^J D_\nu X^J \right) \right. \\
& \quad + \frac{1}{6} \varepsilon^{\mu\nu\lambda} (X^{IJK} D_\mu X^I D_\nu X^J D_\lambda X^K) \\
& \quad + \frac{1}{4} \left(X^{IJK} X^{IJL} D^\mu X^K D_\mu X^L - \frac{1}{6} X^{IJK} X^{IJK} D^\mu X^L D_\mu X^L \right) \\
& \quad \left. + \frac{1}{24} \left(X^{IJM} X^{KLM} X^{IKN} X^{JLN} - \frac{1}{12} X^{IJK} X^{IJK} X^{LMN} X^{LMN} \right) \right] \\
\delta X^I &= i\bar{\epsilon} \Gamma^I \Psi \\
\delta \Psi &= \Gamma^\mu \Gamma_I \partial_\mu X^I \epsilon + \frac{1}{3!} \frac{1}{2} \Gamma^{\mu\nu\lambda} H_{\mu\nu\lambda} \epsilon \\
\delta H_{\mu\nu\lambda} &= 3i\bar{\epsilon} \Gamma_{[\mu\nu} \partial_{\lambda]} \Psi \\
\Gamma^\mu \partial_\mu \Psi &= 0, \partial_\mu \partial^\mu X^I = 0, \partial_{[\mu} H_{\nu\lambda\rho]} = 0 \\
D_\mu X_a^I &= \partial_\mu X_a^I - \tilde{A}_{\mu a}^b X_b^I \\
\delta X_a^I &= i\bar{\epsilon} \Gamma^I \Psi_a \\
\delta \Psi_a &= \Gamma^\mu \Gamma^I D_\mu X_a^I \epsilon + \frac{1}{3!} \frac{1}{2} \Gamma_{\mu\nu\lambda} H_a^{\mu\nu\lambda} \epsilon - \frac{1}{2} \Gamma_\lambda \Gamma^{IJ} C_b^\lambda X_c^I X_d^J f^{cdb}{}_a \epsilon \\
\delta H_{\mu\nu\lambda a} &= 3i\bar{\epsilon} \Gamma_{[\mu\nu} D_{\lambda]} \Psi_a + i\bar{\epsilon} \Gamma^I \Gamma_{\mu\nu\lambda\kappa} C_b^K X_c^I \Psi_d f^{cdb}{}_a \\
\delta \tilde{A}_{\mu a}^b &= i\bar{\epsilon} \Gamma_{\mu\lambda} C_c^\lambda \Psi_d f^{cdb}{}_a \\
\delta C_a^\mu &= 0. \\
D^2 X_a^I &= \frac{i}{2} \bar{\Psi}_c C_B^v \Gamma_v \Gamma^I \Psi_d f_a^{cdb} + C_b^v C_{vg} X_c^J X_e^J X_f^I f^{efg}{}_d f^{cdb}{}_a \\
D_{[\mu} H_{\nu\lambda\rho]a} &= -\frac{1}{4} \varepsilon_{\mu\nu\lambda\rho\sigma\tau} C_b^\sigma X_c^I D^\tau X_d^I f^{cdb}{}_a - \frac{i}{8} \varepsilon_{\mu\nu\lambda\rho\sigma\tau} C_b^\sigma \bar{\Psi}_c \Gamma^\tau \Psi_d f^{cdb}{}_a \\
\Gamma^\mu D_\mu \Psi_a &= -X_c^I C_b^v \Gamma_v \Gamma^I \Psi_d f_a^{cdb} \\
\tilde{F}_{\mu\nu}{}^b{}_a &= -C_c^\lambda H_{\mu\nu\lambda d} f_a^{cdb} \\
& \quad C_c^\rho D_\rho X_b^I f_a^{cdb} = 0, \quad D_\mu C_a^v = 0 \\
& \quad C_c^\rho D_\rho \Psi_d f_a^{cdb} = 0, \quad C_c^\mu C_d^v f_a^{bcd} = 0 \\
& \quad C_c^\rho D_\rho H_{\mu\nu\lambda a} f_a^{cdb} = 0. \\
[\epsilon] &= -\frac{1}{2}, [\Psi] = [X] + \frac{1}{2}, [X] \\
T_{\mu\nu} &= D_\mu X_a^I D_\nu X^{Ia} - \frac{1}{2} \eta_{\mu\nu} D_\lambda X_a^I D^\lambda X^{Ia} \\
& \quad + \frac{1}{4} \eta_{\mu\nu} C_b^\lambda X_a^I X_c^J C_{\lambda g} X_f^I X_e^J f^{cdab} f^{efg}{}_d + \frac{1}{4} H_{\mu\lambda\rho} H_{\nu}{}^{\lambda\rho a} \\
& \quad - \frac{i}{2} \bar{\Psi}_a \Gamma_\mu D_\nu \Psi^a + \frac{i}{2} \eta_{\mu\nu} \bar{\Psi}_a \Gamma^\lambda D_\lambda \Psi^a - \frac{i}{2} \eta_{\mu\nu} \bar{\Psi}_a C_b^\lambda X_c^I \Gamma_\lambda \Gamma^I \Psi_d f^{abcd} \\
& \quad f^{+AB}{}_C = f^{AB}{}_C, f^{ABC} = f^{ABC} \\
& \quad \langle C_A^\lambda \rangle = g \delta_5^\lambda \delta_A^+ \\
& \quad \tilde{F}_{\alpha\beta}{}^B{}_A = -g H_{\alpha\beta 5} f^{DB}{}_A,
\end{aligned}$$



$$\begin{aligned}
0 &= \tilde{D}^\alpha \tilde{D}_\alpha X_A^I - g \frac{i}{2} \tilde{\Psi}_C \Gamma_5 \Gamma^I \Psi_D f^{CD}{}_A - g^2 X_C^J X_E^J X_F^I f^{EF}{}_D f^{CD}{}_A \\
0 &= \tilde{D}_{[\alpha} H_{\beta\gamma]5A} \\
0 &= \tilde{D}^\alpha H_{\alpha\beta 5A} + \frac{1}{2} g f^{CD}{}_A \left(X_C^I \tilde{D}_\beta X_D^I + \frac{i}{2} \tilde{\Psi}_C \Gamma_\beta \Psi_D \right) \\
0 &= \Gamma^\mu \tilde{D}_\mu \Psi_A + g X_C^I \Gamma_5 \Gamma^I \Psi_D f^{CD}{}_A \\
0 &= \partial_5 X_D^I = \partial_5 \Psi_D = \partial_5 H_{\mu\nu\lambda D}, \\
&\quad \delta X_A^I = i\bar{\epsilon} \Gamma^I \Psi_A \\
\delta \Psi_A &= \Gamma^\alpha \Gamma^I \tilde{D}_\alpha X_A^I \epsilon + \frac{1}{2} \Gamma_{\alpha\beta} \Gamma_5 H_A^{\alpha\beta 5} \epsilon - \frac{1}{2} \Gamma_5 \Gamma^{IJ} X_C^I X_D^J f^{CD}{}_A \epsilon \\
\delta \tilde{A}_{\alpha A}^B &= i\bar{\epsilon} \Gamma_\alpha \Gamma_5 \Psi_D f^{DB}{}_C. \\
g &= g_{YM}^2, H_{\alpha\beta 5}^A = -\frac{1}{g_{YM}^2} F_{\alpha\beta}^A, \tilde{A}_{\alpha A}^B = A_{\alpha C} f^{CD}{}_A, \\
0 &= \partial^\mu \partial_\mu X_\pm^I \\
0 &= \partial_{[\mu} H_{\nu\lambda\rho]\pm} \\
0 &= \Gamma^\mu \partial_\mu \Psi_\pm \\
\delta X_\pm^I &= i\bar{\epsilon} \Gamma^I \Psi_\pm \\
\delta \Psi_\pm &= \Gamma^\mu \Gamma^I \partial_\mu X_\pm^I \epsilon + \frac{1}{3!} \frac{1}{2} \Gamma_{\mu\nu\lambda} H_\pm^{\mu\nu\lambda} \epsilon \\
\delta H_{\mu\nu\lambda\pm} &= 3i\bar{\epsilon} \Gamma_{[\mu\nu} \partial_{\lambda]} \Psi_\pm.
\end{aligned}$$

3.1. Agujeros negros cuánticos. En cuanto a los agujeros negros de mecánica cuántica, en esta sección calcularemos su densidad, volumen, diámetro y radio, radiación y termodinámica entrópica – Gibbons - Hawking - Penrose, condiciones inflacionarias, pelo y dualidad AdS/CFT, topología (verbigracia, el grupo de Poincaré, grupo de Lorentz, y parametrización hiperbolide, etc), simetrías y métricas riemanniana, euclídea, einsteniana, teoremas de Gubser-Klevanov-Polyakov y Witten, de Kerr, de Schwarzschild, de Klein – Gordon, Brown-York, de Calabi Yau, etc, así como su masa, tensores de stress – energía, coordenadas de movimiento en un espacio curvo, operadores, generadores y osciladores, el horizonte de eventos y la singularidad, ésta última la cual está dada por la transformación de la información que absorbe el microagujero negro, lo que por ósmosis, supone la formación de superespacios y por ende, de supermembranas infinitas en supersimetrías bosón – fermión, todo esto, en relación a un campo cuántico de gauge curvado o deformado geoméricamente en distorsión e indiscriminado y permeado por un campo de gauge gravitónico. Téngase en cuenta, que el agujero negro cuántico, se forma por el colapso de una superpartícula, sea por el salto de masa que provoca un escalar exponencial de energía, esto es, superior a cero respecto del estado de vacío, o por su naturaleza supermasiva, lo que hace que la partícula – estrella, colapse en sí misma, distorsionando así, el espacio – tiempo cuántico.

El modelo matemático, en un espacio Anti-de Sitter, viene dado por lo que sigue:

$$\begin{aligned}
T_H &= \frac{\hbar c^3}{8\pi k_B GM} \\
S &= \frac{k_B A}{\ell_p^2 4} \\
\langle \Psi | \phi(x_1) \phi(x_2) | \Psi \rangle |_{x_1 \rightarrow x_2} &\approx -\frac{\hbar}{\ell_p^2} \\
S(A) &= k_B \left(\frac{\gamma_c A}{\gamma 4\ell_p^2} - \frac{1}{2} \ln \frac{A}{\ell_p^2} \right) \\
&\quad \square \phi(x) = 0 \\
(f, g) &= -i \int_\Sigma d\Sigma^\mu (f \nabla_\mu g^* - g^* \nabla_\mu f)
\end{aligned}$$



$$\begin{aligned}
[\phi(t_0, \vec{x}), \pi(t_0, \vec{x}')] &= i\hbar\delta(\vec{x} - \vec{x}') \\
\mathcal{S} &= \mathcal{S}_0 \oplus \bar{\mathcal{S}}_0 \\
(f, f) &> 0 \quad \forall f \in \mathcal{S}_0 \\
(f, g^*) &= 0 \quad \forall f, g \in \mathcal{S}_0 \\
a(f) &= (\phi, f) \\
a^\dagger(f) &= -a(f^*) \\
[a(f), a^\dagger(g)] &= \hbar(g, f) \\
[a(f), a(g)] &= -\hbar(g^*, f), [a^\dagger(f), a^\dagger(g)] = -\hbar(g, f^*) \\
[a(f), a(g)] &= 0 = [a^\dagger(f), a^\dagger(g)]. \\
a^\dagger(f_1) \dots a^\dagger(f_n) |0\rangle & \\
a^\dagger(u_{i_1}^{in}) \dots a^\dagger(u_{i_n}^{in}) |in\rangle & \\
u_j^{out} &= \sum_i (\alpha_{ji} u_i^{in} + \beta_{ji} u_i^{in*}) \\
(u_i^{in}, u_j^{in}) &= \delta_{ij} \\
\alpha_{ij} &= (u_i^{out}, u_j^{in}), \beta_{ij} = -(u_i^{out}, u_j^{in*}) \\
\sum_k (\alpha_{ik} \alpha_{jk}^* + \beta_{ik} \beta_{jk}^*) &= \delta_{ij} \\
\sum_k (\alpha_{ik} \beta_{jk} - \beta_{ik} \alpha_{jk}) &= 0 \\
u_i^{in} &= \sum_j (\alpha_{ji}^* u_j^{out} - \beta_{ji} u_j^{out*}) \\
a_i^{in} &= \sum_j (\alpha_{ji} a_j^{out} + \beta_{ji}^* a_j^{out\dagger}) \\
a_i^{out} &= \sum_j (\alpha_{ij}^* a_j^{in} - \beta_{ij}^* a_j^{in\dagger}) \\
N_i^{out} &\equiv \hbar^{-1} a_i^{out\dagger} a_i^{out} \\
\langle in | N_j^{out} | in \rangle &= \hbar^{-1} \langle in | a_j^{out\dagger} a_j^{out} | in \rangle = \\
&= \hbar^{-1} \langle in | \sum_i (-\beta_{ji} a_i^{in}) \sum_k (-\beta_{jk}^* a_k^{in\dagger}) | in \rangle = \\
&= \sum_i |\beta_{ji}|^2. \\
&\sum_k \alpha_{ik} \alpha_{jk}^* = \delta_{ij} \\
\langle in | N_{i_1 i_2}^{out} | in \rangle &= \sum_k \beta_{i_1 k} \beta_{i_2 k}^* = -\sum_k (u_{i_1}^{out}, u_k^{in*}) (u_{i_2}^{out*}, u_k^{in}) = \\
&= \sum_k \left(\int_\Sigma d\Sigma_1^\mu u_{i_1}^{out}(x_1) \overleftrightarrow{\partial}_\mu u_k^{in}(x_1) \right) \left(\int_\Sigma d\Sigma_2^\nu u_{i_2}^{out*}(x_2) \overleftrightarrow{\partial}_\nu u_k^{in*}(x_2) \right). \\
\langle in | \phi(x_1) \phi(x_2) | in \rangle &= \hbar \sum_k u_k^{in}(x_1) u_k^{in*}(x_2) \\
\langle in | N_{i_1 i_2}^{out} | in \rangle &= \hbar^{-1} \int_\Sigma d\Sigma_1^\mu d\Sigma_2^\nu \left[u_{i_1}^{out}(x_1) \overleftrightarrow{\partial}_\mu \right] \left[u_{i_2}^{out*}(x_2) \overleftrightarrow{\partial}_\nu \right] \langle in | \phi(x_1) \phi(x_2) | in \rangle \\
\langle in | N_{i_1 i_2}^{out} | in \rangle &= \hbar^{-1} \int_\Sigma d\Sigma_1^\mu d\Sigma_2^\nu \left[u_{i_1}^{out}(x_1) \overleftrightarrow{\partial}_\mu \right] \left[u_{i_2}^{out*}(x_2) \overleftrightarrow{\partial}_\nu \right] \langle in | : \phi(x_1) \phi(x_2) : | in \rangle \\
(\phi_1, \phi_2) &= -i \int_\Sigma d\Sigma^\mu (\phi_1 \partial_\mu \phi_2^* - \phi_2^* \partial_\mu \phi_1) \\
\phi &= \sum_i (a_i^{in} u_i^{in} + a_i^{in\dagger} u_i^{in*}).
\end{aligned}$$



$$\phi = \sum_i (a_i^{\text{out}} u_i^{\text{out}} + a_i^{\text{out} \dagger} u_i^{\text{out} *}) + (a_i^{\text{int}} u_i^{\text{int}} + a_i^{\text{int} \dagger} u_i^{\text{int} *}).$$

$$u_{wlm}^{\text{in}}|_{I^-} = \frac{1}{\sqrt{4\pi w}} \frac{e^{-i w v}}{r} Y_l^m(\theta, \phi)$$

$$u_{wlm}^{\text{out}}|_{I^+} = \frac{1}{\sqrt{4\pi w}} \frac{e^{-i w u}}{r} Y_l^m(\theta, \phi)$$

$$\beta_{wlm, w'l'm'} = -(u_{wlm}^{\text{out}}, u_{w'l'm'}^{\text{in}*}) = i \int_{I^-}^{I^+} dv r^2 d\Omega (u_{wlm}^{\text{out}} \partial_v u_{w'l'm'}^{\text{in}} - u_{w'l'm'}^{\text{in}} \partial_v u_{wlm}^{\text{out}})$$

$$u_{wlm}^{\text{out}}|_{I^-} = \frac{t_l(w)}{\sqrt{4\pi w}} \frac{e^{-i w u(v)}}{r} Y_l^m(\theta, \phi) \Theta(v_H - v)$$

$$u = v_H - \kappa^{-1} \ln \kappa |v_H - v|$$

$$\beta_{wlm, w'l'm'} = \frac{-(-)^m t_l(w)}{2\pi} \sqrt{\frac{w'}{w}} \int_{-\infty}^{v_H} dv e^{-i w (v_H - \kappa^{-1} \ln \kappa |v_H - v|) - i w' v} \delta_{ll'} \delta_{m-m'}$$

$$\beta_{wlm, w'l'm'} = \frac{-(-)^m t_l(w)}{2\pi \kappa} \sqrt{\frac{w'}{w}} \frac{e^{-i(w+w')v_H}}{(-\kappa^{-1} w' i + \epsilon)^{1+\kappa^{-1} w i}} \Gamma(1 + \kappa^{-1} w i) \delta_{ll'} \delta_{m-m'}$$

$$\int_0^{+\infty} dw' \beta_{w_1 w'} \beta_{w_2 w'}^*$$

$$\int_0^{+\infty} \frac{dw'}{w'} e^{-\kappa^{-1} w_1 i \ln(-\kappa^{-1} w' - i\epsilon)} e^{\kappa^{-1} w_2 i \ln(\kappa^{-1} w' - i\epsilon)} = 2\pi \kappa e^{-\pi \kappa^{-1} w_1} \delta(w_1 - w_2)$$

$$\int_0^{+\infty} dw' \beta_{w_1 w'} \beta_{w_2 w'}^* = \frac{|t_l(w_1)|^2}{e^{2\pi \kappa^{-1} w_1} - 1} \delta(w_1 - w_2)$$

$$\frac{dN_{lm}(w)}{dw dt} \equiv \frac{1}{2\pi} \langle in | N_{wlm}^{\text{out}} | in \rangle = \frac{1}{2\pi} \frac{\Gamma_l(w)}{e^{2\pi \kappa^{-1} w} - 1}$$

$$\frac{1}{e^{\hbar w / k_B T} - 1}$$

$$T_H = \frac{\hbar \kappa}{2\pi k_B}$$

$$u_{wlm}^{\text{out}}(t, r, \theta, \phi)|_{I^+} = u_w^{\text{out}}(u) \frac{Y_l^m(\theta, \phi)}{r}$$

$$\langle in | N_{i_1 i_2}^{\text{out}} | in \rangle = \frac{4}{\hbar} \int_{I^-} r_1^2 dv_1 d\Omega_1 \int_{I^-} r_2^2 dv_2 d\Omega_2 u_{w_1}^{\text{out}} u_{w_2}^{\text{out}*} \times$$

$$\frac{Y_{l_1}^{m_1}(\theta_1, \phi_1) Y_{l_2}^{m_2*}(\theta_2, \phi_2)}{r_1 r_2} \partial_{v_1} \partial_{v_2} \langle in | \phi(x_1) \phi(x_2) | in \rangle.$$

$$\langle in | \phi(x_1) \phi(x_2) | in \rangle = \hbar \int_0^{\infty} dw \sum_{l, m} \frac{e^{-i w v_1} Y_l^m(\theta_1, \phi_1)}{\sqrt{4\pi w} r_1} \frac{e^{i w v_2} Y_l^{m*}(\theta_2, \phi_2)}{\sqrt{4\pi w} r_2}$$

$$u_w^{\text{out}}|_{I^-} = t_l(w) \frac{e^{-i w u(v)}}{\sqrt{4\pi w}} \Theta(v_H - v)$$

$$\partial_{v_1} \partial_{v_2} \int_0^{\infty} dw \frac{e^{-i w (v_1 - v_2)}}{4\pi w} = \lim_{\epsilon \rightarrow 0^+} -\frac{1}{4\pi} \frac{1}{(v_1 - v_2 - i\epsilon)^2}$$

$$\langle in | N_{i_1 i_2}^{\text{out}} | in \rangle = -\frac{t_{l_1}(w_1) t_{l_2}^*(w_2)}{4\pi^2 \sqrt{w_1 w_2}} \int_{-\infty}^{v_H} dv_1 dv_2 \frac{e^{-i w_1 u(v_1) + i w_2 u(v_2)}}{(v_1 - v_2 - i\epsilon)^2} \delta_{l_1 l_2} \delta_{m_1 m_2}$$

$$\langle in | N_{i_1 i_2}^{\text{out}} | in \rangle = -\frac{t_{l_1}(w_1) t_{l_2}^*(w_2)}{4\pi^2 \sqrt{w_1 w_2}} \times$$

$$\times \int_{-\infty}^{\infty} du_1 du_2 \frac{\frac{dv}{du}(u_1) \frac{dv}{du}(u_2)}{[v(u_1) - v(u_2) - i\epsilon]^2} e^{-i w_1 u_1 + i w_2 u_2} \delta_{l_1 l_2} \delta_{m_1 m_2}$$



$$\langle in|N_{i_1 i_2}^{out}|in\rangle = \frac{-t_{l_1}(w_1)t_{l_2}^*(w_2)}{4\pi^2\sqrt{w_1 w_2}} \int_{-\infty}^{+\infty} du_1 du_2 \frac{\left(\frac{\kappa}{2}\right)^2 e^{-iw_1 u_1 + iw_2 u_2}}{\left[\sinh \frac{\kappa}{2}(u_1 - u_2 - i\epsilon)\right]^2} \delta_{l_1 l_2} \delta_{m_1 m_2}$$

$$\langle in|N_{i_1 i_2}^{out}|in\rangle = -\frac{t_{l_1}(w_1)t_{l_2}^*(w_2)\delta(w_1 - w_2)}{2\pi\sqrt{w_1 w_2}} \int_{-\infty}^{+\infty} dz e^{-i\frac{w_1+w_2}{2}z} \frac{\left(\frac{\kappa}{2}\right)^2 \delta_{l_1 l_2} \delta_{m_1 m_2}}{\left[\sinh \frac{\kappa}{2}(z - i\epsilon)\right]^2}$$

$$\langle in|N_{wlm}^{out}|in\rangle = \frac{|t_l(w)|^2}{e^{2\pi\kappa^{-1}w} - 1}$$

$$\langle in|N_{i_1 i_2}^{out}|in\rangle = -\frac{t_{l_1}(w_1)t_{l_2}^*(w_2)}{4\pi^2\sqrt{w_1 w_2}} \int_{-\infty}^{v_H} dv_1 dv_2 e^{-iw_1 u(v_1) + iw_2 u(v_2)} \times$$

$$\left[\frac{1}{(v_1 - v_2)^2} - \frac{\frac{du}{dv}(v_1) \frac{du}{dv}(v_2)}{[u(v_1) - u(v_2)]^2} \right] \delta_{l_1 l_2} \delta_{m_1 m_2}$$

$$v = \frac{au + b}{cu + d}$$

$$\langle in|N_{i_1 i_2}^{out}|in\rangle = -\frac{t_{l_1}(w_1)t_{l_2}^*(w_2)}{4\pi^2\sqrt{w_1 w_2}} \int_{I^+} du_1 du_2 e^{-iw_1 u_1 + iw_2 u_2} \times$$

$$\left[\frac{\frac{dv}{du}(u_1) \frac{dv}{du}(u_2)}{(v(u_1) - v(u_2))^2} - \frac{1}{[u_1 - u_2]^2} \right] \delta_{l_1 l_2} \delta_{m_1 m_2}$$

$$\langle in|N_{i_1 i_2}^{out}|in\rangle = -\frac{t_{l_1}(w_1)t_{l_2}^*(w_2)\delta(w_1 - w_2)}{2\pi\sqrt{w_1 w_2}} \int_{-\infty}^{+\infty} dz e^{-i\frac{w_1+w_2}{2}z} \times$$

$$\left[\frac{\left(\frac{\kappa}{2}\right)^2}{\left(\sinh \frac{\kappa}{2}z\right)^2} - \frac{1}{z^2} \right] \delta_{l_1 l_2} \delta_{m_1 m_2}$$

$$\langle in|N_{wlm}^{out}|in\rangle = -\frac{|t_l(w)|^2}{2\pi w} \int_{-\infty}^{+\infty} dz e^{-iwz} \left[\frac{\left(\frac{\kappa}{2}\right)^2}{\left(\sinh \frac{\kappa}{2}z\right)^2} - \frac{1}{z^2} \right]$$

$$= \frac{|t_l(w)|^2}{e^{2\pi w \kappa^{-1}} - 1}$$

$$\gamma^\mu \nabla_\mu \psi = 0$$

$$(\psi_1, \psi_2) = \int_\Sigma d\Sigma^\mu \bar{\psi}_1 \gamma_\mu \psi_2$$

$$\langle in|N_{i_1 i_2}^{out}|in\rangle = \hbar^{-1} \int_\Sigma d\Sigma_1^\mu d\Sigma_2^\nu [\bar{u}_{i_2}^{out}(x_2) \gamma_\nu]_b [\gamma_\mu u_{i_1}^{out}(x_1)]^a \langle in|\bar{\psi}_a(x_1) \psi^b(x_2)|in\rangle$$

$$u_{w\kappa_j m_j}^{out}(t, r, \theta, \phi) \Big|_{I^+} \sim \frac{e^{-iwu}}{\sqrt{4\pi r}} \begin{pmatrix} \eta(\hat{r})_{\kappa_j}^{m_j} \\ (\hat{r}\vec{\sigma})\eta(\hat{r})_{\kappa_j}^{m_j} \end{pmatrix}$$

$$u_{w\kappa_j m_j}^{out}(t, r, \theta, \phi) \Big|_{I^-} \sim t_{\kappa_j}(w) \frac{e^{-iwu(v)}}{\sqrt{4\pi r}} \begin{pmatrix} \eta(\hat{r})_{\kappa_j}^{m_j} \\ -(\hat{r}\vec{\sigma})\eta(\hat{r})_{\kappa_j}^{m_j} \end{pmatrix} \Theta(v_H - v) \sqrt{\frac{du(v)}{dv}},$$

$$\langle in|\bar{\psi}_a(x_1) \psi^b(x_2)|in\rangle = \hbar \sum_k \bar{v}_{k,a}^{in}(x_1) v_k^{in,b}(x_2)$$

$$v_k^{in} \rightarrow v_{w\kappa_j m}^{in}(t, r, \theta, \phi) \Big|_{I^-} \sim \frac{e^{iww}}{\sqrt{4\pi r}} \begin{pmatrix} \eta(\hat{r})_{\kappa_j}^{m_j} \\ -(\hat{r}\vec{\sigma})\eta(\hat{r})_{\kappa_j}^{m_j} \end{pmatrix}$$



$$\begin{aligned}
\langle in|N_{i_1 i_2}^{out}|in\rangle &= -i \frac{t_{\kappa_{j_1}}(w_1)t_{\kappa_{j_2}}^*(w_2)}{4\pi^2} \delta_{m_{j_1} m_{j_2}} \delta_{\kappa_{j_1} \kappa_{j_2}} \times \\
&\times \int_{-\infty}^{v_H} dv_1 dv_2 \sqrt{\frac{du(v_1)}{dv} \frac{du(v_2)}{dv} \frac{e^{-iw_1 u(v_1)+iw_2 u(v_2)}}{(v_1 - v_2 - i\epsilon)}} \\
\langle in|N_{i_1 i_2}^{out}|in\rangle &= -i \frac{t_{\kappa_{j_1}}(w_1)t_{\kappa_{j_2}}^*(w_2)}{4\pi^2} \delta_{m_{j_1} m_{j_2}} \delta_{\kappa_{j_1} \kappa_{j_2}} \times \\
&\times \int_{-\infty}^{\infty} du_1 du_2 e^{-iw_1 u_1 + iw_2 u_2} \frac{\left(\frac{\kappa}{2}\right)}{\sinh\left[\frac{\kappa}{2}(u_1 - u_2 - i\epsilon)\right]}. \\
\langle in|N_{i_1 i_2}^{out}|in\rangle &= -\frac{i}{2\pi} t_{\kappa_{j_1}}(w_1)t_{\kappa_{j_2}}^*(w_2) \delta_{m_{j_1} m_{j_2}} \delta_{\kappa_{j_1} \kappa_{j_2}} \delta(w_1 - w_2) \times \\
&\times \int_{-\infty}^{+\infty} dz e^{-i\frac{w_1+w_2}{2}z} \frac{\left(\frac{\kappa}{2}\right)}{\sinh\left[\frac{\kappa}{2}(z - i\epsilon)\right]}. \\
\frac{-i}{2\pi} \int_{-\infty}^{+\infty} dz e^{-iwz} \frac{\left(\frac{\kappa}{2}\right)}{\sinh\left[\frac{\kappa}{2}(z - i\epsilon)\right]} &= \frac{1}{e^{2\pi w \kappa^{-1}} + 1} \\
\langle in|N_{w m_j \kappa_j}^{out}|in\rangle &= \frac{|t_{\kappa_j}(w)|^2}{e^{2\pi \kappa^{-1} w} + 1} \\
\langle in|N_{i_1 i_2}^{out}|in\rangle &= -i \frac{t_{\kappa_{j_1}}(w_1)t_{\kappa_{j_2}}(w_2)^*}{4\pi^2} \int_{-\infty}^{v_H} dv_1 dv_2 \sqrt{\frac{du(v_1)}{dv} \frac{du(v_2)}{dv}} \times \\
&e^{-iw_1 u(v_1)+iw_2 u(v_2)} \left[\frac{1}{[v_1 - v_2]} - \frac{\sqrt{\frac{du(v_1)}{dv} \frac{du(v_2)}{dv}}}{[u(v_1) - u(v_2)]} \right] \delta_{\kappa_{j_1} \kappa_{j_2}} \delta_{m_{j_1} m_{j_2}} \\
\langle in|N_{i_1 i_2}^{out}|in\rangle &= -i \frac{t_{\kappa_{j_1}}(w_1)t_{\kappa_{j_2}}(w_2)^*}{4\pi^2} \int_{I^+} du_1 du_2 \sqrt{\frac{dv(u_1)}{du} \frac{dv(u_2)}{du}} \times \\
&e^{-iw_1 u_1 + iw_2 u_2} \left[\frac{\sqrt{\frac{dv(u_1)}{du} \frac{dv(u_2)}{du}}}{[v(u_1) - v(u_2)]} - \frac{1}{[u_1 - u_2]} \right] \delta_{\kappa_{j_1} \kappa_{j_2}} \delta_{m_{j_1} m_{j_2}} \\
\langle in|N_{w \kappa_j m_j}^{out}|in\rangle &= -i \frac{|t_{\kappa_j}(w)|^2}{2\pi} \int_{-\infty}^{+\infty} dz e^{-iwz} \left[\frac{\left(\frac{\kappa}{2}\right)}{\sinh\left(\frac{\kappa}{2}z\right)} - \frac{1}{z} \right] \\
\frac{-i}{2\pi} \int_{-\infty}^{+\infty} dz e^{-iwz} \left[\frac{\left(\frac{\kappa}{2}\right)}{\sinh\left(\frac{\kappa}{2}z\right)} - \frac{1}{z} \right] &= \frac{1}{e^{2\pi w \kappa^{-1}} + 1} \\
b^\mu &= a \left[X \left(\frac{\partial}{\partial T} \right)^\mu + T \left(\frac{\partial}{\partial X} \right)^\mu \right], \\
t &= a^{-1} \tanh^{-1} T/X, \xi = a^{-1} \ln a \sqrt{X^2 - T^2}, Y = Y_0, Z = Z_0, \\
T &= \frac{e^{a\xi}}{a} \sinh at, X = \frac{e^{a\xi}}{a} \cosh at, Y = Y_0, Z = Z_0 \\
(e^{-2a\xi} (-\partial_t^2 + \partial_\xi^2) + \partial_Y^2 + \partial_Z^2) \phi(t, \xi, Y, Z) &= 0 \\
[(-\partial_t^2 + \partial_\xi^2) - e^{2a\xi} (k_Y^2 + k_Z^2) \hbar^{-2}] \phi(t, \xi) &= 0
\end{aligned}$$

$$\begin{aligned}
u_{w, \vec{k}_\perp}^R &= \frac{e^{-iwt}}{2\pi^2 \sqrt{a}} \sinh^2 \left(\frac{\pi w}{a} \right) K_{iw/a} \left(\frac{|\vec{k}_\perp|}{a\hbar} e^{a\xi} \right) e^{i\vec{k}_\perp \cdot \vec{x}_\perp} \\
u_{k_X, \vec{k}_\perp}^M &= \frac{1}{\sqrt{2(2\pi)^3 k_0}} e^{-ik_0 T + i(k_X X + \vec{k}_\perp \cdot \vec{x}_\perp)} \\
\beta_{w\vec{k}_\perp, k'_X \vec{k}'_\perp} &= -[2\pi a k'_0 (e^{2\pi w/a} - 1)]^{-1/2} \left(\frac{k'_0 + k'_X}{k'_0 - k'_X} \right)^{-iw/2a} \delta(\vec{k}_\perp - \vec{k}'_\perp) \\
&\int_{-\infty}^{+\infty} d\vec{k}' \beta_{w_1 \vec{k}_\perp, \vec{k}'} \beta_{w_2 \vec{k}_\perp, \vec{k}'}^* \\
&\int_{-\infty}^{+\infty} dk'_X (2\pi a k'_0)^{-1} \left(\frac{k'_0 + k'_X}{k'_0 - k'_X} \right)^{-i(w_1 - w_2)/a} = \delta(w_1 - w_2) \\
&\int_{-\infty}^{+\infty} d\vec{k}' \beta_{w_1 \vec{k}_\perp, \vec{k}'} \beta_{w_2 \vec{k}_\perp, \vec{k}'}^* = \frac{1}{e^{2\pi w_1/a} - 1} \delta(w_1 - w_2) \delta(\vec{k}_\perp - \vec{k}'_\perp) \\
&g \int d\tau m(\tau) \phi(x(\tau)) \\
&g \int d\tau \langle E\psi | m(\tau) \phi(x(\tau)) | E_0 0_M \rangle \\
&g^2 |\langle E | m(0) | E_0 \rangle|^2 F(E - E_0), \\
F(w) &= \int_{-\infty}^{+\infty} d\tau_1 \int_{-\infty}^{+\infty} d\tau_2 e^{-iw(\tau_1 - \tau_2)} \langle 0_M | \phi(x(\tau_1)) \phi(x(\tau_2)) | 0_M \rangle \\
\dot{F}(w) &= \int_{-\infty}^{+\infty} d\Delta \tau e^{-iw\Delta\tau} \langle 0_M | \phi(x(\tau_1)) \phi(x(\tau_2)) | 0_M \rangle \\
\dot{F}(w) &= - \int_{-\infty}^{+\infty} d\Delta \tau e^{-iw\Delta\tau/\hbar} \left[\frac{\hbar}{4\pi^2 (\Delta\tau - i\epsilon)^2} \right] = -\frac{w}{2\pi} \theta(-w) \\
T &= \frac{1}{a} \sinh a\tau, X = \frac{1}{a} \cosh a\tau \\
\langle 0_M | \phi(x(\tau_1)) \phi(x(\tau_2)) | 0_M \rangle &= \frac{-\hbar \left(\frac{a}{2} \right)^2}{4\pi^2 \sinh^2 \frac{a}{2} (\Delta\tau - i\epsilon)} \\
\dot{F}(w) &= \int_{-\infty}^{+\infty} d\Delta \tau e^{-iw\Delta\tau} \frac{-\hbar \left(\frac{a}{2} \right)^2}{4\pi^2 \sinh^2 \frac{a}{2} (\Delta\tau - i\epsilon)} = \frac{1}{2\pi} \frac{w}{e^{2\pi w/\hbar a} - 1} \\
\dot{F}(w) &= \int_{-\infty}^{+\infty} d\Delta \tau e^{-iw\Delta\tau} [\langle 0_M | \Phi(x(\tau_1)) \Phi(x(\tau_2)) | 0_M \rangle - \langle 0_R | \Phi(x(\tau_1)) \Phi(x(\tau_2)) | 0_R \rangle] \\
\dot{F}(w) &= \int_{-\infty}^{+\infty} d\Delta \tau e^{-iw\Delta\tau} \left[\frac{-\hbar \left(\frac{a}{2} \right)^2}{4\pi^2 \sinh^2 \frac{a}{2} \Delta\tau} + \frac{\hbar}{4\pi^2 (\Delta\tau)^2} \right] = \frac{1}{2\pi} \frac{\hbar w}{e^{2\pi w a} - 1} \\
\langle 0_M | N_{i_1, i_2}^R | 0_M \rangle &= \frac{4}{\hbar} \int_{H_0^+} dv_1 d\vec{x}_{\perp 1} dv_2 d\vec{x}_{\perp 2} u_{i_1}^R(x_1) u_{i_2}^{R*}(x_2) \\
&\partial_{v_1} \partial_{v_2} [\langle 0_M | \phi(x_1) \phi(x_2) | 0_M \rangle - \langle 0_R | \phi(x_1) \phi(x_2) | 0_R \rangle]. \\
\langle 0_M | \partial_{V_1} \phi(x_1) \partial_{V_2} \phi(x_2) | 0_M \rangle &= \frac{1}{4\pi^2} \int d\vec{k}_\perp G_{\vec{k}_\perp}^M(x_1, x_2) e^{i\vec{k}_\perp \cdot (\vec{x}_{\perp 1} - \vec{x}_{\perp 2})} \\
G_{\vec{k}_\perp}^M(x_1, x_2) \Big|_{H_0^+} &= \partial_{V_1} \partial_{V_2} \frac{\hbar}{2\pi} K_0 \left(|\vec{k}_\perp| \sqrt{-(V_1 - V_2)(U_1 - U_2)} \right) \Big|_{H_0^+} \\
&= -\frac{\hbar}{4\pi} \frac{1}{(V_1 - V_2)^2} \\
\langle 0_M | \partial_{V_1} \phi(x_1) \partial_{V_2} \phi(x_2) | 0_M \rangle \Big|_{H_0^+} &= -\frac{\hbar}{4\pi} \frac{1}{(V_1 - V_2)^2} \delta(\vec{x}_{\perp 1} - \vec{x}_{\perp 2})
\end{aligned}$$

$$\begin{aligned}
u_{w, \vec{k}_\perp}^R &= \frac{e^{i\gamma(w, |\vec{k}_\perp|)}}{(2\pi)^{3/2} \sqrt{2w}} (e^{-iwu} + re^{-iwv}) e^{i\vec{k}_\perp \vec{x}_\perp} \\
\langle {}_R 0 | \partial_{v_1} \phi(x_1) \partial_{v_2} \phi(x_2) | 0_R \rangle \Big|_{H^+} &= \frac{1}{4\pi^2} \int d\vec{k}_\perp G_{\vec{k}_\perp}^R(x_1, x_2) \Big|_{H^+} e^{-i\vec{k}_\perp \vec{x}_\perp} \\
&= -\frac{\hbar}{4\pi} \frac{1}{(v_1 - v_2)^2} \delta(\vec{x}_{\perp 1} - \vec{x}_{\perp 2}) \\
\langle {}_M 0 | N_{i_1, i_2}^R | 0_M \rangle &= -\frac{|r|^2}{4\pi^2 \sqrt{w_1 w_2}} \int_{v_1, v_2 > 0} e^{-iw_1 v_1 + iw_2 v_2} \\
&\quad \left[\frac{dV_1 dV_2}{(V_1 - V_2)^2} - \frac{dv_1 dv_2}{(v_1 - v_2)^2} \right] \delta(\vec{k}_{\perp 1} - \vec{k}_{\perp 2}) \\
\langle {}_M 0 | N_{i_1, i_2}^R | 0_M \rangle &= -\frac{|r|^2}{4\pi^2 \sqrt{w_1 w_2}} \int_{-\infty}^{\infty} dv_1 dv_2 e^{-iw_1 v_1} e^{iw_2 v_2} \times \\
&\quad \times \left[\frac{(a/2)^2}{\sinh^2 \frac{a}{2} (v_1 - v_2)} - \frac{1}{(v_1 - v_2)^2} \right] \delta(\vec{k}_{\perp 1} - \vec{k}_{\perp 2}) \\
\langle {}_M 0 | N_{i_1, i_2}^R | 0_M \rangle &= -\frac{1}{4\pi^2 w_1} \int_{-\infty}^{\infty} d(\Delta v) e^{-iw_1 \Delta v} \left[\frac{(a/2)^2}{\sinh^2 \frac{a}{2} \Delta v} - \frac{1}{(\Delta v)^2} \right] \times \\
&\quad \times \delta(w_1 - w_2) \delta(\vec{k}_{\perp 1} - \vec{k}_{\perp 2}) = \frac{1}{e^{2\pi w_1/a} - 1} \delta(\vec{k}_1 - \vec{k}_2) \\
\langle {}_M 0 | N_{i_1, i_2}^R | 0_M \rangle &= -\frac{1}{4\pi^2 \sqrt{w_1 w_2}} \int_{u_1, u_2 < 0} e^{-iw_1 u_1 + iw_2 u_2} \\
&\quad \left[\frac{dU_1 dU_2}{(U_1 - U_2)^2} - \frac{du_1 du_2}{(u_1 - u_2)^2} \right] \delta(\vec{k}_{\perp 1} - \vec{k}_{\perp 2}) \\
&\quad \left(\square - \frac{m^2}{\hbar^2} - \xi R \right) \phi = 0 \\
\dot{F}(w) &= \int_{-\infty}^{+\infty} d\Delta \tau e^{-iw\Delta\tau} \langle \phi(x(\tau_1)) \phi(x(\tau_2)) \rangle \\
\langle \phi(x_1) \phi(x_2) \rangle &= \frac{H^2 \eta_1 \eta_2 \hbar}{4\pi^2 [-(\eta_1 - \eta_2 - i\epsilon)^2 + (\vec{x}_1 - \vec{x}_2)^2]} \\
ds^2 &= \frac{1}{H^2 \eta^2} (-d\eta^2 + d\vec{x}^2) \\
t = \tau \vec{x} &= \vec{x}_0 \\
\dot{F}(w) &= \int_{-\infty}^{+\infty} d\Delta \tau e^{-iw\Delta\tau} \frac{-\hbar H^2}{16\pi^2 \sinh^2 \frac{H}{2} (\Delta\tau - i\epsilon)} \\
\langle 0_C | \phi(x(\tau_1)) \phi(x(\tau_2)) | 0_C \rangle &= -\frac{\hbar}{4\pi^2 (\Delta\tau - i\epsilon)^2} \\
\dot{F}(w) &= \int_{-\infty}^{+\infty} d\Delta \tau e^{-i(w)\Delta\tau/\hbar} \left[\frac{-\hbar \left(\frac{H}{2}\right)^2}{4\pi^2 \sinh^2 \frac{H}{2} \Delta\tau} + \frac{\hbar}{4\pi^2 (\Delta\tau)^2} \right] \\
\frac{dN_{lm}(w)}{dw dt} &= \frac{1}{2\pi} \Gamma_{lm}(w) \frac{1}{e^{2\pi\kappa^{-1}w} \pm 1} \\
\int_0^{+\Lambda} \frac{dw'}{w'} e^{-\kappa^{-1}w_1 i \ln(-\kappa^{-1}w' - i\epsilon)} e^{\kappa^{-1}w_2 i \ln(\kappa^{-1}w' - i\epsilon)} &= e^{-\pi\kappa^{-1}w_1} 2\pi \delta_\sigma[\kappa^{-1}(w_1 - w_2)]
\end{aligned}$$

$$\delta_\sigma[\kappa^{-1}(w_1 - w_2)] = \frac{\sin \left[\frac{\kappa^{-1}(w_1 - w_2)}{\sigma} \right]}{\pi \kappa^{-1}(w_1 - w_2)}$$

$$\sigma = \frac{1}{\ln[\kappa^{-1}\Lambda]}$$

$$u_{jn}^{out} = \frac{1}{\sqrt{\epsilon}} \int_{j\epsilon}^{(j+1)\epsilon} dw e^{2\pi i w n / \epsilon} u_{wlm}^{out}$$

$$\beta_{jn, w'} = \frac{1}{\sqrt{\epsilon}} \int_{j\epsilon}^{(j+1)\epsilon} dw e^{2\pi i w n / \epsilon} \beta_{ww'}$$

$$\langle in | N_{jn}^{out, \sigma} | in \rangle \approx \frac{|t_l(w_j)|^2}{e^{2\pi \kappa^{-1} w_j} - 1} \frac{\sin \left[\left(\frac{2\pi n}{\epsilon} - v_H \right) \frac{\pi \kappa \sigma}{2} \right]}{\left[\left(\frac{2\pi n}{\epsilon} - v_H \right) \frac{\pi \kappa \sigma}{2} \right]}$$

$$I^B(w\kappa^{-1}, \alpha\kappa) = -\frac{1}{2\pi w} \int_{-\alpha}^{+\alpha} dz e^{-i w z} \left[\frac{\left(\frac{\kappa}{2} \right)^2}{\left(\sinh \frac{\kappa}{2} z \right)^2} - \frac{1}{z^2} \right]$$

$$I_{i\epsilon}^B(w\kappa^{-1}, \alpha\kappa) \equiv \frac{-1}{2\pi w} \int_{-\alpha}^{+\alpha} dz e^{-i w z} \frac{\left(\frac{\kappa}{2} \right)^2}{\left[\sinh \frac{\kappa}{2} (z - i\epsilon) \right]^2}$$

$$\frac{-1}{2\pi w} \int_{-\alpha}^{+\alpha} dz e^{-i w z} \frac{1}{(z - i\epsilon)^2}$$

$$I^B(w\kappa^{-1}, \alpha\kappa) = -\frac{Si(\alpha w)}{\pi} - \frac{\kappa}{4\pi w} \{ e^{i\alpha w} ({}_2F_1[1, -i w \kappa^{-1}; 1 - i w \kappa^{-1}; e^{-\alpha\kappa}] - {}_2F_1[1, i w \kappa^{-1}; 1 + i w \kappa^{-1}; e^{\alpha\kappa}]) + e^{-i\alpha w} ({}_2F_1[1, i w \kappa^{-1}; 1 + i w \kappa^{-1}; e^{-\alpha\kappa}] - {}_2F_1[1, -i w \kappa^{-1}; 1 - i w \kappa^{-1}; e^{\alpha\kappa}]) \} + \frac{1}{2\pi \alpha w} \cos(\alpha w) \left[\alpha \kappa \frac{(1 + e^{\alpha\kappa})}{(e^{\alpha\kappa} - 1)} - 2 \right]$$

$$I^B(w\kappa^{-1}, \alpha\kappa) = \left(\frac{1}{12\pi} \alpha \kappa - \frac{1}{720\pi} (\alpha \kappa)^3 + O[(\alpha \kappa)^5] \right) \frac{\kappa}{w}$$

$$- \left(\frac{1}{72\pi} (\alpha \kappa)^3 + O[(\alpha \kappa)^5] \right) \frac{w}{\kappa} + (O[(\alpha \kappa)^5]) \left(\frac{w}{\kappa} \right)^3 + \dots$$

$$\lim_{w\kappa^{-1} \rightarrow 0} \frac{I^B(w\kappa^{-1}, \alpha\kappa)}{(e^{2\pi w \kappa^{-1}} - 1)^{-1}} = \frac{\alpha \kappa}{6} \ll 1$$

$$\frac{I^B(w_{\text{typical}} \kappa^{-1}, \alpha\kappa)}{(e^{2\pi w_{\text{typical}} \kappa^{-1}} - 1)^{-1}} \sim 0.3 \alpha \kappa \ll 1$$

$$\frac{I^B(w\kappa^{-1}, \alpha\kappa)}{(e^{2\pi w \kappa^{-1}} - 1)^{-1}} \approx \frac{\alpha \kappa (e^{2\pi w \kappa^{-1}} - 1)}{12\pi w \kappa^{-1}}$$

$$I_{i\epsilon \rightarrow 0}^B(w\kappa^{-1}, \alpha\kappa) = \frac{e^{\alpha \kappa (1 - i w \kappa^{-1})} + e^{i\alpha w}}{2\pi w \kappa^{-1} (e^{\alpha \kappa} - 1)} + \frac{1}{2\pi (i + w \kappa^{-1})} \{ e^{\alpha \kappa (1 - i w \kappa^{-1})} \times {}_2F_1[1, 1 - i w \kappa^{-1}; 2 - i w \kappa^{-1}; e^{\alpha \kappa}] - e^{-\alpha \kappa (1 - i w \kappa^{-1})} \times {}_2F_1[1, 1 - i w \kappa^{-1}; 2 - i w \kappa^{-1}; e^{-\alpha \kappa}] \}$$

$$I_{i\epsilon}^B(w\kappa^{-1}, \alpha\kappa) = \left(\frac{1}{\pi \alpha \kappa} + \frac{\alpha \kappa}{12\pi} - \frac{(\alpha \kappa)^3}{720\pi} + O((\alpha \kappa)^5) \right) \frac{\kappa}{w} - \frac{1}{2}$$

$$+ \left(\frac{\alpha \kappa}{2\pi} - \frac{(\alpha \kappa)^3}{72\pi} + O((\alpha \kappa)^5) \right) \frac{w}{\kappa} - \left(\frac{(\alpha \kappa)^3}{72\pi} + O((\alpha \kappa)^5) \right) \left(\frac{w}{\kappa} \right)^3 + \dots$$

$$\lim_{\kappa^{-1} w \rightarrow 0} \frac{I_{i\epsilon}^B(w\kappa^{-1}, \alpha\kappa)}{(e^{2\pi w \kappa^{-1}} - 1)^{-1}} = \frac{2}{\alpha \kappa} \gg 1$$

$$\begin{aligned}
I^F(w\kappa^{-1}, \alpha\kappa) &\equiv \frac{-i}{2\pi} \int_{-\alpha}^{+\alpha} dz e^{-iwz} \left[\frac{\left(\frac{\kappa}{2}\right)}{\left(\sinh \frac{\kappa}{2} z\right)^2} - \frac{1}{z} \right] = \\
&\frac{Si(\alpha w)}{\pi} + \frac{1}{2\pi(1+4w^2\kappa^{-2})} \left\{ (-i+2w\kappa^{-1})(e^{-\alpha\kappa/2+i\alpha w} \times \right. \\
&{}_2F_1\left[1, \frac{1}{2}-i w\kappa^{-1}; \frac{3}{2}-i w\kappa^{-1}; e^{-\alpha\kappa}\right] - {}_2F_1\left[1, \frac{1}{2}-i w\kappa^{-1}; \frac{3}{2}-i w\kappa^{-1}; e^{\alpha\kappa}\right] \times \\
&e^{\alpha\kappa/2-i\alpha w} + (i+2w\kappa^{-1}) \left(e^{-\alpha\kappa/2-i\alpha w} {}_2F_1\left[1, \frac{1}{2}+i w\kappa^{-1}; \frac{3}{2}+i w\kappa^{-1}; e^{-\alpha\kappa}\right] \right. \\
&\left. \left. - e^{\alpha\kappa/2+i\alpha w} {}_2F_1\left[1, \frac{1}{2}+i w\kappa^{-1}; \frac{3}{2}+i w\kappa^{-1}; e^{\alpha\kappa}\right] \right) \right\} \\
I^F(w\kappa^{-1}, \alpha\kappa) &= \left(\frac{(\alpha\kappa)^3}{72\pi} + O[(\alpha\kappa)^5] \right) \frac{w}{\kappa} + O[(\alpha\kappa)^5] \left(\frac{w}{\kappa} \right)^3 + \dots \\
&\frac{I^F(w\kappa^{-1}, \alpha\kappa)}{(e^{2\pi w\kappa^{-1}} + 1)^{-1}} \sim \frac{(\alpha\kappa)^3 w}{36\pi \kappa} \ll 1 \\
&\frac{I^F(w_{\text{typical}} \kappa^{-1}, \alpha\kappa)}{(e^{2\pi w_{\text{typical}} \kappa^{-1}} + 1)^{-1}} \sim 3 \cdot 10^{-3} (\alpha\kappa)^3 \ll 1 \\
&\frac{I^F(w\kappa^{-1}, \alpha\kappa)}{(e^{2\pi w\kappa^{-1}} + 1)^{-1}} \approx \frac{\alpha^3 w \kappa^2 (e^{2\pi w\kappa^{-1}} + 1)}{72\pi} \\
&(V_1 - V_2)^2 \sim e^{2av_1} (v_1 - v_2)^2 \\
&-\frac{1}{4\pi^2 w} \int_{|\Delta v| > \alpha} d(\Delta v) e^{-i w \Delta v} \left[\frac{(a/2)^2}{\sinh^2 \frac{a}{2} \Delta v} - \frac{1}{(\Delta v)^2} \right] \\
&\approx \frac{1}{e^{2\pi w/a} - 1} - \frac{\alpha a}{12\pi w/a} + O(\alpha^3 a^3) \\
&(U_1 - U_2)^2 \sim e^{-2au_1} (u_1 - u_2)^2.
\end{aligned}$$

$$\begin{aligned}
E^2 f^2(E, p) - p^2 g^2(E, p) &= m^2 \\
U: p_\mu \equiv (-E, p_i) &\rightarrow \Pi_\mu \equiv (-\epsilon, \Pi_i) = (-f(E, p)E, g(E, p)p_i).
\end{aligned}$$

$$L(p^\mu) = [U^{-1} \quad L \quad U](p^\mu)$$

$$E = \frac{\epsilon}{1 + \ell_p \hbar^{-1} \epsilon}$$

$$p_i = \frac{\Pi_i}{1 + \ell_p \hbar^{-1} \epsilon}$$

$$U: (p_\mu, x^\nu) \rightarrow (\Pi_\mu, X^\nu),$$

$$X^\mu = X^\mu(x^\nu, p_\rho)$$

$$\phi_j(x) \rightarrow \left| \frac{\partial x'}{\partial x} \right|^{\Delta_j/d} \phi_j(x'),$$

$$\langle \phi_j(x_1) \phi_j(x_2) \rangle = \left| \frac{\partial x'}{\partial x} \right|_{x_1}^{\Delta_j/d} \left| \frac{\partial x'}{\partial x} \right|_{x_2}^{\Delta_j/d} \langle \phi_j(x'_1) \phi_j(x'_2) \rangle.$$

$$\langle \phi_j(x_1) \phi_j(x_2) \rangle = \frac{C_j}{(x_1 - x_2)^{2\Delta_j}},$$

$$\langle \phi(x_1) \phi(x_2) \rangle = \frac{1}{4\pi^2} \frac{\hbar}{(x_1 - x_2)^2},$$

$$\langle \partial_\pm \phi(x_1) \partial_\pm \phi(x_2) \rangle = -\frac{1}{4\pi} \frac{\hbar}{(x_1^\pm - x_2^\pm)^2}$$

$$U: \langle \Phi_j(x_1) \Phi_j(x_2) \rangle \rightarrow \langle \phi_j(x_1) \phi_j(x_2) \rangle$$



$$\begin{aligned}
U^{-1}: \langle \Phi_j(x_1) \Phi_j(x_2) \rangle &= \frac{\langle \phi_j(x_1) \phi_j(x_2) \rangle}{1 - \ell_p^2 \hbar^{-1} \langle \phi_j(x_1) \phi_j(x_2) \rangle} \\
\langle \Phi_j(x_1) \Phi_j(x_2) \rangle &\rightarrow \frac{\left| \frac{\partial x'}{\partial x} \right|^{\Delta_j} / d(x_1) \left| \frac{\partial x'}{\partial x} \right|^{\Delta_j/d} (x_2) \langle \phi_j(x'_1) \phi_j(x'_2) \rangle}{1 - \ell_p^2 \hbar^{-1} \left| \frac{\partial x'}{\partial x} \right|^{\Delta_j/d} (x_1) \left| \frac{\partial x'}{\partial x} \right|^{\Delta_j/d} (x_2) \langle \phi_j(x'_1) \phi_j(x'_2) \rangle} \\
\langle \Phi(x_1) \Phi(x_2) \rangle &= \frac{1}{4\pi^2(x_1 - x_2)^2 - \ell_p^2} \frac{\left| \frac{\partial x'}{\partial x} \right|^{1/4} (x_1) \left| \frac{\partial x'}{\partial x} \right|^{1/4} (x_2)}{4\pi^2(x_1 - x_2)^2 - \ell_p^2 \left| \frac{\partial x'}{\partial x} \right|^{1/4} (x_1) \left| \frac{\partial x'}{\partial x} \right|^{1/4} (x_2)} \\
\langle \Phi(x_1) \Phi(x_2) \rangle|_{x_1 \rightarrow x_2} &\approx -\frac{\hbar}{\ell_p^2} = -\frac{1}{G} \\
\langle \Psi | \phi(x_1) \phi(x_2) | \Psi \rangle|_{x_1 \rightarrow x_2} &\sim \langle \phi(x_1) \phi(x_2) \rangle|_{x_1 \rightarrow x_2}, \\
\langle \Psi | \phi(x_1) \phi(x_2) | \Psi \rangle|_{x_1 \rightarrow x_2} &\approx -\frac{\hbar}{\ell_p^2} = -\frac{1}{G} \\
\langle \partial_{\pm} \Phi(x_1) \partial_{\pm} \Phi(x_2) \rangle &= \frac{-\hbar}{4\pi(x_1^{\pm} - x_2^{\pm})^2 + \ell_p^2} \\
\langle \partial_{\pm} \Phi(x_1) \partial_{\pm} \Phi(x_2) \rangle &\rightarrow \frac{\frac{dx'^{\pm}}{dx^{\pm}}(x_1) \frac{dx'^{\pm}}{dx^{\pm}}(x_2) \langle \phi_{\pm}(x'_1) \phi_{\pm}(x'_2) \rangle}{1 - \ell_p^2 \hbar^{-1} \frac{dx^{\pm\pm}}{dx^{\pm}}(x_1) \frac{dx^{\pm\pm}}{dx^{\pm}}(x_2) \langle \phi_{\pm}(x'_1) \phi_{\pm}(x'_2) \rangle} \\
&\rightarrow \frac{-\hbar}{4\pi(x_1^{\pm} - x_2^{\pm})^2 + \ell_p^2} \frac{-\hbar \frac{dx^{\pm}}{dx^{\pm}}(x_1) \frac{dx^{\pm}}{dx^{\pm}}(x_2)}{4\pi(x_1^{\pm} - x_2^{\pm\pm})^2 + \ell_p^2 \frac{dx^{\pm}}{dx^{\pm}}(x_1) \frac{dx'^{\pm}}{dx^{\pm}}(x_2)}, \\
&\rightarrow \frac{-\hbar}{4\pi(x_1^{\pm} - x_2^{\pm})^2 + \ell_p^2} \frac{-\hbar e^{\pm 2\chi}}{4\pi(x_1'^{\pm} - x_2'^{\pm})^2 + \ell_p^2 e^{\pm 2\chi}} \\
\langle \phi(x_1) \phi(x_2) \rangle &= \frac{m\hbar}{4\pi^2 \sqrt{(x_1 - x_2)^2}} K_1 \left(m\sqrt{(x_1 - x_2)^2} \right) \\
\langle \Phi(x_1) \Phi(x_2) \rangle &\equiv \frac{\langle \phi_m(x_1) \phi_m(x_2) \rangle}{1 - \ell_p^2 \hbar^{-1} \langle \phi_m(x_1) \phi_m(x_2) \rangle} \\
&= \frac{m\hbar K_1 \left(m\sqrt{(x_1 - x_2)^2} \right)}{4\pi^2 \sqrt{(x_1 - x_2)^2} - m\ell_p^2 K_1 \left(m\sqrt{(x_1 - x_2)^2} \right)} \\
\langle \Phi(x_1) \Phi(x_2) \rangle|_{x_1 \rightarrow x_2} &\approx -\frac{\hbar}{\ell_p^2} \\
\langle \phi(x_1) \phi(x_2) \rangle &= \int_{-\infty}^{+\infty} ds e^{-im^2 s} K(x_1, x_2; s) = \frac{\hbar}{4\pi^2(x_1 - x_2)^2} \\
\langle \Phi(x_1) \Phi(x_2) \rangle &= \int_{-\infty}^{+\infty} ds e^{-im^2 s} e^{i\ell_p^2/(4\pi^2 s)} K(x_1, x_2; s) \\
\langle 0_M | \Phi(x_1) \Phi(x_2) | 0_M \rangle &= \frac{\hbar}{4\pi^2(x_1 - x_2)^2 - \ell_p^2}
\end{aligned}$$

$$\begin{aligned}
\dot{F}_{\ell_P}(w) &= -\frac{\hbar}{4\pi^2} \int_{-\infty}^{+\infty} d\Delta\tau e^{-i w \Delta\tau} \times \\
&\times \left[\frac{1}{(2/a)^2 \sinh^2 \left(\frac{a}{2} \Delta\tau\right) + \ell_P^2/4\pi^2} - \frac{1}{(\Delta\tau)^2 + \ell_P^2/4\pi^2} \right] \\
\dot{F}_{\ell_P}(w) &= \frac{\hbar}{2\pi} \left[\frac{w e^{\pi w/a}}{(e^{2\pi w/a} - 1)} \frac{\sinh \left[\frac{W}{a}(\theta - \pi)\right]}{\frac{W}{a} \sin \theta} + \frac{\pi e^{-w\ell_P}}{\ell_P} \right] \\
\dot{F}_{\ell_P}(w) &\approx \frac{\hbar}{2\pi} \left[\frac{w}{e^{2\pi w/a} - 1} - \frac{\ell_P a^2}{32\pi} + O(\ell_P/a^2)^3 \right] \\
&- \frac{1}{4\pi^2} \int_{-\infty}^{+\infty} d\Delta\tau e^{-i w \Delta\tau} \frac{\hbar}{(2/a)^2 \sinh^2 \left(\frac{a}{2} \Delta\tau\right) + \ell_P^2/4\pi^2} \\
&\frac{w\hbar}{2\pi} \frac{e^{\pi w/a}}{(e^{2\pi w/a} - 1)} \frac{\sinh \left[\frac{W}{a}(\theta - \pi)\right]}{\frac{W}{a} \sin \theta} \\
U^{-1}: \langle \Phi(x_1) \Phi(x_2) \rangle &= \frac{\langle \phi(x_1) \phi(x_2) \rangle}{1 - \ell_P^2 \hbar^{-1} \langle \phi(x_1) \phi(x_2) \rangle} \\
&= \frac{H^2 \eta_1 \eta_2 \hbar}{4\pi^2 [-(\eta_1 - \eta_2 - i\epsilon)^2 + (\vec{x}_1 - \vec{x}_2)^2] - \ell_P^2 H^2 \eta_1 \eta_2} \\
\dot{F}_{\ell_P}(w) &= -\frac{\hbar}{4\pi^2} \int_{-\infty}^{+\infty} d\Delta\tau e^{-i w \Delta\tau} \times \\
&\times \left[\frac{1}{(2/H)^2 \sinh^2 \left(\frac{H}{2} \Delta\tau\right) + \ell_P^2/4\pi^2} - \frac{1}{(\Delta\tau)^2 + \ell_P^2/4\pi^2} \right] \\
\dot{F}_{\ell_P}(w) &= \frac{\hbar}{2\pi} \left[\frac{w e^{\pi w/H}}{(e^{2\pi w/H} - 1)} \frac{\sinh \left[\frac{W}{H}(\theta - \pi)\right]}{\frac{W}{H} \sin \theta} + \frac{\pi e^{-w\ell_P}}{\ell_P} \right] \\
\dot{F}_{\ell_P}(w) &\approx \frac{\hbar}{2\pi} \left[\frac{w}{e^{2\pi w/H} - 1} - \frac{\ell_P H^2}{32\pi} + O(\ell_P/H^2)^3 \right] \\
\phi(x^\mu) &= \sum_{l,m} \frac{\phi_l(t,r)}{r} Y_{lm}(\theta, \varphi) \\
\left(-\frac{\partial^2}{\partial t^2} + \frac{\partial^2}{\partial r^{*2}} - V_l(r) \right) \phi_l(t,r) &= 0 \\
\langle \partial_u \phi_l(x_1) \partial_u \phi_l(x_2) \rangle &= -\frac{\hbar}{4\pi} \frac{\frac{dv}{du}(u_1) \frac{dv}{du}(u_2)}{[v(u_1) - v(u_2)]^2 + \ell_P^2 \frac{dv}{du}(u_1) \frac{dv}{du}(u_2)}, \\
\langle in | N_{wlm}^{out} | in \rangle &= -\frac{|t_l(w)|^2}{2\pi w} \int_{-\infty}^{+\infty} dz e^{-i w z} \left[\frac{\left(\frac{\kappa}{2}\right)^2}{\left(\sinh \frac{\kappa}{2} z\right)^2 + \left(\frac{\kappa}{2}\right)^2 \frac{\ell_P^2}{4\pi^2}} - \frac{1}{z^2 + \frac{\ell_P^2}{4\pi^2}} \right] \\
\langle in | N_{wlm}^{out} | in \rangle &= |t_l(w)|^2 \left[\frac{e^{\pi w/\kappa}}{(e^{2\pi w/\kappa} - 1)} \frac{\sinh \left[\frac{W}{\kappa}(\theta - \pi)\right]}{\frac{W}{\kappa} \sin \theta} + \frac{\pi e^{-w\ell_P}}{w\ell_P} \right] \\
T_{\mu\nu} &:= -\frac{2}{\sqrt{-g}} \delta S^\phi \\
T_{00} &= \rho = 1/2 \dot{\phi}^2 + V(\phi) \\
T_{ii} &= p g_{ii} = g_{ii} (1/2 \dot{\phi}^2 - V(\phi))
\end{aligned}$$

$$\begin{aligned}
ds^2 &= -dt^2 + a^2(t)d\vec{x}^2 \\
h_{\vec{k}}(t, \vec{x}) &= \sqrt{\frac{16\pi G}{2(2\pi)^3 k^3}} e^{i\vec{k}\vec{x}} (H - ike^{-Ht}) e^{i(kH^{-1}e^{-Ht})} \\
P_t(k) &= 4\Delta_h^2(k) = \frac{8}{M_p^2} \left(\frac{H(t_k)}{2\pi} \right)^2 \\
\int_0^\infty \frac{dk}{k} \frac{\sin(k|\Delta\vec{x}|)}{|\Delta\vec{x}|} \Delta_h^2(k, t) &= \langle 0_{as} | h(t, \vec{x}) h(t, \vec{x} + \Delta\vec{x}) | 0_{as} \rangle \\
\int_0^\infty \frac{dk}{k} \Delta_h^2(k, t) &= \langle 0_{as} | h(t, \vec{x}) h(t, \vec{x}) | 0_{as} \rangle \\
P_{\mathcal{R}}(k) &= \frac{1}{2M_p^2 \epsilon(t_k)} \left(\frac{H(t_k)}{2\pi} \right)^2 \\
n_s - 1 &\equiv d \ln P_{\mathcal{R}} / d \ln k, n_t \equiv d \ln P_t / d \ln k \\
ds^2 &= -dt^2 + a^2(t)d\vec{x}^2 \\
\langle 0_{as} | N_i^c | 0_{as} \rangle &= \sum_j |\beta_{ij}|^2 \\
&= \frac{1}{\hbar} \int_{\Sigma} d\Sigma_1^\mu d\Sigma_2^\nu \left(u_i^c(x_1) \vec{\partial}_\mu \right) \left(u_i^{c*}(x_2) \vec{\partial}_\nu \right) \times \\
&\quad [\langle 0_{as} | \phi(x_1) \phi(x_2) | 0_{as} \rangle - \langle 0_c | \phi(x_1) \phi(x_2) | 0_c \rangle] \\
\dot{F}(w) &= \int_{-\infty}^{+\infty} d\Delta\tau e^{-i w \Delta\tau} [\langle 0_{as} | \phi(x(\tau)) \phi(x(\tau + \Delta\tau)) | 0_{as} \rangle \\
&\quad - \langle 0_c | \phi(x(\tau)) \phi(x(\tau + \Delta\tau)) | 0_c \rangle] \\
\int_0^\infty \frac{dk}{k} \Delta_\phi^2(k, t; dS) &= \langle 0_{as} | \phi(t, \vec{x}) \phi(t, \vec{x}) | 0_{as} \rangle \\
\int_0^\infty \frac{dk}{k} \Delta_\phi^2(k, t) &\equiv \langle 0_{as} | \phi(t, \vec{x}) \phi(t, \vec{x}) | 0_{as} \rangle - \langle 0_c | \phi(t, \vec{x}) \phi(t, \vec{x}) | 0_c \rangle \\
u_{\vec{k}}^{dS}(\vec{x}, t) &= \frac{1}{\sqrt{(2\pi)^3 a(t)^3}} v_k(t) e^{i\vec{k}\vec{x}} \\
v_k(t) &= \frac{1}{2} \sqrt{\frac{\pi}{H}} H_\nu^{(1)}(kH^{-1} \exp(-Ht)) \\
\langle 0_{as} | \phi(t, \vec{x}) \phi(t, \vec{x}) | 0_{as} \rangle &= \hbar (4\pi^2 a(t)^3)^{-1} \int_0^\infty |v_k(t)|^2 k^2 dk \\
\Delta_\phi^2(k, t; dS) &= \hbar (4\pi^2 a(t)^3)^{-1} k^3 |v_k(t)|^2 \\
\Delta_\phi^2(\vec{k}; dS) &= \frac{\hbar H^2}{8\pi} \left| H_\nu^{(1)}(\vec{k}H^{-1}) \right|^2 \approx \frac{\hbar H^2}{2\pi} \\
u_w^c(\vec{r}, \vec{t}) &= \frac{e^{-i w \vec{t}} N_\nu(w)}{\sqrt{4\pi} \vec{r}} \left[P_{\nu-1/2}^{i\frac{w}{H}}(H\vec{r}) - \alpha_\nu(w) Q_{\nu-1/2}^{i\frac{w}{H}}(H\vec{r}) \right] \\
u_w^c(\vec{r} = 0) &= e^{-i w \vec{t}} \frac{\tilde{N}_\nu(w/H)}{\sqrt{4\pi w}} H \beta_\nu(w/H) \\
\beta_\nu(z) &= - \frac{2^{1/2-\nu} \sin(\pi\mu_\nu^*) \sin\left(\frac{\pi\mu_\nu}{2}\right)}{\pi^2 \cosh \pi z} \Gamma(\mu_\nu) [z |\Gamma(\mu_{-\nu}/2)|^2 + \\
2\pi \left(\frac{m^2}{\hbar^2 H^2} - 2 \right) &\mathfrak{I} \left(\frac{\Gamma(\mu_{-\nu}) {}_2F_1\left(\frac{3}{2} - \nu, \frac{1}{2} - \nu + iz; 2 + iz; -1\right)}{\Gamma(2 + iz) \cos \frac{\pi\mu_\nu}{2}} \right) \Big] \\
\langle 0_c | \phi(t, \vec{x}) \phi(t, \vec{x}) | 0_c \rangle &= \frac{\hbar H^2}{4\pi} \int_{\frac{m}{\hbar}}^\infty \frac{dw}{w} |\tilde{N}_\nu|^2 |\beta_\nu|^2
\end{aligned}$$

$$\Delta_{\phi}^2(\bar{k}; C) = \hbar \frac{H^2}{4\pi} \frac{\bar{k}^2}{\bar{k}^2 + m^2 \hbar^{-2}} |\tilde{N}_\nu(\bar{k}H^{-1})|^2 |\beta_\nu(\bar{k}H^{-1})|^2 \left] \right.$$

$\frac{m^2}{H^2 \hbar^2}$	10^{-1}	10^{-3}	10^{-5}
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$$\frac{\Delta_{\phi}^2(k)}{\Delta_{\phi}^2(\bar{k}, dS)} \quad 0.2212 \times 10^{-1} \quad 0.2525 \times 10^{-3} \quad 0.2529 \times 10^{-5}$$

$$|\tilde{N}_{3/2}(w/H)|^2 = \frac{|\Gamma(1 - iw/H)|^2}{\pi} \sinh^2 \frac{\pi w}{2H}$$

$$\Delta_{\phi_{m=0}}^2(\bar{k}; dS) = \frac{\hbar H^2}{4\pi^2} \left(1 + \frac{\bar{k}^2}{H^2} \right) = \Delta_{\phi_{m=0}}^2(\bar{k}; C)$$

$$\Delta_{\phi}^2(\bar{k}) \equiv \Delta_{\phi}^2(\bar{k}; dS) - \Delta_{\phi}^2(\bar{k}; C)$$

$$\Delta_{\phi}^2(\bar{k}) = \hbar H^2 \left[\frac{1}{8\pi} |H_\nu^{(1)}(\bar{k}H^{-1})|^2 - \frac{1}{4\pi} \frac{\bar{k}^2}{\bar{k}^2 + m^2 \hbar^{-2}} |\tilde{N}_\nu(\bar{k}H^{-1})|^2 |\beta_\nu(\bar{k}H^{-1})|^2 \right]$$

$$\left. \frac{\Delta_{\phi}^2(\bar{k})}{\Delta_{\phi}^2(\bar{k}, dS)} \right|_{\bar{k}=H} \approx 0.25 \frac{m^2}{H^2 \hbar^2}$$

$$\dot{v}_{\bar{k}} + (\bar{k}^2 - \ddot{a}/a)v_{\bar{k}} = 0.$$

$$\ddot{v}_{\bar{k}} + (k^2 - (2 + 3\epsilon)/\tau^2)v_{\bar{k}} = 0.$$

$$\Delta_h^2(k, t) \equiv 4\pi k^3 |u_{\bar{k}}(\vec{x}, t)|^2$$

$$\Delta_h^2(k) \approx \left(\frac{H}{2\pi} \right)^2$$

$$\langle h^2(\vec{x}, t) \rangle = \int_0^\infty \frac{dk}{k} \Delta_h^2(k, t)$$

$$\langle h^2(\vec{x}, t) \rangle \approx \int_0^\infty \frac{dk}{k} \frac{16\pi G k^3}{4\pi^2 a^3} \left[\frac{a}{k} \left[1 + \frac{(2 + 3\epsilon)}{2k^2 \tau^2} \right] + \dots \right]$$

$$(16\pi G k^3 / 4\pi^2 a^3) [1/w + \dot{a}^2 / 2a^2 w^3 + \ddot{a} / 2aw^3]$$

$$\langle h^2(\vec{x}, t) \rangle = \int_0^\infty \frac{dk}{k} \frac{16\pi G k^3}{4\pi^2 a^3} \left[-a \frac{\tau\pi}{2} |H_\nu^{(1)}(-k\tau)|^2 - \frac{a}{k} \left[1 + \frac{(2 + 3\epsilon)}{2k^2 \tau^2} \right] \right].$$

$$H_{3/2}^{(1)}(x) = i \exp(ix) \sqrt{2/\pi x} (-1 + i/x)$$

$$\Delta_h^2(k) = 4GH^2 \left[\frac{1 + \epsilon}{2} |H_\nu^{(1)}(-k\tau = 1 + \epsilon)|^2 - \frac{(4 - \epsilon)}{2\pi} \right]$$

$$\Delta_h^2(k) = \alpha 16\pi G (H(t_k)/2\pi)^2 \epsilon(t_k) + O(\epsilon^2)$$

$$P_t(k) = \frac{8\alpha}{M_p^2} \left(\frac{H(t_k)}{2\pi} \right)^2 \epsilon(t_k)$$

$$(k^3 / 4\pi^2 a^3) [m^2 \dot{a}^2 / 2a^2 w^5 + m^2 \ddot{a} / 4aw^5 - 5m^4 \dot{a}^2 / 8a^2 w^7].$$

$$\Delta_{\phi}^2(k) = 4\hbar GH^2 \left[\frac{1 + \epsilon}{2} |H_\nu^{(1)}(-k\tau = 1 + \epsilon)|^2 - \frac{(4 - \epsilon - 5/2m^2/H^2)}{2\pi} \right]$$

$$P_{\mathcal{R}} = \frac{1}{2M_p^2 \epsilon(t_k)} \left(\frac{H(t_k)}{2\pi} \right)^2 (\alpha \epsilon(t_k) + 3\beta \eta(t_k))$$



$$n_t = 2(\epsilon - \eta)$$

$$n_s - 1 = -6\epsilon + 2\eta + \frac{4\alpha\epsilon^2 + (6\beta - 2\alpha)\epsilon\eta - 3\beta(n'_t/2 - 4\epsilon(\epsilon - \eta))}{\alpha\epsilon + 3\beta\eta}$$

$$r = 1 - n_s + \frac{96}{25}n_t + \frac{11}{5}\sqrt{(1 - n_s)^2 + \frac{96}{25}n_t^2}$$

$$r = \frac{\alpha p^2}{(3\beta(p - 1)/2 + \alpha p/4)N}, 1 - n_s = \frac{p}{2N}$$

$$S = k_B \frac{A}{4\ell_p^2}$$

$$a_1 + a_2 + \dots + a_N = 0 \pmod{\kappa}.$$

$$2m_I = -a_I, \pmod{\kappa}.$$

$$A = 8\pi\gamma\ell_p^2 \sum_I \sqrt{j_I(j_I + 1)}$$

$$S = k_B \ln n(A_0)$$

$$A(\vec{j}) := 2 \sum_I \sqrt{j_I(j_I + 1)}$$

$$A_0 - \delta \leq A(\vec{j}) \leq A_0 + \delta,$$

$$\sum_I m_I = 0.$$

$$\sum_{I=1}^N 2\sqrt{|m_I|(|m_I| + 1)} \leq A$$

$$\sum_{I=1}^N m_I = 0, A(\vec{m}) := 2 \sum_{i=1}^N \sqrt{|m_i|(|m_i| + 1)} \leq A_0$$

$$N_{DL}(A) = \sum_{A' < A} d_{DL}(A')$$

$$\sum_{I=1}^N m_I = 0, A(\vec{m}) := 2 \sum_{i=1}^N \sqrt{|m_i|(|m_i| + 1)} = A_0$$

$$A_0 - \delta \leq A(\vec{j}) := \sum_I \sqrt{j_I(j_I + 1)} \leq A_0 + \delta,$$

$$\sum_I m_I = 0.$$

$$N(A) = \theta(A - \sqrt{3})(2N(A - \sqrt{3}) + 2N(A - 2\sqrt{2}) + \dots + 2N(A - 2\sqrt{|m_I|(|m_I| + 1)}) + \dots + 2\lfloor\sqrt{A^2 + 1} - 1\rfloor),$$

$$P(s) := \int_0^\infty N(A)e^{-As} dA$$

$$= 2 \int_{\sqrt{3}}^\infty \lfloor\sqrt{A^2 + 1} - 1\rfloor e^{-As} dA + 2 \int_{\sqrt{3}}^\infty \left(\sum_{k=1}^\infty N(A - \sqrt{k(k+2)}) \right) e^{-As} dA$$

$$= \frac{2}{s} \sum_{k=1}^\infty e^{-s\sqrt{k(k+2)}} + 2 \sum_{k=1}^\infty \int_{\sqrt{3}}^\infty N(A - \sqrt{k(k+2)}) e^{-As} dA$$

$$= \frac{2}{s} \sum_{k=1}^\infty e^{-s\sqrt{k(k+2)}} + 2 \sum_{k=1}^\infty e^{-s\sqrt{k(k+2)}} \int_{-\sqrt{k(k+2)}}^\infty N(A) e^{-As} dA$$

$$= \frac{2}{s} \sum_{k=1}^\infty e^{-s\sqrt{k(k+2)}} + 2P(s) \sum_{k=1}^\infty e^{-s\sqrt{k(k+2)}}$$



$$P(s) = \frac{2 \sum_{k=1}^{\infty} e^{-s\sqrt{k(k+2)}}}{s \left(1 - 2 \sum_{k=1}^{\infty} e^{-s\sqrt{k(k+2)}}\right)}$$

$$N(A) = \sum_{s_i} \text{Res}[P(s)e^{sA}]$$

$$\left(1 - 2 \sum_{k=1}^{\infty} e^{-s\sqrt{k(k+2)}}\right) = 0.$$

$$N(A) \approx C e^{\pi\gamma_M A}$$

$$S = k_B \ln n(A) = k_B \ln (N(A) + 1) = k_B \left(\frac{\gamma_M}{\gamma} \frac{A}{4\ell_p^2} + O(A^0) \right)$$

$$N(A, p) = (1 - \delta(0, p)) \theta(A - 2\sqrt{|p|(|p| + 1)})$$

$$+ \sum_{k=1}^{\infty} (N(A - \sqrt{k(k+2)}, p - k/2) + N_{\leq}(A - \sqrt{k(k+2)}, p + k/2))$$

$$P(s, \omega) := \int_0^{\infty} e^{-sA} \left[\sum_{p \in \mathbb{Z}/2} e^{i\omega p} N(A, p) \right] dA$$

$$= \int_0^{\infty} e^{-sA} \left[\sum_{p \in \mathbb{Z}/2} e^{i\omega p} (1 - \delta(0, p)) \theta(A - 2\sqrt{|p|(|p| + 1)}) \right] dA$$

$$+ \int_0^{\infty} e^{-sA} \left[\sum_{p \in \mathbb{Z}/2} e^{i\omega p} \sum_{\ell \in \mathbb{Z}_*} N\left(A - \sqrt{|\ell|(|\ell| + 2)}, p - \frac{\ell}{2}\right) \right] dA$$

$$= \frac{2}{s} \sum_{k=1}^{\infty} e^{-s\sqrt{k(k+2)}} \cos \frac{k\omega}{2} + 2P(s, \omega) \sum_{k=1}^{\infty} e^{-s\sqrt{k(k+2)}} \cos \frac{k\omega}{2}$$

$$P(s, \omega) = \frac{2}{s} \left(\sum_{k=1}^{\infty} e^{-s\sqrt{k(k+2)}} \cos \frac{k\omega}{2} \right) \left(1 - 2 \sum_{k=1}^{\infty} e^{-s\sqrt{k(k+2)}} \cos \frac{k\omega}{2} \right)^{-1}.$$

$$N(A, 0) = \frac{C}{\sqrt{2\pi\beta_M A}} e^{\pi\gamma_M A}$$

$$S = k_B \left(\frac{\gamma_M}{\gamma} \frac{A}{4\ell_p^2} - \frac{1}{2} \ln (A/\ell_p^2) + O(A^0) \right).$$

$$\sum_{k=1}^{k_{\max}} n_k \sqrt{k(k+2)} = A_0$$

$$d_{GM}(n_1, \dots, n_{k_{\max}}) = \frac{(\sum_k n_k)!}{\prod_k n_k!} \prod_k (k+1)^{n_k},$$

$$\hat{n}_k = \frac{n_k}{\sum_k n_k} = (k+1) e^{-\lambda\sqrt{k(k+2)}}$$

$$\sum_{k=1}^{k_{\max}} (k+1) e^{-\lambda\sqrt{k(k+2)}} = 1$$

$$d_{GM}(A) = \frac{c}{\sqrt{A}} e^{\lambda A}$$

$$S(A) = k_B \left(\frac{\lambda}{\pi\gamma} \frac{A}{4\ell_p^2} - \frac{1}{2} \ln \frac{A}{\ell_p^2} + O(A^0) \right).$$

$$\gamma_{GM} = 0.27398563$$

$$\Delta A = \gamma\chi\ell_p^2$$



$$\chi \approx 8.80$$

$$(S_t, p_t) = (S_0 - 2t, p_0 + 3t)$$

$$3S + 2p = K$$

$$\hat{k}_{md} = \sum_k k \hat{n}_k.$$

$$\frac{k_{md}}{p_{md}} = \hat{k}_{md}.$$

$$S(p, K) = \frac{K}{3} - \frac{2}{3}p.$$

$$\frac{k_{md}(p, K)}{p_{md}(K)} = \frac{\frac{K}{3} - \frac{2}{3}p_{md}(K)}{p_{md}(K)} = \hat{k}_{md},$$

$$p_{md}(K) = \frac{K}{3\hat{k}_{md} + 2}.$$

$$\hat{A}_{md} = 4\pi\gamma\ell_P^2 \sum_k \hat{n}_k \sqrt{k(k+2)}$$

$$A_{md}(K) = p_{md}(K) \hat{A}_{md} = \frac{K \hat{A}_{md}}{3\hat{k}_{md} + 2}$$

$$\Delta A = A_{md}(K+2) - A_{md}(K) = \hat{A}_{md}(p_{md}(K+2) - p_{md}(K)) = \frac{2\hat{A}_{md}}{3\hat{k}_{md} + 2}$$

$$\Delta A_{GM} = \chi_{GM} \gamma_{GM} = \frac{8\pi\gamma_{GM}\ell_P^2 \sum_k \sqrt{k(k+2)}(k+1)e^{-\lambda_{GM}\sqrt{k(k+2)}}}{3\left(\sum_k k(k+1)e^{-\lambda_{GM}\sqrt{k(k+2)}}\right) + 2}$$

$$\Delta A_{DL} = \chi_{DL} \gamma_{DL} = \frac{8\pi\gamma_{DL}\ell_P^2 \sum_k 2\sqrt{k(k+2)}e^{-\lambda_{DL}\sqrt{k(k+2)}}}{3\left(\sum_k 2ke^{-\lambda_{DL}\sqrt{k(k+2)}}\right) + 2}$$

$$\chi_{GM} = 8.789242, \chi_{DL} = 8.784286$$

$$\frac{|\chi_{GM} - 8\ln 3|}{8\ln 3} = 0.000039 = 0.004\%$$

$$\frac{|\chi_{DL} - 8\ln 3|}{8\ln 3} = 0.00052 = 0.05\%$$

$$\frac{|\chi_{GM} - \chi_{DL}|}{\chi_{GM}} = 0.00056 = 0.06\%$$

$$A(\vec{m}) := 2 \sum_I^N \sqrt{|m_I|(|m_I| + 1)} = \sum_I^N \sqrt{(k_I + 1)^2 - 1} = \sum_{k_I}^{k_{\max}} n_k \sqrt{(k+1)^2 - 1},$$

$$\sum_{k_I=1}^{k_{\max}} n_{k_I} \sqrt{(k_I + 1)^2 - 1} = \sum_{i=1}^r q_i \sqrt{p_i}$$

$$\sqrt{(k+1)^2 - 1} = y \sqrt{p_i}$$

$$(k_m^i + 1 + y_m^i \sqrt{p_i}) = (k_i^1 + 1 + y_1^i \sqrt{p_i})^m,$$

$$\sum_{i=1}^r \sum_{m=1}^{\infty} n_{k_m^i} y_m^i \sqrt{p_i} = \sum_{i=1}^r q_i \sqrt{p_i}$$

$$\sum_{m=1}^{\infty} y_m^i n_{k_m^i} = q_i, i = 1, \dots, r$$

$$S_{q_i}^i = \left\{ \alpha_i = \left\{ (n_{k_m^i}) \right\}_{m=1}^{M_i} : \sum_{m=1}^{M_i} y_m^i n_{k_m^i} = q_i \right\}$$



$$\begin{aligned}
R(\alpha) &= \frac{(\sum_{i=1}^r \sum_m n_{k_m^i})!}{\prod_{i=1}^r \prod_m n_{k_m^i}!}, \\
P_{DL}(\alpha) &= \frac{2^N}{M} \sum_{l=0}^{M-1} \prod_{i=1}^N \cos(2\pi l k_i / M) \\
&\quad \prod_{i=1}^N (z^{k_i} + z^{-k_i}) \\
P_{DL}(\alpha) &= \frac{1}{2\pi} \int_0^{2\pi} d\theta \prod_k (2\cos k\theta)^{n_k} \\
d_{DL}(A) &= \sum_{\alpha \in \mathcal{S}_A} R(\alpha) P_{DL}(\alpha) \\
R(\alpha) &= \frac{(\sum_{i=1}^r \sum_m n_{k_m^i})!}{\prod_{i=1}^r \prod_m n_{k_m^i}!} \\
\left[\frac{k_1}{2}\right] \otimes \left[\frac{k_2}{2}\right] &= \bigoplus_{k_3=0}^{\infty} N_{k_1 k_2}^{k_3} \left[\frac{k_3}{2}\right] \\
\chi_{k_1} \chi_{k_2} &= \sum_{k_3} N_{k_1 k_2}^{k_3} \chi_{k_3} \\
\langle \chi_{k_i} | \chi_{k_j} \rangle_{SU(2)} &:= \int_0^{2\pi} \frac{d\theta}{\pi} \sin^2 \theta \chi_{k_i} \chi_{k_j} = \delta_{ij} \\
P_{GM}(\alpha) &= \sum_{j=0}^{\infty} \left\langle \prod_k \chi_k^{n_k} | \chi_{k_j} \right\rangle_{SU(2)} \\
&= \sum_{j=0}^{\infty} \int_0^{2\pi} \frac{d\theta}{\pi} \sin^2 \theta \prod_k \left(\frac{\sin(k+1)\theta}{\sin \theta} \right)^{n_k} \frac{\sin(k_j+1)\theta}{\sin \theta} \\
C_{k+2} &= X C_{k+1} - C_k, k = 0, 1, \dots \\
C_k &= U_k(X/2), k = 0, 1, \dots \\
\left[\frac{k_1}{2}\right] \otimes \left[\frac{k_2}{2}\right] \otimes \dots \otimes \left[\frac{k_N}{2}\right] &= \bigoplus_{k=0}^{\infty} (C_{k_2} C_{k_3} \dots C_{k_N})_{k_1 k} \left[\frac{k}{2}\right] \\
&\quad \sum_{k=0}^{\infty} (C_{k_2} C_{k_3} \dots C_{k_N})_{k_1 k} \\
P_{GM}(\alpha) &= \frac{2}{\pi} \int_0^{\pi} d\theta \cos \frac{\theta}{2} \left[\cos \frac{\theta}{2} - \cos \left(K + \frac{3}{2} \right) \theta \right] \prod_k \left(\frac{\sin(k+1)\theta}{\sin \theta} \right)^{n_k} \\
\langle \eta_i | \eta_j \rangle_{U(1)} &= \frac{1}{2\pi} \int_0^{2\pi} d\theta e^{-im_i \theta} e^{im_j \theta} = \delta_{ij} \\
P(\alpha) &= \frac{1}{2\pi} \int_0^{2\pi} d\theta \prod_k \left(\frac{\sin(k+1)\theta}{\sin \theta} \right)^{n_k} \\
\prod_k \left(\sum_{\alpha=0}^k z^{k-2\alpha} \right)^{n_k} &= \prod_k \left(\frac{z^{k+1} - z^{-k-1}}{z - z^{-1}} \right)^{n_k} \\
P_{GM}(\alpha) &= \frac{1}{2\pi} \int_0^{2\pi} d\theta \prod_k \left(\frac{\sin(k+1)\theta}{\sin \theta} \right)^{n_k} \\
d_{GM}(A) &= \sum_{\alpha \in \mathcal{S}_\alpha} R(\alpha) P_{GM}(\alpha)
\end{aligned}$$



$$\begin{aligned}
\Delta A_{DL} &:= \gamma_{DL} \chi_{DL} \\
\Delta A_{GM} &:= \gamma_{GM} \chi_{DL} \\
\sum_{I=0}^N m_I = 0, A(\vec{j}) &:= 2 \sum_{I=1}^N \sqrt{j_I(j_I + 1)} = A \\
R(A) &= \sum_{\alpha \in \mathcal{S}_A} \frac{(\sum_i^r \sum_m n_{k_m^i})!}{\prod_{i=1}^r \prod_m n_{k_m^i}!} \\
\sum_{m=1}^{\infty} n_{k_m^i} y_m^i &= q_i, i = 1, \dots, r \\
&\sum_{\alpha \in \mathcal{S}_A} \frac{1}{\prod_{i=1}^r \prod_m n_{k_m^i}!} \\
&n_{k_1^i} y_1^i + n_{k_2^i} y_2^i + \dots = q_i \\
f(x) &= (x^{0 \cdot y_1^i} + x^{1 \cdot y_1^i} + x^{2 \cdot y_1^i} + \dots) (x^{0 \cdot y_2^i} + x^{1 \cdot y_2^i} + x^{2 \cdot y_2^i} + \dots) \dots, \\
f(x) &= \frac{1}{1 - x y_1^i} \frac{1}{1 - x y_2^i} \dots = \prod_{m=1}^{\infty} \frac{1}{1 - x y_m^i} \\
f(x_1, x_2, \dots) &= \prod_{i=1}^{\infty} \prod_{m=1}^{\infty} \frac{1}{1 - x_i y_m^i} \\
\sum_{\alpha \in \mathcal{S}_A} 1 &= [x_1^{q_1}, x_2^{q_2}, \dots, x_r^{q_r}] f(x_1, x_2, \dots) \\
f(x) \rightarrow g(x) &= \left(\frac{1}{0!} x^{0 \cdot y_1^i} + \frac{1}{1!} x^{1 \cdot y_1^i} + \frac{1}{2!} x^{2 \cdot y_1^i} + \dots \right) \left(\frac{1}{0!} x^{0 \cdot y_2^i} + \frac{1}{1!} x^{1 \cdot y_2^i} + \dots \right) \dots \\
g(x) &= \prod_{m=1}^{\infty} e^{x y_m^i} \\
\sum_{\alpha \in \mathcal{S}_A} \frac{1}{\prod_{i=1}^r \prod_m n_{k_m^i}!} &= [x_1^{q_1}, x_2^{q_2}, \dots, x_r^{q_r}] g(x_1, x_2, \dots) \\
g(x_1, x_2, \dots) &= \prod_{i=1}^{\infty} \prod_{m=1}^{\infty} e^{x_i y_m^i} = \exp \left(\sum_{i=1}^{\infty} \sum_{m=1}^{\infty} x_i y_m^i \right) \\
\exp \left(\sum_{i=1}^{\infty} \sum_{m=1}^{\infty} x_i y_m^i \right) &= 1 + \left(\sum_{i=1}^{\infty} \sum_{m=1}^{\infty} x_i y_m^i \right) + \frac{1}{2!} \left(\sum_{i=1}^{\infty} \sum_{m=1}^{\infty} x_i y_m^i \right)^2 + \dots, \\
G(x_1, x_2, \dots) &= 0! 1 + 1! \left(\sum_{i=1}^{\infty} \sum_{m=1}^{\infty} x_i y_m^i \right) + 2! \frac{1}{2!} \left(\sum_{i=1}^{\infty} \sum_{m=1}^{\infty} x_i y_m^i \right)^2 + \dots \\
&= \frac{1}{1 - \sum_{i=1}^{\infty} \sum_{m=1}^{\infty} x_i y_m^i} \\
R(A) &= [x_1^{q_1}, x_2^{q_2}, \dots, x_r^{q_r}] G(x_1, x_2, \dots) \\
&= [x_1^{q_1}, x_2^{q_2}, \dots, x_r^{q_r}] G_N(x_1, x_2, \dots), \\
G_N(x_1, x_2, \dots) &= \left(\sum_{i=1}^{\infty} \sum_{m=1}^{\infty} x_i y_m^i \right)^N.
\end{aligned}$$

$$\begin{aligned}
G(x_1, x_2, \dots) &= \sum_{j=1}^{\infty} G_j(x_1, x_2, \dots) = 1 + \left(\sum_{i=1}^{\infty} \sum_{m=1}^{\infty} x_i^{y_m^i} \right) + \left(\sum_{i=1}^{\infty} \sum_{m=1}^{\infty} x_i^{y_m^i} \right)^2 + \dots = \\
&= \frac{1}{\left(1 - \sum_{i=1}^{\infty} \sum_{m=1}^{\infty} x_i^{y_m^i} \right)} \\
&\quad \prod_{l=1}^N (z^{k_l} + z^{-k_l}), \\
&\quad \prod_{l=1}^N \sum_{\alpha=0}^{k_l} z^{k_l - 2\alpha} = \prod_{l=1}^N \frac{z^{k_l+1} - z^{-k_l-1}}{z - z^{-1}} \\
G^{DL}(z, x_1, x_2, \dots) &= \left(1 - \sum_{i=1}^{\infty} \sum_{m=1}^{\infty} (z^{k_m^i} + z^{-k_m^i}) x_i^{y_m^i} \right)^{-1} \\
G^{GM}(z, x_1, x_2, \dots) &= \left(1 - \sum_{i=1}^{\infty} \sum_{m=1}^{\infty} \left(\sum_{\alpha=0}^{k_m^i} z^{k_m^i - 2\alpha} \right) x_i^{y_m^i} \right)^{-1} \\
d_{DL}(A) &= [z^0, x_1^{q_1}, x_2^{q_2}, \dots, x_r^{q_r}] G^{DL}(x_1, x_2, \dots) \\
d_{GM}(A) &= [z^0, x_1^{q_1}, x_2^{q_2}, \dots, x_r^{q_r}] G^{GM}(x_1, x_2, \dots) \\
G^{DL}(z, x_1, x_2) &= \\
&= \frac{1}{1 - (z^2 + z^{-2})x_1^2 - (z^{16} + z^{-16})x_1^{12} - (z + z^{-1})x_2 - (z^6 + z^{-6})x_2^4 - (z^{25} + z^{-25})x_2^{15}}. \\
d_{DL} &= 991809938488860909241077458398212 \\
d_{DL}(A) &= \sum_{n \in \mathbb{N}} d_{DL}(A_n) \delta(A - An) \\
n(A) &= 1 + \int_0^A dA d_{DL}(A) \\
\mathcal{L}[\delta(A - A_n); s] &= e^{-A_n s} \\
\mathcal{L}[d_{DL}(A); s] &= \sum_{n \in \mathbb{N}} d_{DL}(A_n) e^{-A_n s} \\
d_{DL}(A) &= \mathcal{L}^{-1} \left[\sum_{n \in \mathbb{N}} d_{DL}(A_n) e^{-A_n s}; A \right] \\
\int_0^A F(A) dA &= \mathcal{L}^{-1} \left[\frac{1}{s} f(s); A \right] \\
\int_0^A d_{DL}(A) dA &= \mathcal{L}^{-1} \left[\frac{1}{s} \sum_{n \in \mathbb{N}} d_{DL}(A_n) e^{-A_n s}; A \right] \\
G^{DL}(1, e^{-s\sqrt{p_1}}, e^{-s\sqrt{p_2}}, \dots) &= \sum_{n \in \mathbb{N}} d_{DL}(A_n) e^{-A_n s} + 1 \\
x_1^{q_1}, x_2^{q_2} \dots x_r^{q_r} &\rightarrow e^{-s(q_1\sqrt{p_1} + q_2\sqrt{p_2} + \dots + q_r\sqrt{p_r})} = e^{-sA_n} \\
G^{DL}(1, e^{-s\sqrt{p_1}}, e^{-s\sqrt{p_2}}, \dots) &= \left(1 - 2 \sum_{i=1}^{\infty} \sum_{m=1}^{\infty} e^{-s y_m^i \sqrt{p_i}} \right)^{-1} \\
y_m^i \sqrt{p_i} &= \sqrt{k_m^i (k_m^i + 2)} \\
GDL(1, e^{-s\sqrt{p_1}}, e^{-s\sqrt{p_2}}, \dots) &= \left(1 - 2 \sum_{k=1}^{\infty} e^{-s \sqrt{k_m^i (k_m^i + 2)}} \right)^{-1}
\end{aligned}$$

$$\int_0^A d_{DL}(A)dA = \mathcal{L}^{-1} \left[\frac{2 \sum_{k=1}^{\infty} e^{-s \sqrt{k_m^i(k_m^i+2)}}}{s \left(1 - 2 \sum_{k=1}^{\infty} e^{-s \sqrt{k_m^i(k_m^i+2)}} \right)} ; A \right]$$

$$n(A) = 1 + \int_0^A d_{DL}(A)dA = \mathcal{L}^{-1} \left[\frac{1}{s \left(1 - 2 \sum_{k=1}^{\infty} e^{-s \sqrt{k_m^i(k_m^i+2)}} \right)} ; A \right]$$

$$[z^0]G_{DL}(z, x_1, x_2, \dots) = \frac{1}{4\pi} \int_0^{4\pi} dw G(e^{iw/2}, x_1, x_2, \dots)$$

$$G(1, x_1, x_2, \dots) \rightarrow [z^0]G(z, x_1, x_2, \dots).$$

$$\int_0^A d_{DL}(A)dA = \mathcal{L}^{-1} \left[\frac{1}{4\pi} \int_0^{4\pi} dw \frac{2 \sum_{k=1}^{\infty} \cos(wk/2) e^{-s \sqrt{k_m^i(k_m^i+2)}}}{s \left(1 - 2 \sum_{k=1}^{\infty} \cos(wk/2) e^{-s \sqrt{k_m^i(k_m^i+2)}} \right)} ; A \right]$$

$$n(A) = \mathcal{L}^{-1} \left[\frac{1}{4\pi} \int_0^{4\pi} dw \frac{1}{s \left(1 - 2 \sum_{k=1}^{\infty} \cos(wk/2) e^{-s \sqrt{k_m^i(k_m^i+2)}} \right)} ; A \right]$$

$$\frac{1}{s \left(1 - 2 \sum_{k=1}^{\infty} e^{-s \sqrt{k_m^i(k_m^i+2)}} \right)}$$

$$\mathcal{L}^{-1} \left[\frac{1}{s} \left(1 - 2 \sum_{k=1}^{\infty} e^{-s \sqrt{k_m^i(k_m^i+2)}} \right)^{-1} ; A \right] \sim \sum_{s_i} \text{Res} \left[\frac{1}{s} \left(1 - 2 \sum_{k=1}^{\infty} e^{-s \sqrt{k_m^i(k_m^i+2)}} \right)^{-1} e^{sa} \right]$$

$$f(s) = \frac{2}{s(e^{s/2} - 3)}.$$

$$\mathcal{L}^{-1}[f(s); A] = 3^{2A},$$

$$\mathcal{L}^{-1}[f(s); A] \propto 3^{2A}$$

$$\lambda_p: z \mapsto \text{diag}(z^p, z^{-p}),$$

$$N^{\mathcal{P}} = \sum_{r_i} N_{j_1 j_2}^{r_1} N_{r_1 j_3}^{r_2} \dots N_{r_{N-2} j_{N-1}}^{r_{N-1}}$$

$$N^{\mathcal{P}} = \langle \chi_{j_1} \dots \chi_{j_N} | \chi_0 \rangle_{SU(2)} = \int_0^{2\pi} \frac{d\theta}{\pi} \sin^2 \theta \prod_{I=1}^N \frac{\sin [(j_I + 1)\theta]}{\sin \theta}$$

$$N_{U(1)}^{\mathcal{P}} = \langle \tilde{\eta}_{p_1} \dots \tilde{\eta}_{p_N} | \eta_0 \rangle_{U(1)} = \frac{1}{2\pi} \int_0^{2\pi} d\theta \prod_I 2 \cos p_I \theta$$

$$\gamma^\mu \nabla_\mu \psi = 0$$

$$\gamma^a V_a^i \partial_i \Phi + \frac{1}{r} \left[\frac{\gamma^2}{\sin^{1/2} \theta} \partial_\theta \sin^{1/2} \theta + \frac{\gamma^3}{\sin \theta} \partial_\phi \right] \Phi = 0$$

$$\partial_t \Phi = -\gamma^0 \gamma^1 \left[\partial_{r^*} + \frac{e^\rho}{r} \left(\frac{\gamma^2 \gamma^1}{\sin^{1/2} \theta} \partial_\theta \sin^{1/2} \theta + \frac{\gamma^3 \gamma^1}{\sin \theta} \partial_\phi \right) \right] \Phi$$

$$\partial_t \Phi = -\gamma^0 \gamma^1 \left[\partial_{r^*} - \frac{e^\rho}{r} \gamma^0 K \right] \Phi$$

$$\begin{aligned}
K &= \gamma^0 \left(\frac{\gamma^1 \gamma^2}{\sin^{1/2} \theta} \partial_\theta \sin^{1/2} \theta + \frac{\gamma^1 \gamma^3}{\sin \theta} \partial_\phi \right) \\
\chi_{m_j \kappa_j}^+ &= \begin{bmatrix} \eta(\hat{r})_{\kappa_j}^{m_j} \\ 0 \end{bmatrix} \\
\chi_{m_j \kappa_j}^- &= \begin{bmatrix} 0 \\ \eta(\hat{r})_{-\kappa_j}^{m_j} \end{bmatrix} \\
\psi_{w\kappa_j m_j}(x) &= \frac{e^{-\rho/2} e^{-i\omega t}}{r} \begin{bmatrix} G_{w\kappa_j}(r) \eta(\hat{r})_{\kappa_j}^{m_j} \\ -i F_{w\kappa_j}(r) \sigma^1 \eta(\hat{r})_{\kappa_j}^{m_j} \end{bmatrix} \\
\partial_r^* G_{w\kappa_j} &= -\frac{e^\rho}{r} \kappa_j G_{w\kappa_j} + w F_{w\kappa_j} \\
\partial_r^* F_{w\kappa_j} &= \frac{e^\rho}{r} \kappa_j F_{w\kappa_j} - w G_{w\kappa_j} \\
\eta(\hat{r})_{\kappa_j < 0}^{m_j} &= \begin{bmatrix} \sqrt{\frac{j+m_j}{2j}} \\ \sqrt{\frac{j-m_j}{2j}} Y_{j-1/2}^{m_j-1/2}(\theta, \phi) \\ Y_{j-1/2}^{m_j+1/2}(\theta, \phi) \end{bmatrix} \\
\eta(\hat{r})_{\kappa_j > 0}^{m_j} &= \begin{bmatrix} \sqrt{\frac{j+1-m_j}{2j+2}} Y_{j+1/2}^{m_j-1/2}(\theta, \phi) \\ -\sqrt{\frac{j+1+m_j}{2j+2}} Y_{j+1/2}^{m_j+1/2}(\theta, \phi) \end{bmatrix} \\
\psi(x) &= \sum_{\kappa_j m_j} \int dw \left[a_{w\kappa_j m_j} u_{w\kappa_j m_j}(x) + b_{w\kappa_j m_j}^\dagger v_{w\kappa_j m_j}(x) \right] \\
\int_{I^+} d\Omega du r^2 \bar{u}^{out} \gamma_+ u^{out} &= \int_{I^-} d\Omega r^2 dv \bar{u}^{out} \gamma_- u^{out} \\
u^{out}(v)|_{I^+} &= \sqrt{du(v)/dv} \Theta(v_H - v) u^{out}(u)|_{I^-} \\
\langle in | N_{i_1 i_2} | in \rangle &= \sum_k \int_{I^-} dv_2 r_2^2 d\Omega_2 \left(\bar{u}_{i_2}^{out,L}(x_2) \frac{[\gamma^0 - \gamma^1]}{2} v_k^{in,L}(x_2) \right) \times \\
&\quad \times \int_{I^-} dv_1 r_1^2 d\Omega_1 \left(\bar{v}_k^{in,L}(x_1) \frac{[\gamma^0 - \gamma^1]}{2} u_{i_1}^{out,L}(x_1) \right) \\
\int d\Omega_2 \bar{u}_{i_2}^{out,L}(x_2) \frac{[\gamma^0 - \gamma^1]}{2} v_k^{in,L}(x_2) &= \frac{t_{j_2}^*(w_2)}{2\pi r_2^2} \sqrt{\frac{du(v)}{dv}} \Theta(v_H - v) \times \\
&\quad \times e^{i w_2 u(v_2) + i w v_2} \delta_{m_{j_2} m_k} \delta_{j_2 j_k} \\
\langle in | N_{i_1 i_2} | in \rangle &= \delta_{m_{j_1} m_{j_2}} \delta_{j_1 j_2} \frac{t_{j_1}(w_1) t_{j_2}^*(w_2)}{4\pi^2} \int_{-\infty}^{v_H} dv_1 dv_2 \\
&\quad \sqrt{\frac{du(v_1)}{dv} \frac{du(v_2)}{dv}} e^{-i w_1 u(v_1) + i w_2 u(v_2)} \int_0^\infty dw e^{-i w (v_1 - v_2)} \\
\int_0^\infty dw e^{-i w (v_1 - v_2)} &= \lim_{\epsilon \rightarrow 0} \frac{-i}{(v_1 - v_2 - i\epsilon)} \\
ds^2 &= -dT^2 + dW^2 + dX^2 + dY^2 + dZ^2
\end{aligned}$$



$$\begin{aligned}
H^{-1}\sinh H\tau &= T \\
H^{-1}\cosh H\tau\cos \xi &= W \\
H^{-1}\cosh Ht\sin \xi\cos \theta &= X \\
H^{-1}\cosh Ht\sin \xi\cos \theta\cos \phi &= Y \\
H^{-1}\cosh Ht\sin \xi\cos \theta\sin \phi &= Z \\
ds^2 &= -d\tau^2 + H^{-2}\cosh^2 H\tau[d\xi^2 + \sin^2 \xi(d\theta^2 + \sin^2 \theta)] \\
ds^2 &= \frac{1}{\cos^2 \tilde{T}}(-d\tilde{T}^2 + d\Omega_3^2) \\
T &= H^{-1}\sinh Ht + \frac{1}{2}He^{Ht}|\vec{x}|^2, X = e^{Ht}x, Y = e^{Ht}y \\
Z &= e^{Ht}z, W = H^{-1}\cosh Ht - \frac{1}{2}He^{Ht}|\vec{x}|^2 \\
ds^2 &= -dt^2 + e^{2Ht}d\vec{x}^2, -\infty < t, x, y, z < \infty \\
T &= (H^{-2} - \tilde{r})^{1/2}\sinh H\tilde{t} \\
X &= (H^{-2} - \tilde{r})^{1/2}\cosh (H\tilde{t}) \\
Y &= \tilde{r}\sin \tilde{\theta}\cos \tilde{\phi} \\
Z &= \tilde{r}\sin \tilde{\theta}\sin \tilde{\phi} \\
W &= \tilde{r}\cos \tilde{\theta} \\
ds^2 &= -(1 - \tilde{r}^2H^2)d\tilde{t}^2 + \frac{1}{(1 - \tilde{r}^2H^2)}d\tilde{r}^2 + \tilde{r}^2(d\tilde{\theta}^2 + \sin^2 \tilde{\theta}d\tilde{\phi}) \\
ds^2 &= -dt^2 + a(t)^2(dx^2 + dy^2 + dz^2) \\
\Box \phi - m^2\phi - \xi R\phi &= 0 \\
\phi(t, x, y, z) &= \frac{1}{(2\pi)^{3/2}} \int d^3k u_k(t) e^{i\vec{k}\vec{x}} \\
\partial_t^2 u_k(t) + 3H\partial_t u_k(t) + \left(\frac{k^2}{a^2} + m^2 + \xi R\right) u_k(t) &= 0 \\
u_k(t) &= a^{-3/2} v_k(t) \\
\partial_t^2 v_k(t) + \left(-\frac{9}{4}H^2 + \frac{k^2}{a^2(t)} + m^2 + 12\xi H^2\right) v_k(t) &= 0 \\
ds^2 &= \frac{1}{H\eta^2}(-d\eta^2 + d\vec{x}^2) \\
\eta^2 \partial_\eta^2 v_k(\eta) + \eta \partial_\eta v_k(\eta) + (k^2\eta^2 - v^2)v_k(\eta) &= 0 \\
v_k(k\eta) &= AJ_v(k\eta) + BY_v(k\eta) \\
v_k(\eta k) &\rightarrow \frac{1}{\sqrt{2\pi k}} e^{ik\eta} \\
v_k(k\eta) &= \frac{1}{2}\sqrt{\frac{\pi}{H}}(J_\nu(k\eta) + iY_\nu(k\eta)) = \frac{1}{2}\sqrt{\frac{\pi}{2H}}H_\nu^{(1)}(k\eta) \\
u_k^{ds} &= \frac{1}{\sqrt{(2\pi)^3 a(t)^3}} v_k(t) \\
G_{ds}(x, x') &= \hbar \frac{H^2(1/4 - \nu)}{16\pi^2 \cos \pi\nu} {}_2F_1\left[3/2 + \nu, 3/2 - \nu; 2; (1 + \cos \sqrt{H^2\sigma})/2\right] \\
G_{ds}(x, x') &= \frac{\hbar^2 H^2 \eta \eta'}{4\pi^2 [-(\eta - \eta' - i\epsilon)^2 + (\vec{x} - \vec{x}')^2]} \\
ds^2 &= -(1 - \tilde{r}^2 H^2) d\tilde{t}^2 + \frac{1}{(1 - \tilde{r}^2 H^2)} d\tilde{r}^2 + \tilde{r}^2 d\Omega^2 \\
\phi(\tilde{t}, \tilde{r}, \tilde{\theta}, \tilde{\phi}) &= \frac{1}{\sqrt{4\pi}} \frac{1}{\tilde{r}} \sum_{l,m} \int dw u_{w,l}^C(\tilde{r}, \tilde{t}) Y_{l,m}(\tilde{\theta}, \tilde{\phi}) \\
(1 - \tilde{r}^2 H^2) \partial_{\tilde{r}}^2 u_{w,l}^C(\tilde{r}) - 2H^2 \tilde{r} \partial_{\tilde{r}} u_{w,l}^C(\tilde{r}) &+ \\
+ \left(2H - \frac{l(l+1)}{\tilde{r}^2} - m^2 - 12\xi H^2 + \frac{w^2}{1 - \tilde{r}^2 H^2}\right) u_{w,l}^C(\tilde{r}) &= 0
\end{aligned}$$



$$\begin{aligned}
& \partial_{r^*}^2 u_{w,l}^C(r^*) + V(r^*)u_{w,l}^C(r^*) = -w^2 u_{w,l}^C(r^*) \\
u_{w,l=0}^C(r^*) &= N_\nu(w) \left[P_{\nu-1/2}^{i\frac{w}{H}}(\tanh Hr^*) - \alpha_\nu(w) Q_{\nu-1/2}^{i\frac{w}{H}}(\tanh Hr^*) \right] \\
\alpha_\nu(w) &= \frac{1}{2} \left[2i + 4(i + e^{-i\pi\nu} e^{\pi w/H})^{-1} \right] \\
P_\nu^{i\frac{w}{H}}(\tanh Hz) &\sim \frac{1}{\Gamma\left[1 - \frac{iw}{H}\right]} e^{i\omega z} \\
Q_\nu^{i\frac{w}{H}}(\tanh Hz) &\sim A(w)e^{i\omega x} + B(w)e^{-i\omega z} \\
A(w) &= \frac{-i\pi}{4\Gamma(1 - iw/H)} \left(\coth \frac{\pi w}{2H} + \tanh \frac{\pi w}{2H} \right) \\
B(w) &= \frac{-\pi}{4\Gamma(1 - iw/H)} \frac{\coth \frac{\pi w}{2H}}{\sinh^2 \frac{\pi w}{2H}} \\
&\quad (u_{w_1,l=0}^C(v), u_{w_2,l=0}^C(v)) \\
&= -i \int_{-\infty}^{\infty} dv d\tilde{\Omega} \tilde{r}^2 (u_{w_1,l=0}^C(v) \partial_v u_{w_2,l=0}^{C*}(v) - u_{w_2,l=0}^{C*}(v) \partial_v u_{w_1,l=0}^C(v)) \\
&= N_\nu^2(w_1 + w_2) \left| \Gamma^{-1} \left(1 - \frac{iw}{H} \right) - \alpha_\nu(w) A(w) \right|^2 \int_{-\infty}^{\infty} dv e^{-i(w_1 - w_2)v} \\
&\quad = \delta(w_1 - w_2) \\
|\tilde{N}_n(w/H)|^2 &= \frac{1}{4\pi} \frac{|\Gamma(1 - iw/H)|^2}{\left| 1 + \frac{i\pi}{4} \alpha_n(w) \left(\coth \frac{\pi w}{2H} + \tanh \frac{\pi w}{2H} \right) \right|^2} \\
u_{w,l=0}^C(\tilde{t}, \tilde{r}) &= \frac{1}{\sqrt{4\pi w}} \frac{1}{\tilde{r}} \tilde{N}_\nu(w/H) e^{-i\omega \tilde{t}} \left[P_{\nu-1/2}^{i\frac{w}{H}}(H\tilde{r}) - \alpha_\nu(w) Q_{\nu-1/2}^{i\frac{w}{H}}(H\tilde{r}) \right] \\
&\quad u_{w,l=0}^C(r^*) = N \sin wr^* \\
&\quad (u_{w_1,l=0}^C(v), u_{w_2,l=0}^C(v)) \\
&= -i \int_{-\infty}^{\infty} dv d\tilde{\Omega} \tilde{r}^2 (u_{w_1,l=0}^C(v) \partial_v u_{w_2,l=0}^{C*}(v) - u_{w_2,l=0}^{C*}(v) \partial_v u_{w_1,l=0}^C(v)) \\
&= N \frac{(w_1 + w_2)}{4} \int_{-\infty}^{\infty} dv e^{-i(w_1 - w_2)v} = \delta(w_1 - w_2) \\
&\quad \langle 0_C(x) | \phi(x) \phi(x + \Delta(x)) | 0_C(x) \rangle \\
&= \lim_{\epsilon \rightarrow 0} \frac{\hbar}{4\pi^2} \int_0^\infty dw e^{-i\omega(\Delta t - i\epsilon)} \frac{1}{\tilde{r}} \sin \left[\frac{w}{H} \tanh^{-1}(\tilde{r}H) \right] \\
&= \lim_{\epsilon \rightarrow 0} \frac{\hbar}{4\pi^2} \frac{1}{\tilde{r}} \frac{\chi(\tilde{r})}{\chi(\tilde{r})^2 - (\Delta\tilde{t} - i\epsilon)^2} \\
\langle 0_C(x) | \phi(x) \phi(x + \Delta(x)) | 0_C(x) \rangle &= \frac{\hbar}{4\pi^2 [(\Delta\tilde{t} - i\epsilon)^2 - \tilde{r}^2]} \\
\langle in | N_{j_1 n_1, j_2 n_2}^{out, \sigma} | in \rangle &= \int_0^\Lambda dw' \beta_{j_1 n_1, w'} \beta_{j_2 n_2, w'}^* = \\
\frac{1}{\epsilon} \int_{j_1 \epsilon}^{(j_1+1)\epsilon} dw_1 \int_{j_2 \epsilon}^{(j_2+1)\epsilon} dw_2 e^{2\pi i w_1 n_1 / \epsilon} e^{-2\pi i w_2 n_2 / \epsilon} \int_0^\Lambda dw' \beta_{w_1 w'} \beta_{w_2 w'}^* \\
\langle in | N_{j_1 n_1, j_2 n_2}^{out, \sigma} | in \rangle &= \frac{1}{\epsilon} \int_{j_1 \epsilon}^{(j_1+1)\epsilon} dw_1 \int_{j_2 \epsilon}^{(j_2+1)\epsilon} dw_2 e^{i \frac{2\pi w_1 n_1}{\epsilon}} e^{-i \frac{2\pi w_2 n_2}{\epsilon}} \\
&\quad t_l(w_1) t_l^*(w_2) \frac{e^{-i(w_1 - w_2)v_H}}{2\pi \sqrt{w_1 w_2}} e^{-\pi \kappa^{-1} \omega_1 i - i\kappa^{-1}(w_1 - w_2)} \\
&\quad \Gamma(1 + i\kappa^{-1} w_1) \Gamma(1 - i\kappa^{-1} w_2) \delta_\sigma[\kappa^{-1}(w_1 - w_2)].
\end{aligned}$$

$$\begin{aligned}
\langle in | N_{j_1 n_1, j_2 n_2}^{out, \sigma} | in \rangle &\approx \delta_{j_1 j_2} \frac{|t_l(w_j)|^2 |\Gamma(1 + i\kappa^{-1} w_j)|^2}{2\pi w_j} e^{-\pi\kappa^{-1} w_j} e^{\frac{2\pi(n_1 - n_2) w_j}{\epsilon}} I_{n_1 n_2}(\sigma) \\
I_{n_1 n_2}(\sigma) &= \frac{1}{\epsilon} \int_{-\epsilon/2}^{\epsilon/2} dx_1 \int_{-\epsilon/2}^{\epsilon/2} dx_2 e^{i[\frac{2\pi n_1}{\epsilon} - v_H] x_1 - i[\frac{2\pi n_2}{\epsilon} - v_H] x_2 - \pi\kappa^{-1}(x_1 + x_2)/2} \delta_\sigma[\kappa^{-1}(x_1 - x_2)] \\
I_{n_1 n_2}(\sigma) &\approx \frac{1}{\epsilon} \int_{-\epsilon/2}^{\epsilon/2} dx_1 e^{\frac{2\pi(n_1 - n_2)x_1}{\epsilon}} \int_{x_1 - \epsilon/2}^{x_1 + \epsilon/2} dy e^{i[\frac{2\pi n_2}{\epsilon} - v_H] y} \delta_\sigma[\kappa^{-1} y] \\
I_{n_1 n_2}(\sigma) &\approx \kappa \delta_{n_1 n_2} \frac{\sin \left[\left(\frac{2\pi n_2}{\epsilon} - v_H \right) \frac{\pi \kappa \sigma}{2} \right]}{\left[\left(\frac{2\pi n_2}{\epsilon} - v_H \right) \frac{\pi \kappa \sigma}{2} \right]} \\
\int_{-\infty}^{\infty} d \log [w/\kappa] e^{-i\kappa^{-1}(w_1 - w_2) \log [w/\kappa]} &\rightarrow \int_{-\log [\Lambda/\kappa]}^{\log [\Lambda/\kappa]} d \log [w/\kappa] e^{-i\kappa^{-1}(w_1 - w_2) \log [w/\kappa]} \\
\int_{-\log [\Lambda/\kappa]}^{\log [\Lambda/\kappa]} d \lambda e^{-i\kappa^{-1}(w_1 - w_2) \lambda} &\rightarrow \int_{-\infty}^{\infty} d \lambda e^{-i\kappa^{-1}(w_1 - w_2) \lambda} e^{-(\lambda/\Lambda)^2} \\
\delta_{\tilde{\sigma}}[\kappa^{-1}(w_1 - w_2)] &= \frac{\exp \left[\frac{\kappa^{-1}(w_1 - w_2)}{2\tilde{\sigma}} \right]^2}{2\tilde{\sigma}\sqrt{\pi}} \\
I_{n_1 n_2}(\tilde{\sigma}) &= \kappa \delta_{n_1 n_2} e^{-\left[\left(\frac{2\pi n_2}{\epsilon} - v_H \right) \kappa \tilde{\sigma} \right]^2} \\
\langle in | N_{j_1 n_1, j_2 n_2}^{out, \sigma} | in \rangle &= \delta_{j_1 j_2} \delta_{n_1 n_2} \frac{|t_l(w_j)|^2}{e^{2\pi\kappa^{-1} w_j} - 1} e^{-\left[\left(\frac{2\pi n_2}{\epsilon} - v_H \right) \kappa \tilde{\sigma} \right]^2} \\
I &= \int_{-\infty}^{\infty} dz e^{-iwz} \left[\frac{(k/2)^2}{\sinh^2 kz/2 + (k/2)^2 \alpha^2} \right] \\
I_1 &= \int_0^{\infty} dx \frac{x^{-iw/k}}{(x-1)^2 + k^2 \alpha^2 x} \\
I_1 &= \frac{2\pi i}{1 - e^{2\pi w/k}} [\text{Res}[a + ib] + \text{Res}[a - ib]] \\
\text{Res}[a \pm ib] &= \pm \frac{(a \pm ib)^{-iw/k}}{2ib} \\
(a + ib) &= e^{w/k\theta} \\
(a - ib) &= (e^{w/k\theta} - e^{w/k(2\pi - \theta)}) \\
I_1 &= \frac{\pi}{1 - e^{2\pi w/k}} \frac{(e^{w/k\theta} - e^{w/k(2\pi - \theta)})}{\alpha \sqrt{1 - \left(\frac{k\alpha}{2}\right)^2}} \\
I_2 &= - \int_{-\infty}^{\infty} dz e^{-iwz} \frac{1}{z^2 + \alpha^2} \\
I_2 &= (-2\pi i \text{Res}[i\alpha])^* \\
I_2 &= -\frac{\pi}{\alpha} e^{-w\alpha} \\
I &= \frac{1}{(1 - e^{2\pi w/k}) \alpha} \left[\frac{(e^{w/k\theta} - e^{w/k(2\pi - \theta)})}{\sqrt{1 - \left(\frac{k\alpha}{2}\right)^2}} + (e^{2\pi w/k} - 1) e^{-w\alpha} \right] \\
I &= w \left[\frac{w e^{\pi w/k} \sinh \left[\frac{w}{a} (\beta - \pi) \right]}{(1 - e^{2\pi w/k}) \frac{w}{a} \sin \beta} - \frac{\pi e^{-w\alpha}}{\alpha} \right] \\
&\quad \frac{1}{w_k(t)} + (W_k(t)^{-1})^{(2)}
\end{aligned}$$

$$\begin{aligned}
(W_k(t)^{-1})^{(2)} &= \frac{m^2 a'^2}{2a^2 w^5} + \frac{m^2 a''}{4aw^5} - \frac{5m^4 a'^2}{8a^2 w^7} + \frac{a'^2}{2a^2 w^3} + \frac{a''}{2aw^3} \\
a' &= \frac{\dot{a}}{a} \\
a'' &= -\frac{\dot{a}^2}{a^3} + \frac{\ddot{a}}{a^2} \\
(W_k(t)^{-1})^{(2)} &= \frac{m^2 \dot{a}^2 a}{2\tilde{w}^5} - \frac{m^2 \dot{a} a}{4\tilde{w}^5} + \frac{m^2 \ddot{a} a^2}{4\tilde{w}^5} - \frac{5m^4 \dot{a}^2 a^3}{8\tilde{w}^7} + \frac{1}{2} \frac{\ddot{a}}{\tilde{w}^3} \\
a &= -\frac{H\tau}{1+\epsilon} \frac{\dot{a}}{\tau}, \frac{\dot{a}}{a} = -\frac{1+\epsilon}{\tau} \frac{\ddot{a}}{a}, \frac{\ddot{a}}{a} = \frac{2+3\epsilon}{\tau^2}, \\
(W_k(t)^{-1})^{(2)} &= \frac{a^3 m^2}{\tilde{w}^5} \left(\frac{3+5\epsilon}{4\tau^2} \right) - \frac{5m^4}{8\tilde{w}^7} \frac{1+2\epsilon}{\tau^2} a^5 + \frac{a}{2\tilde{w}^3} \frac{2+3\epsilon}{\tau^2} \\
\tilde{w} &= a \sqrt{\frac{k}{a^2} + m^2} = \frac{1}{\tau} \sqrt{(1+\epsilon)^2 + \frac{m^2}{H^2} (1+\epsilon)} \approx \frac{1}{\tau} \sqrt{1+2\epsilon + \frac{m^2}{H^2}} \\
w_k^{-1} + (W_k(t)^{-1})^{(2)} &\approx \frac{a^3 m^2}{\tilde{w}^5} \frac{3}{4\tau^2} + \frac{a}{2\tilde{w}^3} \frac{2+3\epsilon}{\tau^2} + \frac{a}{\tilde{w}} \approx \\
&\approx -\frac{3m^2}{4H^3} \left(1 - 5\epsilon - \frac{5m^2}{2H^2} \right) - \frac{2+5\epsilon}{2H} \left(1 - 3\epsilon - \frac{3m^2}{2H^2} \right) \\
&\approx \frac{1}{H} \left(-2 + \frac{\epsilon}{2} + \frac{5m^2}{4H^2} \right)
\end{aligned}$$

Para un espacio – tiempo AdS, tenemos:

$$\begin{aligned}
i[M_{\mu\nu}, M_{\rho\sigma}] &= \eta_{\nu\rho} M_{\mu\sigma} - \eta_{\mu\rho} M_{\nu\sigma} - \eta_{\sigma\mu} M_{\rho\nu} + \eta_{\sigma\nu} M_{\rho\mu} \\
i[P_\mu, M_{\sigma\rho}] &= \eta_{\mu\rho} P_\sigma - \eta_{\mu\sigma} P_\rho \\
[P_\mu, P_\nu] &= 0 \\
i[D, P_\mu] &= P_\mu \\
[M_{\mu\nu}, D] &= 0 \\
K_\mu: x^\mu &\rightarrow \frac{x^\mu + a^\mu x^2}{1 + 2x^\nu a_\nu + a^2 x^2} \\
i[M_{\mu\nu}, K_\rho] &= \eta_{\mu\rho} K_\nu - \eta_{\nu\rho} K_\mu \\
[D, K_\mu] &= iK_\mu \\
[P_\mu, K_\nu] &= 2i(M_{\mu\nu} - \eta_{\mu\nu} D) \\
[K_\mu, K_\nu] &= 0 \\
J_{\mu\nu} = M_{\mu\nu}, J_{\mu d} &= \frac{1}{2}(K_\mu - P_\mu), J_{\mu(d+1)} = \frac{1}{2}(K_\mu + P_\mu), J_{(d+1)d} = D \\
J_{ab} &= \begin{pmatrix} J_{\mu\nu} & J_{\mu d} & J_{\mu(d+1)} \\ -J_{\mu d} & 0 & D \\ -J_{\mu(d+1)} & -D & 0 \end{pmatrix} \\
P_\mu: x_\mu &\rightarrow x_\mu + a_\mu \Rightarrow d \\
M_{\mu\nu}: x_\mu &\rightarrow \Lambda_\mu^\nu x_\nu \Rightarrow \frac{d(d-1)}{2} \\
D: x^\mu &\rightarrow \lambda x^\mu \Rightarrow 1 \\
K_\mu: x_\mu &\rightarrow \frac{x_\mu + a_\mu x^2}{1 + 2x_\nu a^\nu + a^2 x^2} \Rightarrow d \\
x &\rightarrow \lambda x \Rightarrow \phi(x) \rightarrow \phi(x)' = \lambda^\Delta \phi(\lambda x) \\
[D, P_\mu] &= -iP_\mu \Rightarrow D(P_\mu \phi) = -i(\Delta + 1)(P_\mu \phi)
\end{aligned}$$



$$\begin{aligned}
\langle \phi(0)\phi(x) \rangle &\equiv \frac{1}{(x^2)^\Delta} \\
ds^2 &= dX^2 + dY^2 + dZ^2 \\
X^2 + Y^2 + Z^2 &= L^2 \\
ds^2 &= L^2(d\theta^2 + \sin^2 \theta d\varphi^2) \\
ds^2 &= dZ^2 + dX^2 + dY^2 \\
-Z^2 + X^2 + Y^2 &= -L^2 \\
\\
ds^2 &= -dZ^2 + dX^2 + dY^2 \\
-Z^2 + X^2 + Y^2 &= -L^2 \\
X = L \sinh \rho \cos \varphi, Y &= L \sinh \rho \sin \varphi, Z = L \cosh \rho \\
ds^2 &= L^2(d\rho^2 + \sinh^2 \rho d\varphi^2) \\
R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R &= -\Lambda g_{\mu\nu} \\
R_{\mu\nu\theta\sigma} &= -\frac{1}{l^2}(g_{\mu\theta}g_{\nu\sigma} - g_{\mu\sigma}g_{\nu\theta}), R_{\mu\nu} = -\frac{3}{l^2}g_{\mu\nu}, R = -\frac{12}{l^2} \\
ds^2 &= -dX_0^2 - dX_4^2 + dX_1^2 + dX_2^2 + dX_3^2 \\
-X_0^2 - X_4^2 + X_1^2 + X_2^2 + X_3^2 &= -l^2 \\
X_0 &= l \cosh \rho \cos \tau \\
X_4 &= l \cosh \rho \sin \tau \\
X_1 &= l \sinh \rho \sin \theta \sin \varphi \\
X_2 &= l \sinh \rho \sin \theta \cos \varphi \\
X_3 &= l \sinh \rho \cos \theta \\
\frac{ds^2}{l^2} &= -\cosh^2 \rho d\tau^2 + d\rho^2 + \sinh^2 \rho (d\theta^2 + \sin^2 \theta d\varphi^2) \\
\frac{ds^2}{l^2} &\approx -d\tau^2 + d\rho^2 + \rho^2 (d\theta^2 + \sin^2 \theta d\varphi^2) \\
r &:= \sinh \rho, t := l\tau \\
ds^2 &= \left(1 + \frac{r^2}{l^2}\right) dt^2 + \left(1 + \frac{r^2}{l^2}\right)^{-1} dr^2 + r^2 (d\theta^2 + \sin^2 \theta d\varphi^2) \\
\frac{ds^2}{l^2} &= \frac{1}{\cos^2 \chi} [-d\tau^2 + d\chi^2 + \sin^2 \chi (d\theta^2 + \sin^2 \theta d\varphi^2)] \\
X_0 &= \frac{lr}{2} \left(\vec{x}_i^2 - t^2 + \frac{1}{r^2} + 1 \right), \quad X_i = lrx_i \quad (i = 1, 2), \\
X_3 &= \frac{lr}{2} \left(\vec{x}_i^2 - t^2 + \frac{1}{r^2} - 1 \right), \quad X_4 = lrt, \\
\frac{ds^2}{l^2} &= r^2 (-dt^2 + d\vec{x}^2) + \frac{dr^2}{r^2} \\
ds^2 &= \frac{l^2}{z^2} (-dt^2 + d\vec{x}^2 + dz^2) \\
Z_{\text{gauge}} &= Z_{\text{string}} \\
\mathcal{L} &= -\frac{N}{2g_{YM}^2} F_{\mu\nu}^M F_M^{\mu\nu} \\
Z_{\text{CFT}} &= e^{-W} \\
\frac{l^4}{l_s^4} &\sim g_{YM}^2 N \sim g_s N \gg 1 \\
Z_{\text{cuerdas}} &\approx e^{-I_{SUGRA}^E} \\
Z_{\text{cuerdas}} &\approx e^{-I_{SUGRA}^E} = e^{-W} = Z_{\text{CFT}} \\
\left\langle e^{\int d^3x \phi_0(\vec{x}) \mathcal{O}(\vec{x})} \right\rangle_{\text{CFT}} &= e^{-I_{\text{bulk}}^E[\phi |_{\partial \text{AdS}} \rightarrow \phi_0]} \\
\left\langle e^{\int d^3x h_{ab}^0 T^{ab}} \right\rangle_{\text{CFT}} &= e^{-I_{\text{bulk}}^E[h_{\mu\nu} |_{\partial \text{AdS}} \rightarrow h_{ab}^0]}
\end{aligned}$$



$$\begin{aligned}
& (\nabla^\mu \nabla_\mu - m^2) \Phi(z, x^n) = 0 \\
& \Phi(x^n, z) = \int \frac{d^2 \vec{k}}{(2\pi)^2} d\omega f_k(z) e^{ik_\mu x^\mu} \\
& \frac{d^2 f_k}{dz^2} - \frac{2}{z} \frac{df_k}{dz} - \left(k^2 + \frac{m^2 l^2}{z^2} \right) f_k = 0 \\
& f_k(z) = a_1 z^{3/2} K_\nu(kz) + a_2 z^{3/2} I_\nu(kz), \\
& \frac{d^2 f_k}{dz^2} - k^2 f_k = 0 \\
& f_k(z) = a_1 z^{3/2} K_\nu(kz), K_\nu(kz) = \frac{\pi I_{-\nu}(kz) - I_\nu(kz)}{2 \sin(\nu\pi)}. \\
& f_k(z) = a_1 z^{3/2} \frac{\pi}{2 \sin \pi\nu} \left[\frac{1}{\Gamma(1-\nu)} \left(\frac{kz}{2}\right)^{-\nu} - \frac{1}{\Gamma(1+\nu)} \left(\frac{kz}{2}\right)^\nu \right] \\
& f_k(z) = \phi_0 z^{\Delta_-} + \phi_1 z^{\Delta_+} \\
& \Delta_\pm (\Delta_\pm - 3) - ml^2 = z^2 k^2 \xrightarrow{z=0} \Delta_\pm (\Delta_\pm - 3) - ml^2 = 0, \\
& \Delta_\pm = \frac{3}{2} \pm \nu, \nu = \sqrt{\frac{9}{4} + m^2 l^2} \\
& m_{BF}^2 \leq m^2 < m_{BF}^2 + \frac{1}{l^2} \Rightarrow 0 \leq \nu < 1 \\
& \phi(r, x^n) = \frac{\alpha(x^n)}{r^{3-\Delta}} + \frac{\beta(x^n)}{r^\Delta} + \dots, \\
& \Delta = \frac{3}{2} + \nu, \nu = \sqrt{\frac{9}{4} + m^2 l^2} \\
& \phi = \frac{\alpha}{r} + \frac{\beta}{r^2} + \dots \\
& \partial_r \phi = -\frac{\alpha}{r^2} - \frac{2\beta}{r^3} + \dots = \frac{\alpha'}{r^2} + \dots \\
& a\phi + b\partial_r \phi = a \left(\frac{\alpha}{r} + \frac{\beta}{r^2} \right) + b \left(-\frac{\alpha}{r^2} - \frac{2\beta}{r^3} \right) = \frac{\alpha'}{r} + \frac{\beta'}{r^2} + \dots \\
& I_{CFT} \rightarrow I_{CFT} + p \int d^3 x \mathcal{O}(x) \\
& I = \frac{1}{2\kappa} \int d^4 x \sqrt{-g} R + I_B \\
& \delta I = \frac{1}{2\kappa} \int d^4 x \sqrt{-g} G_{\alpha\beta} \delta g^{\alpha\beta} + \int d^4 x \sqrt{-g} g^{\alpha\beta} \delta R_{\alpha\beta} + \delta I_B \\
& G_{\alpha\beta} = R_{\alpha\beta} - \frac{1}{2} g_{\alpha\beta} R \\
& \delta I_B = - \int_{\mathcal{M}} d^4 x \sqrt{-g} g^{\alpha\beta} \delta R_{\alpha\beta} = - \oint_{\partial\mathcal{M}} \epsilon v^\mu n_\nu \sqrt{-h} d^3 x \\
& g^{\alpha\beta} \delta R_{\alpha\beta} = v^\mu{}_{;\mu}, v^\mu = g^{\alpha\beta} \delta \Gamma_{\alpha\beta}^\mu - g^{\alpha\mu} \delta \Gamma_{\alpha\beta}^\beta, \epsilon = n^\mu n_\mu = \pm 1 \\
& I_B = \int_{\partial\mathcal{M}} d^3 x \sqrt{-h} K \\
& I = \frac{1}{2\kappa} \int_{\mathcal{M}} d^4 x \sqrt{-g} R + \frac{1}{\kappa} \int_{\partial\mathcal{M}} d^3 x \sqrt{-h} K \\
& ds^2 = -N(r) dt^2 + H(r) dr^2 + r^2 (d\theta^2 + \sin^2 \theta d\varphi^2) \\
& \text{Krets} = R^{\alpha\beta\gamma\sigma} R_{\alpha\beta\gamma\sigma} \\
& I[g_{\mu\nu}] = \frac{1}{2\kappa} \int_{\mathcal{M}} d^4 x \sqrt{-g} R + \frac{1}{\kappa} \int_{\partial\mathcal{M}} d^3 x K \sqrt{-h} \\
& R_{\mu\nu} = 0
\end{aligned}$$

$$\begin{aligned}
ds^2 &= -\left(1 - \frac{\mu}{r}\right) dt^2 + \left(1 - \frac{\mu}{r}\right)^{-1} dr^2 + r^2(d\theta^2 + \sin^2 \theta d\phi^2) \\
M &= \frac{4\pi\mu}{\kappa} = \frac{\mu}{2G} \\
I[g_{\mu\nu}, A_\mu] &= \frac{1}{2\kappa} \int_{\mathcal{M}} d^4x \sqrt{-g} \left(R - \frac{1}{4} F_{\mu\nu} F^{\mu\nu} \right) + \frac{1}{\kappa} \int_{\partial\mathcal{M}} d^3x K \sqrt{-h} \\
R_{\mu\nu} - \frac{1}{2} R g_{\mu\nu} &= \frac{1}{2} T_{\mu\nu}^{EM} \\
\nabla_\mu F^{\mu\nu} &= 0 \\
T_{\mu\nu}^{EM} &= F_{\mu\alpha} F_\nu^\alpha - \frac{1}{4} g_{\mu\nu} F^2 \\
A &\equiv A_\mu dx^\mu = \left(\frac{q}{r} - \frac{q}{r_+} \right) dt, F = -\frac{q}{r^2} dr \wedge dt \\
ds^2 &= -f(r) dt^2 + f(r)^{-1} dr^2 + r^2(d\theta^2 + \sin^2 \theta d\phi^2) \\
f(r) &= 1 - \frac{\mu}{r} + \frac{q^2}{4r^2} = \frac{(r - r_-)(r - r_+)}{r^2} \\
M &= \frac{4\pi\mu}{\kappa} = \frac{\mu}{2G} \\
Q &\equiv \frac{1}{\kappa} \oint d^2 \star F = -\frac{q}{4G} \\
r_\pm &= G \left(M \pm \sqrt{M^2 - 4Q^2} \right) \\
\Phi &= A_t|_{r=\infty} - A_t|_{r=r_+} = \frac{4GQ}{r_+} \\
ds^2 &= \frac{r^2}{l^2} (-dt^2 + l^2 d\Sigma_k^2) \\
ds^2 &= -N(r) dt^2 + H(r) dr^2 + S(r) d\Sigma_k^2 \\
&\begin{cases} d\theta^2 + \sin^2 \theta d\phi^2 & \text{for } k = +1 \\ \frac{1}{l^2} \sum_{i=1}^2 dx_i^2 & \text{for } k = 0 \\ d\theta^2 + \sinh^2 \theta d\phi^2 & \text{for } k = -1 \end{cases} \\
&\mathbb{R} \times H^2, \mathbb{R} \times \mathbb{R}^2, \mathbb{R} \times S^2 \\
d\Sigma_k^2 &= \frac{dy^2}{1 - ky^2} + (1 - ky^2) dz^2 \\
I[g_{\mu\nu}] &= \frac{1}{2\kappa} \int_{\mathcal{M}} d^4x \sqrt{-g} (R - 2\Lambda) + \frac{1}{\kappa} \int_{\partial\mathcal{M}} d^3x K \sqrt{-h} \\
R_{\mu\nu} - \frac{1}{2} g_{\mu\nu} R &= -\Lambda g_{\mu\nu} \\
ds^2 &= -\left(k - \frac{\mu}{r} + \frac{r^2}{l^2}\right) dt^2 + \left(k - \frac{\mu}{r} + \frac{r^2}{l^2}\right)^{-1} dr^2 + r^2 d\Sigma_k^2 \\
k - \frac{\mu}{r_h} + \frac{r_h^2}{l^2} &= 0 \\
M &= \frac{\sigma_k \mu}{\kappa} \\
I[g_{\mu\nu}, \phi] &= \int d^4x \sqrt{-g} \left[\frac{R}{2\kappa} - \frac{1}{2} \partial_\mu \phi \partial^\mu \phi - V(\phi) \right] \\
\frac{1}{\sqrt{-g}} \partial_\mu (\sqrt{-g} g^{\mu\nu} \partial_\nu \phi) - \frac{\partial V}{\partial \phi} &= 0 \\
\left. \frac{dV}{d\phi} \right|_{\phi=0} &= 0, V(0) = -\frac{3}{\kappa l^2}, \left. \frac{d^2 V}{d\phi^2} \right|_{\phi=0} < 0
\end{aligned}$$



$$\begin{aligned}
V(\phi)_{AdS} &= \left(-\frac{1}{l^2} + \alpha\phi\right)(4 + 2\cosh\phi) - 6\alpha\sinh\phi \\
V(\phi)_{flat} &= 2\alpha\phi(2 + \cosh\phi) - 6\alpha\sinh\phi \\
E_{\mu\nu} &= R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R - \kappa T_{\mu\nu}^\phi \\
T_{\mu\nu}^\phi &= \partial_\mu\phi\partial_\nu\phi - g_{\mu\nu}\left[\frac{1}{2}(\partial\phi)^2 + V(\phi)\right] \\
ds^2 &= \Omega(x)\left[-f(x)dt^2 + \frac{\eta^2 dx^2}{f(x)} + \frac{dy^2}{l^2} + \frac{dz^2}{l^2}\right] \\
E_t^t - E_x^x &= 0 \rightarrow \phi'^2 = \frac{3\Omega'^2 - 2\Omega''\Omega}{\Omega^2} \\
E_t^t - E_y^y &= 0 \rightarrow f'' + \frac{\Omega'f'}{\Omega} = 0 \\
E_t^t + E_y^y &= 0 \rightarrow V(\phi) = -\frac{1}{\Omega^2\eta^2}(f\Omega'' + f'\Omega') \\
\Omega(x) &= \frac{\eta^2 x^{\nu-1}}{\eta^2(x^\nu - 1)^2} \\
\phi'^2 &= \frac{(\nu-1)^2}{x^2} - \frac{4\nu(\nu-1)x^{\nu-2}}{x^\nu - 1} + \frac{4\nu^2 x^{\nu-1}}{(x^\nu - 1)^2} + \frac{2(\nu-1)}{x^2} + \frac{4\nu(1-\nu-x^\nu)x^{\nu-2}x^{\nu-2}}{(x^\nu - 1)^2} \\
\phi'^2 &= \frac{\nu^2 - 1}{2\kappa x^2} \rightarrow \int_\phi^{\phi=0} d\phi = \sqrt{\frac{\nu^2 - 1}{2\kappa}} \int_x^1 \frac{dx}{x} \\
\phi(x) &= l_\nu^{-1} \ln x, l_\nu^{-1} = \sqrt{\frac{\nu^2 - 1}{2\kappa}} \\
(f'\Omega)' &= 0 \\
f(x) &= \frac{c_2\eta^2}{\nu^2} \int \frac{(x^\nu - 1)^2}{x^{\nu-1}} dx + c_1 \\
f(x) &= c_1 + \frac{c_2\eta^2}{\nu^2} \left(\frac{x^{2+\nu}}{2+\nu} + \frac{x^{2-\nu}}{2-\nu} - x^2\right) \\
f(x) &= \frac{1}{l^2} + \alpha \left[\frac{1}{\nu^2 - 4} - \frac{x^2}{\nu^2} \left(1 + \frac{x^{-\nu}}{\nu - 2} - \frac{x^\nu}{\nu + 2}\right)\right] \\
V(\phi) &= \frac{\Lambda(\nu^2 - 4)}{6\kappa\nu^2} \left[\frac{\nu-1}{\nu+2} e^{-\phi l_\nu(\nu+1)} + \frac{\nu+1}{\nu-2} e^{\phi l_\nu(\nu-1)} + 4\frac{\nu^2-1}{\nu^2-4} e^{-\phi l_\nu}\right] \\
&+ \frac{\alpha}{\kappa\nu^2} \left[\frac{\nu-1}{\nu+2} \sinh \phi l_\nu(\nu+1) - \frac{\nu+1}{\nu-2} \sinh \phi l_\nu(\nu-1) + 4\frac{\nu^2-1}{\nu^2-4} \sinh \phi l_\nu\right] \\
V(\phi) &= \frac{\Lambda}{\kappa} - \frac{\phi^2}{l^2} + \frac{\kappa\Lambda(\nu^2-3)}{18} \frac{\phi^4}{\nu^2-1} - \frac{l_\nu^3}{90} (\Lambda\nu^2 - 4\Lambda - 6\alpha)\phi^5 + O(\phi^6) \\
V(0) &= \frac{\Lambda}{\kappa}, \left.\frac{dV}{d\phi}\right|_{\phi=0} = 0, \left.\frac{d^2V}{d\phi^2}\right|_{\phi=0} = -\frac{2}{l^2} \\
ds^2 &= \Omega(x) \left[-f(x)dt^2 + \frac{\eta^2 dx^2}{f(x)} + d\theta^2 + \sin^2\theta d\varphi^2\right] \\
f(x) &= \frac{1}{l^2} + \alpha \left[\frac{1}{\nu^2 - 4} - \frac{x^2}{\nu^2} \left(1 + \frac{x^{-\nu}}{\nu - 2} - \frac{x^\nu}{\nu + 2}\right)\right] + \frac{x}{\Omega(x)} \\
\Omega(x)|_{\nu=1} &= r^2 \Rightarrow x = 1 \pm \frac{1}{\eta r} \\
-g_{tt} &= \Omega(x)f(x) = k - \frac{\mu}{r} + \frac{r^2}{l^2}, \mu = \mp \frac{\alpha + 3\eta^2}{3\eta^3}
\end{aligned}$$

$$\begin{aligned}
I[g_{\mu\nu}, A_\mu, \phi] &= \frac{1}{16\pi G_N} \int d^4x \sqrt{-g} \left[R - \frac{1}{4} e^{\gamma\phi} F^2 - \frac{1}{2} \partial_\mu \phi \partial^\mu \phi - V(\phi) \right] \\
&\quad \nabla_\mu (e^{\gamma\phi} F^{\mu\nu}) = 0 \\
&\quad \frac{1}{\sqrt{-g}} \partial_\mu (\sqrt{-g} g^{\mu\nu} \partial_\nu \phi) - \frac{\partial V}{\partial \phi} - \frac{1}{4} \gamma e^{\gamma\phi} F^2 = 0 \\
&\quad R_{\mu\nu} - \frac{1}{2} g_{\mu\nu} R = \frac{1}{2} [T_{\mu\nu}^\phi + T_{\mu\nu}^{EM}] \\
T_{\mu\nu}^\phi &= \partial_\mu \phi \partial_\nu \phi - g_{\mu\nu} \left[\frac{1}{2} (\partial\phi)^2 + V(\phi) \right], T_{\mu\nu}^{EM} = e^{\gamma\phi} \left(F_{\mu\alpha} F_\nu^\alpha - \frac{1}{4} g_{\mu\nu} F^2 \right) \\
V(\phi) &= \left(\frac{\Lambda}{3} + \alpha\phi \right) (4 + 2\cosh(\phi)) - 6\alpha\sinh(\phi). \\
ds^2 &= \Omega(x) \left[-f(x) dt^2 + \frac{\eta^2 dx^2}{x^2 f(x)} + d\theta^2 + \sin^2 \theta d\varphi^2 \right] \\
f(x) &= \frac{1}{l^2} + \alpha \left[\frac{(x^2 - 1)}{2x} - \ln x \right] + \eta^2 \frac{(x - 1)^2}{x} - \frac{q^2 \eta^2}{2x^2} (x - 1)^3 \\
\Omega(x) &= \frac{x}{\eta^2 (x - 1)^2}, \phi(x) = \ln x \\
A &= q \left(\frac{1}{x} - \frac{1}{x_+} \right) dt, F = -\frac{q}{x^2} dx \wedge dt \\
I_g &= -\frac{1}{8\pi G_N} \int_{\partial\mathcal{M}} d^3x \sqrt{-h} \Xi(l, \mathcal{R}, \nabla\mathcal{R}) \\
ds^2 &= -\left(k + \frac{r^2}{l^2} \right) dt^2 + \left(k + \frac{r^2}{l^2} \right)^{-1} dr^2 + r^2 d\Sigma_k^2 \\
h_{ab} dx^a dx^b &= -\left(k + \frac{R^2}{l^2} \right) dt^2 + R^2 d\Sigma_k^2 \\
ds^2 &= \frac{r_b^2}{l^2} (-dt^2 + l^2 d\Sigma_k^2) \\
ds^2 &= -\left(1 - \frac{\mu}{r} + \frac{r^2}{l^2} \right) dt^2 + \left(1 - \frac{\mu}{r} + \frac{r^2}{l^2} \right)^{-1} dr^2 + r^2 (d\theta^2 + \sin^2 \theta d\varphi^2) \\
T &= \frac{f'}{4\pi} \Big|_{r_+} = \beta^{-1} = \frac{1}{4\pi} \left(\frac{3r_+^2 + l^2}{l^2 r_+} \right) \\
I_{bulk}^E &= \frac{12\pi\beta}{\kappa l^2} \int_{r_+}^R r^2 dr = \frac{4\pi\beta}{\kappa l^2} (R^3 - r_+^3) \\
h_{ab} dx^a dx^b &= -\left(1 - \frac{\mu}{R} + \frac{R^2}{l^2} \right) dt^2 + R^2 (d\theta^2 + \sin^2 \theta d\varphi^2) \\
n_a &= \frac{\delta_a^r}{\sqrt{g^{rr}}}, K_{ab} = \frac{\sqrt{g^{rr}}}{2} \partial_r h_{ab}, K = \frac{1}{l^2 R^2} \left(1 - \frac{\mu}{R} + \frac{R^2}{l^2} \right)^{-1/2} \left(-\frac{3l^2 \mu}{2} + 3R^3 + 2Rl^2 \right) \\
I_{GH}^E &= -\frac{4\pi\beta}{\kappa l^2} \left(-\frac{3l^2 \mu}{2} + 3R^3 + 2Rl^2 \right) \\
I_{bh}^E &= I_{bulk}^E + I_{GH}^E = \frac{4\pi\beta}{\kappa l^2} \left(\frac{3l^2 \mu}{2} - 2R^3 - 2Rl^2 - r_+^3 \right). \\
ds^2 &= -\left(1 + \frac{r^2}{l^2} \right) dt^2 + \left(1 + \frac{r^2}{l^2} \right)^{-1} dr^2 + r^2 (d\theta^2 + \sin^2 \theta d\varphi^2) \\
I_{Ads}^E &= I_{bulk}^E + I_{GH}^E = \frac{4\pi\beta_0}{\kappa l^2} (-2R^3 - 2l^2 R)
\end{aligned}$$

$$\begin{aligned}
\beta_0 \sqrt{1 + \frac{R^2}{l^2}} &= \beta \sqrt{1 + \frac{R^2}{l^2} - \frac{\mu}{R}} \\
I^E = I_{bh}^E - I_{Ads}^E &= \frac{4\pi\beta}{\kappa l^2} \left[\left(\frac{3l^2\mu}{2} - 2R^3 - 2Rl^2 - r_+^3 \right) - \frac{\beta_0}{\beta} (-2R^3 - 2l^2R) \right] \\
F &= \beta^{-1} I^E = \frac{4\pi}{\kappa l^2} \left(\frac{l^2\mu}{2} - r_+^3 \right) \\
I_g &= -\frac{1}{\kappa} \int_{\partial M} d^3x \sqrt{-h} \left(\frac{2}{l} + \frac{l\mathcal{R}}{2} \right). \\
I_{bulk}^E + I_{GH}^E &= \frac{4\pi\beta}{\kappa l^2} \left(\frac{3l^2\mu}{2} - 2r_b^3 - 2r_b l^2 - r_+^3 \right). \\
I_g^E &= \frac{4\pi\beta}{\kappa l^2} \left(1 + \frac{l^2}{R^2} - \frac{\mu l^2}{R^3} \right)^{\frac{1}{2}} (2R^3 + \kappa l^2 R) \Big|_{R=r_b} = \frac{4\pi\beta}{\kappa l^2} (2r_b^3 + 2l^2 r_b - \mu l^2), \\
I^E &= I_{bulk}^E + I_{GH}^E + I_g^E = \frac{4\pi\beta}{\kappa l^2} \left(\frac{l^2\mu}{2} - r_+^3 \right). \\
F_{curve} &= \frac{2\pi r_+}{\kappa}, F_{SAds} = \frac{4\pi}{\kappa l^2} \left(\frac{l^2\mu}{2} - r_+^3 \right), \\
E &= -T^2 \frac{\partial I^E}{\partial T} = \frac{\mu}{2G} \\
T_{curve} &= \frac{1}{4\pi r_h}, T_{SAds} = \frac{1}{4\pi r_+} \left(1 + \frac{3r_+^2}{l^2} \right). \\
ds^2 &= h_{ab} dx^a dx^b = -N(R) dt^2 + S(R) d\Sigma_k^2. \\
\tau^{ab} &\equiv \frac{2}{\sqrt{-h}} \frac{\delta I}{\delta h_{ab}} \\
ds^2 &= - \left(1 - \frac{\mu}{r} + \frac{r^2}{l^2} \right) dt^2 + \left(1 - \frac{\mu}{r} + \frac{r^2}{l^2} \right)^{-1} dr^2 + r^2 d\Omega^2. \\
ds^2 &= - \left(1 - \frac{\mu}{R} + \frac{R^2}{l^2} \right) dt^2 + R^2 d\Omega^2 \\
I &= \frac{1}{2\kappa} \int_M d^4x \sqrt{-g} (R - 2\Lambda) + \frac{1}{\kappa} \int_{\partial M} d^3x \sqrt{-h} K - \frac{1}{\kappa} \int_{\partial M} d^3x \sqrt{-h} \left(\frac{2}{l} + \frac{l\mathcal{R}}{2} \right), \\
\tau_{ab} &= -\frac{1}{8\pi G} \left(K_{ab} - h_{ab} K - \frac{2}{l} h_{ab} + l G_{ab} \right). \\
ds_{borde}^2 &= \frac{R^2}{l^2} (-dt^2 + l^2 d\Omega^2) \\
ds_{dual}^2 &= \gamma_{ab} dx^a dx^b = -dt^2 + l^2 d\Omega^2 \\
\langle \tau_{ab}^{dual} \rangle &= \lim_{R \rightarrow \infty} \frac{R}{l} \tau_{ab} = \frac{\mu}{16\pi G l^2} (3\delta_a^0 \delta_b^0 + \gamma_{ab}) \\
I[g_{\mu\nu}, \phi] &= \int_{\mathcal{M}} d^4x \sqrt{-g} \left[\frac{R}{2\kappa} - \frac{1}{2} (\partial\phi)^2 - V(\phi) \right] + \frac{1}{\kappa} \int_{\partial M} d^3x K \sqrt{-h} \\
\mathcal{H}_\perp &= \frac{2\kappa}{\sqrt{g}} \left[\pi_{ij} \pi^{ij} - \frac{1}{2} (\pi_i^i)^2 \right] - \frac{1}{2\kappa} \sqrt{g} {}^{(3)}R \\
&\quad + \frac{1}{2} \left(\frac{\pi_\phi^2}{\sqrt{g}} + \sqrt{g} g^{ij} \phi_{,i} \phi_{,j} \right) + \sqrt{g} V(\phi) \\
\mathcal{H}_i &= -2\pi_{ij}^j + \pi_\phi \phi_{,i} \\
ds^2 &= -(N^\perp)^2 dt^2 + g_{ij} (dx^i + N^i dt)(dx^j + N^j dt)
\end{aligned}$$

$$\begin{aligned}
H[\xi] &= \int_{\partial\mathcal{M}} d^3x (\xi^\perp \mathcal{H}_\perp + \xi^i \mathcal{H}_i) + Q[\xi] \\
\delta Q[\xi] &= \oint d^2S_l \left[\frac{G^{ijkl}}{2\kappa} (\xi^\perp \delta g_{ij|k} - \xi^\perp{}_{,k} \delta g_{ij}) + 2\xi_k \delta \pi^{kl} \right. \\
&\quad \left. + (2\xi^k \pi^{jl} - \xi^l \pi^{jk}) \delta g_{jk} - (\sqrt{g} \xi^\perp g^{lj} \phi_{,j} + \xi^l \pi_\phi) \delta \phi \right] \\
G^{ijkl} &\equiv \frac{1}{2} \sqrt{g} (g^{ik} g^{jl} + g^{il} g^{jk} - 2g^{ij} g^{kl}) \\
\xi^\perp &= N^\perp \xi^t, \xi^i = {}^{(3)}\xi^i + N^i \xi^t \\
\delta M &\equiv \delta Q[\partial_t] = \oint d^2S_l \left[\frac{G^{ijkl}}{2\kappa} (\xi^\perp \delta g_{ij|k} - \xi^\perp{}_{,k} \delta g_{ij}) - \sqrt{g} \xi^\perp g^{lj} \phi_{,j} \delta \phi \right] \\
\delta M &= \delta M_G + \delta M_\phi \\
\delta M_G &= \oint d^2S_l \frac{G^{ijkl}}{2\kappa} (\xi^\perp \delta g_{ij|k} - \xi^\perp{}_{,k} \delta g_{ij}) \\
\delta M_\phi &= - \oint d^2S_l \sqrt{g} \xi^\perp g^{lj} \phi_{,j} \delta \phi \\
d\bar{s}^2 &= \bar{g}_{\mu\nu} dx^\mu dx^\nu = - \left(k + \frac{r^2}{l^2} \right) dt^2 + \frac{dr^2}{k + \frac{r^2}{l^2}} + r^2 d\Sigma_k^2 \\
h_{ab} dx^a dx^b &= \frac{r^2}{l^2} (-dt^2 + l^2 d\Sigma_k^2) \\
m_{BF}^2 + \frac{1}{l^2} &> m^2 \geq m_{BF}^2, m_{BF}^2 = -\frac{9}{4l^2} \\
ds^2 &= -N(r) dt^2 + H(r) dr^2 + S(r) d\Sigma_k^2 \\
V(\phi) &= -\frac{3}{\kappa l^2} - \frac{\phi^2}{l^2} + O(\phi^4) \\
\phi(r) &= \frac{\alpha}{r} + \frac{\beta}{r^2} + O(r^{-3}) \\
N(r) &= -g_{tt} = \frac{r^2}{l^2} + k - \frac{\mu}{r} + O(r^{-2}) \\
S(r) &= r^2 + O(r^{-2}) \\
NS'^2 H - 2NS''HS + (NH)'S'S - 2\kappa NHS^2 \phi'^2 &= 0 \\
H(r) = g_{rr} &= \frac{l^2}{r^2} + \frac{l^4}{r^4} \left(-k - \frac{\alpha^2 \kappa}{2l^2} \right) + \frac{l^5}{r^5} \left(\frac{\mu}{l} - \frac{4\kappa\alpha\beta}{3l^3} \right) + O(r^{-6}) \\
g_{rr} &= \frac{l^2}{r^2} + \frac{al^4}{r^4} + \frac{bl^5}{r^5} + O(r^{-6}) \\
\beta &= C\alpha^2 \\
V(\phi) &= -\frac{3}{\kappa l^2} - \frac{\phi^2}{l^2} + \lambda\phi^3 + O(\phi^4) \\
\phi(r) &= \frac{\alpha}{r} + \frac{\beta}{r^2} + \frac{\gamma \ln(r)}{r^2} + O(r^{-3}) \\
H(r) = g_{rr} &= \frac{l^2}{r^2} + \frac{l^4}{r^4} \left(-k - \frac{\kappa\alpha^2}{2l^2} \right) + \frac{l^5}{r^5} \left(\frac{\mu}{l} - \frac{4\kappa\alpha\beta}{3l^3} + \frac{2\kappa\alpha\gamma}{9l^3} \right) + \frac{l^5 \ln r}{r^5} \left(-\frac{4\kappa\alpha\gamma}{3l^3} \right) + O \left[\frac{[\ln(r)]^2}{r^6} \right] \\
H(r) &= \frac{l^2}{r^2} + \frac{l^4 a}{r^4} + \frac{l^5 b}{r^5} + \frac{l^5 c \ln r}{r^5} + O \left[\frac{[\ln(r)]^2}{r^6} \right] \\
a = -k - \frac{\alpha^2 \kappa}{2l^2}, b &= \frac{\mu}{l} - \frac{4\kappa\alpha\beta}{3l^3} + \frac{2\kappa\alpha\gamma}{9l^3}, c = -\frac{4\kappa\alpha\gamma}{3l^3} \\
\partial_r \left(\frac{\phi' S \sqrt{N}}{\sqrt{H}} \right) - S \sqrt{NH} \frac{\partial V}{\partial \phi} &= 0 \\
\frac{3\alpha^2 l^2 \lambda + \gamma}{l^2} + O(r^{-1}) &= 0
\end{aligned}$$

$$\begin{aligned}
\xi^r &= r\eta^r(x^m) + O(r^{-1}) \\
\xi^m &= O(1) \\
\phi'(x) &= \phi(x) + \xi^\mu \partial_\mu \phi(x) = \frac{\alpha'}{r} + \frac{\beta'}{r^2} + \frac{\gamma' \ln(r)}{r^2} + O(r^{-3}), \\
\alpha' &= \alpha - \eta^r \alpha + \xi^m \partial_m \alpha \\
\beta' &= \beta - \eta^r (2\beta - \gamma) + \xi^m \partial_m \beta \\
\gamma' &= \gamma - 2\gamma \eta^r + \xi^m \partial_m \gamma \\
\alpha' \frac{\partial \gamma}{\partial \alpha} - \gamma' &= 0 = \alpha \frac{\partial \gamma}{\partial \alpha} - \gamma + \eta^r (2\gamma - \alpha \frac{\partial \gamma}{\partial \alpha}) + \xi^m \left(\frac{\partial \alpha}{\partial x^m} \frac{\partial \gamma}{\partial \alpha} - \frac{\partial \gamma}{\partial x^m} \right) \\
\alpha' \frac{\partial \beta}{\partial \alpha} - \beta' &= 0 = \alpha \frac{\partial \beta}{\partial \alpha} - \beta + \eta^r (2\beta - \gamma - \alpha \frac{\partial \beta}{\partial \alpha}) + \xi^m \left(\frac{\partial \alpha}{\partial x^m} \frac{\partial \beta}{\partial \alpha} - \frac{\partial \beta}{\partial x^m} \right) \\
\beta(\alpha) &= (-C_\gamma \ln(\alpha) + C) \alpha^2 \\
\phi(r) &= \frac{\alpha}{r^{\Delta_-}} + \frac{\beta}{r^{\Delta_+}} + \dots \\
I_{CFT} &\rightarrow I_{CFT} - \int d^3x W[\mathcal{O}(x)] \\
I_g^{ct} &= -\frac{1}{\kappa} \int_{\partial\mathcal{M}} d^3x \sqrt{-h} \left(\frac{2}{l} + \frac{\mathcal{R}l}{2} \right) \\
I &= \int d^4x \sqrt{-g} \left(\frac{R}{2\kappa} - \frac{(\partial\phi)^2}{2} - V(\phi) \right) + \frac{1}{\kappa} \int_{\partial\mathcal{M}} d^3x \sqrt{-h} K - \frac{1}{\kappa} \int_{\partial\mathcal{M}} d^3x \sqrt{-h} \left(\frac{2}{l} + \frac{\mathcal{R}l}{2} \right) + I_\phi \\
I_\phi^{ct} &= \frac{1}{6\kappa} \int_{\partial\mathcal{M}} d^3x \sqrt{-h} \left(\phi n^\nu \partial_\nu \phi - \frac{\phi^2}{2l} \right) \\
I_\phi &= - \int_{\partial\mathcal{M}} d^3x \sqrt{-h} \left[\frac{\phi^2}{2l} + \frac{W(\alpha)}{l\alpha^3} \phi^3 \right] \\
\delta I &= \int d^3x \sqrt{-h} \left[\frac{1}{r} \left(-\sqrt{g^{rr}} \phi' - \frac{\phi}{l} - \frac{3W(\alpha)\phi^2}{l\alpha^3} \right) \left(1 + \frac{1}{r} \frac{d^2W(\alpha)}{d\alpha^2} \right) + \left(\frac{3W(\alpha)}{\alpha} - \beta \right) \frac{\phi^3}{l\alpha^3} \right] \delta\alpha.
\end{aligned}$$

$$\begin{aligned}
\lim_{r \rightarrow \infty} \delta I &= 0 \\
I_\phi &= - \int_{\partial\mathcal{M}} d^3x \sqrt{-h} \left[\frac{\phi^2}{2l} + \frac{\phi^3}{l\alpha^3} \left(W - \frac{\alpha\gamma}{3} \right) - \frac{\phi^3 C_\gamma}{3l} \ln \left(\frac{\phi}{\alpha} \right) \right] \\
ds^2 &= \Omega(x) \left[-f(x) dt^2 + \frac{\eta^2 dx^2}{f(x)} + d\Sigma_k^2 \right] \\
I_{bulk}^E &= \int_0^{1/T} d\tau \int_{x_+}^{x_b} d^3x \sqrt{g^E} V(\phi) = \frac{\sigma_k}{2\eta\kappa T} \frac{d(\Omega f)}{dx} \Big|_{x_+}^{x_b} \\
\Omega(x) \rightarrow S(r), f(x) &\rightarrow \frac{N(r)}{S(r)}, dx \rightarrow \frac{\sqrt{NH}}{\eta S} dr \\
I_{bulk}^E &= \frac{\sigma_k}{2\kappa T} \frac{S}{\sqrt{NH}} \frac{dN}{dr} \Big|_{r_+}^{r_b} \\
ds^2 &= h_{ab} dx^a dx^b = \Omega(x_0) [-f(x_0) dt^2 + d\Sigma_k] \\
n_a &= \frac{\delta_a^x}{\sqrt{g^{xx}}} \Big|_{x=x_0}, K_{ab} = \frac{\sqrt{g^{xx}}}{2} \partial_x g_{ab} \Big|_{x=x_0}, K = \frac{1}{2\eta} \left(\frac{f}{\Omega} \right)^{1/2} \left[\frac{(\Omega f)'}{\Omega f} + \frac{2\Omega'}{\Omega} \right] \Big|_{x_0} \\
I_{GH}^E &= - \frac{\sigma_k}{\kappa T} \frac{\Omega f}{2\eta} \left[\frac{(\Omega f)'}{\Omega f} + \frac{2\Omega'}{\Omega} \right] \Big|_{x_b} = - \frac{\sigma_k}{2T\kappa} \left(\frac{S}{\sqrt{NH}} \frac{dN}{dr} + \frac{2N}{\sqrt{NH}} \frac{dS}{dr} \right) \Big|_{r_b} \\
I_g^{ct} &= \frac{2\sigma_k}{\kappa T l} \left(\Omega^{3/2} f^{1/2} + \frac{l^2 k}{2} f^{1/2} \Omega^{1/2} \right) \Big|_{x_b} = \frac{2\sigma_k}{\kappa T l} S \sqrt{N} \left(1 + \frac{l^2 k}{2S} \right) \Big|_{r_b} \\
T &= \frac{N'}{4\pi \sqrt{NH}} \Big|_{r_+}
\end{aligned}$$



$$\begin{aligned}
I_{bulk}^E + I_{GH}^E + I_g^{ct} &= -\frac{1}{T} \left[\frac{\sigma_k S(r_+) T}{4G} \right] - \frac{\sigma_k}{2\kappa T} \left[\frac{2N}{\sqrt{NH}} \frac{dS}{dr} - \frac{4}{l} S \sqrt{N} \left(1 + \frac{l^2 k}{2S} \right) \right] \Big|_{r_b} \\
I_{bulk}^E + I_{GH}^E + I_g^{ct} &= -\frac{\mathcal{A}}{4G} - \frac{\sigma_k}{T} \left(-\frac{\mu}{\kappa} + \frac{4\alpha\beta}{3l^2} + \frac{r\alpha^2}{2l^2} \right) \Big|_{r_b} \\
I_\phi^{ct} &= \int_{\partial\mathcal{M}} d^3x \sqrt{h^E} \left[\frac{\phi^2}{2l} + \frac{W(\alpha)}{l\alpha^3} \phi^3 \right] = \frac{\sigma_k}{T} \left(\frac{W}{l^2} + \frac{\alpha\beta}{l^2} + \frac{r\alpha^2}{2l^2} \right) \Big|_{r_\infty} \\
I^E = I_{bulk}^E + I_{GH}^E + I_g^{ct} + I_\phi^{ct} &= -\frac{\mathcal{A}}{4G} + \frac{\sigma_k}{T} \left[\frac{\mu}{\kappa} + \frac{1}{l^2} \left(W - \frac{\alpha}{3} \frac{dW}{d\alpha} \right) \right] \\
F = I^E T &= M - TS \\
M = -T^2 \frac{\partial I^E}{\partial T} &= \sigma_k \left[\frac{\mu}{\kappa} + \frac{1}{l^2} \left(W - \frac{\alpha}{3} \frac{dW}{d\alpha} \right) \right] \\
S = -\frac{\partial(I^E T)}{\partial T} &= \frac{\mathcal{A}}{4G} \\
I_{bulk}^E + I_{GH}^E + I_g^{ct} + I_\phi^{ct} &= -\frac{\mathcal{A}}{4G} + \frac{\sigma_k}{T} \left[\frac{\mu}{\kappa} + \frac{1}{l^2} \left[W(\alpha) - \frac{\alpha}{3} \frac{dW}{d\alpha} + \frac{2\alpha\gamma}{9} - \frac{\alpha\gamma}{3} \ln r \right] \right] \\
\bar{I}_\phi^{ct} &= \int_{\partial\mathcal{M}} d^3x \sqrt{h^E} \left\{ \frac{\phi^3 \gamma}{3\alpha^2 l} \left[\ln \left(\frac{\alpha}{\phi} \right) - 1 \right] \right\} = \frac{\sigma_k}{T} \left[-\frac{\alpha\gamma}{3l^2} + \frac{\alpha\gamma \ln r}{3l^2} + O(r^{-1} \ln r) \right] \\
I^E = I_{bulk}^E + I_{GH}^E + I_g^{ct} + I_\phi^{ct} + \bar{I}_\phi^{ct} &= -\frac{\mathcal{A}}{4G} + \frac{\sigma_k}{T} \left[\frac{\mu}{\kappa} + \frac{1}{l^2} \left(W - \frac{\alpha}{3} \frac{dW}{d\alpha} - \frac{\alpha\gamma}{9} \right) \right] \\
M = -T^2 \frac{\partial I^E}{\partial T} &= \sigma_k \left[\frac{\mu}{\kappa} + \frac{1}{l^2} \left(W - \frac{\alpha}{3} \frac{dW}{d\alpha} - \frac{\alpha\gamma}{9} \right) \right] \\
S = -\frac{\partial(I^E T)}{\partial T} &= \frac{\mathcal{A}}{4G} \\
\tau_{ab} &= -\frac{1}{\kappa} \left(K_{ab} - h_{ab} K + \frac{2}{l} h_{ab} - l G_{ab} \right) - \frac{h_{ab}}{l} \left[\frac{\phi^2}{2} + \frac{W(\alpha)}{\alpha^3} \phi^3 \right]. \\
\tau_{tt} &= \frac{l}{R} \left[\frac{\mu}{8\pi G l^2} + \frac{1}{l^4} \left(W - \frac{\alpha\beta}{3} \right) \right] + O(R^{-2}) \\
\tau_{\theta\theta} &= \frac{l}{R} \left[\frac{\mu}{16\pi G} - \frac{1}{l^2} \left(W - \frac{\alpha\beta}{3} \right) \right] + O(R^{-2}) \\
\tau_{\phi\phi} &= \frac{l \sin^2 \theta}{R} \left[\frac{\mu}{16\pi G} - \frac{1}{l^2} \left(W - \frac{\alpha\beta}{3} \right) \right] + O(R^{-2}) \\
\langle \tau_{ab}^{dual} \rangle &= \frac{3\mu}{16\pi G l^2} \delta_a^0 \delta_b^0 + \frac{\gamma_{ab}}{l^2} \left[\frac{\mu}{16\pi G} - \frac{1}{l^2} \left(W(\alpha) - \frac{\alpha\beta}{3} \right) \right] \\
\langle \tau^{dual} \rangle &= -\frac{3}{l^4} \left[W(\alpha) - \frac{\alpha\beta}{3} \right] \\
\tau_{ab} &= -\frac{1}{\kappa} \left(K_{ab} - h_{ab} K + \frac{2}{l} h_{ab} - l G_{ab} \right) - \frac{h_{ab}}{l} \left[\frac{\phi^2}{2} + \frac{\phi^3}{\alpha^3} \left(W - \frac{\alpha\gamma}{3} \right) + \frac{\phi^3 \gamma}{3\alpha^2} \ln \left(\frac{\alpha}{\phi} \right) \right] \\
\tau_{tt} &= \frac{l}{R} \left[\frac{\mu}{8\pi G l^2} + \frac{1}{l^4} \left(W - \frac{\alpha\beta}{3} - \frac{\alpha\gamma}{9} \right) \right] + O \left[\frac{(\ln R)^3}{R^2} \right] \\
\tau_{\theta\theta} &= \frac{l}{R} \left[\frac{\mu}{16\pi G} - \frac{1}{l^2} \left(W - \frac{\alpha\beta}{3} - \frac{\alpha\gamma}{9} \right) \right] + O \left[\frac{(\ln R)^3}{R^2} \right] \\
\tau_{\phi\phi} &= \frac{l \sin^2 \theta}{R} \left[\frac{\mu}{16\pi G} - \frac{1}{l^2} \left(W - \frac{\alpha\beta}{3} - \frac{\alpha\gamma}{9} \right) \right] + O \left[\frac{(\ln R)^3}{R^2} \right] \\
\langle \tau_{ab}^{dual} \rangle &= \frac{3\mu}{16\pi G l^2} \delta_a^0 \delta_b^0 + \frac{\gamma_{ab}}{l^2} \left[\frac{\mu}{16\pi G} - \frac{1}{l^2} \left(W(\alpha) - \frac{\alpha\beta}{3} - \frac{\alpha\gamma}{9} \right) \right] \\
\langle \tau^{dual} \rangle &= -\frac{3}{l^4} \left(W - \frac{\alpha\beta}{3} - \frac{\alpha\gamma}{9} \right) \\
\langle \tau^{dual} \rangle = 0 &\Rightarrow \gamma = -3l^2 \lambda \alpha^2, W(\alpha) = \alpha^3 [C + l^2 \lambda \ln \alpha] \\
&\quad -5/4l^2 > m^2 \geq -9/4l^2
\end{aligned}$$

$$\begin{aligned} \phi(r) &= \frac{\alpha}{r} + \frac{\beta}{r^2} + O(r^{-3}) \\ \delta M_G &= \frac{\sigma_k}{\kappa} [r\delta a + l\delta b + O(1/r)] \\ \delta M_\phi &= \frac{\sigma_k}{l^2} [r\alpha\delta a + \alpha\delta\beta + 2\beta\delta\alpha + O(1/r)] \\ \delta M &= \frac{\sigma_k}{\kappa l^2} [r(l^2\delta a + \kappa\alpha\delta a) + l^3\delta b + \kappa(\alpha\delta\beta + 2\beta\delta\alpha) + O(1/r)] \\ &\quad \frac{k+a}{\kappa} + \frac{\alpha^2}{2l^2} = 0 \\ \delta M &= \frac{\sigma_k}{\kappa l^2} [l^3\delta b + \kappa(\alpha\delta\beta + 2\beta\delta\alpha)] \\ M &= \sigma_k \left[\frac{lb}{\kappa} + \frac{1}{l^2} \left(\alpha \frac{dW(\alpha)}{d\alpha} + W(\alpha) \right) \right] \\ &\quad \frac{lc}{\kappa} - 4\alpha^3\lambda = 0 \\ \delta M_G &= \left\{ \frac{l\delta b}{\kappa} + \frac{\delta a}{\kappa} r + \frac{l\delta c}{\kappa} \ln(r) + O\left(\frac{\ln(r)^2}{r}\right) \right\} \sigma_k \\ \delta M_\phi &= \left[\frac{\alpha\delta\beta + 2\beta\delta\alpha + 3\alpha^2 l^2 \lambda \delta\alpha}{l^2} + r \frac{\alpha\delta\alpha}{l^2} \right. \\ &\quad \left. - 12\lambda\alpha^2 \delta\alpha \ln(r) + O\left(\frac{\ln(r)^2}{r}\right) \right] \sigma_k \\ \delta M &= \left[\frac{l\delta b}{\kappa} + \frac{\alpha\delta\beta + 2\beta\delta\alpha + 3\alpha^2 l^2 \lambda \delta\alpha}{l^2} \right] \sigma_k \\ M &= \left[\frac{lb}{\kappa} + \frac{1}{l^2} \left(\alpha \frac{dW}{d\alpha} + W(\alpha) + \alpha^3 l^2 \lambda \right) \right] \sigma_k \\ M &= \left[\frac{\mu}{\kappa} + \frac{1}{l^2} \left(W(\alpha) - \frac{1}{3} \alpha \frac{dW}{d\alpha} + \frac{1}{3} \alpha^3 l^2 \lambda \right) \right] \sigma_k \\ &\quad W(\alpha) = \alpha^3 [C + l^2 \lambda \ln(\alpha)] \\ d\Sigma_k^2 &= \frac{dy^2}{1-ky^2} + (1-ky^2)d\phi^2 \\ E &= \int d\sigma^i \tau_{ij} \xi^j = \int dy d\phi S u^i \tau_{ij} \xi^j \\ ds^2 &= \sigma_{ij} dx^i dx^j = S d\Sigma_k^2 \\ E &= \sigma_k \left[\frac{\mu}{\kappa} + \frac{1}{l^2} \left(W - \frac{\alpha}{3} \frac{dW}{d\alpha} \right) \right] \\ E &= \sigma_k \left[\frac{\mu}{\kappa} + \frac{1}{l^2} \left(W - \frac{1}{3} \alpha \frac{dW}{d\alpha} - \frac{\alpha\gamma}{9} \right) \right] = \sigma_k \left[\frac{\mu}{\kappa} + \frac{1}{l^2} \left(W - \frac{1}{3} \alpha \frac{dW}{d\alpha} - \frac{\alpha^3 C_\gamma}{9} \right) \right]. \\ V(\phi) &= \frac{\Lambda(v^2-4)}{6\kappa v^2} \left[\frac{v-1}{v+2} e^{-\phi l_\nu(v+1)} + \frac{v+1}{v-2} e^{\phi l_\nu(v-1)} + 4 \frac{v^2-1}{v^2-4} e^{-\phi l_\nu} \right] \\ &+ \frac{\Upsilon}{\kappa v^2} \left[\frac{v-1}{v+2} \sinh \phi l_\nu(v+1) - \frac{v+1}{v-2} \sinh \phi l_\nu(v-1) + 4 \frac{v^2-1}{v^2-4} \sinh \phi l_\nu \right] \\ \phi(x) &= l_\nu^{-1} \ln x \\ f(x) &= \frac{1}{l^2} + \Upsilon \left[\frac{1}{v^2-4} - \frac{x^2}{v^2} \left(1 + \frac{x^{-v}}{v-2} - \frac{x^v}{v+2} \right) \right] + \frac{x}{\Omega(x)} \\ \Omega(x) &= r^2 + O(r^{-3}) \\ x &= 1 + \frac{1}{\eta r} + \frac{m}{r^3} + \frac{n}{r^4} + \frac{p}{r^5} + O(r^{-6}) \\ \Omega(x) &= r^2 - \frac{24m\eta^3 + v^2 - 1}{12\eta^2} - \frac{24n\eta^4 - v^2 + 1}{12\eta^3 r} + \frac{720m^2\eta^6 - 480\eta\eta^5 + v^4 - 20v^2 + 19}{240\eta^4 r^2} \\ &\quad + O(r^{-3}) \end{aligned}$$

$$\begin{aligned}
x &= 1 + \frac{1}{\eta r} - \frac{(v^2 - 1)}{24\eta^3 r^3} \left[1 - \frac{1}{\eta r} - \frac{9(v^2 - 9)}{80\eta^2 r^2} \right] + O(r^{-6}) \\
-g_{tt} &= f(x)\Omega(x) = \frac{r^2}{l^2} + 1 + \frac{Y + 3\eta^2}{3\eta^3 r} + O(r^{-3}) \\
g_{rr} &= \frac{\Omega(x)\eta^2}{f(x)} \left(\frac{dx}{dr} \right) = \frac{l^2}{r^2} - \frac{l^4}{r^4} - \frac{l^2(v^2 - 1)}{4\eta^2 r^4} - \frac{l^2(3\eta^2 l^2 + Yl^2 - v^2 + 1)}{3\eta^3 r^5} + O(r^{-6}) \\
\phi(x) &= l_v^{-1} \ln x = \frac{1}{l_v \eta r} - \frac{1}{2l_v \eta^2 r^2} - \frac{v^2 - 9}{24\eta^3 r^3} + O(r^{-4}) \\
M &= \sigma \left[\frac{\mu}{\kappa} + \frac{1}{l^2} \left(W - \frac{\alpha}{3} \frac{dW}{d\alpha} \right) \right] \\
M &= -\frac{\sigma}{\kappa} \left(\frac{3\eta^2 + Y}{3\eta^3} \right) \\
I_{CFT} &\rightarrow I_{CFT} + \frac{l_v}{6} \int d^3 x \mathcal{O}^3 \\
ds^2 &= -dt^2 + dx^2 + dy^2 + dz^2 \\
t &= \rho \sinh(\alpha\tau), x = \rho \cosh(\alpha\tau) \\
ds^2 &= -a^2 \rho^2 d\tau^2 + d\rho^2 + dy^2 + dz^2 \\
ds^2 &= N(r) dt_E^2 + H(r) dr^2 + S(r) d\Sigma_k^2 \\
ds^2 &= g_{rr} \frac{4N}{[(N)']^2} \left[\rho^2 \frac{[(N)']^2}{4Ng_{rr}} d\tau_E^2 + d\rho^2 \right] \\
T &= \frac{1}{\beta} = \frac{(N^2)'}{4\pi \sqrt{N^2 g_{rr}}} \Big|_H \\
T_{flat} &= \frac{1}{4\pi r_h}, T_{RN-flat} = \frac{1}{4\pi r_+} \left(1 - \frac{q^2}{4r_+^2} \right) \\
T_{Sch-Ads} &= \frac{1}{4\pi r_+} \left(1 + \frac{3r_+^2}{l^2} \right), T_{RN-Ads} = \frac{1}{4\pi r_+} \left(1 + \frac{3r_+^2}{l^2} - \frac{q^2}{4r_+^2} \right) \\
ds^2 &= \Omega(x) \left[-f(x) dt^2 + \frac{\eta^2 dx^2}{f(x)} + d\theta^2 + d\Sigma_k^2 \right] \\
f(x) &= \frac{1}{l^2} + \alpha \left[\frac{1}{v^2 - 4} - \frac{x^2}{v^2} \left(1 + \frac{x^{-v}}{v - 2} - \frac{x^v}{v + 2} \right) \right] + \frac{kx}{\Omega(x)} \\
f'\Omega(x) &= \frac{\alpha}{\eta^2} + 2k + kv \frac{x^v + 1}{x^v - 1} \\
T &= \frac{f'}{4\pi\eta} \Big|_{x_h} = \frac{1}{4\pi\eta\Omega(x_h)} \left(\frac{\alpha}{\eta^2} + 2k + kv \frac{x_h^v + 1}{x_h^v - 1} \right) \\
ds^2 &= \Omega(x) \left(-f(x) dt^2 + \frac{\eta^2 dx^2}{f(x)} + d\theta^2 + \sin^2 \theta d\varphi^2 \right) \\
\Omega(x) &= \frac{1}{\eta^2 (x - 1)^2}, f(x) = \frac{1}{l^2} + \frac{1}{3} \alpha (x - 1)^3 + \eta^2 x (x - 1)^2 \\
x &= 1 + \frac{1}{\eta r}, x = 1 - \frac{1}{\eta r} \\
\Omega(x)f(x) &= F(r) = 1 - \frac{\mu}{r} + \frac{r^2}{l^2}, \mu = \frac{\alpha + 3\eta^2}{3\eta^3} \\
I[g_{\mu\nu}] &= I_{bulk} + I_{GH} - \frac{1}{\kappa} \int_{\partial\mathcal{M}} d^3 x \sqrt{-h} \left(\frac{2}{l} + \frac{\mathcal{R}l}{2} \right) \\
E_t^t - E_\phi^\phi &= 0 \Rightarrow 0 = f'' + \frac{\Omega' f'}{\Omega} + 2\eta^2 \\
E_t^t + E_\phi^\phi &= 0 \Rightarrow 2\kappa V(\phi) = -\frac{(f\Omega'' + f'\Omega')}{\Omega^2 \eta^2} + \frac{2}{\Omega}
\end{aligned}$$

$$I_{bulk}^E = \frac{4\pi\beta}{\eta^3 \kappa l^2} \left[-\frac{1}{(x_b - 1)^3} + \frac{1}{(x_h - 1)^3} \right] = \frac{4\pi\beta}{\kappa l^2} (r_b^3 - r_h^3)$$

$$ds^2 = h_{ab} dx^a dx^b = \Omega(x) [-f(x) dt^2 + d\theta^2 + \sin^2 \theta d\phi^2]$$

$$n_a = \frac{\delta_a^x}{\sqrt{g^{xx}}}, K_{ab} = \frac{\sqrt{g^{xx}}}{2} \partial_x h_{ab}$$

$$I_{GH}^E = -\frac{2\pi\beta}{\kappa} \left[-\frac{6}{l^2 \eta^3 (x-1)^3} - \frac{4}{\eta(x-1)} - \left(\frac{\alpha + 3\eta^2}{\eta^3} \right) \right] \Big|_{x_b} = -\frac{2\pi\beta}{\kappa} \left(\frac{6r_b^3}{l^2} + 4r_b - 3\mu \right)$$

$$I_g^E = \frac{2\pi\beta}{\kappa} \left[\frac{4}{l^2 \eta^3 (x_b - 1)^3} + \frac{4}{\eta(x_b - 1)} - 2\mu \right] = \frac{2\pi\beta}{\kappa} \left(\frac{4r_b^3}{l^2} + 4r_b - 2\mu \right)$$

$$I^E = I_{bulk}^E + I_{GH}^E + I_g^E = \frac{4\pi\beta}{\kappa l^2} \left[\frac{1}{\eta^3 (x_h - 1)^3} + \frac{\mu l^2}{2} \right] = \frac{4\pi\beta}{\kappa l^2} \left(-r_h^3 + \frac{\mu l^2}{2} \right)$$

$$I_\phi^E = \int_{\partial\mathcal{M}} d^3x^E \sqrt{h^E} \left(\frac{\phi^2}{2l} - \frac{l_\nu}{6l} \phi^3 \right) = \frac{4\pi\beta}{\kappa} \left[-\frac{v^2 - 1}{4l^2 \eta^3 (x_b - 1)} + \frac{v^2 - 1}{3l^2 \eta^3} \right]$$

$$I_{bulk}^E + I_{surf}^E + I_g^E = -\frac{1}{T} \left(\frac{AT}{4G} \right) + \frac{4\pi\beta}{\kappa} \left[\frac{v^2 - 1}{4l^2 \eta^3 (x_b - 1)} + \frac{12\eta^2 l^2 + 4\alpha l^2 - 4v^2 + 4}{12l^2 \eta^3} \right]$$

$$I^E = \beta \left(-\frac{AT}{4G} + \frac{4\pi}{\kappa} \frac{3\eta^2 + \alpha}{3\eta^3} \right)$$

$$M = \frac{1}{2G} \left(\frac{\alpha + 3\eta^2}{3\eta^3} \right)$$

$$T = \frac{f'(x)}{4\pi\eta} \Big|_{x=x_h} = \frac{1}{4\pi\eta\Omega(x_h)} \left[\frac{\alpha}{\eta^2} + 2 + v \frac{x_h^\nu + 1}{x_h^\nu - 1} \right], S = \frac{A}{4G} = \frac{4\pi\Omega(x_h)}{4G},$$

$$\frac{\partial M}{\partial \eta} \frac{d\eta}{dx_h} = T \left(\frac{\partial S}{\partial x_h} + \frac{\partial S}{\partial \eta} \frac{d\eta}{dx_h} \right).$$

$$M = -\frac{1}{2G} \left(\frac{\alpha + 3\eta^2}{3\eta^3} \right)$$

$$T = -\frac{1}{4\pi\eta\Omega(x_h)} \left[\frac{\alpha}{\eta^2} + 2 + v \frac{x_h^\nu + 1}{x_h^\nu - 1} \right]$$

$$C = \frac{\partial M}{\partial T}$$

$$I[g_{\mu\nu}] = \int_{\mathcal{M}} d^4x (R - 2\Lambda) \sqrt{-g} + 2 \int_{\partial\mathcal{M}} d^3x K \sqrt{-h} - \int_{\partial\mathcal{M}} d^3x \frac{4}{l} \sqrt{-h}$$

$$ds^2 = -\left(-\frac{\mu_b}{r} + \frac{r^2}{l^2} \right) dt^2 + \left(-\frac{\mu_b}{r} + \frac{r^2}{l^2} \right)^{-1} dr^2 + \frac{r^2}{l^2} (dx_1^2 + dx_2^2)$$

$$I_b^E = \frac{2LL_b\beta_b}{l^4} \left(-r_h^3 + \frac{\mu_b l^2}{2} \right) = -\frac{LL_b\beta_b r_h^3}{l^4}$$

$$T = \beta_b^{-1} = \frac{(-g_{tt})'}{4\pi} \Big|_{r=r_h} = \frac{3r_h}{4\pi l^2}$$

$$E = -T^2 \frac{\partial I_b^E}{\partial T} = \frac{2LL_b\mu_b}{l^2}$$

$$S = -\frac{\partial(I_b^E T)}{\partial T} = \frac{LL_b r_h^2}{4l^2 G} = \frac{\mathcal{A}}{4G}$$

$$ds_{dual}^2 = \frac{l^2}{R^2} ds^2 = \gamma_{ab} dx^a dx^b = -dt^2 + dx_1^2 + dx_2^2$$

$$\langle \tau_{ab}^{dual} \rangle = \lim_{R \rightarrow \infty} \frac{R}{l} \tau_{ab} = \frac{\mu_b}{16\pi G_N l^2} [3\delta_a^0 \delta_b^0 + \gamma_{ab}]$$

$$E = Q_{\xi_t} = \int d\Sigma^i \tau_{ij} \xi^j = \frac{LL_b}{l^2 \kappa} \left[\mu_b + \frac{l^2}{4R} + O(R^{-2}) \right]$$

$$ds^2 = -\frac{r^2}{l^2} d\tau^2 + \left(-\frac{\mu_s}{r} + \frac{r^2}{l^2}\right)^{-1} dr^2 + \left(-\frac{\mu_s}{r} + \frac{r^2}{l^2}\right) d\theta^2 + \frac{r^2}{l^2} dx_2^2$$

$$-\frac{\mu_s}{r_s} + \frac{r_s^2}{l^2} = 0$$

$$L_s = \frac{4\pi\sqrt{g_{\theta\theta}g_{rr}}}{(g_{\theta\theta})'} \Big|_{r=r_s} = \frac{4\pi l^2}{3r_s}$$

$$I_s^E = -\frac{LL_s\beta_s\mu_s}{l^2}$$

$$M = -\frac{LL_s\mu_s}{l^2}$$

$$\Delta I = I_b^E - I_s^E = \frac{L}{2\kappa l^4} \left(\frac{4\pi l^2}{3}\right)^3 L_b\beta_b(L_s^{-3} - \beta_b^{-3}) = \frac{L}{2\kappa l^4} \left(\frac{4\pi l^2}{3}\right)^3 L_b\beta_b \left(\frac{1}{L_s^3} - T^3\right)$$

$$\frac{\mathcal{A}}{Tl^3} = \frac{L}{l} \left(\frac{4\pi}{3}\right)^2 L_s T$$

$$I[g_{\mu\nu}, \phi] = \int_{\mathcal{M}} d^4x \sqrt{-g} \left[R - \frac{(\partial\phi)^2}{2} - V(\phi) \right] + 2 \int_{\partial\mathcal{M}} d^3x K \sqrt{-h}$$

$$V(\phi) = \frac{\Lambda(v^2 - 4)}{3v^2} \left[\frac{v-1}{v+2} e^{-\phi l_\nu(v+1)} + \frac{v+1}{v-2} e^{\phi l_\nu(v-1)} + 4 \frac{v^2-1}{v^2-4} e^{-\phi l_\nu} \right]$$

$$+ \frac{2\alpha}{v^2} \left[\frac{v-1}{v+2} \sinh \phi l_\nu(v+1) - \frac{v+1}{v-2} \sinh \phi l_\nu(v-1) + 4 \frac{v^2-1}{v^2-4} \sinh \phi l_\nu \right]$$

$$V(\phi) = \frac{2\Lambda}{27} \left(5e^{-\phi\sqrt{2}} + 10e^{\phi\sqrt{2}/2} + 16e^{-\phi\sqrt{2}/4} \right)$$

$$+ \frac{4\alpha}{45} [\sinh(\phi\sqrt{2}) - 10\sinh(\phi\sqrt{2}/2) + 16\sinh(\phi\sqrt{2}/4)]$$

$$\Omega(x) = \frac{9x^2}{\eta^2(x^3 - 1)^2}$$

$$\phi(x) = 2\sqrt{2} \ln x$$

$$f(x) = \frac{1}{l^2} + \alpha \left[\frac{1}{5} - \frac{x^2}{9} \left(1 + x^{-3} - \frac{x^3}{5} \right) \right]$$

$$ds^2 = \frac{R^2}{l^2} [-dt^2 + dx_1^2 + dx_2^2]$$

$$R^2 \equiv \frac{1}{\eta^2(x-1)^2}$$

$$ds_{dual}^2 = \frac{l^2}{R^2} ds^2 = \gamma_{ab} dx^a dx^b = -dt^2 + dx_1^2 + dx_2^2$$

$$I_{BH}^E = \beta_b \left(-\frac{\mathcal{A}T}{4G_N} + \frac{2LL_b}{l^2} \frac{\alpha}{3\eta^3} \right) = -\frac{LL_b\alpha\beta_b}{3l^2\eta^3}$$

$$\mathcal{A} = \frac{LL_b\Omega(x_h)}{l^2}, T = \frac{\alpha}{4\pi\eta^3\Omega}$$

$$M_b = \frac{2LL_b\mu_b}{l^2}, \mu_b = \frac{\alpha}{3\eta^3}$$

$$ds^2 = \Omega_s(x) \left[-\frac{d\tau^2}{l^2} + \frac{\lambda^2 dx^2}{f(x)} + f(x) d\theta^2 + \frac{dx_2^2}{l^2} \right]$$

$$\Omega_s(x) = \frac{9x^2}{\lambda^2(x^3 - 1)^2}$$

$$L_s = \frac{4\pi\lambda}{f'} \Big|_{x=x_s} = \frac{4\pi\lambda^3\Omega_s}{\alpha}$$



$$\begin{aligned}
I_{\text{soliton}}^E &= -\frac{L\beta_s\Omega_s(x_s)}{4l^2G_N} + \frac{2LL_s\beta_s}{l^2} \frac{\alpha}{3\lambda^3} = -\frac{LL_s\beta_s}{l^2} \left(\frac{\alpha}{3\lambda^3}\right) \\
M_{\text{soliton}} &= -\frac{LL_s\mu_s}{l^2}, \mu_s = \frac{\alpha}{3\lambda^3} \\
I^\phi &= -\int d^3x\sqrt{-h} \left(\frac{\phi^2}{2l} - \frac{l_\nu}{6l}\phi^3\right), \tau_{ab}^\phi = -\frac{2}{\sqrt{-h}} \frac{\delta I^\phi}{\delta h^{ab}} \\
\tau_{ab} &= -\frac{1}{\kappa} \left(K_{ab} - h_{ab}K + \frac{2}{l}h_{ab} - lE_{ab}\right) - \frac{h_{ab}}{l} \left(\frac{\phi^2}{2} - \frac{l_\nu}{6}\phi^3\right) \\
\tau_{\tau\tau} &= \frac{\alpha(x-1)}{3\lambda^2l} + O[(x-1)^2] \\
\tau_{\theta\theta} &= \frac{2\alpha(x-1)}{3\lambda^2l} + O[(x-1)^2] \\
\tau_{x_2x_2} &= -\frac{\alpha(x-1)}{3\lambda^2l} + O[(x-1)^2] \\
ds_{\text{dual}}^2 &= \frac{l^2}{R^2} ds^2 = -d\tau^2 + d\theta^2 + dx_2^2 \\
\langle \tau_{ab}^{\text{dual}} \rangle &= \lim_{R \rightarrow \infty} \frac{R}{l} \tau_{ab} = \lim_{x \rightarrow x_b} \left[-\frac{1}{\lambda l(x-1)} \right] \tau_{ab} = \frac{1}{l^2} \left(\frac{\alpha}{3\lambda^3}\right) [-3\delta_a^\theta \delta_b^\theta + \gamma_{ab}] \\
\Box \\
M &= \oint_{\Sigma} d^2y \sqrt{\sigma} m^a \tau_{ab} \xi^b = \frac{LL_s f^{1/2} \Omega}{\sqrt{-g_{\tau\tau}}} (\partial_\tau)^i \tau_{ij} (\partial_\tau)^j = -\frac{LL_s}{l^2} \left[\frac{\alpha}{3\lambda^3} + O(x-1) \right] \Box, \\
ds^2 &= \sigma_{ij} dx^i dx^j = \Omega(x) \left[f(x) d\theta^2 + \frac{dx_2^2}{l^2} \right] \\
r_b^2 &= \frac{\Omega(x_h, \eta)}{l^2}, r_s^2 = \frac{\Omega(x_s, \lambda)}{l^2} \\
E &= M_{bh} - M_{\text{soliton}} = \frac{LL_b}{l^2} (2\mu_b + \mu_s) \\
\Delta F &= \beta_b^{-1} (I_{BH}^E - I_{\text{soliton}}^E) = \frac{TL\alpha}{3l^2} \left(\frac{L_s\beta_s}{\lambda^3} - \frac{L_b\beta_b}{\eta^3} \right) \\
\Delta F &= \frac{4\pi LL_s}{3l^2} \left[\frac{\Omega(\lambda, x_s)}{L_s} - T\Omega(\eta, x_h) \right] = \frac{4\pi L}{3l^2} \Omega(\lambda, x_s) \left(1 - \frac{r_b^3}{r_s^3} \right) \\
\frac{\mathcal{A}}{Tl^3} &= \frac{\alpha L}{4\pi l^5} \frac{\beta_b^2 L_s}{\eta^3} = \frac{L\mathcal{L}}{l} \left(\frac{\lambda}{\eta} \right) \\
\mathcal{L} &= \frac{16\pi^2}{\alpha^2 l^4} \left[\frac{9x_h^2}{(x_h^3 - 1)^2} \right]^3 \\
\frac{\mathcal{A}}{Tl^3} &= \frac{L\mathcal{L}(\alpha, l) r_b}{l r_s} \\
I &= \int d^{n+1}x \sqrt{-g} \left[\frac{R}{2\kappa} - \frac{(\partial\phi)^2}{2} - V(\phi) \right] + \frac{1}{\kappa} \int_{\partial\mathcal{M}} d^n x \sqrt{-h} K + I_g + I_\phi \\
G_{\mu\nu} &= \kappa T_{\mu\nu} \\
G_{\mu\nu} &= R_{\mu\nu} - \frac{1}{2} g_{\mu\nu} R, T_{\mu\nu} = \partial_\mu \phi \partial_\nu \phi - g_{\mu\nu} \left[\frac{(\partial\phi)^2}{2} + V \right] \\
G &= -\frac{R(n-1)}{2}, T = -(n-1) \left[\frac{(\partial\phi)^2}{2} + V \frac{(n+1)}{(n-1)} \right] \\
G_{\mu\nu} = \kappa T_{\mu\nu} &\rightarrow \frac{R}{2\kappa} = \frac{(\partial\phi)^2}{2} + V \frac{(n+1)}{(n-1)} \\
I_{\text{bulk}}^E &= -\frac{2}{n-1} \int d^{n+1}x \sqrt{g^E} V(\phi)
\end{aligned}$$

$$\begin{aligned}
ds^2 &= \Omega(x) \left[-f(x)dt^2 + \frac{\eta^2 dx^2}{f(x)} + d\Sigma_k^2 \right] \\
E_t^t - E_x^x &= 0 \Rightarrow 2\kappa\phi'^2 = \frac{D-2}{2\Omega^2} [3(\Omega')^2 - 2\Omega\Omega''] \\
E_t^t - \frac{1}{D-2} g^{ab} E_{ab} &= 0 \Rightarrow f'' + \frac{D-2}{2\Omega} \Omega' f' + 2k\eta^2 = 0 \\
E_t^t + \frac{1}{D-2} g^{ab} E_{ab} &= 0 \Rightarrow 2\kappa V = -\frac{D-2}{2\eta^2\Omega^2} \left[f\Omega'' + \frac{D-4}{2\Omega} f(\Omega')^2 + \Omega' f' \right] + \frac{k(D-2)}{\Omega} \\
&\frac{d}{dx} [\Omega^{(D-2)/2} f'] + 2\eta^2 k \Omega^{(D-2)/2} = 0 \\
-\frac{2\eta^2\Omega^{D/2}(2\kappa V)}{D-2} &= f\Omega''\Omega^{(D-4)/2} + \Omega'(f\Omega^{(D-4)/2})' - 2\eta^2 k \Omega^{(D-2)/2} \\
2\kappa V &= -\frac{D-2}{2\eta^2\Omega^{D/2}} \left[\Omega^{\frac{D-4}{2}} (f\Omega)' \right]' \\
d\Sigma_k^2 &= v_{ij} dx^i dx^j \\
I_{bulk}^E &= \frac{\beta\sigma_{k,n-1}}{2\kappa\eta} \left[\Omega^{\frac{D-4}{2}} (f\Omega)' \right]_{x_h}^{x_b} \\
h_{ab} dx^a dx^b &= \Omega(x) [-f(x)dt^2 + d\Sigma_k^2] \\
n_a &= \frac{\delta_a^x}{\sqrt{g^{xx}}}, K_{ab} = \frac{\sqrt{g^{xx}}}{2} \partial_x h_{ab}, K = \frac{1}{2\eta} \left(\frac{f}{\Omega} \right)^{1/2} \left[\frac{(\Omega f)'}{\Omega f} + (D-2) \frac{\Omega'}{\Omega} \right]. \\
I_{GH}^E &= -\frac{\beta\sigma_{k,n-1}}{2\kappa\eta} \Omega^{(D-2)/2} f \left[\frac{(f\Omega)'}{f\Omega} + (D-2) \frac{\Omega'}{\Omega} \right] \\
I_g &= -\frac{1}{\kappa} \int d^n x \sqrt{-h} \left[\frac{n-1}{l} + \frac{l\mathcal{R}}{2(n-2)} + \frac{l^3}{2(n-4)(n-2)^2} \left(\mathcal{R}_{ab}\mathcal{R}^{ab} - \frac{n\mathcal{R}^2}{4(n-1)} \right) \right] \\
\mathcal{R}_{ij} &= \frac{(n-2)k}{\Omega} \sigma_{ij}, \mathcal{R} = \frac{k(n-2)(n-1)}{\Omega}, \mathcal{R}_{ab}\mathcal{R}^{ab} = \frac{(n-2)^2(n-1)k^2}{\Omega^2} \\
\mathcal{R}_{ab}\mathcal{R}^{ab} - \frac{n\mathcal{R}^2}{4(n-1)} &= -\frac{k^2}{4\Omega^2} (n-2)^2(n-1)(n-4) \\
I_g^E &= \frac{\beta\sigma_{k,n-1}(D-2)}{\kappa l} \sqrt{\Omega^{D-1} f} \left(1 + \frac{l^2 k}{2\Omega} - \frac{l^4 k^2}{8\Omega^2} \right) \\
\beta^{-1} = T &= \frac{f'}{4\pi\eta} \Big|_{x_h}, S = \frac{\mathcal{A}}{4G} \\
I_{bulk}^E + I_{GH}^E + I_g^E &= -\frac{1}{T} \left(\frac{\mathcal{A}T}{4G} \right) \\
-\frac{\sigma_{D-2,k}}{2\kappa T} \Omega^{(D-2)/2} (D-2) &\left[\frac{f\Omega'}{\eta\Omega} - \frac{\sqrt{\Omega}f}{2l} \left(1 + \frac{l^2 k}{2\Omega} - \frac{l^4 k^2}{8\Omega^2} \right) \right] \Big|_{x_b} \\
ds^2 &= -N(r)dt^2 + H(r)dr^2 + S(r)d\Sigma_k^2 \\
\Omega(x) \rightarrow S(r), f(x) &\rightarrow \frac{N(r)}{S(r)}, \frac{\sqrt{NH}}{\eta S} dr \rightarrow dx \\
2\kappa V &= -\frac{D-2}{2\eta^2\Omega^{D/2}} \left[\Omega^{\frac{D-4}{2}} (f\Omega)' \right]' \rightarrow 2\kappa V = -\frac{D-2}{2S^{\frac{D-2}{2}}\sqrt{NH}} \frac{d}{dr} \left(\frac{S^{\frac{D-2}{2}}}{\sqrt{NH}} \frac{dN}{dr} \right) \\
I_{bulk}^E &= \frac{\beta\sigma_{k,n-1}}{2\kappa\eta} \left[\Omega^{\frac{D-4}{2}} (f\Omega)' \right]_{x_h}^{x_b} \rightarrow I_{bulk}^E = \frac{\beta\sigma_{k,n-1}}{2\kappa} \frac{dN}{dr} \frac{S^{(n-1)/2}}{\sqrt{NH}} \Big|_{r_h}^{r_b} \\
h_{ab} dx^a dx^b &= -N(R)dt^2 + S(R)d\Sigma_k^2 \\
n_\mu &= \frac{\delta_\mu^r}{\sqrt{g^{rr}}}, K_{\mu\nu} = \frac{\sqrt{g^{rr}}}{2} \partial_r h_{\mu\nu}, K = \frac{1}{2\sqrt{H}} \left[\frac{N'}{N} + (n-1) \frac{S'}{S} \right]
\end{aligned}$$



$$I_{GH}^E = -\frac{\beta\sigma_{k,n-1}}{2\kappa\eta}\Omega^{\frac{D-2}{2}}f\left[\frac{(f\Omega)'}{f\Omega} + (D-2)\frac{\Omega'}{\Omega}\right]\Big|_{x_b} \rightarrow I_{GH}^E = -\frac{\sigma_{k,n-1}S^{\frac{D-2}{2}}}{2\kappa T\sqrt{NH}}\left[\frac{dN}{dr} + (D-2)\frac{N}{S}\frac{dS}{dr}\right]\Big|_{r_b}$$

$$I_g^E = \frac{\beta\sigma_{k,n-1}(n-1)}{l\kappa}\Omega^{\frac{D-1}{2}}f^{\frac{1}{2}}\left(1 + \frac{l^2k}{2\Omega} - \frac{l^4k^2}{8\Omega^2}\right) \rightarrow I_g^E = \frac{\beta\sigma_{k,n-1}(n-1)}{l\kappa}S^{\frac{D-2}{2}}N^{\frac{1}{2}}\left(1 + \frac{l^2k}{2S} - \frac{l^4k^2}{8S^2}\right)_{r_b}$$

$$\beta^{-1} = T = \frac{N'}{4\pi\sqrt{NH}}\Big|_{r_h}, S = \frac{\mathcal{A}}{4G}$$

$$I_{bulk}^E + I_{GH}^E + I_g^E = -\frac{1}{T}\left(\frac{\mathcal{A}T}{4G}\right)$$

$$-\frac{\sigma_{D-2,k}}{2\kappa T}S^{(D-2)/2}(D-2)\left[\frac{NS'}{S\sqrt{NH}} - \frac{2\sqrt{N}}{l}\left(1 + \frac{l^2k}{2S} - \frac{l^4k^2}{8S^2}\right)\right]\Big|_{r_b}$$

$$2\kappa\phi'^2 = \frac{D-2}{2\Omega^2}[3(\Omega')^2 - 2\Omega\Omega'']; \quad \frac{2\kappa\phi'^2}{D-2} = \frac{1}{2S^2}[S'^2 - 2SS''] + \frac{S'(NH)'}{2S\sqrt{NH}}$$

$$\frac{d}{dx}\left[\Omega^{\frac{D-2}{2}}\frac{df}{dx}\right] = -2\eta^2k\Omega^{\frac{D-2}{2}}; \quad \frac{d}{dr}\left[\frac{S^{D/2}}{\sqrt{NH}}\frac{d}{dr}\left(\frac{N}{S}\right)\right] = -2k\sqrt{NHS}^{\frac{D-4}{2}}$$

$$2\kappa V = -\frac{D-2}{2\eta^2\Omega^{D/2}}\left[\Omega^{\frac{D-4}{2}}(f\Omega)'\right]'; \quad 2\kappa V = -\frac{D-2}{2S^{\frac{D-2}{2}}\sqrt{NH}}\frac{d}{dr}\left(\frac{S^{\frac{D-2}{2}}dN}{\sqrt{NH}dr}\right)$$

$$\partial_x\left[\Omega^{\frac{D-4}{2}}f\phi'\right] = \eta^2\Omega^{D/2}\frac{\partial V}{\partial\phi}; \quad \partial_r\left(S^{\frac{D-2}{2}}\phi'\sqrt{\frac{N}{H}}\right) = \sqrt{NHS}^{\frac{D-2}{2}}\frac{\partial V}{\partial\phi}$$

$$\tau^{ab} \equiv \frac{2}{\sqrt{-h}}\frac{\delta I}{\delta h_{ab}}$$

$$I_{GH} + I_g + I_\phi = \frac{1}{\kappa}\int d^n x\sqrt{-h}K - \frac{1}{\kappa}\int d^n x\sqrt{-h}\left[\frac{n-1}{l} + \frac{l\mathcal{R}}{2(n-2)}\right] - \int d^n x\sqrt{-h}\Psi$$

$$\tau_{ab} = -\frac{1}{\kappa}\left(K_{ab} - h_{ab}K + \frac{n-1}{l}h_{ab} - \frac{l}{n-2}G_{ab}\right) - h_{ab}[\Psi].$$

$$h_{ab}dx^a dx^b = -N(R)dt^2 + S(R)d\Sigma_k^2$$

$$G_{ab} = \mathcal{R}_{ab} - \frac{1}{2}\mathcal{R}h_{ab}, G_{tt} = \frac{(n-2)(n-1)kN}{2S}, G_{ij} = -\frac{(n-2)(n-3)}{2}kv_{ij}$$

$$\tau_{tt} = -\frac{(n-1)}{\kappa}\left[\frac{NS'}{2S\sqrt{H}} - \frac{N}{l}\left(1 + \frac{l^2k}{2S}\right)\right] + N[\Psi]$$

$$\tau_{ij} = \frac{v_{ij}}{\kappa}\left[\frac{S}{2\sqrt{H}}\left(\frac{N'}{N} + \frac{S'}{S}(n-2)\right) - \frac{(n-1)S}{l} - \frac{lk(n-3)}{2}\right] - v_{ij}S[\Psi]$$

$$h_{ab}dx^a dx^b = -L^2dt^2 + \sigma_{ij}(dy^i + L^i dt)(dy^j + L^j dt)$$

$$E = Q_{\frac{\partial}{\partial t}} = \oint_{\Sigma} d^{D-2}y\sqrt{\sigma}u^a\tau_{ab}\xi^b = \left(\oint_{\Sigma} d^2y\sqrt{v}\right)\frac{S^{\frac{D-2}{2}}\tau_{tt}}{\sqrt{N}} = \frac{\sigma_{k,n-1}S^{\frac{D-2}{2}}}{\sqrt{N}}\tau_{tt}$$

$$\sigma_{ij}dx^i dx^j = Sd\Sigma_k^2$$

$$\tau_{tt} = -\frac{(n-1)}{\kappa}\left[\frac{f^{3/2}\Omega'}{2\eta\sqrt{\Omega}} - \frac{\Omega f}{l}\left(1 + \frac{l^2k}{2\Omega}\right)\right] + \frac{\Omega f}{\kappa}[\Psi],$$

$$\tau_{ij} = \frac{v_{ij}}{\kappa}\left[\frac{(\Omega f)'}{2\eta\sqrt{\Omega f}} + \frac{(n-2)\Omega'\sqrt{f}}{2\eta\sqrt{\Omega}} - \frac{(n-1)\Omega}{l} - \frac{lk(n-3)}{2}\right] - \frac{v_{ij}\Omega}{\kappa}[\Psi].$$



$$E = Q \frac{\partial}{\partial t} = \oint_{\Sigma} d^{D-2} y \sqrt{\sigma} u^a \tau_{ab} \xi^b = \frac{\sigma_{k,n-1} \Omega^{\frac{D-2}{2}}}{\sqrt{\Omega} f} \tau_{tt}$$

$$\eta^2 V(\phi) = -\frac{f \Omega''}{\eta^2 \Omega^2} - \frac{f' \Omega'}{\eta^2 \Omega^2}$$

$$\eta^2 V(\phi) = -\frac{1}{x^{2\nu-2} \nu^4} \left[\frac{1}{l^2} + \frac{\alpha}{2} \left(\frac{x^{2+\nu} - 1}{2 + \nu} + \frac{x^{2-\nu} - 1}{2 + \nu} - x^2 + 1 \right) \right] \left[\frac{x^{\nu-1}(\nu-1)^2 \nu^2 - x^{\nu-1}(\nu-1) \nu^2}{x^2(x^\nu-1)^2 \eta^2} \right. \\ \left. - \frac{4x^{2\nu-1}(\nu-1)\nu^3 + 2x^{2\nu-1}\nu^4}{(x^\nu-1)^4 \eta^2 x^2} + \frac{2x^{2\nu-1}\nu^3 - 4x^{2\nu-1}(\nu-1)\nu^3}{x^2 \eta^2 (x^\nu-1)^3} \right] (x^\nu-1)^4 \eta^4 \\ - \frac{1}{2x^{2\nu-2} \nu^4} \left(\frac{x^{\nu-1}(\nu-1)\nu^2}{x \eta^2 (x^\nu-1)^2} - \frac{2x^{2\nu-1}\nu^3}{x \eta^2 (x^\nu-1)^3} \right) (x^{1+\nu} + x^{1-\nu} - 2x)$$

$$V(\phi) = -\frac{f(x)\nu^2}{x^{2\nu-2} \nu^4} \left(\frac{x^{3\nu-3}\nu^2 + 4x^{2\nu-3}\nu^2 + x^{\nu-3}\nu^2 + 3x^{3\nu-3}\nu - 3x^{\nu-3}\nu}{(x^\nu-1)^4 \eta^2} \right) (x^\nu-1)^4 \eta^2 \\ - \frac{f(x)\nu^2}{x^{2\nu-2} \nu^4} \left(\frac{2x^{3\nu-3} - 4x^{2\nu-3} + 2x^{\nu-3}}{(x^\nu-1)^4 \eta^2} \right) (x^\nu-1)^4 \eta^2$$

$$- \frac{\alpha \eta^4 (x^\nu-1)^4}{2x^{2\nu-2} \nu^4} \left(-\frac{\nu^2}{\eta^2 (x^\nu-1)^3} \right) (x^{2\nu-2}(\nu+1) + x^{\nu-2}(\nu-1))(x^{1+\nu} + x^{1-\nu} - 2x)$$

$$V(\phi) = -\frac{f(x)x^{-\nu}}{x\nu^2} (x^{2\nu}(\nu+1)(\nu+2) + 4x^\nu(\nu^2-1) + (\nu-1)(\nu-2))$$

$$+ \frac{\alpha(x^\nu-1)x^{2\nu-2}}{2\nu^2 x^{\nu-1} x^{\nu-1}} (x^{1+\nu} + x^{1-\nu} - 2x)(x^{-\nu}(\nu-1) + (\nu+1))$$

$$V(\phi) = -\frac{f(x)}{2x\nu^2} (2x^\nu(\nu+1)(\nu+2) + 8(\nu^2-1) + 2x^{-\nu}(\nu-1)(\nu-2))$$

$$+ \frac{\alpha}{2\nu^2 x^{\nu-1}} (x^\nu-1)(1+\nu+x^{-\nu}(\nu-1))(1+x^{2\nu}-2x^\nu)$$

$$V(\phi) = -\frac{f(x)}{2x\nu^2} (2x^\nu(\nu+1)(\nu+2) + 8(\nu^2-1) + 2x^{-\nu}(\nu-1)(\nu-2))$$

$$- \frac{\alpha}{2\nu^2 x^{\nu-1}} (x^\nu-1)^2 (2-x^\nu(\nu+1) + x^{-\nu}(\nu-1))$$



$$\begin{aligned}
V(\phi) &= -\frac{e^{-l_v\phi}}{2v^2} \left[\frac{1}{l^2} + \frac{\alpha}{2} \left(\frac{e^{(2+v)l_v\phi} - 1}{2+v} + \frac{e^{(2-v)l_v\phi} - 1}{2-v} - x^2 + 1 \right) \right] [8(v^2 - 1) \\
&\quad + 2(v+1)(v+2)e^{vl_v\phi} + \\
2(v-1)(v-2)e^{-vl_v\phi}] - \frac{\alpha e^{l_v\phi}}{2v^2} \left(\exp \frac{vl_v\phi}{2} - \exp -\frac{vl_v\phi}{2} \right)^2 [2 - e^{vl_v\phi}(v+1) + e^{-vl_v\phi}(v-1)] \\
V(\phi) &= -\frac{(v^2-4)}{l^2v^2} \left[\frac{v-1}{v+2} e^{-\phi l_v(v+1)} + \frac{v+1}{v-2} e^{\phi l_v(v-1)} + 4 \frac{v^2-1}{v^2-4} e^{-\phi l_v} \right] \\
&\quad - \frac{\alpha}{2v^2} \left[\frac{e^{(v+2)l_v\phi}}{2+v} - \frac{e^{(2-v)l_v\phi}}{v-2} - e^{2l_v\phi} + \frac{v^2}{v^2-4} \right] [4(v^2-1)e^{-l_v\phi} + (v+1)(v+2)e^{(v-1)l_v\phi} + \\
&\quad (v-1)(v-2)e^{-(v+1)l_v\phi}] \\
&\quad - \frac{\alpha}{2v^2} (e^{vl_v\phi} - 2 + e^{-vl_v\phi}) [2e^{l_v\phi} - e^{(v+1)l_v\phi}(v+1) + e^{(1-v)l_v\phi}(v-1)] \\
V(\phi) &= V_\Lambda(\phi) - \frac{\alpha}{2(v^2-4)} [v^2(e^{l_v\phi(v-1)} - e^{l_v\phi(v+1)} - e^{l_v\phi(v-1)} + e^{-l_v\phi(v+1)} + 4e^{-l_v\phi} - 4e^{l_v\phi}) + \\
&\quad 3v(e^{\phi l_v(v-1)} + e^{\phi l_v(v+1)} - e^{-\phi l_v(v-1)} - e^{-\phi l_v(v+1)}) \\
&\quad + 2(e^{\phi l_v(v-1)} - e^{\phi l_v(v+1)} - e^{-\phi l_v(v-1)} + e^{-\phi l_v(v+1)}) - 4e^{-l_v\phi} + 4e^{l_v\phi}] \\
V(\phi) &= V_\Lambda(\phi) - \frac{\alpha}{2(v^2-4)} [(2v^2+6v+4)\sinh \phi l_v(v-1) - (2v^2-6v+4)\sinh \phi l_v(v+1) \\
&\quad + 8(1-v^2)\sinh \phi l_v] \\
V(\phi) &= \frac{\Lambda(v^2-4)}{3v^2} \left[\frac{v-1}{v+2} e^{-\phi l_v(v+1)} + \frac{v+1}{v-2} e^{\phi l_v(v-1)} + 4 \frac{v^2-1}{v^2-4} e^{-\phi l_v} \right] \\
&\quad + \alpha \left[\frac{v-1}{v+2} \sinh \phi l_v(v+1) - \frac{v+1}{v-2} \sinh \phi l_v(v-1) + 4 \frac{v^2-1}{v^2-4} \sinh l_v \right] \\
dx^\mu \wedge dx^\nu &= dx^\mu \otimes dx^\nu - dx^\nu \otimes dx^\mu \\
dx^{\mu_1} \wedge \dots \wedge dx^{\mu_p} &= \sum_{\sigma} (-1)^{|\sigma|} dx^{\sigma(\mu_1)} \otimes \dots \otimes dx^{\sigma(\mu_p)} \\
H &= \frac{1}{p!} H_{\mu_1 \dots \mu_p} dx^{\mu_1} \wedge \dots \wedge dx^{\mu_p} \\
A &= A_\mu dx^\mu \\
F &= \frac{1}{2} F_{\mu\nu} dx^\mu \wedge dx^\nu \\
*(dx^{\mu_1} \wedge \dots \wedge dx^{\mu_p}) &= \frac{\sqrt{-g}}{(D-p)!} \epsilon^{\mu_1 \dots \mu_p}{}_{\nu_{p+1} \dots \nu_D} dx^{\nu_{p+1}} \wedge \dots \wedge dx^{\nu_D} \\
dV &= d^D x \sqrt{-g} = \frac{\sqrt{-g}}{D!} \epsilon_{\mu_1 \dots \mu_D} dx^{\mu_1} \wedge \dots \wedge dx^{\mu_D} \\
S &= -m_1 c \int ds_1 - m_2 c \int ds_2 - \frac{c^3}{16\pi G} \int d^4 x \sqrt{-g} R[g] \\
\Delta p_i^\mu &= p_i^\mu(\tau \rightarrow +\infty) - p_i^\mu(\tau \rightarrow -\infty) \\
&= G\Delta p_i^{(1)\mu} + G^2\Delta p_i^{(2)\mu} + \dots + G^5\Delta p_i^{(5)\mu} + \dots \\
\frac{d}{dx} \vec{I}(x, D) &= \hat{M}(x, D) \vec{I}(x, D) \\
\mathcal{J}(\phi_1, \phi_2, \dots, \phi_n; x) &:= \int_1^x dx' \phi_1(x') \mathcal{J}(\phi_2, \dots, \phi_n; x') \\
\left[\left(x \frac{d}{dx} - 1 \right)^4 - x^4 \left(x \frac{d}{dx} + 1 \right)^4 \right] \varpi(x) &= 0. \\
\theta^{(5)} &= \frac{M^5 \Gamma}{b^5 c^{10}} (\theta^{(5,0)} + v\theta^{(5,1)} + v^2\theta^{(5,2)} + v^3\Gamma^{-2}\theta^{(5,3)}) \\
P_{\text{rad}}^\mu &= -\Delta p_1^\mu - \Delta p_2^\mu
\end{aligned}$$

$$S_i = -\frac{m_i}{2} \int d\tau g_{\mu\nu} [x_i(\tau)] \dot{x}_i^\mu(\tau) \dot{x}_i^\nu(\tau)$$

$$\Delta p_1^{(5)\mu} = m_1 m_2 \left(m_2^4 \Delta p_{0SF}^{(5)\mu} + m_1 m_2^3 \Delta p_{1SF}^{(5)\mu} \right. \\ \left. + m_1^2 m_2^2 \Delta p_{2SF}^{(5)\mu} + m_1^3 m_2 \Delta p_{3SF}^{(5)\mu} + m_1^4 \Delta p_{4SF}^{(5)\mu} \right)$$

$$\Delta p_1 = \Gamma + \overset{G^2}{\text{diagram}} + \overset{G^3}{\text{diagram}} + \overset{G^4}{\text{diagram}} + \overset{G^5}{\text{diagram}} + \dots$$

$\begin{aligned} \ell^\mu &= \frac{1}{\ell \cdot v_j + i0^+} \\ \text{wavy line} &= \frac{1}{(\ell^0 + i0^+)^2 - \ell^2} \\ \dots &= \delta(\ell \cdot v_j) \end{aligned}$

~mG^{1/2}

~G^{1/2}

~G

~G^{3/2}

~G^2

$= \int d^D \ell \frac{\delta(\ell \cdot v_2)}{\ell^2 (q - \ell)^2 \ell \cdot v_1}$

$$\mathcal{J}_{\{n\}}^{\{\sigma\}} = \int_{\ell_1 \dots \ell_4} \frac{\delta^{(n_1-1)}(\ell_1 \cdot v_1) \prod_{i=2}^L \delta^{(n_i-1)}(\ell_i \cdot v_2)}{\prod_{i=1}^4 D_i^{n_i}(\sigma_i) \prod_{I < J} D_{IJ}^{n_{IJ}}}$$

$$D_1 = \ell_1 \cdot v_2 + \sigma_1 i0^+, D_i = \ell_i \cdot v_1 + \sigma_i i0^+$$

$$D_{1j} = (\ell_1 - \ell_j)^2 + \sigma_{4+j} \text{sign}(\ell_1^0 - \ell_j^0) i0^+, D_{q1} = (\ell_1 + q)^2$$

$$\frac{d}{dx} \vec{J}(x, \epsilon) = \epsilon \hat{A}(x) \vec{J}(x, \epsilon),$$

$$\vec{J} = \mathcal{P} e^{\int_1^x dx \hat{A}(x)} \vec{J}$$

$$[\hat{\theta}^3 - 2x^2(2 + 4\hat{\theta} + 3\hat{\theta}^2 + \hat{\theta}^3) + x^4(2 + \hat{\theta})^3] I_1|_{\epsilon=0} = 0$$

$$(d - A^u(x)) W^u(x) = 0$$

$$\alpha_1 = \frac{\varpi_0^2}{x(\varpi_0 \varpi_1' - \varpi_0' \varpi_1)}$$

$$\frac{d}{dx} \vec{I}_{CY3} = \sum_{i=-2}^1 \epsilon^i \tilde{M}_{CY3}^i(x) \vec{I}_{CY3}$$

$$\vec{I}_{CY3} = T_{\epsilon\text{-scalings}} (W^{SS})^{-1} \vec{I}_{CY3}$$

$$G_1'(x) = -\frac{96x(x^4 + 1)\varpi_0(x)^2}{(x-1)^2(x+1)^2(x^2+1)^2\alpha_1(x)}$$

$$I_i^{\text{canonicalization}} = f(\epsilon) g(x) \sum_k c_k \tilde{\mathcal{S}}_k$$

$$G_8'(x) = \frac{\varpi_{K3}(x) G_3(x) \varpi_0(x) \alpha_1'(x)}{\alpha_1(x)^2}$$

$$\ell_i^P = (\ell_i^0, \ell_i) \sim (v, 1), \ell_i^R = (\ell_i^0, \ell_i) \sim (v, v)$$

$$\Delta p_{1SF}^{(5)\mu} = \frac{1}{b^5} (\hat{b}^\mu c_b(\gamma) + \check{v}_2^\mu c_v(\gamma) + \check{v}_1^\mu c_v'(\gamma))$$

$$\Delta p_{1SF}^{(5)\mu} = \frac{1}{b^5} (\hat{b}^\mu \bar{c}_b(\gamma) + \check{v}_2^\mu \bar{c}_v(\gamma) + \check{v}_1^\mu \bar{c}_v'(\gamma))$$

$$c_w(\gamma) = c_{w, \text{even}}(\gamma) + c_{w, \text{odd}}(\gamma),$$

$$\bar{c}_w(\gamma) = \bar{c}_{w, \text{even}}(\gamma) + \bar{c}_{w, \text{odd}}(\gamma),$$

$$c_{w,z}(\gamma) = \sum_\alpha d_{w,z}^{(\alpha)}(\gamma) F_{w,z}^{(\alpha)}(\gamma)$$

$$+ \sum_\alpha d_{w,z}^{(\alpha, \text{tail})}(\gamma) F_{w,z}^{(\alpha, \text{tail})}(\gamma) \log(\gamma - 1)$$

$$P_{\text{rad}}^{(5)\mu} = \frac{M^6 v^2}{b^5} \left([r_1(\gamma) \hat{b}^\mu + r_2(\gamma)(v_1^\mu - v_2^\mu)] \frac{m_1 - m_2}{M} + [v_1^\mu + v_2^\mu](r_3(\gamma) + v r_4(\gamma)) \right).$$

$$\mathbf{p}_{\text{out}} = \mathbf{p}_{\text{in}} + \Delta \mathbf{p}_1 + \frac{E_1}{E} \mathbf{P}_{\text{recoil}} + \mathcal{O}(G^6)$$

$$\theta = \Gamma \sum_{n=1}^5 \left(\frac{GM}{b} \right)^n \sum_{m=0}^{\lfloor \frac{n-1}{2} \rfloor} v^m \theta^{(n,m)}(\gamma) + \frac{1}{\Gamma} \sum_{n=4}^5 \left(\frac{GM}{b} \right)^n v^{n-2} \theta^{(n,n-2)}(\gamma) + \mathcal{O}(G^6).$$

$$E_{\text{rad}}^{(5)} = \frac{M^6 v^2}{\Gamma b^5} [(1-\gamma)r_2(\gamma) + (1+\gamma)r_3(\gamma) + \mathcal{O}(v)] = \frac{M^6 v^2 \pi}{5\Gamma b^5 v^3} \left[122 + \frac{3583}{56} v^2 + \frac{297\pi^2}{4} v^3 - \frac{71471}{504} v^4 + \left(\frac{9216}{7} - \frac{24993\pi^2}{224} \right) v^5 + \left(\frac{2904562807}{6899200} + \frac{99\pi^2}{2} - \frac{10593}{70} \log \frac{v}{2} \right) v^6 + \left(\frac{7296}{7} - \frac{2927\pi^2}{28} \right) v^7 + \left(\frac{4924457539}{29429400} + \frac{8301\pi^2}{112} - \frac{491013}{3920} \log \frac{v}{2} \right) v^8 + \left(\frac{99524416}{40425} - \frac{46290891\pi^2}{157696} \right) v^9 + \mathcal{O}(v^{10}, v) \right].$$

$$G(a_1, \dots, a_n; y) = \int_0^y \frac{dy'}{y' - a_1} G(a_2, \dots, a_n; y')$$

$$J\left(\frac{1+x^2}{x} \varpi_{0, \kappa_3}(x); x\right)$$

$$J\left(\frac{1}{(1-x^2)x\varpi_{0, \kappa_3}(x)}, \frac{1+x^2}{x} \varpi_{0, \kappa_3}(x); x\right)$$

$$J\left(\frac{1}{(1-x^2)x\varpi_{0, \kappa_3}(x)}, \frac{1}{(1-x^2)x\varpi_{0, \kappa_3}(x)}, \frac{1+x^2}{x} \varpi_{0, \kappa_3}(x); x\right)$$

$$\theta^{(5,1)} = \frac{4}{5v^8} - \frac{137}{5v^6} + \frac{3008}{45v^5} + \frac{41\pi^2 - 3427}{4v^4} + \frac{84\pi^2 - 4096}{5v^3} + \frac{12544\log(2v)}{45} + \frac{3593\pi^2}{72} - \frac{445867}{432} + \frac{2144536}{11025} + \frac{453\pi^2}{35} + \left(-\frac{7552\log(2v)}{1575} + \frac{246527\pi^2}{1440} - \frac{1111790903}{756000} \right) + \left(\frac{19424344}{363825} + \frac{1787\pi^2}{672} \right) v + \left(-\frac{1762784\log(2v)}{11025} - \frac{184881\pi^2}{2240} + \frac{56424801733}{49392000} \right) v^2 + \left(\frac{16004496043}{104053950} - \frac{835619\pi^2}{59136} \right) v^3 + \mathcal{O}(v^4),$$

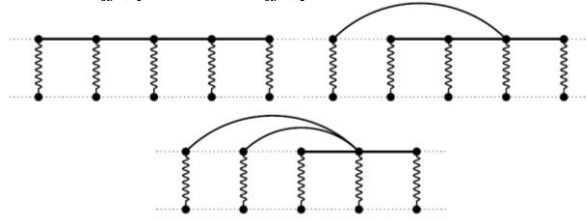
$$\theta^{(5,3)} = \frac{128}{3v} + \left(\frac{14528}{175} + \frac{37\pi^2}{10} \right) v - \frac{22016v^2}{225} + \left(\frac{3262832}{33075} + \frac{893\pi^2}{112} \right) v^3 + \frac{2877184v^4}{7875} + \left(\frac{15803\pi^2}{960} - \frac{4464536}{14553} \right) v^5 + \frac{970766528v^6}{1819125} + \left(\frac{11234077\pi^2}{394240} - \frac{258810752887}{780404625} \right) v^7 + \frac{83694772064v^8}{70945875} + \left(\frac{119425757\pi^2}{2795520} - \frac{5667010769993}{5533778250} \right) v^9 + \frac{11015320038116v^{10}}{5462832375} + \mathcal{O}(v^{11})$$



$$\begin{aligned}
r_1(\gamma) &= c_b(\gamma) - \bar{c}_b(\gamma) \\
r_2(\gamma) &= \frac{c_v(\gamma) - \bar{c}_v(\gamma) - c'_v(\gamma) + \bar{c}'_v(\gamma)}{2(\gamma - 1)} \\
r_3(\gamma) &= -\frac{c_v(\gamma) + \bar{c}_v(\gamma) + c'_v(\gamma) + \bar{c}'_v(\gamma)}{2(\gamma + 1)} \\
(\Gamma &= \sqrt{1 + 2v(\gamma - 1)}, v = \mu/M = m_1 m_2 / M^2, M = m_1 + m_2) \\
E_{\text{rad}}^{(5)} &= \frac{M^6 v^2}{\Gamma b^5} [(1 - \gamma)r_2(\gamma) + (1 + \gamma)r_3(\gamma) + \mathcal{O}(v)] \\
\mathbf{P}_{\text{recoil}}^{(5)} &= \frac{M^5 v^2 (m_1 - m_2)}{b^5} \left(r_1(\gamma) \hat{\mathbf{b}} + (r_2(\gamma) - r_3(\gamma)) \frac{\mathbf{p}_{\text{in}}}{\mu} + \mathcal{O}(v) \right) \\
r_1(\gamma) &= -\frac{64}{3v^2} - \frac{16192}{525} - \frac{37\pi^2}{20} - \frac{30208v}{225} + \left(-\frac{856768}{33075} - \frac{3429\pi^2}{1120} \right) v^2 - \frac{22016v^3}{2625} + \left(-\frac{1117888}{4851} - \frac{80723\pi^2}{13440} \right) v^4 \\
&\quad - \frac{123897344v^5}{606375} + \left(-\frac{36746586176}{780404625} - \frac{22515319\pi^2}{2365440} \right) v^6 - \frac{16343148032v^7}{70945875} + \left(-\frac{169791059264}{869593725} - \frac{33021283\pi^2}{2562560} \right) v^8 \\
&\quad - \frac{584895938048v^9}{1820944125} + \left(-\frac{53344474395584}{517408266375} - \frac{7894087273\pi^2}{492011520} \right) v^{10} + \mathcal{O}(v^{11}), \\
\frac{r_2(\gamma) - r_3(\gamma)}{\pi} &= -\frac{53}{3v^3} + \frac{72997}{5040v} + \frac{1491}{400} - \frac{1509\pi^2}{140} + \frac{20211v}{640} + \left(\frac{75661\pi^2}{4480} - \frac{2678867}{16800} \right) v^2 \\
&\quad + \left(\frac{41053 \log\left(\frac{v}{2}\right)}{2450} - \frac{503\pi^2}{70} - \frac{123069432361}{4346496000} \right) v^3 + \left(\frac{6139957\pi^2}{394240} - \frac{15259259693}{124185600} \right) v^4 \\
&\quad + \left(\frac{223443793 \log\left(\frac{v}{2}\right)}{15523200} - \frac{69203\pi^2}{6720} + \frac{64307545227137}{22375761408000} \right) v^5 + \left(\frac{252585041\pi^2}{5857280} - \frac{16101198460801}{56504448000} \right) v^6 \\
&\quad + \left(\frac{313945836331 \log\left(\frac{v}{2}\right)}{8879270400} - \frac{346561\pi^2}{16896} - \frac{204989896406483131}{1828419360768000} \right) v^7 + \left(\frac{918930349\pi^2}{11714560} - \frac{2985881877364537}{8524099584000} \right) v^8 \\
&\quad + \left(\frac{16292858440697 \log\left(\frac{v}{2}\right)}{307814707200} - \frac{288050119\pi^2}{8785920} - \frac{431015900890630345739}{1676186519175168000} \right) v^9 \\
&\quad + \left(\frac{326238718783\pi^2}{2788065280} - \frac{98473341190358572667}{217581903931392000} \right) v^{10} + \mathcal{O}(v^{11}). \\
\mathcal{S}_{\text{bulk}} &= -\frac{1}{16\pi G} \int d^4x \left[\sqrt{-g} g^{\alpha\beta} g^{\gamma\delta} g^{\mu\nu} (\Gamma_{\alpha\gamma\delta} \Gamma_{\mu\beta\nu} - \Gamma_{\alpha\gamma\mu} \Gamma_{\delta\beta\nu}) + \frac{1}{2} G^\mu G^\nu \eta_{\mu\nu} \right] \\
G^\mu &= \partial_{v_1} h^{\mu\nu_1} - \frac{1}{2} \partial^\mu h + \frac{1}{2} \partial_{v_1} h h^{\mu\nu_1} + \frac{1}{2} \partial^\mu h_{v_1\nu_2} h^{\nu_1\nu_2} - \partial^{\nu_1} h^{\mu\nu_2} h_{\nu_1\nu_2} - \frac{1}{4} \partial^\mu h h \\
&\quad + \partial_{v_1} h^{\nu_3\mu} h^{\nu_1\nu_2} h_{\nu_2\nu_3} + \frac{1}{2} \partial_{v_1} h^{\nu_1\nu_2} h_{\nu_2\nu_3} h^{\nu_3\mu} - \frac{1}{4} \partial_{v_2} h h_{\nu_2\nu_3} h^{\nu_3\mu} + \partial_{v_1} h_{\nu_2\nu_3} h^{\nu_1\nu_2} h^{\nu_3\mu} \\
&\quad - \frac{1}{2} \partial^{\nu_2} h_{\nu_1\nu_2} h^{\mu\nu_1} h + \frac{1}{4} \partial^\mu h_{\nu_1\nu_2} h^{\nu_1\nu_2} h - \frac{1}{2} \partial^{\nu_1} h^{\mu\nu_2} h_{\nu_1\nu_2} h + \frac{1}{8} \partial_{v_1} h^{\mu\nu_1} h^2. \\
\frac{d}{dx} \vec{J}(x; \epsilon) &= \epsilon \hat{A}(x) \vec{J}(x; \epsilon) \\
\varpi_0^{[0]} &= x + \frac{x^5}{16} + \frac{81x^9}{4096} + \dots, \\
\varpi_1^{[0]} &= \log(x) \left(x + \frac{x^5}{16} + \frac{81x^9}{4096} + \dots \right) + \frac{x^5}{16} + \frac{189x^9}{8192} + \dots, \\
\alpha_1^{[0]} &= 1 - \frac{x^4}{4} - \frac{93x^8}{1024} + \dots.
\end{aligned}$$

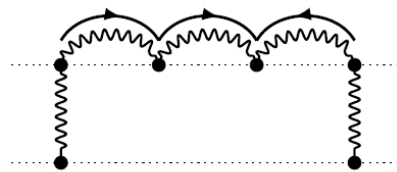
$$\begin{aligned} \varpi_0^{[1]} &= 1 - (1-x) + \frac{(1-x)^3}{3} + \frac{(1-x)^4}{3} + \dots, \\ \varpi_1^{[1]} &= 1 - \frac{(1-x)^2}{2} - \frac{(1-x)^3}{6} + \frac{(1-x)^4}{3} + \dots, \\ \alpha_1^{[1]} &= 1 - 2 \cdot (1-x) + 2 \cdot (1-x)^2 - \frac{(1-x)^3}{3} + \frac{(1-x)^4}{2} + \dots. \\ \varpi_{0,K3}^{[0]} &= \frac{4}{\pi^2} K^2(x^2) = 1 + \frac{x^2}{2} + \frac{11x^4}{32} + \dots, \\ \varpi_{0,K3}^{[1]} &= \frac{4}{\pi^2} K^2(1-x^2) = 1 - (1-x) + \frac{7(1-x)^2}{8} - \frac{3(1-x)^3}{4} + \dots \\ \hat{A} &= \sum_{j=0}^2 \hat{A}_{-1,j} \frac{\log^j(1-x)}{1-x} + \sum_{i,j \geq 0} \hat{A}_{i,j} (1-x)^i \log^j(1-x) \\ \vec{J}(x; \epsilon) &= \mathcal{P} e^{\epsilon \int_1^x dx \hat{A}(x) \vec{J}} \\ \vec{J}(x; \epsilon) &= \sum_{k,m,n} \vec{f}_{k,m,n}(\epsilon) (1-x)^{k+m\epsilon} \log^n(1-x) \\ \frac{d\hat{S}_0(x; \epsilon)}{dx} &= \epsilon \sum_{j=0}^2 \hat{A}_{-1,j} \frac{\log^j(1-x)}{1-x} \hat{S}_0(x; \epsilon) \\ \vec{J} &= \sum_{i,j}^N (\hat{F}_{i,j}(\epsilon) (1-x)^i \log^j(1-x)) \cdot \hat{S}_0(x; \epsilon) \cdot \vec{J} + \mathcal{O}((1-x)^{N+1}) \end{aligned}$$

$$\lim_{x \rightarrow 1} \vec{I}(x; \epsilon) = \lim_{x \rightarrow 1} \hat{T}(x; \epsilon)^{-1} \vec{J}(x; \epsilon)$$

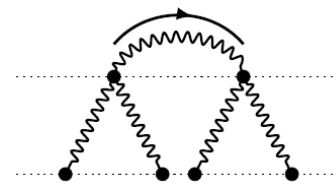


$$-i\delta(x) = \frac{1}{x + i0^+} - \frac{1}{x - i0^+}$$

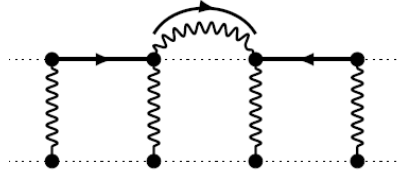
$$M[\dots, \omega_x, \dots] = M[\dots, \omega_+, \dots] - M[\dots, \omega_-, \dots]$$



$$\begin{aligned} M[g_+ g_+ g_-] &= \\ &= -\frac{1}{12288\pi^4} + \frac{12\gamma_E - 41 - 6\log(2) - 12\log(\pi)}{36864\pi^4} \epsilon + \mathcal{O}(\epsilon^2), \end{aligned}$$



$$\begin{aligned} M[0, g_+, 0] &= \\ &= -\frac{1}{65536\pi^2} - \frac{4\gamma_E + 3 + 26\log(2) + 4\log(\pi)}{65536\pi^2} \epsilon + \mathcal{O}(\epsilon^2), \end{aligned}$$

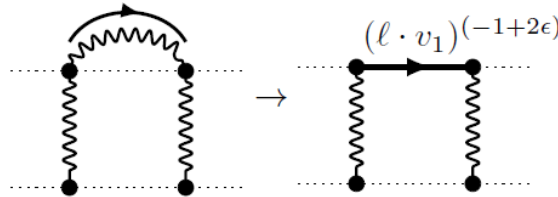


$$M[\omega_+, g_+, \omega_-] = - \frac{1}{8192\pi^4\epsilon^2} - \frac{2\gamma_E - 3 - 3\log(2) - 2\log(\pi)}{4096\pi^4\epsilon} + \mathcal{O}(\epsilon^0)$$

$$\int_{\ell} = \int \frac{d^D \ell}{(2\pi)^D} \frac{\delta(\ell \cdot v_1)}{(\ell^2 + i0)^{\nu_1} ((\ell - q)^2 + i0)^{\nu_2} (\ell \cdot v_2 + i0)^{\nu_3}} = (4\pi)^{\frac{1-D}{2}} (-1)^{\nu_1 + \nu_2 + \nu_3} i^{\nu_3} |q|^{D-1-2\nu_1-2\nu_2-\nu_3} \times \frac{\Gamma(\nu_1 + \nu_2 + \frac{\nu_3}{2} - \frac{D-1}{2}) \Gamma(\frac{\nu_3}{2}) \Gamma(\frac{D-1}{2} - \nu_1 - \frac{\nu_3}{2}) \Gamma(\frac{D-1}{2} - \nu_2 - \frac{\nu_3}{2})}{2\Gamma(\nu_1)\Gamma(\nu_2)\Gamma(\nu_3)\Gamma(D-1-\nu_1-\nu_2-\nu_3)}$$

$$\int_{\ell} \frac{\delta(\ell \cdot v_1)}{(\ell^2 + i0)^{\nu_1} ((\ell - q)^2 + i0)^{\nu_2}} = (4\pi)^{\frac{1-D}{2}} (-1)^{\nu_1 + \nu_2} |q|^{D-1-2\nu_1-2\nu_2} \frac{\Gamma(\nu_1 + \nu_2 - \frac{D-1}{2}) \Gamma(\frac{D-1}{2} - \nu_1) \Gamma(\frac{D-1}{2} - \nu_2)}{\Gamma(\nu_1)\Gamma(\nu_2)\Gamma(D-1-\nu_1-\nu_2)}$$

$$M[\omega_+] = \text{Diagram} = -i(4\pi)^{\epsilon-1} |q|^{-2(\epsilon+1)} \frac{\Gamma(-\epsilon)^2 \Gamma(\epsilon+1)}{2\Gamma(-2\epsilon)}$$



$$\int_{\ell} \frac{\delta(\ell \cdot v)}{(\ell - k)^2 + i0^+ (k - \ell) \cdot v} = -(+i)^{1-2\epsilon} (k \cdot v + i0)^{1-2\epsilon} (4\pi)^{\frac{2\epsilon-3}{2}} \Gamma\left(\frac{2\epsilon-1}{2}\right)$$

$$\frac{1}{(\omega + i0^+)^{\alpha}} = e^{-i\frac{\pi}{2}\alpha} \frac{1}{\Gamma(\alpha)} \int_0^{\infty} du u^{\alpha-1} e^{-u(0-i\omega)}$$

$$\frac{1}{(\omega - i0^+)^{\alpha}} = e^{i\frac{\pi}{2}\alpha} \frac{1}{\Gamma(\alpha)} \int_0^{\infty} du u^{\alpha-1} e^{-u(0+i\omega)}$$

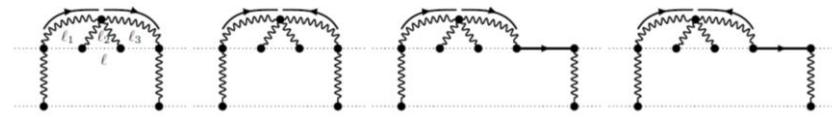
$$\frac{1}{(\ell^2)^{\alpha}} = \frac{1}{\Gamma(\alpha)} \int_0^{\infty} du u^{\alpha-1} e^{-u\ell^2}$$

$$\frac{1}{(\omega + i0^+)^{\alpha} (\omega - i0^+)^{\beta}} = e^{-i\pi\alpha} \frac{\sin[\pi\beta]}{\sin[\pi(\alpha + \beta)]} \frac{1}{(\omega - i0^+)^{\alpha+\beta}} + e^{i\pi\beta} \frac{\sin[\pi\alpha]}{\sin[\pi(\alpha + \beta)]} \frac{1}{(\omega + i0^+)^{\alpha+\beta}}$$

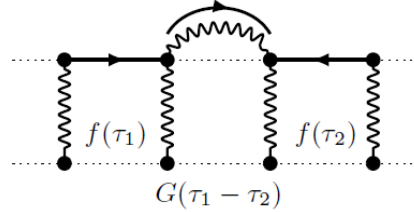
$$\text{Diagram} = \frac{\sin(\pi\alpha_+)}{\sin(\pi[\alpha_+ + \alpha_-])} \text{Diagram}$$

$$+ \frac{\sin(\pi\alpha_-)}{\sin(\pi[\alpha_+ + \alpha_-])}$$

$$\alpha_{\pm} = \sum_{i|s_i=\pm i0^+} \nu_i + \sum_{a|t_a=\pm i0^+} (2n_a - 3 + 2\epsilon)$$



$$\begin{aligned}
 \text{Diagram} &= f(\epsilon)(-i(\omega + i0^+))^{1-6\epsilon}, \\
 &= f(\epsilon)(\omega + i0^+)^{\frac{1-6\epsilon}{2}}(\omega - i0^+)^{\frac{1-6\epsilon}{2}}. \\
 f(\epsilon) &= \int_{\ell_1, \ell_2, \ell_3} \frac{1}{(\ell_1 - \ell_2)^2(\ell_2 - \ell_3)^2(\ell_1^2 + 1)(\ell_3^2 + 1)},
 \end{aligned}$$



$$\begin{aligned}
 A &= M[\omega_+, g_+, \omega_-](|b|, \epsilon) = \int d\tau_1 d\tau_2 f(\tau_1)G(\tau_1 - \tau_2)f(\tau_2) \\
 M[\omega_+, \omega_+, g_+] &= \frac{1}{2}M_1[\omega_+, \omega_+, g_+] \\
 M[\omega_-, \omega_-, g_+] &= \frac{1}{2}M_1[\omega_-, \omega_-, g_+] \\
 M[\omega_+, \omega_+, g_+](|b|, \epsilon) &= \frac{1}{2} \int d\tau_1 d\tau_2 f(\tau_1)f(\tau_1)G(\tau_1 - \tau_2) \\
 M[\omega_-, \omega_-, g_+](|b|, \epsilon) &= \frac{1}{2} \int d\tau_1 d\tau_2 f(-\tau_1)f(-\tau_1)G(\tau_1 - \tau_2) \\
 &= \frac{1}{2} \int d\tau_1 d\tau_2 f(\tau_2)f(\tau_2)G(\tau_1 - \tau_2) \\
 B &= M[\omega_+, \omega_+, g_+](|b|, \epsilon) + M[\omega_-, \omega_-, g_+](|b|, \epsilon) \\
 &= \frac{1}{2} \int d\tau_1 d\tau_2 (f^2(\tau_1) + f^2(\tau_2))G(\tau_1 - \tau_2).
 \end{aligned}$$

$$f(\tau) = \frac{c_0}{\epsilon} + c_1(\tau) + \mathcal{O}(\epsilon)$$

$$\begin{aligned}
 A &= \int d\tau_1 d\tau_2 \left(\frac{c_0^2}{\epsilon^2} + \frac{c_0}{\epsilon}(c_1(\tau_1) + c_1(\tau_2)) + \mathcal{O}(\epsilon) \right) G(\tau_1 - \tau_2) \\
 &= \int d\tau_1 d\tau_2 \left(\frac{c_0^2}{\epsilon^2} + 2\frac{c_0}{\epsilon}c_1(\tau_1) \right) G(\tau_1 - \tau_2) + \mathcal{O}(\epsilon) \\
 B &= \frac{1}{2} \int d\tau_1 d\tau_2 \left(\frac{c_0^2}{\epsilon^2} + 2\frac{c_0}{\epsilon}c_1(\tau_1) + \frac{c_0^2}{\epsilon^2} + 2\frac{c_0}{\epsilon}c_1(\tau_2) + \mathcal{O}(\epsilon) \right) G(\tau_1 - \tau_2) \\
 &= \int d\tau_1 d\tau_2 \left(\frac{c_0^2}{\epsilon^2} + 2\frac{c_0}{\epsilon}c_1(\tau_1) \right) G(\tau_1 - \tau_2) + \mathcal{O}(\epsilon)
 \end{aligned}$$

$$\begin{aligned}
 M[\omega_+, g_+, \omega_-] &= -M[\omega_+, \omega_+, g_+] - M[\omega_-, \omega_-, g_+] + \mathcal{O}(\epsilon^0) \\
 M[\omega_+, g_+, \omega_+, \omega_-] &= -M[\omega_+, \omega_+, g_+, \omega_+] - M[\omega_-, \omega_-, g_+, \omega_+] + \mathcal{O}(\epsilon^0)
 \end{aligned}$$

$$M[0, g_+, 0] = -\frac{1}{65536\pi^2} - \frac{(3 + 26\log(2))\epsilon}{65536\pi^2} + \mathcal{O}(\epsilon^2)$$

$$M[g_+, 0, 0] = -\frac{1}{8192\pi^4} - \frac{(17 + 6\log(2))\epsilon}{8192\pi^4} + \mathcal{O}(\epsilon^2)$$

$$M[0, g_+, \omega_+, \omega_-] = -\frac{1}{4096\pi^4\epsilon^2} - \frac{6 + 3\log(2)}{2048\pi^4\epsilon} + \mathcal{O}(\epsilon^0)$$

$$\begin{aligned}
M[0, g_+ \omega_+, \omega_+] &= \frac{1}{4096\pi^4 \epsilon^2} + \frac{6 + 3\log(2)}{2048\pi^4 \epsilon} + \mathcal{O}(\epsilon^0) \\
M[g_+, \omega_-, \omega_-] &= -\frac{1}{16384\pi^4 \epsilon^2} - \frac{13 + 9\log(2)}{8192\pi^4 \epsilon} + \mathcal{O}(\epsilon^0) \\
M[g_+, \omega_-, \omega_+] &= \frac{1}{16384\pi^4 \epsilon^2} + \frac{19 + 15\log(2)}{8192\pi^4 \epsilon} + \mathcal{O}(\epsilon^0) \\
M[g_+, \omega_+, \omega_-] &= \frac{1}{16384\pi^4 \epsilon^2} + \frac{-1 + 3\log(2)}{8192\pi^4 \epsilon} + \mathcal{O}(\epsilon^0) \\
M[g_+, \omega_+, \omega_+] &= \frac{1}{16384\pi^4 \epsilon^2} + \frac{7 + 3\log(2)}{8192\pi^4 \epsilon} + \mathcal{O}(\epsilon^0) \\
M[\omega_+, g_+, \omega_-] &= \frac{1}{8192\pi^4 \epsilon^2} + \frac{3 + 3\log(2)}{4096\pi^4 \epsilon} + \mathcal{O}(\epsilon^0) \\
M[\omega_+, g_+, \omega_+] &= \frac{1}{8192\pi^4 \epsilon^2} + \frac{7 + 3\log(2)}{4096\pi^4 \epsilon} + \mathcal{O}(\epsilon^0) \\
M[g_+ \omega_+, 0, \omega_-] &= -\frac{1}{4096\pi^4 \epsilon^2} - \frac{6 + 3\log(2)}{2048\pi^4 \epsilon} + \mathcal{O}(\epsilon^0) \\
M[g_+ \omega_+, 0, \omega_+] &= \frac{1}{4096\pi^4 \epsilon^2} + \frac{6 + 3\log(2)}{2048\pi^4 \epsilon} + \mathcal{O}(\epsilon^0) \\
M[g_+ \omega_+, \omega_-, 0] &= \frac{1}{8192\pi^2 \epsilon} + \mathcal{O}(\epsilon^0) \\
M[g_+ \omega_+, \omega_+, 0] &= -\frac{1}{8192\pi^2 \epsilon} + \mathcal{O}(\epsilon^0) \\
M[0, g_+, \omega_-] &= \frac{i}{768\pi^3} + \frac{55i\epsilon}{2304\pi^3} + \mathcal{O}(\epsilon^2) \\
M[0, g_+, \omega_+] &= \frac{i(-1 + 4\log(2))\epsilon}{768\pi^3} + \mathcal{O}(\epsilon^2) \\
M[0, g_+ \omega_+, 0] &= \frac{i}{1024\pi^3} + \frac{i(4 + 2\log(2))\epsilon}{256\pi^3} + \mathcal{O}(\epsilon^2) \\
M[g_+, 0, \omega_-] &= \frac{i}{1024\pi^3} + \frac{i(25 + 2\log(2))\epsilon}{1536\pi^3} + \mathcal{O}(\epsilon^2) \\
M[g_+, 0, \omega_+] &= -\frac{i}{3072\pi^3} + \frac{i(-41 + 30\log(2))\epsilon}{4608\pi^3} + \mathcal{O}(\epsilon^2) \\
M[g_+, \omega_-, 0] &= \frac{i}{2048\pi^3} + \frac{i(2 + 3\log(2))\epsilon}{512\pi^3} + \mathcal{O}(\epsilon^2) \\
M[g_+, \omega_+, 0] &= \frac{i}{2048\pi^3} + \frac{i(4 + 3\log(2))\epsilon}{512\pi^3} + \mathcal{O}(\epsilon^2) \\
M[g_+ \omega_+, 0, 0] &= \frac{i}{1024\pi^3} + \frac{i(7 + 6\log(2))\epsilon}{512\pi^3} + \frac{i(272 + 9\pi^2 + 48\log(2)(7 + 3\log(2)))\epsilon^2}{2048\pi^3} + \mathcal{O}(\epsilon^3) \\
M[\omega_+, g_+ \omega_+, \omega_-] &= -\frac{i}{2048\pi^3 \epsilon^2} - \frac{i(1 + 2\log(2))}{1024\pi^3 \epsilon} + \mathcal{O}(\epsilon^0) \\
M[\omega_+, g_+ \omega_+, \omega_+] &= \frac{i}{2048\pi^3 \epsilon^2} + \frac{i(1 + 2\log(2))}{1024\pi^3 \epsilon} + \mathcal{O}(\epsilon^0) \\
M[g_+ \omega_+, \omega_-, \omega_-] &= \frac{i}{4096\pi^3 \epsilon^2} + \frac{i(1 + 2\log(2))}{2048\pi^3 \epsilon} + \mathcal{O}(\epsilon^0) \\
M[g_+ \omega_+, \omega_-, \omega_+] &= -\frac{3i}{4096\pi^3 \epsilon^2} - \frac{3i(1 + 2\log(2))}{2048\pi^3 \epsilon} + \mathcal{O}(\epsilon^0) \\
M[g_+ \omega_+, \omega_+, \omega_-] &= -\frac{3i}{4096\pi^3 \epsilon^2} - \frac{3i(1 + 2\log(2))}{2048\pi^3 \epsilon} + \mathcal{O}(\epsilon^0) \\
M[g_+ \omega_+, \omega_+, \omega_+] &= \frac{i}{4096\pi^3 \epsilon^2} + \frac{i(1 + 2\log(2))}{2048\pi^3 \epsilon} + \mathcal{O}(\epsilon^0) \\
M[g_- g_+, 0] &= \frac{1}{30720\pi^4 \epsilon} + \frac{58 + 15\log(2)}{115200\pi^4} + \mathcal{O}(\epsilon^1)
\end{aligned}$$

$$\begin{aligned}
M[g_+g_+, 0] &= -\frac{1}{30720\pi^4\epsilon} - \frac{58 + 15\log(2)}{115200\pi^4} + \mathcal{O}(\epsilon^1) \\
M[g_-g_+\omega_+, \omega_+] &= -\frac{1}{4096\pi^4\epsilon} - \frac{7 + 2\log(2)}{2048\pi^4} + \mathcal{O}(\epsilon^1) \\
M[g_+g_+\omega_+, \omega_-] &= \frac{3}{4096\pi^4\epsilon} + \frac{17 + 6\log(2)}{2048\pi^4} + \mathcal{O}(\epsilon^1) \\
M[g_+g_+\omega_+, \omega_+] &= -\frac{1}{4096\pi^4\epsilon} - \frac{3 + 2\log(2)}{2048\pi^4} + \mathcal{O}(\epsilon^1) \\
M[g_-g_+, \omega_+] &= \frac{i}{8192\pi^3\epsilon} + \frac{i(7 + 10\log(2))}{8192\pi^3} + \mathcal{O}(\epsilon^1) \\
M[g_+g_+, \omega_-] &= \frac{i}{8192\pi^3\epsilon} + \frac{i(7 + 10\log(2))}{8192\pi^3} + \frac{i(342 + 5\pi^2 + 60\log(2)(7 + 5\log(2)))\epsilon}{49152\pi^3} + \mathcal{O}(\epsilon^2) \\
M[g_+g_+, \omega_+] &= -\frac{i}{8192\pi^3\epsilon} - \frac{i(7 + 10\log(2))}{8192\pi^3} - \frac{i(342 + 29\pi^2 + 60\log(2)(7 + 5\log(2)))\epsilon}{49152\pi^3} + \mathcal{O}(\epsilon^2) \\
M[g_-g_+\omega_+, 0] &= 0 \\
M[g_+g_+\omega_+, 0] &= -\frac{i\epsilon}{384\pi^3} + \frac{i(-34 + 9\log(2))\epsilon^2}{576\pi^3} + \mathcal{O}(\epsilon^3) \\
M[g_-g_+g_+] &= -\frac{1}{12288\pi^4} - \frac{(41 + 6\log(2))\epsilon}{36864\pi^4} + \mathcal{O}(\epsilon^2) \\
M[g_+g_+g_+] &= \frac{1}{4096\pi^4} + \frac{(41 + 6\log(2))\epsilon}{12288\pi^4} + \frac{(587 - 6\pi^2 + 3\log(2)(41 + 3\log(2)))\epsilon^2}{18432\pi^4} + \mathcal{O}(\epsilon^3) \\
M[\text{tail}_+\text{tail}_+] &= -\frac{3}{4096\pi^4\epsilon} - \frac{3(8 + \log(2))}{2048\pi^4} - \frac{3(90 + \log(2)(16 + \log(2)))\epsilon}{2048\pi^4} + \mathcal{O}(\epsilon^2) \\
M[\text{tail}_+\text{tail}_-] &= -\frac{1}{4096\pi^4\epsilon} - \frac{8 + \log(2)}{2048\pi^4} + \mathcal{O}(\epsilon^1) \\
M[g_-g_+g_+\omega_+] &= \frac{i}{4096\pi^3} + \frac{i(1 + 4\log(2))\epsilon}{1024\pi^3} + \mathcal{O}(\epsilon^2) \\
M[g_+g_+g_+\omega_+] &= -\frac{i}{4096\pi^3} - \frac{i(1 + 4\log(2))\epsilon}{1024\pi^3} + \mathcal{O}(\epsilon^2) \\
M[\text{tail}_+\text{tail}_+\omega_+] &= \frac{1024\pi^3\epsilon}{i} + \frac{i(1 + 2\log(2))}{128\pi^3} + \mathcal{O}(\epsilon^1) \\
M[\text{tail}_+\text{tail}_-\omega_+] &= 0 \\
G'_1(x) &= -\frac{96x(x^4 + 1)\varpi_0(x)^2}{(x - 1)^2(x + 1)^2(x^2 + 1)^2\alpha_1(x)} \\
G'_2(x) &= -\frac{16(7x^{12} + 314x^{10} + 329x^8 + 1340x^6 + 329x^4 + 314x^2 + 7)\varpi_0(x)^2}{3(x - 1)^3x(x + 1)^3(x^2 + 1)^3}
\end{aligned}$$

$$\begin{aligned}
& + \frac{16(7x^8 + 136x^6 + 42x^4 + 136x^2 + 7)\varpi_0(x)^2\alpha_1'(x)}{3(x-1)^2(x+1)^2(x^2+1)^2\alpha_1(x)} - \frac{2(5x^8 + 28x^6 + 262x^4 + 28x^2 + 5)\varpi_0(x)^2\alpha_1'(x)^2}{3(x-1)x(x+1)(x^2+1)\alpha_1(x)^2} \\
& - \frac{xG_1(x)^2\alpha_1(x)^2}{(x-1)(x+1)(x^2+1)\varpi_0(x)^2}, \\
G_3'(x) &= -\frac{xG_1(x)\alpha_1(x)^2}{(x-1)(x+1)(x^2+1)\varpi_0(x)^2}, \\
G_4'(x) &= -\frac{16(7x^8 + 136x^6 + 42x^4 + 136x^2 + 7)\varpi_0(x)^2}{3(x-1)^2x(x+1)^2(x^2+1)^2\alpha_1(x)} + \frac{4(5x^8 + 28x^6 + 262x^4 + 28x^2 + 5)\varpi_0(x)^2\alpha_1'(x)}{3(x-1)x^2(x+1)(x^2+1)\alpha_1(x)^2} \\
& + \frac{G_2(x)}{x\alpha_1(x)}, \\
G_5'(x) &= \varpi_{0,K3}(x)\varpi_0''(x), \\
G_6'(x) &= \varpi_{0,K3}(x)\varpi_0'(x), \\
G_7'(x) &= \frac{\varpi_{0,K3}(x)G_3(x)\varpi_0'(x)}{\alpha_1(x)}, \\
G_8'(x) &= \frac{\varpi_{0,K3}(x)G_3(x)\varpi_0(x)\alpha_1'(x)}{\alpha_1(x)^2}, \\
G_9'(x) &= \frac{\varpi_{0,K3}(x)\varpi_0'(x)}{\alpha_1(x)} - \frac{\varpi_{0,K3}(x)\varpi_0(x)\alpha_1'(x)}{2\alpha_1(x)^2}, \\
G_{10}'(x) &= -\frac{(3x^4 - 2x^2 + 3)\varpi_{0,K3}(x)\varpi_0(x)}{32x^2} + \frac{4x^3(3x^2 - 5)G_5(x)}{32x^2} + \frac{(27x^4 - 10x^2 + 3)G_6(x)}{32x^2}, \\
G_{11}'(x) &= -\frac{3(x-1)(x+1)(x^2+1)\varpi_{0,K3}(x)\varpi_0(x)}{8x^2} + \frac{12x^5G_5(x)}{8x^2} + \frac{(27x^4 - 4x^2 - 3)G_6(x)}{8x^2}, \\
G_{12}'(x) &= -\frac{(x^{12} + 2x^{10} - 73x^8 + 236x^6 - 73x^4 + 2x^2 + 1)\varpi_{0,K3}(x)\varpi_0(x)}{16(x-1)^2x^2(x+1)^2(x^2+1)^2} - \frac{(x^2-3)(3x^2-1)\varpi_{0,K3}(x)G_1(x)\alpha_1(x)}{32(x-1)(x+1)(x^2+1)\varpi_0(x)} \\
& + \left[\frac{3(x-1)(x+1)(x^2+1)\varpi_{0,K3}(x)\varpi_0(x)}{16x^2\alpha_1(x)} - \frac{(x^2-3)(3x^2-1)G_5(x)}{32x\alpha_1(x)} - \frac{(9x^4-10x^2-3)G_6(x)}{32x^2\alpha_1(x)} \right] G_3(x) \\
& + \frac{(x^2-3)(3x^2-1)(x^4-16x^2+1)G_5(x)}{48(x-1)x(x+1)(x^2+1)} + \frac{(9x^{12} - 143x^{10} - 85x^8 + 624x^6 - 551x^4 + 47x^2 + 3)G_6(x)}{24(x-1)^2x^2(x+1)^2(x^2+1)^2} \\
& - \frac{(9x^4 - 10x^2 - 3)G_7(x)}{16x^2} + \frac{(9x^4 - 10x^2 - 3)G_8(x)}{32x^2} - \frac{2(x^4 - 16x^2 + 1)G_{10}(x)}{3(x-1)x(x+1)(x^2+1)}, \\
G_{13}'(x) &= -\frac{15(x^2+1)\varpi_0(x)}{(x-1)^2(x+1)^2} - \frac{1}{\varpi_{0,K3}(x)} \left[\frac{13(x^2-3)(3x^2-1)G_5(x)}{6(x-1)x(x+1)} + \frac{13(9x^4-10x^2-3)G_6(x)}{6(x-1)x^2(x+1)} \right. \\
& \left. - \frac{208G_{10}(x)}{3(x-1)x(x+1)} \right], \\
G_{14}'(x) &= -\frac{1}{\varpi_{0,K3}(x)} \left[-\frac{3(x^2+1)G_5(x)}{4x} - \frac{(9x^4-4x^2+3)G_6(x)}{4(x-1)x^2(x+1)} + \frac{2G_{11}(x)}{(x-1)x(x+1)} \right], \\
G_{15}'(x) &= -\frac{(x^8 - 12x^6 - 2x^4 - 12x^2 + 1)\varpi_{0,K3}(x)\varpi_0(x)}{4(x-1)x^2(x+1)(x^2+1)} - \frac{3\varpi_{0,K3}(x)G_1(x)\alpha_1(x)}{8\varpi_0(x)} \\
& + \left[\frac{(3x^4 - 2x^2 + 3)\varpi_{0,K3}(x)\varpi_0(x)}{4x^2\alpha_1(x)} - \frac{3(x-1)(x+1)(x^2+1)G_5(x)}{8x\alpha_1(x)} - \frac{(9x^4 - 4x^2 + 3)G_6(x)}{8x^2\alpha_1(x)} \right. \\
& \left. + \frac{G_{11}(x)}{x\alpha_1(x)} \right] G_3(x) + \frac{(x^4 - 16x^2 + 1)G_5(x)}{4x} + \frac{(9x^8 - 89x^6 + 32x^4 - 11x^2 + 3)G_6(x)}{6(x-1)x^2(x+1)(x^2+1)} - \frac{3(3x^4 + 1)G_7(x)}{4x^2} \\
& + \frac{3(3x^4 + 1)G_8(x)}{8x^2} - \frac{2(x^4 - 16x^2 + 1)G_{11}(x)}{3(x-1)x(x+1)(x^2+1)},
\end{aligned}$$

$$\begin{aligned}
G'_{16}(x) &= -\frac{3(x-1)(x+1)(x^2+1)\varpi_{0,K3}(x)\varpi_0(x)}{16x^2\alpha_1(x)} + \frac{(x^2-3)(3x^2-1)G_5(x)}{32x\alpha_1(x)} + \frac{(9x^4-10x^2-3)G_6(x)}{32x^2\alpha_1(x)} \\
&\quad + \frac{(9x^4-10x^2-3)G_9(x)}{16x^2} - \frac{G_{10}(x)}{x\alpha_1(x)}, \\
G'_{17}(x) &= -\frac{(3x^4-2x^2+3)\varpi_{0,K3}(x)\varpi_0(x)}{4x^2\alpha_1(x)} - \frac{3(x-1)(x+1)(x^2+1)G_5(x)}{8x\alpha_1(x)} + \frac{(9x^4-4x^2+3)G_6(x)}{8x^2\alpha_1(x)} \\
&\quad + \frac{3(3x^4+1)G_9(x)}{4x^2} - \frac{G_{11}(x)}{x\alpha_1(x)}, \\
G'_{18}(x) &= \frac{\varpi_0(x)}{(x-1)^2}, \\
G'_{19}(x) &= \frac{\varpi_0(x)}{x}, \\
G'_{20}(x) &= \frac{\varpi_0(x)}{(x+1)^2},
\end{aligned}$$

$$\varpi_{0,K3} = \left(\frac{2}{\pi}\right)^2 K^2(1-x^2) \text{ and } f'(x) = \frac{d}{dx}f(x)$$

4. Espacio Cuántico Relativista o Curvo. La curvatura en un campo de gauge indiscriminado, se forma por la gravedad que le resulta permeada, transferida por la interacción del gravitón y su compañera, con las partículas subatómicas de origen bosónico o fermiónico, las mismas que se vuelven superpartículas, lo que distorsiona el espacio – tiempo cuántico. Como ha quedado explícito en este trabajo, el efecto cuántico gravitacional, puede ser intrínseco o endógeno, en tratándose de las partículas supermasivas y exógeno o extrínseco, esto es, por interacción de la partícula de que se trate con el campo gravitónico de gauge. Así es como concluimos que la supergravedad cuántica deriva de una gravedad diseñada en espacios de calibre arbitrarios.

El Modelo Matemático bajo la métrica Friedmann-Lemaître-Robertson-Walker, queda expresado así:

$$\begin{aligned}
S[\phi] &= \frac{1}{2} \int_{\mathbb{R} \times \mathcal{M}} d^d x \sqrt{g} \{g^{\mu\nu} \partial_\mu \phi \partial_\nu \phi + m^2 \phi^2 + \xi R \phi^2\}, \\
Z[J] &\equiv e^{-\frac{1}{\hbar}W[J]} := \int \mathcal{D}\phi e^{-\frac{1}{\hbar}S[\phi] + \frac{1}{\hbar}(J,\phi)} \\
\langle \phi(x_1)\phi(x_2) \dots \phi(x_N) \rangle &= \frac{\hbar^N}{Z[0]} \frac{\delta}{\delta J(x_N)} \dots \frac{\delta}{\delta J(x_2)} \frac{\delta}{\delta J(x_1)} Z[J] \Big|_{J=0} \\
\Gamma[\phi_J] &:= W[J] + (J, \phi_J), \\
\frac{\delta \Gamma[\phi_J]}{\delta \phi_J(x)} &= J(x) \\
Z[J] &= e^{-\frac{1}{\hbar}S[\phi_J] + \frac{1}{\hbar}(J,\phi_J)} \int \mathcal{D}\phi e^{-\frac{1}{\hbar}(\delta S - J,\phi) - \frac{1}{2\hbar}(\phi, (\delta^2 S, \phi))} \\
\delta S(x) &:= \frac{\delta S[\phi]}{\delta \phi(x)} \Big|_{\phi=\phi_J} \quad \delta^2 S(x, y) := \frac{\delta^2 S[\phi]}{\delta \phi(x) \delta \phi(y)} \Big|_{\phi=\phi_J} \\
\Gamma[\phi_J] &= S[\phi_J] - \hbar \log \int \mathcal{D}\phi e^{-\frac{1}{2}(\phi, (\delta^2 S, \phi))} + \mathcal{O}(\hbar^2), \\
\frac{\delta S[\phi]}{\delta \phi(x)} \Big|_{\phi=\phi_{\text{clás}}} &= J(x), \\
\phi(x) &= \sum_n c_n \varphi_n(x), \text{ con } c_n = (\varphi_n, \phi), \\
(\phi, A\phi) &= \sum_{n,m} c_n c_m (\varphi_n, A\varphi_m) = \sum_n \lambda_n c_n^2,
\end{aligned}$$



$$\begin{aligned}
\mathcal{D}\phi &= \prod_n dc_n \\
\int \mathcal{D}\phi e^{-\frac{1}{2}(\phi, A\phi)} &\sim \int \left(\prod_n dc_n \right) e^{-\frac{1}{2} \sum_n \lambda_n c_n^2} = \prod_n \int dc_n e^{-\frac{1}{2} \lambda_n c_n^2} \\
&\sim \prod_n \lambda_n^{-1/2} \\
\Gamma^{(1-\text{loop})}[\phi] &= \frac{1}{2} \log \text{Det}A. \\
\Gamma^{(1-\text{loop})}[\psi] &= -\log \text{Det}A. \\
A &= -\partial_\beta^2 - \Delta + \xi R + m^2, \\
\Delta &= \frac{1}{\sqrt{g}} \partial_i (g^{ij} \sqrt{g} \partial_j) \\
\log \text{Det}A &= -\frac{d}{ds} \zeta_A(s) \Big|_{s=0}, \\
\zeta_A(s) &:= \text{Tr} A^{-s} \\
\zeta_A(s) &= \sum_{n \in \mathbb{N}} \lambda_n^{-s} \\
\lambda^{-s} &= \frac{1}{\Gamma(s)} \int_0^\infty dt t^{s-1} e^{-t} \\
\sum_n f(n) &= \sum_m \tilde{f}(m), \\
\sum_{k \in \mathbb{Z}} e^{-(k+c)^2 t} &= \sqrt{\frac{\pi}{t}} \sum_{k \in \mathbb{Z}} e^{-\pi^2 k^2 / t} e^{2\pi i k c} \\
\zeta_A(s) &= \frac{1}{\Gamma(s)} \int_0^\infty dt t^{s-1} \text{Tr} e^{-tA} \\
K_A(t) &:= \text{Tr} e^{-tA} = \sum_{n \in \mathbb{N}} e^{-\lambda_n t} \\
\log \text{Det}A &= -\int_0^\infty \frac{dt}{t} K_A(t) \\
\zeta_A(s) &= \sum_n (\lambda_n / \mu^2)^{-s} \\
\zeta(s) &= \sum_{\lambda_n > 0} \lambda_n^{-s} + (-1)^{-s} \sum_{\lambda_n < 0} (-\lambda_n)^{-s} \\
\eta(s) &= \sum_{\lambda_n > 0} \lambda_n^{-s} - \sum_{\lambda_n < 0} (-\lambda_n)^{-s}, \\
\zeta'_D(0) &= -i \frac{\pi}{2} [\zeta_{D^2}(0) + \eta_D(0)] + \frac{1}{2} \zeta'_{D^2}(0). \\
a(A, B) &:= \log \frac{\text{Det}AB}{\text{Det}A \text{Det}B}, \\
\zeta_R(s) &:= \sum_{n=1}^\infty n^{-s} \\
\zeta_R(s) &= \sum_{n=1}^\infty \frac{1}{\Gamma(s)} \int_0^\infty dt t^{s-1} e^{-nt} \\
\zeta_R(s) &= \frac{1}{\Gamma(s)} \int_0^\infty dt \frac{t^{s-1}}{e^t - 1}
\end{aligned}$$

$$\frac{t}{e^t - 1} = \sum_{n=0}^{\infty} \frac{B_n}{n!} t^n$$

$$\zeta_R(s) = \frac{1}{\Gamma(s)} \int_1^{\infty} dt \frac{t^{s-1}}{e^t - 1} + \frac{1}{\Gamma(s)} \sum_{n=0}^{\infty} \frac{B_n}{n!} \frac{1}{s+n-1}$$

$$\zeta_R(s) = \frac{1}{s-1} + \mathcal{O}(1)$$

$$\zeta_H(s, q) := \sum_{n=0}^{\infty} (n+q)^{-s}, q \neq 0, -1, -2, \dots$$

$$ds^2 = -dt^2 + a^2(\eta) \left[\frac{dr^2}{1 - \kappa r^2} + r^2(d\theta^2 + \text{sen}^2\theta d\phi^2) \right]$$

$$x^0 = a \cos \chi \cos \theta$$

$$x^2 = a \text{sen} \chi \text{sen} \phi$$

$$x^3 = a \cos \chi \text{sen} \theta$$

$$\frac{ds^2}{a^2} = d\chi^2 + \cos^2 \chi d\theta^2 + \text{sen}^2 \chi d\phi^2$$

$$g^1 = \frac{1}{a^2}; g^2 = \frac{1}{a^2 \cos^2 \chi}; g^3 = \frac{1}{a^2 \text{sen}^2 \chi}$$

$$\Delta_{S^3} = \frac{1}{a^2} \left[\frac{1}{\cos \chi \text{sen} \chi} \partial_{\chi} (\cos \chi \text{sen} \chi \partial_{\chi}) + \frac{1}{\cos^2 \chi} \partial_{\theta}^2 + \frac{1}{\text{sen}^2 \chi} \partial_{\phi}^2 \right]$$

$$\chi_k(h) = \sum_{\alpha=1}^s a_{\alpha} \chi^{(\alpha)}(h).$$

$$a_{\alpha} = \frac{1}{|H|} \sum_{h \in H} \chi_k(h) \chi^{(\alpha)*}(h),$$

$$a_1 = \frac{1}{|H|} \sum_{h \in H} \chi_k(h)$$

$$\chi_k(\theta_L, \theta_R) = \frac{\text{sen}(k\theta_L/2) \text{sen}(k\theta_R/2)}{\text{sen}(\theta_L/2) \text{sen}(\theta_R/2)}$$

$$\chi_k(h) = k \frac{\text{sen}(k\theta_h/2)}{\text{sen}(\theta_h/2)}$$

$$d_k^{(H)} = \frac{k}{|H|} \sum_{h \in H} \frac{\text{sen}(k\theta_h/2)}{\text{sen}(\theta_h/2)}$$

$$f(\chi, \theta, \phi) = X(\chi)\Theta(\theta)\Phi(\phi),$$

$$-\frac{1}{\Phi} \frac{d^2\Phi}{d\phi^2} = m_1^2$$

$$-\frac{1}{\Theta} \frac{d^2\Theta}{d\theta^2} = m_2^2$$

$$-\frac{1}{\cos \chi \text{sen} \chi} \frac{d}{d\chi} \left(\cos \chi \text{sen} \chi \frac{dX}{d\chi} \right) + \left(\frac{m_1^2}{\text{sen}^2 \chi} + \frac{m_2^2}{\cos^2 \chi} \right) X = a^2 \lambda X,$$

$$\Phi(\phi) = e^{im_1\phi}$$

$$\Theta(\theta) = e^{im_2\theta}$$

$$X(\chi) = \cos^{\alpha} \chi \text{sen}^{\beta} \chi \tilde{X}(\chi)$$

$$\tilde{X}'' - \left[\frac{m_1^2 - \frac{1}{4}}{\text{sen}^2 \chi} + \frac{m_2^2 - \frac{1}{4}}{\cos^2 \chi} - (a^2 \lambda + 1) \right] \tilde{X} = 0.$$

$$\tilde{X}_{n, m_1, m_2}(\chi) = (\text{sen} \chi)^{|m_1| + \frac{1}{2}} (\cos \chi)^{|m_2| + \frac{1}{2}} P_n^{(|m_1|, |m_2|)}(\cos 2\chi),$$

$$a^2 \lambda = (2n + |m_1| + |m_2| + 1)^2 - 1$$

$$f_{k,m_1,m_2}(\chi, \theta, \phi) = e^{im_1\phi} e^{im_2\theta} \operatorname{sen}^{|m_1|} \chi \cos^{|m_2|} \chi P_{(k-|m_1|-|m_2|-1)/2}^{(|m_1|, |m_2|)}(\cos 2\chi)$$

$$d_k = 1 + \sum_{n=0}^{\frac{k-1}{2}-1} 4(k-1-2n) = k^2$$

$$d_k = \sum_{n=0}^{\frac{k-2}{2}} 4(k-1-2n) = k^2$$

$$\theta \rightarrow \theta + \frac{2\pi m}{p}, \phi \rightarrow \phi + \frac{2\pi m}{p}$$

$$k = 2n + |m_1| + |m_2| + 1$$

$$lp = m_1 + m_2$$

$$2b = 2n + |m_1| + |m_2|$$

$$(2q+1)l = m_1 + m_2$$

$$d_{2b+1}^{(I)} = \sum_{l=0}^{l_{\max}} [(2q+1)l + 1]$$

$$d_{2b+1}^{(II)} = d_{2b+1}^{(I)}$$

$$d_{2b+1}^{(III)} = b + 2 \sum_{j=1}^b \left\lfloor \frac{j}{2q+1} \right\rfloor$$

$$d_{2b+1}^{(III,-)} = \begin{cases} \left\lfloor \frac{2b}{2q+1} \right\rfloor + 1, & \text{si } r \text{ es par} \\ \left\lfloor \frac{2b}{2q+1} \right\rfloor, & \text{si } r \text{ es impar} \end{cases}$$

$$d_{2b+1}^{(IV)} = \begin{cases} d_{2b+1}^{(III)} - \left\lfloor \frac{2b}{2q+1} \right\rfloor, & \text{si } r \text{ es par} \\ d_{2b+1}^{(III)} - \left\lfloor \frac{2b}{2q+1} \right\rfloor + 1, & \text{si } r \text{ es impar} \end{cases}$$

$$d_k^{k \text{ impar}} = \begin{cases} k \left(\left\lfloor \frac{k-1}{2q+1} \right\rfloor + 1 \right), & \text{si } r \text{ es par} \\ k \left\lfloor \frac{k-1}{2q+1} \right\rfloor, & \text{si } r \text{ es impar} \end{cases}$$

$$d_k^{(p=2q+1)} = \begin{cases} k \left(\frac{k-r}{2q+1} + 1 \right), & \text{si } r \text{ es impar} \\ k \left(\frac{k-r}{2q+1} \right), & \text{si } r \text{ es par} \end{cases}$$

$$k = 2n + |m_1| + |m_2| + 1$$

$$2lq = m_1 + m_2$$

$$2b = 2n + |m_1| + |m_2|,$$

$$2lq = m_1 + m_2.$$

$$d_{2b+1}^{(II)} = \sum_{l=1}^{\lfloor b/q \rfloor} (2lq + 1).$$

$$d_{2b+1}^{(IV)} = b + 2 \sum_{n=0}^{b-1} \left\lfloor \frac{b-n}{q} \right\rfloor$$

$$d_k^{(p=2q)} = \begin{cases} k \left(1 + 2 \left\lfloor \frac{k}{2q} \right\rfloor \right), & \text{si } k \text{ es impar} \\ 0, & \text{si } k \text{ es par} \end{cases}$$

$$\frac{\text{sen}(k(\theta + 2\pi)/2)}{\text{sen}((\theta + 2\pi)/2)} = \frac{\text{sen}(k\pi + k\theta/2)}{\text{sen}(\pi + \theta/2)} = (-1)^{k+1} \frac{\text{sen}(k\theta/2)}{\text{sen}(\theta/2)}$$

$$d_k^{(H^*)} = \frac{k}{|H^*|} [1 + (-1)^{k+1}] \sum_{h \in H} \frac{\text{sen}(k\theta_h/2)}{\text{sen}(\theta_h/2)}$$

$$d_k^{(p)} = \frac{k}{2p} \left[\lim_{\theta \rightarrow 0} \frac{\text{sen}(k\theta/2)}{\text{sen}(\theta/2)} + \sum_{m=1}^{p-1} \frac{\text{sen}(km\pi/p)}{\text{sen}(m\pi/p)} + p \frac{\text{sen}(k\pi/2)}{\text{sen}(\pi/2)} \right]$$

$$\sum_{m=1}^{p-1} \frac{\text{sen}(km\pi/p)}{\text{sen}(m\pi/p)} = p - k + 2p \left\lfloor \frac{k}{2p} \right\rfloor$$

$$d_k^{(p)} = \begin{cases} k \left(\left\lfloor \frac{k}{2p} \right\rfloor + \frac{1 + (-1)^{\lfloor (k-1)/2 \rfloor}}{2} \right), & \text{si } k \text{ es impar} \\ 0, & \text{si } k \text{ es par} \end{cases}$$

$$d_{2n+1}^{(p)} = \begin{cases} (2n+1) \left\lfloor \frac{n}{p} \right\rfloor, & \text{si } n \text{ es impar} \\ (2n+1) \left(\left\lfloor \frac{n}{p} \right\rfloor + 1 \right), & \text{si } n \text{ es par} \end{cases}$$

$$d_{2n+1}^{(T^*)} = (2n+1) \left\{ 2 \left\lfloor \frac{n}{3} \right\rfloor + \left\lfloor \frac{n}{2} \right\rfloor - n + 1 \right\}$$

$$d_{2n+1}^{(O^*)} = (2n+1) \left\{ \left\lfloor \frac{n}{4} \right\rfloor + \left\lfloor \frac{n}{3} \right\rfloor + \left\lfloor \frac{n}{2} \right\rfloor - n + 1 \right\}$$

$$d_{2n+1}^{(I^*)} = (2n+1) \left\{ \left\lfloor \frac{n}{5} \right\rfloor + \left\lfloor \frac{n}{3} \right\rfloor + \left\lfloor \frac{n}{2} \right\rfloor - n + 1 \right\}$$

$$S[\phi] = \frac{1}{2} \int_{\mathbb{R} \times \mathcal{M}} d^4x \sqrt{g} \left(g^{\mu\nu} \partial_\mu \phi \partial_\nu \phi + \frac{1}{a^2} \phi^2 \right),$$

$$A = -\frac{1}{\sqrt{g}} \partial_\mu (g^{\mu\nu} \sqrt{g} \partial_\nu) + \frac{1}{a^2}$$

$$\zeta_{S^1 \times \mathcal{M}}(s) = \mu^{2s} \sum_{l \in \mathbb{Z}} \sum_{k=1}^{\infty} d_k^{(H)} \left[\left(\frac{2\pi l}{\beta} \right)^2 + \left(\frac{k}{a} \right)^2 \right]^{-s}$$

$$\zeta_{S^3}(s) = \mu^{2s} \sum_{l \in \mathbb{Z}} \sum_{k=1}^{\infty} k^2 \left[\left(\frac{2\pi l}{\beta} \right)^2 + \left(\frac{k}{a} \right)^2 \right]^{-s}$$

$$\zeta_{S^3}(s) = (\mu a)^{2s} \sum_{l \in \mathbb{Z}} \sum_{k=1}^{\infty} k^2 \left[\left(\frac{2\pi a l}{\beta} \right)^2 + k^2 \right]^{-s}$$

$$\zeta_{S^3}^{l=0}(s) = (\mu a)^{2s} \sum_{k=1}^{\infty} k^{-(2s-2)} = (\mu a)^{2s} \zeta_R(2s-2).$$

$$\Gamma_{S^3}^{l=0} = \frac{\zeta_R(3)}{4\pi^2}$$

$$\zeta_{S^3}^{l \neq 0}(s) = (\mu a)^{2s} \left\{ \sum_{l=1}^{\infty} \sum_{k \in \mathbb{Z}} \left[k^2 + \left(\frac{2l\pi a}{\beta} \right)^2 \right]^{-(s-1)} - \sum_{l=1}^{\infty} \left(\frac{2l\pi a}{\beta} \right)^2 \sum_{k \in \mathbb{Z}} \left[k^2 + \left(\frac{2l\pi a}{\beta} \right)^2 \right]^{-s} \right\}$$

$$\zeta_1(s) = \frac{(\mu a)^{2s}}{\Gamma(s-1)} \sum_{l=1}^{\infty} \sum_{k \in \mathbb{Z}} \int_0^{\infty} dt t^{s-1-1} e^{-[k^2 + (2\pi a l/\beta)^2]t}$$

$$\zeta_1(s) = \frac{(\mu a)^{2s} \sqrt{\pi}}{\Gamma(s-1)} \sum_{l=1}^{\infty} \sum_{k \in \mathbb{Z}} \int_0^{\infty} dt t^{s-\frac{3}{2}-1} e^{-\pi^2 k^2/t} e^{-(2l\pi a)^2 t/\beta^2}$$

$$\zeta_1(s) = \frac{(\mu a)^{2s}(s-1)\sqrt{\pi}}{\Gamma(s)} \left\{ \Gamma\left(s - \frac{3}{2}\right) \left(\frac{2\pi a}{\beta}\right)^{3-2s} \zeta_R(2s-3) \right. \\ \left. + 2 \sum_{l=1}^{\infty} \sum_{k=1}^{\infty} \left(\frac{2al}{\beta k}\right)^{\frac{3}{2}-s} K_{s-\frac{3}{2}}(4\pi^2 lka/\beta) \right\}$$

$$\zeta_2(s) = -\frac{(\mu a)^{2s}\sqrt{\pi}}{\Gamma(s)} \left\{ \Gamma\left(s - \frac{1}{2}\right) \left(\frac{2\pi a}{\beta}\right)^{3-2s} \zeta_R(2s-3) \right. \\ \left. + 2 \sum_{l=1}^{\infty} \sum_{k=1}^{\infty} \left(\frac{2\pi al}{\beta}\right)^2 \left(\frac{2al}{\beta k}\right)^{\frac{1}{2}-s} K_{s-\frac{1}{2}}(4\pi^2 lka/\beta) \right\}$$

$$\Gamma_{S^3} = \frac{\zeta_R(3)}{4\pi^2} - \frac{\pi^4}{45} \left(\frac{a}{\beta}\right)^3 \\ + \frac{1}{4\pi^2} \sum_{k,l=1}^{\infty} \frac{1}{k^3} \left[2 + 2 \frac{4\pi^2 kla}{\beta} + \left(\frac{4\pi^2 kla}{\beta}\right)^2 \right] e^{-4\pi^2 kla/\beta}.$$

$$\zeta_{S^3}(s) = \left(\frac{\mu\beta}{2\pi}\right)^{2s} \sum_{l \in \mathbb{Z}} \sum_{k=1}^{\infty} k^2 \left[\left(\frac{\beta k}{2\pi a}\right)^2 + l^2 \right]^{-s}$$

$$\zeta_{S^3}(s) = \frac{1}{\Gamma(s)} \left(\frac{\mu\beta}{2\pi}\right)^{2s} \sum_{l \in \mathbb{Z}} \sum_{k=1}^{\infty} k^2 \int_0^{\infty} dt t^{s-1} e^{-(\beta k/2\pi a)^2 + l^2} t}$$

$$\zeta_{S^3}(s) = \frac{\sqrt{\pi}}{\Gamma(s)} \left(\frac{\mu\beta}{2\pi}\right)^{2s} \sum_{k=1}^{\infty} k^2 \sum_{l \in \mathbb{Z}} \int_0^{\infty} dt t^{s-\frac{3}{2}} e^{-(\beta k/2\pi a)^2 t} e^{-(\pi l)^2/t}$$

$$\zeta_{S^3}^{l=0}(s) = \frac{(\mu a)^{2s}\beta \Gamma\left(s - \frac{1}{2}\right)}{2\sqrt{\pi} a \Gamma(s)} \zeta_R(2s-3)$$

$$\zeta_{S^3}^{l \neq 0}(s) = \frac{4\sqrt{\pi}}{\Gamma(s)} \left(\frac{2\pi}{\mu\beta}\right)^{-2s} \sum_{k=1}^{\infty} k^2 \sum_{l=1}^{\infty} \left(\frac{2\pi^2 al}{\beta k}\right)^{s-\frac{1}{2}} K_{\frac{1}{2}-s}(kl\beta/a)$$

$$\Gamma_{S^3} = \frac{\beta}{2a} \zeta_R(-3) - \sqrt{\frac{2\beta}{\pi a}} \sum_{k=1}^{\infty} k^2 \sum_{l=1}^{\infty} \sqrt{\frac{k}{l}} K_{\frac{1}{2}}(kl\beta/a),$$

$$\Gamma_{S^3} = \frac{\beta}{240a} - \sum_{k=1}^{\infty} k^2 \sum_{l=1}^{\infty} \frac{e^{-kl\beta/a}}{l} = \frac{\beta}{240a} + \sum_{k=1}^{\infty} k^2 \log(1 - e^{-k\beta/a}).$$

$$\sum_{n=0}^{\infty} f(n) = \oint_{\mathcal{C}} dz \frac{f(z)}{e^{2\pi iz} - 1}$$

$$\oint_{\mathcal{C}} dz \frac{f(z)}{e^{2\pi iz} - 1} = i \int_{\epsilon}^{\infty} dy \frac{f(iy)}{1 - e^{-2\pi y}} + \frac{1}{2} \text{Res} \left(\frac{f(z)}{e^{2\pi iz} - 1}, 0 \right) \\ + i \int_{\epsilon}^{\infty} dy \frac{f(-iy)}{1 - e^{2\pi y}}$$

$$\int_{\epsilon}^{\infty} dy \frac{f(iy)}{1 - e^{-2\pi y}} = \int_{\epsilon}^{\infty} dy f(iy) \frac{1 - e^{-2\pi y} + e^{-2\pi y}}{1 - e^{-2\pi y}}$$

$$= -i \int_0^{\infty} dx f(x + i\epsilon) + \int_{\epsilon}^{\infty} dy \frac{f(iy)}{e^{2\pi y} - 1}$$

$$\sum_{n=0}^{\infty} f(n) = \frac{1}{2} f(0) + \int_0^{\infty} dx f(x) + i \int_0^{\infty} dy \frac{f(iy) - f(-iy)}{e^{2\pi y} - 1}$$



$$\begin{aligned}
& \int_{C_\epsilon^{(n)}} dz \frac{z^2 \log(1 - e^{-z\beta/a})}{e^{2\pi iz} - 1}, \\
& \int_0^\infty dx x^2 \log(1 - e^{-x\beta/a}) = -\frac{\pi^4 a^3}{45\beta^3} \\
& i \int_0^\infty dy \frac{g(iy) - g(-iy)}{e^{2\pi y} - 1} = \int_0^\infty dy \frac{y^2(\pi - \beta x/a)}{e^{2\pi y} - 1} \\
& + 2\pi \sum_{n=0}^\infty n \int_{2\pi na/\beta}^{2\pi(n+1)a/\beta} dy \frac{y^2}{e^{2\pi y} - 1} \\
& \int_0^\infty dy \frac{y^2(\pi - \beta x/a)}{e^{2\pi y} - 1} = \frac{\zeta_R(3)}{4\pi^2} - \frac{\beta}{240a} \\
& \frac{1}{e^{2\pi x} - 1} = e^{-2\pi x} \frac{1}{1 - e^{-2\pi x}} = e^{-2\pi x} \sum_{k=0}^\infty e^{-2\pi kx}, \\
\Gamma_{S^3} = & -\frac{\pi^4 a^3}{45\beta^3} + \frac{\zeta_R(3)}{4\pi^2} + \frac{1}{4\pi^2} \sum_{l,k=1}^\infty \frac{1}{k^3} \left[2 + 2\frac{4\pi^2 lka}{\beta} + \left(\frac{4\pi^2 lka}{\beta}\right)^2 \right] e^{-4\pi^2 lka/\beta} \\
\zeta_{\mathcal{M}}(s) = & (\mu a)^{2s} \sum_{l \in \mathbb{Z}} \sum_{k=1}^\infty d_k^{(H)} \left[\left(\frac{2\pi a l}{\beta}\right)^2 + k^2 \right]^{-s}, \\
d_k^{(Z_{2q+1})} = & \begin{cases} \frac{k^2}{2q+1} - \frac{rk}{2q+1} + k, & \text{si } r \text{ es impar} \\ \frac{k^2}{2q+1} - \frac{rk}{2q+1}, & \text{si } r \text{ es par} \end{cases} \\
\zeta_{S^3/\mathbb{Z}_{2q+1}}(s) = & \frac{\zeta_{S^3}(s)}{2q+1} + \delta\zeta_{2q+1}(s) \\
\delta\zeta_{2q+1}(s) = & -\frac{(\mu a)^{2s}}{p} \sum_{l \in \mathbb{Z}} \sum_{n=0}^\infty \left\{ \sum_{r=1}^q 2r(np+2r) \left[(np+2r)^2 + \left(\frac{2\pi a l}{\beta}\right)^2 \right]^{-s} \right. \\
& \left. + \sum_{r=0}^{q-1} 2(r-q)(np+2r+1) \left[(np+2r+1)^2 + \left(\frac{2\pi a l}{\beta}\right)^2 \right]^{-s} \right\} \\
\delta\zeta_{2q+1}(s) = & -2 \left(\frac{\mu a}{p}\right)^{2s} \sum_{l \in \mathbb{Z}} \sum_{n=0}^\infty \left\{ \sum_{r=1}^q r \left(n + \frac{2r}{p}\right) \left[\left(n + \frac{2r}{p}\right)^2 + \left(\frac{2\pi a l}{p\beta}\right)^2 \right]^{-s} \right. \\
& \left. - \sum_{r=1}^q r \left(n + 1 - \frac{2r}{p}\right) \left[\left(n + 1 - \frac{2r}{p}\right)^2 + \left(\frac{2\pi a l}{p\beta}\right)^2 \right]^{-s} \right\} \\
\delta\zeta_{2q+1}^{l=0}(s) = & -2 \left(\frac{\mu a}{p}\right)^{2s} \sum_{r=1}^q r \left[\zeta_H\left(2s-1, \frac{2r}{p}\right) - \zeta_H\left(2s-1, 1 - \frac{2r}{p}\right) \right] \\
\delta\zeta_{2q+1}^{l \neq 0}(s) = & -4 \left(\frac{\mu a}{p}\right)^{2s} \frac{\Gamma(2-2s)(2\pi)^{2s-1}}{\Gamma(s)\Gamma(1-s)} \sum_{r=1}^q r \sum_{n=1}^\infty n^{2s-2} \text{sen}\left(\frac{4\pi nr}{p}\right) \\
\delta\zeta_{2q+1}^{l \neq 0}(s) = & -4 \left(\frac{\mu a}{p}\right)^{2s} \sum_{r=1}^q r \sum_{n \in \mathbb{Z}} \sum_{l=1}^\infty \left(n + \frac{r}{p}\right) \left[\left(n + \frac{r}{p}\right)^2 + \left(\frac{2\pi a l}{p\beta}\right)^2 \right]^{-s}
\end{aligned}$$

$$\begin{aligned}
\delta\zeta_{2q+1}^{l\neq 0}(s) &= -\frac{(\mu a)^{2s} p^{1-2s}}{(1-s)} \sum_{r=1}^q r \\
&\times \frac{d}{d\alpha} \sum_{n\in\mathbb{Z}} \sum_{l=1}^{\infty} \left[\left(n + \frac{2r+2\alpha-2}{p} \right)^2 + \left(\frac{2\pi a l}{p\beta} \right)^2 \right]^{-s+1} \\
\delta\zeta_{2q+1}^{l\neq 0}(s) &= -\frac{16(\mu a)^{2s} \pi^{\frac{3}{2}}}{\Gamma(s) p^{2s}} \sum_{r=1}^q r \sum_{l,n=1}^{\infty} n \operatorname{sen} \left(\frac{4\pi n r}{p} \right) \\
&\times \left(\frac{p\beta n}{a l} \right)^{s-\frac{3}{2}} K_{s-\frac{3}{2}}(4\pi^2 a l n / p\beta) \\
\Gamma_{S^3/Z_{p=2q+1}} &= \frac{\zeta_R(3)}{4\pi^2 p} + \frac{1}{\pi} \sum_{r=1}^q r \sum_{n=1}^{\infty} \frac{1}{n^2} \operatorname{sen} \left(\frac{4\pi n r}{p} \right) - \frac{\pi^4}{45p} \left(\frac{a}{\beta} \right)^3 \\
&+ \frac{1}{4\pi^2 p} \sum_{k,l=1}^{\infty} \frac{1}{k^3} \left[2 + 2 \frac{4\pi^2 a k l}{\beta} + \left(\frac{4\pi^2 a k l}{\beta} \right)^2 \right] e^{-4\pi^2 a k l / \beta} \\
&+ \frac{2}{\pi p} \sum_{r=1}^q r \sum_{n,l=1}^{\infty} \frac{1}{n^2} \operatorname{sen} \left(\frac{4\pi n r}{p} \right) \left(1 + \frac{4\pi^2 a l n}{p\beta} \right) e^{-4\pi^2 a n l / p\beta} \\
\zeta_{S^3/Z_2}(s) &= (\mu a)^{2s} \sum_{l\in\mathbb{Z}} \sum_{n=0}^{\infty} (2n+1)^2 [\tilde{\omega}_l^2 + (2n+1)^2]^{-s} \\
\zeta_{S^3/Z_2}(s) &= (\mu a)^{2s} \sum_{l\in\mathbb{Z}} \sum_{k=1}^{\infty} \left\{ k^2 (\tilde{\omega}_l^2 + k^2)^{-s} - (2k)^2 [\tilde{\omega}_l^2 + (2k)^2]^{-s} \right\}. \\
\Gamma_{S^3/Z_2}(\beta) &= \Gamma_{S^3}(\beta) - 4\Gamma_{S^3}(2\beta). \\
d_{2k+1}^{(Z_{2q})} &= \left(2 \left\lfloor \frac{2k+1}{2q} \right\rfloor + 1 \right) (2k+1) \\
d_{2k+1}^{(Z_{2q})} &= \left(2 \frac{k-r}{q} + 1 \right) (2k+1) = \frac{(2k+1)^2}{q} + \frac{q-2r-1}{q} (2k+1) \\
\zeta_{S^3/Z_{2q}}(s) &= \frac{\zeta_{S^3/Z_2}(s)}{q} + \delta\zeta_{2q}(s) \\
\delta\zeta_{2q}(s) &= 2 \left(\frac{\mu a}{2q} \right)^{2s} \sum_{r=0}^{q-1} [(q-r-1) - r] \\
&\times \sum_{l\in\mathbb{Z}} \sum_{n=0}^{\infty} \left(n + \frac{2r+1}{2q} \right) \left[\left(n + \frac{2r+1}{2q} \right)^2 + \left(\frac{\pi a l}{q\beta} \right)^2 \right]^{-s} \\
\delta\zeta_{2q}(s) &= 2 \left(\frac{\mu a}{2q} \right)^{2s} \sum_{r=1}^{q-1} r \sum_{l\in\mathbb{Z}} \sum_{n=0}^{\infty} \left\{ \left(n + 1 - \frac{2r+1}{2q} \right) \left[\left(n + 1 - \frac{2r+1}{2q} \right)^2 + \left(\frac{\pi a l}{q\beta} \right)^2 \right]^{-s} \right. \\
&\quad \left. - \left(n + \frac{2r+1}{2q} \right) \left[\left(n + \frac{2r+1}{2q} \right)^2 + \left(\frac{\pi a l}{q\beta} \right)^2 \right]^{-s} \right\} \\
\delta\zeta_{2q}^{l=0}(s) &= -2 \left(\frac{\mu a}{2q} \right)^{2s} \sum_{r=1}^{q-1} r \left\{ \zeta_H(2s-1, \frac{2r+1}{2q}) \right. \\
&\quad \left. - \zeta_H \left(2s-1, 1 - \frac{2r+1}{2q} \right) \right\}
\end{aligned}$$

$$\delta\zeta_{2q}^{l=0}(s) = -2 \left(\frac{\mu a}{2q}\right)^{2s} \left[\sum_{r=0}^{q-1} (2r+1) \zeta_H \left(2s-1, \frac{2r+1}{2q}\right) - q^{2s} \zeta_H \left(2s-1, \frac{1}{2}\right) \right]$$

$$\delta\zeta_{2q}^{l\neq 0}(s) = -\frac{16\pi^{3/2}}{q\Gamma(s)} \left(\frac{\mu a}{2q}\right)^{2s} \sum_{r=1}^{q-1} r \sum_{l,n=1}^{\infty} n \operatorname{sen} \left(2\pi n \frac{2r+1}{2q}\right) \times \left(\frac{q\beta n}{al}\right)^{s-\frac{3}{2}} K_{s-\frac{3}{2}} \left(\frac{2\pi^2 aln}{q\beta}\right)$$

$$\Gamma_{S^3/Z_{p=2q}} = -\frac{6\zeta_R(3)}{4\pi^2 p} + \frac{1}{\pi} \sum_{r=1}^{q-1} r \sum_{n=1}^{\infty} \frac{1}{n^2} \operatorname{sen} \left(2\pi n \frac{2r+1}{p}\right) - \frac{\pi^4}{45p} \left(\frac{a}{\beta}\right)^3$$

$$+ \frac{1}{2\pi^2 p} \sum_{k,l=1}^{\infty} \frac{1}{k^3} \left[2 + 2 \frac{4\pi^2 ak l}{\beta} + \left(\frac{4\pi^2 ak l}{\beta}\right)^2 \right] e^{-4\pi^2 ak l/\beta}$$

$$- \frac{2}{\pi^2 p} \sum_{k,l=1}^{\infty} \frac{1}{k^3} \left[2 + 2 \frac{2\pi^2 ak l}{\beta} + \left(\frac{2\pi^2 ak l}{\beta}\right)^2 \right] e^{-2\pi^2 ak l/\beta}$$

$$+ \frac{4}{\pi p} \sum_{r=1}^{q-1} r \sum_{n,l=1}^{\infty} \frac{1}{n^2} \operatorname{sen} \left(2\pi n \frac{2r+1}{p}\right) \left(1 + \frac{4\pi^2 aln}{\beta p}\right) e^{-4\pi^2 aln/\beta p}.$$

$$d_{2k+1}^{(D_p^*)} = \begin{cases} (2k+1) \left\lfloor \frac{k}{p} \right\rfloor, & \text{si } k \text{ es impar,} \\ (2k+1) \left(\left\lfloor \frac{k}{p} \right\rfloor + 1 \right), & \text{si } k \text{ es par,} \end{cases}$$

$$d_{2k+1}^{(D_p^*)} = \frac{1}{2} d_{2k+1}^{(Z_{2p})} + \frac{(-1)^k}{2} (2k+1)$$

$$\zeta_{S^3/D_p^*}(s) = \frac{1}{2} \zeta_{S^3/Z_{2p}}(s) + \delta\zeta(s),$$

$$\delta\zeta(s) = \frac{(\mu a)^{2s}}{2} \sum_{l \in \mathbb{Z}} \sum_{k=0}^{\infty} (-1)^k (2k+1) \left[(2k+1)^2 + \left(\frac{2\pi al}{\beta}\right)^2 \right]^{-s}.$$

$$\delta\zeta(s) = \frac{(\mu a)^{2s}}{2} \sum_{l \in \mathbb{Z}} \left\{ \sum_{k=0}^{\infty} (4k+1) \left[(4k+1)^2 + \left(\frac{2\pi al}{\beta}\right)^2 \right]^{-s} - \sum_{k=0}^{\infty} (4k+3) \left[(4k+3)^2 + \left(\frac{2\pi al}{\beta}\right)^2 \right]^{-s} \right\}$$

$$\delta\zeta^{l=0}(s) = \frac{(\mu a)^{2s} 4^{1-2s}}{2} \left[\zeta_H \left(2s-1, \frac{1}{4}\right) - \zeta_H \left(2s-1, \frac{3}{4}\right) \right].$$

$$\delta\zeta^{l\neq 0}(s) = (\mu a)^{2s} \sum_{l=1}^{\infty} \sum_{k \in \mathbb{Z}} (4k+1) \left[(4k+1)^2 + \left(\frac{2\pi al}{\beta}\right)^2 \right]^{-s}$$

$$\delta\zeta^{l\neq 0}(s) = \frac{(\mu a)^{2s}}{2(1-s)} \sum_{l=1}^{\infty} \sum_{k \in \mathbb{Z}} \frac{d}{d\alpha} \left[(4k+\alpha)^2 + \left(\frac{2\pi al}{\beta}\right)^2 \right]^{-s+1} \Big|_{\alpha=1}$$

$$\delta\zeta^{l\neq 0}(s) = -\frac{\sqrt{\pi}(\mu a)^{2s}}{8\Gamma(s)} \frac{d}{d\alpha} \sum_{l=1}^{\infty} \sum_{k \in \mathbb{Z}} e^{i\frac{\pi}{2} \alpha k} \int_0^{\infty} dt t^{s-\frac{3}{2}-1} e^{-t(2\pi al/\beta)^2 - (\pi k)^2/16t} \Big|_{\alpha=1}$$

$$\delta\zeta^{l\neq 0}(s) = \frac{\pi^{\frac{3}{2}}(\mu a)^{2s}}{4\Gamma(s)} \sum_{\substack{k,l=1 \\ k \text{ impar}}}^{\infty} (-1)^{\lfloor \frac{k}{2} \rfloor} k \left(\frac{8al}{k\beta}\right)^{\frac{3}{2}-s} K_{\frac{3}{2}-s}(\pi^2 kla/\beta)$$

$$\begin{aligned}
\Gamma_{S^3/D_p^*} &= -\frac{3\zeta_R(3)}{16\pi^2 p} + \frac{1}{2\pi} \sum_{r=1}^{p-1} r \sum_{n=1}^{\infty} \frac{1}{n^2} \operatorname{sen} \left(2\pi n \frac{2r+1}{2p} \right) - \frac{G}{\pi} \\
&\quad - \frac{\pi^4}{180p} \left(\frac{a}{\beta} \right)^3 - \frac{2}{\pi} \sum_{\substack{k,l=1 \\ k \text{ impar}}}^{\infty} (-1)^{\lfloor \frac{k}{2} \rfloor} \frac{1}{k^2} \left(1 + \frac{\pi^2 akl}{\beta} \right) e^{-\pi^2 kla/\beta} \\
&\quad + \frac{1}{2\pi^2 p} \sum_{k,l=1}^{\infty} \frac{1}{k^3} \left[2 + 2 \frac{4\pi^2 akl}{\beta} + \left(\frac{4\pi^2 kla}{\beta} \right)^2 \right] e^{-4\pi^2 akl/\beta} \\
&\quad - \frac{1}{2\pi^2 p} \sum_{k,l=1}^{\infty} \frac{1}{k^3} \left[2 + 2 \frac{2\pi^2 akl}{\beta} + \left(\frac{2\pi^2 kla}{\beta} \right)^2 \right] e^{-2\pi^2 akl/\beta} \\
&\quad + \frac{1}{\pi p} \sum_{r=1}^{p-1} r \sum_{n,l=1}^{\infty} \frac{1}{n^2} \operatorname{sen} \left(2\pi n \frac{2r+1}{2p} \right) \left(1 + \frac{2\pi^2 aln}{\beta p} \right) e^{-2\pi^2 tna/\beta p} \\
d_{2n+1}^{(T^*)} &= d_{2n+1}^{(Z_6)} + \frac{1}{2} d_{2n+1}^{(Z_4)} - \frac{1}{2} d_{2n+1}^{(Z_2)} \\
d_{2n+1}^{(O^*)} &= \frac{1}{2} d_{2n+1}^{(Z_8)} + \frac{1}{2} d_{2n+1}^{(Z_6)} + \frac{1}{2} d_{2n+1}^{(Z_4)} - \frac{1}{2} d_{2n+1}^{(Z_2)} \\
d_{2n+1}^{(I^*)} &= \frac{1}{2} d_{2n+1}^{(Z_{10})} + \frac{1}{2} d_{2n+1}^{(Z_6)} + \frac{1}{2} d_{2n+1}^{(Z_4)} - \frac{1}{2} d_{2n+1}^{(Z_2)} \\
\Gamma &= \beta E_0 + \sum_{k=1}^{\infty} d_k \log(1 - e^{-\beta \lambda_k}) \\
\zeta_{S^3/H}(s) &= \left(\frac{\mu\beta}{2\pi} \right)^{2s} \sum_{l \in \mathbb{Z}} \sum_{k=1}^{\infty} d_k \left[\left(\frac{\lambda_k \beta}{2\pi} \right)^2 + l^2 \right]^{-s} \\
\zeta_{S^3/H}(s) &= \frac{1}{\Gamma(s)} \left(\frac{\mu\beta}{2\pi} \right)^{2s} \sum_{l \in \mathbb{Z}} \sum_{k=1}^{\infty} d_k \int_0^{\infty} dt t^{s-1} e^{-t[l^2 + (\lambda_k \beta / 2\pi)^2]} \\
\zeta_{S^3/H}^{l=0}(s) &= \frac{\beta \mu^{2s}}{2\sqrt{\pi}} \frac{\Gamma(s - \frac{1}{2})}{\Gamma(s)} \sum_{k=1}^{\infty} d_k \lambda_k^{1-2s}, \\
\zeta_{S^3/H}^{l \neq 0}(s) &= \frac{2\sqrt{\pi}}{\Gamma(s)} \left(\frac{2\pi}{\mu\beta} \right)^{-2s} \sum_{k=0}^{\infty} d_k \sum_{l=1}^{\infty} 2 \left(\frac{2\pi^2 l}{\lambda_k \beta} \right)^{s-\frac{1}{2}} K_{\frac{1}{2}-s}(\lambda_k \beta l). \\
E_0 &= \frac{1}{2} \lim_{s \rightarrow 0} \sum_{k=1}^{\infty} d_k (\lambda_k^2)^{\frac{1}{2}-s}, \\
\Gamma_{S^3/H} &= \beta E_0 - 2\sqrt{\pi} \sum_{k=1}^{\infty} d_k \sum_{l=1}^{\infty} \left(\frac{2\pi^2 l}{\lambda_k \beta} \right)^{-1/2} K_{\frac{1}{2}}(\lambda_k \beta l). \\
E_{0,S^3/Z_{2q+1}} &= \frac{E_{0,S^3}}{2q+1} + \delta E_{0,2q+1}, \\
\delta E_{0,2q+1} &= \frac{1}{2} \lim_{s \rightarrow 0} a^{2s-1} \sum_{n=0}^{\infty} \left\{ \sum_{r=0}^{q-1} \left(1 - \frac{2r+1}{2q+1} \right) [n(2q+1) + 2r+1]^{2-2s} \right. \\
&\quad \left. - \sum_{r=0}^q \frac{2r}{2q+1} [n(2q+1) + 2r]^{2-2s} \right\} \\
\delta E_{0,2q+1} &= \frac{2q+1}{a\pi^3} \sum_{r=1}^q r \sum_{n=1}^{\infty} \frac{1}{n^3} \operatorname{sen} \left(\frac{4\pi nr}{2q+1} \right).
\end{aligned}$$

$$\begin{aligned}
E_{0,S^3/Z_{2q+1}} &= -\frac{(2q+1)^4 + 10(2q+1)^2 - 14}{720a(2q+1)}. \\
E_{0,S^3/Z_2} &= \frac{1}{2} \lim_{s \rightarrow 0} a^{2s-1} \sum_{k=0}^{\infty} (2k+1)^{3-2s} = -\frac{7}{240a} \\
E_{0,S^3/Z_{2q}} &= \frac{E_{0,S^3/Z_2}}{q} + \delta E_{0,2q} \\
\delta E_{0,2q} &= \frac{1}{2} \lim_{s \rightarrow 0} a^{2s-1} \sum_{r=0}^{q-1} [(q-r-1) - r] \sum_{n=0}^{\infty} (2nq+2r+1)^{2-2s} \\
E_{0,S^3/Z_{2q}} &= -\frac{(2q)^4 + 10(2q)^2 - 14}{720a(2q)} \\
E_{0,S^3/Z_p} &= -\frac{p^4 + 10p^2 - 14}{720ap} \\
\Gamma_{S^3/Z_p} &= \frac{14 - 10p^2 - p^4 \beta}{720p a} + \mathcal{O}(e^{-\beta/a}) \\
E_{0,S^3/D_p^*} &= \frac{1}{2} E_{0,S^3/Z_{2p}} + \delta E_0 \\
\delta E_0 &= \frac{1}{4} \lim_{s \rightarrow 0} a^{2s-1} \sum_{k=0}^{\infty} (-1)^k (2k+1)^{2-2s} \\
E_{0,S^3/D_p^*} &= -\frac{8p^4 + 20p^2 + 180p - 7}{1440ap} \\
E_{0,S^3/\Gamma^*} &= -\frac{3761}{360 \times 24a}, E_{0,S^3/O^*} = -\frac{11321}{360 \times 48a} \\
E_{0,S^3/I^*} &= -\frac{43553}{360 \times 120a} \\
E_{S^3/H} &\sim \frac{\pi^4 a^3}{15|H|\beta^4} \\
S_{S^3/H} &\sim \frac{4\pi^4}{45|H|} \left(\frac{a}{\beta}\right)^3 + S_{0,S^3/H}, \\
S_{\text{top},S^3/H} &= \lim_{\beta \rightarrow 0} \left[S_{S^3/H} - \frac{1}{|H|} S_{S^3} \right]. \\
S &= \frac{1}{2} \int_{\mathbb{R} \times S^3/H} d^4x \sqrt{g} \left\{ g^{\mu\nu} \partial_\mu \phi \partial_\nu \phi + \frac{R}{6} \phi^2 + \left[m^2 + \left(\xi - \frac{1}{6} \right) R \right] \phi^2 \right\}, \\
\zeta_{S^3/H}(s) &= \mu^{2s} \sum_{k=1}^{\infty} \sum_{l \in \mathbb{Z}} d_k^{(H)} \left[\left(\frac{k}{a} \right)^2 + m^2 + \left(\frac{2\pi l}{\beta} \right)^2 \right]^{-s}, \\
\zeta_{S^3}(s) &= \left(\frac{\mu\beta}{2\pi} \right)^{2s} \sum_{k=1}^{\infty} \sum_{l \in \mathbb{Z}} k^2 \left[\left(\frac{k\beta}{2\pi a} \right)^2 + \left(\frac{\beta m}{2\pi} \right)^2 + l^2 \right]^{-s}, \\
\zeta_{S^3}(s) &= \left(\frac{\mu\beta}{2\pi} \right)^{2s} \frac{\pi^{\frac{1}{2}}}{\Gamma(s)} \sum_{k=1}^{\infty} \sum_{l \in \mathbb{Z}} k^2 \int_0^\infty dt t^{s-\frac{1}{2}-1} e^{-\left[\left(\frac{k\beta}{2\pi a} \right)^2 + \left(\frac{\beta m}{2\pi} \right)^2 \right] t - (\pi l)^2/t} \\
\zeta_{S^3}(s) &= \left(\frac{\mu\beta}{2\pi} \right)^{2s} \frac{\pi^{\frac{1}{2}}}{\Gamma(s)} \left\{ \Gamma\left(s - \frac{1}{2}\right) \sum_{k=1}^{\infty} k^2 \left[\left(\frac{k\beta}{2\pi a} \right)^2 + \left(\frac{\beta m}{2\pi} \right)^2 \right]^{\frac{1}{2}-s} \right. \\
&\quad \left. + 4 \sum_{k,l=1}^{\infty} k^2 (\pi l)^{s-\frac{1}{2}} \left[\left(\frac{k\beta}{2\pi a} \right)^2 + \left(\frac{\beta m}{2\pi} \right)^2 \right]^{\frac{1}{4}-\frac{s}{2}} K_{\frac{1}{2}-s} \left(\sqrt{k^2 + (am)^2} \beta l/a \right) \right\}
\end{aligned}$$

$$\begin{aligned}
\Gamma_{S^3}^{l \neq 0}(\beta) &= \sum_{k=1}^{\infty} k^2 \log \left(1 - e^{-\frac{\beta}{a} \sqrt{k^2 + (am)^2}} \right) \\
\zeta_{S^3}^{l=0}(s) &= \frac{(\mu a)^{2s} \beta \Gamma\left(s - \frac{1}{2}\right)}{4\pi^{\frac{1}{2}} a \Gamma(s)} \sum_{k \in \mathbb{Z}} \left\{ [k^2 + (am)^2]^{\frac{3}{2}-s} - (am)^2 [k^2 + (am)^2]^{\frac{1}{2}-s} \right\} \\
\zeta_{S^3}^{l=0}(s) &= \frac{(\mu a)^{2s} \beta}{\Gamma(s)} \frac{1}{4a} \left\{ \frac{1}{2} (am)^{4-2s} \Gamma(s-2) \right. \\
&\quad + 4 \left(s - \frac{3}{2} \right) \sum_{k=1}^{\infty} \left(\frac{\pi k}{am} \right)^{s-2} K_{s-2}(2\pi am k) \\
&\quad \left. - 4(am)^2 \sum_{k=1}^{\infty} \left(\frac{\pi k}{am} \right)^{s-1} K_{s-1}(2\pi am k) \right\} \\
\Gamma_{S^3}^{l=0}(\beta) &= -\frac{(am)^4 \beta}{32} \frac{1}{a} \left(\frac{3}{2} + 2 \log \frac{\mu}{m} \right) \\
&\quad + \frac{(am)^3 \beta}{2\pi} \frac{1}{a} \sum_{k=1}^{\infty} \frac{1}{k} K_1(2\pi am k) + \frac{3(am)^2 \beta}{4\pi^2} \frac{1}{a} \sum_{k=1}^{\infty} \frac{1}{k^2} K_2(2\pi am k). \\
\Gamma_{S^3}(\beta) &= \frac{3(am)^2 \beta}{4\pi^2} \frac{1}{a} \sum_{k=1}^{\infty} \frac{1}{k^2} K_2(2\pi am k) + \frac{(am)^3 \beta}{2\pi} \frac{1}{a} \sum_{k=1}^{\infty} \frac{1}{k} K_1(2\pi am k) \\
&\quad + \sum_{k=1}^{\infty} k^2 \log \left(1 - e^{-\frac{\beta}{a} \sqrt{k^2 + (am)^2}} \right) \\
E_0 &= \frac{3(am)^2}{4\pi^2 a} \sum_{k=1}^{\infty} \frac{1}{k^2} K_2(2\pi am k) + \frac{(am)^3}{2\pi a} \sum_{k=1}^{\infty} \frac{1}{k} K_1(2\pi am k) \\
\zeta_{S^3}^{l=0}(s) &= (\mu a)^{2s} \frac{\beta}{2\pi^{\frac{1}{2}} a} \frac{\Gamma\left(s - \frac{1}{2}\right)}{\Gamma(s)} \sum_{k=1}^{\infty} k^{3-2s} \left[1 + \left(\frac{am}{k} \right)^2 \right]^{\frac{1}{2}-s} \\
\zeta_{S^3}^{l=0}(s) &= (\mu a)^{2s} \frac{\beta}{2\pi^{\frac{1}{2}} a} \frac{\Gamma\left(s - \frac{1}{2}\right)}{\Gamma(s)} \sum_{n=0}^{\infty} \frac{\Gamma\left(\frac{3}{2} - s\right) (am)^{2n}}{n! \Gamma\left(\frac{3}{2} - n - s\right)} \zeta_R(2s + 2n - 3) \\
\Gamma_{S^3}^{l=0} &= -\frac{(am)^4 \beta}{16} \frac{1}{a} [\log(\mu a/2) + 1 + \gamma] + \frac{\beta}{2a} \sum_{\substack{n=0 \\ n \neq 2}}^{\infty} \frac{\Gamma\left(\frac{3}{2}\right) (am)^{2n}}{n! \Gamma\left(\frac{3}{2} - n\right)} \zeta_R(2n - 3) \\
\Gamma_{S^3}(\beta) &= \frac{\beta}{240a} - \frac{(am)^2 \beta}{48} \frac{1}{a} - \frac{(am)^4 \beta}{16} \frac{1}{a} \left[\log\left(\frac{ma}{2}\right) + \frac{1}{4} + \gamma \right] \\
&\quad + \frac{\beta}{2a} \sum_{n=3}^{\infty} \frac{\Gamma\left(\frac{3}{2}\right) (am)^{2n}}{n! \Gamma\left(\frac{3}{2} - n\right)} \zeta_R(2n - 3) + \sum_{k=1}^{\infty} k^2 \log \left(1 - e^{-\frac{\beta}{a} \sqrt{k^2 + (am)^2}} \right) \\
\zeta_{S^3}(s) &= \frac{\mu^{2s} a^2}{2} \sum_{k, l \in \mathbb{Z}} \left[\left(\frac{k}{a} \right)^2 + m^2 + \left(\frac{2\pi l}{\beta} \right)^2 \right]^{-s+1} \\
&\quad - \frac{\mu^{2s} a^2}{2} \sum_{k, l \in \mathbb{Z}} \left[m^2 + \left(\frac{2\pi l}{\beta} \right)^2 \right] \left[\left(\frac{k}{a} \right)^2 + m^2 + \left(\frac{2\pi l}{\beta} \right)^2 \right]^{-s},
\end{aligned}$$

$$\begin{aligned}
\zeta_{S^3}(s) &= \frac{\pi^{\frac{1}{2}}(\mu a)^{2s}}{2} \sum_{l \in \mathbb{Z}} \left[(a^2 \lambda_{l,0})^{-s+\frac{3}{2}} \frac{\Gamma(s-\frac{3}{2})}{2\Gamma(s)} \right. \\
&+ \frac{4}{\Gamma(s-1)} \sum_{k=1}^{\infty} (k\pi)^{s-\frac{3}{2}} (a^2 \lambda_{l,0})^{-\frac{s+3}{2}} K_{s-\frac{3}{2}}(2k\pi \sqrt{a^2 \lambda_{l,0}}) \\
&\left. - \frac{4}{\Gamma(s)} \sum_{k=1}^{\infty} (k\pi)^{s-\frac{1}{2}} (a^2 \lambda_{l,0})^{-\frac{s+5}{2}} K_{s-\frac{1}{2}}(2k\pi \sqrt{a^2 \lambda_{l,0}}) \right] \\
\zeta_{S^3}^{k=0}(s) &= (\mu a)^{2s} \frac{\pi^{\frac{1}{2}} \Gamma(s-\frac{3}{2})}{4 \Gamma(s)} [(am)^{3-2s} \\
&+ 2 \left(\frac{2\pi a}{\beta} \right)^{3-2s} \sum_{n=0}^{\infty} \frac{\Gamma(\frac{5}{2}-s)}{n! \Gamma(\frac{5}{2}-n-s)} \left(\frac{m\beta}{2\pi} \right)^{2n} \zeta_R(2s+2n-3) \Big] \\
\Gamma_{S^3}^{k=0}(\beta) &= -\frac{\pi}{6} (am)^3 - \frac{\pi}{3} \left(\frac{2\pi a}{\beta} \right)^3 \sum_{\substack{n=0 \\ n \neq 2}}^{\infty} \frac{\Gamma(\frac{5}{2})}{n! \Gamma(\frac{5}{2}-n)} \left(\frac{m\beta}{2\pi} \right)^{2n} \zeta_R(2n-3) \\
&- \frac{(am)^4 \beta}{32 a} [2 \log(\mu\beta/2\pi) + 2\gamma - 2 \log 2] \\
\Gamma_{S^3} &= -\frac{\pi}{6} (am)^3 - \frac{\pi^4}{45} \left(\frac{a}{\beta} \right)^3 + \frac{\pi^2}{12} (am)^2 \frac{a}{\beta} \\
&+ \frac{1}{4\pi^2} \sum_{k=1}^{\infty} \frac{1}{k^3} \left[1 + 2\pi m a k + \frac{1}{2} (2\pi m a k)^2 \right] e^{-2\pi m a k} + \mathcal{O}(m\beta) \\
\zeta_{S^3}^{k=0}(s) &= \frac{\pi(\mu a)^{2s}}{4\Gamma(s)} \left(\frac{2\pi a}{\beta} \right)^{3-2s} \left[\left(\frac{m\beta}{2\pi} \right)^{4-2s} \Gamma(s-2) + 4 \sum_{l=1}^{\infty} \left(\frac{2\pi^2 l}{m\beta} \right)^{s-2} K_{2-s}(m\beta l) \right] \\
\Gamma_{S^3}^{k=0}(\beta) &= -\frac{(am)^4 \beta}{32 a} \left[\frac{3}{2} + 2 \log \left(\frac{\mu}{m} \right) \right] - (am)^2 \frac{a}{\beta} \sum_{l=1}^{\infty} \frac{1}{l^2} K_2(m\beta l). \\
\Gamma_{S^3}(\beta) &= -(am)^2 \frac{a}{\beta} \sum_{l=1}^{\infty} \frac{1}{l^2} K_2(m\beta l) + \mathcal{O}(e^{-a/\beta}) \\
\zeta_{S^3/Z_{2q+1}}(s) &= \mu^{2s} \sum_{l \in \mathbb{Z}} \sum_{n=0}^{\infty} \left[\sum_{\substack{r=0 \\ r \text{ par}}}^{p-1} (np+r)n + \sum_{\substack{r=1 \\ r \text{ impar}}}^{p-2} (np+r)(n+1) \right] \lambda_{l,np+r}^{-s} \\
&= \frac{1}{p} \zeta_{S^3}(s) + \delta \zeta_{2q+1}(s) \\
\delta \zeta_{2q+1}(s) &= -\frac{\mu^{2s}}{p} \sum_{l \in \mathbb{Z}} \sum_{n=0}^{\infty} \left\{ \sum_{r=1}^q 2r(np+2r) \lambda_{l,np+2r}^{-s} \right. \\
&\quad \left. + \sum_{r=0}^{q-1} 2(r-q)(np+2r+1) \lambda_{l,np+2r+1}^{-s} \right\}
\end{aligned}$$

$$\begin{aligned}
\delta\zeta_{2q+1}(s) &= -2\left(\frac{\mu a}{p}\right)^{2s} \sum_{l \in \mathbb{Z}} \sum_{n=0}^{\infty} \sum_{r=1}^q r \\
&\quad \times \left\{ \left(n + \frac{2r}{p}\right) \left[\left(n + \frac{2r}{p}\right)^2 + \left(\frac{am}{p}\right)^2 + \left(\frac{2\pi al}{p\beta}\right)^2 \right]^{-s} \right. \\
&\quad \left. - \left(n + 1 - \frac{2r}{p}\right) \left[\left(n + 1 - \frac{2r}{p}\right)^2 + \left(\frac{am}{p}\right)^2 + \left(\frac{2\pi al}{p\beta}\right)^2 \right]^{-s} \right\} \\
\delta\zeta_{2q+1}^{l=0}(s) &= -2\sqrt{\pi} \frac{\Gamma\left(s - \frac{1}{2}\right)}{\Gamma(s)} \left(\frac{\mu a}{p}\right)^{2s} \sum_{n=0}^{\infty} \sum_{r=1}^q r \left\{ \left(n + \frac{2r}{p}\right) \left[\left(n + \frac{2r}{p}\right)^2 + \left(\frac{am}{p}\right)^2 \right]^{\frac{1}{2}-s} \right. \\
&\quad \left. - \left(n + 1 - \frac{2r}{p}\right) \left[\left(n + 1 - \frac{2r}{p}\right)^2 + \left(\frac{am}{p}\right)^2 \right]^{\frac{1}{2}-s} \right\} \\
\delta\zeta_{2q+1}^{l=0}(s) &= \sqrt{\pi} \frac{\Gamma\left(s - \frac{3}{2}\right)}{\Gamma(s)} \left(\frac{\mu a}{p}\right)^{2s} \sum_{n \in \mathbb{Z}} \sum_{r=1}^q r \frac{d}{d\alpha} \left[\left(n + \frac{2r}{p} + \alpha\right)^2 + \left(\frac{am}{p}\right)^2 \right]^{\frac{3}{2}-s} \Bigg|_{\alpha=0} \\
\delta\zeta_{2q+1}^{l=0}(s) &= \frac{8(am)^2}{p^2\Gamma(s)} \left(\frac{\pi\mu^2}{pm^2}\right)^s \sum_{n=1}^{\infty} \frac{1}{n^{1-s}} \sum_{r=1}^q r \operatorname{sen}(4\pi nr/p) K_{2-s}(2\pi amn/p) \\
\delta\zeta_{2q+1}^{l=0}(s) &= \frac{1}{s-1} \left(\frac{\mu a}{p}\right)^{2s} \sum_{n \in \mathbb{Z}} \sum_{r=1}^q r \frac{d}{d\alpha} \left[\left(n + \frac{2r}{p} + \alpha\right)^2 + \left(\frac{am}{p}\right)^2 \right]^{-s+1} \Bigg|_{\alpha=0} \\
\delta\zeta_{2q+1}^{l=0}(s) &= \frac{-8}{\Gamma(s)} \left(\frac{am}{p}\right)^{\frac{3}{2}} \left(\frac{\pi\mu^2 a}{pm}\right)^s \sum_{r=1}^q r \sum_{n=1}^{\infty} n^{s-\frac{1}{2}} \operatorname{sen}(4\pi rn/p) K_{\frac{3}{2}-s}(2\pi amn/p) \\
\delta\Gamma_{S^3/Z_{p=2q+1}}^{l=0} &= \frac{1}{\pi} \sum_{r=1}^q r \sum_{n=1}^{\infty} \frac{1}{n^2} \left(\frac{2\pi amn}{p} + 1\right) \operatorname{sen}(4\pi rn/p) e^{-2\pi amn/p} \\
\zeta_{S^3/Z_2}(s) &= \mu^{2s} \sum_{l \in \mathbb{Z}} \left\{ \sum_{k=1}^{\infty} k^2 \left[\left(\frac{k}{a}\right)^2 + m^2 + \left(\frac{2\pi l}{\beta}\right)^2 \right]^{-s} \right. \\
&\quad \left. - \sum_{k=1}^{\infty} (2k)^2 \left[\left(\frac{2k}{a}\right)^2 + m^2 + \left(\frac{2\pi l}{\beta}\right)^2 \right]^{-s} \right\} \\
&= \zeta_{S^3}(s; \beta, a, m) - 2^{2-2s} \zeta_{S^3}(s; 2\beta, a, m/2), \\
\Gamma_{S^3/Z_2}(\beta, a, m) &= \Gamma_{S^3}(\beta, a, m) - 4\Gamma_{S^3}(2\beta, a, m/2) \\
\zeta_{S^3/Z_{2q}}(s) &= \frac{1}{q} \zeta_{S^3/Z_2}(s) + \delta\zeta_{2q}(s), \\
\delta\zeta_{2q}(s) &= \mu^{2s} \sum_{n=0}^{\infty} \sum_{r=0}^{q-1} \left(\frac{q-r-1}{q} - \frac{r}{q}\right) (2nq + 2r + 1) \\
&\quad \times \sum_{l \in \mathbb{Z}} \left[\left(\frac{2nq + 2r + 1}{a}\right)^2 + m^2 + \left(\frac{2\pi l}{\beta}\right)^2 \right]^{-s}
\end{aligned}$$

$$\begin{aligned}
\delta\zeta_{2q}(s) &= 2q \left(\frac{\mu a}{2q}\right)^{2s} \sum_{l \in \mathbb{Z}} \sum_{n=0}^{\infty} \left\{ \sum_{r=0}^{q-1} \frac{r}{q} \left(n+1 - \frac{2r+1}{2q}\right) \right. \\
&\quad \times \left[\left(n+1 - \frac{2r+1}{2q}\right)^2 + \left(\frac{am}{2q}\right)^2 + \left(\frac{2\pi al}{2q\beta}\right)^2 \right]^{-s} \\
&\quad \left. - \sum_{r=0}^{q-1} \frac{r}{q} \left(n + \frac{2r+1}{2q}\right) \left[\left(n + \frac{2r+1}{2q}\right)^2 + \left(\frac{am}{2q}\right)^2 + \left(\frac{2\pi al}{2q\beta}\right)^2 \right]^{-s} \right\} \\
\delta\zeta_{2q}^{l=0}(s) &= \sqrt{\pi} \frac{\Gamma\left(s - \frac{1}{2}\right)}{\Gamma(s)} \left(\frac{\mu a}{2q}\right)^{2s} \sum_{n \in \mathbb{Z}} \sum_{r=1}^{q-1} r \frac{d}{d\alpha} \left[\left(n + \frac{2r+1}{2q} + \alpha\right)^2 + \left(\frac{am}{2q}\right)^2 \right]^{\frac{1}{2}-s} \Bigg|_{\alpha=0} \\
\delta\zeta_{2q}^{l=0}(s) &= \frac{2\pi am}{\Gamma(s)} \left(\frac{\pi a \mu^2}{2qm}\right)^s \sum_{r=1}^{q-1} r \sum_{n=1}^{\infty} n^s \operatorname{sen}\left(2\pi n \frac{2r+1}{2q}\right) K_{1-s}(\pi amn/q) \\
\delta\zeta_{2q}^{l=0}(s) &= \frac{2}{s-1} \left(\frac{\mu a}{2q}\right)^{2s} \sum_{n \in \mathbb{Z}} \sum_{r=1}^{q-1} r \frac{d}{d\alpha} \left[\left(n + \frac{2r+1}{2q} + \alpha\right)^2 + \left(\frac{am}{2q}\right)^2 \right]^{-s+1} \Bigg|_{\alpha=0} \\
\delta\zeta_{2q}^{l=0}(s) &= \frac{-8}{\Gamma(s)} \left(\frac{am}{2q}\right)^{\frac{3}{2}} \left(\frac{\pi \mu^2 a}{2qm}\right)^s \sum_{r=1}^{q-1} r \sum_{n=1}^{\infty} n^{s-\frac{1}{2}} \operatorname{sen}\left(2\pi n \frac{2r+1}{2q}\right) K_{\frac{3}{2}-s}(\pi amn/q) \\
\delta\Gamma_{S^3/Z_{2q}}^{l=0} &= \frac{1}{\pi} \sum_{r=1}^{q-1} r \sum_{n=1}^{\infty} \frac{1}{n^2} \left(\frac{2\pi amn}{2q} + 1\right) \operatorname{sen}\left(2\pi n \frac{2r+1}{2q}\right) e^{-\pi amn/q} \\
\zeta_{S^3/D_p^*}(s) &= \frac{1}{2} \zeta_{S^3/Z_{2p}}(s) + \delta\zeta(s) \\
\delta\zeta(s) &= \frac{\mu^{2s}}{2} \sum_{l \in \mathbb{Z}} \sum_{k=0}^{\infty} [(4k+1)\lambda_{l,4k+1}^{-s} - (4k+3)\lambda_{l,4k+3}^{-s}] \\
\delta\zeta(s) &= \frac{(\mu a)^{2s}}{4(-s+1)} \sum_{k,l \in \mathbb{Z}} \frac{d}{d\alpha} \left[\left(\frac{2\pi al}{\beta}\right)^2 + (am)^2 + (4k+1)^2 \right]^{-s+1} \Bigg|_{\alpha=1} \\
\delta\zeta^{l=0}(s) &= -\frac{2(am)^2}{\Gamma(s)} \left(\frac{\pi \mu^2 a}{4m}\right)^s \sum_{k=1}^{\infty} \frac{1}{k^{1-s}} \operatorname{sen}(\pi k/2) K_{2-s}(\pi amk/2) \\
\delta\zeta^{l=0}(s) &= \frac{\pi^{\frac{3}{2}}(\mu a)^{2s}}{16\Gamma(s)} \sum_{k=1}^{\infty} k \operatorname{sen}(\pi k/2) \left(\frac{4am}{\pi k}\right)^{\frac{3}{2}-s} K_{\frac{3}{2}-s}(\pi amk/2). \\
\delta\Gamma_{S^3/D_p^*}^{l=0} &= -\frac{1}{\pi} \sum_{k=1}^{\infty} \operatorname{sen}(\pi k/2) \left(1 + \frac{\pi amk}{2}\right) e^{-\pi amk/2}, \\
\Gamma &= \beta E_0^{(H)} + \sum_{k=1}^{\infty} d_k^{(H)} \log \left(1 - e^{-\beta \sqrt{\lambda_k^2 + m^2}}\right), \\
S_{S^3/H} &= \sum_{k=1}^{\infty} d_k^{(H)} \left[\frac{\frac{\beta}{a} \sqrt{k^2 + (am)^2}}{e^{\frac{\beta}{a} \sqrt{k^2 + (am)^2}} - 1} - \log \left(1 - e^{-\frac{\beta}{a} \sqrt{k^2 + (am)^2}}\right) \right], \\
S_{S^3/H} &\sim \frac{4\pi^4}{45|H|} \left(\frac{a}{\beta}\right)^3 - \frac{\pi^2}{6|H|} (am)^2 \frac{a}{\beta} + S_{0,S^3/H}(am), \\
\langle T_{\mu}^{\mu} \rangle &= -\frac{cR}{12} \\
\langle T_{\mu}^{\mu} \rangle &= cW_{\mu\nu\rho\sigma} W^{\mu\nu\rho\sigma} - aE_4
\end{aligned}$$

$$\begin{aligned}
W^2 &= \frac{1}{16\pi^2} \left(R^{\mu\nu\rho\sigma} R_{\mu\nu\rho\sigma} - 2R^{\mu\nu} R_{\mu\nu} + \frac{1}{3} R^2 \right) \\
E_4 &= \frac{1}{16\pi^2} \left(R^{\mu\nu\rho\sigma} R_{\mu\nu\rho\sigma} - 4R^{\mu\nu} R_{\mu\nu} + R^2 \right) \\
\langle T_\mu^\mu \rangle &= \sum_i c_i I_i - (-1)^{d/2} a E_d \\
a &= \frac{(-1)^{d/2}}{d} \frac{dF}{d \log r} \\
\frac{dg^i}{dt} &= -\mu \frac{dg^i(\mu)}{d\mu} =: -\beta^i(g(\mu)), \\
\frac{\partial}{\partial t} &= -\beta^i(g) \frac{\partial}{\partial g_i} \\
\dot{C}(g(t)) &= -\beta^i(g) \frac{\partial}{\partial g_i} C(g(t)) \leq 0, \\
\dot{C}(g(t)) &= -G_{ij}(g) \beta^i(g) \beta^j(g), \\
S_{\text{ent}} &= \lim_{q \rightarrow 1} S_q := \lim_{q \rightarrow 1} \frac{q \Gamma_{S^3} - \Gamma_{qS^3}}{1 - q} \\
S_{\text{ren}} &= am \frac{dS_{\text{ent}}}{d(am)} - S_{\text{ent}} \\
\tilde{F} &= \Gamma - \frac{n}{3} g \frac{d\Gamma}{dg} \\
S_{\text{hol}, S^3/H} &:= \lim_{\beta \rightarrow 0} \left[S_{S^3/H}(am, \beta/a) - \frac{1}{|H|} S_{S^3}(am, \beta/a) \right] \\
S_{\text{hol}, S^3/H} &= \frac{1}{|H|} \Gamma_{S^3} - \Gamma_{S^3/H} \\
S_{\text{hol}, S^3/Z_{2q+1}} &= -\frac{1}{\pi} \sum_{r=1}^q r \sum_{n=1}^{\infty} \frac{1}{n^2} \left(\frac{2\pi amn}{2q+1} + 1 \right) \text{sen} \left(\frac{4\pi rn}{2q+1} \right) e^{-\frac{2\pi amn}{2q+1}}, \\
S_{\text{hol}, S^3/Z_{2q}} &= -\frac{1}{8\pi^2 q} \sum_{k=1}^{\infty} \frac{1}{k^3} \left[1 + 2\pi amk + \frac{1}{2} (2\pi amk)^2 \right] e^{-2\pi amk} \\
&\quad + \frac{1}{\pi^2 q} \sum_{k=1}^{\infty} \frac{1}{k^3} \left[1 + \pi mak + \frac{1}{2} (\pi amk)^2 \right] e^{-\pi amk} \\
&\quad - \frac{1}{\pi} \sum_{r=1}^{q-1} r \sum_{n=1}^{\infty} \frac{1}{n^2} \left(\frac{2\pi amn}{2q} + 1 \right) \text{sen} \left(2\pi \frac{2r+1}{2q} n \right) e^{-\frac{2\pi amn}{2q}} \\
S_{\text{hol}, S^3/D_p^*} &= \frac{1}{2} S_{\text{hol}, S^3/Z_{2q}} + \frac{1}{\pi} \sum_{\substack{k=1 \\ k \text{ impar}}}^{\infty} \frac{(-1)^{\lfloor k/2 \rfloor}}{k^2} \left(1 + \frac{\pi amk}{2} \right) e^{-\frac{\pi amk}{2}}. \\
\zeta_{S^3}(s) &= (\mu a)^{2s} \sum_{k=1}^{\infty} k^2 \left[k^2 - \frac{1}{4} + (am)^2 \right]^{-s}. \\
\zeta_{S^3}(s) &= (\mu a)^{2s} \sum_{n=0}^{\infty} \frac{\Gamma(-s+1)}{n! \Gamma(-s+1-n)} \left[(am)^2 - \frac{1}{4} \right]^n \zeta_R(2s+2n-2), \\
\zeta_{S^3}(s) &= \frac{(\mu a)^{2s}}{\Gamma(s)} \sum_{n=0}^{\infty} \frac{(-1)^n \Gamma(s+n)}{n!} \left[(am)^2 - \frac{1}{4} \right]^n \zeta_R(2s+2n-2), \\
\Gamma_{S^3}(am) &= \frac{\zeta_R(3)}{4\pi^2} + \frac{1}{2} \sum_{n=0}^{\infty} \frac{(-1)^n}{n+1} \left[(am)^2 - \frac{1}{4} \right]^{n+1} \zeta_R(2n).
\end{aligned}$$

$$\begin{aligned}
\Gamma_{S^3}(0) &= \frac{\zeta_R(3)}{4\pi^2} - \frac{1}{8} \sum_{n=0}^{\infty} \frac{\zeta_R(2n)}{(n+1)2^{2n}} \\
\sum_{n=1}^{\infty} \frac{\zeta_R(2n)}{(n+1)2^{2n}} &= \frac{1}{2} + \log\left(\frac{B^{14}}{2}\right) \\
\Gamma_{S^3}(0) &= -\frac{3\zeta_R(3)}{16\pi^2} - \frac{\log 2}{8} \approx 0,0638 \\
\frac{d\Gamma_{S^3}}{d(am)^2} &= \frac{1}{2} \sum_{n=0}^{\infty} (-1)^n \left[(am)^2 - \frac{1}{4} \right]^n \zeta_R(2n) \\
&\quad \sum_{n=1}^{\infty} \frac{\zeta_R(2n)}{2^{2n-1}} = 1 \\
\zeta_{S^3/Z_2}(s) &= (\mu a)^{2s} \sum_{\substack{k=1 \\ k \text{ impar}}}^{\infty} k^2 \left[k^2 - \frac{1}{4} + (am)^2 \right]^{-s} \\
\zeta_{S^3/Z_2}(s) &= \zeta_{S^3}(s) - 4^{1-s} (\mu a)^{2s} \sum_{k=1}^{\infty} k^2 \left[k^2 - \frac{1}{16} + \frac{(am)^2}{4} \right]^{-s} \\
&= \zeta_{S^3}(s) \Big|_{am} - 2^{2-2s} \zeta_{S^3}(s) \Big|_{\tilde{m}}, \\
\Gamma_{S^3/Z_2}(am) &= \Gamma_{S^3}(am) - 4\Gamma_{S^3}(\tilde{m}) \\
S_{\text{hol},S^3/Z_2}(am) &= 4\Gamma_{S^3}(\tilde{m}) - \frac{1}{2}\Gamma_{S^3}(am) \\
\frac{d}{d(am)^2} S_{\text{hol},S^3/Z_2}(am) &= \frac{d}{d\tilde{m}^2} \Gamma_{S^3}(\tilde{m}) \\
\frac{d}{d(am)^2} S_{\text{hol},S^3/Z_2}(am) \Big|_{am=0} &= \frac{d}{d\tilde{m}^2} \Gamma_{S^3}(\tilde{m}) \Big|_{\tilde{m}^2=3/16} \\
\frac{d}{d\tilde{m}^2} \Gamma_{S^3}(\tilde{m}) \Big|_{\tilde{m}^2=3/16} &= \frac{1}{2} \sum_{n=0}^{\infty} \frac{\zeta_R(2n)}{4^{2n}} = -\frac{\pi}{16}. \\
\frac{d}{d(am)^2} S_{\text{hol},S^3/Z_2}(am) \Big|_{am=0} &= -\frac{\pi}{4} \\
\zeta_{\pm}(s) &= (\mu a)^{2s} \sum_{k=1}^{\infty} k^2 (k \pm M)^{-s} \\
\zeta_{\pm}(s) &= (\mu a)^{2s} [\zeta_H(s-2, 1 \pm M) \mp 2M\zeta_H(s-1, 1 \pm M) \\
&\quad + M^2\zeta_H(s, 1 \pm M)] \\
\Gamma_{S^3}(m) &= -\frac{1}{2} \{ \zeta'_H(-2, 1 + M) + \zeta'_H(-2, 1 - M) - 2M[\zeta'_H(-1, 1 + M) \\
&\quad - \zeta'_H(-1, 1 - M)] + M^2[\zeta'_H(0, 1 + M) + \zeta'_H(0, 1 - M)] \} \\
\partial_p S_{\text{top},S^3/Z_p}^{(4)} \Big|_{p=1} &= 2F_{S^3}^{(3)} \\
S_{\text{top},S^3/Z_p}^{(4)} &= \frac{1}{4} \int_0^{\infty} dx \Re \frac{p \sinh z + \cosh z \operatorname{senh} p z}{z \sinh^2 z \operatorname{senh} \frac{pz}{2}} \\
&\quad - \frac{1}{4p} \int_0^{\infty} dx \Re \frac{\operatorname{senh} z (1 + \cosh z)}{z \sinh^2 z \operatorname{senh}^2 \frac{z}{2}}
\end{aligned}$$

$$\begin{aligned}
\partial_p S_{\text{top}, S^3/Z_p}^{(4)} \Big|_{p=1} &= \frac{1}{4} \int_0^\infty dx \Re \frac{\sinh z}{z \sinh^2 z \sinh^2 \frac{z}{2}} + \frac{1}{4} \int_0^\infty dx \Re \frac{\sinh z (1 + \cosh z)}{z \sinh^2 z \sinh^2 \frac{z}{2}} \\
&+ \frac{1}{4} \int_0^\infty dx \Re \frac{\cosh^2 z - \coth \frac{z}{2} \sinh z (1 + \cosh z)}{\sinh^2 z \sinh^2 \frac{z}{2}} \\
&= I_1 + I_2 + I_3 \\
\cosh^2 z - \coth \frac{z}{2} \sinh z (1 + \cosh z) &= -1 - 2 \cosh z \\
I_3 &= -\frac{1}{4} \int_0^\infty dx \Re \frac{1 + 2 \cosh z}{\sinh^2 z \sinh^2 \frac{z}{2}} \\
I_3 &= -\frac{1}{4} \int_0^\infty dx \Re \frac{1 + 2 \cosh \left(x + i \frac{\pi}{2}\right)}{\sinh^2 \left(x + i \frac{\pi}{2}\right) \sinh^2 \left(\frac{x}{2} + i \frac{\pi}{4}\right)} \\
\Re \frac{1 + 2 \cosh \left(x + i \frac{\pi}{2}\right)}{\sinh^2 \left(x + i \frac{\pi}{2}\right) \sinh^2 \left(\frac{x}{2} + i \frac{\pi}{4}\right)} &= 2 \left(\frac{3}{\cosh^4 x} - \frac{2}{\cosh^2 x} \right), \\
\int_0^\infty \frac{dx}{\cosh^2 x} &= 1, \int_0^\infty \frac{dx}{\cosh^4 x} = \frac{2}{3} \\
\partial_p S_{\text{top}, S^3/Z_p}^{(4)} \Big|_{p=1} &= I_1 + I_2 = \frac{1}{4} \int_0^\infty dx \Re \frac{2 + \cosh z}{z \sinh z \sinh^2 \frac{z}{2}} \\
\partial_p S_{\text{top}, S^3/Z_p}^{(4)} \Big|_{p=1} &= \frac{1}{4} \int_0^\infty dx \frac{2\pi - 6x \sinh x - \pi \sinh^2 x}{\left(x^2 + \frac{\pi^2}{4}\right) \cosh^3 x} \\
F^{(3)} &= -\frac{1}{4} \int_0^\infty dx \Re \frac{\cosh \frac{z}{2} (1 + \cosh z)}{z \sinh^2 z \sinh^2 \frac{z}{2}} \\
F^{(3)} &= -\frac{\pi}{16} \int_0^\infty dx \frac{1 - \cosh 2x}{\left(x^2 + \frac{\pi^2}{4}\right) \cosh^3 x} \\
\frac{1}{2} \int_0^\infty dx \frac{4 - 2 \cosh^2 \frac{\pi x}{2}}{(x^2 + 1) \cosh^3 \frac{\pi x}{2}} - \frac{1}{2} \int_0^\infty dx \frac{3x \sinh \frac{\pi x}{2}}{(x^2 + 1) \cosh^3 \frac{\pi x}{2}} &= \tilde{I}_1 + \tilde{I}_2 \\
\frac{1}{1 + x^2} &= \int_0^\infty dt e^{-t} \cos tx, \frac{x}{1 + x^2} = \int_0^\infty dt e^{-t} \sin tx \\
\tilde{I}_1 &= \int_0^\infty dt e^{-t} \int_0^\infty dx \frac{2 - \cosh^2 \frac{\pi x}{2}}{\cosh^3 \frac{\pi x}{2}} \cos tx = \frac{3\zeta_R(3)}{2\pi^2} \\
\tilde{I}_2 &= -\frac{3}{2} \int_0^\infty dt e^{-t} \int_0^\infty dx \frac{\sinh \frac{\pi x}{2}}{\cosh^3 \frac{\pi x}{2}} \sin tx = -\frac{3\zeta_R(3)}{2\pi^2} \\
\Delta^{(d)} + \frac{d(d-2)}{4a^2} + m^2 & \\
a^2 \lambda_n &= \left(n + \frac{d-1}{2}\right)^2 - \frac{1}{4} + (am)^2 \quad n \in \mathbb{N} \\
d_n &= \frac{(2n + d - 1)(n + d - 2)!}{n!(d-1)!}, \\
d_n &= \frac{2n(n+k-1)!}{(2k)!(n-k)!} \\
\frac{(n+k-1)!}{(n-k)!} &= (n^2 - (k-1)^2)(n^2 - (k-2)^2) \cdots (n^2 - 1)n
\end{aligned}$$



$$\frac{(n+k-1)!}{(n-k)!} = \sum_{l=1}^k c_l^{(2k+1)} n^{2l-1},$$

$$\zeta_{S^{2k+1}}(s) = \frac{2(\mu a)^{2s}}{(2k)!} \sum_{n=1}^{\infty} \sum_{l=1}^k c_l^{(2k+1)} n^{2l} [n^2 + (am)^2 - 1/4]^{-s}$$

$$\Gamma_{S^{2k+1}} = -\frac{1}{2} \frac{d}{ds} [\zeta_+(s) + \zeta_-(s)] \Big|_{s=0}$$

$$\zeta_{\pm}(s) = \frac{2(\mu a)^{2s}}{(2k)!} \sum_{n=1}^{\infty} \sum_{l=1}^k c_l^{(2k+1)} n^{2l} (n \pm M)^{-s}$$

$$n^{2l} = \sum_{p=0}^{2l} \frac{(2l)!}{p!(2l-p)!} (n \pm M)^{2l-p} (\mp M)^p$$

$$\zeta_{\pm}(s) = \frac{2(\mu a)^{2s}}{(2k)!} \sum_{l=1}^k c_l^{(2k+1)} \sum_{p=0}^{2l} \frac{(2l)!}{p!(2l-p)!} (\mp M)^p \zeta_H(s-2l+p, 1 \pm M)$$

$$\Gamma_{S^{2k+1}}(m) = -\frac{1}{(2k)!} \sum_{l=1}^k c_l^{(2k+1)} \sum_{p=0}^{2l} \frac{(2l)!}{p!(2l-p)!} M^p$$

$$\times [\zeta'_H(p-2l, 1-M) + (-1)^p \zeta'_H(p-2l, 1+M)]$$

$$\zeta_{S^{2k+1}/Z_2}(s) = \frac{1 + (-1)^{k+1}}{2} \zeta_{S^{2k+1}}(s) + (-1)^k \Delta \zeta(s)$$

$$\zeta_{\pm}(s) = \frac{2^{1-2s} (\mu a)^{2s}}{(2k)!} \sum_{l=1}^k c_l^{(2k+1)} 2^{2l} \sum_{p=0}^{2l} \frac{(2l)!}{p!(2l-p)!} \frac{(\mp M)^p}{2^p} \zeta_H\left(s-2l+p, 1 \pm \frac{M}{2}\right)$$

$$S_{\text{hol}, S^{2k+1}/Z_2} = \frac{(-1)^{k+1}}{(2k)!} \sum_{l=1}^k c_l^{(2k+1)} \sum_{p=0}^{2l} \frac{(2l)!}{p!(2l-p)!} M^p$$

$$\times \left\{ \frac{1}{2} [\zeta'_H(p-2l, 1-M) + (-1)^p \zeta'_H(p-2l, 1+M)] \right.$$

$$\left. - 2^{2l-p} \left[\zeta'_H\left(p-2l, 1 - \frac{M}{2}\right) + (-1)^p \zeta'_H\left(p-2l, 1 + \frac{M}{2}\right) \right] \right\}$$

$$d_n = \frac{2}{(2k-1)!} \left(n + k - \frac{1}{2}\right) \frac{(n+2k-2)!}{n!}$$

$$\frac{(n+2k-2)!}{n!} = \frac{(\tilde{n} + k - \frac{3}{2})!}{(\tilde{n} - k + \frac{1}{2})!} = \prod_{l=0}^{k-2} \left[\tilde{n}^2 - \left(l + \frac{1}{2}\right)^2 \right]$$

$$\frac{(n+2k-2)!}{n!} = \sum_{l=0}^{k-1} c_l^{(2k)} \tilde{n}^{2l}$$

$$\zeta_{S^{2k}}(s) = \frac{2(\mu a)^{2s}}{(2k-1)!} \sum_{l=0}^{k-1} c_l^{(2k)} \sum_{n=k}^{\infty} \left(n - \frac{1}{2}\right)^{1+2l-2s} \left[1 + \frac{(am)^2 - 1/4}{\left(n - \frac{1}{2}\right)^2} \right]^{-s},$$

$$\zeta_{S^{2k}}(s) = \frac{2(\mu a)^{2s}}{(2k-1)!} \sum_{l=0}^{k-1} c_l^{(2k)} \sum_{p=0}^{\infty} \frac{\Gamma(-s+1)}{p! \Gamma(-s+1-p)} \left[(am)^2 - \frac{1}{4} \right]^p$$

$$\times \zeta_H\left(2s+2p-2l-1, \frac{3}{2}\right)$$

$$\zeta_{S^{2k}}(0) = \frac{2}{(2k-1)!} \sum_{l=0}^{k-1} c_l^{(2k)} \left\{ \zeta_H\left(-2l-1, \frac{3}{2}\right) + \frac{1}{2(l+1)} \left[\frac{1}{4} - (am)^2 \right]^{l+1} \right\}$$

$$\zeta_{S^{2k}/Z_2}(s) = \frac{4(\mu a)^{2s}}{(2k-1)!} \sum_{l=0}^{k-1} c_l^{(2k)} 2^{-2s+2l} \sum_{n=0}^{\infty} \left(n - \frac{1}{4}\right)^{1+2l-2s} \left[1 + \frac{(am)^2 - 1/4}{4\left(n - \frac{1}{4}\right)^2}\right]^{-s}$$

$$\zeta_{S^{2k}/Z_2}(s) = \frac{(\mu a)^{2s}}{(2k-1)!} \sum_{l=0}^{k-1} c_l^{(2k)} \sum_{p=0}^{\infty} \frac{\Gamma(s+p)}{p! \Gamma(s)} 2^{2+2l-2s-2p} \left[\frac{1}{4} - (am)^2\right]^p$$

$$\times \zeta_H\left(2s+2p-2l-1, \frac{3}{4}\right)$$

$$\zeta_{S^{2k}/Z_2}(0) = \frac{1}{(2k-1)!} \sum_{l=0}^{k-1} c_l^{(2k)} \left\{ 2^{2+2l} \zeta_H\left(-2l-1, \frac{3}{4}\right) + \frac{1}{2(l+1)} \left[\frac{1}{4} - (am)^2\right]^{l+1} \right\}$$

$$\zeta^{\text{hol}}(s) := \zeta_{S^{2k}/Z_2}(s) - \frac{1}{2} \zeta_{S^{2k}}(s)$$

$$S_{\text{hol}, S^{2k}/Z_2} = \frac{(-1)^{k+1}}{(2k-1)!} \sum_{l=0}^{k-1} c_l^{(2k)} \left\{ \frac{\left[\frac{1}{4} - (am)^2\right]^{l+1}}{2(l+1)} [\log 2 + \psi(3/4) - \psi(3/2)] \right.$$

$$- 2^{2+2l} \left[-\log 2 \zeta_H\left(-2l-1, \frac{3}{4}\right) + \zeta'_H\left(-2l-1, \frac{3}{4}\right) \right]$$

$$- \sum_{\substack{p=1 \\ p \neq l+1}}^{\infty} \frac{\left[\frac{1}{4} - (am)^2\right]^p}{2p} \left[2^{2+2l-2p} \zeta_H\left(2p-2l-1, \frac{3}{4}\right) \right.$$

$$\left. - \zeta_H\left(2p-2l-1, \frac{3}{2}\right) \right] + \zeta'_H\left(-2l-1, \frac{3}{2}\right) \left. \right\}$$

$$S[\psi] = \frac{1}{2} \int d^d x \sqrt{g} \bar{\psi} (\not{\partial} + m) \psi,$$

$$\lambda_n^{\pm} = \pm \frac{i}{a} \left(n + \frac{d}{2}\right) + m, d_n \equiv d_n^{\pm} = 2 \left\lfloor \frac{d}{2} \right\rfloor \binom{d+n-1}{n},$$

$$\lambda_n^+ = \frac{i}{a} \left(n + \frac{d}{2}\right) + m, n \text{ par}, \lambda_n^- = -\frac{i}{a} \left(n + \frac{d}{2}\right) + m, n \text{ impar},$$

$$\zeta_{\text{hol}}^{(+)}(s) := \zeta_{S^3/Z_2}(s) - \frac{1}{2} \zeta_{S^3}(s).$$

$$\zeta_{\text{hol}}^{(+)}(s) = \frac{\mu^s}{2} \left\{ \sum_{\substack{n=0 \\ n \text{ par}}}^{\infty} d_n [(\lambda_n^+)^{-s} - (\lambda_n^-)^{-s}] - \sum_{\substack{n=0 \\ n \text{ impar}}}^{\infty} d_n [(\lambda_n^+)^{-s} - (\lambda_n^-)^{-s}] \right\}$$

$$\zeta_{\text{hol}}^{(+)}(s) = (\mu a)^s 2^{1-s} \left\{ e^{-i\frac{\pi}{2}s} \delta \zeta_H(s, am) - e^{i\frac{\pi}{2}s} \delta \zeta_H(s, -am) \right\}$$

$$\delta \zeta_H(s, am) := \tilde{\delta} \zeta_H(s-2, am) + iam \tilde{\delta} \zeta_H(s-1, am)$$

$$- \frac{(am)^2 + 1/4}{4} \tilde{\delta} \zeta_H(s, am)$$

$$\text{con} \tilde{\delta} \zeta_H(s, am) := \zeta_H\left(s, \frac{3}{4} - \frac{iam}{2}\right) - \zeta_H\left(s, \frac{5}{4} - \frac{iam}{2}\right)$$

$$\Gamma_{\text{hol}}^{(+)}(am) = -i\frac{\pi}{8} + 2\delta \zeta'_H(0, am) - 2\delta \zeta'_H(0, -am)$$

$$\Gamma_{\text{hol}}^{(+)}(0) = -i\frac{\pi}{8}$$

$$\Gamma_{\text{hol}}^{(-)}(am) = -i\frac{\pi}{8} - 2\delta \zeta'_H(0, am) + 2\delta \zeta'_H(0, -am)$$

$$Z^{S^3/Z_2} := \frac{1}{2} (Z^{(+)} + Z^{(-)})$$

$$\begin{aligned}
\Gamma^{S^3/Z_2} &\equiv -\log Z^{S^3/Z_2} = -\log \left(\frac{e^{-\Gamma^{(+)}} + e^{-\Gamma^{(-)}}}{2} \right) \\
S_{\text{hol}, S^3/Z_2} &= \log \cos \Im(\Gamma^{S^3/Z_2, (+)}) \\
\Gamma^{S^d/Z_2} &:= -\log \left(\frac{e^{-\Gamma^{(+)}} + e^{-\Gamma^{(-)}}}{2} \right) \\
\lambda_n^\pm &= \pm \frac{i}{a} n + m, n = k + 1, k + 2, \dots, \\
d_n &= \frac{2^{k+1}}{(2k + 1)!} \frac{(n + k)!}{(n - k - 1)!} \\
\frac{(n + k)!}{(n - k - 1)!} &= (n + k)(n + k - 1)(n + k - 2) \dots (n - k + 1)(n - k) \\
&= n(n^2 - 1)(n^2 - 2^2) \dots (n^2 - k^2) \\
d_n &= \frac{2^{k+1}}{(2k + 1)!} \sum_{l=0}^k c_l^{(2k+2)} n^{1+2l} \\
\zeta_{S^{2k+2}}(s) &= (\mu a)^s \frac{2^{k+1}}{(2k + 1)!} \sum_{l=0}^k c_l^{(2k+2)} \left\{ \sum_{n=1}^{\infty} n^{1+2l} (in + am)^{-s} + \right. \\
&\quad \left. + (-1)^{-s} \sum_{n=0}^{\infty} n^{1+2l} (in - am)^{-s} \right\} \\
\sum_{n=1}^{\infty} n^{1+2l} (in + am)^{-s} &= i^{-s} \sum_{n=1}^{\infty} (n - iam + iam)^{1+2l} (n - iam)^{-s} \\
&= i^{-s} \sum_{p=0}^{2l+1} \binom{2l+1}{p} (iam)^{2l+1-p} \zeta_H(s - p, 1 - iam) \\
\zeta_{S^{2k+2}}(s) &= (\mu a)^s \frac{2^{k+1}}{(2k + 1)!} \sum_{l=0}^k c_l^{(2k+2)} \sum_{p=0}^{2l+1} \binom{2l+1}{p} (iam)^{2l+1-p} \\
&\quad \times \{ i^{-s} \zeta_H(s - p, 1 - iam) + (-i)^{-s} (-1)^{1-s-p} \zeta_H(s - p, iam) \} \\
\Gamma^{S^{2k+2}} &= \frac{2^{k+1}}{(2k + 1)!} \sum_{l=0}^k c_l^{(2k+2)} \sum_{p=0}^{2l+1} \binom{2l+1}{p} (iam)^{2l+1-p} \\
&\quad \times \{ \log(\mu a) [\zeta_H(-p, 1 - iam) + (-1)^{1+p} \zeta_H(-p, iam)] \\
&\quad + \zeta'_H(-p, 1 - iam) + (-1)^{1+p} \zeta'_H(-p, iam) \\
&\quad - i\pi [\zeta_H(-p, 1 - iam) - (-1)^{1+p} \zeta_H(-p, iam)] \} \\
\Gamma^{S^{2k+2}} &= \frac{2^{k+1}}{(2k + 1)!} \sum_{l=0}^k c_l^{(2k+2)} \sum_{p=0}^{2l+1} \binom{2l+1}{p} (iam)^{2l+1-p} \\
&\quad \times \{ 2(-1)^{1+p} \zeta_H(-p, iam) \log(\mu a) + \zeta'_H(-p, 1 - iam) \\
&\quad + (-1)^{1+p} \zeta'_H(-p, iam) \} \\
\Gamma^{S^{2k+2}/Z_2, (\pm)} &= \frac{2^{k+2}}{(2k + 1)!} \sum_{l=0}^k c_l^{(2k+2)} \sum_{p=0}^{2l+1} \binom{2l+1}{p} 2^p (iam)^{2l+1-p} \\
&\quad \times \left\{ \log \frac{\mu a}{2} \left[\zeta_H \left(-p, \frac{1}{4} \mp \frac{1}{4} - \frac{iam}{2} \right) - (-1)^p \zeta_H \left(-p, \frac{1}{4} \pm \frac{1}{4} + \frac{iam}{2} \right) \right] \right. \\
&\quad + \zeta'_H \left(-p, \frac{1}{4} \mp \frac{1}{4} - \frac{iam}{2} \right) - (-1)^p \zeta'_H \left(-p, \frac{1}{4} \pm \frac{1}{4} + \frac{iam}{2} \right) \\
&\quad \left. - \frac{i\pi}{2} \left[\zeta_H \left(-p, \frac{1}{4} \mp \frac{1}{4} - \frac{iam}{2} \right) + (-1)^p \zeta_H \left(-p, \frac{1}{4} \pm \frac{1}{4} + \frac{iam}{2} \right) \right] \right\}
\end{aligned}$$

$$\begin{aligned}
S_{\text{hol}, S^{2k+2}/Z_2} &= \log \cos \mathfrak{S} \left(\Gamma^{S^{2k+2}/Z_2, (+)} \right). \\
\mathcal{A}^{(2k+2)}(m) &= \frac{2^{k+2}}{(2k+1)!} \sum_{l=0}^k c_l^{(2k+2)} \sum_{p=0}^{2l+1} \binom{2l+1}{p} (iam)^{2l+1-p} \\
&\quad \times (-1)^{1+p} \zeta_H(-p, iam) \\
(iam)^{2l+1-p} \zeta_H(-p, iam) &= -\frac{1}{p+1} \sum_{q=0}^{p+1} \binom{p+1}{q} (iam)^{2l+2-q} B_q \\
-\frac{1}{p+1} \sum_{q=0}^{p+1} \binom{p+1}{q} (iam)^{2l+2-q} B_q &= -\frac{1}{p+1} \binom{p+1}{1} (iam)^{2l+1} B_1 \\
\frac{2^{k+1}}{(2k+1)!} \sum_{l=0}^k c_l^{(2k+2)} \sum_{p=0}^{2l+1} \binom{2l+1}{p} (-1)^{1+p} (iam)^{2l+1} \\
&\quad - \frac{2^{k+1}}{(2k+1)!} \sum_{l=0}^k c_l^{(2k+2)} (iam)^{2l+1} (1-1)^{2l+1} \\
\mathcal{A}^{(2k+2)} &= \frac{2^{k+2}}{(2k+1)!} \sum_{l=0}^k c_l^{(2k+2)} \zeta_R(-2l-1), \\
\mathcal{A}^{(2k+2)} &= -\frac{2^{k+2}}{(2k+1)!} \sum_{l=0}^k c_l^{(2k+2)} \frac{B_{2l+2}}{2l+2} \\
\lambda_n^\pm &= \pm \frac{i}{a} \left(n + \frac{1}{2} \right) + m, n = k, k+1, k+2, \dots, \\
d_n &= \frac{2^k (n+k)!}{(2k)! (n-k)!} \\
\frac{(n+k)!}{(n-k)!} &= (n+k)(n+k-1)(n+k-2) \dots (n-k+1) \\
&= \left(n + \frac{1}{2} + k - \frac{1}{2} \right) \left(n + \frac{1}{2} + k - \frac{3}{2} \right) \dots \left(n + \frac{1}{2} - k + \frac{1}{2} \right) \\
&= \left[\left(n + \frac{1}{2} \right)^2 - \left(k - \frac{1}{2} \right)^2 \right] \left[\left(n + \frac{1}{2} \right)^2 - \left(k - \frac{3}{2} \right)^2 \right] \dots \left[\left(n + \frac{1}{2} \right)^2 - \frac{1}{4} \right] \\
d_n &= \frac{2^k}{(2k)!} \sum_{l=0}^k c_l^{(2k+1)} \left(n + \frac{1}{2} \right)^{2l} \\
\zeta_{S^{2k+1}}(s) &= (\mu a)^s \frac{2^k}{(2k)!} \sum_{l=0}^k c_l^{(2k+1)} \left\{ i^{-s} \sum_{n=0}^{\infty} \left(n + \frac{1}{2} \right)^{2l} \left(n + \frac{1}{2} - iam \right)^{-s} \right. \\
&\quad \left. + (-i)^{-s} \sum_{n=0}^{\infty} \left(n + \frac{1}{2} \right)^{2l} \left(n + \frac{1}{2} + iam \right)^{-s} \right\} \\
\zeta_{S^{2k+1}}(s) &= (\mu a)^s \frac{2^k}{(2k)!} \sum_{l=0}^k c_l^{(2k+1)} \sum_{p=0}^{2l} \binom{2l}{p} (iam)^{2l-p} \\
&\quad \times \left\{ i^{-s} \zeta_H \left(s - p, \frac{1}{2} - iam \right) + (-i)^{-s} (-1)^p \zeta_H \left(s - p, \frac{1}{2} + iam \right) \right\}
\end{aligned}$$

$$\begin{aligned}
\Gamma^{S^{2k+1}} &= \frac{2^k}{(2k)!} \sum_{l=0}^k c_l^{(2k+1)} \sum_{p=0}^{2l} \binom{2l}{p} (iam)^{2l-p} \\
&\times \left\{ \zeta'_H \left(-p, \frac{1}{2} - iam \right) + (-1)^p \zeta'_H \left(-p, \frac{1}{2} + iam \right) \right. \\
&\quad \left. - i\pi \zeta_H \left(-p, \frac{1}{2} - iam \right) \right\} \\
\Gamma^{S^{2k+1}/Z_2, (\pm)} &= \frac{2^k}{(2k)!} \sum_{l=0}^k c_l^{(2k+1)} \sum_{p=0}^{2l} \binom{2l}{p} 2^p (iam)^{2l-p} \\
&\times \left\{ \zeta'_H \left(-p, \frac{1}{2} - \frac{iam}{2} \mp \frac{1}{4} \right) + (-1)^p \zeta'_H \left(-p, \frac{1}{2} + \frac{iam}{2} \pm \frac{1}{4} \right) \right. \\
&\quad \left. - i\pi \zeta_H \left(-p, \frac{1}{2} - \frac{iam}{2} \mp \frac{1}{4} \right) \right\}
\end{aligned}$$

Bajo la métrica de Yukawa para curvatura, obtenemos:

$$\begin{aligned}
m_* &\lesssim g M_P^{D/2-1} \\
Y_{ijk} &= e^{\phi_4/2} \text{Vol}_X^{1/4} \Theta_{ijk}^{1/4} W_{ijk} \\
|Y_{ijk}| &\simeq e^{-\phi_4} \left(\frac{m_{\text{bos/fer}}^i}{M_P} \frac{m_{\text{bos/fer}}^j}{M_P} \frac{m_{\text{bos/fer}}^k}{M_P} \right)^{1/2} \\
Y_* &\sim \frac{1}{u^r} \\
m_{\text{bos/fer}} &\sim g_* M_s \sim g_*^2 M_P, Y_* \sim g_*, g_* \sim e^{\phi_4} \\
\sum_{\alpha} N_{\alpha} ([\Pi_{\alpha}] + [\Pi_{\alpha^*}]) &= 4[\Pi_{06}]. \\
J_c &\equiv B + ij = (b^a + it^a) \omega_a \\
\Omega_c &\equiv C_3 + ie^{-\phi} \text{Re} \Omega = (\zeta^K + iu^K) \alpha_K \\
K_K &\equiv -\log(\text{Vol}_{X_6}) = -\log \left(\frac{i}{48} \mathcal{K}_{abc} (T^a - \bar{T}^a)(T^b - \bar{T}^b)(T^c - \bar{T}^c) \right) \\
K_Q &\equiv -2 \log \mathcal{H} = -2 \log \left(\frac{i}{8\ell_s^6} \int_{X_6} e^{-2\phi} \Omega \wedge \bar{\Omega} \right) = 4\phi_4 \\
\ell_K &\equiv -\frac{1}{2} \frac{\partial K}{\partial u^K} = -\frac{1}{2\mathcal{H}} \int_{\Sigma_{\bar{K}}} e^{-\phi} \text{Im} \Omega \\
[\Pi_{\alpha}] &= P_{\alpha J} [\Sigma_+^J] + Q_{\alpha}^K [\Sigma_{\bar{K}}], [\Pi_{\alpha^*}] = P_{\alpha J} [\Sigma_-^J] - Q_{\alpha}^K [\Sigma_{\bar{K}}] \\
2\pi i f_{\alpha\alpha} &= P_{\alpha K} U^K \\
\frac{1}{2} M_P^2 \int_{X_4} G_{KL} (d\zeta^K - 2Q_{\alpha}^K A^{\alpha}) \wedge * (d\zeta^L - 2Q_{\beta}^L A^{\beta}) \\
M_{\alpha\beta}^2 &= 4G_{KL} Q_{\alpha}^K Q_{\beta}^L g_{\alpha} g_{\beta} M_P^2 \\
V_D &= \frac{1}{2} \sum_{\alpha} g_{\alpha}^2 \left(\xi_{\alpha} + \sum_i q_{\alpha}^i K_{i\bar{i}} \Phi_i \bar{\Phi}_{\bar{i}} \right)^2, \text{ with } \pi \xi_{\alpha} = Q_{\alpha}^K \ell_K M_P^2 \\
&\quad q_{\alpha}^i g_{\alpha}^2 \xi_{\alpha} + q_{\beta}^i g_{\beta}^2 \xi_{\beta}. \\
&\quad \theta_{\alpha\beta}^1 + \theta_{\alpha\beta}^2 + \theta_{\alpha\beta}^3 \in 2\mathbb{Z}, \\
m_{\alpha\beta}^2 (\text{scalar})_{\bar{k}, \pm r} &= 2\pi \left[\lambda_{\alpha\beta}^{\bar{k}} \pm |\theta_{\alpha\beta}^r| \right] e^{2\phi_4} M_P^2, \\
\lambda_{\alpha\beta}^{\bar{k}} &= \sum_r k_r |\theta_{\alpha\beta}^r| + \frac{1}{2} \sum_p |\theta_{\alpha\beta}^p|, k_r \in \mathbb{Z}_{>0}
\end{aligned}$$



$$\begin{aligned}
m_{\alpha\beta}^2(\text{vector})_{\bar{k}} &= 2\pi\lambda_{\alpha\beta}^{\bar{k}} e^{2\phi_4} M_{\text{P}}^2 \\
Y_{ijk} &= h_{qu} \sum_{\bar{n}} d_{\bar{n}} e^{2\pi i \left[\int_{D_{\bar{n}}} J_c - \int_{\partial D_{\bar{n}}} v \right]} \\
Y_{ijk} &= e^{\phi_4/2} \prod_{r=1}^3 (\text{Im} T^r)^{1/4} [\Theta^{(r)}]^{1/4} e^{\frac{1}{2}H^{(r)}} W_{ijk}^{(r)} \\
\Theta^{(r)} &= 2\pi \frac{\Gamma(1 - |\chi_{ab}^r|) \Gamma(1 - |\chi_{bc}^r|) \Gamma(1 - |\chi_{ca}^r|)}{\Gamma(|\chi_{ab}^r|) \Gamma(|\chi_{bc}^r|) \Gamma(|\chi_{ca}^r|)} \\
W_{ijk}^{(r)} &= \vartheta \begin{bmatrix} \delta_{ijk}^r \\ 0 \end{bmatrix} (v^r, |I_{ab}^r I_{bc}^r I_{ca}^r| T^r) \\
H^{(r)} &= 2\pi i \frac{v^r \text{Im} v^r}{\text{Im} T^r} |I_{ab}^r I_{bc}^r I_{ca}^r|^{-1} \\
Y_{ijk} &= e^{\phi_4/2} \text{Vol}_X^{1/4} \Theta_{ijk}^{1/4} W_{ijk} \\
Y_{ijk} &= e^{K/2} (K_{i\bar{i}} K_{j\bar{j}} K_{k\bar{k}})^{-1/2} W_{ijk} \\
K_{i\bar{i}} K_{j\bar{j}} K_{k\bar{k}} &= e^{3(K/2 - \phi_4)} \Theta_{ijk}^{-1/2} \\
K_{i\bar{i}} &= \frac{e^{K/2 - \phi_4}}{\sqrt{2\pi}} \prod_{r=1}^3 \left(\frac{\Gamma(|\chi_i^r|)}{\Gamma(1 - |\chi_i^r|)} \right)^{1/2} \\
K_{i\bar{i}} &= \frac{e^{K/2 - \phi_4}}{\sqrt{2\pi}} \left[\frac{\Gamma(|\theta_{\alpha\beta}^1|)}{\Gamma(1 - |\theta_{\alpha\beta}^1|)} \frac{\Gamma(|\theta_{\alpha\beta}^2|)}{\Gamma(1 - |\theta_{\alpha\beta}^2|)} \frac{\Gamma(1 - |\theta_{\alpha\beta}^3|)}{\Gamma(|\theta_{\alpha\beta}^3|)} \right]^{1/2} |\theta_{\alpha\beta}^r| \ll 1 \approx \left[\frac{e^{K - 2\phi_4} |\theta_{\alpha\beta}^3|}{2\pi |\theta_{\alpha\beta}^1| |\theta_{\alpha\beta}^2|} \right]^{1/2} \\
K &= -2\log(\text{Vol}_X^{1/2} \mathcal{H}) = -2\log(\hat{V}_X^{1/2} \hat{\mathcal{H}} - f(\Phi^i, \bar{\Phi}^{\bar{i}})) = -2\log\left(\hat{V}_X^{1/2} \hat{\mathcal{H}} - \frac{1}{2} k_{i\bar{i}} \Phi^i \bar{\Phi}^{\bar{i}} + \dots\right), \\
K &= -2\log(\hat{V}_X^{1/2} \hat{\mathcal{H}}) + \frac{k_{i\bar{i}}}{\hat{V}_X^{1/2} \hat{\mathcal{H}}} \Phi^i \bar{\Phi}^{\bar{i}} + \dots \\
k_{i\bar{i}} &= \frac{e^{-\phi_4}}{\sqrt{2\pi}} \prod_{r=1}^3 \left(\frac{\Gamma(|\chi_i^r|)}{\Gamma(1 - |\chi_i^r|)} \right)^{1/2} \\
k_{i\bar{i}} &\simeq \frac{M_{\text{P}}}{m_{\text{bos/fer}}^i} \\
K_{i\bar{i}} &= e^{K/2} h_{i\bar{i}} \prod_{r=1}^2 \left(\frac{\Gamma(|\chi_i^r|)}{\Gamma(1 - |\chi_i^r|)} \right)^{1/2} \simeq e^{K/2} h_{i\bar{i}} \\
h_{i\bar{i}} &\lesssim g_{\alpha}^{-1} g_{\beta}^{-1} \\
|Y_{ijk}|^2 &= A \text{Vol}_X^{1/2} |W_{ijk}|^2 \\
|Y_{ijk}|^2 &= B e^{\phi_4} \Theta_{ijk}^{1/2} \\
m_{\text{bos/fer}} &\sim |Y|^2 M_{\text{P}} \\
|Y_{ijk}|^2 &\simeq B e^{-2\phi_4} \frac{m_{\text{bos/fer}}^i}{M_{\text{P}}} \frac{m_{\text{bos/fer}}^j}{M_{\text{P}}} \frac{m_{\text{bos/fer}}^k}{M_{\text{P}}} \\
\left(\frac{m_{\text{bos/fer}}^i}{M_{\text{P}}}\right)^{1/w_i} &\sim \left(\frac{m_{\text{bos/fer}}^j}{M_{\text{P}}}\right)^{1/w_j} \sim \left(\frac{m_{\text{bos/fer}}^k}{M_{\text{P}}}\right)^{1/w_k} \sim e^{\phi_4}, w_i, w_j, w_k \geq 1 \\
m_{\text{bos/fer}}^{ijk} &\sim |Y_{ijk}|^{\frac{2w_{ijk}}{w_i + w_j + w_k - 2}} M_{\text{P}} \\
|Y_{ijk}|^2 &\simeq B e^{-2\phi_4} h_{i\bar{i}}^{-1} \frac{m_{\text{bos/fer}}^j}{M_{\text{P}}} \frac{m_{\text{bos/fer}}^k}{M_{\text{P}}} = B h_{i\bar{i}}^{-1} \frac{m_{\text{bos/fer}}^j}{M_{\text{S}}} \frac{m_{\text{bos/fer}}^k}{M_{\text{S}}} \\
m_{\text{mono}} &\simeq g^{-2} m_{\text{ele}} \\
T_{\text{string}} &= 2\pi |\xi_{\alpha}| \simeq g_{\alpha}^{-2} |m_{\text{ele}}^2| \\
m_{\alpha\beta,+}^2 &= g_{\alpha}^2 T_{\alpha}^+ - g_{\beta}^2 T_{\beta}^-, m_{\alpha\beta,-}^2 = -g_{\alpha}^2 T_{\alpha}^- + g_{\beta}^2 T_{\beta}^+
\end{aligned}$$



$$\begin{aligned}
m_{\alpha\beta,\min}^2 &= \min_K (|g_\alpha^2 Q_\alpha^K - g_\beta^2 Q_\beta^K| T_K) \\
\nabla g_\alpha^{-2} \cdot \nabla T_K &\neq 0, \text{ y/o } \nabla g_\beta^{-2} \cdot \nabla T_K \neq 0 \\
m_{\alpha\beta,\min} &\lesssim m_{\text{bos/fer}}^{\min} \lesssim m_{\alpha\beta,\pm} \\
u^l &\sim \lambda \rightarrow \infty \\
u^K \ell_K &= 2 \\
K_Q &= -n \log \lambda + \dots \\
M_s &\sim M_P \lambda^{-n/4} \\
e^{-\phi} \text{Re} \Omega &= s \alpha_0 + u^1 \alpha_1 + u^2 \alpha_2 \\
K_Q &= -\log s - \log \kappa(u^1, u^2) \\
\kappa(u^1, u^2) &\equiv \frac{1}{6} [\kappa_{111}(u^1)^3 + 3\kappa_{112}(u^1)^2 u^2 + 3\kappa_{122} u^1 (u^2)^2 + \kappa_{222}(u^2)^3] \\
\ell_0 &= \frac{1}{2s}, \ell_1 = \frac{\partial_{u^1} \kappa(u^1, u^2)}{2\kappa(u^1, u^2)}, \ell_2 = \frac{\partial_{u^2} \kappa(u^1, u^2)}{2\kappa(u^1, u^2)} \\
[\Pi_\alpha] &= \frac{1}{2} [A] + \frac{1}{2} f_\alpha^i [D_i] - \left(\frac{1}{2} \kappa_{ijk} f_\alpha^j f_\alpha^k + \frac{1}{24} c_{2i} \right) [C^i] + \left(\frac{1}{6} \kappa_{ijk} f_\alpha^i f_\alpha^j f_\alpha^k + \frac{1}{24} c_{2i} f_\alpha^i \right) [B] \\
\frac{2\pi}{g_\alpha^2} &= s - \left(\frac{1}{2} \kappa_{ijk} f_\alpha^j f_\alpha^k + \frac{1}{24} c_{2i} \right) u^i, \\
\ell_i f_\alpha^i &= \frac{1}{2s} \left(\frac{1}{6} \kappa_{ijk} f_\alpha^i f_\alpha^j f_\alpha^k + \frac{1}{24} c_{2i} f_\alpha^i \right) \\
M_{\alpha\alpha}^2 &\sim g_\alpha^2 Q_\alpha^K Q_\alpha^L \frac{\partial_K \partial_L (s\kappa)}{s\kappa} M_P^2 \sim \frac{M_P^2}{s u^2}. \\
|Y|^2 &\sim 1/s \\
\text{I) } s &\sim u \sim \lambda, \text{ II) } s \sim u^3 \sim \lambda. \\
\text{I) } \frac{m_{\text{Stu}}}{M_P} &\sim \left(\frac{m_{\text{KK}}}{M_P} \right)^{3/2} \sim \left(\frac{m_{\text{bos/fer}}}{M_P} \right)^{3/2} \sim \left(\frac{M_s}{M_P} \right)^{3/2} \sim \left(\frac{T_{\text{D4}}^{1/2}}{M_P} \right)^3 m_{\text{bos/fer}} \sim |Y|^2 M_P \\
\text{II) } \frac{m_{\text{Stu}}}{M_P} &\sim \left(\frac{m_{\text{KK}}}{M_P} \right)^{5/4} \sim \left(\frac{m_{\text{bos/fer}}}{M_P} \right)^{5/4} \sim \left(\frac{T_{\text{D4}}^{1/2}}{M_P} \right)^{5/3} \sim \left(\frac{M_s}{M_P} \right)^{5/3}, m_{\text{bos/fer}} \sim |Y|^{4/3} M_P. \\
\frac{m_*^2}{M_P^2} &\sim \left(\frac{T_{\text{D4}}}{M_P^2} \right)^w \\
\frac{m_{\text{KK/w}}}{M_P} &\sim \frac{T_{\text{D4}}^{1/2}}{M_P} \sim 1/\sqrt{\lambda} \\
\frac{m_{\text{bos/fer}}^{\min}}{M_P} &\sim 1/\sqrt{\lambda} \\
\frac{m_{\text{KK/w}}}{M_P} &\sim \frac{T_{\text{D4}}}{M_P^2} \sim 1/\lambda \\
m_{\text{bos/fer}}^{\min} &\sim M_P/\lambda, \text{ or } m_{\text{bos/fer}}^{\min} \sim M_P/1\sqrt{\lambda} \\
M_{\alpha\beta}^2 &= g_\alpha g_\beta Q_\alpha^K Q_\beta^L \frac{\partial_K \partial_L \mathcal{H}}{\mathcal{H}} M_P^2 \\
\frac{m_{\text{bos/fer}}^{\min}}{M_P} &\sim \left(\frac{m_{\text{Stu}}}{M_P} \right)^{2/3} \sim 1/\lambda \\
\frac{m_{\text{bos/fer}}^{\min}}{M_P} &\sim \left(\frac{m_{\text{Stu}}}{M_P} \right)^{1/2} \sim 1/\sqrt{\lambda} \\
\frac{m_{\text{bos/fer}}^{\min}}{M_P} &\sim \frac{m_{\text{Stu}}}{M_P} \sim 1/\sqrt{\lambda} \\
|Y| &\sim \frac{1}{u^r}, \text{ with } r = \frac{1}{4}, \frac{1}{2}, \frac{3}{4}, 1
\end{aligned}$$



$$\begin{aligned} \Omega &= \ell_s^3 R_1 R_3 R_5 dz_1 \wedge dz_2 \wedge dz_3 \\ &= \ell_s^3 (R_1 dy_1 + iR_2 dy_2) \wedge (R_3 dy_3 + iR_4 dy_4) \wedge (R_5 dy_5 + iR_6 dy_6) \\ \alpha_0 &= 4dy_1 \wedge dy_3 \wedge dy_5 \ell_s^3, \alpha_1 = -4dy_1 \wedge dy_4 \wedge dy_6 \ell_s^3 \\ \alpha_2 &= -4dy_2 \wedge dy_3 \wedge dy_6 \ell_s^3, \alpha_3 = -4dy_2 \wedge dy_4 \wedge dy_5 \ell_s^3 \\ \beta^0 &= -2dy_2 \wedge dy_4 \wedge dy_6 \ell_s^3, \beta^1 = 2dy_2 \wedge dy_3 \wedge dy_5 \ell_s^3 \\ \beta^2 &= 2dy_1 \wedge dy_4 \wedge dy_5 \ell_s^3, \beta^3 = 2dy_1 \wedge dy_3 \wedge dy_6 \ell_s^3 \\ s &\equiv u^{(0)} = \frac{1}{4} e^{-\phi} R_1 R_3 R_5, u^{(1)} = \frac{1}{4} e^{-\phi} R_1 R_4 R_6 \\ u^{(2)} &= \frac{1}{4} e^{-\phi} R_2 R_3 R_6, u^{(3)} = \frac{1}{4} e^{-\phi} R_2 R_4 R_5 \\ \mathcal{H} &= \frac{i}{32 \ell_s^6} \int_{T^6} e^{-2\phi} \Omega \wedge \bar{\Omega} = 4 \sqrt{u^{(0)} u^{(1)} u^{(2)} u^{(3)}} \end{aligned}$$

$$\ell_K = \frac{1}{2u^K}$$

$$m_{KK,i} = \frac{2M_s}{R_{2i-1}} = \frac{M_P}{A_i^{1/2} 2\sqrt{su^{(i)}}}, m_{w,i} = \frac{1}{2} R_{2i} M_s = \frac{A_i^{1/2} M_P}{2\sqrt{su^{(i)}}}$$

$$M_s^2 = \frac{M_P^2}{4\sqrt{su^{(1)}u^{(2)}u^{(3)}}}$$

$$\begin{aligned} [\Pi_\alpha] &= 2[(n_\alpha^1, m_\alpha^1)(n_\alpha^2, m_\alpha^2)(n_\alpha^3, m_\alpha^3)], n_\alpha^i, m_\alpha^i \in \mathbb{Z} \\ &= n_\alpha^1 n_\alpha^2 n_\alpha^3 \Sigma_+^0 - n_\alpha^1 m_\alpha^2 m_\alpha^3 \Sigma_+^1 - m_\alpha^1 n_\alpha^2 m_\alpha^3 \Sigma_+^2 - m_\alpha^1 m_\alpha^2 n_\alpha^3 \Sigma_+^3 \\ &+ \frac{1}{2} (m_\alpha^1 m_\alpha^2 m_\alpha^3 \Sigma_0^- - m_\alpha^1 n_\alpha^2 n_\alpha^3 \Sigma_1^- - n_\alpha^1 m_\alpha^2 n_\alpha^3 \Sigma_2^- - n_\alpha^1 n_\alpha^2 m_\alpha^3 \Sigma_3^-) \end{aligned}$$

$$[\Pi_{\alpha^*}] = 2[(n_\alpha^1, -m_\alpha^1)(n_\alpha^2, -m_\alpha^2)(n_\alpha^3, -m_\alpha^3)].$$

$$\Pi_a = 2(k, 1)(k, 1)(k, -1),$$

$$\Pi_b = 2(0, 1)(1, 0)(0, -1),$$

$$\Pi_c = 2(0, 1)(0, -1)(1, 0)$$

$$\frac{\pi \xi_a}{M_P^2} = -\frac{1}{4s} - \frac{k^2}{4u^{(1)}} - \frac{k^2}{4u^{(2)}} + \frac{k^2}{4u^{(3)}}$$

$$2s = e^{-\phi_4} \frac{1}{\sqrt{\tau_1 \tau_2 \tau_3}}, 2u^{(i)} = e^{-\phi_4} \sqrt{\frac{\tau_j \tau_k}{\tau_i}}$$

$$\theta_1 + \theta_2 = \theta_3, \text{ donde } \pi \theta_i = \tan^{-1} \frac{\tau_i}{k}$$

$$\frac{2\pi}{g_a^2} = 2(k^3 s + ku^{(1)} + ku^{(2)} - ku^{(3)}), \frac{2\pi}{g_b^2} = 2u^{(2)}, \frac{2\pi}{g_c^2} = 2u^{(3)}$$

$$q_a g_a^2 \xi_a = q_a \tan[\pi(\theta_3 - \theta_2 - \theta_1)] M_s^2$$

$$M_{aa}^2 = \frac{1}{2} g_a^2 \left(\frac{1}{s^2} + \frac{k^4}{(u^{(1)})^2} + \frac{k^4}{(u^{(2)})^2} + \frac{k^4}{(u^{(3)})^2} \right) M_P^2$$

$$\frac{m_{\text{bos/fer},i}}{M_P} \simeq \frac{1}{2\sqrt{ksu^{(i)}}}$$

Sector	$m_{\text{bos/fer},1}/M_P$	$m_{\text{bos/fer},2}/M_P$	$m_{\text{bos/fer},3}/M_P$
aa^*	$1/\sqrt{ksu^{(1)}}$	$1/\sqrt{ksu^{(2)}}$	$1/\sqrt{ksu^{(3)}}$
ab	e^{ϕ_4}	$1/\sqrt{ksu^{(2)}}$	e^{ϕ_4}
bc	e^{ϕ_4}	e^{ϕ_4}	e^{ϕ_4}
ca	e^{ϕ_4}	e^{ϕ_4}	$1/\sqrt{ksu^{(3)}}$



$$m_{\text{KKD}_{ai}} = \frac{2M_s}{R_{2i-1}\sqrt{k^2 + \tau_i^2}} \simeq \frac{M_{\text{P}}}{A_i^{1/2}2k\sqrt{su^{(i)}}}, m_{\text{wD}_{ai}} = \frac{R_{2i}}{2\sqrt{k^2 + \tau_i^2}}M_s \simeq \frac{A_i^{1/2}M_{\text{P}}}{2k\sqrt{su^{(i)}}}.$$

$$s, u^{(2)} \sim \lambda, u^{(1)}, u^{(3)} \\ T_{\text{D}4,0}^{1/2}, T_{\text{D}4,2}^{1/2} \sim \lambda^{-1/2}M_{\text{P}}, M_s \sim \lambda^{-1/2}M_{\text{P}}$$

$$K_{ca} = \frac{e^{K/2-\phi_4}}{\sqrt{2\pi}} \left[\frac{\Gamma(\frac{1}{2}-\theta_1)\Gamma(\frac{1}{2}-\theta_2)}{\Gamma(\frac{1}{2}+\theta_1)\Gamma(\frac{1}{2}+\theta_2)} \frac{\Gamma(\theta_3)}{\Gamma(1-\theta_3)} \right]^{1/2}$$

$$K_{ab} = \frac{e^{K/2-\phi_4}}{\sqrt{2\pi}} \left[\frac{\Gamma(\frac{1}{2}+\theta_1)}{\Gamma(\frac{1}{2}-\theta_1)} \frac{\Gamma(\theta_2)}{\Gamma(1-\theta_2)} \frac{\Gamma(\frac{1}{2}-\theta_3)}{\Gamma(\frac{1}{2}+\theta_3)} \right]^{1/2}$$

$$K_{bc} = e^{K/2-\phi_4}\sqrt{\tau_1/2\pi} = 2e^{K/2}\sqrt{u^{(2)}u^{(3)}/2\pi}$$

$$Y_{abc} = B\sqrt{2\pi g_b g_c} \left[\frac{\Gamma(1-\theta_2)\Gamma(1-\theta_3)}{\Gamma(\theta_2)\Gamma(\theta_3)} \frac{\Gamma(\frac{1}{2}+\theta_2)\Gamma(\frac{1}{2}+\theta_3)}{\Gamma(\frac{1}{2}-\theta_2)\Gamma(\frac{1}{2}-\theta_3)} \right]^{1/4}$$

$$Y_{abc} \sim \left(\frac{\theta_2\theta_3}{u^{(2)}u^{(3)}} \right)^{1/4} \sim \frac{e^{-\phi_4}}{\sqrt{su^{(2)}u^{(3)}}} \sim \lambda^{-1/3}$$

$$Y_{abc} \sim \left(\frac{\theta_2}{u^{(2)}} \right)^{1/4} \sim \frac{e^{-\phi_4/2}}{\sqrt{s^{1/2}u^{(2)}}} \sim \lambda^{-1/2}$$

Limit	Scenario	g_a	g_b	g_c	Y_{abc}
$s \sim (u^{(i)})^3 \sim \lambda$	STU-like II)	$\lambda^{-1/2}$	$\lambda^{-1/6}$	$\lambda^{-1/6}$	$\lambda^{-1/3}$
4.20	quasi-EFT IIa)	$\lambda^{-1/2}$	$\lambda^{-1/2}$	const.	$\lambda^{-1/2}$
(4.20)'	quasi-EFT IIa)	$\lambda^{-1/2}$	const.	const.	const.

$$\Pi_a = 8(1,0)(k,1)(k,-1),$$

$$\Pi_b = 2(0,1)(1,0)(0,-1),$$

$$\Pi_c = 2(0,1)(0,-1)(1,0),$$

$$\frac{2\pi}{g_a^2} = 2(k^2s + u^{(1)})$$

$$\frac{\pi\xi_a}{M_{\text{P}}^2} = \frac{k}{4} \left(-\frac{1}{u^{(2)}} + \frac{1}{u^{(3)}} \right)$$

$$q_a g_a^2 \xi_a = q_a \tan[\pi(\theta_3 - \theta_2)] M_s^2$$

$$M_{aa}^2 = \frac{1}{2} g_a^2 \left(\frac{k^2}{(u^{(2)})^2} + \frac{k^2}{(u^{(3)})^2} \right) M_{\text{P}}^2$$

Sector	$m_{\text{bos/fer},1}/M_{\text{P}}$	$m_{\text{bos/fer},2}/M_{\text{P}}$	$m_{\text{bos/fer},3}/M_{\text{P}}$
aa^*	e^{ϕ_4}	e^{ϕ_4} or $1/\sqrt{ksu^{(2)}}$	e^{ϕ_4} or $1/\sqrt{ksu^{(3)}}$
ab	e^{ϕ_4}	$1/\sqrt{ksu^{(2)}}$	e^{ϕ_4}
bc	e^{ϕ_4}	e^{ϕ_4}	e^{ϕ_4}



$$ca \quad e^{\phi_4} \quad e^{\phi_4} \quad 1/\sqrt{ksu^{(3)}}$$

$$\begin{aligned} s \sim \lambda, u^{(i)} = \text{const.} &\rightarrow e^{\phi_4} \sim \lambda^{-1/4}, \\ s, u^{(2)} = u^{(3)} \sim \lambda, u^{(1)} = \text{const.} &\Rightarrow e^{\phi_4} \sim \lambda^{-3/4}, \\ u^{(2)} = u^{(3)} \sim \lambda, s, u^{(1)} = \text{const.} &\rightarrow e^{\phi_4} \sim \lambda^{-1/2}, \\ s, u^{(1)} \sim \lambda, u^{(2)} = u^{(3)} = \text{const.} &\rightarrow e^{\phi_4} \sim \lambda^{-1/2}, \\ s \sim \lambda, u^{(2)} = u^{(3)} \sim \lambda^{1/2}, u^{(1)} = \text{const.} &\Rightarrow e^{\phi_4} \sim \lambda^{-1/2}. \end{aligned}$$

Limit	Scenario	$\#T_{D4}^{1/2}$	$\#m_{\text{KK}}^{\text{bulk}}$	$\#m_{\text{KK}}^{\text{brane}}$	$\#m_{\text{bos/fer}}$	g_a	g_b	g_c	Y_{abc}
(4.33)	EFT I)	1	6	6	2 or 4	$\lambda^{-1/2}$	const	const	$\lambda^{-1/4}$
(4.34)	EFT IIa)	-	6	4	2 or 4	$\lambda^{-1/2}$	$\lambda^{-1/2}$	$\lambda^{-1/2}$	$\lambda^{-3/4}$
(4.35)	EFT IIb)	-	2	2	-	const	$\lambda^{-1/2}$	$\lambda^{-1/2}$	$\lambda^{-1/2}$
(4.36)	EFT IIc)	-	2	2	-	$\lambda^{-1/2}$	const	const	const
(4.37)	non-EFT	-	4	4	2 or 4	$\lambda^{-1/2}$	$\lambda^{-1/4}$	$\lambda^{-1/4}$	$\lambda^{-1/2}$

$$K_{ca} = \frac{e^{K/2-\phi_4}}{\sqrt{2\pi}} \left[\frac{\Gamma(\frac{1}{2}-\theta_2)}{\Gamma(\frac{1}{2}+\theta_2)} \frac{\Gamma(\theta_3)}{\Gamma(1-\theta_3)} \right]^{1/2}$$

$$K_{ab} = \frac{e^{K/2-\phi_4}}{\sqrt{2\pi}} \left[\frac{\Gamma(\theta_2)}{\Gamma(1-\theta_2)} \frac{\Gamma(\frac{1}{2}-\theta_3)}{\Gamma(\frac{1}{2}+\theta_3)} \right]^{1/2}$$

$$K_{bc} = e^{K/2-\phi_4} \sqrt{\tau_1/2\pi} = 2e^{K/2} \sqrt{u^{(2)}u^{(3)}/2\pi} = e^{K/2} g_b^{-1} g_c^{-1}$$

$$Y_{abc} = B \sqrt{2\pi g_b g_c} \left[\frac{\Gamma(1-\theta_2)}{\Gamma(\theta_2)} \frac{\Gamma(1-\theta_3)}{\Gamma(\theta_3)} \frac{\Gamma(\frac{1}{2}+\theta_2)}{\Gamma(\frac{1}{2}-\theta_2)} \frac{\Gamma(\frac{1}{2}+\theta_3)}{\Gamma(\frac{1}{2}-\theta_3)} \right]^{1/4}$$

$p_{\text{KK}}/p_{\text{bos/fer}}$	$p_{\text{bos/fer}} = 0$	$p_{\text{bos/fer}} = 2$	$p_{\text{bos/fer}} = 4$	$p_{\text{bos/fer}} = 6$
$p = 2$	EFT IIb), IIc) (PS)	quasi-EFT IIa) (TM)	X	X
$p = 4$	$\mathcal{N} = 2$ sectors	non-EFT (PS)	non-EFT (PS)*	X
$p = 6$	$\mathcal{N} = 2$ sectors	EFT IIa) (PS)	EFT IIa) (PS)*	STU-like II) (TM)

$$\begin{aligned} m_n^{(t)} &= n_t^{1/p_t} m_0^{(t)}, \\ \Lambda_{QG} &= N_t^{1/p_t} m_0^{(t)}, \end{aligned}$$



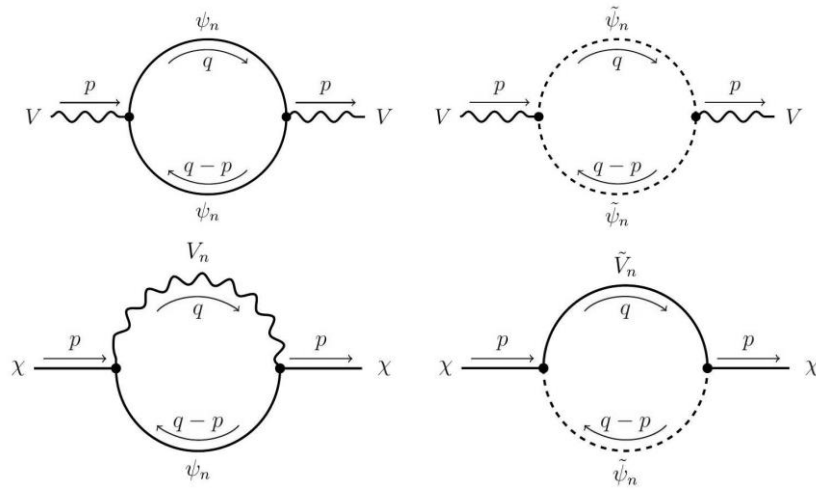
$$\Lambda_{QG} \lesssim \frac{M_P^{(D)}}{N^{1/(D-2)}}$$

$$\Lambda_{QG} \simeq (m_0^{(t)})^{p_t/(D-2+p_t)}$$

$$N_g \simeq \prod_{i=1}^{p_{\text{bos/fer}}/2} \left(\frac{M_s}{m_{\text{bos/fer}}^i} \right)^2$$

$$M_s^{(p_{\text{bos/fer}}+2)} \simeq \prod_{i=1}^{p_{\text{bos/fer}}/2} (m_{\text{bos/fer}}^i)^2$$

$$N_{aa^*} \simeq \prod_i \left(\frac{M_s}{m_{\text{bos/fer}}^i} \right)^2 \simeq k^3 s^2 e^{2\phi_4} \simeq k^3 s$$



$$\delta(g_{\alpha\beta}^{-2}) \simeq q_{\alpha\beta}^2 \sum_{k=1}^{N_g} \log(\Lambda_{QG}^2/m_k^2) \simeq q_{\alpha\beta}^2 N_g = q_{\alpha\beta}^2 \prod_{i=1}^{p_{\text{bos/fer}}/2} \left(\frac{M_s}{m_{\text{bos/fer}}^i} \right)^2$$

$$\delta(g_{\alpha\beta}^{-2}) \simeq q_{\alpha\beta}^2 \prod_{i=1}^{p_{\text{bos/fer}}/2} \frac{1}{\theta_i}$$

$$g_{\alpha\beta}^{-2} \simeq q_{\alpha\beta}^2 N \simeq q_{\alpha\beta}^2 \left(\frac{M_P}{M_s} \right)^2 = q_{\alpha\beta}^2 e^{-2\phi_4},$$

$$g_{\alpha\beta}^{-2} \simeq \prod_{i=1}^{p_{\text{bos/fer}}/2} \left(\frac{M_P}{m_{\text{bos/fer}}^i} \right)^{4/(2+p_{\text{bos/fer}})}$$

$$\delta(g_a^{-2}) \simeq (2)^2 N_g \simeq 4k^3 s$$

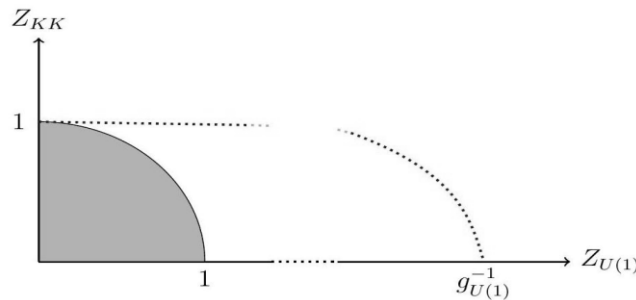
$$\delta(g_a^{-2}) \simeq \left(\frac{M_s}{m_{\text{bos/fer},2}} \right)^2 \simeq ku, \delta(g_a^{-2}) \simeq \left(\frac{M_s}{m_{\text{bos/fer},3}} \right)^2 \simeq ku.$$

$$K_{\alpha\beta} \simeq \prod_{i=1}^{p_{\text{bos/fer}}/2} \left(\frac{M_s}{m_{\text{bos/fer}}^i} \right)^2$$

$$K_{\alpha\beta} \simeq \prod_{i=1}^{p_{\text{bos/fer}}/2} \left(\frac{M_P}{m_{\text{bos/fer}}^i} \right)^{4/(2+p_{\text{bos/fer}})}$$

$$K_{\alpha\beta} \simeq g_{\alpha\beta}^{-2} \simeq N_g \simeq N \simeq \left(\frac{M_P}{M_s} \right)^2 = e^{-2\phi_4}$$

$$\begin{aligned}
m_{KK}(\lambda) &\simeq \frac{1}{\lambda^{w/2}} \\
K_\lambda &\simeq \frac{1}{m_{KK}^2} (\partial_\lambda m_{KK})^2 \simeq \frac{1}{\lambda^2} \\
K_Q(\lambda) &\sim -\log \lambda \\
Y_{ijk} &\simeq \left(\frac{m_i}{M_S}\right)^{\gamma_i} \left(\frac{m_j}{M_S}\right)^{\gamma_j} \left(\frac{m_k}{M_S}\right)^{\gamma_k} \\
Y_{ijk} &\simeq e^{\phi_4/2} \left(\frac{m_{\text{bos/fer}}^i}{M_S}\right)^{1/2} \left(\frac{m_{\text{bos/fer}}^j}{M_S}\right)^{1/2} \left(\frac{m_{\text{bos/fer}}^k}{M_S}\right)^{1/2} \\
Y_{ijk} &\simeq e^{\phi_4/2} \left(\frac{m_{\text{bos/fer}}^i}{M_S}\right)^{1/2} \\
Y_{ijk} &\simeq e^{\phi_4/2} N^{-1/2} \simeq e^{\phi_4} \simeq g_{\alpha\beta}
\end{aligned}$$



$$\alpha' m_{\alpha\beta}^2 = \frac{L^2}{4\pi\alpha'} + N_{\text{bos}}(\theta_{\alpha\beta}^r) + \frac{(r + v_\theta)^2}{2} - \frac{1}{2} + E_{\alpha\beta}$$

$$E_{\alpha\beta} = \sum_{r=1}^3 \frac{1}{2} |\theta_{\alpha\beta}^r| (1 - |\theta_{\alpha\beta}^r|).$$

$$\chi = r + v_\theta = \left(\theta_{\alpha\beta}^1 - \frac{1}{2}, \theta_{\alpha\beta}^2 - \frac{1}{2}, \theta_{\alpha\beta}^3 + \frac{1}{2}, -\frac{1}{2} \right),$$

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$$\chi = (\theta_{\alpha\beta}^1 - 1, \theta_{\alpha\beta}^2, \theta_{\alpha\beta}^3, 0) \quad \alpha' m_{\alpha\beta}^2 = \frac{1}{2} (-|\theta_{\alpha\beta}^1| + |\theta_{\alpha\beta}^2| + |\theta_{\alpha\beta}^3|)$$

$$\chi = (\theta_{\alpha\beta}^1, \theta_{\alpha\beta}^2 - 1, \theta_{\alpha\beta}^3, 0) \quad \alpha' m_{\alpha\beta}^2 = \frac{1}{2} (|\theta_{\alpha\beta}^1| - |\theta_{\alpha\beta}^2| + |\theta_{\alpha\beta}^3|)$$

$$\chi = (\theta_{\alpha\beta}^1, \theta_{\alpha\beta}^2, \theta_{\alpha\beta}^3 + 1, 0) \quad \alpha' m_{\alpha\beta}^2 = \frac{1}{2} (|\theta_{\alpha\beta}^1| + |\theta_{\alpha\beta}^2| - |\theta_{\alpha\beta}^3|)$$

$$\chi = (\theta_{\alpha\beta}^1 + 1, \theta_{\alpha\beta}^2, \theta_{\alpha\beta}^3, 0) \quad \alpha' m_{\alpha\beta}^2 = \frac{1}{2} (3|\theta_{\alpha\beta}^1| + |\theta_{\alpha\beta}^2| + |\theta_{\alpha\beta}^3|)$$

$$\chi = (\theta_{\alpha\beta}^1, \theta_{\alpha\beta}^2 + 1, \theta_{\alpha\beta}^3, 0) \quad \alpha' m_{\alpha\beta}^2 = \frac{1}{2} (|\theta_{\alpha\beta}^1| + 3|\theta_{\alpha\beta}^2| + |\theta_{\alpha\beta}^3|)$$

$$\chi = (\theta_{\alpha\beta}^1, \theta_{\alpha\beta}^2, \theta_{\alpha\beta}^3 - 1, 0) \quad \alpha' m_{\alpha\beta}^2 = \frac{1}{2} (|\theta_{\alpha\beta}^1| + |\theta_{\alpha\beta}^2| + 3|\theta_{\alpha\beta}^3|)$$

$$\left(\alpha - |\theta_{\alpha\beta}^1|\right)^{k_1} \left(\alpha - |\theta_{\alpha\beta}^2|\right)^{k_2} \left(\alpha - |\theta_{\alpha\beta}^3|\right)^{k_3}$$

$$\chi = (\theta_{\alpha\beta}^1, \theta_{\alpha\beta}^2, \theta_{\alpha\beta}^3, \pm 1) \quad \alpha' m_{\alpha\beta}^2 = \frac{1}{2} (|\theta_{\alpha\beta}^1| + |\theta_{\alpha\beta}^2| + |\theta_{\alpha\beta}^3|)$$

$$K_Q = -\log s - \log \frac{1}{6} \kappa_{ijk} u^i u^j u^k$$

$$s = e^{-\phi V_Y}, u^i = e^{-\phi t^i}$$

$$[\Pi_\alpha] = \frac{1}{2} [A] + \frac{1}{2} f_\alpha^i [D_i] - \left(\frac{1}{2} \kappa_{ijk} f_\alpha^j f_\alpha^k + \frac{1}{24} c_{2i} \right) [C^i] + \left(\frac{1}{6} \kappa_{ijk} f_\alpha^i f_\alpha^j f_\alpha^k + \frac{1}{24} c_{2i} f_\alpha^i \right) [B]$$

$$[\Pi_\alpha] \cdot [\Pi_\beta] = \frac{1}{6} \kappa_{ijk} (f_\beta^i - f_\alpha^i) (f_\beta^j - f_\alpha^j) (f_\beta^k - f_\alpha^k) + \frac{1}{12} c_{2i} (f_\beta^i - f_\alpha^i),$$



$$\begin{aligned}
r_i^f + r_j^f + r_k^b &= (0,0,0, \pm 1) \\
\sum_i r_i^b &= (1,1,1,0) \\
\theta_{ab} &= (\theta_{ab}^1, \theta_{ab}^2, \theta_{ab}^3, 0) \\
\theta_{bc} &= (\theta_{bc}^1, \theta_{bc}^2, \theta_{bc}^3, 0) \\
\theta_{ca} &= (\theta_{ca}^1, \theta_{ca}^2, \theta_{ca}^3, 0) \\
\theta_{ab} &= (|\theta_{ab}^1|, |\theta_{ab}^2|, -|\theta_{ab}^3|, 0) \\
\theta_{bc} &= (|\theta_{bc}^1|, -|\theta_{bc}^2|, |\theta_{bc}^3|, 0) \\
\theta_{ca} &= (-|\theta_{ca}^1|, |\theta_{ca}^2|, |\theta_{ca}^3|, 0) \\
r_{ab}^b &= (0,0,1,0), r_{bc}^b = (0,1,0,0), r_{ca}^b = (1,0,0,0), \\
r_{ab}^f &= \frac{1}{2}(-, -, +, -), r_{bc}^f = \frac{1}{2}(-, +, -, -), r_{ca}^f = \frac{1}{2}(+, -, -, -). \\
K_{ab} &= \frac{e^{K/2-\phi_4}}{\sqrt{2\pi}} \left(\frac{\Gamma(|\theta_{ab}^1|)\Gamma(|\theta_{ab}^2|)\Gamma(1-|\theta_{ab}^3|)}{\Gamma(1-|\theta_{ab}^1|)\Gamma(1-|\theta_{ab}^2|)\Gamma(|\theta_{ab}^3|)} \right)^{1/2} \\
K_{bc} &= \frac{e^{K/2-\phi_4}}{\sqrt{2\pi}} \left(\frac{\Gamma(|\theta_{bc}^1|)\Gamma(1-|\theta_{bc}^2|)\Gamma(|\theta_{bc}^3|)}{\Gamma(1-|\theta_{bc}^1|)\Gamma(|\theta_{bc}^2|)\Gamma(1-|\theta_{bc}^3|)} \right)^{1/2} \\
K_{ca} &= \frac{e^{K/2-\phi_4}}{\sqrt{2\pi}} \left(\frac{\Gamma(1-|\theta_{ca}^1|)\Gamma(|\theta_{ca}^2|)\Gamma(|\theta_{ca}^3|)}{\Gamma(|\theta_{ca}^1|)\Gamma(1-|\theta_{ca}^2|)\Gamma(1-|\theta_{ca}^3|)} \right)^{1/2} \\
K_{\alpha\beta} &\simeq \left[\frac{e^{K-2\phi_4} |\theta_{\alpha\beta}^i|}{2\pi |\theta_{\alpha\beta}^j| |\theta_{\alpha\beta}^k|} \right]^{1/2} \\
K_{\alpha\beta} &\simeq e^{K/2} \frac{M_{\text{P}}}{m_{\text{bos/fer}}} \\
\Omega^{(3,0)} &= \frac{1}{3!} \Omega_{\mu\nu\rho} dz^\mu \wedge dz^\nu \wedge dz^\rho \\
\frac{1}{2!} \omega_{\mu\nu\bar{\sigma}} dz^\mu \wedge dz^\nu \wedge dz^{\bar{\sigma}} &\leftrightarrow dz^{\bar{\sigma}} (\omega_{\bar{\sigma}}^\mu = \bar{\Omega}^{\mu\nu\rho} \omega_{\nu\rho\bar{\sigma}}) \partial_\mu \\
\kappa &= \int_X \Omega \wedge \omega^\mu \wedge \omega^\nu \wedge \omega^\rho \Omega_{\mu\nu\rho} \\
\mathbf{248} = \text{Adj}_{E_8} &\rightarrow (\text{Adj}_H, \mathbf{1}) \oplus (\mathbf{1}, \text{Adj}_G) \bigoplus_i (\mathbf{R}_H^{(i)}, \mathbf{R}_G^{(i)}) \\
\kappa(a^{(i)}, b^{(j)}, c^{(k)}) &= \int_X \Omega \wedge \tilde{\Omega}(a^{(i)}, b^{(j)}, c^{(k)}) \\
S_{\text{eff}} = \int d^4x \left[\int d^4\theta K(\Phi^a, \bar{\Phi}^{\bar{b}}) + \frac{1}{4g^2} \left(\int d^2\theta \text{tr} \mathcal{W}_\alpha \mathcal{W}^\alpha + \int d^2\theta W(\Phi^a) \right) + \text{h.c.} \right] \\
K(\Phi^a, \bar{\Phi}^{\bar{b}}) &\supset N_{a\bar{b}} \Phi^a \bar{\Phi}^{\bar{b}} + \dots \\
N_{a\bar{b}} \sim (a, b) &= \int_X a \wedge \bar{\star}_V b \\
Y(a^{(i)'}, b^{(j)'}, c^{(k)'}) &= \frac{\int_X \Omega \wedge \tilde{\Omega}(a^{(i)'}, b^{(j)'}, c^{(k)'})}{\int_X \Omega \wedge \bar{\Omega}} \\
\mathbf{248} &\rightarrow (78, \mathbf{1}) \oplus (1, 8) \oplus (\bar{27}, \bar{3}) \oplus (27, 3) \\
[\alpha] &\mapsto [\Omega(\alpha)] \\
(a, b) &= \int_X a \wedge \bar{\star}_g b \stackrel{\text{def}}{=} \langle a, b \rangle_{\text{WP}} \\
\tilde{\Omega}(a, b, c) &\stackrel{\text{def}}{=} a^\mu \wedge b^\nu \wedge c^\rho \Omega_{\mu\nu\rho}
\end{aligned}$$

$$\begin{aligned}
Y_{ijk} &= \frac{\int_X \Omega \wedge \bar{\Omega}(a_i, a_j, a_k)}{\sqrt{N_i N_j N_k} \int_X \Omega \wedge \bar{\Omega}} \\
\langle a, b \rangle_{\text{WP}} &= \int_{X_t} \rho(a) \wedge \overline{g_t \mathcal{H} \rho(b)}, \\
\rho\left(\frac{\partial}{\partial t}\right) &= \left[\left\{ \frac{\partial f_{jk}^\mu(z_k, t)}{\partial t} \frac{\partial}{\partial z_j^\mu} \right\} \right], \text{ donde } z_k = f_{kj}(z_j, t) \\
\langle a, b \rangle_{\text{WP}} &= - \frac{\int_{X_t} \Omega(\mathcal{H}\rho(a)) \wedge \overline{\Omega(\mathcal{H}\rho(b))}}{\int_{X_t} \Omega \wedge \bar{\Omega}} \int_{X_t} \text{vol}_{g_t} \\
&\quad \zeta_t: X \xrightarrow{\cong} X_t \\
&\quad [\{\bar{\partial}\xi_j\}_j] \in \check{H}^1(X, T_X) \simeq H^1(X, T_X) \\
\left. \frac{d\Omega_t}{dt} \right|_{t=0} &= \Omega' + \Omega(\phi) \in \Gamma(X, \Omega^{n,0}) \oplus \Gamma(X, \Omega^{n-1,1}) \\
(\alpha, \beta) &= \int_X \alpha \wedge \bar{\beta} \\
\frac{\langle a, a \rangle_{\text{WP}}}{\int_X \text{vol}_g} &= \frac{\left(\left. \frac{d\Omega_t}{dt} \right|_{t=0}, \left. \frac{d\Omega_t}{dt} \right|_{t=0} \right)}{(\Omega, \Omega)} + \frac{\left| \left(\Omega, \left. \frac{d\Omega_t}{dt} \right|_{t=0} \right) \right|^2}{(\Omega, \Omega)^2} \\
&\quad \partial\Omega' + \bar{\partial}\Omega' + \partial(\Omega(\phi)) = 0 \\
&\quad \Omega(\phi) + \bar{\partial}\psi = \mathcal{H}\Omega(\phi) \\
\bar{\partial}\Omega' + \partial(\Omega(\phi)) + \partial\bar{\partial}\psi - \bar{\partial}\partial\psi &= \bar{\partial}(\Omega' + \partial\psi) + \partial\mathcal{H}(\Omega(\phi)) = 0 \\
\left(\left. \frac{d\Omega_t}{dt} \right|_{t=0}, \left. \frac{d\Omega_t}{dt} \right|_{t=0} \right) &= (\Omega', \Omega') + (\Omega(\phi), \Omega(\phi)) \\
(\Omega', \Omega') &= (c\Omega - \partial\psi, c\Omega - \partial\psi) = |c|^2(\Omega, \Omega) + (\partial\psi, \partial\psi) \\
(\Omega(\phi), \Omega(\phi)) &= (\mathcal{H}\Omega(\phi) - \bar{\partial}\psi, \mathcal{H}\Omega(\phi) - \bar{\partial}\psi) = (\mathcal{H}\Omega(\phi), \mathcal{H}\Omega(\phi)) + (\bar{\partial}\psi, \bar{\partial}\psi) \\
\left(\left. \frac{d\Omega_t}{dt} \right|_{t=0}, \left. \frac{d\Omega_t}{dt} \right|_{t=0} \right) &= |c|^2(\Omega, \Omega) + (\mathcal{H}\Omega(\phi), \mathcal{H}\Omega(\phi)) \\
\left(\left. \frac{d\Omega_t}{dt} \right|_{t=0}, \Omega \right) &= (\Omega', \Omega) = (c\Omega - \partial\psi, \Omega) = c(\Omega, \Omega) \\
- \frac{\left(\left. \frac{d\Omega_t}{dt} \right|_{t=0}, \left. \frac{d\Omega_t}{dt} \right|_{t=0} \right)}{(\Omega, \Omega)} + \frac{\left| \left(\Omega, \left. \frac{d\Omega_t}{dt} \right|_{t=0} \right) \right|^2}{(\Omega, \Omega)^2} &= - \frac{(\mathcal{H}\Omega(\phi), \mathcal{H}\Omega(\phi))}{(\Omega, \Omega)} \\
X &= \begin{bmatrix} n_1 & q_1^1 & \dots & q_K^1 \\ \vdots & \vdots & \ddots & \vdots \\ n_m & q_1^N & \dots & q_K^N \end{bmatrix} \\
H^1(T_X) &\simeq \bigoplus_{l=1}^K H^0(X, \mathcal{O}_X(q_l)) / \sim \ni F \\
\frac{\partial p_j}{\partial Z^k} \xi^k + F_j &= 0, \text{ for all } j \in \{1, \dots, K\} \\
\text{Jac}(p)^+ &= \overline{\text{Jac}(p)} \cdot (\text{Jac}(p) \cdot \overline{\text{Jac}(p)})^{-1} \\
\xi^k &= (\text{Jac}(p)^+)^{kj} F_j + \sum_i c_i \chi_i^k, \text{ for } c_i \in \mathbb{C} \\
\rho\left(\frac{\partial}{\partial t}\right) &= [\bar{\partial}\xi] \\
\kappa_{ijk} &= \int_X \Omega_{\alpha\beta\gamma} \frac{\partial \xi_i^\alpha}{\partial \bar{z}^\mu} \frac{\partial \xi_j^\beta}{\partial \bar{z}^\nu} \frac{\partial \xi_k^\gamma}{\partial \bar{z}^\delta} \Omega \wedge d\bar{z}^\mu \wedge d\bar{z}^\nu \wedge d\bar{z}^\delta
\end{aligned}$$

$$\begin{array}{ccc} A & \xrightarrow{\check{\zeta}_t} & A_t \\ \uparrow & & \uparrow \\ X & \xrightarrow{\zeta_t} & X_t \end{array}$$

$$\begin{aligned} \Omega_t \wedge dp_t^1 \wedge \dots \wedge dp_t^K &= dZ_t^1 \wedge \dots \wedge dZ_t^M \\ \frac{d}{dt} (dZ_t^1 \wedge \dots \wedge dZ_t^M) \Big|_{t=0} &= \frac{d}{dt} \det \left(\delta_k^j + t \frac{\partial \xi^j}{\partial Z^k} \right) \Big|_{t=0} dZ^1 \wedge \dots \wedge dZ^M + \\ &+ \sum_{j,k=1}^M \frac{\partial \xi^j}{\partial Z^k} dZ^1 \wedge \dots \wedge \widehat{dZ^j} \wedge dZ^k \wedge \dots \wedge dZ^M \\ \frac{d}{dt} (dp_t^j) \Big|_{t=0} &= d \left(\frac{\partial p^j}{\partial Z^k} \xi^k + F^j \right) = 0 \\ \frac{d\Omega_t}{dt} \Big|_{t=0} \wedge dp^1 \wedge \dots \wedge dp^K \Big|_X &= \text{TrJac}(\xi) dZ^1 \wedge \dots \wedge dZ^K \Big|_X + \\ &+ \sum_{j,k=1}^M \frac{\partial \xi^j}{\partial Z^k} dZ^1 \wedge \dots \wedge \widehat{dZ^j} \wedge dZ^k \wedge \dots \wedge dZ^M \Big|_X \end{aligned}$$

$$\alpha = f(Z) dZ^1 \wedge \dots \wedge dZ^n, \beta = \sum_{j,k=1}^n g_k^j(Z) dZ^1 \wedge \dots \wedge \widehat{dZ^j} \wedge \dots \wedge dZ^n \wedge d\bar{Z}^k$$

$$\begin{aligned} f(Z) &= \text{TrJac}(\xi) / \det \left(\frac{\partial p^j}{\partial Z^{n+k}} \right)_{j,k=1}^K \\ g_k^j(Z) &= \frac{(-1)^{n-j}}{\det \left(\frac{\partial p^{j'}}{\partial Z^{n+k'}} \right)_{j',k'=1}^K} \left[\frac{\partial \xi^j}{\partial Z^k} + \sum_{l=1}^K \frac{\partial \xi^j}{\partial Z^{n+l}} \frac{\partial \bar{Z}^{n+l}}{\partial Z^k} \right] \\ i^* d\bar{Z}^{n+l} &= \sum_{k=1}^n \frac{\partial \bar{Z}^{n+l}}{\partial Z^k} dZ^k \\ \int_{X_t} \text{vol}_{g_t} f &\simeq \frac{1}{N} \sum_{k=1}^N f(p_k^{(t)}) \\ x_0^3 + x_1^3 + x_2^3 - 3\psi x_3 x_4 x_5 = 0, x_3^3 + x_4^3 + x_5^3 - 3\psi x_0 x_1 x_2 = 0 \\ \kappa &= \left[\frac{\Gamma(3/5)^3 \Gamma(1/5)}{\Gamma(2/5)^3 \Gamma(4/5)} \right]^{1/2} \approx 1.09236 \end{aligned}$$

$$\mathbb{Z}_5: z_j \rightarrow \alpha^j z_j$$

$$\mathbb{Z}'_5: z_j \rightarrow z_{j+1}$$

$$Y_{5,5,5} = 1.550 \pm 0.002$$

$$X = \begin{bmatrix} 3 & 3 & 1 & 0 \\ 3 & 0 & 1 & 3 \end{bmatrix}$$

$$p^1 = \frac{1}{3}(x_0^3 + x_1^3 + x_2^3 + x_3^3) = 0$$

$$p^2 = x_0 y_0 + x_1 y_1 + x_2 y_2 + x_3 y_3 = 0$$

$$p^3 = \frac{1}{3}(y_0^3 + y_1^3 + y_2^3 + y_3^3) = 0$$



$$\begin{aligned}
& \begin{cases} (x_0, x_1, x_2, x_3) \rightarrow (x_0, \alpha^2 x_1, \alpha x_2, \alpha x_3) \\ (y_0, y_1, y_2, y_3) \rightarrow (y_0, \alpha y_1, \alpha^2 y_2, \alpha^2 y_3) \end{cases} \\
& \kappa_{ijk} = \int_{X/\mathbb{Z}_3} \Omega \wedge \Omega(\lambda_i, \lambda_j, \lambda_k) \\
p_2 &= x_0 y_0 + x_1 y_1 + (1 + \epsilon)(x_2 y_2 + x_3 y_3) = 0, \epsilon \in \mathbb{R} \\
\eta &= \phi + \bar{\partial}_E \varsigma \in \mathcal{H}_{\bar{\partial}}^{0,1}(X, T_X), \varsigma \in C^\infty(T_X) \\
\ell(\theta) &:= (\bar{\partial}_E^\dagger \eta_\theta, \bar{\partial}_E^\dagger \eta_\theta) = \int_{X_t} \bar{\partial}_E^\dagger \eta_\theta \wedge \bar{\kappa}_E \bar{\partial}_E^\dagger \eta_\theta \\
&= \int_{X_t} \text{vol}_{g_t} g_{\mu\bar{\nu}} (\bar{\partial}_E^\dagger \eta_\theta)^\mu \overline{(\bar{\partial}_E^\dagger \eta_\theta)^\nu} \\
&= 25 \frac{\Gamma(4/5)^5 \Gamma(2/5)^5}{\Gamma(1/5)^5 \Gamma(3/5)^5} \approx 0.1922 \\
& \quad (')[\alpha] \mapsto [\Omega(\alpha)] \\
H^1(T_X) &\simeq H^{0,1}(T_X) \simeq H^{n,n-1}(T_X^*)^* \simeq H^{n-1}(\Omega^n \otimes T_X^*)^* \\
H^{n-1}(\Omega^n \otimes T_X^*)^* &\simeq H^{1,n-1}(\mathcal{O}_X)^* \simeq H^{n-1,1}(\mathcal{O}_X^*) \simeq H^{n-1,1}(\mathcal{O}_X) \\
\Omega(\bar{\partial}\phi) &= \sum_{\mu=1}^n (-1)^{\mu-1} f \alpha_{\bar{\nu}}^\mu d\bar{z}^\nu \wedge \hat{\mu} = \sum_{\mu=1}^n (-1)^{\mu-1} f \frac{\partial \phi}{\partial \bar{z}^\nu} d\bar{z}^\nu \wedge \hat{\mu} \\
\bar{\partial}\Omega(\alpha) &= \sum_{\mu=1}^n (-1)^{\mu-1} f \frac{\partial \alpha_{\bar{\nu}}^\mu}{\partial \bar{z}^\sigma} d\bar{z}^\sigma \wedge d\bar{z}^\nu \wedge \hat{\mu} = 0 \\
(')\langle a, b \rangle_{\text{WP}} &= - \frac{\int_{X_t} \Omega(\mathcal{H}\rho(a)) \wedge \overline{\Omega(\mathcal{H}\rho(b))}}{\int_{X_t} \Omega \wedge \bar{\Omega}} \int_{X_t} \text{vol}_{g_t} \\
\Omega(\alpha) &= \sum_{\mu,\nu=1}^n (-1)^{\mu-1} \alpha_{\bar{\nu}}^\mu f d\bar{z}^\nu \wedge \hat{\mu} \\
\Omega(\alpha) \wedge \overline{\Omega(\beta)} &= \sum_{\mu,\mu',\nu,\nu'=1}^n (-1)^{\mu+\mu'} |f|^2 \alpha_{\bar{\nu}}^\mu \overline{\beta_{\bar{\nu}'}^{\mu'}} d\bar{z}^\nu \wedge \hat{\mu} \wedge d\bar{z}^{\nu'} \wedge \widehat{\mu'} \\
\Omega(\alpha) \wedge \overline{\Omega(\beta)} &= \left(\sum_{\mu,\nu=1}^n (-1)^{\mu+\nu} \alpha_{\bar{\nu}}^\mu \overline{\beta_{\bar{\mu}}^\nu} (-1)^{\mu+\nu-1} \right) \Omega \wedge \bar{\Omega} = -\alpha_{\bar{\nu}}^\mu \overline{\beta_{\bar{\mu}}^\nu} \Omega \wedge \bar{\Omega} \\
\int_{X_t} \Omega(\alpha) \wedge \overline{\Omega(\beta)} &= -\frac{1}{\kappa} \int_{X_t} \alpha \wedge \overline{\beta} \\
\partial\Omega(\alpha) &= \partial \left(\sum_{\mu=1}^n (-1)^{\mu-1} f(z) \alpha_{\bar{\nu}}^\mu d\bar{z}^\nu \wedge \hat{\mu} \right) \\
&= \sum_{\mu=1}^n \frac{\partial}{\partial z^\mu} (f(z) \alpha_{\bar{\nu}}^\mu) d\bar{z}^\nu \wedge dz^1 \wedge \dots \wedge dz^n = 0 \\
& \quad g_{\mu\bar{\rho}} g^{\sigma\bar{\nu}} \alpha_{\bar{\nu}}^\mu = \alpha_{\bar{\rho}}^\sigma \\
\frac{\partial}{\partial z^\mu} \left(g_{\sigma\bar{\nu}} g^{\mu\bar{\rho}} \alpha_{\bar{\rho}}^\sigma \det g \right) d\bar{z}^\nu &= \kappa \frac{\partial}{\partial z^\mu} \left(|f(z)|^2 g_{\sigma\bar{\nu}} g^{\mu\bar{\rho}} \alpha_{\bar{\rho}}^\sigma \right) d\bar{z}^\nu \\
&= \kappa \overline{f(z)} \frac{\partial}{\partial z^\mu} \left(f(z) \alpha_{\bar{\nu}}^\mu \right) d\bar{z}^\nu = 0
\end{aligned}$$

5. La brecha de masa y la curvatura en campos y supercampos cuánticos de Yang – Mills. La brecha de masa, que se constituye esencialmente en un cambio de estado de energía de una partícula, respecto del estado de vacío, de tal suerte que, es ésta excitación, es decir, esta modificación de energía

superior a cero, la que puede curvar la geometría del espacio – tiempo cuántico en el que interactúe la partícula excitada.

El modelo matemático, queda expuesto de la siguiente manera:

$$\begin{aligned}
 u \wedge^* v &= \langle u, v \rangle_k d\omega \\
 \langle u, v \rangle &= \int_M u \wedge^* v = \int_M \langle u, v \rangle_k d\omega. \\
 \langle A, B \rangle &= -\text{Tr}_{\text{Mat}(\bar{N}, \mathbb{C})}[AB] \\
 A \cdot \Omega &:= A^\Omega = \Omega^{-1} d\Omega + \Omega^{-1} A \Omega \\
 S_{\text{YM}}(A) &= \int_M |dA + A \wedge A|^2 d\omega \\
 \{dx^0 \wedge dx^1, dx^0 \wedge dx^2, dx^0 \wedge dx^3, dx^1 \wedge dx^2, dx^3 \wedge dx^1, dx^2 \wedge dx^3\} \\
 * (dx^0 \wedge dx^1) &= dx^2 \wedge dx^3, \quad * (dx^0 \wedge dx^2) = dx^3 \wedge dx^1, \quad * (dx^0 \wedge dx^3) = dx^1 \wedge dx^2 \\
 * (dx^2 \wedge dx^3) &= dx^0 \wedge dx^1, \quad * (dx^3 \wedge dx^1) = dx^0 \wedge dx^2, \quad * (dx^1 \wedge dx^2) = dx^0 \wedge dx^3 \\
 \vec{x} \cdot \vec{y} &= -x^0 y^0 + \sum_{i=1}^3 x^i y^i \\
 -(x^0 - y^0)^2 + \sum_{i=1}^3 (x^i - y^i)^2 &\geq 0 \\
 \mathbb{H}_{\text{YM}}(\mathfrak{g}) &= \{1\} \oplus \bigoplus_{n \geq 1} \mathcal{H}(\rho_n) \\
 \mathbb{H}_{\text{YM}}(\mathfrak{g}) &:= \bigoplus_{n=0}^{\infty} \mathcal{H}(\rho_n) \\
 \left\langle \sum_{n=0}^{\infty} v_n, \sum_{n=0}^{\infty} u_n \right\rangle &:= \sum_{n=0}^{\infty} \langle v_n, u_n \rangle \\
 \{(a^0 + sb^0, a^1, a^2 + sb^2, a^3 + tb^3)^T \in \mathbb{R}^4: s, t \in I\} \\
 \inf_{\vec{0} \neq \vec{v} \in TS} \frac{|v|^2}{v^{0,2}} &< 1 \\
 \inf_{\vec{0} \neq \vec{v} \in TS} \frac{|v|^2}{v^{0,2}} &> 1 \\
 \lambda(S, f_\alpha \otimes \rho(E^\alpha), \{\hat{f}_a\}_{a=0}^3) + \mu(\tilde{S}, g_\alpha \otimes \rho(E^\alpha), \{\hat{f}_a\}_{a=0}^3) \\
 \left\langle (S, f_\alpha \otimes \rho(E^\alpha), \{\hat{f}_a\}_{a=0}^3), (\tilde{S}, g_\beta \otimes \rho(E^\beta), \{\hat{g}_a\}_{a=0}^3) \right\rangle \\
 := \int_{S \cap \tilde{S}} [f_\alpha \overline{g_\beta}] \cdot d|\hat{\rho}| \cdot \text{Tr}[-\rho(E^\alpha)\rho(E^\beta)] \\
 \equiv \sum_{\alpha=1}^N C(\rho) \int_{I^2} [f_\alpha \cdot \overline{g_\alpha}](\sigma(\hat{s})) \left| \sum_{0 \leq a < b \leq 3} \hat{\rho}_\sigma^{ab}(\hat{s}) [\det \tilde{J}_{ab}^\sigma(\hat{s})] \right| d\hat{s}, \\
 \{(S_0 + \vec{a}, \rho(E^\alpha), \{e_a\}_{a=0}^3): \vec{a} \in \mathbb{R}\} \\
 \langle (S_0 + \vec{a}, \rho(E^\alpha), \{e_a\}_{a=0}^3), (S_0 + \vec{b}, \rho(E^\alpha), \{e_a\}_{a=0}^3) \rangle = 0 \\
 \{\vec{a}_1, A_1\} \{\vec{a}_2, A_2\} = \{\vec{a}_1 + A_1 \vec{a}_2, A_1 A_2\} \\
 \{\vec{a}_1, \Lambda_1\} \{\vec{a}_2, \Lambda_2\} = \{\vec{a}_1 + Y(\Lambda_1) \vec{a}_2, \Lambda_1 \Lambda_2\} \\
 (S, f_\alpha \otimes \rho(E^\alpha), \{\hat{f}_a\}_{a=0}^3) \mapsto U(\vec{a}, \Lambda) (S, f_\alpha \otimes \rho(E^\alpha), \{\hat{f}_a\}_{a=0}^3), \\
 U(\vec{a}, \Lambda) (S, f_\alpha \otimes \rho(E^\alpha), \{\hat{f}_a\}_{a=0}^3) \\
 := (\Lambda S + \vec{a}, e^{-i[\vec{a} \cdot \hat{H}(\rho_n) \Lambda \hat{f}_0 + \hat{P}(\rho_n) \Lambda \hat{f}_1]} f_\alpha(\Lambda^{-1}(\cdot - \vec{a})) \otimes \rho(E^\alpha), \{\Lambda \hat{f}_a\}_{a=0}^3)
 \end{aligned}$$

$$\begin{aligned}
& e^{-i[\vec{a} \cdot (\hat{H}(\rho_n)Y(\Lambda)\hat{f}_0 + \hat{P}(\rho_n)Y(\Lambda)\hat{f}_1)]} f_\alpha [Y(\Lambda^{-1})(\vec{x} - \vec{a})] \otimes \rho(E^\alpha) \\
& \equiv e^{-i[\vec{a} \cdot (\hat{H}(\rho_n)Y(\Lambda)\hat{f}_0 + \hat{P}(\rho_n)Y(\Lambda)\hat{f}_1)]} f_\alpha [\sigma(\hat{s})] \otimes \rho(E^\alpha) \\
& U(\vec{a}, 1) \left(S, f_\alpha \otimes \rho(E^\alpha), \{\hat{f}_a\}_{a=0}^3 \right) \\
& \equiv e^{-i[\vec{a} \cdot (\hat{H}(\rho_n)\hat{f}_0 + \hat{P}(\rho_n)\hat{f}_1)]} \left(S + \vec{a}, f_\alpha(\cdot - \vec{a}) \otimes \rho(E^\alpha), \{\hat{f}_a\}_{a=0}^3 \right) \\
& \left(U(\vec{a}, \Lambda) \left(S, f_\alpha \otimes \rho(E^\alpha), \{\hat{f}_a\}_{a=0}^3 \right), U(\vec{a}, \Lambda) \left(\tilde{S}, g_\beta \otimes \rho(E^\beta), \{\hat{f}_a\}_{a=0}^3 \right) \right) \\
& := \int_{[\Lambda S + \vec{a}] \cap [\Lambda \tilde{S} + \vec{a}]} d|\hat{\rho}| e^{-i[\vec{a} \cdot (\hat{H}(\rho_n)\Lambda\hat{f}_0 + \hat{P}(\rho_n)\Lambda\hat{f}_1)]} e^{i[\vec{a} \cdot (\hat{H}(\rho_n)\Lambda\hat{f}_0 + \hat{P}(\rho_n)\Lambda\hat{f}_1)]} \\
& \times [f_\alpha \overline{g_\beta}] (\Lambda^{-1}(\cdot - \vec{a})) \cdot \text{Tr}[-\rho(E^\alpha)\rho(E^\beta)] \\
& = \int_{\Lambda(S \cap \tilde{S}) + \vec{a}} [f_\alpha \overline{g_\beta}] (\Lambda^{-1}(\cdot - \vec{a})) \cdot d|\hat{\rho}| \cdot \text{Tr}[-\rho(E^\alpha)\rho(E^\beta)] \\
& = \int_{S \cap \tilde{S}} [f_\alpha \overline{g_\beta}] (\cdot) \cdot d|\hat{\rho}| \cdot \text{Tr}[-\rho(E^\alpha)\rho(E^\beta)] \\
& \sum_{u=1}^{\infty} \left(S_u, f_\alpha^u \otimes \rho(E^\alpha), \{\hat{f}_a^u\}_{a=0}^3 \right) \\
& \sum_{u=1}^{\infty} \left(S, f_\alpha^u \otimes \rho_n(E^\alpha), \{\hat{f}_a^u\}_{a=0}^3 \right) \{ \phi^{\alpha, n}(g): 1 \leq \alpha \leq N \} \{ \hat{H}(\rho_n), \hat{P}(\rho_n) \}_{n \geq 1}, \vec{a} = a^0 \hat{f}_0 + a^1 \hat{f}_1. \\
& \{ \hat{H}(\rho), \hat{P}(\rho) \} \hat{H}(\rho)^2 - \hat{P}(\rho)^2 = m(\rho)^2, m(\rho) \geq 0.1, m_0 := \inf_{n \in \mathbb{N}} m(\rho_n) > 0 \\
& \frac{\hat{P}(\rho_n)^2}{\hat{H}(\rho_n)^2} - 1 = -\frac{m(\rho_n)^2}{\hat{H}(\rho_n)^2} \\
& 0 > \frac{\hat{P}(\rho_n)^2}{\hat{H}(\rho_n)^2} - 1 \rightarrow 0 \\
& \oplus_{n=0}^{\infty} \mathcal{H}(\rho_n), \text{SL}(2, \mathbb{C}) D^{(j,k)}, j, k D^{(j,k)}(-1) = (-1)^{2(j+k)} \text{SU}(2), \Lambda \in \text{SU}(2) \mapsto D^{(j,0)}(\Lambda) \\
& \hat{f}_a = \Lambda e_a \equiv Y(\Lambda) e_a \{ \hat{g}_a \}_{a=0}^3 \hat{f}_a = \hat{g}_a Y(\Lambda) e_a = Y(\tilde{\Lambda}) e_a \tilde{\Lambda} \in \text{SL}(2, \mathbb{C}). Y(\Lambda) = Y(\tilde{\Lambda}), \Lambda = \pm \tilde{\Lambda}. j + k, \\
& D^{(j,k)}(\pm 1) = 1, D^{(j,k)} \\
& \vec{k} = (k^0, k^1, k^2, k^3), k^a \in \{0\} \cup \mathbb{N} \\
& D^{\vec{k}} = \left(\frac{\partial}{\partial x^0} \right)^{k_0} \left(\frac{\partial}{\partial x^1} \right)^{k_1} \left(\frac{\partial}{\partial x^2} \right)^{k_2} \left(\frac{\partial}{\partial x^3} \right)^{k_3}, \vec{x}^{\vec{k}} = (x^0)^{k_0} (x^1)^{k_1} (x^2)^{k_2} (x^3)^{k_3} \\
& |\vec{k}| = \sum_{a=0}^3 |k^a| \\
& \|f\|_{r,s} := \sum_{|\vec{k}| \leq r} \sum_{|\vec{l}| \leq s} \sup_{\vec{x} \in \mathbb{R}^4} |\vec{x}^{\vec{k}} D^{\vec{l}} f(\vec{x})|. \\
& \hat{s} = (s, \bar{s}) \mapsto \sigma(s, \bar{s}) = s \hat{f}_0 + \bar{s} \hat{f}_1, s, \bar{s} \in \mathbb{R} \\
& d\hat{s} = ds d\bar{s}, \tilde{\eta}: \vec{v} \in \mathbb{R}^4 \mapsto \vec{v}, (\hat{H}(\rho_n)\hat{f}_0 + \hat{P}(\rho_n)\hat{f}_1) \in \mathbb{R}, \tilde{f} \{ \hat{f}_0, \hat{f}_1 \}: \mathbb{R}^4 \rightarrow \mathbb{C} \\
& \vec{x} \in \mathbb{R}^4 \mapsto \tilde{f}^{\{ \hat{f}_0, \hat{f}_1 \}} (\hat{H}(\rho_n), \hat{P}(\rho_n)) (\vec{x}) := \int_{S^b} \frac{e^{-i\tilde{\eta}(\cdot)}}{2\pi} \tilde{f}(\vec{x} + \cdot) d|\hat{\rho}| \\
& = \int_{S \in \mathbb{R}^2} \frac{e^{-i[\sigma(\hat{s}) \cdot (\hat{H}(\rho_n)\hat{f}_0 + \hat{P}(\rho_n)\hat{f}_1)]}}{2\pi} \tilde{f}(\vec{x} + \sigma(\hat{s})) \cdot |\hat{\rho}_\sigma|(\hat{s}) d\hat{s} \\
& \tilde{f}^{\{ \hat{f}_0, \hat{f}_1 \}} \notin \mathcal{P}, \tilde{f} \equiv 0. \\
& \hat{H}(\rho_n)\hat{f}_0 + \hat{P}(\rho_n)\hat{f}_1 S^b. \\
& \vec{x} = \sum_{a=0}^3 x^a \hat{f}_a \\
& \tilde{f}^{\{ \hat{f}_0, \hat{f}_1 \}} (\hat{H}(\rho_n), \hat{P}(\rho_n)) (\vec{x}) \\
& = e^{-i[x^0 \hat{H}(\rho_n) - x^1 \hat{P}(\rho_n)]} \tilde{f}^{\{ \hat{f}_0, \hat{f}_1 \}} (\hat{H}(\rho_n), \hat{P}(\rho_n)) (x^2 \hat{f}_2 + x^3 \hat{f}_3)
\end{aligned}$$

$$\begin{aligned}
\phi^{\alpha,n}(\tilde{f})1 &:= (S_0, \tilde{f}^{\{e_0, e_1\}} \otimes \rho_n(F^\alpha), \{e_a\}_{a=0}^3) \\
&\equiv (S_0, \tilde{f}^{\{e_0, e_1\}} (\hat{H}(\rho_n), \hat{P}(\rho_n)) \otimes \rho_n(F^\alpha), \{e_a\}_{a=0}^3) \in \mathcal{H}(\rho_n), \\
\left\{ a_0 1 + \sum_{n,u=1}^{\infty} (S_{n,u}, f_{n,\alpha}^u \otimes \rho_n(E^\alpha), \{\hat{f}_a^{n,u}\}_{a=0}^3) : a_0 \in \mathbb{C}, f_{n,\alpha}^u \in \mathcal{P}_{S_{n,u}}, S_{n,u} \in \mathcal{S} \right\} \\
&\int_{I^2} |f \circ \sigma|^2(\hat{s}) \left| \sum_{0 \leq a < b \leq 3} \hat{\rho}_{\sigma}^{ab}(\hat{s}) [\det \hat{f}_{ab}^{\sigma}(\hat{s})] \right| d\hat{s} < \infty \\
\phi^{\alpha,n}(f) &\sum_{u=1}^{\infty} (S_u, g_\beta^u \otimes \rho_m(E^\beta), \{\hat{f}_a^u\}_{a=0}^3) \\
&:= \begin{cases} \sum_{u=1}^{\infty} (S_u, f_n^{\{\hat{f}_0^u, \hat{f}_1^u\}} A(\Lambda^u)_\gamma^\alpha \cdot g_\beta^u \otimes \rho_n([F^\gamma, E^\beta]), \{\hat{f}_a^u\}_{a=0}^3), & m = n \\ 0, & m \neq n \end{cases} \\
&\{\phi^{\alpha,n}(f) : 1 \leq \alpha \leq \underline{N}\} \\
\phi^{\alpha,n}(f) \sum_{m=0}^{\infty} v_m &:= \sum_{m=0}^{\infty} \phi^{\alpha,n}(f) v_m \\
&= a_0 (S_0, f_n^{\{e_0, e_1\}} \otimes \rho_n(E^\alpha), \{e_a\}_{a=0}^3) + \phi^{\alpha,n}(f) v_n \\
\phi^{\alpha,n}(f) (S, g_\beta \otimes \rho_n(E^\beta), \{\hat{f}_a\}_{a=0}^3) \\
&:= (S, f_n^{\{\hat{f}_0, \hat{f}_1\}} \cdot A(\Lambda)_\gamma^\alpha g_\beta \otimes \rho_n([F^\gamma, E^\beta]), \{\hat{f}_a\}_{a=0}^3) \\
\phi^{\alpha,n}(g)^* (S, f_\beta \otimes \rho_n(E^\beta), \{\hat{f}_a\}_{a=0}^3) \\
&= - (S, \overline{g^{\{\hat{f}_0, \hat{f}_1\}}} \overline{A(\Lambda)_\gamma^\alpha} \cdot f_\beta \otimes \rho_n([F^\gamma, E^\beta]), \{\hat{f}_a\}_{a=0}^3) \\
&\quad + \left\langle (S, f_\beta \otimes \rho_n(E^\beta), \{\hat{f}_a\}_{a=0}^3), \phi^{\alpha,n}(g) 1 \right\rangle 1, \\
\langle \text{ad}(\rho(E^\alpha)) \rho(E^\beta), \rho(E^\gamma) \rangle &= - \langle \rho(E^\beta), \text{ad}(\rho(E^\alpha)) \rho(E^\gamma) \rangle. \\
\lambda_1 &= \begin{pmatrix} 0 & 1 & 0 \\ -1 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}, \quad \lambda_2 = \begin{pmatrix} 0 & i & 0 \\ i & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}, \quad \lambda_3 = \begin{pmatrix} i & 0 & 0 \\ 0 & -i & 0 \\ 0 & 0 & 0 \end{pmatrix} \\
\lambda_4 &= \begin{pmatrix} 0 & 0 & i \\ 0 & 0 & 0 \\ i & 0 & 0 \end{pmatrix}, \quad \lambda_5 = \begin{pmatrix} 0 & 0 & 1 \\ 0 & 0 & 0 \\ -1 & 0 & 0 \end{pmatrix}, \\
\lambda_6 &= \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & i \\ 0 & i & 0 \end{pmatrix}, \quad \lambda_7 = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 1 \\ 0 & -1 & 0 \end{pmatrix}, \quad \lambda_8 = \frac{1}{\sqrt{3}} \begin{pmatrix} i & 0 & 0 \\ 0 & i & 0 \\ 0 & 0 & -2i \end{pmatrix} \\
[\lambda_7, \lambda_4] &= \lambda_2, [\lambda_7, \lambda_5] = \lambda_1, \text{ad}(\lambda_4) \text{ad}(\lambda_5) \lambda_7 = -\lambda_6 \\
\text{ad}(\lambda_7) \text{ad}(\lambda_4) \text{ad}(\lambda_5) \lambda_7 &= \lambda_3 - \sqrt{3} \lambda_8, [\lambda_5, \lambda_4] = \lambda_3 + \sqrt{3} \lambda_8 \\
\check{\varphi}_\delta(t) &:= \begin{cases} \frac{1}{\delta} t, & 0 \leq t \leq \delta \\ 1, & \delta < t \leq 1 - \delta \\ \frac{1}{\delta} (1 - t), & 1 - \delta < t \leq 1 \end{cases} \\
\|f - \check{g}_\epsilon\|_{L^2} &= \left[\int_{I^2} |1 - g_\delta^{n-1}|^2(\hat{t}) |f|^2(\hat{t}) d\hat{t} \right]^{1/2} \\
&\leq M \|1 - g_\delta^{n-1}\|_{L^2} < \epsilon \\
\psi^{\alpha_1, m_1}(g_1) \dots \psi^{\alpha_k, m_k}(g_k) \mathcal{D} &\subset \mathcal{D} \\
\text{span}\{\text{ad}(F^{\alpha_1}) \dots \text{ad}(F^{\alpha_k}) F^\beta : 1 \leq \beta \leq \underline{N}, 1 \leq \alpha_i \leq \underline{N}, i = 1, 2, \dots, k\} &= \mathfrak{g}.
\end{aligned}$$

$$E^\gamma = \sum_{\beta=1}^{N(\gamma)} d_{m,\beta}^\gamma \text{ad}(F^{\alpha_1^{\gamma,\beta}}) \dots \text{ad}(F_{m-1}^{\alpha_{m-1}^{\gamma,\beta}}) F_m^{\alpha_m^{\gamma,\beta}}$$

$$F_i \in \mathcal{P}: \vec{x} \in \mathbb{R}^4 \mapsto \frac{1}{c_n} p_1(x^0) \frac{1}{d_n} p_1(x^1) f_i(x^2, x^3), i = 1, \dots, m.$$

$$F_i^{\{e_0, e_1\}}(\hat{H}(\rho_n), \hat{P}(\rho_n))(0, 0, x^2, x^3) = \frac{1}{c_n} \hat{p}_1(\hat{H}(\rho_n)) \frac{1}{d_n} \hat{p}_1(\hat{P}(\rho_n)) f_i(x^2, x^3) = f_i(x^2, x^3)$$

$$[F_1^{\{e_0, e_1\}} \dots F_m^{\{e_0, e_1\}}](\hat{H}(\rho_n), \hat{P}(\rho_n))(0, 0, x^2, x^3) = \prod_{i=1}^m f_i(x^2, x^3)$$

$$\sum_{\beta=1}^{N(\gamma)} d_{m,\beta}^\gamma \phi^{\alpha_1^{\gamma,\beta}, n}(F_1) \dots \phi^{\alpha_m^{\gamma,\beta}, n}(F_m) 1 = \left(S, \prod_{i=1}^m f_i \otimes \rho_n(E^\gamma), \{e_a\}_{a=0}^3 \right)$$

$$C_m = \text{span}\{\phi^{\alpha_1, n}(F_1) \dots \phi^{\alpha_m, n}(F_m) 1: F_i \in \mathcal{P}, 1 \leq \alpha_i \leq \underline{N}\}$$

$$\text{span}\{\text{ad}(F^{\alpha_1}) \dots \text{ad}(F^{\alpha_{\tilde{n}-1}}) F^{\alpha_{\tilde{n}}}: 1 \leq \alpha_i \leq \underline{N}, 1 \leq \tilde{n} \leq \underline{n}\} = g$$

$$\tilde{E}^\gamma = \sum_{\xi=1}^{N(\gamma)} d_{\tilde{n}}(\gamma, \xi) \text{ad}(F^{\alpha_1(\gamma, \xi)}) \dots \text{ad}(F^{\alpha_{\tilde{n}-1}(\gamma, \xi)}) F^{\alpha_{\tilde{n}}(\gamma, \xi)}$$

$$\|f - g_1 \dots g_{\tilde{n}}\|_{L^2} < \frac{\epsilon}{C(\rho_n)}$$

$$\sum_{\xi=1}^{N(\gamma)} d_{\tilde{n}}(\gamma, \xi) \phi^{\alpha_1(\gamma, \xi), n}(G_1) \dots \phi^{\alpha_{\tilde{n}-1}(\gamma, \xi), n}(G_{\tilde{n}-1}) \phi^{\alpha_{\tilde{n}}(\gamma, \xi), n}(G_{\tilde{n}}) 1$$

$$= \left(S, \prod_{i=1}^{\tilde{n}} g_i \otimes \rho_n(\tilde{E}^\gamma), \{e_a\}_{a=0}^3 \right)$$

$$\left| \left(S, f \otimes \rho(\tilde{E}^\gamma), \{e_a\}_{a=0}^3 \right) - \left(S, \prod_{i=1}^{\tilde{n}} g_i \otimes \rho_n(\tilde{E}^\gamma), \{e_a\}_{a=0}^3 \right) \right| < \epsilon$$

$$c_\alpha^{\gamma, \beta} = A(\Lambda)_\delta^\alpha \text{Tr} \left[- \left[\text{ad}(\rho_n(F^\delta)) \rho_n(E^\gamma) \right] \rho_n(E^\beta) \right]$$

$$T(f) := \left\langle \phi^{\alpha, n}(f) \left(\tilde{S}, \tilde{g}_\gamma \otimes \rho_n(E^\gamma), \{\hat{f}_a\}_{a=0}^3 \right), \left(S, g_\beta \otimes \rho_n(E^\beta), \{\hat{f}_a\}_{a=0}^3 \right) \right\rangle$$

$$= c_\alpha^{\gamma, \beta} \int_{I^2} d\hat{t} \left[f^{\{\hat{f}_0, \hat{f}_1\}} \cdot \tilde{g}_\gamma \cdot \overline{g}_\beta \right] (\sigma(\hat{t})) \cdot |\hat{\rho}_\sigma|(\hat{t})$$

$$h = [\tilde{g}_\gamma \cdot \overline{g}_\beta] \circ \sigma \cdot |\hat{\rho}_\sigma| c_\alpha^{\gamma, \beta}.$$

$$T(f) = \int_{I^2} f^{\{\hat{f}_0, \hat{f}_1\}}(\sigma(\hat{t})) h(\hat{t}) d\hat{t}$$

$$\vec{x} \mapsto f^{\{\hat{f}_0, \hat{f}_1\}}(\vec{x}) \equiv f^{\{\hat{f}_0, \hat{f}_1\}}(\hat{H}(\rho_n), \hat{P}(\rho_n))(\vec{x})$$

$$= \int_{\hat{s} \in \mathbb{R}^2} \frac{e^{-i[\tilde{\sigma}(\hat{s}) \cdot (\hat{H}(\rho_n) \hat{f}_0 + \hat{P}(\rho_n) \hat{f}_1)]}}{2\pi} f(\vec{x} + \tilde{\sigma}(\hat{s})) |\hat{\rho}_{\tilde{\sigma}}|(\hat{s}) d\hat{s}$$

$$T(f) = \int_{I^2} \{ \hat{f}_0, \hat{f}_1 \}(\sigma(\hat{t})) h(\hat{t}) d\hat{t}$$

$$= \int_{\hat{s} \in \mathbb{R}^2, \hat{t} \in I^2} d\hat{t} d\hat{s} f(\sigma(\hat{t}) + \tilde{\sigma}(\hat{s})) |\hat{\rho}_{\tilde{\sigma}}|(\hat{s}) |\hat{\rho}_\sigma|(\hat{t}) \cdot \frac{e^{-i[\tilde{\sigma}(\hat{s}) \cdot \vec{\alpha}]}}{2\pi} [\tilde{g}_\gamma \cdot \overline{g}_\beta] \circ \sigma(\hat{t}) \cdot c_\alpha^{\gamma, \beta}$$

$$(\hat{s}, \hat{t}) \in \mathbb{R}^2 \times I^2 \mapsto |\hat{\rho}_{\tilde{\sigma}}|(\hat{s}) |\hat{\rho}_\sigma|(\hat{t}) \cdot \frac{e^{-i[\tilde{\sigma}(\hat{s}) \cdot \vec{\alpha}]}}{2\pi} [\tilde{g}_\gamma \cdot \overline{g}_\beta] \circ \sigma(\hat{t}) \cdot c_\alpha^{\gamma, \beta},$$

$$\begin{aligned}
& f \in \mathcal{P} \\
& \mapsto \int_{\hat{s} \in \mathbb{R}^2, \hat{t} \in I^2} d\hat{t} d\hat{s} f(\sigma(\hat{t}) + \tilde{\sigma}(\hat{s})) |\hat{\rho}_{\hat{\sigma}}|(\hat{s}) |\hat{\rho}_{\sigma}|(\hat{t}) \cdot \frac{e^{-i[\tilde{\sigma}(\hat{s}) \cdot \vec{a}]}}{2\pi} [\tilde{g}_\gamma \cdot \overline{g}_\beta] \circ \sigma(\hat{t}) \\
& \times \text{Tr}[-[\text{ad}(\cdot)\rho_n(E^\gamma)]\rho_n(E^\beta)] \\
T(f) & \\
& = \int_{\mathbb{R}^4} f(\vec{x}) \frac{e^{i[x^0 \hat{H}(\rho_n) - x^1 \hat{P}(\rho_n)]}}{2\pi} [\tilde{g}_\gamma \cdot \overline{g}_\beta](x^2, x^3) d\vec{x} \cdot \text{Tr}[-[\text{ad}(\rho_n(F^\alpha))\rho_n(E^\gamma)]\rho_n(E^\beta)] \\
& \quad \vec{x} = (x^0, x^1, x^2, x^3) \mapsto \frac{e^{i[x^0 \hat{H}(\rho_n) - x^1 \hat{P}(\rho_n)]}}{2\pi} [\tilde{g}_\gamma \cdot \overline{g}_\beta](x^2, x^3), \\
& \quad U(\vec{a}, \Lambda) \phi^{\alpha, n}(f) U(\vec{a}, \Lambda)^{-1} (S, g_\beta \otimes \rho_n(E^\beta), \{\hat{f}_a\}_{a=0}^3) \\
& \quad = A(\Lambda^{-1})_\gamma^\alpha \phi^{\gamma, n}(f(\Lambda^{-1}(\cdot - \vec{a}))) (S, g_\beta \otimes \rho_n(E^\beta), \{\hat{f}_a\}_{a=0}^3) \\
& U(\vec{a}, \Lambda) \phi^{\alpha, n}(f) U(\vec{a}, \Lambda)^{-1} 1 \\
& = (\Lambda S_0 + \vec{a}, e^{-i[\vec{a} \cdot (\hat{H}(\rho_n) \hat{g}_0 + \hat{P}(\rho_n) \hat{g}_1)]} f^{\{\hat{g}_0, \hat{g}_1\}}(\Lambda^{-1}(\cdot - \vec{a})) \otimes \rho_n(F^\alpha), \{\hat{g}_a\}_{a=0}^3) \\
& \quad T(\rho_n, \vec{a}) = e^{-i[\vec{a} \cdot (\hat{H}(\rho_n) \hat{f}_0 + \hat{P}(\rho_n) \hat{f}_1)]}, T(\rho_n, \vec{a})^{-1} = e^{i[\vec{a} \cdot (\hat{H}(\rho_n) \hat{f}_0 + \hat{P}(\rho_n) \hat{f}_1)]} \\
U(\vec{a}, \Lambda) \phi^{\alpha, n}(f) U(\vec{a}, \Lambda)^{-1} (S, g_\beta \otimes \rho_n(E^\beta), \{\hat{f}_a\}_{a=0}^3) \\
& = U(\vec{a}, \Lambda) \phi^{\alpha, n}(f) (\Lambda^{-1}(S - \vec{a}), T(\rho_n, \vec{a})^{-1} g_\beta(\Lambda \cdot + \vec{a}) \otimes \rho_n(E^\beta), \mathcal{D}) \\
& = U(\vec{a}, \Lambda) (\Lambda^{-1}(S - \vec{a}), [T(\rho_n, \vec{a})^{-1} f^e](\cdot) d_\gamma^\alpha \cdot g_\beta(\Lambda \cdot + \vec{a}) \otimes \text{ad}(\rho_n(F^\gamma)) \rho_n(E^\beta), \mathcal{D}) \\
& = (S, T(\rho_n, \vec{a}) T(\rho_n, \vec{a})^{-1} f^e(\Lambda^{-1}(\cdot - \vec{a})) d_\gamma^\alpha g_\beta(\cdot) \otimes \text{ad}(\rho_n(F^\gamma)) \rho_n(E^\beta), \{\hat{f}_a\}_{a=0}^3). \\
& \quad f^e(\Lambda^{-1}(\vec{x} - \vec{a})) \equiv f^{\{\Lambda^{-1} \hat{f}_0, \Lambda^{-1} \hat{f}_1\}}(\Lambda^{-1}(\vec{x} - \vec{a})), \\
& \quad f^{\{\hat{g}_0, \hat{g}_1\}}(\hat{H}(\rho_n), \hat{P}(\rho_n))(\Lambda^{-1}(\vec{x} - \vec{a})) \\
& := \int_{\hat{s} \in \mathbb{R}^2} \frac{e^{-i[\hat{\sigma}(\hat{s}) \cdot (\hat{H}(\rho_n) \hat{g}_0 + \hat{P}(\rho_n) \hat{g}_1)]}}{2\pi} f(\vec{y} + \hat{\sigma}(\hat{s})) |\hat{\rho}_{\hat{\sigma}}|(\hat{s}) d\hat{s} \\
& = \int_{\hat{s} \in \mathbb{R}^2} \frac{e^{-i[\sigma(\hat{s}) \cdot (\hat{H}(\rho_n) \hat{f}_0 + \hat{P}(\rho_n) \hat{f}_1)]}}{2\pi} f(\vec{y} + \Lambda^{-1} \sigma(\hat{s})) |\hat{\rho}_\sigma|(\hat{s}) d\hat{s} \\
& = \int_{\hat{s} \in \mathbb{R}^2} \frac{e^{-i[\sigma(\hat{s}) \cdot (\hat{H}(\rho_n) \hat{f}_0 + \hat{P}(\rho_n) \hat{f}_1)]}}{2\pi} f(\Lambda^{-1}(\vec{x} + \sigma(\hat{s}) - \vec{a})) |\hat{\rho}_\sigma|(\hat{s}) d\hat{s} \\
& = f(\Lambda^{-1}(\cdot - \vec{a}))^{\{\hat{f}_0, \hat{f}_1\}}(\hat{H}(\rho_n), \hat{P}(\rho_n))(\vec{x}) \\
& (S, f(\Lambda^{-1}(\cdot - \vec{a}))^{\{\hat{f}_0, \hat{f}_1\}} A(\Lambda^{-1})_\gamma^\alpha A(\vec{\Lambda})_\delta^\gamma \cdot g_\beta \otimes \text{ad}(\rho_n(F^\delta)) \rho_n(E^\beta), \{\tilde{\Lambda} e_a\}_{a=0}^3) \\
& = A(\Lambda^{-1})_\gamma^\alpha \phi^{\gamma, n}[f(\Lambda^{-1}(\cdot - \vec{a}))] (S, g_\beta \otimes \rho_n(E^\beta), \{\tilde{\Lambda} e_a\}_{a=0}^3) \\
U(\vec{a}, \Lambda) \phi^{\alpha, n}(f) U(\vec{a}, \Lambda)^{-1} 1 = U(\vec{a}, \Lambda) \phi^{\alpha, n}(f) 1 \\
& = U(\vec{a}, \Lambda) (S_0, f^{\{e_0, e_1\}} \otimes \rho_n(F^\alpha), \{e_a\}_{a=0}^3) \\
& = (\Lambda S_0 + \vec{a}, e^{-i[\vec{a} \cdot (\hat{H}(\rho_n) \hat{g}_0 + \hat{P}(\rho_n) \hat{g}_1)]} f^{\{e_0, e_1\}}(\Lambda^{-1}(\cdot - \vec{a})) \otimes \rho_n(F^\alpha), \{\hat{g}_a\}_{a=0}^3) \\
& \quad f(\cdot)^{\{\hat{g}_0, \hat{g}_1\}} \mapsto f(\cdot - \Lambda \vec{a})^{\{\hat{g}_0, \hat{g}_1\}}, \\
& (\vec{x} - \vec{y}) \cdot (\vec{x} - \vec{y}) = -(x^0 - y^0)^2 + \sum_{i=1}^3 (x^i - y^i)^2 \leq (\geq) 0. \\
& (S, f_\alpha \otimes \rho(E^\alpha)) \equiv (S, f_\alpha \otimes \rho(E^\alpha), \{\hat{f}_a\}_{a=0}^3) \\
& [\phi^{\alpha, n}(f), \phi^{\beta, n}(g)] := \phi^{\alpha, n}(f) \phi^{\beta, n}(g) - \phi^{\alpha, n}(g) \phi^{\beta, n}(f), \\
& [\phi^{\alpha, n}(f)^*, \phi^{\beta, n}(g)^*] := \phi^{\alpha, n}(f)^* \phi^{\beta, n}(g)^* - \phi^{\alpha, n}(g)^* \phi^{\beta, n}(f)^*. \\
& [\phi^{\alpha, n}(f), \phi^{\beta, n}(g)] (S, h_\gamma \otimes \rho_n(E^\gamma)) = 0, \\
& [\phi^{\alpha, n}(f)^*, \phi^{\beta, n}(g)^*] 1 = 0, [\phi^{\alpha, n}(f)^*, \phi^{\beta, n}(g)^*] (S, h_\gamma \otimes \rho_n(E^\gamma)) = 0,
\end{aligned}$$

$$\begin{aligned}
& [\phi^{\alpha,n}(f), \phi^{\beta,n}(g)^*]_{\pm} := \phi^{\alpha,n}(f)\phi^{\beta,n}(g)^* \pm \phi^{\alpha,n}(g)\phi^{\beta,n}(f)^* \\
& [\phi^{\alpha,n}(f)^*, \phi^{\beta,n}(g)]_{\pm} := \phi^{\alpha,n}(f)^*\phi^{\beta,n}(g) \pm \phi^{\alpha,n}(g)^*\phi^{\beta,n}(f) \\
& [\phi^{\alpha,n}(f), \phi^{\beta,n}(g)^*]_{\pm} (S, h_{\gamma} \otimes \rho_n(E^{\gamma})) \\
& = -A(\Lambda)_{\delta}^{\alpha} \overline{A(\Lambda)_{\mu}^{\beta}} (S, B^{\pm}[f^e \cdot \overline{g^e} \pm g^e \cdot \overline{f^e}] \cdot h_{\gamma} \otimes \text{ad}(\rho_n(F^{\delta})) \text{ad}(\rho_n(F^{\mu}))\rho_n(E^{\gamma})) \\
& + \left\langle (S, h_{\gamma} \otimes \rho_n(E^{\gamma})), \phi^{\beta,n}(g)1 \right\rangle \phi^{\alpha,n}(f)1 \\
& \pm \left\langle (S, h_{\gamma} \otimes \rho_n(E^{\gamma})), \phi^{\beta,n}(f)1 \right\rangle \phi^{\alpha,n}(g)1 \\
& [\phi^{\alpha,n}(f)^*, \phi^{\beta,n}(g)]_{\pm} (S, h_{\gamma} \otimes \rho_n(E^{\gamma})) \\
& = -\overline{A(\Lambda)_{\delta}^{\alpha}} A(\Lambda)_{\mu}^{\beta} (S, B^{\pm}[\overline{f^e} \cdot g^e \pm \overline{g^e} \cdot f^e] \cdot h_{\gamma} \otimes \text{ad}(\rho_n(F^{\delta})) \text{ad}(\rho_n(F^{\mu}))\rho_n(E^{\gamma})) \\
& + \left\langle (S, g^e A(\Lambda)_{\mu}^{\beta} \cdot h_{\gamma} \otimes \text{ad}(\rho_n(F^{\mu}))\rho_n(E^{\gamma})), \phi^{\alpha,n}(f)1 \right\rangle 1 \\
& \pm \left\langle (S, f^e A(\Lambda)_{\mu}^{\beta} \cdot h_{\gamma} \otimes \text{ad}(\rho_n(F^{\mu}))\rho_n(E^{\gamma})), \phi^{\alpha,n}(g)1 \right\rangle 1 \\
& \phi^{\alpha,n}(f)\phi^{\beta,n}(g)^* (S, h_{\gamma} \otimes \rho_n(E^{\gamma})) \\
& = \phi^{\alpha,n}(f)A(\Lambda)_{\mu}^{\beta} \left[(S, -\overline{g^e} \cdot h_{\gamma} \otimes \text{ad}(\rho_n(F^{\mu}))\rho_n(E^{\gamma})) \right. \\
& \quad \left. + \left\langle (S, h_{\gamma} \otimes \rho_n(E^{\gamma})), \phi^{\beta,n}(g)1 \right\rangle 1 \right] \\
& = -A(\Lambda)_{\delta}^{\alpha} \overline{A(\Lambda)_{\mu}^{\beta}} (S, f^e \cdot \overline{g^e} \cdot h_{\gamma} \otimes \text{ad}(\rho_n(F^{\delta})) \text{ad}(\rho_n(F^{\mu}))\rho_n(E^{\gamma})) \\
& \quad + \left\langle (S, h_{\gamma} \otimes \rho_n(E^{\gamma})), \phi^{\beta,n}(g)1 \right\rangle \phi^{\alpha,n}(f)1 \\
& \phi^{\alpha,n}(g)\phi^{\beta,n}(f)^* (S, h_{\gamma} \otimes \rho_n(E^{\gamma})) \\
& = -A(\Lambda)_{\delta}^{\alpha} \overline{A(\Lambda)_{\mu}^{\beta}} (S, g^e \cdot \overline{f^e} \cdot h_{\gamma} \otimes \text{ad}(\rho_n(F^{\delta})) \text{ad}(\rho_n(F^{\mu}))\rho_n(E^{\gamma})) \\
& \quad + \left\langle (S, h_{\gamma} \otimes \rho_n(E^{\gamma})), \phi^{\beta,n}(f)1 \right\rangle \phi^{\alpha,n}(g)1 \\
& -A(\Lambda)_{\delta}^{\alpha} \overline{A(\Lambda)_{\mu}^{\beta}} (S, [f^e \cdot \overline{g^e} \pm g^e \cdot \overline{f^e}] \cdot h_{\gamma} \otimes \text{ad}(\rho_n(F^{\delta})) \text{ad}(\rho_n(F^{\mu}))\rho_n(E^{\gamma})) \\
& + \left\langle (S, h_{\gamma} \otimes \rho_n(E^{\gamma})), \phi^{\beta,n}(g)1 \right\rangle \phi^{\alpha,n}(f)1 \pm \left\langle (S, h_{\gamma} \otimes \rho_n(E^{\gamma})), \phi^{\beta,n}(f)1 \right\rangle \phi^{\alpha,n}(g)1 \\
& g^{\{\hat{f}_0, \hat{f}_1\}}(\hat{H}, \hat{P})(\vec{x}) = \int_{\hat{s} \in \mathbb{R}^2} \frac{e^{-i[\vec{y}(\hat{s}) \cdot (\hat{H}\hat{f}_0 + \hat{P}\hat{f}_1)]}}{2\pi} g(\vec{x} + \vec{y}(\hat{s})) |\hat{\rho}_{\vec{y}}|(\hat{s}) d\hat{s} \\
& \overline{f^{\{\hat{f}_0, \hat{f}_1\}}}(\hat{H}, \hat{P})(\vec{x}) = \int_{\hat{t} \in \mathbb{R}^2} \frac{e^{i[\vec{y}(\hat{t}) \cdot (\hat{H}\hat{f}_0 + \hat{P}\hat{f}_1)]}}{2\pi} \bar{f}(\vec{x} + \vec{y}(\hat{t})) |\hat{\rho}_{\vec{y}}|(\hat{t}) d\hat{t} \\
& [g^{\{\hat{f}_0, \hat{f}_1\}} \cdot \overline{f^{\{\hat{f}_0, \hat{f}_1\}}}] (\hat{H}, \hat{P})(\vec{x}) \\
& = \int_{\hat{s}, \hat{t} \in \mathbb{R}^2} \frac{e^{-i[\vec{y}(\hat{s}) \cdot (\hat{H}\hat{f}_0 + \hat{P}\hat{f}_1)]}}{(2\pi)^2} g_{\vec{x}}(\vec{y}(\hat{s}) + \vec{y}(\hat{t})) \bar{f}_{\vec{x}}(\vec{y}(\hat{t})) |\hat{\rho}_{\vec{y}}|(\hat{s}) |\hat{\rho}_{\vec{y}}|(\hat{t}) d\hat{s} d\hat{t} \\
& = \int_{\hat{t} \in \mathbb{R}^2, \hat{s} \in D} \frac{e^{-i[\vec{y}(\hat{s}) \cdot (\hat{H}\hat{f}_0 + \hat{P}\hat{f}_1)]}}{(2\pi)^2} g_{\vec{x}}(\vec{y}(\hat{s}) + \vec{y}(\hat{t})) \bar{f}_{\vec{x}}(\vec{y}(\hat{t})) |\hat{\rho}_{\vec{y}}|(\hat{s}) |\hat{\rho}_{\vec{y}}|(\hat{t}) d\hat{s} d\hat{t} \\
& [f^{\{\hat{f}_0, \hat{f}_1\}} \cdot g^{\{\hat{f}_0, \hat{f}_1\}}] (\hat{H}, \hat{P})(\vec{x}) \\
& = \int_{\hat{t} \in \mathbb{R}^2, \hat{s} \in -D} \frac{e^{-i[\vec{y}(\hat{s}) \cdot (\hat{H}\hat{f}_0 + \hat{P}\hat{f}_1)]}}{(2\pi)^2} f_{\vec{x}}(\vec{y}(\hat{s}) + \vec{y}(\hat{t})) \bar{g}_{\vec{x}}(\vec{y}(\hat{t})) |\hat{\rho}_{\vec{y}}|(\hat{s}) |\hat{\rho}_{\vec{y}}|(\hat{t}) d\hat{s} d\hat{t}
\end{aligned}$$



$$\begin{aligned}
& \frac{e^{-i[(u-v)\cdot(\hat{H}(\rho_n)\hat{f}_0+\hat{P}(\rho_n)\hat{f}_1)]}}{(2\pi)^2} g_{\vec{x}}(u) \bar{f}_{\vec{x}}(v) \mp \frac{e^{-i[(v-u)\cdot(\hat{H}(\rho_n)\hat{f}_0+\hat{P}(\rho_n)\hat{f}_1)]}}{(2\pi)^2} f_{\vec{x}}(v) \bar{g}_{\vec{x}}(u) = 0 \\
& (u-v)\cdot(\hat{H}(\rho_n)\hat{f}_0+\hat{P}(\rho_n)\hat{f}_1) = c(\hat{P}(\rho_n)\hat{f}_0+\hat{H}(\rho_n)\hat{f}_1)\cdot(\hat{H}(\rho_n)\hat{f}_0+\hat{P}(\rho_n)\hat{f}_1) = 0 \\
& e^{-i[(u-v)\cdot(\hat{H}(\rho_n)\hat{f}_0+\hat{P}(\rho_n)\hat{f}_1)]} = \cos[(u-v)\cdot(\hat{H}(\rho_n)\hat{f}_0+\hat{P}(\rho_n)\hat{f}_1)] = 1 \\
& \left[\frac{e^{-i[(u-v)\cdot(\hat{H}(\rho_n)\hat{f}_0+\hat{P}(\rho_n)\hat{f}_1)]}}{(2\pi)^2} - \frac{e^{-i[(v-u)\cdot(\hat{H}(\rho_n)\hat{f}_0+\hat{P}(\rho_n)\hat{f}_1)]}}{(2\pi)^2} \right] f_{\vec{x}}(v) \bar{g}_{\vec{x}}(u) \\
& \left[-\frac{e^{-i[(u-v)\cdot(\hat{H}(\rho_n)\hat{f}_0+\hat{P}(\rho_n)\hat{f}_1)]}}{(2\pi)^2} + \frac{e^{-i[(v-u)\cdot(\hat{H}(\rho_n)\hat{f}_0+\hat{P}(\rho_n)\hat{f}_1)]}}{(2\pi)^2} \right] f_{\vec{x}}(v) \bar{g}_{\vec{x}}(u), \\
& (f, g) \in \mathcal{P} \times \mathcal{P} \\
& \mapsto \left\langle \phi^{\alpha, n}(f) \phi^{\beta, n}(g)^* (S, h_\gamma \otimes \rho_n(E^\gamma)), (\tilde{S}, \tilde{h}_\gamma \otimes \rho_n(E^\gamma)) \right\rangle \\
& \quad - \left\langle (S, h_\gamma \otimes \rho_n(E^\gamma)), \phi^{\beta, n}(g) 1 \right\rangle \left\langle \phi^{\alpha, n}(f) 1, (\tilde{S}, \tilde{h}_\gamma \otimes \rho_n(E^\gamma)) \right\rangle \\
& \int_{\vec{x} \in \mathbb{R}^4} \int_{\vec{y} \in \mathbb{R}^4} W(\vec{x}, \vec{y}) f(\vec{x}) \otimes_{\mathbb{R}} g(\vec{y}) d\vec{x} d\vec{y} \\
& = \left\langle \phi^{\alpha, n}(f) \phi^{\beta, n}(g)^* (S, h_\gamma \otimes \rho_n(E^\gamma)), (\tilde{S}, \tilde{h}_\gamma \otimes \rho_n(E^\gamma)) \right\rangle \\
& \quad - \left\langle (S, h_\gamma \otimes \rho_n(E^\gamma)), \phi^{\beta, n}(g) 1 \right\rangle \left\langle \phi^{\alpha, n}(f) 1, (\tilde{S}, \tilde{h}_\gamma \otimes \rho_n(E^\gamma)) \right\rangle. \\
& (x^0 \hat{f}_0 + x^1 \hat{f}_1, y^0 \hat{f}_0 + y^1 \hat{f}_1) \mapsto \\
& \int_{(x^2, x^3) \in \mathbb{R}^2} \int_{(y^2, y^3) \in \mathbb{R}^2} W(\vec{x}, \vec{y}) f(\vec{x}) \otimes_{\mathbb{R}} g(\vec{y}) dx^2 dx^3 dy^2 dy^3, \\
& \int_{\vec{x} \in \mathbb{R}^4} \int_{\vec{y} \in \mathbb{R}^4} \operatorname{Re} W(\vec{x}, \vec{y}) [f(\vec{x}) \otimes_{\mathbb{R}} g(\vec{y})] d\vec{x} d\vec{y} \\
& \quad := \int_{\vec{x} \in \mathbb{R}^4} \int_{\vec{y} \in \mathbb{R}^4} W(\vec{x}, \vec{y}) [f(\vec{x}) \otimes_{\mathbb{R}} g(\vec{y})] d\vec{x} d\vec{y} \\
& \int_{\vec{x} \in \mathbb{R}^4} \int_{\vec{y} \in \mathbb{R}^4} \operatorname{Im} W(\vec{x}, \vec{y}) [f(\vec{x}) \otimes_{\mathbb{R}} g(\vec{y})] d\vec{x} d\vec{y} \\
& \quad := \int_{\vec{x} \in \mathbb{R}^4} \int_{\vec{y} \in \mathbb{R}^4} W(\vec{x}, \vec{y}) [f(\vec{x}) \otimes_{\mathbb{R}} g(\vec{y})] d\vec{x} d\vec{y} \\
& \int_{\vec{x} \in \mathbb{R}^4} \int_{\vec{y} \in \mathbb{R}^4} W(\vec{x}, \vec{y}) [f(\vec{x}) \otimes_{\mathbb{R}} g(\vec{y})] d\vec{x} d\vec{y} \\
& = \int_{\vec{x} \in \mathbb{R}^4} \int_{\vec{y} \in \mathbb{R}^4} \operatorname{Re} W(\vec{x}, \vec{y}) [\underline{f}(\vec{x}) \underline{g}(\vec{y}) + \bar{f}(\vec{x}) \bar{g}(\vec{y})] d\vec{x} d\vec{y} \\
& \quad + i \int_{\vec{x} \in \mathbb{R}^4} \int_{\vec{y} \in \mathbb{R}^4} \operatorname{Im} W(\vec{x}, \vec{y}) [\bar{f}(\vec{x}) \underline{g}(\vec{y}) - \underline{f}(\vec{x}) \bar{g}(\vec{y})] d\vec{x} d\vec{y} \\
& \quad \operatorname{Re} W(\vec{x}, \vec{y}) - \operatorname{Re} W(\vec{y}, \vec{x}) = 0, \\
& \quad \operatorname{Im} W(\vec{x}, \vec{y}) + \operatorname{Im} W(\vec{y}, \vec{x}) = 0, \\
& \phi^{\alpha, n}(f) \phi^{\beta, n}(g)^* (S, h_\gamma \otimes \rho_n(E^\gamma)) - \left\langle (S, h_\gamma \otimes \rho_n(E^\gamma)), \phi^{\beta, n}(g) 1 \right\rangle \phi^{\alpha, n}(f) 1 \\
& = -A(\Lambda)_\delta^\alpha A(\Lambda)_\mu^\beta \left(S, \left[f^{\{\hat{f}_0, \hat{f}_1\}} \cdot \overline{g^{\{\hat{f}_0, \hat{f}_1\}}} \right] \cdot h_\gamma \otimes \operatorname{ad}(\rho_n(F^\delta)) \operatorname{ad}(\rho_n(F^\mu)) \rho_n(E^\gamma) \right). \\
& \quad -A(\Lambda)_\delta^\alpha A(\Lambda)_\mu^\beta \left(S, A^\pm \cdot h_\gamma \otimes \operatorname{ad}(\rho_n(F^\delta)) \operatorname{ad}(\rho_n(F^\mu)) \rho_n(E^\gamma) \right), \\
& A^\pm(\vec{z}) = \int_{\hat{s}, \hat{t} \in \mathbb{R}^2} \frac{e^{-i[(\vec{y}(\hat{s}) - \vec{y}(\hat{t})) \cdot (\hat{H}\hat{f}_0 + \hat{P}\hat{f}_1)]}}{(2\pi)^2} f_{\vec{z}}(\vec{y}(\hat{s})) \bar{g}_{\vec{z}}(\vec{y}(\hat{t})) |\hat{\rho}_{\vec{y}}(\hat{s})| |\hat{\rho}_{\vec{y}}(\hat{t})| d\hat{s} d\hat{t} \\
& \quad \pm \int_{\hat{s}, \hat{t} \in \mathbb{R}^2} \frac{e^{-i[(\vec{y}(\hat{t}) - \vec{y}(\hat{s})) \cdot (\hat{H}\hat{f}_0 + \hat{P}\hat{f}_1)]}}{(2\pi)^2} g_{\vec{z}}(\vec{y}(\hat{t})) \bar{f}_{\vec{z}}(\vec{y}(\hat{s})) |\hat{\rho}_{\vec{y}}(\hat{s})| |\hat{\rho}_{\vec{y}}(\hat{t})| d\hat{s} d\hat{t}
\end{aligned}$$



$$\begin{aligned}
& \operatorname{Re}W(\vec{x}, \vec{y}) - \operatorname{Re}W(\vec{y}, \vec{x}) = 0, \\
& \operatorname{Im}W(\vec{x}, \vec{y}) + \operatorname{Im}W(\vec{y}, \vec{x}) = 0 \\
(f, g) \in \mathcal{P} \times \mathcal{P} & \mapsto \left\langle \phi^{\alpha, n}(f) * \phi^{\beta, n}(g) \left(S, h_\gamma \otimes \rho_n(E^\gamma) \right), \left(\tilde{S}, \tilde{h}_\gamma \otimes \rho_n(E^\gamma) \right) \right\rangle \\
& \int_{\vec{x} \in \mathbb{R}^4} \int_{\vec{y} \in \mathbb{R}^4} \tilde{W}(\vec{x}, \vec{y}) f(\vec{x}) \otimes_{\mathbb{R}} g(\vec{y}) d\vec{x} d\vec{y} \\
& = \left\langle \phi^{\alpha, n}(f) * \phi^{\beta, n}(g) \left(S, h_\gamma \otimes \rho_n(E^\gamma) \right), \left(\tilde{S}, \tilde{h}_\gamma \otimes \rho_n(E^\gamma) \right) \right\rangle \\
& \int_{\vec{x} \in \mathbb{R}^4} \int_{\vec{y} \in \mathbb{R}^4} \tilde{W}(\vec{x}, \vec{y}) [f(\vec{x}) \otimes_{\mathbb{R}} g(\vec{y})] d\vec{x} d\vec{y} \\
& = \int_{\vec{x} \in \mathbb{R}^4} \int_{\vec{y} \in \mathbb{R}^4} \operatorname{Re} \tilde{W}(\vec{x}, \vec{y}) [\underline{f}(\vec{x}) \underline{g}(\vec{y}) + \bar{f}(\vec{x}) \bar{g}(\vec{y})] d\vec{x} d\vec{y} \\
& \quad + i \int_{\vec{x} \in \mathbb{R}^4} \int_{\vec{y} \in \mathbb{R}^4} \operatorname{Im} \tilde{W}(\vec{x}, \vec{y}) [\underline{f}(\vec{x}) \bar{g}(\vec{y}) - \bar{f}(\vec{x}) \underline{g}(\vec{y})] d\vec{x} d\vec{y} \\
& \quad \operatorname{Re} \tilde{W}(\vec{x}, \vec{y}) - \operatorname{Re} \tilde{W}(\vec{y}, \vec{x}) = 0 \\
& \quad \operatorname{Im} \tilde{W}(\vec{x}, \vec{y}) + \operatorname{Im} \tilde{W}(\vec{y}, \vec{x}) = 0 \\
(f, g) \in \mathcal{P} \times \mathcal{P} & \mapsto \left\langle \phi^{\alpha, n}(f) * \phi^{\beta, n}(g) 1, \left(\tilde{S}, \tilde{h}_\gamma \otimes \rho_n(E^\gamma) \right) \right\rangle. \\
& \int_{\vec{x} \in \mathbb{R}^4} \int_{\vec{y} \in \mathbb{R}^4} \check{W}(\vec{x}, \vec{y}) f(\vec{x}) \otimes_{\mathbb{R}} g(\vec{y}) d\vec{x} d\vec{y} \\
& = \left\langle \phi^{\alpha, n}(f) * \phi^{\beta, n}(g) 1, \left(\tilde{S}, \tilde{h}_\gamma \otimes \rho_n(E^\gamma) \right) \right\rangle \\
& \int_{\vec{x} \in \mathbb{R}^4} \int_{\vec{y} \in \mathbb{R}^4} \check{W}(\vec{x}, \vec{y}) [f(\vec{x}) \otimes_{\mathbb{R}} g(\vec{y})] d\vec{x} d\vec{y} \\
& = \int_{\vec{x} \in \mathbb{R}^4} \int_{\vec{y} \in \mathbb{R}^4} \operatorname{Re} \check{W}(\vec{x}, \vec{y}) [\underline{f}(\vec{x}) \underline{g}(\vec{y}) + \bar{f}(\vec{x}) \bar{g}(\vec{y})] d\vec{x} d\vec{y} \\
& \quad + i \int_{\vec{x} \in \mathbb{R}^4} \int_{\vec{y} \in \mathbb{R}^4} \operatorname{Im} \check{W}(\vec{x}, \vec{y}) [\underline{f}(\vec{x}) \bar{g}(\vec{y}) - \bar{f}(\vec{x}) \underline{g}(\vec{y})] d\vec{x} d\vec{y} \\
& \quad \operatorname{Re} \check{W}(\vec{x}, \vec{y}) - \operatorname{Re} \check{W}(\vec{y}, \vec{x}) = 0, \\
& \quad \operatorname{Im} \check{W}(\vec{x}, \vec{y}) + \operatorname{Im} \check{W}(\vec{y}, \vec{x}) = 0, \\
& \quad \phi^{\alpha, n}(f) \phi^{\beta, n}(g) * 1 \pm \phi^{\alpha, n}(g) \phi^{\beta, n}(f) * 1 = 0. \\
W(\vec{y}, \vec{x}) = \bar{W}(\vec{x}, \vec{y}), & \check{W}(\vec{y}, \vec{x}) = \bar{W}(\vec{x}, \vec{y}), \check{W}(\vec{y}, \vec{x}) = \bar{W}(\vec{x}, \vec{y}), \\
& \phi^{\alpha, n}(f) \phi^{\beta, n}(g) \pm \phi^{\beta, n}(g) \phi^{\alpha, n}(f) = 0 \\
& \phi^{\alpha, n}(f) * \phi^{\beta, n}(g) \pm \phi^{\beta, n}(g) \phi^{\alpha, n}(f) * = 0 \\
& \langle f \sqrt{\phi_\kappa}, g \sqrt{\phi_\kappa} \rangle = \int_{\mathbb{R}^4} f g \cdot \phi_\kappa d\lambda \\
& \left\{ \frac{h_i(\kappa x^0) h_j(\kappa x^1) h_k(\kappa x^2) h_l(\kappa x^3)}{\sqrt{i! j! k! l!}} \sqrt{\phi_\kappa(\vec{x})} \mid \vec{x} = (x^0, x^1, x^2, x^3) \in \mathbb{R}^4, i, j, k, l \geq 0 \right\} \\
& \left\langle \sum_{0 \leq a < b \leq 3} f_{ab} \otimes dx^a \wedge dx^b, \sum_{0 \leq a < b \leq 3} \hat{f}_{ab} \otimes dx^a \wedge dx^b \right\rangle = \sum_{0 \leq a < b \leq 3} \langle f_{ab}, \hat{f}_{ab} \rangle. \\
df & = \sum_{i=1}^3 \partial_0 f_i \otimes dx^0 \wedge dx^i + \sum_{1 \leq i < j \leq 3} (\partial_i f_j - \partial_j f_i) dx^i \wedge dx^j \\
c_\gamma^{\alpha\beta} & = -\operatorname{Tr} [E^\gamma [E^\alpha, E^\beta]], E^\alpha, E^\beta, E^\gamma \in \mathfrak{g}.
\end{aligned}$$



$$\begin{aligned}
dA + A \wedge A &= \sum_{\gamma=1}^N \left[\sum_{j=1}^3 a_{0;j,\gamma} \otimes dx^0 \wedge dx^j + \sum_{1 \leq i < j \leq 3} a_{i;j,\gamma} \otimes dx^i \wedge dx^j \right. \\
&\quad \left. + \sum_{1 \leq i < j \leq 3} \sum_{1 \leq \alpha, \beta \leq N} a_{i,\alpha} a_{j,\beta} c_\gamma^{\alpha\beta} \otimes dx^i \wedge dx^j \right] \otimes E^\gamma \\
Z &= \int_{\{dA: A \in \mathcal{S}_\kappa(\mathbb{R}^4) \otimes \Lambda^1(\mathbb{R}^3) \otimes \mathfrak{g}\}} e^{-\frac{1}{2} \int_{\mathbb{R}^4} |dA + A \wedge A|^2 d\omega} D[dA] \\
\int_{\mathbb{R}^4} |dA + A \wedge A|^2 d\omega &= \sum_{1 \leq i < j \leq 3} \int_{\mathbb{R}^4} \left[\sum_{\alpha=1}^N a_{i;j,\alpha}^2 + \sum_{\gamma=1}^N \sum_{\substack{\alpha, \beta \\ \alpha, \beta}} a_{i,\alpha} a_{j,\beta} a_{i,\hat{\alpha}} a_{j,\hat{\beta}} c_\gamma^{\alpha\beta} c_\gamma^{\hat{\alpha}\hat{\beta}} \right. \\
&\quad \left. + 2 \sum_{\gamma=1}^N \sum_{\alpha, \beta} a_{i;j,\gamma} a_{i,\alpha} a_{j,\beta} c_\gamma^{\alpha\beta} \right] d\omega + \sum_{j=1}^3 \int_{\mathbb{R}^4} \sum_{\alpha=1}^N a_{0;j,\alpha}^2 d\omega. \\
\exp \left[-\frac{1}{2} \sum_{\alpha=1}^N \int_{\mathbb{R}^4} d\omega \sum_{1 \leq i < j \leq 3} a_{i;j,\alpha}^2 + \sum_{j=1}^3 a_{0;j,\alpha}^2 \right] D[dA] \\
\langle z^r, z^{r'} \rangle &= \frac{1}{\pi} \int_{\mathbb{C}} z^r \cdot \overline{z^{r'}} e^{-|z|^2} dx dp, z = x + \sqrt{-1}p \\
&\quad \left\{ \frac{z^n}{\sqrt{n!}} : n \geq 0 \right\} \\
\Psi_\kappa: \frac{h_i(\kappa \cdot)}{\sqrt{i!}} \frac{h_j(\kappa \cdot)}{\sqrt{j!}} \frac{h_k(\kappa \cdot)}{\sqrt{k!}} \frac{h_l(\kappa \cdot)}{\sqrt{l!}} \sqrt{\phi_\kappa} &\mapsto \frac{z_0^i}{\sqrt{i!}} \frac{z_1^j}{\sqrt{j!}} \frac{z_2^k}{\sqrt{k!}} \frac{z_3^l}{\sqrt{l!}} \\
\Psi_\kappa: f_{i,\alpha} \otimes dx^i \otimes E^\alpha &\mapsto \Psi_\kappa(f_{i,\alpha}) \otimes dx^i \otimes E^\alpha. \\
\left\langle \sum_{0 \leq a < b \leq 3} f_{ab} \otimes dx^a \wedge dx^b, \sum_{0 \leq a < b \leq 3} \hat{f}_{ab} \otimes dx^a \wedge dx^b \right\rangle &= \sum_{0 \leq a < b \leq 3} \langle f_{ab}, \hat{f}_{ab} \rangle. \\
\mathfrak{d}_a \left[z_a^p \prod_{\substack{b \neq a \\ b=0, \dots, 3}} z_b^{q_b} \right] &= \left[\frac{p}{2} z_a^{p-1} - \frac{1}{2} z_a^{p+1} \right] \cdot \prod_{\substack{b \neq a \\ b=0, \dots, 3}} z_b^{q_b} \\
\mathfrak{d}: H^2(\mathbb{C}^4) \otimes \Lambda^1(\mathbb{R}^3) &\rightarrow H^2(\mathbb{C}^4) \otimes \Lambda^2(\mathbb{R}^4) \\
\mathfrak{d} \sum_{i=1}^3 f_i \otimes dx^i &= \sum_{i=1}^3 [\mathfrak{d}_0 f_i] \otimes dx^0 \wedge dx^i + \sum_{1 \leq i < j \leq 3} [\mathfrak{d}_i f_j - \mathfrak{d}_j f_i] \otimes dx^i \wedge dx^j \\
\Psi_\kappa: \mathcal{S}_\kappa(\mathbb{R}^4) \otimes \Lambda^2(\mathbb{R}^4) \otimes \mathfrak{g} &\rightarrow H^2(\mathbb{C}^4) \otimes \Lambda^2(\mathbb{R}^4) \otimes \mathfrak{g} \\
\Psi_\kappa: \sum_{1 \leq i < j \leq 3} f_{i,j,\alpha} \otimes dx^i \wedge dx^j \otimes E^\alpha &\mapsto \sum_{1 \leq i < j \leq 3} \Psi_\kappa[f_{i,j,\alpha}] \otimes dx^i \wedge dx^j \otimes E^\alpha. \\
\langle dA, dA \rangle &= \kappa^2 \langle \mathfrak{d} \Psi_\kappa[A], \mathfrak{d} \Psi_\kappa[A] \rangle \\
&= \frac{1}{Z} e^{-\frac{1}{2} \int_{\mathbb{C}^4} |\kappa \mathfrak{d} A + A \wedge A|^2 d\lambda_4} D[\mathfrak{d}A] \\
Z &= \int_{\{\mathfrak{d}A: A \in H^2(\mathbb{C}^4) \otimes \Lambda^1(\mathbb{R}^3) \otimes \mathfrak{g}\}} e^{-\frac{1}{2} \int_{\mathbb{C}^4} |\kappa \mathfrak{d} A + A \wedge A|^2 d\lambda_4} D[\mathfrak{d}A] \\
\mathbb{H} &:= \{[\mathfrak{d}_0 H^2(\mathbb{C}^4)] \otimes [* \Lambda^2(\mathbb{R}^3)]\} \oplus \{H^2(\mathbb{C}^4) \otimes \Lambda^2(\mathbb{R}^3)\} \\
&\subset H^2(\mathbb{C}^4) \otimes \Lambda^2(\mathbb{R}^4) \\
\frac{1}{Z} e^{-\frac{1}{2} \int_{\mathbb{C}^4} |\kappa \mathfrak{d} A + A \wedge A|^2 d\lambda_4} D[\mathfrak{d}A] &:= \frac{y^\kappa d\tilde{\mu}_{\kappa^2}^{X^N}}{\int_{\mathbb{B} \otimes \mathfrak{g}} y^\kappa d\tilde{\mu}_{\kappa^2}^{NN}} = \frac{y^\kappa d\tilde{\mu}_{\kappa^2}^{X^N}}{\mathbb{E}[y^\kappa]},
\end{aligned}$$

$$\begin{aligned}
\mathbb{E}_{\text{YM}}^\kappa[F] &:= \frac{1}{\int_{\mathbb{B} \otimes \mathfrak{g}} y^\kappa d\tilde{\mu}_{\kappa^2}^{\times N}} \int_{\mathbb{B} \otimes \mathfrak{g}} F y^\kappa d\tilde{\mu}_{\kappa^2}^{\times N} \\
\beta(c) &= \frac{\partial c}{\partial [\ln \tilde{N}_n]}, \\
\text{Tr}[\rho(E^\alpha)\rho(E^\beta)] &= C(\rho)\text{Tr}[E^\alpha E^\beta] \\
\mathcal{E}(\rho) &:= -\sum_{\alpha=1}^N \rho(E^\alpha)\rho(E^\alpha) \\
&\sum_{\alpha=1}^N \sum_{j=1}^3 \mathfrak{d}_0 A_{j,\alpha} \otimes dx^0 \wedge dx^j \otimes E^\alpha \in \mathbb{B} \otimes \mathfrak{g} \\
\frac{1}{\kappa} \sum_{\alpha=1}^N \frac{\kappa^2}{4} \int_{\hat{s} \in [-\delta, 1+\delta]^2} d\hat{s} \sum_{j=1}^3 |J_{0j}^\sigma|(\hat{s}) \kappa [\psi \cdot \mathfrak{d}_0 A_{j,\alpha}] (\kappa \sigma(\hat{s})/2) \otimes \rho(E^\alpha) \\
&c \sum_{\alpha=1}^N \int_{\hat{s} \in I_\delta^2} d\hat{s} \sum_{j=1}^3 |J_{0j}^\sigma|(\hat{s}) [\mathfrak{d}_0 A_{j,\alpha}] (\sigma(\hat{s})) \otimes \rho(E^\alpha) \\
\left\langle v_{R[a,T]}^{\kappa,\rho} \right\rangle^2 &= -\mathbb{E} \left[\left(\cdot, v_{R[a,T]}^{\kappa,\rho} \right)_\# y^\kappa \right] \\
&-\mathbb{E} \left[\left(\cdot, v_{R[a,T]}^{\kappa,\rho} \right)_\# \left(\cdot, v_{R_\delta[a,T]}^{\kappa,\rho} \right)_\# y^\kappa \right], \\
A^\rho &:= \sum_{\alpha=1}^N \sum_{i=1}^3 a_{i,\alpha} \otimes dx^i \otimes \rho(E^\alpha) \in \mathcal{S}_\kappa(\mathbb{R}^4) \otimes \Lambda^1(\mathbb{R}^3) \otimes \rho(\mathfrak{g}) \\
\frac{1}{Z} \int_{\{dA \in \mathcal{S}_\kappa(\mathbb{R}^4) \otimes \Lambda^2(\mathbb{R}^4) \otimes \mathfrak{g}\}} \exp \left[c \int_{R[a]} d[A^\rho] \right] e^{-\frac{1}{2} S_{\text{YM}}(A)} D[dA] \\
&:= \mathbb{E}_{\text{YM}}^\kappa \left[\exp \left[\left(\cdot, v_{R[a]}^{\kappa,\rho} \right)_\# \right] \right] \\
Z &= \int_{\{dA \in \mathcal{S}_\kappa(\mathbb{R}^4) \otimes \Lambda^2(\mathbb{R}^4) \otimes \mathfrak{g}\}} e^{-\frac{1}{2} S_{\text{YM}}(A)} DA \\
&-\mathbb{E} \left[\left(\cdot, v_{R[a]}^{\kappa,\rho_n} \right)_\# \left(\cdot, v_{R_\delta[a]}^{\kappa,\rho_n} \right)_\# y^\kappa \right] = \frac{|a|}{4} \otimes \mathcal{E}(\rho_n) - \epsilon(n, \kappa), \\
\frac{\bar{c}}{\kappa^4} C(\rho_n) &\leq \text{Tr} \epsilon(n, \kappa) \leq \frac{\bar{c}}{\kappa^4} C(\rho_n) \\
-\frac{1}{C(\rho_n)} \mathbb{E} \left[\left(\cdot, v_{R[a]}^{\kappa,\rho_n} \right)_\# \left(\cdot, v_{R_\delta[a]}^{\kappa,\rho_n} \right)_\# y^\kappa \right] &= \frac{|a|}{4} \otimes \frac{C_2(\rho_n)}{C(\rho_n)} \mathbb{I}_n - \frac{1}{C(\rho_n)} \epsilon(n, \kappa), \\
\tilde{N}_n \frac{C_2(\rho_n)}{C(\rho_n)} - \frac{1}{C(\rho_n)} \text{Tr} \epsilon(n, \kappa) &\equiv N - \frac{1}{C(\rho_1)} \text{Tr} \epsilon(1, \kappa) \\
G_n^{(2)}(c, e) &:= \frac{\tilde{N}_n}{e} - \text{Tr} \epsilon(1/c) \\
&= \frac{\tilde{N}_n}{e} - c^4 \bar{\lambda} + f(c^5) \\
\left[e \frac{\partial}{\partial e} + \beta(c) \frac{\partial}{\partial c} + 2\gamma(c) \right] G_n^{(2)}(c, e) &= 0. \\
G_n^{(2)}(c, e) &= \frac{\tilde{N}_n}{e} - c^4 \bar{\lambda} + f(c^5) \\
\frac{\partial}{\partial e} G_n^{(2)}(c, e) &= -\frac{\tilde{N}_n}{e^2}, \frac{\partial}{\partial c} G_n^{(2)}(c, e) = -4c^3 \bar{\lambda} + \tilde{f}(c^4), \\
-\frac{\tilde{N}_n}{e} - 4\beta(c)c^3 \bar{\lambda} + 2\gamma(c)G_n^{(2)}(c, e) + \beta(c)\tilde{f}(c^4) &= 0
\end{aligned}$$

$$\begin{aligned}
& -\frac{c}{4}\tilde{f}(c^4) + f(c^5) - 4c^3\lambda(c)\bar{\lambda} + \lambda(c)\tilde{f}(c^4) = 0 \\
& \frac{1}{c^4}|\tilde{f}(c^4)| + \frac{1}{c^5}|f(c^5)| \leq \tilde{C}_3 \\
& \lambda(c) = \frac{1}{-4c^3\bar{\lambda} + \tilde{f}(c^4)} \left[\frac{c}{4}\tilde{f}(c^4) - f(c^5) \right] \\
& \frac{\hat{p}^2}{\hat{H}^2} - 1 = -\frac{m^2}{\hat{H}^2} \\
& E^\alpha = \sum_{\beta=1}^l a_{\alpha,\beta} H_\beta, 1 \leq \alpha \leq l \\
& \lambda_\rho \equiv (\lambda_\rho(H_1), \dots, \lambda_\rho(H_l)), \\
& \mathcal{E}(\rho) = -\sum_{\alpha=1}^N \rho(E^\alpha) \rho(E^\alpha) = C_2(\rho)I, C_2(\rho) \geq 0 \\
& \langle -\rho(E^\alpha) \rho(E^\alpha) v, v \rangle = \langle \rho(E^\alpha) v, \rho(E^\alpha) v \rangle \geq 0 \\
& \langle C_2(\rho) v, v \rangle = \sum_{\alpha=1}^N \langle \rho(E^\alpha) v, \rho(E^\alpha) v \rangle \geq \sum_{\alpha=1}^l \left| \sum_{\beta=1}^l a_{\alpha,\beta} \lambda_\rho(H_\beta) v \right|_2^2 \\
& = \sum_{\beta=1}^l \sum_{\gamma=1}^l \sum_{\alpha=1}^l \lambda_\rho(H_\beta) a_{\alpha,\beta} a_{\alpha,\gamma} \lambda_\rho(H_\gamma) \geq c |\lambda_\rho|_2^2 \\
& \{\rho_n: \mathfrak{g} \rightarrow \text{End}(\mathbb{C}^{\tilde{N}_n})\}_{n=1}^\infty, \text{ such that } 0 < C_2(\rho_n) \leq C_2(\rho_{n+1}) \\
& \hat{H}(\rho_n)^2 := \frac{\tilde{N}_n}{4} C_2(\rho_n) = \frac{N}{4} C(\rho_n) > 0 \\
& \frac{\partial c}{\partial [\ln \tilde{N}]} = -\frac{c}{4} + \lambda(c), |\lambda(c)| \leq \tilde{C}_4 c^2 \\
& \frac{dc}{d[\ln \tilde{N}]} = -\frac{c}{4} + \lambda(c) \Rightarrow \frac{dc}{c - 4\lambda(c)} = -\frac{d[\ln \tilde{N}]}{4} \\
& \frac{1}{c} \frac{dc}{1 + \mu(c)} = -\frac{d[\ln \tilde{N}]}{4} \Rightarrow \left[\frac{1}{c} \sum_{k=0}^\infty (-1)^k \mu(c)^k \right] dc = -\frac{d[\ln \tilde{N}]}{4} \\
& \ln c + \tilde{\mu}(c) = -\frac{1}{4} \ln \tilde{N} + C \\
& ce^{\tilde{\mu}(c)} = \frac{1}{\hat{C}} \tilde{N}^{-1/4} \Rightarrow c(1 + \tilde{\mu}(c)) = \frac{1}{\hat{C}} \tilde{N}^{-1/4} \\
& \left| \frac{1}{\hat{C}} \frac{1}{1 + \tilde{\mu}(c)} \right| \leq \frac{1}{\hat{C}} [1 + \tilde{C}_8 |\tilde{\mu}(c)|] \leq \frac{1}{\hat{C}} [1 + \tilde{C}_8 \tilde{C}_7 c] \\
& \leq \frac{1}{\hat{C}} \left[1 + \tilde{C}_7 \tilde{C}_8 \frac{1}{\hat{C}} \tilde{N}^{-1/4} (1 + \tilde{C}_8 \tilde{C}_7) \right] = \frac{1}{\hat{C}} + \frac{\tilde{N}^{-1/4}}{\hat{C}^2} \tilde{C}_7 \tilde{C}_8 (1 + \tilde{C}_7 \tilde{C}_8) \\
& \tilde{N}^{1/4} |\tilde{\mu}(c)| \leq c \tilde{N}^{1/4} \tilde{C}_7 \leq \frac{\tilde{C}_7}{\hat{C}} + \frac{\tilde{C}_7}{\hat{C}^2} \tilde{C}_7 \tilde{C}_8 (1 + \tilde{C}_7 \tilde{C}_8) \\
& \kappa = \frac{1}{c} = \hat{C} \tilde{N}^{1/4} (1 + \tilde{\mu}(c)) \\
& = \hat{C} \tilde{N}^{1/4} + \tilde{R}(c) \\
& |\tilde{R}(c)| \leq \tilde{C}_7 + \frac{\tilde{C}_7}{\hat{C}} \tilde{C}_7 \tilde{C}_8 (1 + \tilde{C}_7 \tilde{C}_8) = \tilde{C}_7 + \frac{1}{\hat{C}} \tilde{C}_9 \\
& 0 < m_1^2 := \text{Tr} \epsilon(1, \kappa_1) = \frac{\tilde{N}_1}{4} C_2(\rho_1) + \text{Tr} \mathbb{E} \left[\left(\cdot, v_{R[a]}^{\kappa_1, \rho_1} \right)_\# \left(\cdot, v_{R\delta[a]}^{\kappa_1, \rho_1} \right)_\# y^\kappa \right] \\
& -\text{Tr} \mathbb{E} \left[\left(\cdot, v_{R[a]}^{\kappa, \rho_n} \right)_\# \left(\cdot, v_{R\delta[a]}^{\kappa, \rho_n} \right)_\# y^\kappa \right] = \frac{\tilde{N}_n}{4} C_2(\rho_n) - \text{Tr} \epsilon(n, \kappa)
\end{aligned}$$

$$\begin{aligned}
& \frac{4}{NC(\rho_n)} \text{TrE} \left[- \left(\cdot, \nu_{R[a]}^{\kappa, \rho_n} \right)_{\#} \left(\cdot, \nu_{R\delta[a]}^{\kappa, \rho_n} \right)_{\#} y^{\kappa} \right] - 1 = - \frac{4 \text{Tr}\epsilon(n, \kappa)}{NC(\rho_n)}. \\
& m_n^2 := \text{Tr}[\epsilon(n, \kappa_n)] \\
& 0 < \frac{c}{\kappa_n^4} < \frac{1}{C(\rho_n)} \text{Tr}[\epsilon(n, \kappa_n)] \leq \frac{\bar{c}}{\kappa_n^4} \\
& \hat{P}(\rho_n)^2 := -\text{TrE} \left[\left(\cdot, \nu_{R[a]}^{\kappa_n, \rho_n} \right)_{\#} \left(\cdot, \nu_{R\delta[a]}^{\kappa_n, \rho_n} \right)_{\#} y^{\kappa} \right] \\
& = \frac{\tilde{N}_n}{4} C_2(\rho_n) - \text{Tr}\epsilon(n, \kappa_n) = \frac{\tilde{N}_n}{4} C_2(\rho_n) - m_n^2 > 0, \\
& \frac{\hat{P}(\rho_n)^2}{\hat{H}(\rho_n)^2} = 1 - \frac{m_n^2}{\hat{H}(\rho_n)^2} \rightarrow 1 \\
& \hat{H}(\rho_n)^2 - \hat{P}(\rho_n)^2 = m_n^2 > 0, \\
& U(\vec{a}, 1) \left(S, f_{\alpha} \otimes \rho(E^{\alpha}), \{\hat{f}_a\}_{a=0}^3 \right) \\
& := e^{-i[f_0 \cdot \hat{H}(\vec{a}, \rho) + \hat{f}_1 \cdot \hat{P}(\vec{a}, \rho)]} \left(S + \vec{a}, f_{\alpha}(\cdot - \vec{a}) \otimes \rho(E^{\alpha}), \{\hat{f}_a\}_{a=0}^3 \right) \\
& = e^{i[a^0 \hat{H}(\rho) - a^1 \hat{P}(\rho)]} \left(S + \vec{a}, f_{\alpha}(\cdot - \vec{a}) \otimes \rho(E^{\alpha}), \{\hat{f}_a\}_{a=0}^3 \right) \\
& \hat{P} \sum_{n=0}^{\infty} v_n := \sum_{n=1}^{\infty} \hat{P}(\rho_n) v_n, \hat{H} \sum_{n=0}^{\infty} v_n := \sum_{n=1}^{\infty} \hat{H}(\rho_n) v_n, \\
& \hat{f}(p) = \mathcal{F}[f](p) := \frac{1}{\sqrt{2\pi}} \int_{\mathbb{R}} e^{-ipx} f(x) dx \\
& \hat{f}(\vec{p}) = \mathcal{F}[f](\vec{p}) := \frac{1}{(2\pi)^2} \int_{\mathbb{R}^4} e^{-i\vec{p} \cdot \vec{x}} f(\vec{x}) d\vec{x} \\
& \left(\frac{1}{\sqrt{2\pi}} \right)^2 \int_{S_0} e^{-i[q^2 x^2 + q^3 x^3]} F_{\alpha}(\hat{H}(\rho), \hat{P}(\rho), x^2, x^3) dx^2 dx^3 \otimes \rho(E^{\alpha}) \\
& = \hat{F}_{\alpha}(\hat{H}(\rho), \hat{P}(\rho), q^2, q^3) \otimes \rho(E^{\alpha}) \\
& \hat{F}(\hat{H}(\rho)\hat{f}_0 + \hat{P}(\rho)\hat{f}_1 + q^2\hat{f}_2 + q^3\hat{f}_3) \\
& := \frac{e^{-i[a^0 \hat{H}(\rho) - a^1 \hat{P}(\rho)]}}{2\pi} \int_{\hat{s} \in \mathbb{R}^2} e^{-i(sq^2 + \bar{s}q^3)} f\{\hat{f}_0, \hat{f}_1\}(\hat{H}(\rho), \hat{P}(\rho))(s\hat{f}_2 + \bar{s}\hat{f}_3) d\hat{s} \\
& \equiv e^{-i[a^0 \hat{H}(\rho) - a^1 \hat{P}(\rho)]} \hat{f}(\hat{H}(\rho)\hat{f}_0 + \hat{P}(\rho)\hat{f}_1 + q^2\hat{f}_2 + q^3\hat{f}_3), \\
& \hat{F}_{\alpha}(\hat{H}(\rho)\hat{f}_0 + \hat{P}(\rho)\hat{f}_1 + q^2\hat{f}_2 + q^3\hat{f}_3) \otimes \rho(E^{\alpha}) \left(S, f_{\alpha} \otimes \rho(E^{\alpha}), \{\hat{f}_a\}_{a=0}^3 \right) \\
& e^{-i\vec{a} \cdot \vec{\alpha}} e^{-i[q^2 a^2 + q^3 a^3]} \hat{F}_{\alpha}(\hat{H}(\rho), \hat{P}(\rho), q^2, q^3) \otimes \rho(E^{\alpha}) \\
& \vec{\alpha} = (\hat{H}(\rho), \hat{P}(\rho), 0, 0)
\end{aligned}$$

$$\begin{aligned}
\langle \phi^{\alpha, n}(f)1, \phi^{\beta, n}(g)1 \rangle &:= C(\rho_n) \text{Tr}[-F^{\alpha} F^{\beta}] \int_{S_0} \left[f^{\{e_0, e_1\}} \overline{g^{\{e_0, e_1\}}} \right] (\hat{H}(\rho_n), \hat{P}(\rho_n))(\hat{s}) d\hat{s} \\
\phi^{\alpha, n}(f)1 &= \int_{\vec{x} \in \mathbb{R}^4} d\vec{x} f(\vec{x}) \phi^{\alpha, n}(\vec{x})1 \\
U(\vec{a}, 1) \phi^{\alpha, n}(\vec{x}) U(\vec{a}, 1)^{-1} &= \phi^{\alpha, n}(\vec{x} + \vec{a}) \\
\frac{1}{2\pi} e^{i[x^0 \hat{H}(\rho_n) - x^1 \hat{P}(\rho_n)]} \delta(\cdot - (x^2, x^3)) &\otimes \rho_n(F^{\alpha}) \\
\int_{\vec{z} \in \mathbb{R}^4} p_{\kappa}^{\vec{x}}(\vec{z}) p_{\kappa}^{\vec{y}}(\vec{z}) d\vec{z} &= \frac{\kappa^2}{4(2\pi)} \exp[-\kappa^2 |\vec{x} - \vec{y}|^2 / 8] \\
\left\langle \frac{1}{2\pi} e^{i[x^0 \hat{H} - x^1 \hat{P}]} p_{\kappa}^{x^+} \otimes \rho_n(F^{\alpha}), \frac{1}{2\pi} e^{i[y^0 \hat{H} - y^1 \hat{P}]} p_{\kappa}^{y^+} \otimes \rho_n(F^{\beta}) \right\rangle
\end{aligned}$$



$$\begin{aligned}
& \frac{1}{(2\pi)^2} e^{i[\hat{H}(x^0-y^0)-\hat{P}(x^1-y^1)]} \int_{\vec{z} \in \mathbb{R}^4} p_k^{x^+}(\vec{z}) p_k^{y^+}(\vec{z}) d\vec{z} \cdot \langle \rho_n(F^\alpha), \rho_n(F^\beta) \rangle \\
&= \frac{\kappa^2}{4(2\pi)^3} e^{i[\hat{H}(x^0-y^0)-\hat{P}(x^1-y^1)]} \exp[-\kappa^2|x^+ - y^+|^2/8] \cdot \langle \rho_n(F^\alpha), \rho_n(F^\beta) \rangle \\
& \frac{1}{(2\pi)^2} e^{i[\hat{H}(x^0-y^0)-\hat{P}(x^1-y^1)]} \cdot \delta(x^+ - y^+) \cdot C(\rho_n) \text{Tr}[-F^\alpha F^\beta], x^+, y^+ \in S_0 \\
& \frac{1}{(2\pi)^2} \int_{\vec{x}, \vec{y} \in \mathbb{R}^4} f(\vec{x}) \overline{g(\vec{y})} \left\langle e^{i(x^0 \hat{H} - x^1 \hat{P})} p_k^{x^+} \otimes \rho_n(F^\alpha), e^{i(y^0 \hat{H} - y^1 \hat{P})} p_k^{y^+} \otimes \rho_n(F^\beta) \right\rangle d\vec{x} d\vec{y} \\
&= \frac{\kappa^2}{4} \int_{\hat{s}, \hat{t} \in S_0} f^{\{e_0, e_1\}}(\hat{s}) \overline{g^{\{e_0, e_1\}}(\hat{t})} \frac{1}{2\pi} \exp[-\kappa^2|\hat{s} - \hat{t}|^2/8] d\hat{s} d\hat{t} \cdot \langle \rho_n(F^\alpha), \rho_n(F^\beta) \rangle \\
& \int_{\hat{s} \in S_0} \left[f^{\{e_0, e_1\}} \overline{g^{\{e_0, e_1\}}} \right](\hat{s}) d\hat{s} \cdot \langle \rho_n(F^\alpha), \rho_n(F^\beta) \rangle
\end{aligned}$$

$$\lim_{\kappa \rightarrow \infty} \int_{\vec{z} \in \mathbb{R}^4} p_k^{\vec{x}}(\vec{z}) f(\vec{z}) d\vec{z} = 0$$

$$A_r^n = \psi^{\alpha_1, n}(f_1) \dots \psi^{\alpha_r, n}(f_r), B_s^n = \psi^{\beta_1, n}(g_1) \dots \psi^{\beta_s, n}(g_s)$$

$$\langle A_r^n P_0 B_s^n 1, 1 \rangle \equiv \langle A_r^n B_s^n 1, 1 \rangle - \langle A_r^n 1, 1 \rangle \langle B_s^n 1, 1 \rangle$$

$$h_\theta = \begin{cases} f_\theta, & \text{if } \psi^{\alpha_\theta, n}(f_\theta) = \phi^{\alpha_\theta, n}(f_\theta) \\ -\bar{f}_\theta, & \text{if } \psi^{\alpha_\theta, n}(f_\theta) = \phi^{\alpha_\theta, n}(f_\theta)^* \end{cases}$$

$$h_\theta = \begin{cases} g_{\theta-r}, & \text{if } \psi^{\beta_{\theta-r}, n}(g_{\theta-r}) = \phi^{\beta_{\theta-r}, n}(g_{\theta-r}) \\ -\bar{g}_{\theta-r}, & \text{if } \psi^{\beta_{\theta-r}, n}(g_{\theta-r}) = \phi^{\beta_{\theta-r}, n}(g_{\theta-r})^* \end{cases}$$

$$\tilde{h}_\theta = \begin{cases} g_\theta, & \text{if } \psi^{\alpha_\theta, n}(g_\theta) = \phi^{\alpha_\theta, n}(g_\theta) \\ -\bar{g}_\theta, & \text{if } \psi^{\alpha_\theta, n}(g_\theta) = \phi^{\alpha_\theta, n}(g_\theta)^* \end{cases}$$

$$\tilde{h}_\theta = \begin{cases} f_{\theta-s}, & \text{if } \psi^{\alpha_\theta, n}(f_{\theta-s}) = \phi^{\alpha_\theta, n}(f_{\theta-s}) \\ -\bar{f}_{\theta-s}, & \text{if } \psi^{\alpha_\theta, n}(f_{\theta-s}) = \phi^{\alpha_\theta, n}(f_{\theta-s})^* \end{cases}$$

$$\tilde{h}_\theta = \begin{cases} f_{\theta-s}, & \text{if } \psi^{\beta_{\theta-r}, n}(f_{\theta-s}) = \phi^{\beta_{\theta-r}, n}(f_{\theta-s}) \\ -\bar{f}_{\theta-s}, & \text{if } \psi^{\beta_{\theta-r}, n}(f_{\theta-s}) = \phi^{\beta_{\theta-r}, n}(f_{\theta-s})^* \end{cases}$$

$$\psi^{\alpha_\theta, n}(h) = \phi^{\alpha_\theta, n}(h)^*, 1 \leq \theta \leq r$$

$$\psi^{\beta_{\theta-r}, n}(h) = \phi^{\beta_{\theta-r}, n}(h)^*, r+1 \leq \theta \leq r+s$$

$$\chi(\theta) = \begin{cases} -1, & \theta \\ 1, & \text{otherwise} \end{cases}$$

$$\begin{aligned}
\int_Q \{h_\theta\}_{\theta=1}^{r+s} &:= \prod_{l=1}^{n(Q)} \left\{ \int_{S_0} \left[\prod_{\theta \in A_l} \int_{y_\theta^- \in \mathbb{R}^2} \frac{e^{i\chi(\theta)[y_\theta^0 \hat{H} - y_\theta^1 \hat{P}]} }{2\pi} h_\theta(y_\theta^-, y^+) dy_\theta^- \right] dy^+ \right\} \\
&= \prod_{l=1}^{n(Q)} \left\{ \int_{S_0} \left[\prod_{\theta \in A_l} \int_{y_\theta^0, y_\theta^1 \in \mathbb{R}} \frac{e^{i\chi(\theta)[y_\theta^0 \hat{H} - y_\theta^1 \hat{P}]} }{2\pi} h_\theta(y_\theta^0, y_\theta^1, y^2, y^3) dy_\theta^0 dy_\theta^1 \right] dy^2 dy^3 \right\}
\end{aligned}$$

$$\langle A_r^n P_0 B_s^n 1, 1 \rangle = \sum_{Q \in \Gamma} c_Q \int_Q \{h_\theta\}_{\theta=1}^{r+s},$$

$$C_r^n = \psi^{\alpha_1, n}(g_1) \dots \psi^{\alpha_s, n}(g_s) \psi^{\alpha_{s+1}, n}(f_1) \dots \psi^{\alpha_r, n}(f_{r-s})$$

$$D_s^n = \psi^{\beta_1, n}(f_{r-s+1}) \dots \psi^{\beta_s, n}(f_r)$$

$$\mathcal{W}^n: f_1 \otimes_{\mathbb{R}} \dots \otimes_{\mathbb{R}} f_r \otimes_{\mathbb{R}} g_1 \otimes_{\mathbb{R}} \dots \otimes_{\mathbb{R}} g_s \mapsto \langle A_r^n P_0 B_s^n 1, 1 \rangle.$$

$$\mathcal{W}^n(\{\vec{x}_\tau\}_{\tau=1}^r, \{\vec{x}_\theta\}_{\theta=r+1}^{r+s}) \equiv \mathcal{W}^n(\{\vec{x}_\tau\}_{\tau=1}^s, \{\vec{x}_\theta\}_{\theta=s+1}^{s+r})$$

$$:= \mathcal{W}^n(\vec{x}_1, \dots, \vec{x}_r, \vec{x}_{r+1}, \dots, \vec{x}_{r+s})$$



$$\begin{aligned}
& \langle A_r^n B_s^n 1, 1 \rangle - \langle A_r^n 1, 1 \rangle \langle B_s^n 1, 1 \rangle \equiv \langle A_r^n P_0 B_s^n 1, 1 \rangle \\
& = \int_{\mathbb{R}^4 \times \dots \times \mathbb{R}^4} \mathcal{W}^n(\{\vec{x}_\tau\}_{\tau=1}^r, \{\vec{x}_\theta\}_{\theta=r+1}^{r+s}) \bigotimes_{\tau=1}^{r+s} p_\tau(\vec{x}_\tau) \cdot \prod_{\tau=1}^{r+s} d\vec{x}_\tau \\
& = c_R \int_{S_0} \left[\prod_{\tau=1}^r \int_{x_\tau^- \in \mathbb{R}^2} E(x_\tau^-) h_\tau(x_\tau^-, x^+) dx_\tau^- \cdot \prod_{\theta=r+1}^{r+s} \int_{x_\theta^- \in \mathbb{R}^2} E(x_\theta^-) h_\theta(x_\theta^-, x^+) dx_\theta^- \right] dx^+ \\
& + \sum_{\substack{Q \neq R \\ Q \in \Gamma}} c_Q \int_Q \{h_\theta\}_{\theta=1}^{r+s}
\end{aligned}$$

$$E(x_\tau^-) = E(x_\tau^0, x_\tau^1) = \frac{e^{i\chi(\tau)[x_\tau^0 \hat{H} - x_\tau^1 \hat{P}]}}{2\pi}$$

$$\mathcal{W}^n: g_1 \otimes_{\mathbb{R}} \dots \otimes_{\mathbb{R}} g_s \otimes_{\mathbb{R}} f_1 \otimes_{\mathbb{R}} \dots \otimes_{\mathbb{R}} f_r \mapsto \langle C_r^n P_0 D_s^n 1, 1 \rangle,$$

$$\begin{aligned}
& \langle C_r^n D_s^n 1, 1 \rangle - \langle C_r^n 1, 1 \rangle \langle D_s^n 1, 1 \rangle \equiv \langle C_r^n P_0 D_s^n 1, 1 \rangle \\
& = \int_{\mathbb{R}^4 \times \dots \times \mathbb{R}^4} \mathcal{W}^n(\{\vec{x}_\tau\}_{\tau=1}^s, \{\vec{x}_\theta\}_{\theta=s+1}^{s+r}) \bigotimes_{\tau=1}^{s+r} \tilde{p}_\tau(\vec{x}_\tau) \cdot \prod_{\tau=1}^{s+r} d\vec{x}_\tau \\
& = c_R \int_{S_0} \left[\prod_{\tau=1}^s \int_{x_\tau^- \in \mathbb{R}^2} E(x_\tau^-) \tilde{h}_\tau(x_\tau^-, x^+) dx_\tau^- \cdot \prod_{\theta=s+1}^{s+r} \int_{x_\theta^- \in \mathbb{R}^2} E(x_\theta^-) \tilde{h}_\theta(x_\theta^-, x^+) dx_\theta^- \right] dx^+ \\
& + \sum_{\substack{Q \neq R \\ Q \in \Gamma}} c_Q \int_Q \{\tilde{h}_\theta\}_{\theta=1}^{s+r}.
\end{aligned}$$

$$\begin{aligned}
& \psi^{\beta_1, n}(g_1(\cdot - \vec{a})) \dots \psi^{\beta_{s-1}, n}(g_{s-1}(\cdot - \vec{a})) \left(S_0 + \vec{a}, g_s^{\{e_0, e_1\}}(\cdot - \vec{a}) \otimes \rho_n(F^{\beta_s}), \{e_a\}_{a=0}^3 \right), \\
& P_0 U(\vec{a}, 1) B_s^n 1 = P_0 \psi^{\beta_1, n}(g_1)_{U(\vec{a})} \psi^{\beta_2, n}(g_2)_{U(\vec{a})} \dots \psi^{\beta_{s-1}, n}(g_{s-1})_{U(\vec{a})} U(\vec{a}, 1) \psi^{\beta_s, n}(g_s) 1 \\
& = e^{i[a^0 \hat{H} - a^1 \hat{P}]} P_0 B_s^{n, \vec{a}} 1 = e^{i[a^0 \hat{H} - a^1 \hat{P}]} [B_s^{n, \vec{a}} 1 - \langle B_s^n 1, 1 \rangle 1]
\end{aligned}$$

$$\vec{a} = \sum_{b=0}^3 a^b e_b$$

$$\langle A_r^n P_0 U(\vec{a}, 1) B_s^n 1, 1 \rangle = e^{i[a^0 \hat{H} - a^1 \hat{P}]} \left[\langle A_r^n B_s^{n, \vec{a}} 1, 1 \rangle - \langle A_r^n 1, 1 \rangle \langle B_s^n 1, 1 \rangle \right]$$

$$\begin{aligned}
& e^{i[a^0 \hat{H} - a^1 \hat{P}]} \int_{\mathbb{R}^2} \frac{e^{i[s \hat{H} - t \hat{P}]} }{2\pi} f(s, t, x^2, x^3) ds dt \\
& = \int_{\mathbb{R}^2} \frac{e^{i[(s+a^0) \hat{H} - (t+a^1) \hat{P}]} }{2\pi} f(s, t, x^2, x^3) ds dt \\
& = \int_{\mathbb{R}^2} \frac{e^{i[s \hat{H} - t \hat{P}]} }{2\pi} f(s - a^0, t - a^1, x^2, x^3) ds dt \\
& = f(\cdot - (a^0, a^1, 0, 0))^{\{e_0, e_1\}}(\hat{H}, \hat{P})(0, 0, x^2, x^3)
\end{aligned}$$

$$\mathcal{W}^n(\{\vec{x}_\tau\}_{\tau=1}^r, \{\vec{x}_\theta + \vec{a}\}_{\theta=r+1}^{r+s}) := \mathcal{W}^n(\vec{x}_1, \dots, \vec{x}_r, \vec{x}_{r+1} + \vec{a}, \dots, \vec{x}_{r+s} + \vec{a})$$



$H^n(\vec{a})$

$$\begin{aligned}
 &:= \int_{\mathbb{R}^4 \times \dots \times \mathbb{R}^4} \mathcal{W}^n(\{\vec{x}_\tau\}_{\tau=1}^r, \{\vec{x}_\theta\}_{\theta=r+1}^{r+s}) \bigotimes_{\tau=1}^r p_\tau(\vec{x}_\tau) \cdot \bigotimes_{\theta=r+1}^{r+s} p_\theta(\vec{x}_{r+\theta} - \vec{a}) \cdot \prod_{\tau=1}^{r+s} d\vec{x}_\tau \\
 &\equiv \int_{\mathbb{R}^4 \times \dots \times \mathbb{R}^4} \mathcal{W}^n(\{\vec{x}_\tau\}_{\tau=1}^r, \{\vec{x}_\theta + \vec{a}\}_{\theta=r+1}^{r+s}) \bigotimes_{\tau=1}^{r+s} p_\tau(\vec{x}_\tau) \cdot \prod_{\tau=1}^{r+s} d\vec{x}_\tau \\
 &= c_R \int_{S_0} \left[\prod_{\tau=1}^r \int_{x_\tau^- \in \mathbb{R}^2} E(x_\tau^-) h_\tau(x_\tau^-, x^+) dx_\tau^- \cdot \prod_{\theta=r+1}^{r+s} \int_{x_\theta^- \in \mathbb{R}^2} E(x_\theta^-) h_\theta^{\vec{a}}(x_\theta^-, x^+) dx_\theta^- \right] dx^+ \\
 &+ \sum_{\substack{Q \neq R \\ Q \in \Gamma}} c_Q \int_Q \{h_\theta^{\vec{a}}\}_{\theta=1}^{r+s}
 \end{aligned}$$

$$E(x_\tau^-) = E(x_\tau^0, x_\tau^1) = \frac{e^{i\chi(\tau)[x_\tau^0 \hat{H} - x_\tau^1 \hat{P}]}}{2\pi}$$

$$\mathcal{W}^n(\{\vec{x}_\tau\}_{\tau=1}^r, \{\vec{x}_\theta + \vec{a}\}_{\theta=r+1}^{r+s}) = \mathcal{W}^n(\vec{\xi}_1, \dots, \vec{\xi}_{r-1}, \vec{\xi}_r - \vec{a}, \vec{\xi}_{r+1}, \dots, \vec{\xi}_{r+s-1})$$

$$\widehat{W}^n(\vec{p}_1, \dots, \vec{p}_{r+s})$$

$$= (2\pi)^4 \delta\left(\sum_{\tau=1}^{r+s} \vec{p}_\tau\right) \widehat{W}^n(\vec{p}_1, \vec{p}_1 + \vec{p}_2, \dots, \vec{p}_1 + \vec{p}_2 + \dots + \vec{p}_{r+s-1}).$$

$$\vec{x}_{\delta_i} - \vec{x}_{\epsilon_i} = \vec{\xi}_{\delta_i} + \vec{\xi}_{\delta_i+1} + \dots + \vec{\xi}_{\epsilon_i-1}$$

$$\vec{x}_{\delta_i} - \vec{x}_{\epsilon_i} = -(\vec{\xi}_{\epsilon_i} + \vec{\xi}_{\epsilon_i+1} + \dots + \vec{\xi}_{\delta_i-1})$$

$$\sum_{\tau=1}^r \chi(\tau)[x_\tau^0 \hat{H} - x_\tau^1 \hat{P}] + \sum_{\theta=r+1}^{r+s} \chi(\theta)[x_\theta^0 \hat{H} - x_\theta^1 \hat{P}] = - \sum_{i=1}^{r+s-1} c_i [\xi_i^0 \hat{H} - \xi_i^1 \hat{P}]$$

$$\prod_{\tau=1}^r E(x_\tau^-) \prod_{\theta=r+1}^{r+s} E(x_\theta^-) = \frac{1}{(2\pi)^{r+s}} \exp\left[-i \sum_{j=1}^{r+s-1} c_j [\xi_j^0 \hat{H} - \xi_j^1 \hat{P}]\right]$$

$$\mathcal{W}^n(\vec{\xi}_1, \dots, \vec{\xi}_{r+s-1}) := \mathcal{W}^n(\{\vec{x}_\tau\}_{\tau=1}^r, \{\vec{x}_\theta\}_{\theta=r+1}^{r+s})$$

$$\mathcal{W}^n(\{\vec{x}_\tau\}_{\tau=1}^r, \{\vec{x}_\theta + \vec{a}\}_{\theta=r+1}^{r+s}) = \mathcal{W}^n(\vec{\xi}_1, \dots, \vec{\xi}_{r-1}, \vec{\xi}_r - \vec{a}, \vec{\xi}_{r+1}, \dots, \vec{\xi}_{r+s-1}),$$

$$\int_{\mathbb{R}^2} e^{-ia^- \cdot q^-} \mathcal{W}^n(\vec{\xi}_1, \dots, \vec{\xi}_t - a^-, \vec{\xi}_{t+1}, \dots, \vec{\xi}_{r+s-1}) da^-$$

$$\mathcal{W}^n(\vec{\xi}_1, \dots, \vec{\xi}_t - a^-, \vec{\xi}_{t+1}, \dots, \vec{\xi}_{r+s-1}) = e^{-ic_t a^- \cdot \vec{a}^-} \mathcal{W}^n(\vec{\xi}_1, \dots, \vec{\xi}_t, \vec{\xi}_{t+1}, \dots, \vec{\xi}_{r+s-1})$$

$$\int_{\mathbb{R}^2} e^{-iq^- \cdot a^-} \mathcal{W}^n(\vec{\xi}_1, \dots, \vec{\xi}_t - a^-, \vec{\xi}_{t+1}, \dots, \vec{\xi}_{r+s-1}) da^-$$

$$= \int_{\mathbb{R}^2} e^{-ia^- \cdot q^-} e^{-ic_t a^- \cdot \vec{a}^-} \mathcal{W}^n(\vec{\xi}_1, \dots, \vec{\xi}_t, \vec{\xi}_{t+1}, \dots, \vec{\xi}_{r+s-1}) da^-$$

$$= 2\pi \mathcal{W}^n(\vec{\xi}_1, \dots, \vec{\xi}_t, \dots, \vec{\xi}_{r+s-1}) \cdot \delta(q^- + c_t(\hat{H}(\rho_n), \hat{P}(\rho_n), 0, 0))$$

$$a^- \mapsto F^-(a^-) \equiv \mathcal{W}^n(\vec{\xi}_1, \dots, \vec{\xi}_r - a^-, \vec{\xi}_{r+1}, \dots, \vec{\xi}_{r+s-1}),$$

$$a^- \mapsto F^+(a^-) = \mathcal{W}^n(\{\vec{x}_\theta + a^-\}_{\theta=1}^s, \{\vec{x}_\tau\}_{\tau=s+1}^{s+r})$$

$$\equiv \mathcal{W}^n(\vec{\xi}_1, \dots, \vec{\xi}_s + a^-, \vec{\xi}_{s+1}, \dots, \vec{\xi}_{r+s-1})$$

$$\vec{k} + \vec{k} = \vec{l} + \vec{l}.$$

$$G(a) = \begin{cases} a + \vec{k}, & 1 \leq a \leq \vec{l} \\ a + \vec{k} + \vec{k}, & \vec{l} + 1 \leq a \leq \vec{k} \\ a + \vec{k} + \vec{k} - \vec{l} \equiv a + \vec{l}, & r + 1 \leq a \leq r + \vec{k} \end{cases}$$

$$\tilde{G}(a) = \begin{cases} a + \vec{k}, & 1 \leq a \leq \vec{k} \\ a + \vec{k} - \vec{l}, & r + 1 \leq a \leq r + \vec{l} - \vec{k} \\ a + \vec{k} - \vec{l} + \vec{k} \equiv a + \vec{l}, & r + \vec{l} - \vec{k} + 1 \leq a \leq r + \vec{k} \end{cases}$$



$$\begin{aligned} \mathcal{W}^n(\{\vec{x}_\tau\}_{\tau=1}^r, \{\vec{x}_\theta\}_{\theta=r+1}^{r+s}) &= \pm \mathcal{W}^n(\{\vec{x}_\theta\}_{\theta=r+1}^{r+s}, \{\vec{x}_\tau\}_{\tau=1}^r) \\ \mathcal{W}_{\kappa_r, \pi_s}^n: f_1 \otimes_{\mathbb{R}} \cdots \otimes_{\mathbb{R}} f_r \otimes_{\mathbb{R}} g_1 \otimes_{\mathbb{R}} \cdots \otimes_{\mathbb{R}} g_s &\mapsto \langle A_r^n P_0 B_s^n 1, 1 \rangle, \\ \mathcal{W}_{\pi_s, \kappa_r}^n: g_1 \otimes_{\mathbb{R}} \cdots \otimes_{\mathbb{R}} g_s \otimes_{\mathbb{R}} f_1 \otimes_{\mathbb{R}} \cdots \otimes_{\mathbb{R}} f_r &\mapsto \langle C_r^n P_0 D_s^n 1, 1 \rangle, \end{aligned}$$

$$\begin{aligned} &\int_{\mathbb{R}^4 \times \cdots \times \mathbb{R}^4} \mathcal{W}^n(\{\vec{x}_\tau\}_{\tau=1}^r, \{\vec{x}_\theta\}_{\theta=r+1}^{r+s}) \bigotimes_{\tau=1}^{r+s} p_\tau(\vec{x}_\tau) \cdot \prod_{\tau=1}^{r+s} d\vec{x}_\tau \\ &= \sum_{\substack{(\kappa_r, \bar{\kappa}_r) \in \Omega_r \\ (\pi_s, \bar{\pi}_s) \in \Omega_s}} \int_{\mathbb{R}^4 \times \cdots \times \mathbb{R}^4} \mathcal{W}_{\kappa_r, \pi_s}^n(\{\vec{x}_\tau\}_{\tau=1}^r, \{\vec{x}_\theta\}_{\theta=r+1}^{r+s}) \prod_{\tau=1}^{r+s} q_\tau^{\{\kappa_r, \pi_s\}}(\vec{x}_\tau) \cdot \prod_{\tau=1}^{r+s} d\vec{x}_\tau \\ &\quad \mathcal{W}_{\kappa_r, \pi_s}^n(\{\vec{x}_\tau\}_{\tau=1}^r, \{\vec{x}_\theta\}_{\theta=r+1}^{r+s}) = \pm \mathcal{W}_{\pi_s, \kappa_r}^n(\{\vec{x}_\theta\}_{\theta=r+1}^{r+s}, \{\vec{x}_\tau\}_{\tau=1}^r) \\ T_1^n &:= \int_{\mathbb{R}^4 \times \cdots \times \mathbb{R}^4} \prod_{\tau=1}^r d\vec{x}_\tau \cdot \varphi_1(\{\vec{x}_\tau\}_{\tau=1}^r) \cdot \prod_{\tau=1}^r \psi^{\alpha_\tau, n}(\vec{x}_\tau) \\ T_2^n &:= \int_{\mathbb{R}^4 \times \cdots \times \mathbb{R}^4} \prod_{\theta=r+1}^{r+s} d\vec{x}_\theta \cdot \varphi_2(\{\vec{x}_\theta\}_{\theta=r+1}^{r+s}) \cdot \prod_{\theta=r+1}^{r+s} \psi^{\beta_{\theta-r}, n}(\vec{x}_\theta) \\ \mathcal{T}_1^n(\{\vec{x}_\tau\}_{\tau=s+1}^r) &:= \int_{\mathbb{R}^4 \times \cdots \times \mathbb{R}^4} \prod_{\theta=r+1}^{r+s} d\vec{x}_\theta \cdot \varphi_1(\{\vec{x}_\tau\}_{\tau=s+1}^r, \{\vec{x}_\theta\}_{\theta=r+1}^{r+s}) \cdot \prod_{\theta=r+1}^{r+s} \psi^{\beta_{\theta-r}, n}(\vec{x}_\theta) \\ \mathcal{T}_2^n(\{\vec{x}_\tau\}_{\tau=s+1}^r) &:= \int_{\mathbb{R}^4 \times \cdots \times \mathbb{R}^4} \prod_{\tau=1}^s d\vec{x}_\tau \cdot \varphi_2(\{\vec{x}_\tau\}_{\tau=1}^s) \cdot \prod_{\tau=1}^r \psi^{\alpha_\tau, n}(\vec{x}_\tau) \\ &\quad T_1^n(\vec{a}) := U(\vec{a}, 1) T_1^n U(\vec{a}, 1)^{-1} \\ &\quad T_2^n(\vec{a}) := U(\vec{a}, 1) T_2^n U(\vec{a}, 1)^{-1} \\ &\quad \mathcal{T}_1^n(\{\vec{x}_\tau\}_{\tau=s+1}^r, \vec{a}), \mathcal{T}_2^n(\{\vec{x}_\tau\}_{\tau=s+1}^r, \vec{a}) \\ T_1 &:= \sum_{n=1}^{\infty} c_{1,n} T_1^n, T_2 := \sum_{n=1}^{\infty} c_{2,n} T_2^n \\ \mathcal{T}_1(\{\vec{x}_\tau\}_{\tau=s+1}^r) &:= \sum_{n=1}^{\infty} c_{1,n} \mathcal{T}_1^n(\{\vec{x}_\tau\}_{\tau=s+1}^r) \\ \mathcal{T}_2(\{\vec{x}_\tau\}_{\tau=s+1}^r) &:= \sum_{n=1}^{\infty} c_{2,n} \mathcal{T}_2^n(\{\vec{x}_\tau\}_{\tau=s+1}^r) \\ |\hat{H}^k T_j(\cdot, \vec{a}) 1|^2 &\leq \sum_{n=1}^{\infty} C(\rho_n)^{2k} |c_{j,n}|^2 |T_j^n(\cdot) 1|^2 < \infty \\ \langle T_1(\vec{y}) T_2(\vec{x}) 1, 1 \rangle &= \sum_{n=1}^{\infty} c_{1,n} c_{2,n} \langle U(\vec{y}, 1) T_1^n U(\vec{y}, 1)^{-1} U(\vec{x}, 1) T_2^n U(\vec{x}, 1)^{-1} 1, 1 \rangle \\ &= \sum_{n=1}^{\infty} c_{1,n} c_{2,n} \langle T_1^n U(\vec{x} - \vec{y}, 1) T_2^n 1, 1 \rangle \\ &\quad \langle \mathcal{T}_2(\{\vec{x}_\tau\}_{\tau=s+1}^r, \vec{x}) \mathcal{T}_1(\{\vec{x}_\tau\}_{\tau=s+1}^r, \vec{y}) 1, 1 \rangle \\ &= \sum_{n=1}^{\infty} c_{1,n} c_{2,n} \langle \mathcal{T}_2^n(\{\vec{x}_\tau\}_{\tau=s+1}^r, \vec{x}) \mathcal{T}_1^n(\{\vec{x}_\tau\}_{\tau=s+1}^r, \vec{y}) 1, 1 \rangle \end{aligned}$$



$$\begin{aligned}
h_{12}^n(\vec{a}) &:= \langle T_1^n(\vec{y})P_0T_2^n(\vec{x})1,1 \rangle \\
&= \langle T_1^nU(\vec{y},1)^{-1}U(\vec{x},1)T_2^n1,1 \rangle - \langle T_1^n1,1 \rangle \langle T_2^n1,1 \rangle \\
&= \langle T_1^nU(\vec{a},1)T_2^n1,1 \rangle - \langle T_1^n1,1 \rangle \langle T_2^n1,1 \rangle. \\
h_{21}^n(\vec{a}) &:= \int_{\mathbb{R}^4 \times \dots \times \mathbb{R}^4} \prod_{\tau=s^+}^r d\vec{x}_\tau \langle \mathcal{T}_2^n(\{\vec{x}_\tau\}_{\tau=s^+}^r, \vec{x})P_0\mathcal{T}_1^n(\{\vec{x}_\tau\}_{\tau=s^+}^r, \vec{y})1,1 \rangle \\
&= \left\langle \int_{\mathbb{R}^{4(r-s)}} \prod_{\tau=s^+}^r d\vec{x}_\tau \mathcal{T}_2^n(\{\vec{x}_\tau\}_{\tau=s^+}^r)U(-\vec{a},1)\mathcal{T}_1^n(\{\vec{x}_\tau\}_{\tau=s^+}^r)1,1 \right\rangle \\
&\quad - \int_{\mathbb{R}^{4(r-s)}} \prod_{\tau=s^+}^r d\vec{x}_\tau \langle \mathcal{T}_1^n(\{\vec{x}_\tau\}_{\tau=s^+}^r)1,1 \rangle \langle \mathcal{T}_2^n(\{\vec{x}_\tau\}_{\tau=s^+}^r)1,1 \rangle \\
h_{12}^n(\vec{a}) &:= \int_{\mathbb{R}^4 \times \dots \times \mathbb{R}^4} \prod_{\tau=1}^{r+s} d\vec{x}_\tau \mathcal{W}^n(\{\vec{x}_\tau\}_{\tau=1}^r, \{\vec{x}_\theta + \vec{a}\}_{\theta=r^+}^{r+s})\varphi_1(\{\vec{x}_\tau\}_{\tau=1}^r)\varphi_2(\{\vec{x}_\theta\}_{\theta=r^+}^{r+s}) \\
h_{21}^n(\vec{a}) &:= \int_{\mathbb{R}^4 \times \dots \times \mathbb{R}^4} \prod_{\tau=1}^{r+s} d\vec{x}_\tau \mathcal{W}^n(\{\vec{x}_\theta + \vec{a}\}_{\theta=r^+}^{r+s}, \{\vec{x}_\tau\}_{\tau=1}^r)\varphi_1(\{\vec{x}_\tau\}_{\tau=1}^r)\varphi_2(\{\vec{x}_\theta\}_{\theta=r^+}^{r+s}) \\
&\quad \langle T_1^nU(\vec{a},1)T_2^n1,1 \rangle - \langle T_1^n1,1 \rangle \langle T_2^n1,1 \rangle \text{ and} \\
&\quad \left\langle \int_{\mathbb{R}^{4(r-s)}} \prod_{\tau=s^+}^r d\vec{x}_\tau \mathcal{T}_2^n(\{\vec{x}_\tau\}_{\tau=s^+}^r)U(-\vec{a},1)\mathcal{T}_1^n(\{\vec{x}_\tau\}_{\tau=s^+}^r)1,1 \right\rangle \\
&\quad - \int_{\mathbb{R}^{4(r-s)}} \prod_{\tau=s^+}^r d\vec{x}_\tau \langle \mathcal{T}_1^n(\{\vec{x}_\tau\}_{\tau=s^+}^r)1,1 \rangle \langle \mathcal{T}_2^n(\{\vec{x}_\tau\}_{\tau=s^+}^r)1,1 \rangle \\
h_{12}(\vec{a}) &:= \sum_{n=1}^{\infty} c_{1,n}c_{2,n}h_{12}^n(\vec{a}), h_{21}(\vec{a}) := \sum_{n=1}^{\infty} c_{1,n}c_{2,n}h_{21}^n(\vec{a}). \\
&\quad \frac{\hat{H}(\rho_n)}{m_n}e_0 + \frac{\hat{P}(\rho_n)}{m_n}e_1, \frac{\hat{P}(\rho_n)}{m_n}e_0 + \frac{\hat{H}(\rho_n)}{m_n}e_1 \\
&\quad \mathcal{W}^n(\{\vec{x}_\tau\}_{\tau=1}^r, \{\vec{x}_\theta + \vec{a}\}_{\theta=r^+}^{r+s}) = 0 \\
&\quad \mathcal{W}^n(\{\vec{x}_\tau\}_{\tau=1}^r, \{\vec{x}_\theta + \vec{a}\}_{\theta=r^+}^{r+s}) \\
&= e^{-i\vec{a} \cdot c_r m_n \int_0^1} \prod_{\theta=1}^{r+s} E(x_\theta^-) \cdot \mathcal{W}_0^n(\{x_\tau^+\}_{\tau=1}^r, \{x_\theta^+ + a^+\}_{\theta=r^+}^{r+s}), \\
&\quad E(x_\theta^-) = E(x_\theta^0, x_\theta^1) = \frac{e^{i\chi(\theta)[x_\theta^0 \hat{H} - x_\theta^1 \hat{P}]}}{2\pi} (0^- \equiv (0,0)) \\
&\quad \mathcal{W}_0^n(\{x_\tau^+\}_{\tau=1}^r, \{x_\theta^+ + a^+\}_{\theta=r^+}^{r+s}) = \mathcal{W}^n(\{0^-, x_\tau^+\}_{\tau=1}^r, \{0^-, x_\theta^+ + a^+\}_{\theta=r^+}^{r+s}) \\
&\quad \mathcal{W}_{\kappa_r, \pi_s}^n(\{\vec{x}_\tau\}_{\tau=1}^r, \{\vec{x}_\theta + \vec{a}\}_{\theta=r^+}^{r+s}) = \pm \mathcal{W}_{\pi_s, \kappa_r}^n(\{\vec{x}_\theta + \vec{a}\}_{\theta=r^+}^{r+s}, \{\vec{x}_\tau\}_{\tau=1}^r) \\
&\quad \mathcal{W}_{\kappa_r, \pi_s}^n(\{\vec{x}_\tau\}_{\tau=1}^r, \{\vec{x}_\theta + \vec{a}\}_{\theta=r^+}^{r+s}) = e^{ia^0 c_r m_n} \mathcal{W}_{\kappa_r, \pi_s}^n(\{\vec{x}_\tau\}_{\tau=1}^r, \{\vec{x}_\theta + a^+\}_{\theta=r^+}^{r+s}), \\
&\quad \mathcal{W}_{\pi_s, \kappa_r}^n(\{\vec{x}_\theta + \vec{a}\}_{\theta=r^+}^{r+s}, \{\vec{x}_\tau\}_{\tau=1}^r) = e^{-ia^0 c_s m_n} \mathcal{W}_{\pi_s, \kappa_r}^n(\{\vec{x}_\theta + a^+\}_{\theta=r^+}^{r+s}, \{\vec{x}_\tau\}_{\tau=1}^r) \\
|h_{12}(\vec{a})| &= \left| \sum_{n=1}^{\infty} c_{1,n}c_{2,n}h_{12}^n(\vec{a}) \right| \\
&\leq \frac{C}{m_0^l + |a^+|^k} (\|\varphi_1\|_{p,q} \|\varphi_2\|_{p,q}), \\
&\quad f_1 \otimes_{\mathbb{R}} f_2 \otimes_{\mathbb{R}} \dots \otimes_{\mathbb{R}} f_r \text{ and } g_1 \otimes_{\mathbb{R}} \dots \otimes_{\mathbb{R}} g_s \\
&\quad \varphi(\{\vec{x}_\theta\}_{\theta=1}^{r+s}) = \varphi_1(\{\vec{x}_\tau\}_{\tau=1}^r) \otimes_{\mathbb{R}} \varphi_2(\{\vec{x}_\theta\}_{\theta=r^+}^{r+s})
\end{aligned}$$

$$\begin{aligned}
\tilde{\varphi}_1(\{q_\tau^-, x_\tau^+\}_{\tau=1}^r) &:= \int_{\mathbb{R}^{2r}} \prod_{\tau=1}^r \frac{e^{-i\chi(\tau)x_\tau^- \cdot q_\tau^-}}{2\pi} \cdot \varphi_1(\{x_\tau^-, x_\tau^+\}_{\tau=1}^r) \prod_{\tau=1}^{r+s} dx_\tau^- \\
\tilde{\varphi}_2(\{q_\theta^-, x_\theta^+\}_{\theta=r+1}^{r+s}) &:= \int_{\mathbb{R}^{2s}} \prod_{\theta=r+1}^{r+s} \frac{e^{-i\chi(\theta)x_\theta^- \cdot q_\theta^-}}{2\pi} \cdot \varphi_2(\{x_\theta^-, x_\theta^+\}_{\theta=r+1}^{r+s}) \prod_{\theta=r+1}^{r+s} dx_\theta^- \\
\tilde{\varphi}(\{q_\tau^-, x_\tau^+\}_{\tau=1}^{r+s}) &:= \tilde{\varphi}_1(\{q_\tau^-, x_\tau^+\}_{\tau=1}^r) \tilde{\varphi}_2(\{q_\theta^-, x_\theta^+\}_{\theta=r+1}^{r+s}) \\
h_{12}^n(\vec{a}) &= e^{-ic_r m_n \vec{a} \cdot \vec{f}_0^n} \int_{\mathbb{R}^{2(r+s)}} \prod_{\tau=1}^{r+s} dx_\tau^+ \mathcal{W}_0^n(\{x_\tau^+\}_{\tau=1}^r, \{x_\theta^+ + a^+\}_{\theta=r+1}^{r+s}) \tilde{\varphi}(\{H_n^-, x_\tau^+\}_{\tau=1}^{r+s}) \\
|\tilde{\varphi}(\{q_\theta^-, x_\theta^+\}_{\theta=1}^{r+s})| &\leq C(\rho_n)^{\tilde{n}} \frac{\|\varphi\|_{\tilde{k}, \tilde{l}}}{\sum_{\theta=1}^{r+s} (|q_\theta^0|^2 + |q_\theta^1|^2)^{l/2} + (|x_\theta^2|^2 + |x_\theta^3|^2)^{k/2}} \\
|\tilde{\varphi}(\{H_n^-, x_\theta^+\}_{\theta=1}^{r+s})| &\leq C(\rho_n)^{\tilde{n}} \frac{\|\varphi\|_{\tilde{k}, \tilde{l}}}{\sum_{\theta=1}^{r+s} m_n^l + (x_\theta^{2,2} + x_\theta^{3,2})^{k/2}} \\
e^{-ic_r m_n \vec{a} \cdot \vec{f}_0^n} \sum_{|\vec{m}| \leq \mathcal{N}} \int_{\mathbb{R}^{2(r+s)}} \prod_{\tau=1}^{r+s} dx_\tau^+ D^{\vec{m}} G_{\vec{m}}^n(\{x_\tau^+\}_{\tau=1}^{r+s}; a^+) \tilde{\varphi}(\{H_n^-, x_\theta^+\}_{\theta=1}^{r+s}) \\
&= h_{12}^n(\vec{a}) \\
&\sum_{|\vec{m}| \leq \mathcal{N}} D^{\vec{m}} G_{\vec{m}}^n(\{x_\tau^+\}_{\tau=1}^{r+s}; a^+) \equiv \mathcal{W}_0^n(\{x_\tau^+\}_{\tau=1}^r, \{x_\theta^+ + a^+\}_{\theta=r+1}^{r+s}) \\
&\sum_{|\vec{m}| \leq \mathcal{N}} |G_{\vec{m}}^n(\{x_\tau^+\}_{\tau=1}^{r+s}; a^+) \leq C(\rho_n)^{\tilde{k}} \left[|a^+|^\alpha + \left(\sum_{\tau=1}^{r+s} |x_\tau^+|^2 \right)^{\gamma/2} \right] \\
\{R > R_0 - \epsilon\} &:= \left\{ \{x_\tau^+\}_{\tau=1}^{r+s} : \sum_{\tau=1}^{r+s} |x_\tau^+|^2 > (R_0 - \epsilon)^2 \right\} \subset \mathbb{R}^{2(r+s)} \\
&\sum_{|\vec{m}| \leq \mathcal{N}} \int_{\mathbb{R}^{2(r+s)}} \prod_{\tau=1}^{r+s} dx_\tau^+ D^{\vec{m}} G_{\vec{m}}^n(\{x_\tau^+\}_{\tau=1}^{r+s}; a^+) \tilde{\varphi}(\{q_\theta^-, x_\theta^+\}_{\theta=1}^{r+s}) \\
&= \sum_{|\vec{m}| \leq \mathcal{N}} \int_{\{R > R_0 - \epsilon\}} \prod_{\tau=1}^{r+s} dx_\tau^+ D^{\vec{m}} G_{\vec{m}}^n(\{x_\tau^+\}_{\tau=1}^{r+s}; a^+) \tilde{\varphi}(\{q_\theta^-, x_\theta^+\}_{\theta=1}^{r+s}) \\
|h_{12}^n(\vec{a})| &\leq \sum_{|\vec{m}| \leq \mathcal{N}} \int_{\{R > R_0 - \epsilon\}} \prod_{\tau=1}^{r+s} dx_\tau^+ |G_{\vec{m}}^n(\{x_\tau^+\}_{\tau=1}^{r+s}; a^+) D^{\vec{m}} \tilde{\varphi}(\{H_n^-, x_\theta^+\}_{\theta=1}^{r+s})| \\
&\leq C(\rho_n)^{\tilde{K}} \int_{\{R > R_0 - \epsilon\}} \frac{[|a^+|^\alpha + R^\gamma]}{(r+s)|m_n|^l + R^{k+l}} \frac{\|\varphi\|_{p,q}}{R} R^{2(r+s)} dR d\Omega \\
&\quad C(\rho_n)^{\tilde{K}} \frac{\|\varphi_1\|_{p,q} \|\varphi_2\|_{p,q}}{m_n^l + |a^+|^k} \\
&\quad \mathcal{W}_{k_r, \pi_s}^n(\{\vec{x}_\tau\}_{\tau=1}^r, \{\vec{x}_\theta + \vec{a}\}_{\theta=r+1}^{r+s}) \rightarrow 0 \\
|h_{12}(\vec{a})| &\leq \sum_{n=1}^{\infty} |c_{1,n} c_{2,n}| |h_{12}^n(\vec{a})| \\
&\leq \frac{C}{m_0^l + |a^+|^k} (\|\varphi_1\|_{p,q} \|\varphi_2\|_{p,q}) \\
\{y_j^+\}^{|A_i|} &= \underbrace{(H_n^-, y_i^+, H_n^-, y_i^+, \dots, H_n^-, y_i^+)}_{|A_i| \text{ copies of } (H_n^-, y_i^+)}.
\end{aligned}$$



$$\begin{aligned}
& \tilde{\varphi}_1(\{H_n^-, x_\tau^+\}_{\tau=1}^r) \text{ and } \tilde{\varphi}_2(\{H_n^-, x_\theta^+\}_{\theta=r+1}^{r+s}). \\
& \tilde{\varphi}_{Q,1}^n(\{x_\theta^+\}_{\theta=n_{z-1}+1}^r) \\
& := \int_{\mathbb{R}^{2(z-1)}} \prod_{i=1}^{z-1} dy_i^+ \tilde{\varphi}_1(\{y_j^+\}^{|A_1|}, \{y_j^+\}^{|A_2|}, \dots, \{y_j^+\}^{|A_{z-1}|}, \{H_n^-, x_\theta^+\}_{\theta=n_{z-1}+1}^r) \\
& \tilde{\varphi}_{Q,2}^n(\{x_\theta^+\}_{\theta=r+1}^{n_z}) \\
& := \int_{\mathbb{R}^{2(n(Q)-z)}} \prod_{i=z^+}^{n(Q)} dy_i^+ \tilde{\varphi}_2(\{H_n^-, x_\theta^+\}_{\theta=r+1}^{n_z}, \{y_j^+\}^{|A_{z+1}|}, \{y_j^+\}^{|A_{z+2}|}, \dots, \{y_j^+\}^{|A_{n(Q)}|}) \\
& \tilde{\varphi}_{Q,1}^n(x_r^+) := \tilde{\varphi}_{Q,1}^n(\{x_r^+\}_{\theta=n_{z-1}+1}^r), \tilde{\varphi}_{Q,2}^n(x_{r+1}^+) := \tilde{\varphi}_{Q,2}^n(\{x_{r+1}^+\}_{\theta=r+1}^{n_z}). \\
h_{12}(\vec{a}) &= e^{-ic_r m_n \vec{a} \cdot \vec{f}_0^n} \sum_{Q \in \Omega_r} c_Q \int_{\mathbb{R}^2} dx^+ \tilde{\varphi}_{Q,1}^n(x^+) \tilde{\varphi}_{Q,2}^n(x^+ - a^+) \\
&= e^{-ic_r m_n \vec{a} \cdot \vec{f}_0^n} \sum_{Q \in \Omega_r} c_Q \int_{\mathbb{R}^2} dx_r^+ \int_{\mathbb{R}^2} dx_{r+1}^+ \tilde{\varphi}_{Q,1}^n(x_r^+) \delta(a^+ + x_{r+1}^+ - x_r^+) \tilde{\varphi}_{Q,2}^n(x_{r+1}^+),
\end{aligned}$$

$$|h_{12}(\vec{a})| = \left| \sum_{n=1}^{\infty} c_{1,n} c_{2,n} h_{12}^n(\vec{a}) \right|$$

$$\leq C \frac{e^{-m_0 |a^+|}}{|a^+| - \epsilon} \|\varphi_1\|_{p_1, q_1} \|\varphi_2\|_{p_2, q_2}$$

$$D(\varphi_1, \varphi_2) = \{\vec{x}_\tau - \vec{y}_\theta \in \mathbb{R}^4 \mid \{\vec{x}_\tau\}_{\tau=1}^r \in \text{supp} \varphi_1, \{\vec{y}_\theta\}_{\theta=r+1}^{r+s} \in \text{supp} \varphi_2\}$$

$$\left(\frac{\partial^2}{\partial a^{0,2}} - \hat{P} \right) \Psi_n(a^0, a) := \left(\sum_{i=0}^3 \frac{\partial^2}{\partial a^{i,2}} \right) \Psi_n(a^0, a) = h_{12}^n(a^0, a^+) g(a^1).$$

$$\Psi_n(a^0, a) = e^{ic_r m_n a^0} \tilde{G}(n; a)$$

$$\tilde{G}(n; a)$$

$$= \sum_{Q \in \Omega_r} c_Q \int_{\mathbb{R}} d\xi^1 \int_{\mathbb{R}^{2 \times 2}} dx_r^+ dx_{r+1}^+ \tilde{\varphi}_{Q,1}^n(x_r^+) G_n(a^1 - \xi^1; x_r^+, x_{r+1}^+ + a^+) \tilde{\varphi}_{Q,2}^n(x_{r+1}^+) g(\xi^1),$$

$$\left(-(c_r m_n)^2 + \sum_{i=1}^3 \frac{\partial^2}{\partial a^{i,2}} \right) G_n(a^1 - \xi^1; x_r^+, x_{r+1}^+ + a^+) = \delta(a^1 - \xi^1, x_r^+ - x_{r+1}^+ - a^+)$$

$$\equiv \delta(a - \xi_r)$$

$$G_n(a^1 - \xi^1; x_r^+, x_{r+1}^+ + a^+) = -\frac{1}{(2\pi)^{3/2}} \int_{\mathbb{R}^3} dq \frac{e^{iq \cdot (a - \xi_r)}}{\omega^2 + |q|^2}$$

$$G_n(a^1 - \xi^1; x_r^+, x_{r+1}^+ + a^+) = -\frac{1}{(2\pi)^{1/2}} \frac{1}{iR} \int_0^\infty \lambda \frac{e^{iR\lambda} - e^{-iR\lambda}}{\omega^2 + \lambda^2} d\lambda$$

$$= -\frac{1}{(2\pi)^{1/2}} \frac{1}{iR} \int_{-\infty}^\infty \frac{\lambda e^{iR\lambda}}{(\lambda - i\omega)(\lambda + i\omega)} d\lambda$$

$$G_n(a^1 - \xi^1; x_r^+, x_{r+1}^+ + a^+) = -\frac{2\pi i}{(2\pi)^{1/2}} \frac{1}{iR} \frac{i\omega e^{-R\omega}}{2i\omega} = -\sqrt{2\pi} \frac{e^{-R|c_r m_n|}}{2R}$$

$$R^2 = |a^1 - \xi^1|^2 + |a^2 + \xi_r^2|^2 + |a^3 + \xi_r^3|^2$$

$$(G_n(a - \xi) := G_n(a^1 - \xi^1; x^-, x^+ + a^+))$$



$$\begin{aligned}
|h_{12}^n(\vec{a})g(a^1)| &= \left| \left[\frac{\partial^2}{\partial a^{0,2}} - \hat{\mathcal{P}} \right] \Psi_n(a^0, a) \right| \\
&= \left| \sum_{Q \in \Omega_r} c_Q \int_{\mathbb{R}} d\xi^1 \int_{\mathbb{R}^4} d\vec{x} \tilde{\varphi}_{Q,1}^n(x^-) G_n(a - \xi) [(c_r m_n)^2 + \hat{\mathcal{P}}] \tilde{\varphi}_{Q,2}^n(x^+) g(\xi^1) \right| \\
&\leq \sum_{Q \in \Omega_r} |c_Q| \left\{ \int_{\mathbb{R}} d\xi^1 \left[|g(\xi^1)| + \frac{1}{m_0^2} |g''(\xi^1)| \right] \cdot \frac{\sqrt{\pi} e^{-m_0|a^+ + x^+ - y^+|}}{\sqrt{2|a^+ + x^+ - y^+|}} \right. \\
&\quad \times \left. \int_{\mathbb{R}^2} dy^+ |\tilde{\varphi}_{Q,1}^n(y^+)| \cdot \int_{\mathbb{R}^2} dx^+ |[(c_r m_n)^2 + \hat{\mathcal{P}}] \tilde{\varphi}_{Q,2}^n(x^+)| \right\} \\
&\leq C(\rho_n)^{\bar{k}} e^{m_0 \epsilon} \|g\|_{p_3, q_3} \|\varphi_1\|_{p_1, q_1} \|\varphi_2\|_{p_2, q_2} \cdot \frac{e^{-m_0|a^+|}}{|a^+| - \epsilon}
\end{aligned}$$

$$|h_{12}^n(\vec{a})| \leq C(\rho_n)^{\bar{k}} e^{m_0 \epsilon} \|\varphi_1\|_{p_1, q_1} \|\varphi_2\|_{p_2, q_2} \cdot \frac{e^{-m_0|a^+|}}{|a^+| - \epsilon}$$

$$|\mathcal{W}^n(\{\vec{x}_\tau\}_{\tau=1}^r, \{\vec{x}_\theta + \vec{a}\}_{\theta=r+}^{r+s})| \leq C(n) e^{-m_0|a^+|}$$

$$J_{ab}^\sigma(s, t) = \begin{pmatrix} \sigma'_a(s, t) & \dot{\sigma}_a(s, t) \\ \sigma'_b(s, t) & \dot{\sigma}_b(s, t) \end{pmatrix}, a \neq b$$

$$\rho_\sigma^{ab} = \frac{1}{\sqrt{\det[1 + W_{ab}^{cd, T} W_{ab}^{cd}]}} \equiv \frac{|J_{ab}^\sigma|}{\sqrt{\det[J_{ab}^{\sigma, T} J_{ab}^\sigma + J_{cd}^{\sigma, T} J_{cd}^\sigma]}}$$

$$\int_S d\rho := \sum_{0 \leq a < b \leq 3} \int_{I^2} \rho_\sigma^{ab}(s, t) |J_{ab}^\sigma|(s, t) ds dt$$

$$j_{ab}^\sigma(s, t) = \begin{cases} J_{ab}^\sigma(s, t), & a \neq 0 \\ \begin{pmatrix} i\sigma'_a(s, t) & i\dot{\sigma}_a(s, t) \\ \sigma'_b(s, t) & \dot{\sigma}_b(s, t) \end{pmatrix}, & a = 0 \end{cases}$$

$$\hat{\rho}_\sigma^{ab} := \frac{\det \hat{j}_{ab}^\sigma}{\sqrt{\det[\hat{j}_{ab}^{\sigma, T} \hat{j}_{ab}^\sigma + \hat{j}_{cd}^{\sigma, T} \hat{j}_{cd}^\sigma]}}$$

$$\int_S d\hat{\rho} := \sum_{0 \leq a < b \leq 3} \int_{I^2} \rho_\sigma^{ab}(s, t) [\det J_{ab}^\sigma(\hat{s})] d\hat{s}$$

$$\int_S d|\hat{\rho}| := \int_{I^2} \left| \sum_{0 \leq a < b \leq 3} \rho_\sigma^{ab}(\hat{s}) [\det J_{ab}^\sigma(\hat{s})] \right| d\hat{s}$$

$$\rho_\sigma(\hat{s}) = \sum_{0 \leq a < b \leq 3} \rho_\sigma^{ab}(\hat{s}) [\det J_{ab}^\sigma(\hat{s})]$$

$$|\rho_\sigma|(\hat{s}) = \left| \sum_{0 \leq a < b \leq 3} \rho_\sigma^{ab}(\hat{s}) [\det J_{ab}^\sigma(\hat{s})] \right|$$

$$\sum_{0 \leq a < b \leq 3} \rho_\sigma^{ab} [\det J_{ab}^{\hat{\sigma}}] = \sqrt{[\sigma' \cdot \sigma'] [\dot{\sigma} \cdot \dot{\sigma}] - [\sigma' \cdot \dot{\sigma}]^2}$$

$$\sum_{0 \leq a < b \leq 3} \hat{\rho}_\sigma^{ab} [\det \hat{j}_{ab}^{\hat{\sigma}}] = \sum_{0 \leq a < b \leq 3} \hat{\rho}_\sigma^{ab} [\det j_{ab}^\sigma]$$

$$\Lambda: s e_0 + \bar{s} e_1 \mapsto [\operatorname{sgn}(\bar{s}) \sqrt{\bar{s}^2 - s^2}] e_1,$$



$$\begin{aligned}
\vec{x} + \vec{v} &= s\hat{f}_0 + \bar{s}\hat{f}_1 + x^2\hat{f}_2 + x^3\hat{f}_3 \\
&\rightarrow_{\bar{\Lambda}} (s, \bar{s}, x^2, x^3) \equiv se_0 + \bar{s}e_1 + x^2e_2 + x^3e_3 \rightarrow_{\Lambda} \left(0, \operatorname{sgn}(\bar{s})\sqrt{\bar{s}^2 - s^2}, x^2, x^3\right) \\
&\rightarrow_{R_3} \left(0, \operatorname{sgn}(\bar{s})\sqrt{x^{2,2} + \bar{s}^2 - s^2}, 0, x^3\right) \rightarrow_{R_2} \left(0, \operatorname{sgn}(\bar{s})\sqrt{x^{2,2} + \bar{s}^2 - s^2}, 0, x^3\right) \\
&\rightarrow_{R_3^{-1}} \left(0, \operatorname{sgn}(\bar{s})\sqrt{\bar{s}^2 - s^2}, x^2, x^3\right) \rightarrow_{\Lambda^{-1}} (s, \bar{s}, x^2, x^3) \\
&\rightarrow_{R_1} (-s, -\bar{s}, x^2, x^3) \rightarrow_{\bar{\Lambda}^{-1}} -s\hat{f}_0 - \bar{s}\hat{f}_1 + x^2\hat{f}_2 + x^3\hat{f}_3. \\
\tilde{\Lambda}(\vec{x}, \vec{v}) &:= \bar{\Lambda}^{-1}R_1\Lambda^{-1}R_3^{-1}R_2R_3\Lambda\bar{\Lambda}, \\
v \cdot u &= \sinh(\phi - \theta) = \begin{pmatrix} \cosh(\phi - \theta) \\ \sinh(\phi - \theta) \end{pmatrix} \cdot \begin{pmatrix} 0 \\ 1 \end{pmatrix} \\
v \cdot u &= -\cosh(\phi - \theta) = \begin{pmatrix} \cosh(\phi - \theta) \\ \sinh(\phi - \theta) \end{pmatrix} \cdot \begin{pmatrix} 1 \\ 0 \end{pmatrix} \\
\Lambda(\theta) &= \begin{pmatrix} \cosh \theta & \sinh \theta \\ \sinh \theta & \cosh \theta \end{pmatrix} \\
\Lambda(\theta) \begin{pmatrix} 0 \\ 1 \end{pmatrix} \cdot \Lambda(\phi) \begin{pmatrix} 1 \\ 0 \end{pmatrix} &= \begin{pmatrix} \sinh \theta \\ \cosh \theta \end{pmatrix} \cdot \begin{pmatrix} \cosh \phi \\ \sinh \phi \end{pmatrix} \\
&= -\sinh \theta \cosh \phi + \cosh \theta \sinh \phi = \sinh(\phi - \theta) \\
&= \begin{pmatrix} \cosh(\phi - \theta) \\ \sinh(\phi - \theta) \end{pmatrix} \cdot \begin{pmatrix} 0 \\ 1 \end{pmatrix} \\
\Lambda(\theta) \begin{pmatrix} 1 \\ 0 \end{pmatrix} \cdot \Lambda(\phi) \begin{pmatrix} 1 \\ 0 \end{pmatrix} &= -\cosh(\theta - \phi) = \begin{pmatrix} \cosh(\phi - \theta) \\ \sinh(\phi - \theta) \end{pmatrix} \cdot \begin{pmatrix} 1 \\ 0 \end{pmatrix}.
\end{aligned}$$

6. Las hiperpartículas. En relación a este tipo de partículas, se refiere a aquellas partículas subatómicas las mismas que, alcanzar, igual o superan la velocidad de la luz (energía cinética), con o sin masa, supermasiva o no, de cuyo resultado, se distorsiona el espacio – tiempo. El efecto cuántico gravitacional inherente a esta especie de partículas, es híbrido, es decir, deriva de la permeabilidad del campo de gauge gravitónico, esto es, por interacción con el gravitón, o en su defecto, deriva de su propio movimiento.

Las ecuaciones de movimiento son las siguientes:

$$\begin{aligned}
S(\mathcal{P}_{|0\rangle,t}^{(n)}, \mathcal{P}_{|1\rangle,t}^{(n)}) &= \frac{2 \left| \langle X \rangle_{\mathcal{P}_{|0\rangle,t}^{(n)}} - \langle X \rangle_{\mathcal{P}_{|1\rangle,t}^{(n)}} \right|}{\Delta X_{\mathcal{P}_{|0\rangle,t}^{(n)}} + \Delta X_{\mathcal{P}_{|1\rangle,t}^{(n)}}} \\
S(\mathcal{P}_{|0\rangle,t}^{(1)}, \mathcal{P}_{|1\rangle,t}^{(1)}) &= \frac{2\sqrt{t}|\mu_0 - \mu_1|}{\sqrt{\mu_0} + \sqrt{\mu_1}} \\
\mathcal{P}_{|j\rangle,t}^{(N)}(k) &= (L_{\mu_j t}^{*N})(k) = L_{N\mu_j t}(k) = \mathcal{P}_{|j\rangle, Nt}^{(1)}(k) \\
t_S^{(1)}/t_S^{(N)} &= N \\
S(\mathcal{P}_{|0\rangle,t'}^{(N)}, \mathcal{P}_{|1\rangle,t}^{(N)}) &= \sqrt{N}S(\mathcal{P}_{|0\rangle,t'}^{(1)}, \mathcal{P}_{|1\rangle,t}^{(1)}) \sim t\sqrt{N} \\
\widetilde{C}_p(|00\rangle\langle 00|) &= |00\rangle\langle 00| \\
\widetilde{C}_p(|10\rangle\langle 10|) &= p|00\rangle\langle 00| + (1-p)|11\rangle\langle 11| \\
\mathcal{P}_{|0\rangle,t}^{(N)}(k) &= L_{N\mu_0 t}(k) \\
\mathcal{P}_{|1\rangle,t}^{(N)}(k) &= \sum_{q=0}^N \mathcal{J}_p^{(N)}(q)(L_{\mu_1 t}^{*q} * L_{\mu_0 t}^{*N-q})(k) \\
&= \sum_{q=0}^N \mathcal{J}_p^{(N)}(q)L_{(q\mu_1 + (N-q)\mu_0)t}(k)
\end{aligned}$$

$$\mathcal{W}_{\mu_1, \mu_0, \lambda, t}(k) = e^{-\lambda t} e^{-\mu_1 t} \frac{(\mu_1 t)^k}{k!} + \int_0^t dt' \lambda e^{-\lambda t'} e^{-(\mu_1 t' + \mu_0(t-t'))} \frac{(\mu_1 t' + \mu_0(t-t'))^k}{k!}$$

$$\mathcal{P}_{|0\rangle, t}^{(N)}(k) = L_{N\mu_0 t}(k)$$

$$\mathcal{P}_{|1\rangle, t}^{(N)}(k) = \sum_{q=0}^N \mathcal{J}_p^{(N)}(q) \left(\mathcal{W}_{\mu_1, \mu_0, \lambda, t}^{*q} * L_{(N-q)\mu_0 t} \right) (k)$$

$$\mathcal{S}^{(N)}(P_{|0\rangle, t}, P_{|1\rangle, t}) = \frac{2 \left| \langle K \rangle_{P_{|1\rangle, t}} - \langle K \rangle_{P_{|0\rangle, t}} \right|}{\sqrt{\Delta K_{\mathcal{P}_{|0\rangle, t}}^{(N)}} + \sqrt{\Delta K_{\mathcal{P}_{|1\rangle, t}}^{(N)}}} \times \left| \langle Q \rangle_{\mathcal{J}_{|0\rangle}^{(N)}} + \langle Q \rangle_{\mathcal{J}_{|1\rangle}^{(N)}} - N \right|$$

$$\Delta K_{\mathcal{P}_{|j\rangle, t}^{(N)}}^2 = \langle Q \rangle_{\mathcal{J}_{|j\rangle}^{(N)}} \Delta K_{\mathcal{P}_{|j\rangle, t}^2} + (N - \langle Q \rangle_{\mathcal{J}_{|j\rangle}^{(N)}}) \Delta K_{\mathcal{P}_{|j\rangle, t}^2 + \left(\langle K \rangle_{P_{|j\rangle, t}} - \langle K \rangle_{P_{|\bar{j}\rangle, t}} \right)^2 \Delta Q_{\mathcal{J}_{|j\rangle}^{(N)}}^2$$

$$\mathcal{J}_{curve, p}^{(N)}(q) = (1-p)^q p \theta(0 \leq q \leq N-2) + (1-p)^{N-1} \theta(q = N)$$

$$\mathcal{J}_{curve, p}^{(N=2\tilde{N})}(q) = (1-p) \left(\mathcal{J}_{flat, p}^{(\tilde{N})} * \mathcal{J}_{flat, p}^{(\tilde{N})} \right) (q) + p \theta(q = 0)$$

$$= (1-p)^{q+1} p^2 \left((q+1) \theta(0 \leq q \leq \tilde{N}-2) + (2\tilde{N}-3-q) \theta(\tilde{N}-1 \leq q \leq 2(\tilde{N}-2)) \right) + 2(1-p)^q p \theta(\tilde{N} \leq q \leq 2\tilde{N}-2) + (1-p)^{2\tilde{N}-1} \theta(q = 2\tilde{N}) + p \theta(q = 0)$$

$$\mathcal{J}_{curve, p}^{(N=2\tilde{N}+1)}(q) = (1-p) \left(\mathcal{J}_{flat, p}^{(\tilde{N}+1)} * \mathcal{J}_{flat, p}^{(\tilde{N})} \right) (q) + p \theta(q = 0)$$

$$= (1-p)^{q+1} p^2 \left((q+1) \theta(0 \leq q \leq \tilde{N}-2) + (2\tilde{N}-3-q) \theta(\tilde{N}-1 \leq q \leq 2(\tilde{N}-2)) + \theta(\tilde{N}-1 \leq q \leq 2\tilde{N}-3) \right) + (1-p)^q p \left(2\theta(\tilde{N}+1 \leq q \leq 2\tilde{N}-1) + \theta(q = \tilde{N}) \right) + (1-p)^{2\tilde{N}} \theta(q = 2\tilde{N}+1) + p \theta(q = 0)$$

$$\mathcal{P}_{|j\rangle, t}^{(N)}(k) = \sum_{q=0}^N \mathcal{J}_{|j\rangle}^{(N)}(q) \left(P_{|j\rangle, t}^{*q} * P_{|\bar{j}\rangle, t}^{*(N-q)} \right) (k)$$

$$\langle K \rangle_{\mathcal{P}_{|j\rangle, t}^{(N)}} = \sum_{q=0}^N \mathcal{J}_{|j\rangle}^{(N)}(q) \langle K \rangle_{L_{|j\rangle, q, t}^{(N)}} = \sum_{q=0}^N \mathcal{J}_{|j\rangle}^{(N)}(q) \left(q \langle K \rangle_{P_{|j\rangle, t}} + (N-q) \langle K \rangle_{P_{|\bar{j}\rangle, t} \right) = \langle Q \rangle_{\mathcal{J}_{|j\rangle}^{(N)}} \langle K \rangle_{P_{|j\rangle, t}} + \left(N - \langle Q \rangle_{\mathcal{J}_{|j\rangle}^{(N)}} \right) \langle K \rangle_{P_{|\bar{j}\rangle, t}}$$



$$\begin{aligned}
\Delta K_{\mathcal{P}_{(j),t}}^2 &= \sum_{q=0}^N \mathcal{J}_{|j}^{(N)}(q) \langle K^2 \rangle_{\mathcal{L}_{|j},q,t}^{(N)} - \left(\sum_{q=0}^N \mathcal{J}_{|j}^{(N)}(q) \langle K \rangle_{\mathcal{L}_{|j},q,t}^{(N)} \right)^2 \\
&= \sum_{q=0}^N \mathcal{J}_{|j}^{(N)}(q) \left(\langle K^2 \rangle_{\mathcal{L}_{|j},q,t}^{(N)} - \langle K \rangle_{\mathcal{L}_{|j},q,t}^{(N)2} \right) + \sum_{q=0}^N \mathcal{J}_{|j}^{(N)}(q) \langle K \rangle_{\mathcal{L}_{|j},q,t}^{(N)2} - \left(\sum_{q=0}^N \mathcal{J}_{|j}^{(N)}(q) \langle K \rangle_{\mathcal{L}_{|j},q,t}^{(N)} \right)^2 \\
&= \sum_{q=0}^N \mathcal{J}_{|j}^{(N)}(q) \Delta K_{\mathcal{L}_{|j},q,t}^2 + \sum_{q=0}^N \mathcal{J}_{|j}^{(N)}(q) \left(q \langle K \rangle_{P_{|j},t} + (N-q) \langle K \rangle_{P_{|j},t} \right)^2 \\
&\quad - \left[\langle Q \rangle_{\mathcal{T}_{|j}^{(N)}} \langle K \rangle_{P_{|j},t} + \left(N - \langle Q \rangle_{\mathcal{T}_{|j}^{(N)}} \right) \langle K \rangle_{P_{|j},t} \right]^2 \\
\Delta K_{\mathcal{P}_{|j},t}^2 &= \langle Q \rangle_{\mathcal{T}_{|j}^{(N)}} \Delta K_{P_{|j},t}^2 + \left(N - \langle Q \rangle_{\mathcal{T}_{|j}^{(N)}} \right) \Delta K_{P_{|j},t}^2 + \left(\langle K \rangle_{P_{|j},t} - \langle K \rangle_{P_{|j},t} \right)^2 \Delta Q_{\mathcal{T}_{|j}^{(N)}}^2 \\
\mathcal{S}^{(N)}(P_{|0},t, P_{|1},t) &= \frac{2 \left| \langle K \rangle_{P_{|1},t} - \langle K \rangle_{P_{|0},t} \right|}{\sqrt{\Delta K_{\mathcal{P}_{|0},t}^2} + \sqrt{\Delta K_{\mathcal{P}_{|1},t}^2}} \left| \langle Q \rangle_{\mathcal{T}_{|0}^{(N)}} + \langle Q \rangle_{\mathcal{T}_{|1}^{(N)}} - N \right| \\
\mathcal{F}_\eta(P_{|0}, P_{|1}) &= \frac{\epsilon_{|0} + \epsilon_{|1}}{2} = \frac{P_{|0}(x \geq \eta) + P_{|1}(x < \eta)}{2} \\
e^{-\mu_0 N t} \frac{(\mu_0 N t)^{k_\eta(t)}}{k_\eta(t)!} \leq \epsilon_{|0} &\leq e^{-\mu_0 N t} \frac{(\mu_0 N t)^{k_\eta(t)}}{k_\eta(t)!} \frac{k_\eta(t)}{k_\eta(t) - \mu_0 N t} \\
e^{-\mu_1 N t} \frac{(\mu_1 N t)^{k_\eta(t)-1}}{(k_\eta(t) - 1)!} \leq \epsilon_{|1} &\leq e^{-\mu_1 N t} \frac{(\mu_1 N t)^{k_\eta(t)-1}}{(k_\eta(t) - 1)!} \frac{\mu_1 N t}{\mu_1 N t - (k_\eta(t) - 1)} \\
L_{\mu_0 N t}(k \geq \eta(t)) &\approx e^{-\mu_0 N t} \frac{(\mu_0 N t)^{k_\eta(t)}}{k_\eta(t)!} \approx e^{-\mu_0 N t \gamma_0}, \\
L_{\mu_1 N t}(k < \eta(t)) &\approx e^{-\mu_1 N t} \frac{(\mu_1 N t)^{k_\eta(t)}}{(k_\eta(t) - 1)!} \approx e^{-\mu_1 N t \gamma_1}, \\
\mathcal{P}_{|j},t^{(1)}(\theta) &= G_{z_j(t), \sqrt{|z_j(t)|}}(\theta) \\
\mathcal{L} &= \bar{\psi} (i\gamma^\mu \partial_\mu - m) \psi - \kappa \bar{\psi} \gamma^0 \gamma^5 \psi + \frac{(\bar{\psi} \psi)^2}{M^2}. \\
\mathcal{L} &= -\frac{1}{4} M^2 \Delta^2 + \int \frac{d^4 k}{(2\pi)^4} \ln \det \mathcal{G}^{-1} \\
\mathcal{L} &= -\frac{1}{4} M^2 (\Delta + m)^2 + \int \frac{d^4 k}{(2\pi)^4} \ln U_+ U_- \\
V_{\text{eff}}(\Delta) &= \frac{1}{4} M^2 (\Delta + m)^2 + \frac{\Delta^4}{32\pi^2} (1 + 4 \ln \Delta / \Lambda_{UV}) \\
&\quad - \frac{\kappa^2 \Delta^2}{4\pi^2} (1 - 2 \ln \Delta / \Lambda_{UV}) + V_0 \\
\Delta_0 &= \Lambda_{UV} \exp(-\pi^2 M^2 / 2\kappa^2) \\
\mathcal{L}_1(\Delta, T) &= \int \frac{d^3 k}{(2\pi)^3} \frac{1}{\beta} \sum_{n=-\infty}^{\infty} \ln \beta^4 U_{n+} U_{n-} \\
V_{\text{eff}}(\Delta, T) &= V_{\text{eff}}(T=0) - \frac{2}{\pi^2 \beta^4} [I_2(\beta\Delta) + \kappa^2 \beta^2 I_0(\beta\Delta)] \\
I_n(x) &= \int_0^\infty ds s^n \ln [1 + e^{-\sqrt{s^2+x^2}}] \\
(1+z)^3 &= \frac{w_0(1+w_\Delta)}{w_\Delta(1+w_0)} e^{\frac{1}{w_0} - \frac{1}{w_\Delta}}
\end{aligned}$$

$$\begin{aligned}
V(\Delta) &\simeq V(\Delta_0) + \frac{1}{2}V''(\Delta_0)(\Delta - \Delta_0)^2 = \frac{1}{2}(\kappa\delta\Delta/\pi)^2 \\
\mathcal{M} &= -\frac{1}{\beta} \sum_{i\omega_n} \int \frac{d^3p}{(2\pi)^3} \text{Tr}[i\mathcal{G}(p)i\mathcal{G}(p+k)] \\
\xi^{00} &= \left. \frac{d\mathcal{M}}{dk_0^2} \right|_{k=0} \quad \xi^{33} = \left. \frac{d\mathcal{M}}{d|k|^2} \right|_{k=0} \\
n_A &= -\frac{\partial V(\Delta_0, T)}{\partial \kappa} = \frac{\kappa\Delta_0^2}{2\pi^2} \left(1 + \frac{\pi^2 M^2}{\kappa^2} \right) + \frac{\kappa T^2}{3} \\
n_\Delta &= \left. \frac{\partial V(\Delta, T)}{\partial \Delta} \right|_{\Delta_0} = \frac{1}{2} M^2 \Delta_0. \\
S_E &= \frac{27\pi^2 \left[\int_0^{\Delta_m} d\Delta \sqrt{2V(\Delta)} \right]^4}{2(V(0) - V(\Delta_m))^3} \simeq 2.8 \times 10^7
\end{aligned}$$

CONCLUSIONES.

En mérito a lo anterior, se concluye que la teoría hipotética contenida en este manuscrito, reúne los requisitos de rigor científico.

Los postulados que instituyen la presente teoría son:

La existencia hipotética del gravitón y su partícula compañera.

La existencia hipotética de un campo de gauge gravitónico, el cual, permea el espacio – tiempo cuántico, dotándolo de gravedad, esto, a través de las interacciones del gravitón, ésta última, entendiéndose como partícula fundamental.

La existencia hipotética de agujeros negros de mecánica cuántica, formados por el colapso de aquellas partículas, entendidas como supermasivas.

La distorsión del espacio – tiempo cuántico, deriva del efecto cuántico gravitacional endógeno o exógeno, en relación a la partícula entendida como supermasiva.

La brecha de masa de una partícula, respecto del estado de vacío, superior a cero, comporta energía de potencial y en consecuencia, supone la distorsión del campo de calibre o de gauge en el que interactúa.

Las hiperpartículas, son partículas subatómicas hipotéticas, las cuales y sin perjuicio de que estén dotadas o no de masa, la energía cinética de éstas, es decir, el movimiento acelerado que le permita alcanzar, igual o superar la velocidad de la luz, verbigracia, el taquión, causa la deformación del espacio – tiempo cuántico.

La teoría cuántica de campos relativistas, permite el acoplamiento de las ecuaciones de campo de Einstein a campos de gauge indiscriminados.

Las partículas distorsionadas, es decir, aquellas que se afectan por el efecto cuántico gravitacional de la superpartícula, no se desplazan en trayectorias orbitales de carácter cosmológicas, al contrario, éstas, se difuminan en torno a la partícula deformante, siguiendo la tipología de la distorsión, esto es, los pliegues y el diámetro de la curvatura.

El desarrollo matemático contenido en los manuscritos que preceden, aplican por acoplamiento a esta teoría.

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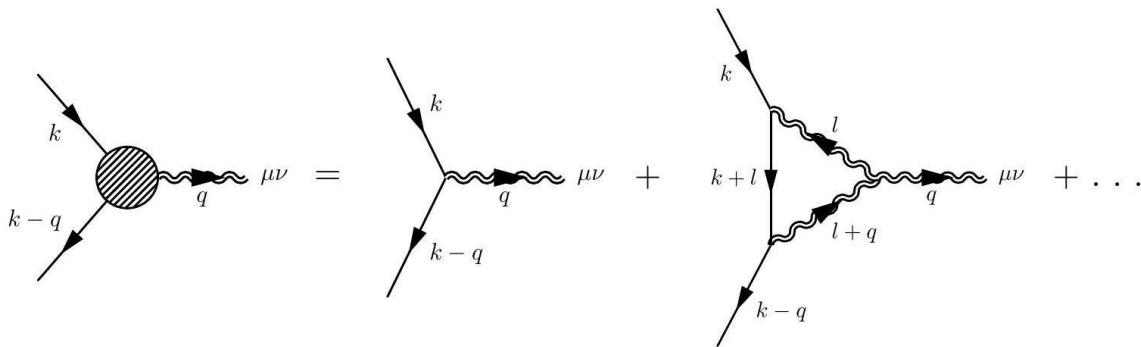


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APÉNDICE A.

Espacios de gauge, reglas de Feynman y métrica de Schwarzschild-Tangherlini para campos cuánticos relativistas o curvos. Fenómenos gravitacionales y sistemas de referencia bajo Relatividad General, a propósito de la interacción y permeo del campo gravitónico y el gravitón en espacios cuánticos deformados, tanto bosónicos como fermiónicos y en relación a las partículas y antipartículas supermasivas propiamente dichas y las hiperpartículas.



$$[G_N] = [\text{mass}]^{-(D-2)}.$$

$$\kappa^2 = 32\pi G_N$$

$$R_{\mu\nu\rho\sigma} = \Gamma_{\mu\nu\sigma,\rho} - \Gamma_{\mu\nu\rho,\sigma} - \Gamma_{\beta\mu\rho}\Gamma_{\nu\sigma}^{\beta} + \Gamma_{\beta\mu\sigma}\Gamma_{\nu\rho}^{\beta}$$

$$\Gamma_{\rho\mu\nu} = \frac{1}{2}(g_{\rho\mu,\nu} + g_{\rho\nu,\mu} - g_{\mu\nu,\rho})$$

$$R_{\mu\nu} = g^{\rho\sigma}R_{\rho\mu\nu\sigma}$$

$$R = g^{\mu\nu}R_{\mu\nu}$$

$$S_{EH} = \int d^D x \sqrt{-g} R$$

$$S_{\phi} = \int d^D x \sqrt{-g} \mathcal{L}_{\phi}$$

$$\mathcal{L}_{\phi} = \frac{1}{2}(g^{\mu\nu}\phi_{,\mu}\phi_{,\nu} - m^2\phi^2)$$

$$S_{(c)} = \frac{2}{\kappa^2}S_{EH} + S_{\phi}$$

$$G^{\mu\nu} = -\frac{\kappa^2}{4} T^{\mu\nu}$$

$$G^{\mu\nu} = R^{\mu\nu} - \frac{1}{2} R g^{\mu\nu}$$

$$\delta S_{EH} = - \int d^D x \sqrt{-g} G^{\mu\nu} \delta g_{\mu\nu}$$

$$D_\mu G^{\mu\nu} = 0$$

$$\delta S_\phi = -\frac{1}{2} \int d^D x \sqrt{-g} T^{\mu\nu} \delta g_{\mu\nu}$$

$$d\tau^2 = \left(1 - \frac{\mu}{r^n}\right) dt^2 - \frac{1}{1 - \frac{\mu}{r^n}} dr^2 - r^2 d\Omega_{D-2}^2$$

$$\mu = \frac{16\pi G_N M}{(D-2)\Omega_{D-2}}$$

$$\Omega_d = \frac{2\sqrt{\pi}^{d+1}}{\Gamma((d+1)/2)}$$

$$S_c = \int d^D x \sqrt{-g} \left(\frac{2R}{\kappa^2} + \frac{1}{2} (g^{\mu\nu} \partial_\mu \phi \partial_\nu \phi - m^2 \phi^2) \right)$$

$$Z_\omega = \int \mathcal{D}g_{\mu\nu} \mathcal{D}\phi \det\left(\frac{\delta G}{\delta \epsilon}\right) \delta(G_\sigma - \omega_\sigma) e^{iS_c}$$

$$\begin{aligned} Z &= \int \mathcal{D}\omega_\sigma Z_\omega \exp\left(i \int d^D x \frac{1}{\kappa^2 \xi} \eta^{\sigma\rho} \omega_\sigma \omega_\rho\right) \\ &= \int \mathcal{D}g_{\mu\nu} \mathcal{D}\phi \mathcal{D}c \mathcal{D}\bar{c} e^{iS_{(c)} + i\frac{2}{\kappa^2} S_{(gf)} + iS_{(gh)}} \end{aligned}$$

$$S_{gf} = \frac{1}{2\xi} \int d^D x \eta^{\sigma\rho} G_\sigma G_\rho$$

$$\begin{aligned} S &= S_{(c)} + \frac{2}{\kappa^2} S_{(gf)} \\ &= \frac{2}{\kappa^2} (S_{EH} + S_{(gf)}) + S_\phi \end{aligned}$$

$$\tilde{\phi}(q) = \int d^D x e^{iqx} \phi,$$

$$\tilde{h}(q) = \int d^D x e^{iqx} \tilde{h}.$$

$$\begin{aligned} \kappa^2 \int d^D x \eta^{\mu\alpha} \eta^{\nu\beta} h_{\mu\nu} h_{\alpha\beta} \eta^{\rho\sigma} \phi_{,\rho} \phi_{,\sigma} &= -\kappa^2 \int \frac{d^D l_{(1)}}{(2\pi)^D} \frac{d^D l_{(2)}}{(2\pi)^D} \frac{d^D p}{(2\pi)^D} \frac{d^D q}{(2\pi)^D} \\ &\times (2\pi)^D \delta^D(l_{(1)} + l_{(2)} + p + q) \eta^{\mu\alpha} \eta^{\nu\beta} \tilde{h}_{\mu\nu}(l_1) \tilde{h}_{\alpha\beta}(l_2) \tilde{\phi}(p) \tilde{\phi}(q) \eta^{\rho\sigma} p_\sigma q_\rho \end{aligned}$$



$$V^{\mu\nu\alpha\beta} = -i4\kappa^2 \frac{1}{2} (\eta^{\mu\alpha}\eta^{\nu\beta} + \eta^{\mu\beta}\eta^{\nu\alpha}) p_\sigma q^\sigma$$

$$p^\mu = \hbar l^\mu$$

$$h_{\mu\nu}(x) = \int \frac{d^D q}{(2\pi)^D} e^{-iqx} \tilde{h}_{\mu\nu}(q)$$

$$\tilde{h}_{\mu\nu}(q) = \int d^D x e^{iqx} h_{\mu\nu}(x)$$

$$\eta_{\mu\nu}^{\parallel} = \frac{k_\mu k_\nu}{m^2}$$

$$\eta_{\mu\nu}^{\perp} = \eta_{\mu\nu} - \frac{k_\mu k_\nu}{m^2}.$$

$$q_{\parallel}^{\mu} = \eta_{\parallel}^{\mu} q^{\nu}$$

$$q_{\perp}^{\mu} = \eta_{\perp}^{\mu} q^{\nu}$$

$$q_{\parallel} = \sqrt{q_{\parallel}^{\mu} q_{\parallel}^{\mu}}$$

$$\hat{g}_{\mu\nu}(\hat{x}) = g_{\alpha\beta}(x) \frac{\partial x^\alpha}{\partial \hat{x}^\mu} \frac{\partial x^\beta}{\partial \hat{x}^\nu}$$

$$x^\mu = \hat{x}^\mu + \hat{\epsilon}^\mu(\hat{x})$$

$$\hat{x}^\mu = x^\mu + \epsilon^\mu(x)$$

$$\epsilon(x) + \hat{\epsilon}(\hat{x}) = 0$$

$$\frac{\partial x^\alpha}{\partial \hat{x}^\mu} = \delta_\mu^\alpha + \frac{\partial \hat{\epsilon}^\alpha(\hat{x})}{\partial \hat{x}^\mu}$$

$$\begin{aligned} \hat{g}_{\mu\nu}(\hat{x}) &= g_{\mu\nu}(x) + g_{\alpha\nu}(x) \frac{\partial \hat{\epsilon}^\alpha(\hat{x})}{\partial \hat{x}^\mu} + g_{\mu\beta}(x) \frac{\partial \hat{\epsilon}^\beta(\hat{x})}{\partial \hat{x}^\nu} + g_{\alpha\beta}(x) \frac{\partial \hat{\epsilon}^\alpha(\hat{x})}{\partial \hat{x}^\mu} \frac{\partial \hat{\epsilon}^\beta(\hat{x})}{\partial \hat{x}^\nu} \\ &= g_{\alpha\beta}(x) \left(\delta_\mu^\alpha \delta_\nu^\beta + \frac{\partial \hat{\epsilon}^\alpha(\hat{x})}{\partial \hat{x}^\mu} \delta_\nu^\beta + \delta_\mu^\alpha \frac{\partial \hat{\epsilon}^\beta(\hat{x})}{\partial \hat{x}^\nu} + \frac{\partial \hat{\epsilon}^\alpha(\hat{x})}{\partial \hat{x}^\mu} \frac{\partial \hat{\epsilon}^\beta(\hat{x})}{\partial \hat{x}^\nu} \right) \end{aligned}$$

$$g_{\mu\nu}(x) = g_{\mu\nu}(\hat{x} + \hat{\epsilon}(\hat{x}))$$

$$= \sum_{n=0.. \infty} \frac{1}{n!} \hat{\epsilon}^\sigma(\hat{x}) \dots \hat{\partial}_\sigma \dots g_{\mu\nu}(\hat{x})$$

$$= g_{\mu\nu}(\hat{x}) + \hat{\epsilon}^\sigma(\hat{x}) \hat{\partial}_\sigma g_{\mu\nu}(\hat{x}) + \frac{1}{2} \hat{\epsilon}^\sigma(\hat{x}) \hat{\epsilon}^\rho(\hat{x}) \hat{\partial}_\sigma \hat{\partial}_\rho g_{\mu\nu}(\hat{x}) + \dots$$

$$= \sum_{n=0.. \infty} \frac{1}{n!} \left(\hat{\epsilon}^\sigma(\hat{x}) \hat{\partial}_\sigma^{(g)} \right)^n g_{\mu\nu}(\hat{x})$$

$$\hat{g}_{\mu\nu} = \left(\delta_\mu^\alpha \delta_\nu^\beta + \partial_\mu \hat{\epsilon}^\alpha \delta_\nu^\beta + \delta_\mu^\alpha \partial_\nu \hat{\epsilon}^\beta + \partial_\mu \hat{\epsilon}^\alpha \partial_\nu \hat{\epsilon}^\beta \right) \sum_{n=0.. \infty} \frac{1}{n!} \left(\hat{\epsilon}^\sigma \partial_\sigma^{(g)} \right)^n g_{\alpha\beta}$$



$$\hat{h}_{\mu\nu} = \partial_\mu \hat{\epsilon}_\nu + \partial_\nu \hat{\epsilon}_\mu + \partial_\mu \hat{\epsilon}_\alpha \partial_\nu \hat{\epsilon}^\alpha + \left(\delta_\mu^\alpha \delta_\nu^\beta + \partial_\mu \hat{\epsilon}^\alpha \delta_\nu^\beta + \delta_\mu^\alpha \partial_\nu \hat{\epsilon}^\beta + \partial_\mu \hat{\epsilon}^\alpha \partial_\nu \hat{\epsilon}^\beta \right) \sum_{n=0..∞} \frac{1}{n!} \left(\hat{\epsilon}^\sigma \partial_\sigma^{(g)} \right)^n h_{\alpha\beta}$$

$$\hat{h}_{\mu\nu} \approx h_{\mu\nu} + \partial_\nu \hat{\epsilon}_\mu + \partial_\mu \hat{\epsilon}_\nu$$

$$S = \int d^D x \sqrt{-g} \left(\frac{2R}{\kappa^2} + \mathcal{L}_\phi \right) + \int d^D x \frac{1}{\kappa^2 \xi} \eta^{\sigma\rho} G_\sigma G_\rho$$

$$g^{\mu\nu} \Gamma_{\mu\nu}^\sigma = 0$$

$$\partial_\mu \left(h_\sigma^\mu - \frac{1}{2} \eta_\sigma^\mu h_\nu^\nu \right) = 0$$

$$g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu}$$

$$g_{\mu\nu} = e^{\pi_{\mu\nu}}$$

$$g^{\mu\nu} = e^{-\pi^{\mu\nu}}$$

$$\sqrt{-g} g_{\mu\nu} = \eta_{\mu\nu} + h'_{\mu\nu}$$

$$G_\sigma = (1 - \alpha) \partial_\mu \left(h_\sigma^\mu - \frac{1}{2} \eta_\sigma^\mu h_\nu^\nu \right) + \alpha g^{\mu\nu} \Gamma_{\sigma\mu\nu}$$

$$\eta^{\mu\nu} \partial_\mu g_{\sigma\nu} - \frac{1}{2} \partial_\sigma \eta^{\mu\nu} g_{\mu\nu} = 0$$

$$I_{\alpha\beta}^{\mu\nu} = \frac{1}{2} \left(\delta_\alpha^\mu \delta_\beta^\nu + \delta_\beta^\mu \delta_\alpha^\nu \right)$$

$$\mathcal{P}_{\alpha\beta}^{\mu\nu} = I_{\alpha\beta}^{\mu\nu} - \frac{1}{2} \eta^{\mu\nu} \eta_{\alpha\beta}$$

$$R = \sum_{n=1..∞} R_{h^n}$$

$$R = \sum_{n=1..∞} R_{G^n}$$

$$g^{\mu\nu} = \eta^{\mu\nu} - h^{\mu\nu} + h_\rho^\mu h^{\rho\nu} - h_\rho^\mu h_\sigma^\rho h^{\sigma\nu} + \dots$$

$$\frac{1}{1+x} = \sum_{n=0..∞} (-x)^n = 1 - x + x^2 - x^3 + \dots$$

$$\hat{h}^{\mu\nu} = g^{\mu\nu} - \eta^{\mu\nu}.$$

$$\hat{h}_{\mu\nu} = -h_{\mu\nu} - \hat{h}_\mu^\sigma h_{\sigma\nu}$$

$$(g^{\mu\nu})_{h^3} = -h_\rho^\mu h_\sigma^\rho h^{\sigma\nu}$$



$$\begin{aligned}\sqrt{-g} &= \exp\left(\frac{1}{2}\text{trln}(\eta_\nu^\mu + h_\nu^\mu)\right) \\ &= \exp\left(\frac{1}{2}\left(h_\mu^\mu - \frac{1}{2}h_\nu^\mu h_\mu^\nu + \frac{1}{3}h_\nu^\mu h_\rho^\nu h_\mu^\rho - \frac{1}{4}h_\nu^\mu h_\rho^\nu h_\sigma^\rho h_\mu^\sigma + \dots\right)\right) \\ &= 1 + \frac{1}{2}h - \frac{1}{4}\mathcal{P}_{\mu\nu}^{\rho\sigma} h^{\mu\nu} h_{\rho\sigma} + \dots\end{aligned}$$

$$\sqrt{-g}_{h^2} = -\frac{1}{4}\mathcal{P}_{\mu\nu}^{\rho\sigma} h^{\mu\nu} h_{\rho\sigma}$$

$$\sqrt{-g}_{h^2}(h_{\mu\nu}^{(a)}, h_{\mu\nu}^{(b)}) = -\frac{1}{4}\mathcal{P}_{\mu\nu}^{\rho\sigma} h^{(a)\mu\nu} h_{\rho\sigma}^{(b)}.$$

$$G^{\mu\nu} = \sum_{n=1..\infty} G_h^{\mu\nu n}$$

$$\begin{aligned}G_h^{\mu\nu}(h_{\mu\nu}) &= G_h^{\mu\nu}\left(\sum_{n=0..\infty} h_{\mu\nu}^{G^n}\right) \\ &= \sum_{n=1..\infty} G_h^{\mu\nu}(h_{\mu\nu}^{G^n})\end{aligned}$$

$$\begin{aligned}G_{h^2}^{\mu\nu}(h_{\mu\nu}, h_{\mu\nu}) &= G_{h^2}^{\mu\nu}\left(\sum_{n=0..\infty} h_{\mu\nu}^{G^n}, \sum_{m=0..\infty} h_{\mu\nu}^{G^m}\right) \\ &= \sum_{n=1..\infty} \sum_{m=1..\infty} G_{h^2}^{\mu\nu}(h_{\mu\nu}^{G^n}, h_{\mu\nu}^{G^m}) \\ &= G_{h^2}^{\mu\nu}(h_{\mu\nu}^G, h_{\mu\nu}^G) + 2G_{h^2}^{\mu\nu}(h_{\mu\nu}^G, h_{\mu\nu}^{G^2}) + \dots\end{aligned}$$

$$(x_1 + x_2 + \dots + x_n)^n$$

$$\begin{aligned}G^{\mu\nu} &\approx G_h^{\mu\nu}(h_{\mu\nu}^G) \\ &+ G_h^{\mu\nu}(h_{\mu\nu}^{G^2}) + G_{h^2}^{\mu\nu}(h_{\mu\nu}^G, h_{\mu\nu}^G) \\ &+ G_h^{\mu\nu}(h_{\mu\nu}^{G^3}) + 2G_{h^2}^{\mu\nu}(h_{\mu\nu}^G, h_{\mu\nu}^{G^2}) + G_{h^3}^{\mu\nu}(h_{\mu\nu}^G, h_{\mu\nu}^G, h_{\mu\nu}^G)\end{aligned}$$

$$S_{EH} + S_{(gf)} = \int d^D x \left(\sqrt{-g} R + \frac{1}{2\xi} \eta^{\rho\sigma} G_\rho G_\sigma \right)$$

$$\begin{aligned}S_{EH} &= \int d^D x \sqrt{-g} R \\ &= \int d^D x \sqrt{-g} g^{\mu\nu} (\Gamma_{\rho\mu,\nu}^\rho - \Gamma_{\mu\nu,\rho}^\rho - \Gamma_{\mu\nu}^\rho \Gamma_{\rho\sigma}^\sigma + \Gamma_{\mu\sigma}^\rho \Gamma_{\nu\rho}^\sigma) \\ &= \int d^D x \sqrt{-g} g^{\mu\nu} (\Gamma_{\mu\nu}^\rho \Gamma_{\rho\sigma}^\sigma - \Gamma_{\mu\sigma}^\rho \Gamma_{\nu\rho}^\sigma)\end{aligned}$$

$$S_{EH} = \int d^D x \sqrt{-g} \frac{1}{4} (2g^{\sigma\gamma} g^{\rho\delta} g^{\alpha\beta} - g^{\gamma\delta} g^{\alpha\beta} g^{\rho\sigma} - 2g^{\sigma\alpha} g^{\gamma\rho} g^{\delta\beta} + g^{\rho\sigma} g^{\alpha\gamma} g^{\beta\delta}) g_{\alpha\beta,\rho} g_{\gamma\delta,\sigma}$$

$$G_\sigma \approx \mathcal{P}_{\rho\sigma}^{\mu\nu} h_{\mu\nu}^\rho - \alpha \Gamma_{\sigma\mu\nu}^{\rho\alpha\beta} h^{\mu\nu} h_{\alpha\beta,\rho}$$



$$\Gamma_{\rho\mu\nu} = \Gamma_{\rho\mu\nu}^{\sigma\alpha\beta} g_{\alpha\beta,\sigma}$$

$$\Gamma_{\sigma\mu\nu}^{\rho\alpha\beta} = I_{\sigma\kappa}^{\alpha\beta} I_{\mu\nu}^{\rho\kappa} - \frac{1}{2} I_{\mu\nu}^{\alpha\beta} \delta_{\sigma}^{\rho}$$

$$G_{\sigma} G^{\sigma} \approx h_{\mu\nu}^{\rho} \mathcal{P}_{\rho\kappa}^{\mu\nu} \mathcal{P}_{\alpha\beta}^{\kappa\sigma} h_{,\sigma}^{\alpha\beta} - 2\alpha \mathcal{P}_{\gamma\delta}^{\rho\kappa} h_{,\rho}^{\gamma\delta} \Gamma_{\kappa\mu\nu}^{\sigma\alpha\beta} h^{\mu\nu} h_{\alpha\beta,\sigma}$$

$$\begin{aligned} (S_{\text{EH}})_{h^2} &= \int d^D x \frac{1}{4} (2\eta^{\sigma\gamma} \eta^{\rho\delta} \eta^{\alpha\beta} - \eta^{\gamma\delta} \eta^{\alpha\beta} \eta^{\rho\sigma} - 2\eta^{\sigma\alpha} \eta^{\gamma\rho} \eta^{\delta\beta} + \eta^{\rho\sigma} \eta^{\alpha\gamma} \eta^{\beta\delta}) h_{\alpha\beta,\rho} h_{\gamma\delta,\sigma} \\ &= \int d^D x \frac{1}{4} (2\eta^{\sigma\gamma} \eta^{\rho\delta} \eta^{\alpha\beta} - \eta^{\gamma\delta} \eta^{\alpha\beta} \eta^{\rho\sigma} - 2\eta^{\sigma\alpha} \eta^{\gamma\rho} \eta^{\delta\beta} + \eta^{\rho\sigma} \eta^{\alpha\gamma} \eta^{\beta\delta}) h_{\alpha\beta,\sigma} h_{\gamma\delta,\rho} \end{aligned}$$

$$2\mathcal{P}_{\kappa\rho}^{\alpha\beta} \mathcal{P}_{\gamma\delta}^{\kappa\sigma} h_{\alpha\beta}^{\rho} h_{,\sigma}^{\gamma\delta} = \left(2\eta^{\sigma\alpha} \eta^{\gamma\rho} \eta^{\delta\beta} - 2\eta^{\sigma\gamma} \eta^{\rho\delta} \eta^{\alpha\beta} + \frac{1}{2} \eta^{\gamma\delta} \eta^{\alpha\beta} \eta^{\rho\sigma} \right) h_{\alpha\beta,\sigma} h_{\gamma\delta,\rho}$$

$$(S_{\text{EH}})_{h^2} = \frac{1}{4} \int d^D x h_{\gamma\delta}^{\rho} \left(\delta_{\sigma}^{\rho} \mathcal{P}_{\alpha\beta}^{\gamma\delta} - 2\mathcal{P}_{\rho\kappa}^{\gamma\delta} \mathcal{P}_{\alpha\beta}^{\sigma\kappa} \right) h_{,\sigma}^{\alpha\beta}$$

$$(S_{\text{EH}})_{h^3} = \frac{1}{2} \int d^D x U_{(c)}^{\mu\nu} \alpha\beta\rho\gamma\delta\sigma h_{\mu\nu} h_{\alpha\beta,\rho} h_{\gamma\delta,\sigma}$$

$$\begin{aligned} U_{(c)}^{\mu\nu} \alpha\beta\rho\gamma\delta\sigma h_{\alpha\beta,\rho} h_{\gamma\delta,\sigma} &= \\ 2I_{\phi\epsilon}^{\mu\nu} \mathcal{P}_{\rho\sigma}^{\alpha\beta} \mathcal{P}_{\gamma\delta}^{\sigma\phi} h_{\alpha\beta}^{\epsilon} h^{\gamma\delta,\rho} - \mathcal{P}_{\alpha\beta}^{\mu\rho} \mathcal{P}_{\gamma\delta}^{\nu\sigma} \eta_{\rho\sigma} h_{,\kappa}^{\alpha\beta} h^{\gamma\delta,\kappa} \\ &+ \mathcal{P}_{\rho\sigma}^{\mu\nu} \left(h_{,\beta}^{\rho\alpha} h_{,\alpha}^{\sigma\beta} - \frac{1}{2} h_{\beta}^{\alpha,\rho} h_{\alpha}^{\beta,\sigma} - h_{,\alpha}^{\rho\sigma} h_{,\beta}^{\alpha\beta} \right) \end{aligned}$$

$$U_{(c)}^{\mu\nu} \alpha\beta\rho\gamma\delta\sigma h_{\mu\nu} h_{\alpha\beta,\rho} h_{\gamma\delta,\sigma} = \frac{1}{2} h_{\nu}^{\mu} h_{,\mu} h^{\nu} - \frac{1}{4} h h_{,\rho} h^{\rho} + h_{\nu}^{\mu} h_{\mu,\rho}^{\nu} h^{\rho} - h_{\nu}^{\mu} h_{\mu}^{\nu,\sigma} h_{\sigma,\rho}^{\rho} + \frac{1}{4} h h_{\nu,\rho}^{\mu} h_{\mu}^{\nu,\rho}$$

$$\begin{aligned} &- h_{\mu}^{\nu} h_{\sigma,\nu}^{\mu} h^{\sigma} - h_{\nu}^{\mu} h_{,\mu}^{\nu} h_{\mu,\rho}^{\rho} + \frac{1}{2} h h_{\sigma,\rho}^{\rho} h^{\sigma} - h_{\nu}^{\mu} h_{\mu,\sigma}^{\rho} h_{\rho}^{\nu,\sigma} - \frac{1}{2} h h_{\nu,\mu}^{\rho} h_{\rho}^{\mu,\nu} \\ &+ h^{\mu\nu} h_{\mu,\rho}^{\sigma} h_{\nu,\sigma}^{\rho} - \frac{1}{2} h_{\nu}^{\mu} h_{\sigma,\mu}^{\rho} h_{\rho}^{\sigma,\nu} + 2h_{\nu}^{\mu} h_{\rho}^{\sigma,\nu} h_{\mu,\sigma}^{\rho} \end{aligned}$$

$$(S_{(\text{gf})})_{h^3} = \frac{1}{2\xi} \int d^D x U_{(\text{gf})}^{\mu\nu} \alpha\beta\rho\gamma\delta\sigma h_{\mu\nu} h_{\alpha\beta,\rho} h_{\gamma\delta,\sigma}$$

$$\begin{aligned} U_{\text{gf}}^{\mu\nu\alpha\beta\rho\gamma\delta\sigma} h_{\alpha\beta,\rho} h_{\gamma\delta,\sigma} &= -2\alpha \mathcal{P}_{\rho\sigma}^{\alpha\beta} h_{\alpha\beta}^{\sigma} \beta_{\kappa\gamma\delta}^{\rho\mu\nu} h^{\gamma\delta,\kappa} \\ &= \alpha \mathcal{P}_{\alpha\beta}^{\rho\sigma} h_{,\sigma}^{\alpha\beta} (-h_{\rho}^{\mu,\nu} - h_{\rho}^{\nu,\mu} + h_{,\rho}^{\mu\nu}) \end{aligned}$$

$$U^{\mu\nu} \alpha\beta\rho\gamma\delta\sigma = U_{(c)}^{\mu\nu} \alpha\beta\rho\gamma\delta\sigma + \frac{1}{\xi} U_{(\text{gf})}^{\mu\nu} \alpha\beta\rho\gamma\delta\sigma,$$

$$S_{\text{EH}} + S_{(\text{gf})} \approx \frac{1}{4} \int d^D x h_{\mu\nu}^{\rho} \left(\delta_{\sigma}^{\rho} \mathcal{P}_{\alpha\beta}^{\mu\nu} - 2 \left(1 - \frac{1}{\xi} \right) \mathcal{P}_{\rho\kappa}^{\mu\nu} \mathcal{P}_{\alpha\beta}^{\sigma\kappa} \right) h_{,\sigma}^{\alpha\beta} + \frac{1}{2} \int d^D x U^{\mu\nu} \alpha\beta\rho\gamma\delta\sigma h_{\mu\nu} h_{\alpha\beta,\rho} h_{\gamma\delta,\sigma}$$

$$\begin{aligned} \alpha \mathcal{P}_{\alpha\beta}^{\rho\sigma} h_{,\sigma}^{\alpha\beta} h_{,\rho}^{\mu\nu} &= \alpha \mathcal{P}^{\alpha\beta\rho\sigma} \Gamma^{\gamma\delta\mu\nu} h_{\alpha\beta,\rho} h_{\mu\nu,\sigma} \\ &= \frac{1}{2} \alpha (\mathcal{P}^{\alpha\beta\rho\sigma} \Gamma^{\gamma\delta\mu\nu} + \mathcal{P}^{\gamma\delta\rho\sigma} \Gamma^{\alpha\beta\mu\nu}) h_{\alpha\beta,\rho} h_{\mu\nu,\sigma} \end{aligned}$$



$$\frac{1}{2}\alpha(\mathcal{P}^{\alpha\beta\rho\sigma}I\gamma\delta\mu\nu + \mathcal{P}\gamma\delta\rho\sigma I\alpha\beta\mu\nu)$$

$$U_{(\text{gf})}^{\mu\nu\ \alpha\beta\rho\gamma\delta\sigma} = -\alpha\left(I^{\mu\nu\rho\kappa}I_{\kappa\lambda}^{\alpha\beta}\mathcal{P}^{\lambda\sigma\gamma\delta} + I^{\mu\nu\sigma\kappa}I_{\kappa\lambda}^{\gamma\delta}\mathcal{P}^{\lambda\rho\alpha\beta}\right) + \frac{1}{2}\alpha(\mathcal{P}^{\alpha\beta\rho\sigma}I\gamma\delta\mu\nu + \mathcal{P}\gamma\delta\rho\sigma I\alpha\beta\mu\nu)$$

$$S_{\text{EH}} = -\int d^D x h_{\mu\nu} \sum_{n=1.. \infty} \frac{1}{(n+1)} \mathcal{G}_{h^n}^{\mu\nu}(h, h, \dots, h)$$

$$S_{(\text{gf})} = -\frac{1}{\xi} \int d^D x h_{\mu\nu} \sum_{n=1.. \infty} \frac{1}{(n+1)} \mathcal{H}_{h^n}^{\mu\nu}(h, h, \dots, h)$$

$$\mathcal{G}_{h^n}^{\mu\nu}(h, h, \dots, h)$$

$$\int d^D x h_{\mu\nu} \frac{1}{(n+1)} \mathcal{G}_{h^n}^{\mu\nu}(h, h, \dots, h)$$

$$\delta \int d^D x h_{\mu\nu} \frac{1}{(n+1)} \mathcal{G}_{h^n}^{\mu\nu}(h, h, \dots, h) = \int d^D x \mathcal{G}_{h^n}^{\mu\nu}(h, h, \dots, h) \delta h_{\mu\nu}$$

$$(S_{\text{EH}})_{h^3} = -\frac{1}{3} \int d^D x h_{\mu\nu} \mathcal{G}_{h^2}^{\mu\nu}(h, h)$$

$$= \frac{1}{2} \int d^D x h_{\mu\nu} U_{(\text{c})}^{\mu\nu\ \alpha\beta\rho\gamma\delta\sigma} h_{\alpha\beta, \rho} h_{\gamma\delta, \sigma}$$

$$= \frac{1}{6} \int d^D x \left(h_{\mu\nu} U_{(\text{c})}^{\mu\nu\ \alpha\beta\rho\gamma\delta\sigma} h_{\alpha\beta, \rho} h_{\gamma\delta, \sigma} + h_{\mu\nu, \sigma} U_{(\text{c})}^{\alpha\beta\gamma\delta\rho\mu\nu\sigma} h_{\alpha\beta} h_{\gamma\delta, \rho} + h_{\mu\nu, \rho} U_{(\text{c})}^{\gamma\delta\mu\nu\rho\ \alpha\beta\sigma} h_{\alpha\beta, \sigma} h_{\gamma\delta} \right)$$

$$= \frac{1}{6} \int d^D x h_{\mu\nu} \left(U_{(\text{c})}^{\mu\nu\ \alpha\beta\rho\gamma\delta\sigma} h_{\alpha\beta, \rho} h_{\gamma\delta, \sigma} - U_{(\text{c})}^{\alpha\beta\ \gamma\delta\rho\ \mu\nu\sigma} \partial_\sigma (h_{\alpha\beta} h_{\gamma\delta, \rho}) - U_{(\text{c})}^{\gamma\delta\mu\nu\alpha\beta\sigma} \partial_\rho (h_{\alpha\beta, \sigma} h_{\gamma\delta}) \right)$$

$$\mathcal{G}_{h^2}^{\mu\nu}(h, h) = -\frac{1}{2} \left(U_{(\text{c})}^{\mu\nu\ \alpha\beta\rho\gamma\delta\sigma} h_{\alpha\beta, \rho} h_{\gamma\delta, \sigma} - 2U_{(\text{c})}^{\alpha\beta\ \gamma\delta\rho\ \mu\nu\sigma} h_{\alpha\beta, \sigma} h_{\gamma\delta, \rho} - 2U_{(\text{c})}^{\alpha\beta\ \gamma\delta\rho\ \mu\nu\sigma} h_{\alpha\beta} h_{\gamma\delta, \rho\sigma} \right)$$

$$\delta S_{\text{EH}} = -\int d^D x \delta h_{\mu\nu} \sum_{n=1.. \infty} \mathcal{G}_{h^n}^{\mu\nu}(h, h, \dots, h)$$

$$\delta S_{\text{EH}} = -\int d^D x \sqrt{-g} G^{\mu\nu} \delta h_{\mu\nu}$$

$$\sum_{n=1.. \infty} \mathcal{G}_{h^n}^{\mu\nu}(h, h, \dots, h) = \sqrt{-g} G^{\mu\nu}$$

$$\mathcal{G}_{h^n}^{\mu\nu} = (\sqrt{-g} G^{\mu\nu})_{h^n}$$

$$= \sum_{j=0..n} G_{h^{n-j}}^{\mu\nu} (\sqrt{-g})_{h^j}$$

$$\mathcal{G}_{h^1}^{\mu\nu} = G_{h^1}^{\mu\nu}$$

$$\mathcal{G}_{h^2}^{\mu\nu} = G_{h^2}^{\mu\nu} + \frac{1}{2} h G_{h^1}^{\mu\nu}$$

$$\mathcal{G}_{h^3}^{\mu\nu} = G_{h^3}^{\mu\nu} + \frac{1}{2} h G_{h^2}^{\mu\nu} - \frac{1}{4} \mathcal{P}^{\rho\sigma} h^{\alpha\beta} h_{\rho\sigma} G_{h^1}^{\mu\nu}$$

$$\mathcal{G}^{\mu\nu} = \sqrt{-g}G^{\mu\nu}$$

$$\delta S_{(\text{gf})} = \frac{1}{\xi} \int d^D x \eta^{\rho\sigma} G_\rho \delta G_\sigma = -\frac{1}{\xi} \int d^D x \sqrt{-g} H^{\mu\nu} \delta h_{\mu\nu}$$

$$\mathcal{H}_{h^n}^{\mu\nu} = (\sqrt{-g}H^{\mu\nu})_{h^n}$$

$$\mathcal{H}^{\mu\nu} = \sqrt{-g}H^{\mu\nu}$$

$$\begin{aligned} G^{\mu\nu} &= \frac{1}{2}(g^{\mu\alpha}g^{\nu\beta} + g^{\mu\beta}g^{\nu\alpha} - g^{\mu\nu}g^{\alpha\beta})R_{\alpha\beta} \\ &= \frac{1}{4}(g^{\mu\alpha}g^{\nu\beta} + g^{\mu\beta}g^{\nu\alpha} - g^{\mu\nu}g^{\alpha\beta})g^{\gamma\delta}F_{\alpha\beta}^{\rho\sigma}{}_{\gamma\delta}(g_{\phi\epsilon,\rho\sigma} + g^{\kappa\lambda}\Gamma_{\kappa\rho\sigma}\Gamma_{\lambda\phi\epsilon}) \end{aligned}$$

$$F_{\mu\nu\rho\sigma}^{\alpha\beta\gamma\delta} = I_{\mu\nu}^{\alpha\beta}I_{\rho\sigma}^{\gamma\delta} + I_{\rho\sigma}^{\alpha\beta}I_{\mu\nu}^{\gamma\delta} - 2I_{\epsilon\zeta}^{\alpha\beta}I_{\mu\nu}^{\zeta\eta}I_{\eta\theta}^{\gamma\delta}I_{\rho\sigma}^{\theta\epsilon}.$$

$$G_a^{\mu\nu} = \frac{1}{4}(g^{\mu\alpha}g^{\nu\beta} + g^{\mu\beta}g^{\nu\alpha} - g^{\mu\nu}g^{\alpha\beta})g^{\gamma\delta}F_{\alpha\beta}^{\rho\sigma}{}_{\gamma\delta}\frac{\phi^\epsilon}{\gamma\delta}\partial_\rho\partial_\sigma g_{\phi\epsilon},$$

$$G_b^{\mu\nu} = \frac{1}{4}(g^{\mu\alpha}g^{\nu\beta} + g^{\mu\beta}g^{\nu\alpha} - g^{\mu\nu}g^{\alpha\beta})g^{\gamma\delta}F_{\alpha\beta}^{\rho\sigma}{}_{\gamma\delta}g^{\kappa\lambda}\Gamma_{\kappa\rho\sigma}\Gamma_{\lambda\phi\epsilon}.$$

$$G_h^{\mu\nu} = (G_a^{\mu\nu})_h = \frac{1}{2}\mathcal{P}^{\mu\nu\alpha\beta}\eta^{\gamma\delta}F_{\alpha\beta\gamma\delta}^{\rho\sigma\phi\epsilon}h_{\rho\sigma,\phi\epsilon} = \frac{1}{2}Q^{\mu\nu\alpha\beta\gamma\delta}h_{\alpha\beta,\gamma\delta}$$

$$\begin{aligned} Q^{\mu\nu\alpha\beta\gamma\delta} &= \mathcal{P}^{\mu\nu\rho\sigma}\eta^{\phi\epsilon}F_{\rho\sigma}^{\alpha\beta\gamma\delta} \\ &= \eta^{\mu\nu}\mathcal{P}^{\alpha\beta\gamma\delta} - 2I_{\sigma\rho}^{\mu\nu}\mathcal{P}^{\rho\phi\alpha\beta}\eta_{\phi\epsilon}\mathcal{P}^{\epsilon\sigma\gamma\delta} \end{aligned}$$

$$(G_a^{\mu\nu})_{h^2} = -\left(\mathcal{P}_{\rho\kappa}^{\mu\nu}h_\sigma^\rho\mathcal{P}_{\alpha\beta}^{-1\sigma\kappa} + \frac{1}{2(D-2)}h^{\mu\nu}\eta_{\alpha\beta}\right)Q_{\gamma\delta}^{\alpha\beta}h_\epsilon h^{\gamma\delta,\phi\epsilon} - \frac{1}{2}F_{\gamma\delta\phi}^{\mu\nu}\rho_{\rho\sigma}h_{\rho\sigma}h^{\gamma\delta,\phi\epsilon}$$

$$\begin{aligned} (G_b^{\mu\nu})_{h^2} &= \frac{1}{2}Q^{\mu\nu\alpha\beta\gamma\delta} \\ &\quad \eta^{\rho\sigma}\Gamma_{\rho\alpha\beta}\Gamma_{\sigma\gamma\delta} \end{aligned}$$

$$(S_{\text{EH}} + S_{(\text{gf})})_{h^2} = \frac{1}{4} \int d^D x h_{\mu\nu}^\rho \left(\delta_\sigma^\rho \mathcal{P}_{\alpha\beta}^{\mu\nu} - 2 \left(1 - \frac{1}{\xi} \right) \mathcal{P}_{\rho\kappa}^{\mu\nu} \mathcal{P}_{\alpha\beta}^{\sigma\kappa} \right) h_{,\sigma}^{\alpha\beta}$$

$$\frac{2}{\kappa^2} (S_{\text{EH}} + S_{(\text{gf})})_{h^2} = \frac{1}{2} \int d^D x h_{\mu\nu}^{\rho\rho} \left(\delta_\sigma^\rho \mathcal{P}_{\alpha\beta}^{\mu\nu} - 2 \left(1 - \frac{1}{\xi} \right) \mathcal{P}_{\rho\kappa}^{\mu\nu} \mathcal{P}_{\alpha\beta}^{\sigma\kappa} \right) h_{,\sigma}^{\alpha\beta}$$

$$\frac{2}{\kappa^2} (S_{\text{EH}} + S_{(\text{gf})})_{h^2} = \frac{1}{2} \int \frac{d^D p}{(2\pi)^D} \tilde{h}_{\mu\nu}^\dagger p^2 \left(\mathcal{P}_{\alpha\beta}^{\mu\nu} - 2 \left(1 - \frac{1}{\xi} \right) \mathcal{P}_{\rho\kappa}^{\mu\nu} \frac{p^\rho p_\sigma}{p^2} \mathcal{P}_{\alpha\beta}^{\sigma\kappa} \right) \tilde{h}^{\alpha\beta}$$

$$\Delta_{\alpha\beta}^{\mu\nu} = \mathcal{P}_{\alpha\beta}^{\mu\nu} - 2 \left(1 - \frac{1}{\xi} \right) \mathcal{P}_{\rho\kappa}^{\mu\nu} \frac{p^\rho p_\sigma}{p^2} \mathcal{P}_{\alpha\beta}^{\sigma\kappa}$$



$$\begin{aligned}
I_{\alpha\beta}^{\mu\nu} &= \frac{1}{2}(\delta_{\nu}^{\mu}\delta_{\beta}^{\alpha} + \delta_{\beta}^{\mu}\delta_{\alpha}^{\nu}) \\
T_{\alpha\beta}^{\mu\nu} &= \frac{1}{4}\eta^{\mu\nu}\eta_{\alpha\beta} \\
C_{\alpha\beta}^{\mu\nu} &= \frac{1}{2}\left(\eta^{\mu\nu}\frac{p_{\alpha}p_{\beta}}{p^2} + \frac{p^{\mu}p^{\nu}}{p^2}\eta_{\alpha\beta}\right) \\
J_{\alpha\beta}^{\mu\nu} &= I_{\rho\kappa}^{\mu\nu}\frac{p_{\sigma}p^{\rho}}{p^2}I_{\alpha\beta}^{\sigma\kappa} \\
K_{\alpha\beta}^{\mu\nu} &= \frac{p^{\mu}p^{\nu}}{p^2}\frac{p_{\alpha}p_{\beta}}{p^2}
\end{aligned}$$

$$\mathcal{P} = I - 2T$$

$$\Delta = \mathcal{P} + 2\left(1 - \frac{1}{\xi}\right)\mathcal{P}\mathcal{J}\mathcal{P}$$

$$G = \alpha_1 I + \alpha_2 T + \alpha_3 C + \alpha_4 J + \alpha_5 K$$

$$\Delta G = I$$

$$A_{\alpha\beta}^{\mu\nu} = \frac{1}{2}\left(\eta^{\mu\nu}\frac{p_{\alpha}p_{\beta}}{p^2} - \frac{p^{\mu}p^{\nu}}{p^2}\eta_{\alpha\beta}\right)$$

$$TT = \frac{D}{4}T$$

$$JJ = \frac{1}{2}J + \frac{1}{2}K$$

$$JJJ = \frac{1}{4}K$$

$$T_{\alpha\beta}^{\mu\nu}C_{\gamma\delta}^{\alpha\beta} = \frac{1}{2}T_{\gamma\delta}^{\mu\nu} + \frac{D}{8}C_{\gamma\delta}^{\mu\nu} + \frac{D}{8}A_{\gamma\delta}^{\mu\nu}$$

$$\begin{aligned}
\Delta_{\alpha\beta}^{\mu\nu}G_{\gamma\delta}^{\alpha\beta} &= \alpha_1 I_{\gamma\delta}^{\mu\nu} \\
&+ \left(\alpha_1\left(-4 + \frac{2}{\xi}\right) + \alpha_2\left(-1 + \frac{1}{2\xi}\right)(D-2) - \alpha_3\frac{1}{\xi}\right)T_{\gamma\delta}^{\mu\nu} \\
&+ \left(\alpha_1\left(2 - 2\frac{1}{\xi}\right) + \alpha_2\frac{D-2}{4}\left(1 - \frac{1}{\xi}\right) + \alpha_3\left(-\frac{D-2}{2} + \frac{D}{4\xi}\right) - (\alpha_4 + \alpha_5)\frac{1}{2\xi}\right)C_{\gamma\delta}^{\mu\nu} \\
&+ \left(\alpha_1 2\left(-1 + \frac{1}{\xi}\right) + \alpha_4\frac{1}{\xi}\right)J_{\gamma\delta}^{\mu\nu} \\
&+ \left(\alpha_3(D-2)\left(\frac{1}{2} - \frac{1}{\xi}\right) + \alpha_5\frac{1}{\xi}\right)K_{\gamma\delta}^{\mu\nu} \\
&+ \left(\alpha_2\frac{D-2}{4}\left(-1 + \frac{1}{\xi}\right) + \alpha_3\left(-\frac{D-2}{2} + \frac{D-4}{4\xi}\right) - (\alpha_4 + \alpha_5)\frac{1}{2\xi}\right)A_{\gamma\delta}^{\mu\nu} \\
&2\left(-1 + \frac{1}{\xi}\right) + \alpha_4\frac{1}{\xi} = 0
\end{aligned}$$

$$\begin{aligned}\alpha_1 &= 1 \\ \alpha_2 &= -\frac{4}{D-2} \\ \alpha_3 &= 0 \\ \alpha_4 &= -2(1-\xi) \\ \alpha_5 &= 0\end{aligned}$$

$$G = \mathcal{P}^{-1} - 2(1-\xi)J$$

$$\mathcal{P}^{-1} = I - \frac{4}{D-2}T$$

$$\Delta = \mathcal{P} - 2\left(1 - \frac{1}{\xi}\right)\mathcal{P}J\mathcal{P}$$

$$G = \mathcal{P}^{-1} - 2(1-\xi)J$$

$$\Delta = \Delta_{(c)} + \frac{1}{\xi}\Delta_{(gf)}$$

$$\Delta_{(c)} = \mathcal{P} - 2\mathcal{P}J\mathcal{P}$$

$$\Delta_{(gf)} = 2\mathcal{P}J\mathcal{P}$$

$$G = G_{(c)} + \xi G_{(gf)}$$

$$G_{(c)} = \mathcal{P}^{-1} - 2J$$

$$G_{(gf)} = 2J$$

$$\Delta_{(c)}^{\mu\nu}{}_{\alpha\beta} = Q_{\alpha\beta}^{\mu\nu} \frac{\sigma\rho}{p_\sigma p_\rho}$$

$$p_\mu \Delta_{(c)}^{\mu\nu}{}_{\alpha\beta} = p_\mu Q_{\alpha\beta}^{\mu\nu} \frac{\sigma\rho}{p^2} = 0$$

$$p_\mu \Delta_{(c)}^{\mu\nu}{}_{\alpha\beta} = p_\mu \mathcal{P}_{\alpha\beta}^{\mu\nu} - 2p_\mu \mathcal{P}_{\rho\kappa}^{\mu\nu} \frac{p^\rho p_\sigma}{p^2} \mathcal{P}_{\alpha\beta}^{\sigma\kappa}$$

$$2p_\mu \mathcal{P}_{\rho\kappa}^{\mu\nu} \frac{p^\rho p_\sigma}{p^2} \mathcal{P}_{\alpha\beta}^{\sigma\kappa} = p_\mu \mathcal{P}_{\alpha\beta}^{\mu\nu}$$

$$\Delta G = \Delta_{(c)} G_{(c)} + \Delta_{(gf)} G_{(gf)} + \xi \Delta_{(c)} G_{(gf)} + \frac{1}{\xi} \Delta_{(gf)} G_{(c)}$$

$$J\mathcal{P}J = J^2 - 2JTJ = \frac{1}{2}J$$

$$\begin{aligned}\Delta_{(gf)} G_{(c)} &= 2\mathcal{P}J\mathcal{P}(\mathcal{P}^{-1} - 2J) \\ &= 2\mathcal{P}(J - 2J\mathcal{P}J) = 0\end{aligned}$$

$$\begin{aligned}\Delta_{(c)} G_{(c)} + \Delta_{(gf)} G_{(gf)} &= (\mathcal{P} - 2\mathcal{P}J\mathcal{P})(\mathcal{P}^{-1} - 2J) + 4\mathcal{P}J\mathcal{P}J \\ &= I - 4\mathcal{P}J + 8\mathcal{P}J\mathcal{P}J = I\end{aligned}$$



$$\begin{array}{c} \text{wavy line} \\ \mu\nu \end{array} \begin{array}{c} \text{arrow} \\ p \end{array} \quad \alpha\beta = \frac{i}{p^2 + i\epsilon} G_{\alpha\beta}^{\mu\nu}$$

$$G_{\alpha\beta}^{\mu\nu} = I_{\alpha\beta}^{\mu\nu} - \frac{1}{D-2} \eta^{\mu\nu} \eta_{\alpha\beta} - 2(1-\xi) I_{\rho\kappa}^{\mu\nu} \frac{p^\rho p_\sigma}{p^2} I_{\alpha\beta}^{\kappa\sigma}$$

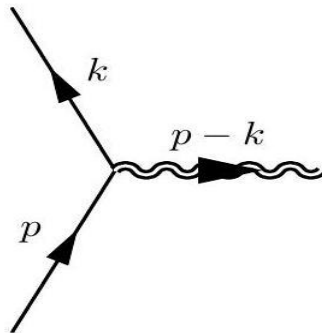
$$S_\phi = \frac{1}{2} \int d^D x \sqrt{-g} (g^{\mu\nu} \phi_{,\mu} \phi_{,\nu} - m^2 \phi^2)$$

$$(S_\phi)_h = -\frac{1}{2} \int d^D x h^{\mu\nu} \phi_{,\mu} \phi_{,\nu} + \frac{1}{4} \int d^D x h_\mu^\mu (\phi^{,\nu} \phi_{,\nu} - m^2 \phi^2)$$

$$\begin{aligned} S_\phi &\approx \frac{1}{2} \int d^D x (\eta^{\mu\nu} \phi_{,\mu} \phi_{,\nu} - m^2 \phi^2) \\ &\quad - \frac{1}{2} \int d^D x \left(h^{\mu\nu} \phi_{,\mu} \phi_{,\nu} - \frac{1}{2} h_\mu^\mu (\phi^{,\nu} \phi_{,\nu} - m^2 \phi^2) \right) \\ &\quad + \frac{1}{2} \int d^D x \left(h_{\mu\nu} h^{\alpha\gamma} \mathcal{P}_{\alpha\beta}^{\mu\nu} \phi_{,\gamma} \phi^{,\beta} - \frac{1}{4} h_{\mu\nu} h^{\alpha\beta} \mathcal{P}_{\alpha\beta}^{\mu\nu} (\phi^{,\nu} \phi_{,\nu} - m^2 \phi^2) \right) \end{aligned}$$

$$V_{\phi^2 h}^{\mu\nu}(p, k, m) = I_{\alpha\beta}^{\mu\nu} p^\alpha k^\beta - \frac{pk - m^2}{2} \eta^{\mu\nu}$$

$$\begin{aligned} V_{\phi^2 h^2}^{\mu\beta}(p, k, m) &= \left(I^{\mu\nu} \frac{\eta}{\lambda} I^{\alpha\beta} \frac{\sigma\rho\kappa}{\kappa} \eta - \frac{1}{4} (\eta^{\mu\nu} I^{\alpha\beta\sigma\rho} + \eta^{\alpha\beta} I^{\mu\nu\sigma\rho}) \right) p_\sigma k_\rho - \frac{pk - m^2}{4} \mathcal{P}^{\mu\nu\alpha\beta} \\ &\quad \frac{i}{p^2 - m^2 + i\epsilon} \end{aligned}$$



$$\mu\nu = -i\kappa V_{\phi^2 h}^{\mu\nu}(p, k, m)$$

$$= 2i\kappa^2 V_{\phi^2 h^2}^{\mu\nu}(p, k, m)$$

$$S_{h^3} = \frac{1}{\kappa^2} \int d^D x h_{\mu\nu} U^{\mu\nu} \alpha\beta\rho\gamma\delta\sigma h_{\alpha\beta,\rho} h_{\gamma\delta,\sigma}$$

$$\kappa_{\mu\nu} = \int \frac{d^D l}{(2\pi)^D} e^{-ilx} \tilde{h}_{\mu\nu}(l)$$

$$S_{h^3} = -\kappa \int \frac{d^D l_{(1)}}{(2\pi)^D} \frac{d^D l_{(2)}}{(2\pi)^D} \frac{d^D l_{(3)}}{(2\pi)^D} (2\pi)^D \delta^D(l_{(1)} + l_{(2)} + l_{(3)}) \tilde{f}_{\mu\nu}^{(1)} U^{\mu\nu} \alpha\beta\rho\gamma\delta\sigma \tilde{h}_{\alpha\beta}^{(2)} \tilde{h}_{\gamma\delta}^{(3)} l_{(2)\rho} l_{(3)\sigma}$$

$$S_{h^3} = -\frac{\kappa}{3} \int \frac{d^D l_{(1)}}{(2\pi)^D} \frac{d^D l_{(2)}}{(2\pi)^D} \frac{d^D l_{(3)}}{(2\pi)^D} (2\pi)^D \delta^D(l_{(1)} + l_{(2)} + l_{(3)}) (U^{\mu\nu} \alpha\beta\rho\gamma\delta\sigma l_{(2)\rho} l_{(3)\sigma} + U^{\alpha\beta} \gamma\delta\rho\mu\nu\sigma l_{(3)\rho} l_{(1)\sigma} + U^{\gamma\delta} \mu\nu\rho\alpha\beta\sigma l_{(1)\rho} l_{(2)\sigma}) \tilde{h}_{\mu\nu}^{(1)} \tilde{h}_{\alpha\beta}^{(2)} \tilde{h}_{\gamma\delta}^{(3)}$$

$$\gamma\delta = 2i\kappa V_{h^3}^{\mu\nu} \alpha\beta \gamma\delta(l_1, l_2, l_3)$$

$$V_{h^3}^{\alpha\beta\gamma\delta\sigma}(l_{(1)}, l_{(2)}, l_{(3)}) = -(U^{\mu\nu} \alpha\beta\rho\gamma\delta\sigma l_{(2)\rho} l_{(3)\sigma} + U^{\alpha\beta} \gamma\delta\rho\mu\nu\sigma l_{(3)\rho} l_{(1)\sigma} + U^{\gamma\delta} \mu\nu\rho\alpha\beta\sigma l_{(1)\rho} l_{(2)\sigma})$$

$$S_{\text{EH}} + S_{(\text{gf})} = - \int d^D x h_{\mu\nu} \sum_{n=1.. \infty} \frac{1}{(n+1)} \left(\mathcal{G}_{h^n}^{\mu\nu}(h, h, \dots, h) + \frac{1}{\xi} \mathcal{H}_{h^n}^{\mu\nu}(h, h, \dots, h) \right)$$

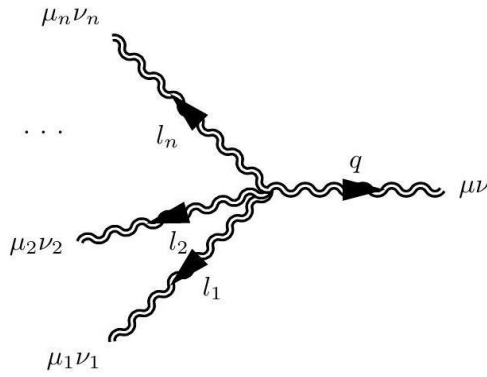
$$\mathcal{G}_{h^3}^{\mu\nu}(h_{\mu\nu}, h_{\mu\nu}, h_{\mu\nu}) = \mathcal{G}_{h^3}^{\mu\nu} \left(\int \frac{d^D l_{(1)}}{(2\pi)^D} e^{-ixl_{(1)}} \tilde{h}_{\mu\nu}^{(1)}, \int \frac{d^D l_{(2)}}{(2\pi)^D} e^{-ixl_{(2)}} \tilde{h}_{\mu\nu}^{(2)}, \int \frac{d^D l_{(3)}}{(2\pi)^D} e^{-ixl_{(3)}} \tilde{h}_{\mu\nu}^{(3)} \right)$$

$$\mathcal{G}_{h^3}^{\mu\nu}(h_{\mu\nu}, h_{\mu\nu}, h_{\mu\nu}) = \int \frac{d^D l_{(1)}}{(2\pi)^D} \frac{d^D l_{(2)}}{(2\pi)^D} \frac{d^D l_{(3)}}{(2\pi)^D} \mathcal{G}_{h^3}^{\mu\nu} \left(e^{-ixl_{(1)}} \tilde{h}_{\mu\nu}^{(1)}, e^{-ixl_{(2)}} \tilde{h}_{\mu\nu}^{(2)}, e^{-ixl_{(3)}} \tilde{h}_{\mu\nu}^{(3)} \right)$$

$$\mathcal{G}_{\tilde{h}^3}^{\mu\nu}(h_{\mu\nu}, h_{\mu\nu}, h_{\mu\nu}) = \int \frac{d^D l_{(1)}}{(2\pi)^D} \frac{d^D l_{(2)}}{(2\pi)^D} \frac{d^D l_{(3)}}{(2\pi)^D} e^{-ix(l_{(1)}+l_{(2)}+l_{(3)})} \tilde{\mathcal{G}}_{\tilde{h}^3}^{\mu\nu}(\tilde{h}_{\mu\nu}^{(1)}, \tilde{h}_{\mu\nu}^{(2)}, \tilde{h}_{\mu\nu}^{(3)})$$

$$\int d^D x e^{iqx} \mathcal{G}_{\tilde{h}^3}^{\mu\nu}(h_{\mu\nu}, h_{\mu\nu}, h_{\mu\nu}) = \int \frac{d^D l_{(1)}}{(2\pi)^D} \frac{d^D l_{(2)}}{(2\pi)^D} \frac{d^D l_{(3)}}{(2\pi)^D} (2\pi)^D \delta^D(l_{(1)} + l_{(2)} + l_{(3)}) \tilde{\mathcal{G}}_{\tilde{h}^3}^{\mu\nu}(\tilde{h}_{\mu\nu}^{(1)}, \tilde{h}_{\mu\nu}^{(2)}, \tilde{h}_{\mu\nu}^{(3)})$$

$$\sum_{n=1.. \infty} -\frac{2\kappa^{n-1}}{n+1} \tilde{\mathcal{G}}_{\mu\nu}^{(0)}\left(\tilde{\mathcal{G}}_{\mu\nu}^{\mu\nu}(\tilde{\mathcal{G}}_{\mu\nu}^{(1)}, \dots, \tilde{\mathcal{G}}_{\mu\nu}^{(n)}) + \frac{1}{\xi} \tilde{\mathcal{H}}_{\tilde{h}^n}^{\mu\nu}(\tilde{\mathcal{G}}_{\mu\nu}^{(1)}, \dots, \tilde{\mathcal{G}}_{\mu\nu}^{(n)})\right)$$



$$= in! \kappa^{n-1} V_{\tilde{h}^{n+1}}^{\mu\nu\mu_1\nu_1\mu_2\nu_2\dots\mu_n\nu_n}(q, l_1, l_2, \dots, l_n)$$

$$V_{\tilde{h}^{n+1}}^{\mu\nu\mu_1\nu_1\mu_2\nu_2\dots\mu_n\nu_n}(q, l_{(1)}, l_{(2)}, \dots, l_{(n)}) \tilde{h}_{\mu_1\nu_1}^{(1)} \tilde{h}_{\mu_2\nu_2}^{(2)} \dots \tilde{h}_{\mu_n\nu_n}^{(n)} = -2 \left(\tilde{\mathcal{G}}_{\tilde{h}^n}^{\mu\nu}(\tilde{h}_{\mu_1\nu_1}^{(1)}, \tilde{h}_{\mu_2\nu_2}^{(2)}, \dots, \tilde{h}_{\mu_n\nu_n}^{(n)}) + \frac{1}{\xi} \tilde{\mathcal{H}}_{\tilde{h}^n}^{\mu\nu}(\tilde{h}_{\mu_1\nu_1}^{(1)}, \tilde{h}_{\mu_2\nu_2}^{(2)}, \dots, \tilde{h}_{\mu_n\nu_n}^{(n)}) \right)$$

$$S = \frac{2}{\kappa^2} (S_{EH} + S_{gf}) + \int d^D x \sqrt{-g} \mathcal{L}_\phi$$

$$= \int d^D x \sqrt{-g} \left(\frac{2R}{\kappa^2} + \mathcal{L}_\phi \right) + \int d^D x \frac{1}{\kappa^2 \xi} \eta^{\sigma\rho} G_\sigma G_\rho$$

$$\delta S_{(gf)} = \frac{1}{\xi} \int d^D x \eta^{\rho\sigma} G_\rho \delta G_\sigma = -\frac{1}{\xi} \int d^D x \sqrt{-g} H^{\mu\nu} \delta h_{\mu\nu}$$

$$G_\sigma = (1 - \alpha) \partial_\mu \left(h_\sigma^\mu - \frac{1}{2} \eta_\sigma^\mu h_\nu^\nu \right) + \alpha g^{\mu\nu} \Gamma_{\sigma\mu\nu}$$

$$G_\sigma = (1 - \alpha) \eta^{\mu\nu} \Gamma_{\sigma\mu\nu}^{\rho\alpha\beta} h_{\alpha\beta,\rho} + \alpha g^{\mu\nu} \Gamma_{\sigma\mu\nu}^{\rho\alpha\beta} h_{\alpha\beta,\rho}$$

$$= (\eta^{\mu\nu} + \alpha (g^{\mu\nu} - \eta^{\mu\nu})) \Gamma_{\sigma\mu\nu}^{\rho\alpha\beta} h_{\alpha\beta,\rho}$$

$$= (\eta^{\mu\nu} + \alpha \hat{h}^{\mu\nu}) \Gamma_{\sigma\mu\nu}^{\rho\alpha\beta} h_{\alpha\beta,\rho}$$

$$\delta \hat{h}^{\mu\nu} = \delta (g^{\mu\nu}) = -g^{\mu\alpha} g^{\nu\beta} \delta g_{\alpha\beta}$$

$$\delta G_\sigma = (\eta^{\mu\nu} + \alpha \hat{h}^{\mu\nu}) \Gamma_{\sigma\mu\nu}^{\rho\alpha\beta} \delta h_{\alpha\beta,\rho} - \alpha \Gamma_{\sigma\mu\nu}^{\rho\alpha\beta} h_{\alpha\beta,\rho} g^{\mu\gamma} g^{\nu\delta} \delta h_{\gamma\delta}$$



$$\begin{aligned}
\delta S_{gf} &= \frac{1}{\xi} \int d^D x \eta^{\rho\sigma} G_\rho \delta G_\sigma \\
&= \frac{1}{\xi} \int d^D x G^\sigma \left((\eta^{\mu\nu} + \alpha \hat{h}^{\mu\nu}) \Gamma_{\sigma\mu\nu}^{\rho\gamma\delta} \delta h_{\gamma\delta,\rho} - \alpha \Gamma_{\sigma\mu\nu}^{\rho\alpha\beta} h_{\alpha\beta,\rho} g^{\mu\gamma} g^{\nu\delta} \delta h_{\gamma\delta} \right) \\
&= -\frac{1}{\xi} \int d^D x \left(\Gamma_{\sigma\mu\nu}^{\rho\gamma\delta} \partial_\rho \left(G^\sigma (\eta^{\mu\nu} + \alpha \hat{h}^{\mu\nu}) \right) + \alpha G^\sigma \Gamma_{\sigma\mu\nu}^{\rho\alpha\beta} h_{\alpha\beta,\rho} g^{\mu\gamma} g^{\nu\delta} \right) \delta h_{\gamma\delta} \\
\sqrt{-g} H^{\mu\nu} = \mathcal{H}^{\mu\nu} &= \alpha G^\rho \Gamma_{\rho\alpha\beta} g^{\alpha\mu} g^{\beta\nu} + \Gamma_{\rho\alpha\beta}^{\sigma\mu\nu} \partial_\sigma \left(G^\rho (\eta^{\alpha\beta} + \alpha \hat{h}^{\alpha\beta}) \right)
\end{aligned}$$

$$G^{\mu\nu} + \frac{1}{\xi} H^{\mu\nu} = -\frac{\kappa^2}{4} T^{\mu\nu}$$

$$G^{\mu\nu} = -\frac{\kappa^2}{4} T^{\mu\nu}$$

$$D_\mu H^{\mu\nu} = 0$$

$$\sqrt{-g} D_\mu H^{\mu\nu} = \partial_\mu (\mathcal{H}^{\mu\nu}) + \Gamma_{\alpha\beta}^\nu \mathcal{H}^{\alpha\beta}$$

$$\sqrt{-g} D_\mu H^{\mu\nu} \approx \partial_\mu \mathcal{H}^{\mu\nu}$$

$$H^{\mu\nu} \approx \Gamma_{\rho\alpha\beta}^{\sigma\mu\nu} \eta^{\alpha\beta} \partial_\sigma G^\rho$$

$$D_\mu H^{\mu\nu} \approx \Gamma_{\rho\alpha\beta}^{\sigma\mu\nu} \eta^{\alpha\beta} \partial_\sigma \partial_\mu G^\rho$$

$$D_\mu H^{\mu\nu} \approx \frac{1}{2} \partial_\sigma \partial^\sigma G^\nu$$

$$G^{\mu\nu} = -\frac{\kappa^2}{4} T^{\mu\nu}$$

$$G_\sigma = 0.$$

$$G^{\mu\nu} + \frac{1}{\xi} H^{\mu\nu} = -\frac{\kappa^2}{4} T^{\mu\nu}$$

$$\begin{aligned}
G^{\mu\nu} &= \sum_{n=1.. \infty} (G^{\mu\nu})_{h^n} \\
&= G_h^{\mu\nu} + G_{\text{non-linear}}^{\mu\nu}
\end{aligned}$$

$$G_h^{\mu\nu} + \frac{1}{\xi} H_h^{\mu\nu} = -\frac{\kappa^2}{4} T^{\mu\nu} - G_{\text{non-linear}}^{\mu\nu} - \frac{1}{\xi} H_{\text{non-linear}}^{\mu\nu}$$

$$\tau^{\mu\nu} = T^{\mu\nu} + \frac{4}{\kappa^2} G_{\text{non-linear}}^{\mu\nu}$$



$$G_h^{\mu\nu} = -\frac{\kappa^2}{4} \tau^{\mu\nu}$$

$$Q^{\mu\nu\alpha\beta\gamma\delta} h_{\alpha\beta,\gamma\delta} = -\frac{\kappa^2}{2} \tau^{\mu\nu}$$

$$0 = \partial_\mu Q^{\mu\nu\alpha\beta\gamma\delta} h_{\alpha\beta,\gamma\delta} = -\frac{\kappa^2}{2} \partial_\mu \tau^{\mu\nu}$$

$$H_h^{\mu\nu} = -H_{\text{non-linear}}^{\mu\nu}$$

$$H_h^{\mu\nu} = \mathcal{P}_{\sigma\kappa}^{\mu\nu} \mathcal{P}_{\gamma\delta}^{\rho\kappa} \partial^\sigma \partial_\rho h^{\gamma\delta}$$

$$G_{(c)}^{\mu\nu}{}_{\alpha\beta} H_h^{\alpha\beta} = 0$$

$$G_{(c)}^{\mu\nu}{}_{\alpha\beta} H_{\text{non-linear}}^{\alpha\beta} = 0$$

$$G_{(gf)}^{\alpha\beta}{}_{\mu\nu} \tau^{\alpha\beta} = 0$$

$$G_{(c)}^{\mu\nu}{}_{\alpha\beta} H_{\text{non-linear}}^{\alpha\beta} = 0$$

$$\mathcal{G}^{\mu\nu} + \frac{1}{\xi} \mathcal{H}^{\mu\nu} = -\frac{\kappa^2}{4} \sqrt{-g} T^{\mu\nu}$$

$$\mathcal{G}^{\mu\nu} = -\frac{\kappa^2}{4} \sqrt{-g} T^{\mu\nu}$$

$$\mathcal{H}^{\mu\nu} = 0$$

$$\mathcal{G}_h^{\mu\nu} = G_h^{\mu\nu}$$

$$\mathcal{H}_h^{\mu\nu} = H_h^{\mu\nu}$$

$$\mathcal{H}_h^{\mu\nu} = -\mathcal{H}_{\text{non-linear}}^{\mu\nu},$$

$$\mathcal{H}_{\text{non-linear}}^{\mu\nu} = H_{\text{non-linear}}^{\mu\nu}$$

$$\tau^{\mu\nu} = \sqrt{-g} T^{\mu\nu} + \frac{4}{\kappa^2} \mathcal{G}_{\text{non-linear}}^{\mu\nu}.$$

$$\left(\eta^{\rho\sigma} \mathcal{P}^{\mu\nu\alpha\beta} - 2 \left(1 - \frac{1}{\xi} \right) \mathcal{P}^{\mu\nu\rho\phi} \eta_{\phi\epsilon} \mathcal{P}^{\alpha\beta\sigma\epsilon} \right) h_{\alpha\beta,\rho\sigma} = -\frac{\kappa^2}{2} \tau^{\mu\nu} - 2 \frac{1}{\xi} H_{\text{non-linear}}^{\mu\nu}$$

$$q^2 \Delta^{\mu\nu\alpha\beta} \tilde{h}_{\alpha\beta} = \frac{\kappa^2}{2} \tilde{\tau}^{\mu\nu} + 2 \frac{1}{\xi} \tilde{H}_{\text{non-linear}}^{\mu\nu}$$

$$\tilde{h}_{\alpha\beta} = \frac{G_{\alpha\beta\mu\nu}}{q^2} \left(\frac{\kappa^2}{2} \tilde{\tau}^{\mu\nu} + \frac{2}{\xi} \tilde{H}_{\text{non-linear}}^{\mu\nu} \right)$$



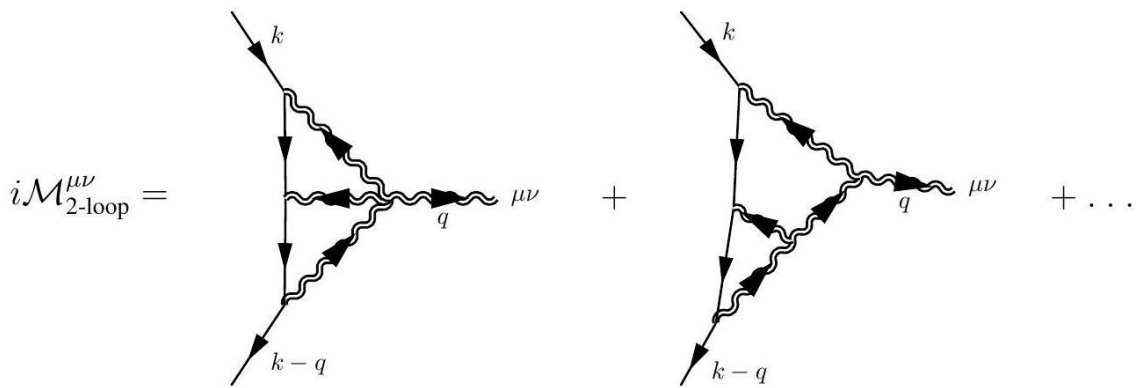
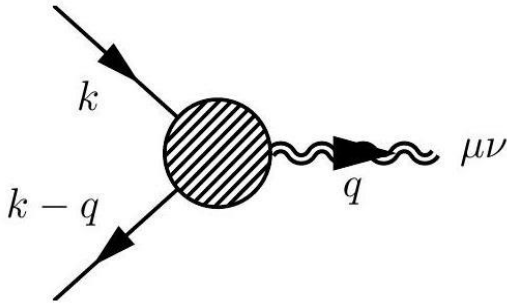
$$\begin{aligned}
\tilde{h}_{\alpha\beta} &= \frac{1}{q^2} \left(\frac{\kappa^2}{2} G_{(c_{\alpha\beta\mu\nu})} \tilde{\tau}^{\mu\nu} + 2G_{(\text{gf})} \alpha\beta\mu\nu \tilde{H}_{\text{non-linear}}^{\mu\nu} \right) \\
&= \frac{\mathcal{P}^{-1} \alpha\beta\mu\nu}{q^2} \left(\frac{\kappa^2}{2} \tilde{\tau}^{\mu\nu} + 2\tilde{H}_{\text{non-linear}}^{\mu\nu} \right) \\
\sum_{n=1.. \infty} \tilde{h}_{\alpha\beta}^{G^n} &= \frac{G_{\alpha\beta\mu\nu}}{q^2} \left(\frac{\kappa^2}{2} \sum_{n=0.. \infty} \tilde{\tau}_{G^n}^{\mu\nu} + 2\frac{1}{\xi} \sum_{n=2.. \infty} (\tilde{H}_{\text{non-linear}}^{\mu\nu})_{G^n} \right) \\
\tilde{h}_{\alpha\beta}^{G^1}(q) &= \frac{G_{\alpha\beta\mu\nu} \kappa^2}{q^2} \tilde{\tau}^{\mu\nu}(q) \\
\tilde{h}_{\alpha\beta}^{G^n} &= \frac{G_{\alpha\beta\mu\nu}}{q^2} \left(\frac{\kappa^2}{2} \tilde{\tau}_{G^{n-1}}^{\mu\nu} + \frac{2}{\xi} (\tilde{H}_{\text{non-linear}}^{\mu\nu})_{G^n} \right) \\
\tilde{h}_{\alpha\beta}^{G^n} &= 2\frac{G_{\alpha\beta\mu\nu}}{q^2} \left(\tilde{G}_{\text{non-linear}}^{\mu\nu} + \frac{1}{\xi} \tilde{H}_{\text{non-linear}}^{\mu\nu} \right)_{G^n} \\
\left(G^{\mu\nu} + \frac{1}{\xi} H^{\mu\nu} \right)_{G^n} &= 0 \\
h_{\mu\nu}^G &= \frac{\kappa^2}{2} \int \frac{d^D q}{(2\pi)^D} e^{-iqx} \frac{G_{\mu\nu\alpha\beta}}{q^2} \tilde{\tau}^{\alpha\beta} \\
\tilde{h}_{\alpha\beta}^{G^2} &= 2\frac{G_{\alpha\beta\mu\nu}}{q^2} \left(\tilde{G}_{\text{non-linear}}^{\mu\nu} + \frac{1}{\xi} \tilde{H}_{\text{non-linear}}^{\mu\nu} \right)_{G^2} \\
(G_{\text{non-linear}}^{\mu\nu})_{G^2} &= G_{h^2}^{\mu\nu}(h_{\mu\nu}^G, h_{\mu\nu}^G) \\
G_{h^2}^{\mu\nu}(h_{\mu\nu}^G, h_{\mu\nu}^G) &= G_{h^2}^{\mu\nu} \left(\int \frac{d^D l_1}{(2\pi)^D} e^{-il_1 x} \tilde{h}_{\mu\nu}^G(l_1), \int \frac{d^D l_2}{(2\pi)^D} e^{-il_2 x} \tilde{h}_{\mu\nu}^G(l_2) \right) \\
G_{h^2}^{\mu\nu}(h_{\mu\nu}^G, h_{\mu\nu}^G) &= \int \frac{d^D l_{(1)}}{(2\pi)^D} \frac{d^D l_{(2)}}{(2\pi)^D} G_{h^2}^{\mu\nu} \left(e^{-il_1 x} \tilde{h}_{\mu\nu}^G(l_{(1)}), e^{-il_2 x} \tilde{h}_{\mu\nu}^G(l_{(2)}) \right) \\
G_{h^2}^{\mu\nu}(h_{\mu\nu}^G, h_{\mu\nu}^G) &= \int \frac{d^D l_{(1)}}{(2\pi)^D} \frac{d^D l_{(2)}}{(2\pi)^D} e^{-i(l_{(1)}+l_{(2)})x} \tilde{G}_{h^2}^{\mu\nu} \left(\tilde{h}_{\mu\nu}^G(l_{(1)}), \tilde{h}_{\mu\nu}^G(l_{(2)}) \right) \\
&= \int \frac{d^D q}{(2\pi)^D} e^{-iqx} \int \frac{d^D l}{(2\pi)^D} \tilde{G}_{h^2}^{\mu\nu} \left(\tilde{h}_{\mu\nu}^G(l), \tilde{h}_{\mu\nu}^G(q-l) \right) \\
(\tilde{G}_{\text{non-linear}}^{\mu\nu})_{G^2} &= \int \frac{d^D l}{(2\pi)^D} \tilde{G}_{h^2}^{\mu\nu} \left(\tilde{h}_{\mu\nu}^G(l), \tilde{h}_{\mu\nu}^G(q-l) \right) \\
\tilde{h}_{\mu\nu}^{G^2} &= 2\frac{G_{\mu\nu\alpha\beta}}{q^2} \int \frac{d^D l}{(2\pi)^D} \left(\tilde{G}_{h^2}^{\alpha\beta} \left(\tilde{h}_{\mu\nu}^G(l), \tilde{h}_{\mu\nu}^G(q-l) \right) + \frac{1}{\xi} \tilde{H}_{h^2}^{\alpha\beta} \left(\tilde{h}_{\mu\nu}^G(l), \tilde{h}_{\mu\nu}^G(q-l) \right) \right)
\end{aligned}$$

$$\begin{aligned}
& \int \frac{d^D l}{(2\pi)^D} \tilde{G}_{h^2}^{\alpha\beta}(\tilde{h}_{\mu\nu}^G(l), \tilde{h}_{\mu\nu}^G(q-l)) \\
&= \int \frac{d^D l}{(2\pi)^D} \tilde{G}_{h^2}^{\alpha\beta} \left(\frac{\kappa^2 G_{\mu\nu\alpha\beta} \tilde{T}^{\alpha\beta}(l)}{2 l^2}, \frac{\kappa^2 G_{\mu\nu\alpha\beta} \tilde{T}^{\alpha\beta}(l-q)}{2 (l-q)^2} \right) \\
&= \frac{\kappa^4}{4} \int \frac{d^D l}{(2\pi)^D} \frac{1}{l^2 (l-q)^2} \tilde{G}_{h^2}^{\alpha\beta} (G_{\mu\nu\alpha\beta} \tilde{T}^{\alpha\beta}(l), G_{\mu\nu\alpha\beta} \tilde{T}^{\alpha\beta}(l-q)) \\
& \quad (G_{\text{non-linear}}^{\mu\nu})_{G^3} = 2G_{h^2}^{\mu\nu}(h_{\mu\nu}^G, h_{\mu\nu}^{G^2}) + G_{h^3}^{\mu\nu}(h_{\mu\nu}^G, h_{\mu\nu}^G, h_{\mu\nu}^G) \\
& \quad G_{h^2}^{\mu\nu}(h_{\mu\nu}^G, h_{\mu\nu}^{G^2}) = \int \frac{d^D q}{(2\pi)^D} e^{-iqx} \int \frac{d^D l}{(2\pi)^D} \tilde{G}_{h^2}^{\mu\nu}(\tilde{h}_{\mu\nu}^G(l), \tilde{h}_{\mu\nu}^{G^2}(q-l)) \\
& \quad G_{h^3}^{\mu\nu}(h_{\mu\nu}^G, h_{\mu\nu}^G, h_{\mu\nu}^G) \\
&= \int \frac{d^D l_{(1)}}{(2\pi)^D} \frac{d^D l_{(2)}}{(2\pi)^D} \frac{d^D l_{(3)}}{(2\pi)^D} e^{-i(l_{(1)}+l_{(2)}+l_{(3)})x} \tilde{G}_{h^3}^{\mu\nu}(\tilde{h}_{\mu\nu}^G(l_{(1)}), \tilde{h}_{\mu\nu}^G(l_{(2)}), \tilde{h}_{\mu\nu}^G(l_{(3)})) \\
&= \int \frac{d^D q}{(2\pi)^D} e^{-iqx} \int \frac{d^D l_{(1)}}{(2\pi)^D} \frac{d^D l_{(2)}}{(2\pi)^D} \tilde{G}_{h^3}^{\mu\nu}(\tilde{h}_{\mu\nu}^G(l_{(1)}), \tilde{h}_{\mu\nu}^G(l_{(2)}-l_{(1)}), \tilde{h}_{\mu\nu}^G(q-l_{(2)})) \\
& \quad \tilde{h}_{\alpha\beta}^{G^3} = 2 \frac{G_{\alpha\beta\mu\nu}}{q^2} \left((\tilde{G}_{\text{non-linear}}^{\mu\nu})_{G^3} + \frac{1}{\xi} (\hat{H}_{\text{non-linear}}^{\mu\nu})_{G^3} \right). \\
& \quad (\tilde{G}_{\text{non-linear}}^{\mu\nu})_{G^3} = 2 \int \frac{d^D l}{(2\pi)^D} \tilde{G}_{h^2}^{\mu\nu}(\tilde{h}_{\mu\nu}^G(l), \tilde{h}_{\mu\nu}^{G^2}(q-l)) \\
& \quad + \int \frac{d^D l_{(1)}}{(2\pi)^D} \frac{d^D l_{(2)}}{(2\pi)^D} \tilde{G}_{h^3}^{\mu\nu}(\tilde{h}_{\mu\nu}^G(l_{(1)}), \tilde{h}_{\mu\nu}^G(l_{(2)}-l_{(1)}), \tilde{h}_{\mu\nu}^G(q-l_{(2)})) \\
& \quad T^{\mu\nu} = \frac{k^\mu k^\nu}{m} \delta^{D-1}(x_\perp) \\
& \quad = m \eta_{\parallel}^{\mu\nu} \delta^{D-1}(x_\perp) \\
& \quad \tilde{T}^{\mu\nu}(q) = \int d^D x e^{iqx} T^{\mu\nu} = 2\pi \delta(q_{\parallel}) m \eta_{\parallel}^{\mu\nu} \\
& \quad \tilde{h}_{\alpha\beta}^{G^1}(q) = \frac{G_{\alpha\beta\mu\nu} \kappa^2}{q^2} \frac{1}{2} 2\pi \delta(q_{\parallel}) m \eta_{\parallel}^{\mu\nu} \\
& \quad = \frac{\kappa^2 m}{2} \frac{2\pi \delta(q_{\parallel})}{q_\perp^2} G_{\alpha\beta\mu\nu} \eta_{\parallel}^{\mu\nu} \\
& \quad G_{\alpha\beta\mu\nu} \tilde{T}^{\mu\nu} = \mathcal{P}^{-1}{}_{\alpha\beta\mu\nu} \tilde{T}^{\mu\nu} \\
& \quad \mathcal{P}^{-1}{}_{\alpha\beta\mu\nu} \tilde{\eta}_{\parallel}^{\mu\nu} = \eta_{\parallel}^{\mu\nu} - \frac{1}{D-2} \eta^{\mu\nu} \\
& \quad = \frac{D-3}{D-2} \left(\eta_{\mu\nu}^{\parallel} - \frac{1}{D-3} \eta_{\mu\nu}^{\perp} \right)
\end{aligned}$$

$$\tilde{h}_{\mu\nu}^G = \frac{\kappa^2 m}{2} \frac{2\pi\delta(q_{\parallel})}{q_{\perp}^2} \frac{D-3}{D-2} \left(\eta_{\mu\nu}^{\parallel} - \frac{1}{D-3} \eta_{\mu\nu}^{\perp} \right)$$

$$h_{\mu\nu}^G = -\frac{\mu}{\sqrt{-x_{\perp}^2}^{D-3}} \left(\eta_{\mu\nu}^{\parallel} - \frac{1}{D-3} \eta_{\mu\nu}^{\perp} \right)$$

$$(G_{\sigma})^G = G_{\sigma}^h(h_{\mu\nu}^G) = \partial^{\rho} \mathcal{P}_{\rho\sigma}^{\mu\nu} h_{\mu\nu}^G$$

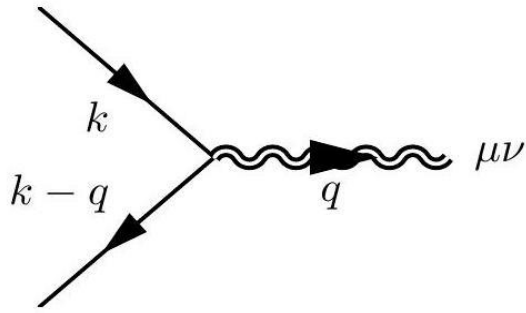


$$2\pi\delta(kq)\mathcal{M}_{\text{vertex}}^{\mu\nu} = -\kappa\tilde{\tau}^{\mu\nu} - \frac{4}{\kappa\xi}\tilde{H}_{\text{non-linear}}^{\mu\nu}$$

$$g_{\mu\nu} = \eta_{\mu\nu} - \frac{\kappa}{2} \int \frac{d^D q \delta(kq) e^{-iqx}}{(2\pi)^{D-1}} \frac{G_{\mu\nu\alpha\beta}}{q^2} \mathcal{M}_{\text{vertex}}^{\alpha\beta}$$

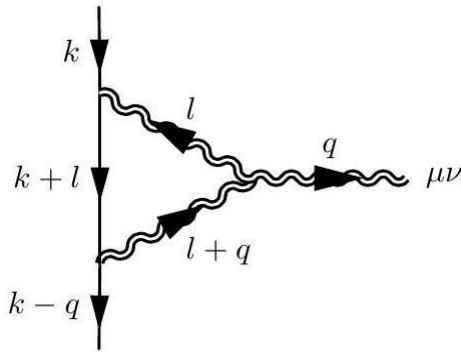
$$\tilde{h}_{\mu\nu} = \frac{\mathcal{P}_{\mu\nu\alpha\beta}^{-1}}{q^2} \left(\frac{\kappa^2}{2} \tilde{\tau}^{\alpha\beta} + 2\tilde{H}_{\text{non-linear}}^{\alpha\beta} \right)$$

$$\begin{aligned} V_{\phi^2 h}^{\mu\nu}(k, k-q, m) &= I_{\alpha\beta}^{\mu\nu} k^{\alpha} (k-q)^{\beta} - \frac{k(k-q) - m^2}{2} \eta^{\mu\nu} \\ &= I_{\alpha\beta}^{\mu\nu} k^{\alpha} (k-q)^{\beta} + \frac{kq}{2} \eta^{\mu\nu} \end{aligned}$$



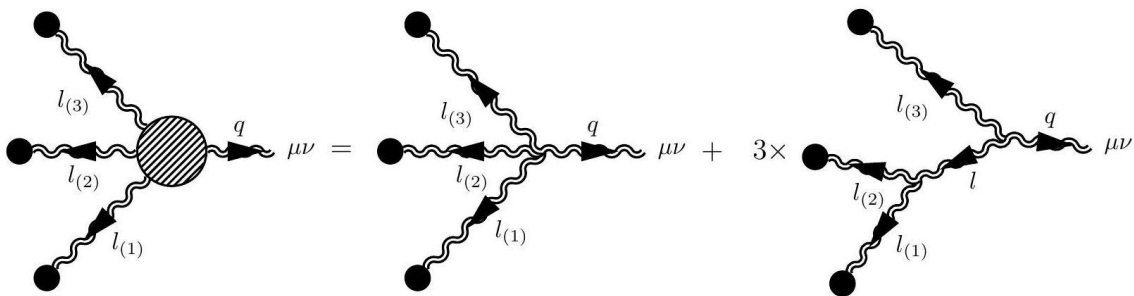
$$\mathcal{M}_{\text{tree}}^{\mu\nu} = -\kappa k^\mu k^\nu = -\kappa m^2 \eta_{\parallel}^{\mu\nu}$$

$$\begin{aligned} \tilde{\tau}^{\mu\nu} &= -\frac{2\pi\delta(kq)}{\kappa} \mathcal{M}_{\text{tree}}^{\mu\nu} \\ &= 2\pi\delta(q_{\parallel}) m \eta_{\parallel}^{\mu\nu} \end{aligned}$$



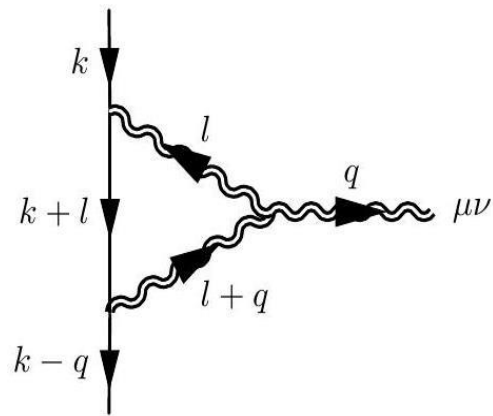
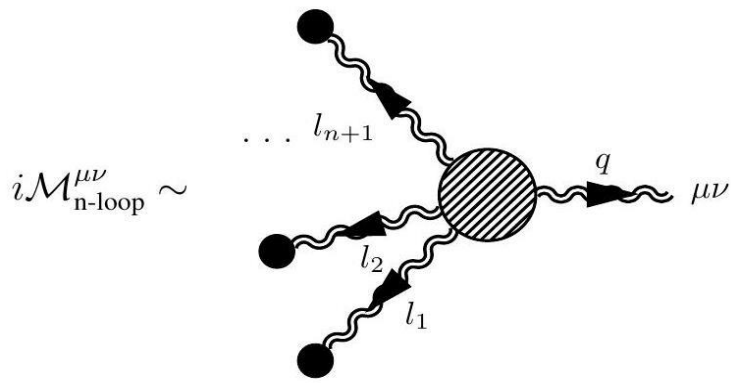
$$i2\kappa V_{h^3}^{\mu\nu\alpha\beta\gamma\delta}(q, l, l+q) \tilde{f}_{\alpha\beta}(l) \tilde{f}_{\gamma\delta}(l+q) = -i4\kappa \left(\tilde{\mathcal{G}}_{h^2}^{\mu\nu}(\tilde{f}_{\alpha\beta}^{(l)}, \tilde{f}_{\gamma\delta}^{(l+q)}) + \frac{1}{\xi} \tilde{\mathcal{H}}_{h^2}^{\mu\nu}(\tilde{f}_{\alpha\beta}^{(l)}, \tilde{f}_{\gamma\delta}^{(l+q)}) \right)$$

$$\mathcal{M}_{1\text{-loop}}^{\mu\nu} \sim \int \frac{d^D l}{(2\pi)^D} V_{h^3}^{\mu\nu\alpha\beta\gamma\delta}(q, l, l+q) \tilde{h}_{\alpha\beta}^G(l) \tilde{h}_{\gamma\delta}^G(l+q)$$



$$\int \frac{d^D l}{(2\pi)^D} \left(\tilde{\mathcal{G}}_{h^2}^{\mu\nu}(\tilde{h}_{\mu\nu}^G(l), \tilde{h}_{\mu\nu}^{G^2}(q-l)) + \frac{1}{\xi} \tilde{\mathcal{H}}_{h^2}^{\mu\nu}(\tilde{h}_{\mu\nu}^G(l), \tilde{h}_{\mu\nu}^{G^2}(q-l)) \right)$$

$$\int \frac{d^D l_{(1)}}{(2\pi)^D} \frac{d^D l_{(2)}}{(2\pi)^D} \tilde{\mathcal{G}}_{h^3}^{\mu\nu}(\tilde{h}_{\mu\nu}^G(l_{(1)}), \tilde{h}_{\mu\nu}^G(l_{(2)} - l_{(1)}), \tilde{h}_{\mu\nu}^G(q - l_{(2)}))$$



$$I = \int \frac{d^D l}{(2\pi)^D} \frac{1}{l^2 + i\epsilon} \frac{1}{(l + q_\perp)^2 + i\epsilon} \frac{1}{(l + k)^2 - m^2 + i\epsilon}$$

$$\frac{1}{(l + k)^2 - m^2 + i\epsilon} \approx \frac{1}{2kl + i\epsilon}$$

$$= \frac{1}{2m} \frac{1}{l_\parallel + i\epsilon}$$

$$\frac{1}{l_\parallel + i\epsilon} = \frac{1}{l_\parallel} - i\pi\delta(l_\parallel)$$

$$I = -\frac{i}{4m} \int \frac{d^D l}{(2\pi)^D} 2\pi\delta(l_\parallel) \frac{1}{l_\perp^2} \frac{1}{(l_\perp + q_\perp)^2}$$

$$N_{D-1} = \int \frac{d^D l}{(2\pi)^D} 2\pi\delta(l_\parallel) \frac{1}{l_\perp^2} \frac{1}{(l_\perp + q_\perp)^2}$$

$$= \int \frac{d^{D-1} l_\perp}{(2\pi)^{D-1}} \frac{1}{l_\perp^2} \frac{1}{(l_\perp + q_\perp)^2}$$

$$N_d = -\frac{\Omega_{d-2} \sqrt{-q_\perp^2}^{d-4}}{4(4\pi)^{d-2} \sin\left(\frac{\pi}{2}d\right)}$$



$$I^\mu = \int \frac{d^D l}{(2\pi)^D} \frac{l^\mu}{(l^2 + i\epsilon)((l + q_\perp)^2 + i\epsilon)((l + k)^2 - m^2 + i\epsilon)}$$

$$I^{\mu\nu} = \int \frac{d^D l}{(2\pi)^D} \frac{l^\mu l^\nu}{(l^2 + i\epsilon)((l + q_\perp)^2 + i\epsilon)((l + k)^2 - m^2 + i\epsilon)}$$

$$I^\mu = Aq_\perp^\mu + Bk^\mu$$

$$I_\perp^\mu = Aq_\perp^\mu$$

$$= -\frac{i}{4m} \int \frac{d^D l}{(2\pi)^D} 2\pi\delta(l_\parallel) \frac{l_\perp^\mu}{l_\perp^2(l_\perp + q_\perp)^2}$$

$$N_{D-1}^\mu = \int \frac{d^D l}{(2\pi)^D} 2\pi\delta(l_\parallel) \frac{l_\perp^\mu}{l_\perp^2(l_\perp + q_\perp)^2}$$

$$2q_\perp^\mu N_{D-1}^\mu = \int \frac{d^D l}{(2\pi)^D} 2\pi\delta(l_\parallel) \frac{1}{l_\perp^2} - \int \frac{d^D l}{(2\pi)^D} 2\pi\delta(l_\parallel) \frac{1}{(l_\perp + q_\perp)^2} - q_\perp^2 \int \frac{d^D l}{(2\pi)^D} 2\pi\delta(l_\parallel) \frac{1}{l_\perp^2(l_\perp + q_\perp)^2} = -q_\perp^2 N_{D-1}$$

$$N_{D-1}^\mu = -\frac{1}{2} q_\perp^\mu N_{D-1}$$

$$I_\parallel = mB$$

$$= \frac{1}{2m} \int \frac{d^D l}{(2\pi)^D} \frac{1}{l^2(l + q_\perp)^2} \frac{l_\parallel}{l_\parallel + i\epsilon}$$

$$I_\parallel = i \frac{N_D}{2m}$$

$$I^\mu = -\frac{i}{4m} N_{D-1}^\mu + \frac{i}{2m^2} N_D k^\mu$$

$$= \frac{i}{8m} N_{D-1} q_\perp^\mu + \frac{i}{2m^2} N_D k^\mu$$

$$I^{\mu\nu} = -\frac{i}{4m} N_{D-1}^{\mu\nu} - \frac{i}{4m^2} N_D (q_\perp^\mu k^\nu + k^\mu q_\perp^\nu)$$

$$N_{D-1}^{\mu\nu} = \int \frac{d^D l}{(2\pi)^D} 2\pi\delta(l_\parallel) \frac{l_\perp^\mu l_\perp^\nu}{l_\perp^2(l_\perp + q_\perp)^2}$$

$$= \frac{q_\perp^2}{4(D-2)} N_{D-1} \left((D-1) \frac{q_\perp^\mu q_\perp^\nu}{q_\perp^2} - \eta_\perp^{\mu\nu} \right)$$

$$2\pi\delta(kq) i\mathcal{M}_{1\text{-loop}}^{\mu\nu} = -i \frac{4}{\kappa} \left(\tilde{G}_{1\text{-loop}} + \frac{1}{\xi} \tilde{H}_{1\text{-loop}} \right)$$

$$= 2\pi\delta(kq) m^4 (2i\kappa) (-i\kappa)^2 i^3 \int \frac{d^D l}{(2\pi)^D} \frac{1}{l^2(l + q_\perp)^2((l + k)^2 - m^2 + i\epsilon)}$$

$$\times f_{\alpha\beta} f_{\gamma\delta} V_{h^3}^{\alpha\beta} \gamma\delta\mu\nu(l, -l - q, q)$$

$$f_{\alpha\beta} = \frac{k^\mu k^\nu}{m^2} \mathcal{P}^{-1}{}_{\mu\nu\alpha\beta} = \eta_\parallel^{\mu\nu} \mathcal{P}^{-1}{}_{\mu\nu\alpha\beta}$$



$$-2(1 - \xi)I_{\rho\kappa}^{\mu\nu} \frac{l^\rho l_\sigma}{l^2} I_{\alpha\beta}^{\kappa\sigma}$$

$$\int \frac{d^D l}{(2\pi)^D} \frac{lk}{l^2 l^2 (l+q)^2 ((l+k)^2 - m^2 + i\epsilon)} = \frac{1}{2} \int \frac{d^D l}{(2\pi)^D} \frac{((l+k)^2 - m^2) - l^2}{l^2 l^2 (l+q)^2 ((l+k)^2 - m^2 + i\epsilon)}$$

$$= \frac{1}{2} \int \frac{d^D l}{(2\pi)^D} \frac{1}{l^2 l^2 (l+q)^2} - \frac{1}{2} \int \frac{d^D l}{(2\pi)^D} \frac{1}{l^2 (l+q)^2 ((l+k)^2 - m^2 + i\epsilon)}$$

$$f_{\mu\nu} = \frac{D-3}{D-2} \eta_{\mu\nu}^\parallel - \frac{1}{D-2} \eta_{\mu\nu}^\perp$$

$$f_{\alpha\beta} f_{\gamma\delta} \tau_{h^3}^{\alpha\beta\gamma\delta\nu} (l, -l-q, q) = f_{\alpha\beta} f_{\gamma\delta} (U^{\mu\nu\alpha\beta\rho\gamma\delta\sigma} l_\rho (l+q)_\sigma + U^{\alpha\beta\gamma\delta\mu\nu\sigma} (l+q)_\rho q_\sigma - U^{\gamma\delta\mu\nu\rho\alpha\beta\sigma} q_\rho l_\sigma)$$

$$f_{\alpha\beta} f_{\gamma\delta} \tau_{h^3}^{\alpha\beta\gamma\delta\mu\nu} (l, -l-q, q) = f_{\alpha\beta} f_{\gamma\delta} (U^{\mu\nu\alpha\beta\rho\gamma\delta\sigma} l_\rho (l+q)_\sigma + U^{\alpha\beta\gamma\delta\rho\mu\nu\sigma} q_\rho q_\sigma)$$

$$2\pi\delta(kq) i \mathcal{M}_{1\text{-loop}}^{\mu\nu} = -2\pi\delta(kq) 2m^4 \kappa^3 f_{\alpha\beta} f_{\gamma\delta} (U^{\mu\nu\alpha\beta\rho\gamma\delta\sigma} (I_{\rho\sigma} + I_\rho q_\sigma) + U^{\alpha\beta\gamma\delta\rho\mu\nu\sigma} q_\rho q_\sigma I)$$

$$I_{\rho\sigma} + I_\rho q_\sigma = i \frac{q^2 N_{D-1}}{16(D-2)m} \left((D-3) \frac{q_\rho^\perp q_\sigma^\perp}{q_\perp^2} + \eta_{\rho\sigma}^\perp \right) + i \frac{N_D}{4m^2} (k_\rho q_\sigma^\perp - q_\rho^\perp k_\sigma)$$

$$2\pi\delta(kq) \mathcal{M}_{1\text{-loop}}^{\mu\nu} = 2\pi\delta(q_\parallel) \frac{m^2 \kappa^3 q_\perp^2 N_{D-1}}{2} f_{\alpha\beta} f_{\gamma\delta} \left(-\frac{1}{4(D-2)} U^{\mu\nu\alpha\beta\rho\gamma\delta\sigma} M_{\rho\sigma}^\perp + U^{\alpha\beta\gamma\delta\rho\mu\nu\sigma} \frac{q_\rho^\perp q_\sigma^\perp}{q_\perp^2} \right)$$

$$M_{\rho\sigma}^\perp = (D-3) \frac{q_\rho^\perp q_\sigma^\perp}{q_\perp^2} + \eta_{\rho\sigma}^\perp$$

$$U_{gf}^{\mu\nu\alpha\beta\rho\gamma\delta\sigma} f_{\alpha\beta} f_{\gamma\delta} M_{\rho\sigma}^\perp$$

$$-\mathcal{P}_{\alpha\beta}^{\rho\sigma} h_{,\sigma}^{\alpha\beta} h_\rho^{\mu,\nu} \rightarrow -\mathcal{P}_{\alpha\beta}^{\rho\sigma} f^{\alpha\beta} f_\rho^\mu M_{\sigma}^{\perp\nu}$$

$$U_{gf}^{\mu\nu\alpha\beta\rho\gamma\delta\sigma} f_{\alpha\beta} f_{\gamma\delta} M_{\rho\sigma}^\perp = \alpha \mathcal{P}_{\alpha\beta}^{\rho\sigma} f^{\alpha\beta} (-f_\rho^\mu M_\sigma^{\perp\nu} - f_\rho^\nu M_\sigma^{\perp\mu} + f^{\mu\nu} M_{\rho\sigma}^\perp)$$

$$= \alpha \eta_{\parallel}^{\rho\sigma} (-f_\rho^\mu M_\sigma^{\perp\nu} - f_\rho^\nu M_\sigma^{\perp\mu} + f^{\mu\nu} M_{\rho\sigma}^\perp) = 0$$

$$U_{gf}^{\alpha\beta\gamma\delta\rho\mu\nu\sigma} f_{\alpha\beta} f_{\gamma\delta} \frac{q_\rho^\perp q_\sigma^\perp}{q_\perp^2}$$

$$-\alpha (I^{\alpha\beta\rho\kappa} I_{\kappa\lambda}^{\gamma\delta} \mathcal{P}^{\lambda\sigma\mu\nu} + I^{\alpha\beta\sigma\kappa} I_{\kappa\lambda}^{\mu\nu} \mathcal{P}^{\lambda\rho\gamma\delta}) f_{\alpha\beta} f_{\gamma\delta} \frac{q_\rho^\perp q_\sigma^\perp}{q_\perp^2} = -\alpha \left(f^{\rho\kappa} f_{\kappa\lambda} \mathcal{P}^{\lambda\sigma\mu\nu} \frac{q_\rho^\perp q_\sigma^\perp}{q_\perp^2} + f^{\sigma\kappa} I_{\kappa\lambda}^{\mu\nu} \eta_{\parallel}^{\lambda\rho} \frac{q_\rho^\perp q_\sigma^\perp}{q_\perp^2} \right)$$

$$= -\alpha \frac{1}{(D-2)^2} \mathcal{P}^{\rho\sigma\mu\nu} \frac{q_\rho^\perp q_\sigma^\perp}{q_\perp^2}$$

$$\frac{1}{2} \alpha (\mathcal{P}^{\gamma\delta\rho\sigma} I^{\mu\nu\alpha\beta} + \mathcal{P}^{\mu\nu\rho\sigma} I^{\gamma\delta\alpha\beta}) f_{\alpha\beta} f_{\gamma\delta} \frac{q_\rho^\perp q_\sigma^\perp}{q_\perp^2} = \frac{1}{2} \alpha \left(\eta_{\parallel}^{\rho\sigma} f^{\mu\nu} \frac{q_\rho^\perp q_\sigma^\perp}{q_\perp^2} + \mathcal{P}^{\mu\nu\rho\sigma} f_{\alpha\beta} f^{\alpha\beta} \frac{q_\rho^\perp q_\sigma^\perp}{q_\perp^2} \right)$$

$$= \frac{1}{2} \alpha f_{\alpha\beta} f^{\alpha\beta} \mathcal{P}^{\mu\nu\rho\sigma} \frac{q_\rho^\perp q_\sigma^\perp}{q_\perp^2}$$



$$\begin{aligned}
U_{gf}^{\alpha\beta} \gamma\delta\rho\mu\nu\sigma f_{\alpha\beta} f_{\gamma\delta} \frac{q_{\rho}^{\perp} q_{\sigma}^{\perp}}{q_{\perp}^2} &= \alpha \left(-\frac{1}{(D-2)^2} + \frac{1}{2} f_{\alpha\beta} f^{\alpha\beta} \right) \mathcal{P}^{\mu\nu\rho\sigma} \frac{q_{\rho}^{\perp} q_{\sigma}^{\perp}}{q_{\perp}^2} \\
&= \alpha \frac{D-3}{2(D-2)} \mathcal{P}^{\mu\nu\rho\sigma} \frac{q_{\rho}^{\perp} q_{\sigma}^{\perp}}{q_{\perp}^2} \\
&= \alpha \frac{D-3}{4(D-2)} \left(2 \frac{q_{\perp}^{\mu} q_{\perp}^{\nu}}{q_{\perp}^2} - \eta^{\mu\nu} \right)
\end{aligned}$$

$$\tilde{H}_{1-\text{loop}}^{\mu\nu} = -\alpha 2\pi\delta(q_{\parallel}) \frac{\kappa^4 m^2 q^2 N_{D-1}}{16} \frac{D-3}{D-2} \mathcal{P}^{\mu\nu\rho\sigma} \frac{q_{\rho}^{\perp} q_{\sigma}^{\perp}}{q_{\perp}^2}.$$

$$U_{(c)}^{\mu\nu} \alpha\beta\rho\gamma\delta\sigma f_{\alpha\beta} f_{\gamma\delta} M_{\rho\sigma}^{\perp}$$

$$U_{(c)}^{\alpha\beta} \gamma\delta\rho\mu\nu\sigma f_{\alpha\beta} f_{\gamma\delta} \frac{q_{\rho}^{\perp} q_{\sigma}^{\perp}}{q_{\perp}^2}$$

$$\begin{aligned}
U_{(c)}^{\mu\nu} \alpha\beta\rho\gamma\delta\sigma h_{\alpha\beta,\rho} h_{\gamma\delta,\sigma} &= 2I_{\phi\epsilon}^{\mu\nu} \mathcal{P}_{\rho\sigma}^{\alpha\beta} \mathcal{P}_{\gamma\delta}^{\sigma\phi} h_{\alpha\beta}^{\epsilon} h^{\gamma\delta,\rho} - \mathcal{P}_{\alpha\beta}^{\mu\rho} \mathcal{P}_{\gamma\delta}^{\nu\sigma} \eta_{\rho\sigma} h_{,\kappa}^{\alpha\beta} h^{\gamma\delta,\kappa} \\
&\quad + \mathcal{P}_{\rho\sigma}^{\mu\nu} \left(h_{,\beta}^{\rho\alpha} h_{,\alpha}^{\sigma\beta} - \frac{1}{2} h_{\beta}^{\alpha,\rho} h_{\alpha}^{\beta,\sigma} - h_{,\alpha}^{\rho\sigma} h_{,\beta}^{\alpha\beta} \right)
\end{aligned}$$

$$2I_{\phi\epsilon}^{\mu\nu} \mathcal{P}_{\rho\sigma}^{\alpha\beta} \mathcal{P}_{\gamma\delta}^{\sigma\phi} f_{\alpha\beta} f^{\gamma\delta} M_{\perp}^{\epsilon\rho} = 2I_{\phi\epsilon}^{\mu\nu} \eta_{\rho\sigma}^{\parallel} \eta_{\parallel}^{\sigma\phi} M_{\perp}^{\epsilon\rho} = 0$$

$$-\mathcal{P}_{\alpha\beta}^{\mu\rho} \mathcal{P}_{\gamma\delta}^{\nu\sigma} \eta_{\rho\sigma} f^{\alpha\beta} f^{\gamma\delta} M_{\kappa}^{\perp\kappa} = -\eta_{\parallel}^{\mu\rho} \eta_{\parallel}^{\nu\sigma} \eta_{\rho\sigma} M_{\kappa}^{\perp\kappa} = -\eta_{\parallel}^{\mu\nu} M_{\kappa}^{\perp\kappa}$$

$$M_{\kappa}^{\perp\kappa} = D-3 + D-1 = 2(D-2)$$

$$-\mathcal{P}_{\alpha\beta}^{\mu\rho} \mathcal{P}_{\gamma\delta}^{\nu\sigma} \eta_{\rho\sigma} f^{\alpha\beta} f^{\gamma\delta} M_{\kappa}^{\perp\kappa} = -2(D-2) \eta_{\parallel}^{\mu\nu}$$

$$f^{\rho\alpha} f^{\sigma\beta} M_{\alpha\beta}^{\perp} = \frac{1}{(D-2)^2} M_{\perp}^{\rho\sigma}$$

$$-\frac{1}{2} f^{\alpha\beta} f_{\alpha\beta} M_{\perp}^{\rho\sigma} = -\frac{1}{2} \frac{(D-3)^2 + D-1}{(D-2)^2} M_{\perp}^{\rho\sigma}$$

$$-f^{\rho\sigma} f^{\alpha\beta} M_{\alpha\beta}^{\perp} = \frac{1}{D-2} f^{\rho\sigma} M_{\kappa}^{\perp\kappa} = 2f^{\rho\sigma}$$

$$f^{\rho\alpha} f^{\sigma\beta} M_{\alpha\beta} - \frac{1}{2} f^{\alpha\beta} f_{\alpha\beta} M^{\rho\sigma} - f^{\rho\sigma} f^{\alpha\beta} M_{\alpha\beta} = -\frac{1}{2} \frac{D-3}{D-2} M^{\rho\sigma} + 2f^{\rho\sigma}$$

$$\mathcal{P}_{\rho\sigma}^{\mu\nu} M^{\rho\sigma} = -(D-2) \eta_{\parallel}^{\mu\nu} + (D-3) \left(\frac{q_{\perp}^{\mu} q_{\perp}^{\nu}}{q_{\perp}^2} - \eta_{\perp}^{\mu\nu} \right)$$

$$\mathcal{P}_{\rho\sigma}^{\mu\nu} \left(f^{\rho\alpha} f^{\sigma\beta} M_{\alpha\beta} - \frac{1}{2} f^{\alpha\beta} f_{\alpha\beta} M^{\rho\sigma} - f^{\rho\sigma} f^{\alpha\beta} M_{\alpha\beta} \right)$$

$$= \frac{D+1}{2} \eta_{\parallel}^{\mu\nu} - \frac{1}{2} \frac{(D-3)^2}{D-2} \left(\frac{q_{\perp}^{\mu} q_{\perp}^{\nu}}{q_{\perp}^2} - \eta_{\perp}^{\mu\nu} \right)$$



$$U^{\mu\nu\alpha\beta\rho\gamma\delta\sigma} f_{\alpha\beta} f_{\gamma\delta} M_{\rho\sigma} = -\frac{3(D-3)}{2} \eta_{\parallel}^{\mu\nu} - \frac{1}{2} \frac{(D-3)^2}{D-2} \left(\frac{q_{\perp}^{\mu} q_{\perp}^{\nu}}{q_{\perp}^2} - \eta_{\perp}^{\mu\nu} \right)$$

$$U^{\alpha\beta\gamma\delta\rho\mu\nu\sigma} f_{\alpha\beta} f_{\gamma\delta} \frac{q_{\rho}^{\perp} q_{\sigma}^{\perp}}{q_{\perp}^2} = \frac{1}{2} \frac{D-4}{D-2} \eta_{\parallel}^{\mu\nu} - \frac{1}{2} \frac{D-1}{D-2} \left(\eta_{\perp}^{\mu\nu} - \frac{q_{\perp}^{\mu} q_{\perp}^{\nu}}{q_{\perp}^2} \right)$$

$$\tilde{G}_{1\text{-loop}}^{\mu\nu} = 2\pi\delta(q_{\parallel}) \frac{\kappa^4 m^2 q^2 N_{D-1}}{64} \left(\frac{D-7}{D-2} \eta_{\parallel}^{\mu\nu} - \frac{(D-3)(3D-5)}{(D-2)^2} \left(\eta_{\perp}^{\mu\nu} - \frac{q^{\mu} q^{\nu}}{q^2} \right) \right).$$

$$2\pi\delta(kq) \mathcal{M}_{1\text{-loop}}^{\mu\nu} = -\frac{4}{\kappa} \left(\tilde{G}_{1\text{-loop}}^{\mu\nu} + \frac{1}{\xi} \tilde{H}_{1\text{-loop}}^{\mu\nu} \right)$$

$$\tilde{G}_{1\text{-loop}}^{\mu\nu} = 2\pi\delta(q_{\parallel}) \frac{\kappa^4 m^2 q^2 N_{D-1}}{64} \left(\frac{D-7}{D-2} \eta_{\parallel}^{\mu\nu} - \frac{(D-3)(3D-5)}{(D-2)^2} \left(\eta_{\perp}^{\mu\nu} - \frac{q^{\mu} q^{\nu}}{q^2} \right) \right)$$

$$\tilde{H}_{1\text{-loop}}^{\mu\nu} = -\alpha 2\pi\delta(q_{\parallel}) \frac{\kappa^4 m^2 q^2 N_{D-1}}{16} \frac{D-3}{D-2} \mathcal{P}^{\mu\nu\rho\sigma} \frac{q_{\rho}^{\perp} q_{\sigma}^{\perp}}{q_{\perp}^2}$$

$$G_{(c)}^{\mu\nu\alpha\beta} \tilde{H}_{\text{non-linear}}^{\alpha\beta} = 0$$

$$G_{c_{\mu\nu}}^{\alpha\beta} \mathcal{P}^{\mu\nu\rho\sigma} \frac{q_{\rho}^{\perp} q_{\sigma}^{\perp}}{q_{\perp}^2} = \frac{q_{\mu}^{\perp} q_{\nu}^{\perp}}{q_{\perp}^2} - 2J_{\mu\nu}^{\alpha\beta} \mathcal{P}^{\rho\sigma} \frac{q_{\rho}^{\perp} q_{\sigma}^{\perp}}{q_{\perp}^2}$$

$$\begin{aligned} J_{\mu\nu}^{\alpha\beta} \mathcal{P}^{\rho\sigma} \frac{q_{\rho}^{\perp} q_{\sigma}^{\perp}}{q_{\perp}^2} &= J_{\mu\nu}^{\alpha\beta} \mathcal{P}^{\rho\sigma} J_{\rho\sigma}^{\gamma\delta} \eta_{\gamma\delta} \\ &= \frac{1}{2} J_{\mu\nu}^{\gamma\delta} \eta_{\gamma\delta} \\ &= \frac{1}{2} \frac{q_{\mu}^{\perp} q_{\nu}^{\perp}}{q_{\perp}^2} \end{aligned}$$

$$\tilde{h}_{\mu\nu}^{G^2} = 2 \frac{G_{\mu\nu\alpha\beta}^{(c)}}{q^2} \tilde{G}_{1\text{-loop}}^{\alpha\beta} + 2 \frac{G_{\mu\nu\alpha\beta}^{(gf)}}{q^2} \tilde{H}_{1\text{-loop}}^{\alpha\beta}$$

$$\frac{G_{\mu\nu\alpha\beta}^{(c)}}{q^2} 2\tilde{G}_{1\text{-loop}}^{\alpha\beta} = 2\pi\delta(q_{\parallel}) \frac{\kappa^4 m^2 N_{D-1}}{32} \left(4 \frac{(D-3)^2}{(D-2)^2} \eta_{\parallel}^{\mu\nu} + \frac{(D-3)(3D-5)}{(D-2)^2} \frac{q_{\mu} q_{\nu}}{q^2} - \frac{D-7}{(D-2)^2} \eta_{\perp}^{\mu\nu} \right)$$

$$\frac{G_{\mu\nu\alpha\beta}^{(gf)}}{q^2} 2\tilde{H}_{1\text{-loop}}^{\alpha\beta} = -2\pi\delta(q_{\parallel}) \alpha \frac{\kappa^4 m^2 N_{D-1}}{8} \frac{D-3}{D-2} \frac{q_{\mu} q_{\nu}}{q^2}$$

$$\mathcal{P}^{\mu\nu\rho\sigma} \frac{q_{\rho}^{\perp} q_{\sigma}^{\perp}}{q_{\perp}^2}$$

$$\tilde{h}_{\mu\nu}^{G^2} = 2\pi\delta(q_{\parallel}) \frac{\kappa^4 m^2 N_{D-1}}{32} \frac{(D-3)^2}{(D-2)^2} \left(4\eta_{\parallel}^{\mu\nu} + \left(\frac{3D-5}{D-3} - 4\alpha \frac{D-2}{D-3} \right) \frac{q_{\mu} q_{\nu}}{q^2} - \frac{D-7}{(D-3)^2} \eta_{\perp}^{\mu\nu} \right)$$

$$N_{D-1} = \frac{\Omega_{D-3} \sqrt{-q_{\perp}^2}^{D-5}}{4(4\pi)^{D-3} \cos\left(\frac{\pi}{2} D\right)}$$



$$\int \frac{d^d q_{\perp}}{(2\pi)^d} e^{-ix_{\perp} q_{\perp}} (-q_{\perp}^2)^{\frac{n}{2}} = \frac{2^n \Gamma\left(\frac{d+n}{2}\right)}{\sqrt{\pi}^d \Gamma\left(-\frac{n}{2}\right)} \frac{1}{(-x_{\perp}^2)^{\frac{d+n}{2}}}$$

$$\int \frac{d^{D-1} q_{\perp} e^{-iq_{\perp} x_{\perp}}}{(2\pi)^{D-1}} N_{D-1} = \left(\frac{1}{\Omega_{D-2}(D-3)\sqrt{-x_{\perp}^2}} \right)^2$$

$$\frac{\partial}{\partial x_{\perp}^{\mu}} \frac{\partial}{\partial x_{\perp}^{\nu}} \int \frac{d^d q_{\perp}}{(2\pi)^d} e^{-ix_{\perp} q_{\perp}} (-q_{\perp}^2)^{\frac{n}{2}} = - \int \frac{d^d q_{\perp}}{(2\pi)^d} e^{-ix_{\perp} q_{\perp}} q_{\mu}^{\perp} q_{\nu}^{\perp} (-q_{\perp}^2)^{\frac{n}{2}}$$

$$\int \frac{d^{D-1} q_{\perp} e^{-iq_{\perp} x_{\perp}}}{(2\pi)^{D-1}} N_{D-1} \frac{q_{\mu}^{\perp} q_{\nu}^{\perp}}{q_{\perp}^2} = \left(\frac{1}{\Omega_{D-2}(D-3)\sqrt{-x_{\perp}^2}^{D-3}} \right)^2 \frac{1}{D-5} \left(2(D-3) \frac{x_{\mu}^{\perp} x_{\nu}^{\perp}}{x_{\perp}^2} - \eta_{\mu\nu}^{\perp} \right)$$

$$h_{\mu\nu}^{G^2} = \frac{\mu^2}{r^{2(D-3)}} \left(\frac{1}{2} \eta_{\mu\nu}^{\parallel} - \frac{(4\alpha-3)D-8\alpha+5}{4(D-5)} \frac{x_{\mu}^{\perp} x_{\nu}^{\perp}}{x_{\perp}^2} - \frac{2(1-\alpha)D^2 - (13-10\alpha)D + 25 - 12\alpha}{4(D-3)^2(D-5)} \eta_{\mu\nu}^{\perp} \right)$$

$$G_{\sigma}^{G^2} = \mathcal{P}_{\rho\sigma}^{\mu\nu} h_{G^2}{}^{\rho}{}_{\mu\nu} - \alpha \Gamma_{\sigma\alpha\beta}^{\rho\mu\nu} h_G^{\alpha\beta} h_{\mu\nu,\rho}^G$$

$$N_{D-1} = \frac{(-1)^n (-q_{\perp}^2)^n}{8(16\pi)^n n! \sqrt{-q_{\perp}^2}}$$

$$N_{D-1} = \frac{1}{\epsilon} \frac{(-1)^{n+1} \Omega_{2+2n} (-q_{\perp}^2)^n}{(4\pi)^{3+2n}} + \frac{(-1)^{n+1} n!}{16\pi^2 (4\pi)^n (1+2n)!} (-q_{\perp}^2)^n \ln(-q_{\perp}^2 r_0^2) + \dots$$

$$N_{D-1} = \frac{(-1)^{n+1} n!}{16\pi^2 (4\pi)^n (1+2n)!} (-q_{\perp}^2)^n \ln(-q_{\perp}^2 r_0^2)$$

$$\ln(-q^2) = \frac{1}{\epsilon} ((-q^2)^{\epsilon} - 1)$$

$$\int \frac{d^{D-1} q_{\perp} e^{-iq_{\perp} x_{\perp}}}{(2\pi)^{D-1}} \frac{N_{D-1}}{q_{\perp}^2} = \frac{(-1)^n n!}{16\pi^2 (4\pi)^n (1+2n)!} \int \frac{d^{D-1} q_{\perp} e^{-iq_{\perp} x_{\perp}}}{(2\pi)^{D-1}} (-q_{\perp}^2)^{n-1} \ln(-q_{\perp}^2)$$

$$= \frac{(-1)^n n!}{16\pi^2 (4\pi)^n (1+2n)!} \int \frac{d^{D-1} q_{\perp} e^{-iq_{\perp} x_{\perp}}}{(2\pi)^{D-1}} \frac{1}{\epsilon} ((-q_{\perp}^2)^{n-1+\epsilon} - (-q_{\perp}^2)^{n-1})$$

$$\int \frac{d^{D-1} q_{\perp} e^{-iq_{\perp} x_{\perp}}}{(2\pi)^{D-1}} \frac{N_{D-1}}{q_{\perp}^2} = \frac{(-1)^n n!}{16\pi^2 (4\pi)^n (1+2n)!} \frac{4^{n-1}}{\pi^{2+n} (-x_{\perp}^2)^{2n+1}} \frac{1}{\epsilon} \left(\frac{4^{\epsilon} \Gamma(2n+1+\epsilon)}{\Gamma(1-n-\epsilon)(-x_{\perp}^2)^{\epsilon}} - \frac{\Gamma(2n+1)}{\Gamma(1-n)} \right)$$

$$\frac{1}{\epsilon} \left(\frac{4^{\epsilon} \Gamma(2n+1+\epsilon)}{\Gamma(1-n-\epsilon)(-x_{\perp}^2)^{\epsilon}} - \frac{\Gamma(2n+1)}{\Gamma(1-n)} \right) = \Gamma(2n+1) \left(\frac{\psi(1-n)}{\Gamma(1-n)} + \frac{\psi(2n+1) - \ln\left(-\frac{x_{\perp}^2}{4}\right)}{\Gamma(1-n)} \right)$$

$$\frac{1}{\epsilon} \left(\frac{4^{\epsilon} \Gamma(1+\epsilon)}{\Gamma(1-\epsilon)(-x_{\perp}^2)^{\epsilon}} - 1 \right) = 2\psi(1) - \ln\left(-\frac{x_{\perp}^2}{4}\right)$$

$$\frac{1}{\epsilon} \left(\frac{4^\epsilon \Gamma(1 + \epsilon) r_0^{2\epsilon}}{\Gamma(1 - \epsilon) (-x_\perp^2)^\epsilon} - 1 \right) = -\ln \left(-\frac{x_\perp^2 e^{2\gamma}}{4r_0^2} \right)$$

$$\int \frac{d^5 q \delta(q_\parallel) e^{-iqx}}{(2\pi)^4} N_4 \frac{q_\mu q_\nu}{q^2} = -\frac{1}{32\pi^4 \sqrt{-x_\perp^2}} \left(\eta_{\mu\nu}^\perp - 6 \frac{x_\mu^\perp x_\nu^\perp}{x_\perp^2} - \left(\eta_{\mu\nu}^\perp - 4 \frac{x_\mu^\perp x_\nu^\perp}{x_\perp^2} \right) \ln \left(-\frac{x_\perp^2 e^{2\gamma}}{4r_0^2} \right) \right)$$

$$h_{\mu\nu}^{(2)} = \frac{\mu^2}{r^4} \left(\frac{1}{2} \eta_{\mu\nu}^\parallel - \frac{2(6\alpha - 5) \ln \frac{r}{r_0} - 1}{16} \eta_{\mu\nu}^\perp + \frac{(6\alpha - 5) (4 \ln \frac{r}{r_0} - 1)}{8} \frac{x_\mu^\perp x_\nu^\perp}{x_\perp^2} \right)$$

$$x^\mu \rightarrow x^\mu + \beta \frac{\mu^2}{r^4} x_\perp^\mu + \dots$$

$$h_{\mu\nu}^{G^2} \rightarrow h_{\mu\nu}^{G^2} - \epsilon_{\mu,\nu} - \epsilon_{\nu,\mu}$$

$$G_\sigma^{G^2} \rightarrow G_\sigma^{G^2} + \partial^2 \epsilon_\sigma$$

$$\partial^2 \frac{1}{r^{D-3}} \propto \delta^{D-1}(x_\perp)$$

$$\partial^2 \frac{x_\mu^\perp}{r^{D-1}} \propto \partial_\mu \partial^2 \frac{1}{r^{D-3}} \propto \partial_\mu \delta^{D-1}(x_\perp)$$

$$h_{\mu\nu}^{G^2} \rightarrow h_{\mu\nu}^{G^2} - 2\beta \frac{\mu^2}{r^4} \left(-4 \frac{x_\mu^\perp x_\nu^\perp}{x_\perp^2} + \eta_{\mu\nu}^\perp \right)$$

$$h_{\mu\nu}^{G^2} \rightarrow h_{\mu\nu}^{G^2} + \ln(\gamma) \frac{6\alpha - 5}{8} \frac{\mu^2}{r^4} \left(-4 \frac{x_\mu^\perp x_\nu^\perp}{x_\perp^2} + \eta_{\mu\nu}^\perp \right)$$

$$\epsilon^\sigma = \beta \frac{x_\perp^\sigma}{r^{D-1}}$$

$$-D + 1 = m(-D + 3),$$

$$d\tau^2 = \left(1 - \frac{\mu}{R^n} \right) dt^2 - \frac{1}{1 - \frac{\mu}{R^n}} dR^2 - R^2 d\Omega_{D-2}^2$$

$$x^0 = t$$

$$x^1 = r \cos(\theta_1)$$

$$x^2 = r \sin(\theta_1) \cos(\theta_2)$$

$$x^3 = r \sin(\theta_1) \sin(\theta_2) \cos(\theta_3)$$

...

$$x^{D-1} = r \sin(\theta_1) \sin(\theta_2) \dots \sin(\theta_{D-2})$$

$$dR^2 = \frac{dR^2}{dr^2} \left(\frac{x_\perp dx_\perp}{r} \right)^2$$

$$d\Omega_{D-2}^2 = -\frac{1}{r^2} \left(dx_\perp^2 + \left(\frac{x_\perp dx_\perp}{r} \right)^2 \right)$$



$$d\tau^2 = \left(1 - \frac{\mu}{R^n}\right) dt^2 - \frac{1}{1 - \frac{\mu}{R^n}} \frac{dR^2}{dr^2} \left(\frac{x_\perp dx_\perp}{r}\right)^2 + \frac{R^2}{r^2} \left(dx_\perp^2 + \left(\frac{x_\perp dx_\perp}{r}\right)^2\right)$$

$$g_{\mu\nu} = \left(1 - \frac{\mu}{R^n}\right) \eta_{\mu\nu}^\parallel + \frac{1}{1 - \frac{\mu}{R^n}} \frac{dR^2}{dr^2} \frac{x_\mu^\perp x_\nu^\perp}{x_\perp^2} + \frac{R^2}{r^2} \left(\eta_{\mu\nu}^\perp - \frac{x_\mu^\perp x_\nu^\perp}{x_\perp^2}\right)$$

$$R = r \left(1 + a \frac{\mu}{r^n} + b \left(a \frac{\mu}{r^n}\right)^2 + \dots\right)$$

$$g_{\mu\nu} \approx \eta_{\mu\nu} - \frac{\mu}{r^n} \eta_{\mu\nu}^\parallel + 2a \frac{\mu}{r^n} \eta_{\mu\nu}^\perp - (2na - 1) \frac{\mu}{r^n} \frac{x_\mu^\perp x_\nu^\perp}{x_\perp^2}$$

$$h_{\mu\nu}^G = -\frac{\mu}{r^n} \left(\eta_{\mu\nu}^\parallel - 2a \eta_{\mu\nu}^\perp + (2na - 1) \frac{x_\mu^\perp x_\nu^\perp}{x_\perp^2}\right)$$

$$\mathcal{P}_{\rho\sigma}^{\mu\nu} h_{\mu\nu}^{G,\rho} = 0$$

$$\partial_\sigma h_{\mu\nu}^G = -n \frac{\mu}{r^{n+1}} \frac{x_\sigma^\perp}{r} \left(\eta_{\mu\nu}^\parallel - 2a \eta_{\mu\nu}^\perp + (2na - 1) \frac{x_\mu^\perp x_\nu^\perp}{x_\perp^2}\right) - \frac{1}{r^{n+1}} (2na - 1) \lambda_{\sigma\mu\nu}$$

$$\begin{aligned} \lambda_{\sigma\mu\nu} &= r \partial_\sigma \frac{x_\mu^\perp x_\nu^\perp}{x_\perp^2} \\ &= \left(-\eta_{\sigma\mu}^\perp + \frac{x_\sigma^\perp x_\mu^\perp}{x_\perp^2}\right) \frac{x_\nu^\perp}{r} + \left(-\eta_{\sigma\nu}^\perp + \frac{x_\sigma^\perp x_\nu^\perp}{x_\perp^2}\right) \frac{x_\mu^\perp}{r} \end{aligned}$$

$$G_\sigma^G = \frac{\mu}{r^{n+1}} (2na - 1) \frac{x_\sigma^\perp}{r}$$

$$1 - \frac{\mu}{R^n} = 1 - \frac{\mu}{r^n} + \frac{1}{2} \left(\frac{\mu}{r^n}\right)^2 + \dots$$

$$\frac{R^2}{r^2} = 1 + \frac{1}{n} \frac{\mu}{r^n} + (2b + 1) \frac{1}{4n^2} \left(\frac{\mu}{r^n}\right)^2 + \dots$$

$$\frac{dR^2}{dr^2} = 1 - \frac{n-1}{n} \frac{\mu}{r^n} - (2(2n-1)b - (n-1)^2) \frac{1}{4n^2} \left(\frac{\mu}{r^n}\right)^2 + \dots$$

$$\frac{1}{1 - \frac{\mu}{R^n}} = 1 + \frac{\mu}{r^n} + \frac{1}{2} \left(\frac{\mu}{r^n}\right)^2 + \dots$$

$$\frac{1}{1 - \frac{\mu}{R^n}} \frac{dR^2}{dr^2} = 1 + \frac{1}{n} \frac{\mu}{r^n} + (1 - n(n-2) - 2(2n-1)b) \frac{1}{4n^2} \left(\frac{\mu}{r^n}\right)^2 + \dots$$

$$g_{\mu\nu} \approx \eta_{\mu\nu} - \frac{\mu}{r^n} \left(\eta_{\mu\nu}^\parallel - \frac{1}{n} \eta_{\mu\nu}^\perp\right) + \left(\frac{\mu}{r^n}\right)^2 \left(\frac{1}{2} \eta_{\mu\nu}^\parallel + \frac{2b+1}{4n^2} \eta_{\mu\nu}^\perp - \frac{4b+n-2}{4n} \frac{x_\mu^\perp x_\nu^\perp}{x_\perp^2}\right)$$

$$h_{\mu\nu}^{G^2} = \frac{\mu^2}{r^{2n}} \left(\frac{1}{2} \eta_{\mu\nu}^\parallel + \frac{2b+1}{4n^2} \eta_{\mu\nu}^\perp - \frac{4b+n-2}{4n} \frac{x_\mu^\perp x_\nu^\perp}{x_\perp^2}\right)$$

$$G_\sigma^{G^2} = \mathcal{P}_{\rho\sigma}^{\mu\nu} h_{\mu\nu}^{G^2,\rho} - \alpha \Gamma_{\sigma\alpha\beta}^{\rho\mu\nu} h_G^{\alpha\beta} h_{\mu\nu,\rho}^G$$



$$\partial_\rho h_{\mu\nu}^{G^2} = \frac{2n\mu^2}{r^{2n+1}} \frac{x_\rho^\perp}{r} \left(\frac{1}{2} \eta_{\mu\nu}^\parallel + \frac{2b+1}{4n^2} \eta_{\mu\nu}^\perp - \frac{4b+n-2}{4n} \frac{x_\mu^\perp x_\nu^\perp}{x_\perp^2} \right) - \frac{\mu^2}{r^{2n+1}} \frac{4b+n-2}{4n} \lambda_{\rho\mu\nu}$$

$$\alpha \Gamma_{\sigma\alpha\beta}^{\rho\mu\nu} h_G^{\alpha\beta} h_{\mu\nu,\rho}^G = -\alpha \frac{\mu^2}{r^{2n+1}} \frac{n+1}{2} \frac{x_\sigma^\perp}{r}$$

$$\mathcal{P}_{\rho\sigma}^{\mu\nu} h^{G^2,\rho} = -\frac{\mu^2}{r^{2n+1}} \frac{n^2+1+(n-2)b}{2n} \frac{x_\sigma^\perp}{r}$$

$$G_\sigma^{G^2} = -\frac{\mu^2}{2nr^{2n+1}} \frac{x_\sigma^\perp}{r} (n^2+1+(n-2)b - \alpha n(n+1))$$

$$b = \frac{-(1-\alpha)n^2 + \alpha n - 1}{n-2}$$

$$R = r \left(1 + \frac{\mu}{2nr^n} + \left(b_0 + b_1 \ln \left(\frac{r}{r_0} \right) \right) \left(\frac{\mu}{2nr^n} \right)^2 + \dots \right)$$

$$\frac{1}{1 - \frac{\mu}{R^n}} \frac{dR^2}{dr^2} = (\dots) + 2b_1 \left(\frac{\mu}{2nr^n} \right)^2$$

$$h_{\mu\nu}^{G^2} = (\dots) + 2b_1 \frac{x_\mu^\perp x_\nu^\perp}{x_\perp^2} \left(\frac{\mu}{2nr^n} \right)^2$$

$$\partial_\rho h_{\mu\nu}^{G^2} = (\dots) + b_1 \frac{\mu^2}{4n^2 r^{2n+1}} \left(8n \frac{x_\sigma^\perp x_\mu^\perp x_\nu^\perp}{r x_\perp^2} - 2 \frac{x_\sigma^\perp}{r} \eta_{\mu\nu}^\perp + 2\lambda_{\sigma\mu\nu} \right)$$

$$h_{\mu\nu}^{G^2} = \frac{\mu^2}{r^{2n}} \left(\frac{1}{2} \eta_{\mu\nu}^\parallel + \frac{2b+1}{4n^2} \eta_{\mu\nu}^\perp + \left(-\frac{4b+n-2}{4n} + \frac{b_1}{2n^2} \right) \frac{x_\mu^\perp x_\nu^\perp}{x_\perp^2} \right)$$

$$\mathcal{P}_{\rho\sigma}^{\mu\nu} h^{G^2,\rho} = \frac{\mu^2}{r^{2n+1}} \left(-\frac{n^2+1+(n-2)b}{2n} + \frac{3n-2}{4n^2} b_1 \right) \frac{x_\sigma^\perp}{r}$$

$$G_\sigma^{G^2} = -\frac{\mu^2}{2nr^{2n+1}} \left((n-2)b - \frac{3n-2}{2n} b_1 - \alpha n(n+1) + n^2 + 1 \right)$$

$$b_1 = 5 - 6\alpha$$

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