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SUPERGRAVEDAD CUÁNTICA RELATIVISTA. TEORIZACIÓN INICIAL

**RELATIVISTIC QUANTUM SUPERGRAVITY. INITIAL
THEORIZATION**

Manuel Ignacio Albuja Bustamante
Investigador independiente



Supergravedad cuántica relativista. Teorización inicial

Manuel Ignacio Albuja Bustamante¹

ignaciomanuelaljubustamante@gmail.com

Investigador independiente

RESUMEN

La supergravedad cuántica, en contraposición a la gravedad cuántica, es un estado del espacio – tiempo cuántico, en el que, una partícula oscura o una partícula estrella, deforman hipergeométricamente el referido espacio, provocando agujeros negros cuánticos, a razón de su aniquilación o colapso gravitacional, o en su defecto o simultáneamente, provocando D – dimensiones, en los que es perfectamente posible la transformación y desplazamiento de materia y energía, de un punto a otro, en dimensiones distintas, es decir, por fuera de la dimensión en $\mathbb{R}^4$. En este punto, la dimensión del tiempo es maleable a razón de la supergravedad. La dualidad holográfica, es una de las principales características de este fenómeno, en la medida en que, si bien se tratan de desdoblamiento espaciales, las leyes de la física no son las mismas, lo que al contrario ocurre con la gravedad cuántica. En este artículo, me propongo sentar bases formales de la supergravedad cuántica, la misma que, puede ser endógena o exógena, al tenor de las premisas fundamentales contenidas en la Teoría Cuántica de Campos Relativistas o Curvos, formulada por este autor en trabajos anteriores. Este trabajo, al igual que los anteriores y paralelos, se conciben como un intento de unificación.

Palabras clave: supergravedad cuántica, partícula estrella, partícula oscura, supermembranas, supersimetría de gauge, supersimetría de Yang – Mills

¹ Autor principal

Correspondencia: ignaciomanuelaljubustamante@gmail.com



Relativistic quantum supergravity. Initial theorization

ABSTRACT

Quantum supergravity, as opposed to quantum gravity, is a state of quantum space-time, in which a dark particle or a star particle hypergeometrically deforms the aforementioned space, causing quantum black holes, due to their annihilation or gravitational collapse, or failing that or simultaneously, causing D-dimensions, in which the transformation and displacement of matter and energy is perfectly possible. from one point to another, in different dimensions, that is, outside the dimension in $\mathbb{R}^4$. At this point, the dimension of time is malleable due to supergravity. Holographic duality is one of the main characteristics of this phenomenon, insofar as, although they are spatial splittings, the laws of physics are not the same, which is the opposite of quantum gravity. In this article, I propose to lay the formal foundations of quantum supergravity, which can be endogenous or exogenous, according to the fundamental premises contained in the Quantum Theory of Relativistic or Curved Fields, formulated by this author in previous works. This work, like the previous and parallel ones, is conceived as an attempt at unification.

Keywords: quantum supergravity, star particle, dark particle, supermembranes, gauge supersymmetry, yang-mills supersymmetry.

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INTRODUCCIÓN.

Existe una diferencia sustancial entre gravedad cuántica y supergravedad cuántica, pues la primera (GC) consiste esencialmente en la deformación del espacio – tiempo cuántico, a propósito de las interacciones y simetría fija o corregida de una partícula supermasiva que a razón de su masa y energía, esto es, de su centro de materia – energía, curva geoméricamente el plano cuántico en el que se propaga vectorialmente y lo torsiona o tuerce, cuya matematización es puramente tensorial y geométrica, esto, a raíz de un espacio de Hilbert – Einstein y una superficie de Riemann, sin que por esto, se produzcan pluridimensiones de gauge. Asimismo, la gravedad cuántica (GC), se desarrolla o cumple los parámetros anteriores, cuando una hiperpartícula, con o sin masa, se aproxima, iguala o supera la velocidad de la luz (momentum), por lo que, en cualquiera de los casos anteriores, el resultado es la formación de la curvatura de Dirac del espacio – tiempo cuántico en el que interactúa, lo que sería el caso del taquión, partícula hipotética que cumple esta descripción. En cuanto a la gravedad cuántica se refiere, la dualidad holográfica se tiene por inexistente, lo que implica, que para el espacio – tiempo cuántico deformado geoméricamente, operan las mismas leyes de la física cuántica, pues se trata de la misma dimensión en $\mathbb{R}^4$. Ahora bien, la supergravedad cuántica, en contrario a la gravedad cuántica, comporta las siguientes características: 1) el espacio – tiempo cuántico se deforma hipergeoméricamente; 2) comporta la existencia de D – dimensiones, esto es, supermembranas y por ende, superespacios; 3) comporta la existencia de agujeros negros cuánticos masivos o supermasivos, según sea el caso; provocados por aniquilación entre dos partículas supermasivas, una partícula supermasiva y otra partícula repercutida o por colisión, según las categorías antes referidas; 4) comporta la existencia de supersimetrías de gauge en invariancia o covariancia; 5) comporta la existencia de dualidad holográfica; 6) la deformación del espacio – tiempo cuántico opera por una partícula oscura o una partícula estrella, siendo la primera, aquella que es extremadamente densa pero su energía – momentum es inferior, en tanto que la segunda, es aquella cuya masa es extremadamente densa al igual que su energía – momentum, por lo que, su colapso o aniquilación, genera radiación. Sin perjuicio de lo anterior, ambos tipos de gravedad, surgen de forma endógena o exógena, según corresponda. En este artículo, abordaremos específicamente la supergravedad cuántica.



RESULTADOS Y DISCUSIÓN.

Supergravedad cuántica. Modelo Matemático.

Cálculos preliminares para campos superplanckianos.

$$m^2 = \frac{J}{\alpha'}$$

$$s = -(p_1 + p_2)^2, t = -(p_2 + p_3)^2, u = -(p_1 + p_3)^2$$

$$A_J(s, t) \sim \frac{(-s)^J}{t - M^2}$$

$$A(s, t) = \frac{\Gamma(-\alpha(s))\Gamma(-\alpha(t))}{\Gamma(-\alpha(s) - \alpha(t))}$$

$$\alpha(s) = \alpha(0) + \alpha' s$$

$$A(s, t) = - \sum_{n=0}^{\infty} \frac{(\alpha(s) + 1) \dots (\alpha(s) + n)}{n!} \frac{1}{\alpha(t) - n}$$

$$\frac{1}{M_{\text{Planck}}^4} \int_0^\Lambda dE E^3 \sim \frac{\Lambda^4}{M_{\text{Planck}}^4}$$

Partícula Supermasiva. Comportamiento de campo.

$$S = m \int_{s_i}^{s_f} ds = m \int_{\tau_0}^{\tau_1} d\tau \sqrt{-\eta_{\mu\nu} \dot{x}^\mu \dot{x}^\nu}$$

$$p_\mu = - \frac{\delta L}{\delta \dot{x}^\mu} = \frac{m \dot{x}_\mu}{\sqrt{-\dot{x}^2}}$$

$$\partial_\tau \left(\frac{m \dot{x}_\mu}{\sqrt{-\dot{x}^2}} \right) = 0$$

$$p^2 + m^2 = 0$$

$$H_{can} = \frac{\partial L}{\partial \dot{x}^\mu} \dot{x}^\mu - L$$

$$H = \frac{N}{2m} (p^2 + m^2)$$

$$\dot{x}^\mu = \{x^\mu, H\} = \frac{N}{m} p^\mu = \frac{N \dot{x}^\mu}{\sqrt{-\dot{x}^2}}$$

$$\dot{x}^2 = -N^2$$



$$S = -\frac{1}{2} \int d\tau e(\tau) (e^{-2}(\tau) (\dot{x}^\mu)^2 - m^2)$$

$$S = -\frac{1}{2} \int d\tau \sqrt{\det g_{\tau\tau}} (g^{\tau\tau} \partial_\tau x \cdot \partial_\tau x - m^2)$$

$$\delta x^\mu(\tau) = x^\mu(\tau + \xi(\tau)) - x^\mu(\tau) = \xi(\tau) \dot{x}^\mu + \mathcal{O}(\xi^2)$$

$$\delta S = \frac{1}{2} \int d\tau \left(\frac{1}{e^2(\tau)} (\dot{x}^\mu)^2 + m^2 \right) \delta e(\tau)$$

$$e^{-2}x^2 + m^2 = 0 \rightarrow e = \frac{1}{m} \sqrt{-\dot{x}^2}$$

$$\delta S = \frac{1}{2} \int d\tau e(\tau) (e^{-2}(\tau) 2\dot{x}^\mu) \partial_\tau \delta x^\mu$$

$$\partial_\tau (e^{-1} \dot{x}^\mu) = 0$$

$$\langle x | x' \rangle = N \int_{x(0)=x}^{x(1)=x'} D e D x^\mu \exp \left(\frac{1}{2} \int_0^1 \left(\frac{1}{e} (\dot{x}^\mu)^2 - e m^2 \right) d\tau \right)$$

$$\delta e = \partial_\tau (\xi e)$$

$$L = \int_0^1 d\tau \sqrt{\det g_{\tau\tau}} = \int_0^1 d\tau e$$

$$\langle x | x' \rangle = N \int_0^\infty dL \int_{x(0)=x}^{x(1)=x'} D x^\mu \exp \left(-\frac{1}{2} \int_0^1 \left(\frac{1}{L} \dot{x}^2 + L m^2 \right) d\tau \right)$$

$$x^\mu(\tau) = x^\mu + (x'^\mu - x^\mu)\tau + \delta x^\mu(\tau),$$

$$\|\delta x\|^2 = \int_0^1 d\tau e (\delta x^\mu)^2 = L \int_0^1 d\tau (\delta x^\mu)^2$$

$$Dx^\mu \sim \prod_\tau \sqrt{L} d\delta x^\mu(\tau)$$

$$\langle x | x' \rangle = N \int_0^\infty dL \int \prod_\tau \sqrt{L} d\delta x^\mu(\tau) e^{-\frac{(x'-x)^2}{2L} - m^2 L/2} e^{-\frac{1}{2L} \int_0^1 (\delta x^\mu)^2}$$

$$\int \prod_\tau \sqrt{L} d\delta x^\mu(\tau) e^{-\frac{1}{L} \int_0^1 (\delta x^\mu)^2} \sim \left(\det \left(-\frac{1}{L} \partial_\tau^2 \right) \right)^{\frac{D}{2}}$$

$$-\frac{1}{L} \partial_\tau^2 \psi(\tau) = \lambda \psi(\tau)$$



$$\psi_n(\tau) = C_n \sin(n\pi\tau), \lambda_n = \frac{n^2}{L}, n = 1, 2, \dots$$

$$\det\left(-\frac{1}{L}\partial_t^2\right) = \prod_{n=1}^{\infty} \frac{n^2}{L}$$

$$\prod_{n=1}^{\infty} L^{-1} = L^{-\zeta(0)} = L^{1/2}, \prod_{n=1}^{\infty} n^a = e^{-a\zeta'(0)} = (2\pi)^{a/2}$$

$$\begin{aligned} \langle x | x' \rangle &= \frac{1}{2(2\pi)^{D/2}} \int_0^\infty dL L^{-\frac{D}{2}} e^{-\frac{(x'-x)^2}{2L} - m^2 L/2} = \\ &= \frac{1}{(2\pi)^{D/2}} \left(\frac{|x-x'|}{m} \right)^{(2-D)/2} K_{(D-2)/2}(m|x-x'|) \end{aligned}$$

$$\begin{aligned} |p\rangle &= \int d^D x e^{ip \cdot x} |x\rangle \\ \langle p | p' \rangle &= \int d^D x e^{-ip \cdot x} \int d^D x' e^{ip' \cdot x'} \langle x | x' \rangle \\ &= \frac{1}{2} \int d^D x' e^{i(p'-p) \cdot x'} \int_0^\infty dL e^{-\frac{L}{2}(p^2+m^2)} \\ &= (2\pi)^D \delta(p-p') \frac{1}{p^2+m^2} \end{aligned}$$

Campos cuánticos relativistas o curvos. Propagadores, operadores, números fantasma, supersimetrías, antisimetrías, gravedad supermembranas y superespacios y pluridimensiones en un campo de Hilbert – Einstein. Métrica de Nambu – Goto.

$$S_{NG} = -T \int dA$$

$$ds^2 = G_{\mu\nu}(X) dX^\mu dX^\nu = G_{\mu\nu} \frac{\partial X^\mu}{\partial \xi^i} \frac{\partial X^\nu}{\partial \xi^j} d\xi^i d\xi^j = G_{ij} d\xi^i d\xi^j$$

$$G_{ij} = G_{\mu\nu} \partial_i X^\mu \partial_j X^\nu$$

$$S_{NG} = -T \int \sqrt{-\det G_{ij}} d^2 \xi = -T \int \sqrt{(\dot{X} \cdot X')^2 - (\dot{X}^2)(X'^2)} d^2 \xi$$

$$\partial_\tau \left(\frac{\delta L}{\delta \dot{X}^\mu} \right) + \partial_\sigma \left(\frac{\delta L}{\delta X^\mu} \right) = 0$$

$$X^\mu(\sigma + \bar{\sigma}) = X^\mu(\sigma)$$

- Neumann :

$$\left. \frac{\delta L}{\delta X'^\mu} \right|_{\sigma=0, \bar{\sigma}} = 0$$



- Dirichlet :

$$\left. \frac{\delta L}{\delta \dot{X}^\mu} \right|_{\sigma=0, \bar{\sigma}} = 0$$

$$\Pi^\mu = \frac{\delta L}{\delta \dot{X}^\mu} = -T \frac{(\dot{X} \cdot X') X'^\mu - (X')^2 \dot{X}^\mu}{[(X' \cdot \dot{X})^2 - (\dot{X})^2 (X')^2]^{1/2}}$$

$$\Pi \cdot X' = 0, \Pi^2 + T^2 X'^2 = 0$$

$$H = \int_0^{\bar{\sigma}} d\sigma (\dot{X} \cdot \Pi - L)$$

$$S_P = -\frac{T}{2} \int d^2 \xi \sqrt{-\det g} g^{\alpha\beta} \partial_\alpha X^\mu \partial_\beta X^\nu \eta_{\mu\nu}$$

$$T_{\alpha\beta} \equiv -\frac{2}{T} \frac{1}{\sqrt{-\det g}} \frac{\delta S_P}{\delta g^{\alpha\beta}} = \partial_\alpha X \cdot \partial_\beta X - \frac{1}{2} g_{\alpha\beta} g^{\gamma\delta} \partial_\gamma X \cdot \partial_\delta X$$

$$g_{\alpha\beta} = \partial_\alpha X \cdot \partial_\beta X$$

$$\frac{1}{\sqrt{-\det g}} \partial_\alpha (\sqrt{-\det g} g^{\alpha\beta} \partial_\beta X^\mu) = 0$$

$$\lambda_1 \int \sqrt{-\det g}$$

$$\lambda_2 \int \sqrt{-\det g} R^{(2)}$$

- Poincaré:

$$\delta X^\mu = \omega_\nu^\mu X^\nu + \alpha^\mu, \delta g_{\alpha\beta} = 0$$

$$\delta g_{\alpha\beta} = \xi^\gamma \partial_\gamma g_{\alpha\beta} + \partial_\alpha \xi^\gamma g_{\beta\gamma} + \partial_\beta \xi^\gamma g_{\alpha\gamma} = \nabla_\alpha \xi_\beta + \nabla_\beta \xi_\alpha$$

$$\begin{aligned} \delta X^\mu &= \xi^\alpha \partial_\alpha X^\mu \\ \delta(\sqrt{-\det g}) &= \partial_\alpha (\xi^\alpha \sqrt{-\det g}) \end{aligned}$$

- Weyl:

$$\delta X^\mu = 0, \delta g_{\alpha\beta} = 2\Lambda g_{\alpha\beta}$$

$$\delta g_{\alpha\beta} = 2\Lambda(x) g_{\alpha\beta}, \delta \phi^i = d_i \Lambda(x) \phi^i$$

$$0 = \delta S = \int d^2 \xi \left[2 \frac{\delta S}{\delta g^{\alpha\beta}} g^{\alpha\beta} + \sum_i d_i \frac{\delta S}{\delta \phi_i} \phi_i \right] \Lambda$$



$$T_{\alpha}^{\alpha} \sim \frac{\delta S}{\delta g^{\alpha\beta}} g^{\alpha\beta} = 0$$

$$g_{\alpha\beta} = e^{2\Lambda(\xi)} \eta_{\alpha\beta}$$

$$\xi_+ = \tau + \sigma, \xi_- = \tau - \sigma$$

$$ds^2 = -d\xi_+ d\xi_-$$

$$g_{++} = g_{--} = 0, g_{+-} = g_{-+} = -\frac{1}{2}$$

$$\partial_{\pm} = \frac{1}{2}(\partial_{\tau} \pm \partial_{\sigma})$$

- Polyakov:

$$S_P \sim T \int d^2\xi \partial_+ X^\mu \partial_- X^\nu \eta_{\mu\nu}$$

$$\xi_+ \rightarrow f(\xi_+), \xi_- \rightarrow g(\xi_-)$$

$$\frac{d(d+1)}{2} - d - 1$$

$$\delta S = T \int d^2\xi (\delta X^\mu \partial_+ \partial_- X_\mu) - T \int_{\tau_0}^{\tau_1} d\tau X'_\mu \delta X^\mu$$

$$X'^\mu|_{\sigma=0,\bar{\sigma}} = 0$$

$$\partial_+ \partial_- X^\mu = 0$$

$$T_{\alpha\beta} = 0$$

$$T_{10} = T_{01} = \frac{1}{2} \dot{X} \cdot X' = 0, T_{00} = T_{11} = \frac{1}{4} (\dot{X}^2 + X'^2) = 0$$

$$(\dot{X} \pm X')^2 = 0$$

$$T_{++} = \frac{1}{2} \partial_+ X \cdot \partial_+ X, T_{--} = \frac{1}{2} \partial_- X \cdot \partial_- X, T_{+-} = T_{-+} = 0$$

$$\partial_- T_{++} + \partial_+ T_{-+} = \partial_+ T_{--} + \partial_- T_{+-} = 0$$

$$\partial_- T_{++} = \partial_+ T_{--} = 0$$

$$Q_f = \int_0^{\bar{\sigma}} f(\xi^+) T_{++}(\xi^+)$$

$$0 = \int d\sigma \partial_- (f(\xi^+) T_{++}) = \partial_\tau Q_f + f(\xi^+) T_{++}|_0^{\bar{\sigma}}$$



$$P_\mu^\alpha = -T\sqrt{\det g} g^{\alpha\beta} \partial_\beta X_\mu$$

$$J_{\mu\nu}^\alpha = -T\sqrt{\det g} g^{\alpha\beta} (X_\mu \partial_\beta X_\nu - X_\nu \partial_\beta X_\mu)$$

$$P_\mu = \int_0^{\bar{\sigma}} d\sigma P_\mu^\tau, J_{\mu\nu} = \int_0^{\bar{\sigma}} d\sigma J_{\mu\nu}^\tau$$

$$\frac{\partial P_\mu}{\partial \tau} = T \int_0^{\bar{\sigma}} d\sigma \partial_\tau^2 X_\mu = T \int_0^{\bar{\sigma}} d\sigma \partial_\sigma^2 X_\mu = T(\partial_\sigma X_\mu(\sigma = \bar{\sigma}) - \partial_\sigma X_\mu(\sigma = 0))$$

Expansiones oscilatorias.

$$\partial_+ \partial_- X^\mu = 0$$

$$X^\mu(\tau, \sigma + 2\pi) = X^\mu(\tau, \sigma)$$

$$X^\mu(\tau, \sigma) = X_L^\mu(\tau + \sigma) + X_R^\mu(\tau - \sigma)$$

$$X_L^\mu(\tau + \sigma) = \frac{x^\mu}{2} + \frac{p^\mu}{4\pi T}(\tau + \sigma) + \frac{i}{\sqrt{4\pi T}} \sum_{k \neq 0} \frac{\bar{\alpha}_k^\mu}{k} e^{-ik(\tau + \sigma)}$$

$$X_R^\mu(\tau - \sigma) = \frac{x^\mu}{2} + \frac{p^\mu}{4\pi T}(\tau - \sigma) + \frac{i}{\sqrt{4\pi T}} \sum_{k \neq 0} \frac{\alpha_k^\mu}{k} e^{-ik(\tau - \sigma)}$$

$$(\alpha_k^\mu)^* = \alpha_{-k}^\mu \text{ and } (\bar{\alpha}_k^\mu)^* = \bar{\alpha}_{-k}^\mu$$

$$\partial_- X_R^\mu = \frac{1}{\sqrt{4\pi T}} \sum_{k \in \mathbb{Z}} \alpha_k^\mu e^{-ik(\tau - \sigma)},$$

$$\partial_+ X_L^\mu = \frac{1}{\sqrt{4\pi T}} \sum_{k \in \mathbb{Z}} \bar{\alpha}_k^\mu e^{-ik(\tau + \sigma)}.$$

$$X'^\mu(\tau, \sigma)|_{\sigma=0, \pi} = 0$$

$$X'^\mu|_{\sigma=0} = \frac{p^\mu - \bar{p}^\mu}{\sqrt{4\pi T}} + \frac{1}{\sqrt{4\pi T}} \sum_{k \neq 0} e^{ik\tau} (\bar{\alpha}_k^\mu - \alpha_k^\mu)$$

$$p^\mu = \bar{p}^\mu \text{ and } \alpha_k^\mu = \bar{\alpha}_k^\mu$$

$$X^\mu(\tau, \sigma) = x^\mu + \frac{p^\mu \tau}{\pi T} + \frac{i}{\sqrt{\pi T}} \sum_{k \neq 0} \frac{\alpha_k^\mu}{k} e^{-ik\tau} \cos(k\sigma)$$

$$\partial_\pm X^\mu = \frac{1}{\sqrt{4\pi T}} \sum_{k \in \mathbb{Z}} \alpha_k^\mu e^{-ik(\tau \pm \sigma)}$$

$$X_{CM}^\mu \equiv \frac{1}{\bar{\sigma}} \int_0^{\bar{\sigma}} d\sigma X^\mu(\tau, \sigma) = x^\mu + \frac{p^\mu \tau}{\pi T}$$



$$p_{CM}^\mu = T \int_0^{\bar{\sigma}} d\sigma \dot{X}^\mu = \frac{T}{\sqrt{4\pi T}} \int d\sigma \sum_k (\alpha_k^\mu + \bar{\alpha}_k^\mu) e^{-ik(\tau \pm \sigma)} = \frac{2\pi T}{\sqrt{4\pi T}} (\alpha_0^\mu + \bar{\alpha}_0^\mu) = p^\mu$$

$$J^{\mu\nu} = T \int_0^{\bar{\sigma}} d\sigma (X^\mu \dot{X}^\nu - X^\nu \dot{X}^\mu) = l^{\mu\nu} + E^{\mu\nu} + \bar{E}^{\mu\nu}$$

$$l^{\mu\nu} = x^\mu p^\nu - x^\nu p^\mu$$

$$E^{\mu\nu} = -i \sum_{n=1}^{\infty} \frac{1}{n} (\alpha_{-n}^\mu \alpha_n^\nu - \alpha_{-n}^\nu \alpha_n^\mu)$$

$$\bar{E}^{\mu\nu} = -i \sum_{n=1}^{\infty} \frac{1}{n} (\bar{\alpha}_{-n}^\mu \bar{\alpha}_n^\nu - \bar{\alpha}_{-n}^\nu \bar{\alpha}_n^\mu)$$

$$\{X^\mu(\sigma, \tau), \dot{X}^\nu(\sigma', \tau)\}_{PB} = \frac{1}{T} \delta(\sigma - \sigma') \eta^{\mu\nu}$$

$$\begin{aligned} \{\alpha_m^\mu, \alpha_n^\nu\} &= \{\bar{\alpha}_m^\mu, \bar{\alpha}_n^\nu\} = -im \delta_{m+n,0} \eta^{\mu\nu} \\ \{\bar{\alpha}_m^\mu, \alpha_n^\nu\} &= 0, \{x^\mu, p^\nu\} = \eta^{\mu\nu}. \end{aligned}$$

$$H = \int d\sigma (\dot{X}\Pi - L) = \frac{T}{2} \int d\sigma (\dot{X}^2 + X'^2)$$

$$H = \frac{1}{2} \sum_{n \in \mathbb{Z}} (\alpha_{-n} \alpha_n + \bar{\alpha}_{-n} \bar{\alpha}_n)$$

$$H = \frac{1}{2} \sum_{n \in \mathbb{Z}} \alpha_{-n} \alpha_n$$

$$L_m = 2T \int_0^{2\pi} d\sigma T_{--} e^{im(\tau-\sigma)}, \bar{L}_m = 2T \int_0^{2\pi} d\sigma T_{++} e^{im(\sigma+\tau)}$$

$$L_m = \frac{1}{2} \sum_n \alpha_{m-n} \alpha_n, \bar{L}_m = \frac{1}{2} \sum_n \bar{\alpha}_{m-n} \bar{\alpha}_n$$

$$L_m^* = L_{-m} \text{ and } \bar{L}_m^* = \bar{L}_{-m}$$

$$H = L_0 + \bar{L}_0.$$

$$L_m = 2T \int_0^\pi d\sigma \{T_{--} e^{im(\tau-\sigma)} + T_{++} e^{im(\sigma+\tau)}\}$$

$$L_m = \frac{1}{2} \sum_n \alpha_{m-n} \alpha_n$$

$$H = L_0$$



$$\begin{aligned}\{L_m, L_n\}_{PB} &= -i(m-n)L_{m+n} \\ \{\bar{L}_m, \bar{L}_n\}_{PB} &= -i(m-n)\bar{L}_{m+n} \\ \{L_m, \bar{L}_n\}_{PB} &= 0\end{aligned}$$

$$\{, \}_{PB} \rightarrow -i[,].$$

Cuantización canónica- Bosonificación.

$$\begin{aligned}[x^\mu, p^\nu] &= i\eta^{\mu\nu} \\ [\alpha_m^\mu, \alpha_n^\nu] &= m\delta_{m+n,0}\eta^{\mu\nu}\end{aligned}$$

$$[\alpha_m^\mu, \alpha_n^{\nu\dagger}] = \delta_{m,n}\eta^{\mu\nu}$$

$$\alpha_m|p\rangle = 0 \quad \forall m > 0.$$

$$|p\rangle, \alpha_{-1}^\mu|p^\mu\rangle, \alpha_{-1}^\mu\alpha_{-1}^\nu\alpha_{-2}^\nu|p^\mu\rangle, \text{ etc.}$$

$$|\alpha_{-1}^0|p\rangle\Box|^2 = \langle p|\alpha_1^0\alpha_{-1}^0|p\rangle = -1,$$

$$L_m = \frac{1}{2} \sum_{n \in \mathbb{Z}} : \alpha_{m-n} \cdot \alpha_n :$$

$$L_0 = \frac{1}{2} \alpha_0^2 + \sum_{n=1}^{\infty} \alpha_{-n} \cdot \alpha_n$$

$$[L_m, L_n] = (m-n)L_{m+n} + \frac{c}{12} m(m^2-1)\delta_{m+n,0}$$

$$0 = \langle \phi | [L_m, L_{-m}] | \phi \rangle = 2m \langle \phi | L_0 | \phi \rangle + \frac{d}{12} m(m^2-1) \langle \phi | \phi \rangle \neq 0$$

$$L_{m>0} | \text{phys} \rangle = 0, (L_0 - a) | \text{phys} \rangle = 0$$

$$\alpha' m^2 = 4(N - a)$$

$$N = \sum_{m=1}^{\infty} \alpha_{-m} \cdot \alpha_m$$

$$\xi'_+ = f(\xi_+), \xi'_- = g(\xi_-)$$

$$X^+ = x^+ + \alpha' p^+ \tau$$

$$X^\pm = X^0 \pm X^1$$

$$\alpha_n^+ = \bar{\alpha}_n^+ = \sqrt{\frac{\alpha'}{2}} p^+ \delta_{n,0}$$

$$\alpha_n^- = \frac{1}{\sqrt{2\alpha'p^+}} \left\{ \sum_{m \in \mathbb{Z}} : \alpha_{n-m}^i \alpha_m^i : - 2a \delta_{n,0} \right\}$$



$$\alpha_{-1}^i \bar{\alpha}_{-1}^j |p\rangle$$

$$\alpha_{-1}^i \bar{\alpha}_{-1}^j |p\rangle = \alpha_{-1}^{[i} \bar{\alpha}_{-1}^{j]} |p\rangle + \left[\alpha_{-1}^{\{i} \bar{\alpha}_{-1}^{j\}} - \frac{1}{d-2} \delta^{ij} \alpha_{-1}^k \bar{\alpha}_{-1}^k \right] |p\rangle + \\ + \frac{1}{d-2} \delta^{ij} \alpha_{-1}^k \bar{\alpha}_{-1}^k |p\rangle$$

$$\alpha' m^2 = 4(1-a)$$

$$\alpha_{-1}^i |p\rangle,$$

$$\alpha_{-2}^i |p\rangle, \alpha_{-1}^i \alpha_{-1}^j |p\rangle,$$

$$j^{\max} = \alpha' m^2 + 1$$

$$\Omega|p,i,j\rangle=\epsilon|p,j,i\rangle$$

$$Z=\int \frac{\mathcal{D}g\mathcal{D}X^\mu}{V_{\rm gauge}}e^{iS_p(g,X^\mu)}$$

$$\|\delta g\| = \int d^2\sigma \sqrt{g} g^{\alpha\beta} g^{\delta\gamma} \delta g_{\alpha\gamma} \delta g_{\beta\delta}$$

$$\|\delta X^\mu\| = \int d^\sigma \sqrt{g} \delta X^\mu \delta X^\nu \eta_{\mu\nu}$$

$$g_{\alpha\beta}=e^{2\phi}h_{\alpha\beta}$$

$$\delta g_{\alpha\beta} = \nabla_\alpha \xi_\beta + \nabla_\beta \xi_\alpha + 2\Lambda g_{\alpha\beta} = (\hat{P}\xi)_{\alpha\beta} + 2\tilde{\Lambda} g_{\alpha\beta}$$

$$\mathcal{D}g = \mathcal{D}(\hat{P}\xi)\mathcal{D}(\tilde{\Lambda}) = \mathcal{D}\xi\mathcal{D}\Lambda \left| \frac{\partial(P\xi, \tilde{\Lambda})}{\partial(\xi, \Lambda)} \right|,$$

$$\left|\frac{\partial(P\xi,\tilde{\Lambda})}{\partial(\xi,\Lambda)}\right|=\left|\det\begin{pmatrix}\hat{P}&0*_&1\end{pmatrix}\right|=|\det P|=\sqrt{\det\hat{P}\hat{P}^\dagger}$$

$$Z=\int \mathcal{D}X^\mu \sqrt{\det PP^\dagger} e^{iS_p(\hat{h}_{\alpha\beta},X^\mu)}$$

- Faddeev-Popov:

$$\sqrt{\det PP^\dagger} = \int \mathcal{D}c\mathcal{D}b e^{i\int d^2\sigma \sqrt{g} g^{\alpha\beta} b_{\alpha\gamma} \nabla_\beta c^\alpha}$$

$$Z = \int \mathcal{D}X \mathcal{D}c \mathcal{D}b e^{i(S_p[X] + S_{gh}[c,b])}$$

$$S_p[X] \, = \, T \int \, d^2\sigma \partial_+ X^\mu \partial_- X_\mu$$

$$S_{gh}[b,c] \, = \, \int \, b_{++} \partial_- c^+ + b_{--} \partial_+ c^-$$



Geometría topológica de un campo de gauge. Supergeometría e hipergeometrización tensorial.

$$\delta g_{\alpha\beta} = \nabla_\alpha \xi_\beta + \nabla_\beta \xi_\alpha + 2\Lambda g_{\alpha\beta} = (\hat{P}\xi)_{\alpha\beta} + 2\tilde{\Lambda}g_{\alpha\beta}$$

$$\hat{P}\xi^* = 0$$

$$(V_\alpha, W_\alpha) = \int d^2\xi \sqrt{\det g} g^{\alpha\beta} V_\alpha W_\beta$$

$$(T_{\alpha\beta}, S_{\alpha\beta}) = \int d^2\xi \sqrt{\det g} g^{\alpha\gamma} g^{\beta\delta} T_{\alpha\beta} S_{\gamma\delta}$$

$$(\hat{P}^\dagger t)_\alpha = -2\nabla^\beta t_{\alpha\beta}.$$

$$\hat{P}^\dagger t^* = 0$$

$$[\delta_\alpha, \delta_\beta] = f_{\alpha\beta}^\gamma \delta_\gamma$$

$$F^A(\phi_i)=0$$

$$\int \frac{\mathcal{D}\phi}{V_{\text{gauge}}} e^{-S_0} \sim \int \mathcal{D}\phi \delta(F^A(\phi)=0) \mathcal{D}b_A \mathcal{D}c^\alpha e^{-S_0 - \int b_A (\delta_\alpha F^A) c^\alpha}$$

$$\sim \int \mathcal{D}\phi \mathcal{D}B_A \mathcal{D}b_A \mathcal{D}c^\alpha e^{-S_0 - i \int B_A F^A(\phi) - \int b_A (\delta_\alpha F^A) c^\alpha} = \int \mathcal{D}\phi \mathcal{D}B_A \mathcal{D}b_A \mathcal{D}c^\alpha e^{-S}$$

$$S=S_0+S_1+S_2, S_1=i\int B_A F^A(\phi), S_2=\int b_A (\delta_\alpha F^A) c^\alpha$$

$$\begin{aligned}\delta_{BRST}\phi_i &= -i\epsilon c^\alpha \delta_\alpha \phi_i \\ \delta_{BRST}b_A &= -\epsilon B_A \\ \delta_{BRST}c^\alpha &= -\frac{1}{2}\epsilon c^\beta c^\gamma f_{\beta\gamma}^\alpha \\ \delta_{BRST}B_A &= 0\end{aligned}$$

$$\delta_{BRST}(b_AF^A)=\epsilon[B_AF^A(\phi)+b_Ac^\alpha\delta_\alpha F^A(\phi)].$$

$$\epsilon\delta_F\langle\psi\mid\psi'\rangle=-i\langle\psi\mid\delta_{BRST}(b_A\delta F^A)\mid\psi'\rangle=\langle\psi\mid\{Q_B,b_A\delta F^A\}\mid\psi'\rangle,$$

$$Q_B\mid\text{phys}\rangle=0$$

$$\begin{aligned}0 &= [Q_B, \delta H] = [Q_B, \delta_B (b_A \delta F^A)] \\ &= [Q_B, \{Q_B, b_A \delta F^A\}] = [Q_B^2, b_A \delta F^A]\end{aligned}$$

$$Q_B^2=0$$

$$|\psi'\rangle=|\psi\rangle+Q_B|\chi\rangle$$

$$\begin{aligned}Q_B\mid\text{phys}\rangle &= 0 \\ \text{and } \mid\text{phys}\rangle &\neq Q_B\mid\text{something}\rangle.\end{aligned}$$



$$\begin{aligned}\delta_B X^\mu &= i\epsilon(c^+ \partial_+ + c^- \partial_-) X^\mu \\ \delta_B c^\pm &= \pm i\epsilon(c^+ \partial_+ + c^- \partial_-) c^\pm \\ \delta_B b_\pm &= \pm i\epsilon(T_\pm^X + T_\pm^{gh})\end{aligned}$$

$$S_{gh} = \int d^2\sigma (b_{++} \partial_- c^+ + b_{--} \partial_+ c^-)$$

$$\begin{aligned}T_{++}^{gh} &= i(2b_{++} \partial_+ c^+ + \partial_+ b_{++} c^+) \\ T_{--}^{gh} &= i(2b_{--} \partial_- c^- + \partial_- b_{--} c^-)\end{aligned}$$

$$\partial_- T_{++}^{gh} = \partial_+ T_{--}^{gh} = 0$$

$$\partial_- b_{++} = \partial_+ b_{--} = \partial_- c^+ = \partial_+ c^- = 0$$

$$\begin{aligned}c^+ &= \sum \bar{c}_n e^{-in(\tau+\sigma)}, c^- = \sum c_n e^{-in(\tau-\sigma)} \\ b_{++} &= \sum \bar{b}_n e^{-in(\tau+\sigma)}, b_{--} = \sum c_n e^{-in(\tau-\sigma)}\end{aligned}$$

$$\{b_m, c_n\} = \delta_{m+n,0}, \{b_m, b_n\} = \{c_m, c_n\} = 0$$

$$L_m^{gh} = \sum_n (m-n):b_{m+n}c_{-n}:, \bar{L}_m^{gh} = \sum_n (m-n):\bar{b}_{m+n}\bar{c}_{-n}:$$

- Virasoro: Campos Fantasma.

$$[L_m^{gh}, L_n^{gh}] = (m-n)L_{m+n}^{gh} + \frac{1}{6}(m-13m^3)\delta_{m+n,0}$$

$$L_m = L_m^X + L_m^{gh} - a\delta_m$$

$$[L_m, L_n] = (m-n)L_{m+n} + A(m)\delta_{m+n}$$

$$A(m) = \frac{d}{12}m(m^2-1) + \frac{1}{6}(m-13m^3) + 2am$$

$$j_B = cT^X + \frac{1}{2}:cT^{gh}: = cT^X + :bc\partial c: ,$$

$$Q_B = \int d\sigma j_B$$

$$Q_B = \sum_n c_n L_{-n}^X + \sum_{m,n} \frac{m-n}{2}:c_m c_n b_{-m-n}: - c_0$$

$$b_0|\text{phys}\rangle = 0$$

$$b_{n>0} | \text{ghost vacuum} \rangle = c_{n>0} | \text{ghost vacuum} \rangle = 0.$$

$$\begin{aligned}b_0|\downarrow\rangle &= 0, & b_0|\uparrow\rangle &= |\downarrow\rangle, \\ c_0|\uparrow\rangle &= 0, & c_0|\downarrow\rangle &= |\uparrow\rangle.\end{aligned}$$



$$0 = Q_B |\downarrow, p\rangle = (L_0^X - 1) c_0 |\downarrow, p\rangle.$$

$$|\psi\rangle = (\zeta \cdot \alpha_{-1} + \xi_1 c_{-1} + \xi_2 b_{-1}) |\downarrow, p\rangle,$$

$$0 = Q_B |\psi\rangle = 2(p^2 c_0 + (p \cdot \zeta) c_{-1} + \xi_1 p \cdot \alpha_{-1}) |\downarrow, p\rangle.$$

$$Q_B |\chi\rangle = 2(p \cdot \zeta' c_{-1} + \xi_1' p \cdot \alpha_{-1}) |\downarrow, p\rangle.$$

Interacciones y amplitudes.

$$g_{\mu\nu} \rightarrow g'_{\mu\nu}(x') = \frac{\partial x^\alpha}{\partial x'^\mu} \frac{\partial x^\beta}{\partial x'^\nu} g_{\alpha\beta}(x)$$

$$g_{\mu\nu}(x) \rightarrow g'_{\mu\nu}(x') = \Omega(x) g_{\mu\nu}(x)$$

$$ds'^2 = ds^2 - (\partial_\mu \epsilon_\nu + \partial_\nu \epsilon_\mu) dx^\mu dx^\nu$$

$$\partial_\mu \epsilon_\nu + \partial_\nu \epsilon_\mu = \frac{2}{d} (\partial \cdot \epsilon) \eta_{\mu\nu}$$

$$\square \epsilon_\nu + \left(1 - \frac{2}{d}\right) \partial_\nu (\partial \cdot \epsilon) = 0$$

$$\partial_\mu \square \epsilon_\nu + \partial_\nu \square \epsilon_\mu = \frac{2}{d} \eta_{\mu\nu} \square (\partial \cdot \epsilon)$$

$$[\eta_{\mu\nu} \square + (d-2) \partial_\mu \partial_\nu] \partial \cdot \epsilon = 0$$

$$\epsilon^\mu = a^\mu \text{translationes}$$

$$\epsilon^\mu = \omega^\mu_\nu x^\nu \text{ rotations } (\omega_{\mu\nu} = -\omega_{\nu\mu}),$$

$$\epsilon^\mu = \lambda x^\mu \text{ transformationes escalares}$$

$$\epsilon^\mu = b^\mu x^2 - 2x^\mu (b \cdot x)$$

$$d + \frac{1}{2}d(d-1) + 1 + d = \frac{1}{2}(d+2)(d+1)$$

$$\partial_1 \epsilon_1 = \partial_2 \epsilon_2, \partial_1 \epsilon_2 = -\partial_2 \epsilon_1$$

$$\partial \bar{\epsilon} = 0, \bar{\partial} \epsilon = 0$$

$$z \rightarrow f(z) \text{ and } \bar{z} \rightarrow \bar{f}(\bar{z})$$

$$\epsilon(z) = - \sum a_n z^{n+1}$$

$$\ell_n = -z^{n+1} \partial_z$$

$$[\ell_m, \ell_n] = (m-n) \ell_{m+n}, [\bar{\ell}_m, \bar{\ell}_n] = (m-n) \bar{\ell}_{m+n}$$

$$z \rightarrow \frac{az+b}{cz+d}$$



Invariancia y covariancia de gauge.

$$z \rightarrow f(z) = \frac{az + b}{cz + d}, \bar{z} \rightarrow \bar{f}(\bar{z}) = \frac{\bar{a}\bar{z} + \bar{b}}{\bar{c}\bar{z} + \bar{d}}$$

$$\Phi(z, \bar{z}) \rightarrow \left(\frac{\partial f}{\partial z} \right)^h \left(\frac{\partial \bar{f}}{\partial \bar{z}} \right)^{\bar{h}} \Phi(f(z), \bar{f}(\bar{z}))$$

$$\Phi(z, \bar{z}) dz^h d\bar{z}^{\bar{h}}$$

$$\left\langle \prod_{i=1}^N \Phi_i(z_i, \bar{z}_i) \right\rangle = \prod_{i=1}^N \left(\frac{\partial f}{\partial z} \right)_{z \rightarrow z_i}^{h_i} \left(\frac{\partial \bar{f}}{\partial \bar{z}} \right)_{\bar{z} \rightarrow \bar{z}_i}^{\bar{h}_i} \left\langle \prod_{j=1}^N \Phi_j(f(z_j), \bar{f}(\bar{z}_j)) \right\rangle.$$

$$\delta_{\epsilon, \bar{\epsilon}} \Phi(z, \bar{z}) = [(h\partial\epsilon + \epsilon\partial) + (\bar{h}\bar{\partial}\bar{\epsilon} + \bar{\epsilon}\bar{\partial})] \Phi(z, \bar{z})$$

$$\delta_{\epsilon, \bar{\epsilon}} G^{(2)}(z_i, \bar{z}_i) = \langle \delta_{\epsilon, \bar{\epsilon}} \Phi_1, \Phi_2 \rangle + \langle \Phi_1, \delta_{\epsilon, \bar{\epsilon}} \Phi_2 \rangle = 0$$

$$\left[\left(\epsilon(z_1) \partial_{z_1} + h_1 \partial \epsilon(z_1) + \epsilon(z_2) \partial_{z_2} + h_2 \partial \epsilon(z_2) \right) + (\text{ barred terms }) \right] G^{(2)}(z_i, \bar{z}_i) = 0$$

$$G^{(2)}(z_i, \bar{z}_i) = \frac{C_{12}}{z_{12}^{2h} \bar{z}_{12}^{2\bar{h}}}$$

$$G^{(3)}(z_i, \bar{z}_i) = \frac{C_{123}}{z_{12}^{\Delta_{12}} z_{23}^{\Delta_{23}} z_{31}^{\Delta_{31}} \bar{z}_{12}^{\bar{\Delta}_{12}} \bar{z}_{12}^{\bar{\Delta}_{12}} \bar{z}_{12}^{\bar{\Delta}_{12}}}$$

$$G^{(4)}(z_i, \bar{z}_i) = f(x, \bar{x}) \prod_{i < j} z_{ij}^{-(h_i + h_j) + h/3} \prod_{i < j} \bar{z}_{ij}^{-(\bar{h}_i + \bar{h}_j) + \bar{h}/3}$$

$$G^N(z_1, \bar{z}_1, \dots, z_N, \bar{z}_N) = \left\langle \prod_{i=1}^N \Phi_i(z_i, \bar{z}_i) \right\rangle$$

$$\sum_{i=1}^N \partial_i G^N = 0$$

$$\sum_{i=1}^N (z_i \partial_i + h_i) G^N = 0$$

$$\sum_{i=1}^N (z_i^2 \partial_i + 2z_i h_i) G^N = 0$$



Cuantización radial.

$$z = e^{\tau+i\sigma}, \bar{z} = e^{\tau-i\sigma}$$

$$H = \ell_0 + \bar{\ell}_0$$

$$T_{\mu}{}^{\mu} = 0$$

$$T_{z\bar{z}} = T_{\bar{z}z} = \frac{1}{4}(T_{00} + T_{11}) = \frac{1}{4}T_{\mu}{}^{\mu}$$

$$\partial_z T_{\bar{z}\bar{z}} = 0 \text{ and } \partial_{\bar{z}} T_{zz} = 0$$

$$T(z) \equiv T_{zz} \text{ and } \bar{T}(\bar{z}) \equiv T_{\bar{z}\bar{z}}$$

$$Q_{\epsilon} = \frac{1}{2\pi i} \oint dz \epsilon(z) T(z), Q_{\bar{\epsilon}} = \frac{1}{2\pi i} \oint d\bar{z} \bar{\epsilon}(\bar{z}) \bar{T}(\bar{z}).$$

$$z \rightarrow z + \epsilon(z), \bar{z} \rightarrow \bar{z} + \bar{\epsilon}(\bar{z})$$

$$\delta_{\epsilon, \bar{\epsilon}} \Phi(z, \bar{z}) = [Q_{\epsilon} + Q_{\bar{\epsilon}}, \Phi(z, \bar{z})]$$

$$R(A(z)B(w)) = \begin{cases} A(z)B(w) & |z| > |w| \\ (-1)^F B(w)A(z) & |z| < |w| \end{cases}$$

$$\left[\int d\sigma B, A \right] = \oint dz R(B(z)A(w))$$

$$\begin{aligned} \delta_{\epsilon, \bar{\epsilon}} \Phi(z, \bar{z}) &= \frac{1}{2\pi i} \oint (dz \epsilon(z) R(T(z) \Phi(w, \bar{w})) + d\bar{z} \bar{\epsilon}(\bar{z}) R(\bar{T}(\bar{z}) \Phi(w, \bar{w}))) \\ &= [(h\partial\epsilon(w) + \epsilon(w)\partial) + (\bar{h}\bar{\partial}\bar{\epsilon}(\bar{w}) + \bar{\epsilon}(\bar{w})\bar{\partial})] \Phi(w, \bar{w}) \end{aligned}$$

$$R(T(z)\Phi(w, \bar{w})) = \frac{h}{(z-w)^2} \Phi(w, \bar{w}) + \frac{1}{z-w} \partial_w \Phi(w, \bar{w}) + \dots$$

$$R(\bar{T}(\bar{z})\Phi(w, \bar{w})) = \frac{\bar{h}}{(\bar{z}-\bar{w})^2} \Phi(w, \bar{w}) + \frac{1}{\bar{z}-\bar{w}} \partial_{\bar{w}} \Phi(w, \bar{w}) + \dots$$

$$F^N(z, z_i, \bar{z}_i) = \left\langle T(z) \prod_{i=1}^N \Phi_i(z_i, \bar{z}_i) \right\rangle,$$

$$F^N(z, z_i, \bar{z}_i) = \sum_{i=1}^N \left(\frac{h_i}{(z-z_i)^2} + \frac{\partial_{z_i}}{z-z_i} \right) \left\langle \prod_{i=1}^N \Phi_i(z_i, \bar{z}_i) \right\rangle.$$

$$\Phi_i(z, \bar{z}) \Phi_j(w, \bar{w}) = \sum_k C_{ijk} (z-w)^{h_k-h_i-h_j} (\bar{z}-\bar{w})^{\bar{h}_k-\bar{h}_i-\bar{h}_j} \Phi_k(w, \bar{w})$$



Bosón libre.

$$S = \frac{1}{4\pi} \int d^2z \partial X \bar{\partial} X$$

$$\langle X(z, \bar{z}) X(w, \bar{w}) \rangle = -\log(|z - w|^2 \mu^2)$$

$$\partial_z X(z) \partial_w X(w) = \partial_z \partial_w \langle XX \rangle + : \partial_z X \partial_w X : = -\frac{1}{(z - w)^2} + : \partial_z X \partial_w X :$$

$$T(z) = -\frac{1}{2} : \partial X \partial X : = -\frac{1}{2} \lim_{z \rightarrow w} \left[\partial_z X \partial_w X + \frac{1}{(z - w)^2} \right]$$

$$\bar{T}(\bar{z}) = -\frac{1}{2} : \bar{\partial} X \bar{\partial} X : = -\frac{1}{2} \lim_{\bar{z} \rightarrow \bar{w}} \left[\partial_{\bar{z}} X \partial_{\bar{w}} X + \frac{1}{(\bar{z} - \bar{w})^2} \right]$$

• Wick:

$$T(z) \partial X(w) = -\frac{1}{2} : \partial X(z) \partial X(z) : \partial X(w) = -\partial X(z) \langle \partial X(z) \partial X(w) \rangle + \dots = \partial X(z) \frac{1}{(z - w)^2} + \dots$$

$$= \frac{\partial X(w)}{(z - w)^2} + \frac{1}{z - w} \partial^2 X(w) + \dots$$

$$T(z) V_a(w, \bar{w}) = -\frac{1}{2} : \partial X(z) \partial X(z) : \sum_{n=0}^{\infty} \frac{i^n a^n}{n!} : X^n(w, \bar{w}) :$$

$$T(z) V_a(w) = -\frac{1}{2} [ia \partial \langle XX \rangle]^2 e^{iaX(w)} - \frac{1}{2} 2ia : \partial X(z) \partial \langle XX \rangle e^{iaX(w)} : + \dots$$

$$= \frac{a^2/2}{(z - w)^2} e^{iaX(w)} + \frac{ia \partial X(z)}{z - w} e^{iaX(w)} + \dots = \frac{a^2/2}{(z - w)^2} V_a(w) + \frac{1}{z - w} \partial V_a(w) + \dots$$

$$G^N = \left\langle \prod_{i=1}^N V_{a_i}(z_i, \bar{z}_i) \right\rangle = \exp \left[\frac{1}{2} \sum_{i,j=1; i \neq j}^N a_i a_j \langle X(z_i, \bar{z}_i) X(z_j, \bar{z}_j) \rangle \right]$$

$$\sum_i a_i = 0$$

$$\langle V_a(z) V_{-a}(w) \rangle = \langle : e^{iaX(z)} : : e^{-iaX(w)} : \rangle = e^{-a^2 \log |z - w|^2} = \frac{1}{|z - w|^{2a^2}}$$

$$i \partial_z X V_a(w, \bar{w}) = a \frac{V_a(w, \bar{w})}{(z - w)} + \text{finito}$$



Carga central – OPE.

$$T(z)T(w) = \frac{c/2}{(z-w)^4} + 2 \frac{T(w)}{(z-w)^2} + \frac{\partial T(w)}{z-w} + \dots$$

$$T(z)\bar{T}(\bar{w}) = \text{regular}$$

$$\begin{aligned} T(z)T(w) &= \frac{1}{4} \{ 2(\partial\partial\langle XX \rangle)^2 + 4: \partial X(z)\partial X(w): \partial\partial\langle XX \rangle + \dots \} \\ &= \frac{1/2}{(z-w)^4} + \frac{2}{(z-w)^2} T(w) + \frac{1}{z-w} \partial T(w) + \dots \end{aligned}$$

$$\tilde{T} = -\frac{1}{2} : \partial X \partial X : + iQ \partial^2 X$$

Fermión libre.

$$c = 1 - 12Q^2$$

$$\partial = \sigma^1 \partial_1 + \sigma^2 \partial_2 = \begin{pmatrix} 0 & \partial_1 - i\partial_2 \\ \partial_1 + i\partial_2 & 0 \end{pmatrix} \sim \begin{pmatrix} 0 & \partial \\ \bar{\partial} & 0 \end{pmatrix}.$$

- Majorana espinor $\begin{pmatrix} \psi \\ \bar{\psi} \end{pmatrix}$:

$$S = -\frac{1}{8\pi} \int d^2z (\psi \bar{\partial} \psi + \bar{\psi} \partial \bar{\psi})$$

$$\bar{\partial} \psi = \partial \bar{\psi} = 0$$

$$\psi(z)\psi(w) = \frac{1}{z-w}, \bar{\psi}(\bar{z})\bar{\psi}(\bar{w}) = \frac{1}{\bar{z}-\bar{w}}$$

$$T(z) = -\frac{1}{2} : \psi(z) \partial \psi(z) : , \bar{T}(\bar{z}) = -\frac{1}{2} : \bar{\psi}(\bar{z}) \bar{\partial} \bar{\psi}(\bar{z}) : .$$

$$T(z)T(w) = \frac{1/4}{(z-w)^4} + \frac{2}{(z-w)^2} T(w) + \frac{1}{z-w} \partial T(w)$$

$$T(z) = \sum_{n \in \mathbb{Z}} z^{-n-2} L_n, \bar{T}(\bar{z}) = \sum_{n \in \mathbb{Z}} \bar{z}^{-n-2} \bar{L}_n$$

$$w = \tau + i\sigma \rightarrow z = e^w$$

$$\Phi_{\text{cyl}}(w) = \sum_{n \in \mathbb{Z}} \phi_n e^{-nw} = \sum_{n \in \mathbb{Z}} \phi_n e^{in(\tau - \sigma)} = \sum_{n \in \mathbb{Z}} \phi_n z^{-n}$$

$$\Phi(z) = \sum_{n \in \mathbb{Z}} \phi_n z^{-n-h}$$

$$T(z) \rightarrow (f')^2 T(f(z)) + \frac{c}{12} \left[\frac{f'''}{f'} - \frac{3}{2} \left(\frac{f''}{f'} \right)^2 \right]$$



$$\begin{aligned}
L_n &= \oint \frac{dz}{2\pi i} z^{n+1} T(z), \bar{L}_n = \oint \frac{d\bar{z}}{2\pi i} \bar{z}^{n+1} \bar{T}(\bar{z}) \\
[L_n, L_m] &= \left(\oint \frac{dz}{2\pi i} \oint \frac{dw}{2\pi i} - \oint \frac{dw}{2\pi i} \oint \frac{dz}{2\pi i} \right) z^{n+1} T(z) w^{m+1} T(w) \\
&= \oint \frac{dw}{2\pi i} \oint \frac{dz}{2\pi i} z^{n+1} w^{m+1} \left(\frac{c/2}{(z-w)^4} + \frac{2T(w)}{(z-w)^2} + \frac{\partial T(w)}{z-w} + \dots \right) \\
&= \oint \frac{dw}{2\pi i} \left(\frac{c}{12} (n+1)n(n-1) w^{n-2} w^{m+1} \right. \\
&\quad \left. + 2(n+1) w^n w^{m+1} T(w) + w^{n+1} w^{m+1} \partial T(w) \right) \\
[L_n, L_m] &= (n-m) L_{n+m} + \frac{c}{12} (n^3 - n) \delta_{n+m,0} \\
[\bar{L}_n, \bar{L}_m] &= (n-m) \bar{L}_{n+m} + \frac{\bar{c}}{12} (n^3 - n) \delta_{n+m,0} \\
[L_n, \bar{L}_m] &= 0
\end{aligned}$$

$$T_\alpha^\alpha = \frac{c}{96\pi^3} \sqrt{g} R^{(2)}$$

$$\int [DX]_g e^{-S[\vartheta_{\alpha\beta}, X]} = e^{-cS_L[g_{\alpha\beta}, \phi]} \int [DX]_g e^{-S[g_{\alpha\beta}, X]}$$

$$S_L[g_{\alpha\beta}, \phi] = \frac{1}{96\pi} \int \sqrt{\det g} g^{\alpha\beta} \partial_\alpha \phi \partial_\beta \phi + \frac{1}{48\pi} \int \sqrt{\det g} R^{(2)} \phi$$

Espacio de Hilbert – Einstein.

$$|A_{\rm in}\rangle = \lim_{\tau\rightarrow-\infty} A(\tau,\sigma)|0\rangle = \lim_{z\rightarrow 0} A(z,\bar{z})|0\rangle.$$

$$\tilde{A}(w,\bar{w})=A(f(w),\bar{f}(\bar{w}))(\partial f(w))^h(\bar{\partial}\bar{f}(\bar{w}))^{\bar{h}}$$

$$\tilde{A}(w,\tilde{w})=A\left(\frac{1}{w},\frac{1}{\tilde{w}}\right)(-w^{-2})^h(-\tilde{w}^{-2})^{\bar{h}}$$

$$\langle A_{\rm out} | = \lim_{w,\bar{w} \rightarrow 0} \langle 0 | \tilde{A}(w,\bar{w})$$

$$[A(z,\bar{z})]^\dagger = A\left(\frac{1}{\bar{z}},\frac{1}{z}\right)\bar{z}^{-2h}z^{-2\bar{h}}$$

$$\langle A_{\rm out} | = \lim_{w\rightarrow 0} \langle 0 | \tilde{A}(w,\bar{w}) = \lim_{z\rightarrow 0} \langle 0 | A\left(\frac{1}{z},\frac{1}{\bar{z}}\right)\bar{z}^{-2h}z^{-2\bar{h}} = \lim_{z\rightarrow 0} \langle 0 | [A(z,\bar{z})]^\dagger = |A_{\rm in}\rangle^\dagger$$

$$T^{\dagger}(z)=\sum_m\frac{L_m^{\dagger}}{\bar{z}^{m+2}}\equiv\sum_m\frac{L_m}{\bar{z}^{-m-2}}\frac{1}{\bar{z}^4}$$

$$L_m^\dagger=L_{-m}$$



$$T(z)|0\rangle = \sum_{m\in\mathbb{Z}} L_m z^{-m-2}|0\rangle$$

$$L_m|0\rangle = 0, m \geq -1$$

$$\langle 0|L_m = 0, m \leq 1$$

$$\Phi_{n>-h}|0\rangle = 0$$

$$[L_n, \Phi(w)] = \oint \frac{dz}{2\pi i} z^{n+1} T(z) \Phi(w) = h(n+1)w^n \Phi(w) + w^{n+1} \partial \Phi(w)$$

$$|h\rangle \equiv \Phi(0)|0\rangle.$$

$$L_{m>0}|h\rangle = L_{m>0}\Phi(0)|0\rangle = [L_m, \Phi(0)]|0\rangle + \Phi(0)L_{m>0}|0\rangle = 0$$

$$|\chi\rangle = L_{-n_1}L_{-n_2}\dots L_{-n_k}|h\rangle,$$

$$(L_{-1}\Phi)(z)\equiv\oint_{c_z}\frac{dw}{2\pi i}T(w)\Phi_h(z)$$

$$\Phi_\chi(z)=\prod_{i=1}^k\oint\frac{dw_i}{2\pi i}(w_i-z)^{-n_i+1}T(w_i)\Phi_h(z)$$

$$[L_{-1},O(z,\bar{z})]=\partial_zO(z,\bar{z})$$

$$\chi_h(q)\equiv\mathrm{Tr}\left[q^{L_0-\frac{c}{24}}\right]$$

$$\chi_h(q)=\frac{q^{h-c/24}}{\prod_{n=1}^{\infty}(1-q^n)}$$

$$\chi_0(q)=\frac{q^{-c/24}}{\prod_{n=2}^{\infty}(1-q^n)}$$

$$\|L_{-n}|0\rangle\|^2=\langle 0|L_{-n}^{\dagger}L_{-n}|0\rangle=\langle 0|\left[\frac{c}{12}(n^3-n)+2nL_0\right]|0\rangle=\frac{c}{12}(n^3-n)$$

$$c=1-\frac{6}{m(m+1)}$$

$$|\chi\rangle=\left(L_{-2}-\frac{3}{4}L_{-1}^2\right)|1/2\rangle.$$

$$J^a(z)J^b(w)=\frac{G^{ab}}{(z-w)^2}+\frac{if^{ab}{}_cJ^c(w)}{z-w}+\text{finito}$$

$$[J_m^a,J_n^b]=mG^{ab}\delta_{m+n,0}+if^{ab}{}_cJ_{m+n}^c.$$

$$T(z)J^a(w)=\frac{J^a(w)}{(z-w)^2}+\frac{\partial J^a(w)}{z-w}$$



$$S = \frac{1}{4\lambda^2} \int_{M_2} d^2 \xi \text{Tr}(\partial_\mu g \partial^\mu g^{-1}) + \frac{ik}{8\pi} \int_{B; \partial B = M_2} d^3 \xi \text{Tr}(\epsilon_{\alpha\beta\gamma} U^\alpha U^\beta U^\gamma)$$

$$T_G(z)=\frac{1}{2(k+\tilde{h})}:J^a(z)J^a(z):$$

$$c_G=\frac{kD_G}{k+\tilde{h}}$$

$$J^a_{m>0}|R_i\rangle=0\,,J^a_0|R_i\rangle=i(T^a_R)_{ij}|R_j\rangle$$

$$J^a(z)R_i(w,\bar{w})=i\frac{(T^a_R)_{ij}}{(z-w)}R_j(w,\bar{w})+\cdots$$

$$h_R=\frac{C_R}{k+\tilde{h}},$$

$$T_{\mathrm{G/H}}(z)J^{\mathrm{H}}(w)=\text{ regular } , T_{\mathrm{G/H}}(z)T_{\mathrm{H}}(w)=\text{ regular }$$

Supersimetrías. Métrica de Ramond – Clifford.

$$S=-\frac{1}{8\pi}\int d^2z\psi^i\bar{\partial}\psi^i$$

$$J^{ij}(z)=i\colon\psi^i(z)\psi^j(z)\colon,\,i<j$$

$$\psi^i(z)\psi^j(w)=\frac{\delta^{ij}}{z-w}$$

$$J^{ij}(z)J^{kl}(w)=\frac{G^{ij,kl}}{(z-w)^2}+if^{ij,kl}_{mn}\frac{J^{mn}(w)}{(z-w)}+\cdots$$

$$2f^{ij,kl}{}_{mn}=(\delta^{ik}\delta^{ln}-\delta^{il}\delta^{kn})\delta^{jm}+(\delta^{jl}\delta^{kn}-\delta^{jk}\delta^{ln})\delta^{im}-(m\leftrightarrow n)$$

$$T(z)=\frac{1}{2(N-1)}\sum_{i<j}^N:J^{ij}(z)J^{ij}(z):.$$

$$T(z)T(w)=\frac{c_G/2}{(z-w)^4}+\frac{2T(w)}{(z-w)^2}+\frac{\partial T(w)}{z-w}$$

$$c_G=\frac{kD}{k+\tilde{h}}$$

$$c_G=\frac{N(N-1)/2}{1+N-2}=\frac{N}{2},$$

$$T(z)=-\frac{1}{2}\sum_{i=1}^N:\psi^i\partial\psi^i:$$



$$h_V = \frac{(N-1)/2}{1+N-2} = \frac{1}{2}$$

$$J^{ij}(z)\psi^k(w)=i\frac{T^{ij}_{kl}}{z-w}\psi^l(w)+\cdots,$$

$$\psi^i\rightarrow -\psi^i.$$

$$\psi(\tau+i\sigma)=\sum_n\psi_ne^{-n(\tau+i\sigma)}$$

$$\psi^i(z)=\sum_n\psi_n^iz^{-n-h}=\sum_n\psi_n^iz^{-n-1/2}$$

$$\{\psi_m^i,\psi_n^j\}=\delta^{ij}\delta_{m+n,0}$$

$$\psi^i_{n>0}|0\rangle=0$$

$$|i\rangle=\psi^i_{-\frac{1}{2}}|0\rangle$$

$$\{(-1)^F,\psi_n^i\}=0$$

$$J^{ij}_{-1}|0\rangle=i\psi^i_{-\frac{1}{2}}\psi^j_{-\frac{1}{2}}|0\rangle$$

$$J^{ij}_{-1}|k\rangle=i\left[\delta^{jk}\psi^i_{-\frac{3}{2}}-\delta^{ik}\psi^j_{-\frac{3}{2}}+\psi^i_{-\frac{1}{2}}\psi^j_{-\frac{1}{2}}\psi^k_{-\frac{1}{2}}\right]|0\rangle.$$

$$\mathrm{Tr}_{NS}[q^{L_0-c/24}]=q^{-\frac{N}{48}}\prod_{n=1}^{\infty}\left(1+q^{n-\frac{1}{2}}\right)^N$$

$$\mathrm{Tr}_{NS}[q^{L_0-c/24}]=\left[\frac{\vartheta_3}{\eta}\right]^{N/2}$$

$$\mathrm{Tr}_{NS}[(-1)^Fq^{L_0-c/24}]=q^{-\frac{N}{48}}\prod_{n=1}^{\infty}\left(1-q^{n-\frac{1}{2}}\right)^N=\left[\frac{\vartheta_4}{\eta}\right]^{N/2}$$

$$\begin{aligned}\chi_0 &= \mathrm{Tr}_{NS}\left[\frac{(1+(-1)^F)}{2}q^{L_0-c/24}\right]=\frac{1}{2}\left(\left[\frac{\vartheta_3}{\eta}\right]^{N/2}+\left[\frac{\vartheta_4}{\eta}\right]^{N/2}\right) \\ \chi_V &= \mathrm{Tr}_{NS}\left[\frac{(1-(-1)^F)}{2}q^{L_0-c/24}\right]=\frac{1}{2}\left(\left[\frac{\vartheta_3}{\eta}\right]^{N/2}-\left[\frac{\vartheta_4}{\eta}\right]^{N/2}\right).\end{aligned}$$

$$\chi_R(v_i)=\mathrm{Tr}_R\left[q^{L_0-c/24}e^{2\pi i\sum_iv_iJ_0^i}\right]$$

$$\chi_0(v_i) = \frac{1}{2} \left[\prod_{i=1}^{N/2} \frac{\vartheta_3(v_i)}{\eta} + \prod_{i=1}^{N/2} \frac{\vartheta_4(v_i)}{\eta} \right]$$

$$\chi_V(v_i) = \frac{1}{2} \left[\prod_{i=1}^{N/2} \frac{\vartheta_3(v_i)}{\eta} - \prod_{i=1}^{N/2} \frac{\vartheta_4(v_i)}{\eta} \right]$$

$$\{\psi_0^i,\psi_0^j\}=\delta^{ij}$$

$$\psi_{m>0}^i|\hat{S}_\alpha\rangle=0\,,\psi_0^i|\hat{S}_\alpha\rangle=\gamma_{\alpha\beta}^i|S_\beta\rangle$$

$$\gamma^{N+1}=\prod_{i=1}^N\left(\psi_0^i/\sqrt{2}\right),\{\gamma^{N+1},\psi_0^i\}=0\,,[\gamma^{N+1}]^2=1$$

$$(-1)^F=\gamma^{N+1}(-1)^{\sum_{n=1}^{\infty}\psi_{-n}^i\psi_n^i}$$

$$(-1)^F|S\rangle=|S\rangle\,,(-1)^F|C\rangle=-|C\rangle.$$

$$G_R^{ij}(z,w)=\langle \hat{S}|\psi^i(z)\psi^j(w)|\hat{S}\rangle$$

$$G_R^{ij}(z,w)=\delta^{ij}\frac{z+w}{2\sqrt{zw}}\frac{1}{z-w}$$

$$\langle X|T(z)|X\rangle=\frac{h}{z^2}$$

$$T(w)=\lim_{z\rightarrow w}\left[-\frac{1}{2}\sum_{i=1}^N\psi^i(z)\partial_w\psi^i(w)+\frac{N}{2(z-w)^2}\right]$$

$$\langle \hat{S}|T(z)|\hat{S}\rangle=\frac{N}{16z^2}$$

$$\mathrm{Tr}_R[q^{L_0-c/24}]=2^{N/2}q^{\frac{N}{16}-\frac{N}{48}}\prod_{n=1}^{\infty}(1+q^n)^N=\left[\frac{\vartheta_2}{\eta}\right]^{N/2}$$

$$\chi_S=\chi_C=\frac{1}{2}\left[\frac{\vartheta_2}{\eta}\right]^{N/2}.$$

$$\chi_S(v_i)=\frac{1}{2}\left[\prod_{i=1}^{N/2}\frac{\vartheta_2(v_i)}{\eta}+\prod_{i=1}^{N/2}\frac{\vartheta_1(v_i)}{\eta}\right],$$

$$\chi_C(v_i)=\frac{1}{2}\left[\prod_{i=1}^{N/2}\frac{\vartheta_2(v_i)}{\eta}-\prod_{i=1}^{N/2}\frac{\vartheta_1(v_i)}{\eta}\right]$$

$$|\hat{S}_\alpha\rangle=\lim_{z\rightarrow 0}\hat{S}_\alpha(z)|0\rangle$$

$$\begin{aligned}\psi^i(z)\hat{S}_\alpha(w) &= \gamma_{\alpha\beta}^i \frac{\hat{S}_\beta(w)}{\sqrt{z-w}} + \dots \\ J^{ij}(z)\hat{S}_\alpha(w) &= \frac{i}{2} [\gamma^i, \gamma^j]_{\alpha\beta} \frac{\hat{S}_\beta(w)}{(z-w)} + \dots \\ \hat{S}_\alpha(z)\hat{S}_\beta(w) &= \frac{\delta_{\alpha\beta}}{(z-w)^{N/8}} + \gamma_{\alpha\beta}^i \frac{\psi^i(w)}{(z-w)^{N/8-1/2}} + \frac{i}{2} [\gamma^i, \gamma^j]_{\alpha\beta} \frac{J^{ij}(w)}{(z-w)^{N/8-1}} + \dots\end{aligned}$$

Supersimetría isotrópica. Métrica de Majorana.

$$S = \frac{1}{2\pi} \int d^2z \partial X \bar{\partial} X + \frac{1}{2\pi} \int d^2z (\psi \bar{\partial} \psi + \bar{\psi} \partial \bar{\psi})$$

$$\delta X = \epsilon(z)\psi, \delta\psi = -\epsilon(z)\partial X, \delta\bar{\psi} = 0$$

$$\delta X = \bar{\epsilon}(\bar{z})\bar{\psi}, \delta\bar{\psi} = -\bar{\epsilon}(\bar{z})\bar{\partial} X, \delta\psi = 0$$

$$G(z) = i\psi\partial X, \bar{G}(\bar{z}) = i\bar{\psi}\bar{\partial} X$$

$$\begin{aligned}G(z)G(w) &= \frac{1}{(z-w)^3} + 2\frac{T(w)}{z-w} + \dots \\ T(z)G(w) &= \frac{3}{2}\frac{G(w)}{(z-w)^2} + \frac{\partial G(w)}{z-w} + \dots\end{aligned}$$

$$T(z) = -\frac{1}{2}:\partial X\partial X:-\frac{1}{2}:\psi\partial\psi:.$$

$$\begin{aligned}G(z)G(w) &= \frac{\hat{c}}{(z-w)^3} + 2\frac{T(w)}{z-w} + \dots \\ T(z)G(w) &= \frac{3}{2}\frac{G(w)}{(z-w)^2} + \frac{\partial G(w)}{z-w} + \dots\end{aligned}$$

$$\begin{aligned}\{G_r, G_s\} &= \frac{\hat{c}}{2}\left(r^2 - \frac{1}{4}\right)\delta_{r+s,0} + 2L_{r+s} \\ [L_m, G_r] &= \left(\frac{m}{2} - r\right)G_{m+r}\end{aligned}$$

$$\chi_{N=1}^{NS} = \text{Tr}[q^{L_0 - c/24}] = q^{h-c/24} \prod_{n=1}^{\infty} \frac{1+q^{n-\frac{1}{2}}}{1-q^n}.$$

$$\{G_0, G_0\} = 2L_0 - \frac{\hat{c}}{8}$$

$$\chi_{N=1}^R = \text{Tr}[q^{L_0 - c/24}] = q^{h-c/24} \prod_{n=1}^{\infty} \frac{1+q^n}{1-q^n}.$$

$$D_\theta = \frac{\partial}{\partial \theta} + \theta \partial_z, \bar{D}_{\bar{\theta}} = \frac{\partial}{\partial \bar{\theta}} + \bar{\theta} \partial_{\bar{z}}$$

$$\hat{X}(z,\bar{z},\theta,\bar{\theta})=X+\theta\psi+\bar{\theta}\bar{\psi}+\theta\bar{\theta}F$$



$$S = \frac{1}{2\pi} \int d^2z \int d\theta d\bar{\theta} D_{\theta} \hat{X} \bar{D}_{\bar{\theta}} \hat{X}$$

$$G^+(z)G^-(w) = \frac{2c}{3} \frac{1}{(z-w)^3} + \left(\frac{2J(w)}{(z-w)^2} + \frac{\partial J(w)}{z-w} \right) + \frac{2}{z-w} T(w) + \dots,$$

$$G^+(z)G^+(w) = \text{regular}, G^-(z)G^-(w) = \text{regular},$$

$$T(z)G^{\pm}(w) = \frac{3}{2} \frac{G^{\pm}(w)}{(z-w)^2} + \frac{\partial G^{\pm}(w)}{z-w} + \dots,$$

$$J(z)G^{\pm}(w) = \pm \frac{G^{\pm}(w)}{z-w} + \dots,$$

$$T(z)J(w) = \frac{J(w)}{(z-w)^2} + \frac{\partial J(w)}{z-w} + \dots,$$

$$J(z)J(w) = \frac{c/3}{(z-w)^2} + \dots$$

$$G^{\pm}(e^{2\pi i}z) = e^{\mp 2\pi i\alpha} G^{\pm}(z)$$

$$J_n^{\alpha} = J_n - \alpha \frac{c}{3} \delta_{n,0}, L_n^{\alpha} = L_n - \alpha J_n + \alpha^2 \frac{c}{6} \delta_{n,0}$$

$$G_{r+\alpha}^{\alpha,+} = G_r^+, G_{r-\alpha}^{\alpha,-} = G_r^-$$

$$G_0^{\pm}G_0^{\pm} = 0, \{G_0^+, G_0^-\} = 2\left(L_0 - \frac{c}{24}\right)$$

$$J_0^{R\pm} = J_0^{NS} \mp \frac{c}{6}, L_0^{R\pm} - \frac{c}{24} = L_0^{NS} - \frac{1}{2}J_0^{NS}$$

$$O_{q_1}(z)O_{q_2}(z) = O_{q_1+q_2}(z)$$

$$J^a(z)J^b(w) = \frac{k}{2} \frac{\delta^{ab}}{(z-w)^2} + i\epsilon^{abc} \frac{J^c(w)}{(z-w)} + \dots$$

$$J^a(z)G^\alpha(w) = \frac{1}{2}\sigma_{\beta\alpha}^a \frac{G^\beta(w)}{(z-w)} + \dots, J^a(z)\bar{G}^\alpha(w) = -\frac{1}{2}\sigma_{\alpha\beta}^a \frac{\bar{G}^\beta(w)}{(z-w)} + \dots$$

$$G^\alpha(z)\bar{G}^\beta(w) = \frac{4k\delta^{\alpha\beta}}{(z-w)^3} + 2\sigma_{\beta\alpha}^a \left[\frac{2J^a(w)}{(z-w)^2} + \frac{\delta J^a(w)}{(z-w)} \right] + 2\delta^{\alpha\beta} \frac{T(w)}{(z-w)} + \dots$$

$$G^\alpha(z)G^\beta(w) = \text{regular}, \bar{G}^\alpha(z)\bar{G}^\beta(w) = \text{regular}$$

Campos fantasmas.

$$S_{\lambda} = \frac{1}{\pi} \int d^2z b \bar{\partial} c$$

$$c(z)b(w) = \frac{1}{z-w}, b(z)c(w) = \frac{\epsilon}{z-w}$$

$$c(z) = \sum_{n \in \mathbb{Z}} z^{-n-(1-\lambda)} c_n, c_n^{\dagger} = c_{-n}$$

$$b(z) = \sum_{n \in \mathbb{Z}} z^{-n-\lambda} b_n, b_n^{\dagger} = \epsilon b_{-n}$$



$$c_m b_n + \epsilon b_n c_m = \delta_{m+n,0}, c_m c_n + \epsilon c_n c_m = b_m b_n + \epsilon b_n b_m = 0$$

$$\text{NS: } b_n, n \in \mathbb{Z} - \lambda, c_n, n \in \mathbb{Z} + \lambda$$

$$\text{R: } b_n, n \in \frac{1}{2} + \mathbb{Z} - \lambda, c_n, n \in \frac{1}{2} + \mathbb{Z} + \lambda$$

$$T=-\lambda b\partial c+(1-\lambda)(\partial b)c$$

$$c=-2\epsilon(6\lambda^2-6\lambda+1)=\epsilon(1-3Q^2), Q=\epsilon(1-2\lambda)$$

$$J(z)=-:b(z)c(z):=\sum_{n\in\mathbb{Z}}z^{-n-1}J_n$$

$$J(z)J(w)=\frac{\epsilon}{(z-w)^2}+\cdots$$

$$J(z)b(w)=-\frac{b(w)}{z-w}+\cdots, J(z)c(w)=\frac{c(w)}{z-w}\cdots$$

$$T(z)J(w)=\frac{Q}{(z-w)^3}+\frac{J(w)}{(z-w)^2}+\frac{\partial_wJ(w)}{z-w}+\cdots$$

$$[L_m,J_n]=-nJ_{m+n}+\frac{Q}{2}m(m+1)\delta_{m+n,0}$$

$$[L_1,J_{-1}]=J_0+Q, Q+J_0^\dagger=[L_1,J_{-1}]^\dagger=-J_0$$

$$\mathbb{Z}_{\text{zero modes of } c}-\mathbb{Z}_{\# \text{ zero modes of } b}=-\frac{\epsilon}{2}Q\chi$$

$$b_{n>-\lambda}|0\rangle=c_{n>\lambda-1}|0\rangle=0$$

$$J(z)=-\partial\phi\,,\langle\phi(z)\phi(w)\rangle=-\log\,(z-w)$$

$$\hat{T}=\frac{1}{2}:J^2:+\frac{1}{2}Q\partial J=\frac{1}{2}(\partial\phi)^2-\frac{Q}{2}\partial^2\phi.$$

$$S_Q=\frac{1}{2\pi}\int\,d^2z\left[\partial\phi\bar{\partial}\phi-\frac{Q}{4}\sqrt{g}R^{(2)}\phi\right]$$

$$T=\hat{T}+T_{\eta\xi}$$

$$T(z):e^{q\phi(w)}:=\left[-\frac{q(q+Q)}{(z-w)^2}+\frac{1}{z-w}\partial_w\right]:e^{q\phi(w)}:+\ldots$$

$$J(z):e^{q\phi(w)}:=\frac{q}{z-w}:e^{q\phi(w)}:\ldots\rightarrow [J_0,:e^{q\phi(w)}:] = q:e^{q\phi(w)}:..$$

$$c(z)=e^{\phi(z)}\eta(z)\,,b(z)=e^{-\phi(z)}\partial\xi(z)$$

$$g_{ij}=\frac{1}{\tau_2}\begin{pmatrix}1&\tau_1\\ \tau_1&|\tau|^2\end{pmatrix}$$



$$ds^2 = g_{ij}d\sigma_i d\sigma_j = \frac{1}{\tau_2} |d\sigma_1 + \tau d\sigma_2|^2 = \frac{dw d\bar{w}}{\tau_2}$$

$$w = \sigma_1 + \tau \sigma_2, \bar{w} = \sigma_1 + \bar{\tau} \sigma_2$$

$$w \rightarrow w + 1, \bar{w} \rightarrow \bar{w} + \tau$$

$$T: \tau \rightarrow \tau + 1$$

$$TST: \tau \rightarrow \frac{\tau}{\tau + 1}.$$

$$S: \tau \rightarrow -\frac{1}{\tau}, S^2 = 1, (ST)^3 = 1$$

$$\tau' = \frac{a\tau + b}{c\tau + d} \leftrightarrow A = \begin{pmatrix} a & b \\ c & d \end{pmatrix}$$

$$L_0^{cyl} = L_0 - \frac{c}{24}, \bar{L}_0^{cyl} = \bar{L}_0 - \frac{\bar{c}}{24},$$

$$\begin{aligned} \int e^{-S} &= \text{Tr}[e^{-2\pi\tau_2 H} e^{2\pi i \tau_1 P}] = \text{Tr}[e^{2\pi i \tau L_0^{cyl}} e^{-2\pi i \bar{\tau} \bar{L}_0^{cyl}}] \\ &= \text{Tr}[q^{L_0 - c/24} \bar{q}^{\bar{L}_0 - \bar{c}/24}] \end{aligned}$$

Escalares compactos.

$$S = \frac{1}{4\pi} \int d^2\sigma \sqrt{g} g^{ij} \partial_i X \partial_j X = \frac{1}{4\pi} \int_0^1 d\sigma_1 \int_0^1 d\sigma_2 \frac{1}{\tau_2} |\tau \partial_1 X - \partial_2 X|^2 = -\frac{1}{4\pi} \int d^2\sigma X \square X$$

$$\square = \frac{1}{\tau_2} |\tau \partial_1 - \partial_2|^2.$$

$$Z(R) = \int DX e^{-S}$$

$$X_{\text{class}} = 2\pi R(n\sigma_1 + m\sigma_2), m, n \in \mathbb{Z}$$

$$X_{\text{class}}(\sigma_1 + 1, \sigma_2) = X(\sigma_1, \sigma_2) + 2\pi n R, X_{\text{class}}(\sigma_1, \sigma_2 + 1) = X(\sigma_1, \sigma_2) + 2\pi m R$$

$$S_{m,n} = \frac{\pi R^2}{\tau_2} |m - n\tau|^2$$

$$\begin{aligned} Z(R) &= \sum_{m,n \in \mathbb{Z}} \int D\chi e^{-S_{m,n} - S(\chi)} \\ &= \sum_{m,n \in \mathbb{Z}} e^{-S_{m,n}} \int D\chi e^{-S(\chi)} \end{aligned}$$

$$\square \psi_i = -\lambda_i \psi_i$$



$$\psi_{m_1, m_2} = e^{2\pi i(m_1 \sigma_1 + m_2 \sigma_2)}, \lambda_{m_1, m_2} = \frac{4\pi^2}{\tau_2} |m_1 \tau - m_2|^2$$

$$\int d^2\sigma \psi_{m_1, m_2} \psi_{n_1, n_2} = \delta_{m_1+n_1, 0} \delta_{m_2+n_2, 0}$$

$$\delta\chi = \sum_{m_1, m_2 \in \mathbb{Z}} 'A_{m_1, m_2} \psi_{m_1, m_2}$$

$$S(\chi) = \frac{1}{4\pi} \sum_{m_1, m_2} ' \lambda_{m_1, m_2} |A_{m_1, m_2}|^2$$

$$\|\delta X\| = \int d^2\sigma \sqrt{\det G} (d\chi)^2 = \sum_{m_1, m_2} ' |dA_{m_1, m_2}|^2$$

$$\int D\chi = \int_0^{2\pi R} d\chi_0 \prod_{m_1, m_2} \frac{dA_{m_1, m_2}}{2\pi}$$

$$\int D\chi e^{-S(\chi)} = \frac{2\pi R}{\prod'_{m_1, m_2} \lambda_{m_1, m_2}^{1/2}} = \frac{2\pi R}{\sqrt{\det' \square}}$$

$$\det' \square = 4\pi^2 \tau_2 \eta^2(\tau) \bar{\eta}^2(\bar{\tau})$$

$$Z(R) = \frac{R}{\sqrt{\tau_2} |\eta|^2} \sum_{m, n \in \mathbb{Z}} e^{-\frac{\pi R^2}{\tau_2} |m - n\tau|^2}$$

$$Z(R) = \sum_{m, n \in \mathbb{Z}} \frac{q^{\frac{P_L^2}{2}} \bar{q}^{\frac{P_R^2}{2}}}{\eta \bar{\eta}}$$

$$P_L = \frac{1}{\sqrt{2}} \left(\frac{m}{R} + nR \right), P_R = \frac{1}{\sqrt{2}} \left(\frac{m}{R} - nR \right)$$

$$J(z) = i\partial X, \bar{J}(\bar{z}) = i\bar{\partial} X$$

$$J(z)J(w) = \frac{1}{(z-w)^2} + \text{finite}, \bar{J}(\bar{z})\bar{J}(\bar{w}) = \frac{1}{(\bar{z}-\bar{w})^2} + \text{finito}$$

- Sugawara:

$$T(z) = -\frac{1}{2}(\partial X)^2 = \frac{1}{2}:J^2:, \bar{T}(\bar{z}) = -\frac{1}{2}(\bar{\partial} X)^2 = \frac{1}{2}:\bar{J}^2:.$$

$$\Delta = \frac{1}{2}Q_L^2, \bar{\Delta} = \frac{1}{2}Q_R^2$$

$$\chi_{Q_L, Q_R}(q, \bar{q}) = \text{Tr}[q^{L_0-1/24} \bar{q}^{\bar{L}_0-1/24}] = \frac{q^{Q_L^2/2} \bar{q}^{Q_R^2/2}}{\eta \bar{\eta}}$$



$$J_0|m,n\rangle = P_L|m,n\rangle, \bar{J}_0|m,n\rangle = P_R|m,n\rangle$$

$$V_{m,n} =: \exp\left[ip_L X + ip_R \bar{X} \right]:,$$

$$J(z)V_{m,n}(w,\bar{w})=p_L\frac{V_{m,n}(w,\bar{w})}{z-w}+\cdots$$

$$\bar{J}(\bar{z})V_{m,n}(w,\bar{w})=p_R\frac{V_{m,n}(w,\bar{w})}{\bar{z}-\bar{w}}+\cdots$$

$$\left\langle \prod_{i=1}^N V_{m_i,n_i}(z_i,\bar{z}_i) \right\rangle = \prod_{i<j}^N z_{ij}^{p_L^i p_L^j} \bar{z}_{ij}^{p_R^i p_R^j},$$

$$:e^{ia\phi(z)}::e^{ib\phi(w)}:=(z-w)^{ab}:e^{ia\phi(z)+ib\phi(w)}:=(z-w)^{ab}[:e^{i(a+b)\phi(w)}: +\mathcal{O}(z-w)]$$

$$\left[V_{m_1,n_1}\right]\cdot\left[V_{m_2,n_2}\right]\sim\left[V_{m_1+m_2,n_1+n_2}\right]$$

$$\Delta-\bar{\Delta}-\frac{c-\bar{c}}{24}=\mathfrak{I}_{\mathrm{integer}}$$

$$\lim_{R\rightarrow\infty}\frac{Z(R)}{R}=\frac{1}{\sqrt{\tau_2}\eta\bar\eta}$$

$$\Delta(\sigma_1,\sigma_2)\equiv \langle \delta\chi(\sigma_1,\sigma_2)\delta\chi(0,0)\rangle=-\sum_{m,n}^{\prime}\frac{1}{|m\tau-n|^2}e^{2\pi i(m\sigma_1+n\sigma_2)}$$

$$\square\,\Delta(\sigma_1,\sigma_2)=\frac{4\pi^2}{\tau_2}[\delta(\sigma_1)\delta(\sigma_2)-1]$$

$$\Delta(\sigma_1,\sigma_2)=-\log\,G(z,\bar{z})\,,G=e^{-2\pi\frac{Imz^2}{\tau_2}}\left|\frac{\vartheta_1(z)}{\vartheta_1'(0)}\right|^2$$

$$S=\frac{1}{4\pi}\int\,d^2\sigma\sqrt{\det g}g^{ab}G_{ij}\partial_aX^i\partial_bX^j+\frac{1}{4\pi}\int\,d^2\sigma\epsilon^{ab}B_{ij}\partial_aX^i\partial_bX^j,$$

$$\mathbb{Z}_{\mathrm{d},\,\mathrm{d}}(G,B)=\frac{\sqrt{\det G}}{(\sqrt{\tau_2}\eta\bar\eta)^N}\sum_{\vec{m},\vec{n}}e^{-\frac{\pi(G_{ij}+B_{ij})}{\tau_2}(m_i+n_i\tau)(m_j+n_j\bar\tau)}$$

$$\mathbb{Z}_{\mathrm{d},\,\mathrm{d}}(G,B)=\frac{\Gamma_{d,d}(G,B)}{\eta^d\bar\eta^d}=\sum_{\vec{m},\vec{n}\in\mathbb{Z}^N}\frac{q^{\frac{1}{2}P_L^2}\bar{q}^{\frac{1}{2}P_R^2}}{\eta^N\bar\eta^N}$$

$$P_{L,R}^2\equiv P_{L,R}^iG_{ij}P_{L,R}^j$$

$$P_L^i=\frac{G^{ij}}{\sqrt{2}}\big(m_j+(B_{jk}+G_{jk})n_k\big),P_R^i=\frac{G^{ij}}{\sqrt{2}}\big(m_j+(B_{jk}-G_{jk})n_k\big)$$

$$J^i(z) = i\partial X^i, \bar{J}^i = i\bar{\partial} X^i$$

$$J^i(z)J^j(w) = \frac{G^{ij}}{(z-w)^2} + \dots$$

$$T(z) = -\frac{1}{2}G_{ij}\partial X^i\partial X^j = \frac{1}{2}G_{ij}:J^iJ^j:.$$

$$\Delta = \frac{1}{2}G_{ij}Q_L^iQ_L^j, \bar{\Delta} = \frac{1}{2}G_{ij}Q_R^iQ_R^j$$

Supersimetría de Higgs.

$$\Delta = \frac{1}{4}(m+n)^2, \bar{\Delta} = \frac{1}{4}(m-n)^2$$

$$J^{\pm}(z)=\frac{1}{\sqrt{2}}:e^{\pm i\sqrt{2}X(z)}:$$

$$J^3(z)=\frac{1}{\sqrt{2}}J(z)=\frac{i}{\sqrt{2}}\partial X(z)$$

$$J^3(z)J^{\pm}(w)=\pm\frac{J^{\pm}(w)}{z-w}+\dots, J^{+}(z)J^{+}(w)=\dots$$

$$J^{-}(z)J^{-}(w)=\dots,$$

$$J^{+}(z)J^{-}(w)=\frac{1/2}{(z-w)^2}+\frac{J^3(w)}{z-w}+\dots,$$

$$J^3(z)J^3(w)=\frac{1/2}{(z-w)^2}+\dots$$

$$a_{-1}^{\mu}\bar{a}_{-1}^{\nu}|0\rangle, a_{-1}^{\mu}\bar{a}_{-1}^{25}|0\rangle, a_{-1}^{25}\bar{a}_{-1}^{\mu}|0\rangle, a_{-1}^{25}\bar{a}_{-1}^{25}|0\rangle,$$

$$|A_{\mu}^{\pm}\rangle=\bar{a}^{\mu}|m=\pm1,n=\pm1\rangle$$

$$|\bar{A}_{\mu}^{\pm}\rangle=a^{\mu}|m=\pm1,n=\mp1\rangle$$

Dualidad T.

$$H=L_0+\bar{L}_0=\frac{1}{2}\left(\frac{m^2}{R^2}+n^2R^2\right), P=L_0-\bar{L}_0=mn$$

$$R\rightarrow\frac{1}{R}, m\leftrightarrow n$$

$$P_L\rightarrow P_L, P_R\rightarrow -P_R$$

$$J(z)\rightarrow J(z), \bar{J}(\bar{z})\rightarrow -\bar{J}(\bar{z})$$

$$\Omega=\begin{pmatrix}A&B\\C&D\end{pmatrix}$$

$$L=\begin{pmatrix}0&\mathbf{1}_N\\ \mathbf{1}_N&0\end{pmatrix}$$



$$\Omega^T L \Omega = L$$

$$E \rightarrow (AE+B)(CE+D)^{-1}, \begin{pmatrix} \vec{m} \\ \vec{n} \end{pmatrix} \rightarrow \Omega \begin{pmatrix} \vec{m} \\ \vec{n} \end{pmatrix}$$

$$G_{ij}=\frac{T_2}{U_2}\begin{pmatrix}1&U_1\\U_1&U_1^2+U_2^2\end{pmatrix},B_{ij}=\begin{pmatrix}0&T_1\\-T_1&0\end{pmatrix},$$

$$\Gamma_{2,2}(T,U)=\sum_{\vec{m},\vec{n}}\exp\left[-\frac{\pi\tau_2}{T_2U_2}\left|-m_1U+m_2+T(n_1+Un_2)\right|^2+2\pi i\tau(m_1n_1+m_2n_2)\right]$$

$$\int e^{-S}=(\det\partial)^{N/2}$$

$$\lambda_{AA}\sim \left(\left(m_1+\frac{1}{2}\right)\tau+\left(m_2+\frac{1}{2}\right)\right), m_{1,2}\in\mathbb{Z}$$

$$(\det\partial)_{AA}=\frac{\vartheta_3(\tau)}{\eta(\tau)}$$

$$\lambda_{AP}\sim \left(\left(m_1+\frac{1}{2}\right)\tau+m_2\right), m_{1,2}\in\mathbb{Z}$$

$$(\det\partial)_{AP}=\frac{\vartheta_4(\tau)}{\eta(\tau)}$$

$$\lambda_{PA}\sim \left(m_1\tau+\left(m_2+\frac{1}{2}\right)\right), m_{1,2}\in\mathbb{Z}$$

$$(\det\partial)_{PA}=\frac{\vartheta_2(\tau)}{\eta(\tau)}$$

$$(\det\partial)\begin{bmatrix}a\\b\end{bmatrix}=\frac{\vartheta\begin{bmatrix}a\\b\end{bmatrix}(\tau)}{\eta(\tau)}$$

$$Z_N^{\rm fermionic}=\frac{1}{2}\sum_{a,b=0}^1\left|\frac{\vartheta\begin{bmatrix}a\\b\end{bmatrix}}{\eta}\right|^N$$

$$Z_N^{\rm fermionic} = |\chi_0|^2 + |\chi_V|^2 + |\chi_S|^2 + |\chi_C|^2$$

- Szegő:

$$\langle \psi^i(z)\psi^j(0)\rangle=\delta^{ij}S\begin{bmatrix}a\\b\end{bmatrix}(z),S\begin{bmatrix}a\\b\end{bmatrix}(z)=\frac{\vartheta\begin{bmatrix}a\\b\end{bmatrix}(z)\vartheta_1'(0)}{\vartheta_1(z)\vartheta\begin{bmatrix}a\\b\end{bmatrix}(0)}$$

$$q^{-N/24}\prod_{n=1}^{\infty}(1-q^n)^N=\eta^N=\left[\frac{1}{2\pi}\frac{\partial_v\vartheta_1(v)|_{v=0}}{\eta}\right]^{N/2}$$



$$\left\langle \prod_{k=1}^N \psi^{i_k}(z_k) \right\rangle_{\text{odd}} = \epsilon^{i_1, \dots, i_N} \eta^N$$

Bosonización Majorana-Weyl $\psi^i(z)$:

$$\psi^i(z)\psi^j(w) = \frac{\delta^{ij}}{z-w} + \dots$$

$$\psi = \frac{1}{\sqrt{2}}(\psi^1 + i\psi^2), \bar{\psi} = \frac{1}{\sqrt{2}}(\psi^1 - i\psi^2)$$

$$J(z) = :\psi\bar{\psi}:, J(z)J(w) = \frac{1}{(z-w)^2} + \dots$$

$$J(z)\psi(w) = \frac{\psi(w)}{z-w} + \dots, J(z)\bar{\psi}(w) = -\frac{\bar{\psi}(w)}{z-w} + \dots$$

$$T(z) = -\frac{1}{2}:\psi^i\partial\psi^i: = \frac{1}{2}:J^2:.$$

$$J(z) = i\partial X, \psi =:e^{iX}:, \bar{\psi} =:e^{-iX}:$$

- Poisson:

$$\begin{aligned} \left| \vartheta \begin{bmatrix} a \\ b \end{bmatrix} \right|^2 &= \frac{1}{\sqrt{2\tau_2}} \sum_{m,n \in \mathbb{Z}} \exp \left[-\frac{\pi}{2\tau_2} |n-b+\tau(m-a)|^2 + i\pi mn \right] \\ &= \frac{1}{\sqrt{2\tau_2}} \sum_{m,n \in \mathbb{Z}} \exp \left[-\frac{\pi}{2\tau_2} |n+\tau m|^2 + i\pi(m+a)(n+b) \right] \end{aligned}$$

$$Z = \frac{1}{2} \sum_{a,b=0}^1 \left| \frac{\vartheta \begin{bmatrix} a \\ b \end{bmatrix}}{\eta} \right|^2 = \frac{1}{2\sqrt{2\tau_2}} \sum_{a,b=0}^1 \sum_{m,n \in \mathbb{Z}} \exp \left[-\frac{\pi}{2\tau_2} |n+\tau m|^2 + i\pi(m+a)(n+b) \right]$$

$$Z_{\text{Dirac}} = \frac{1}{\sqrt{2\tau_2}} \sum_{m,n \in \mathbb{Z}} \exp \left[-\frac{\pi}{2\tau_2} |n+\tau m|^2 \right]$$

Orbifolds.

$$g \left[\prod_{i=1}^N a_{-n_i} \prod_{j=1}^{\bar{N}} \bar{a}_{\bar{n}_j} |m, n\rangle \right] = (-1)^{N+\bar{N}} \prod_{i=1}^N a_{-n_i} \prod_{j=1}^{\bar{N}} \bar{a}_{\bar{n}_j} |-m, -n\rangle.$$

$$Z(R)^{\text{invariant}} = \frac{1}{2} Z(R) + \frac{1}{2} \text{Tr} [g q^{L_0-1/24} \bar{q}^{\bar{L}_0-1/24}].$$

$$\frac{1}{2} \text{Tr} [g q^{L_0-1/24} \bar{q}^{\bar{L}_0-1/24}] = \frac{1}{2} (q\bar{q})^{-1/24} \prod_{n=1}^{\infty} \frac{1}{(1+q^n)(1+\bar{q}^n)} = \left| \frac{\eta}{\vartheta_2} \right|$$



$$Z(R)^{\text{invariant}} = \frac{1}{2}Z(R) + \left| \frac{\eta}{\vartheta_2} \right|.$$

$$X(\sigma,\tau)=x_0+\frac{i}{\sqrt{4\pi T}}\sum_{n\in\mathbb{Z}}\left(\frac{a_{n+1/2}}{n+1/2}e^{i(n+1/2)(\sigma+\tau)}+\frac{\bar{a}_{n+1/2}}{n+1/2}e^{-i(n+1/2)(\sigma-\tau)}\right).$$

$$a_{n+1/2}|H^{0,\pi}\rangle=\bar{a}_{n+1/2}|H^{0,\pi}\rangle=0\; n\geq 0$$

$$Z^{\text{twisted}}=\frac{1}{2}\text{Tr}[(1+g)q^{L_0-1/24}\bar{q}^{\bar{L}_0-1/24}]$$

$$\begin{aligned} &= \frac{1}{2} \frac{1}{(q\bar{q})^{48}} \left[\prod_{n=1}^{\infty} \frac{1}{\left(1-q^{n-\frac{1}{2}}\right)\left(1-\bar{q}^{n-\frac{1}{2}}\right)} + \prod_{n=1}^{\infty} \frac{1}{\left(1+q^{n-\frac{1}{2}}\right)\left(1+\bar{q}^{n-\frac{1}{2}}\right)} \right] \\ &= \left| \frac{\eta}{\vartheta_4} \right| + \left| \frac{\eta}{\vartheta_3} \right| \end{aligned}$$

$$Z^{\text{orb}}\left(R\right)=Z^{\text{untwisted}}+Z^{\text{twisted}}=\frac{1}{2}Z(R)+\left|\frac{\eta}{\vartheta_2}\right|+\left|\frac{\eta}{\vartheta_4}\right|+\left|\frac{\eta}{\vartheta_3}\right|$$

$$Z^{\text{orb}}=\frac{1}{2}\sum_{h,g=0}^1Z\left[\begin{smallmatrix}h\\g\end{smallmatrix}\right]$$

$$Z\left[\begin{smallmatrix}h\\g\end{smallmatrix}\right]=2\left|\frac{\eta}{\vartheta\left[\begin{smallmatrix}1-h\\1-g\end{smallmatrix}\right]}\right|,\left(h,g\right)\neq\left(0,0\right)$$

$$\begin{array}{l} \tau\rightarrow\tau+1\colon Z\left[\begin{smallmatrix}h\\g\end{smallmatrix}\right]\rightarrow Z\left[\begin{smallmatrix}h\\h+g\end{smallmatrix}\right] \\ \tau\rightarrow-\frac{1}{\tau}\colon Z\left[\begin{smallmatrix}h\\g\end{smallmatrix}\right]\rightarrow Z\left[\begin{smallmatrix}h\\g\end{smallmatrix}\right] \end{array}$$

$$[H^0]\cdot[H^0]\sim\sum_{n,m}C^{2m,2n}[V_{2m,2n}^+]+C^{2m,2n+1}[V_{2m,2n+1}^+]$$

$$[H^\pi]\cdot[H^\pi]\sim\sum_{n,m}C^{2m,2n}[V_{2m,2n}^+]-C^{2m,2n+1}[V_{2m,2n+1}^+]$$

$$[H^0]\cdot[H^\pi]\sim\sum_{n,m}C^{2m+1,2n}[V_{2m+1,2n}^+]$$

$$C^{m,n}=\sqrt{2}2^{-2(h_{m,n}+\bar{h}_{m,n})}\,, C_{0,0}=1$$

$$h_{m,n}=(m/R+nR)^2/4\,, \bar{h}_{m,n}=(m/R-nR)^2/4$$

$$\begin{array}{l} (H^0,H^\pi,V_{m,n}^+)\rightarrow(-H^0,H^\pi,(-1)^mV_{m,n}^+)\\ (H^0,H^\pi,V_{m,n}^+)\rightarrow(H^\pi,H^0,(-1)^nV_{m,n}^+) \end{array}$$



$$(H^0, H^\pi, V_{m,n}^+) \rightarrow (-H^0, -H^\pi, V_{m,n}^+)$$

$$Z^{\text{orb}}(R) = Z^{\text{orb}}(1/R)$$

$$\begin{pmatrix} H^0 \\ H^\pi \end{pmatrix} \rightarrow \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix} \begin{pmatrix} H^0 \\ H^\pi \end{pmatrix}$$

$$Z^{\text{Ising}} = \frac{1}{2} \left[\left| \frac{\vartheta_2}{\eta} \right| + \left| \frac{\vartheta_3}{\eta} \right| + \left| \frac{\vartheta_4}{\eta} \right| \right],$$

$$Z \begin{bmatrix} 0 \\ 1 \end{bmatrix} = \sum_{m,n \in \mathbb{Z}} (-1)^m \frac{\exp \left[\frac{i\pi\tau}{2} \left(\frac{m}{R} + nR \right)^2 - \frac{i\pi\bar{\tau}}{2} \left(\frac{m}{R} - nR \right)^2 \right]}{\eta\bar{\eta}}$$

$$Z \begin{bmatrix} 1 \\ 0 \end{bmatrix} = \sum_{m,n \in \mathbb{Z}} \frac{\exp \left[\frac{i\pi\tau}{2} \left(\frac{m}{R} + \left(n + \frac{1}{2}\right)R \right)^2 - \frac{i\pi\bar{\tau}}{2} \left(\frac{m}{R} - \left(n + \frac{1}{2}\right)R \right)^2 \right]}{\eta\bar{\eta}}.$$

$$Z \begin{bmatrix} 1 \\ 1 \end{bmatrix} = \sum_{m,n \in \mathbb{Z}} (-1)^m \frac{\exp \left[\frac{i\pi\tau}{2} \left(\frac{m}{R} + \left(n + \frac{1}{2}\right)R \right)^2 - \frac{i\pi\bar{\tau}}{2} \left(\frac{m}{R} - \left(n + \frac{1}{2}\right)R \right)^2 \right]}{\eta\bar{\eta}}.$$

$$Z \begin{bmatrix} h \\ g \end{bmatrix} = \sum_{m,n \in \mathbb{Z}} (-1)^{gm} \frac{\exp \left[\frac{i\pi\tau}{2} \left(\frac{m}{R} + \left(n + \frac{h}{2}\right)R \right)^2 - \frac{i\pi\bar{\tau}}{2} \left(\frac{m}{R} - \left(n + \frac{h}{2}\right)R \right)^2 \right]}{\eta\bar{\eta}}$$

$$Z \begin{bmatrix} h \\ g \end{bmatrix} = \frac{R}{\sqrt{\tau_2} \eta \bar{\eta}} \sum_{m,n \in \mathbb{Z}} \exp \left[-\frac{\pi R^2}{\tau_2} \left| m + \frac{g}{2} + \left(n + \frac{h}{2}\right) \tau \right|^2 \right]$$

$$g_{\text{translation}} = \exp \left[2\pi i \sum_{i=1}^d (m_i \theta_i + n_i \phi_i) \right]$$

CFT en superficies de Riemann. Amplitudes escalares y operadores de vértice.

$$\langle 1 \rangle_{g=2} = \sum_i q^{h_i - c/24} \bar{q}^{\bar{h}_i - \bar{c}/24} \langle \phi_i \rangle_{g=1} \langle \phi_i \rangle_{g=1}$$

$$S_P = \frac{1}{4\pi\alpha'} \int d^2\sigma \partial X^\mu \bar{\partial} X^\nu \eta_{\mu\nu}$$

$$\langle X^\mu(z,\bar{z})X^\nu(w,\bar{w})\rangle=-\frac{\alpha'}{4}\eta^{\mu\nu}\log\,|z-w|^2$$

$$T=-\frac{2}{\alpha'}\eta_{\mu\nu}\partial X^\mu\partial X^\nu$$

$$T(z)O(w,\bar{w})=-ip^\mu\epsilon_{\mu\nu}\frac{\alpha'}{4}\frac{\bar{\partial}x^\nu V_p}{(z-w)^3}+\left(1+\frac{\alpha'p^2}{4}\right)\frac{O(w,\bar{w})}{(z-w)^2}+\frac{\partial_wO(w,\bar{w})}{z-w}+\cdots$$



$$p^\mu \epsilon_{\mu\nu} = p^\nu \epsilon_{\mu\nu} = 0$$

$$\Lambda^4 = \int_{\mathcal{F}} \frac{d^2\tau}{\tau_2^2} Z_{\text{bosonic}}(\tau, \bar{\tau}) = \int_{\mathcal{F}} \frac{d^2\tau}{\tau_2^2} \frac{1}{(\sqrt{\tau_2} \eta \bar{\eta})^{24}}$$

$$\chi = 2(1-g) - B - C.$$

Acción stress – tensor – energía - momentum. Centro de masa. Partícula Supermasiva u Oscura.

$$S_P = \frac{1}{4\pi\alpha'} \int d^2\xi \left[\sqrt{g} g^{ab} G_{\mu\nu}(X) + \epsilon^{ab} B_{\mu\nu}(X) \right] \partial_a X^\mu \partial_b X_\nu + \frac{1}{8\pi} \int d^2\xi \sqrt{g} R^{(2)} \Phi(X)$$

$$\chi = \frac{1}{4\pi} \int \sqrt{g} R^{(2)}$$

$$\mathcal{G}_{\text{spacetime dimension}} = e^{\Phi_0/2}$$

$$\frac{T_a^a}{\sqrt{g}} = \frac{\beta^\Phi}{96\pi^3} R^{(2)} + \frac{1}{2\pi} (\beta_{\mu\nu}^G g^{ab} + \beta_{\mu\nu}^B \epsilon^{ab}) \partial_a X^\mu \partial_b X_\nu$$

$$\frac{\beta_{\mu\nu}^G}{\alpha'} = R_{\mu\nu} - \frac{1}{4} H_{\mu\rho\sigma} H_\nu^{\rho\sigma} + \nabla_\mu \nabla_\nu \Phi + \mathcal{O}(\alpha')$$

$$\frac{\beta_{\mu\nu}^B}{\alpha'} = \nabla^\mu [e^{-\Phi} H_{\mu\nu\rho}] + \mathcal{O}(\alpha')$$

$$\beta^\Phi = D - 26 + 3\alpha' \left[(\nabla\Phi)^2 - 2 \square \Phi - R + \frac{1}{12} H^2 \right] + \mathcal{O}(\alpha'^2)$$

$$H_{\mu\nu\rho} = \partial_\mu B_{\nu\rho} + \partial_\nu B_{\rho\mu} + \partial_\rho B_{\mu\nu}$$

$$\beta^\Phi = \beta_{\mu\nu}^G = \beta_{\mu\nu}^B = 0$$

$$\alpha'^{D-2} S^{\text{tree}} \sim \int d^D x \sqrt{-\det G} e^{-\Phi} \left[R + (\nabla\Phi)^2 - \frac{1}{12} H^2 + \frac{D-26}{3} \right] + \mathcal{O}(\alpha')$$

$$G_{\mu\nu}^E = e^{-\frac{2\Phi}{D-2}} G_{\mu\nu}$$

$$S_E^{\text{tree}} \sim \frac{1}{\kappa^2} \int d^D x \sqrt{G^E} \left[R - \frac{1}{D-2} (\nabla\Phi)^2 - \frac{e^{-4\Phi/(D-2)}}{12} H^2 + e^{2\Phi/(D-2)} \frac{D-26}{3} \right] + \mathcal{O}(\alpha')$$

$$\kappa = g_{\text{spacetime dimension}} \alpha'^{(D-2)/2}$$

Supersimetría de Einstein – Polyakov – Virasoro – Majorana - Weyl. Superconformidad de gauge.

$$S_P^{\text{II}} = \frac{1}{4\pi\alpha'} \int \sqrt{g} \left[g^{ab} \partial_a X^\mu \partial_b X^\mu + \frac{i}{2} \psi^\mu \partial \psi^\mu + \frac{i}{2} (\chi_a \gamma^b \gamma^a \psi^\mu) \left(\partial_b X^\mu - \frac{i}{4} \chi_b \psi^\mu \right) \right]$$



$$\delta g_{ab} = i\epsilon(\gamma_a \chi_b + \gamma_b \chi_a), \delta \chi_a = 2\nabla_a \epsilon, \\ \delta X^\mu = i\epsilon \psi^\mu, \delta \psi^\mu = \gamma^a \left(\partial_a X^\mu - \frac{i}{2} \chi_a \psi^\mu \right) \epsilon, \delta \bar{\psi}^\mu = 0,$$

$$g_{ab} = e^\phi \delta_{ab}, \chi_a = \gamma_a \zeta$$

$$G_{\text{matter}} = i\psi^\mu \partial X^\mu, \bar{G}_{\text{matter}} = i\bar{\psi}^\mu \bar{\partial} X^\mu$$

$$j_{BRST} = \gamma G_{\text{matter}} + c T_{\text{matter}} + \frac{1}{2} (c T_{\text{ghost}} + \gamma G_{\text{ghost}})$$

$$G_{\text{matter}} = i\psi^\mu \partial X^\mu, T_{\text{matter}} = -\frac{1}{2} \partial X^\mu \partial X^\mu - \frac{1}{2} \psi^\mu \partial \psi^\mu,$$

$$G_{\text{ghost}} = -i \left(c \partial \beta - \frac{1}{2} \gamma b + \frac{3}{2} \partial c \beta \right), T_{\text{ghost}} = T_{bc} - \frac{1}{2} \gamma \partial \beta - \frac{3}{2} \partial \gamma \beta.$$

$$Q = \frac{1}{2\pi i} [\oint dz j_{BRST} + \oint d\bar{z} \bar{j}_{BRST}]$$

$$X^+ = x^+ + p^+ \tau, \psi^+ = \bar{\psi}^+ = 0$$

$$0 = \{G_0, G_0\} = 2 \left(L_0 - \frac{D-2}{16} \right)$$

$$L_0 = \bar{L}_0, L_0 - \frac{1}{2} = 0$$

$$(-1)^F = \exp \left[i\pi \sum_{r \in Z+1/2} \psi_r^i \psi_{-r}^i \right].$$

$$(-1)^F = \prod_{\mu=0}^9 \psi_0^\mu \exp \left[i\pi \sum_{n=1}^{\infty} \psi_n^i \psi_{-n}^i \right] = \Gamma^{11} \exp \left[i\pi \sum_{n=1}^{\infty} \psi_n^i \psi_{-n}^i \right]$$

$$\{(-1)^{F_L}, G_0\} = 0, \{(-1)^{F_R}, \bar{G}_0\} = 0$$

$$\{\Gamma^{11}, \partial\} = 0$$

$$Z^{IIB} = \frac{(\chi_V - \chi_S)(\bar{\chi}_V - \bar{\chi}_S)}{(\sqrt{\tau_2} \eta \bar{\eta})^8}$$

$$Z^{IIB} = \frac{1}{(\sqrt{\tau_2} \eta \bar{\eta})^8} \frac{1}{2} \sum_{a,b=0}^1 (-1)^{a+b+ab} \frac{1}{2} \sum_{\bar{a},\bar{b}=0}^1 (-1)^{\bar{a}+\bar{b}+\bar{a}\bar{b}} \frac{\vartheta^4 \begin{bmatrix} a \\ b \end{bmatrix} \bar{\vartheta}^4 \begin{bmatrix} \bar{a} \\ \bar{b} \end{bmatrix}}{\eta^4 \bar{\eta}^4},$$

$$Z^{IIA} = \frac{1}{(\sqrt{\tau_2} \eta \bar{\eta})^8} \frac{1}{2} \sum_{a,b=0}^1 (-1)^{a+b} \frac{1}{2} \sum_{\bar{a},\bar{b}=0}^1 (-1)^{\bar{a}+\bar{b}+\bar{a}\bar{b}} \frac{\vartheta^4 \begin{bmatrix} a \\ b \end{bmatrix} \bar{\vartheta}^4 \begin{bmatrix} \bar{a} \\ \bar{b} \end{bmatrix}}{\eta^4 \bar{\eta}^4}.$$



Estados supermasivos y pluridimensiones por deformación de superficie. Métrica de Majorana.

$$\{\Gamma^\mu, \Gamma^\nu\} = -2\eta^{\mu\nu}, \eta^{\mu\nu} = (- + + \dots +)$$

$$\Gamma_\mu = \eta_{\mu\nu} \Gamma^\nu, \Gamma^\mu = \eta^{\mu\nu} \Gamma_\nu$$

$$\Gamma^0 \Gamma_\mu^\dagger \Gamma^0 = \Gamma_\mu, \Gamma^0 \Gamma_\mu \Gamma^0 = -\Gamma_\mu^T$$

$$\Gamma_{11} = \Gamma_0 \dots \Gamma_9, (\Gamma_{11})^2 = 1, \{\Gamma_{11}, \Gamma^\mu\} = 0;$$

$$\Gamma^{\mu_1 \dots \mu_k} = \frac{1}{k!} \Gamma^{[\mu_1} \dots \Gamma^{\mu_k]} = \frac{1}{k!} (\Gamma^{\mu_1} \dots \Gamma^{\mu_k} \pm \text{permutations}).$$

$$\Gamma_{11} \Gamma^{\mu_1 \dots \mu_k} = \frac{(-1)^{\lfloor \frac{k}{2} \rfloor}}{(10-k)!} \epsilon^{\mu_1 \dots \mu_{10}} \Gamma_{\mu_{k+1} \dots \mu_{10}}$$

$$\Gamma^{\mu_1 \dots \mu_k} \Gamma_{11} = \frac{(-1)^{\lfloor \frac{k+1}{2} \rfloor}}{(10-k)!} \epsilon^{\mu_1 \dots \mu_{10}} \Gamma_{\mu_{k+1} \dots \mu_{10}}$$

$$\Gamma^\mu \Gamma^{\nu_1 \dots \nu_k} = \Gamma^{\mu \nu_1 \dots \nu_k} - \frac{1}{(k-1)!} \eta^{\mu[\nu_1} \Gamma^{\nu_2 \dots \nu_k]}$$

$$\Gamma^{\nu_1 \dots \nu_k} \Gamma^\mu = \Gamma^{\nu_1 \dots \nu_k \mu} - \frac{1}{(k-1)!} \eta^{\mu[\nu_k} \Gamma^{\nu_1 \dots \nu_{k-1}]}$$

$$F_{\alpha\beta} = S_\alpha(i\Gamma^0)_{\beta\gamma} \tilde{S}_\gamma$$

$$\Gamma_{11} F = F, F \Gamma_{11} = -\xi F$$

$$F_{\alpha\beta} = \sum_{k=0}^{10} \frac{i^k}{k!} F_{\mu_1 \dots \mu_k} (\Gamma^{\mu_1 \dots \mu_k})_{\alpha\beta}$$

$$F^{\mu_1 \dots \mu_k} = \frac{(-1)^{\lfloor \frac{k+1}{2} \rfloor}}{(10-k)!} \epsilon^{\mu_1 \dots \mu_{10}} F_{\mu_{k+1} \dots \mu_{10}}$$

$$F^{\mu_1 \dots \mu_k} = \xi \frac{(-1)^{\lfloor \frac{k}{2} \rfloor + 1}}{(10-k)!} \epsilon^{\mu_1 \dots \mu_{10}} F_{\mu_{k+1} \dots \mu_{10}}$$

$$(p_\mu \Gamma^\mu) F = F (p_\mu \Gamma^\mu) = 0$$

$$p^{[\mu} F^{\nu_1 \dots \nu_k]} = p_\mu F^{\mu \nu_2 \dots \nu_k} = 0$$

$$dF = d^* F = 0$$

$$F_{\mu_1 \dots \mu_k} = \frac{1}{(k-1)!} \partial_{[\mu_1} C_{\mu_2 \dots \mu_k]}$$

$$F_{(k)} = dC_{(k-1)}$$

$$X(\sigma + 2\pi) = X(2\pi - \sigma)$$



$$Z_{\text{compact}}(\bar{q}) = \sum_{L_{16}} \frac{\bar{q}^{\frac{p_R^2}{2}}}{\bar{\eta}^{16}} = \frac{\bar{\Gamma}_{16}(\bar{q})}{\bar{\eta}^{16}}$$

$$Z^{\text{heterotic}} = \frac{1}{(\sqrt{\tau_2}\eta\bar{\eta})^8} \frac{\bar{\Gamma}_{16}}{\bar{\eta}^{16}} \frac{1}{2} \sum_{a,b=0}^1 (-1)^{a+b+ab} \frac{\vartheta\left[\begin{smallmatrix} a \\ b \end{smallmatrix}\right]^4}{\eta^4}$$

$$\tau \rightarrow -\frac{1}{\tau}: \bar{\Gamma}_{16} \rightarrow \bar{\tau}^8 \bar{\Gamma}_{16}$$

$$\bar{\Gamma}_{E_8 \times E_8} = (\bar{\Gamma}_8)^2 = \left[\frac{1}{2} \sum_{a,b=0,1} \bar{\vartheta} \left[\begin{smallmatrix} a \\ b \end{smallmatrix} \right]^8 \right] = 1 + 2 \cdot 240 \bar{q} + \mathcal{O}(\bar{q}^2).$$

$$\bar{\Gamma}_{O(32)/Z_2} = \frac{1}{2} \sum_{a,b=0,1} \bar{\vartheta} \left[\begin{smallmatrix} a \\ b \end{smallmatrix} \right]^{16} = 1 + 480 \bar{q} + \mathcal{O}(\bar{q}^2)$$

$$\bar{J}^{ij} = i\bar{\psi}^i\bar{\psi}^j$$

$$L_0=\frac{1}{2}, \bar{L}_0=1$$

$$G_0=0\,,\bar{L}_0=1$$

$$Z_{\text{fermions}}\left[\begin{smallmatrix} h \\ g \end{smallmatrix}\right] = \frac{1}{2} \sum_{a,b=0}^1 (-1)^{a+b+ab+ag+bh+gh} \frac{\vartheta^4\left[\begin{smallmatrix} a \\ b \end{smallmatrix}\right]}{\eta^4}$$

$$\bar{Z}_{E_8}\left[\begin{smallmatrix} h \\ g \end{smallmatrix}\right] = \frac{1}{2} \sum_{\gamma,\delta=0}^1 (-1)^{\gamma g+\delta h} \frac{\bar{\vartheta}^8\left[\begin{smallmatrix} \gamma \\ \delta \end{smallmatrix}\right]}{\bar{\eta}^8}$$

$$Z^{\text{heterotic}}_{O(16)\times O(16)} = \frac{1}{2} \sum_{h,g=0}^1 \frac{\bar{Z}_{E_8}\left[\begin{smallmatrix} h \\ g \end{smallmatrix}\right]^2}{(\sqrt{\tau_2}\eta\bar{\eta})^8} \frac{1}{2} \sum_{a,b=0}^1 (-1)^{a+b+ab+ag+bh+gh} \frac{\vartheta^4\left[\begin{smallmatrix} a \\ b \end{smallmatrix}\right]}{\eta^4}.$$

$$\hat{X}^\mu(z,\theta)=X^\mu(z)+\theta\psi^\mu(z)$$

$$\int \, dz \int \, d\theta V(z,\theta) = \int \, dz \int \, d\theta (V_0(z) + \theta V_{-1}(z)) = \int \, dz V_{-1}$$

$$V^{\text{boson}}(\epsilon,p,z,\theta)=\epsilon_\mu\colon D\hat{X}^\mu e^{ip\cdot\hat{X}}$$

$$V_0^{\text{boson}}=\epsilon_\mu\psi^\mu e^{ip\cdot X}, V_{-1}^{\text{boson}}(\epsilon,p,z)=\epsilon_\mu\colon(\partial X^\mu+ip\cdot\psi\psi^\mu)e^{ip\cdot X}\colon,$$

$$V_{-1}^{\text{boson}}(\epsilon,p,z) = \left[Q_{\text{BRST}}, \xi(z) e^{-\phi(z)} \epsilon \cdot \psi e^{ip \cdot X} \right]$$

$$V_{-1/2}^{\text{fermion}}(u,p,z)=u^\alpha(p)\colon e^{-\phi(z)/2}S_\alpha(z)e^{ip\cdot X}\colon,$$

$$V_{1/2}^{\text{fermion}}(u, p) = [Q_{\text{BRST}}, \xi(z) V_{-1/2}^{\text{fermion}}(u, p, z)] = u^\alpha(p) e^{\phi/2} S_\alpha e^{ip \cdot X} + \dots$$

$$Q_\alpha = \frac{1}{2\pi i} \oint dz: e^{-\phi(z)/2} S_\alpha(z)$$

$$[Q_\alpha, V_{-1/2}^{\text{fermion}}(u, p, z)] = V_{-1}^{\text{boson}}(\epsilon^\mu = u^\beta \gamma_{\beta\alpha}^\mu, p, z),$$

$$[Q_\alpha, V_0^{\text{boson}}(\epsilon, p, z)] = V_{-1/2}^{\text{fermion}}(u^\beta = ip^\mu \epsilon^\nu (\gamma_{\mu\nu})_\alpha^\beta, p, z)$$

Acciones supersimétricas y de supergravedad. Métrica de Yang – Mills - Chern-Simons.

Superinvariancias y supercovariancias.

$$\delta_\epsilon \phi \sim \phi^m \psi \epsilon, \delta_\epsilon \psi \sim \partial \phi^m \epsilon + \phi^m \psi^2 \epsilon$$

$$L_{\text{YM}} = -\frac{1}{4} F_{\mu\nu}^a F^{a,\mu\nu} - \bar{\chi}^a \Gamma^\mu D_\mu \chi^a$$

$$F_{\mu\nu}^a = \partial_\mu A_\nu^a - \partial_\nu A_\mu^a + g f_{bc}^a A_\mu^b A_\nu^c$$

$$D_\mu \chi^a = \partial_\mu \chi^a + g f_{bc}^a A_\mu^b \chi^c$$

$$L_{\text{SUGRA}}^{N=1} = -\frac{1}{2\kappa^2} R - \frac{3}{4} \phi^{-\frac{3}{2}} H_{\mu\nu\rho} H^{\mu\nu\rho} - \frac{9}{16\kappa^2} \frac{\partial_\mu \phi \partial^\mu \phi}{\phi^2} - \frac{1}{2} \bar{\psi}^\mu \Gamma^{\mu\nu\rho} \nabla_\nu \psi_\rho - \frac{1}{2} \bar{\lambda} \Gamma^\mu \nabla_\mu \lambda$$

$$- \frac{3\sqrt{2}}{8} \frac{\partial_\nu}{\phi} \bar{\psi}^\mu \Gamma^\nu \Gamma^\mu \lambda$$

$$+ \frac{\sqrt{2}\kappa}{16} \phi^{-3/4} H_{\nu\rho\sigma} [\bar{\psi}_\mu \Gamma^{\mu\nu\rho\sigma\tau} \psi_\tau + 6\bar{\psi}^\nu \Gamma^\rho \psi^\sigma - \sqrt{2} \bar{\psi}_\mu \Gamma^{\nu\rho\sigma} \Gamma^\mu \lambda] + (\mathbb{F}_{\text{fermion}})^4$$

$$L_{\text{SUGRA+YM}}^{N=1} = L_{\text{SUGRA}}^{N=1} + \phi^{-3/4} L'_{\text{YM}}$$

$$\hat{H}_{\mu\nu\rho} = H_{\mu\nu\rho} - \frac{\kappa}{\sqrt{2}} \omega_{\mu\nu\rho}^{\text{CS}}$$

$$\omega_{\mu\nu\rho}^{\text{CS}} = A_\mu^a F_{\nu\rho}^a - \frac{g}{3} f_{abc} A_\mu^a A_\nu^b A_\rho^c + \text{cyclic}$$

$$\delta B = \frac{\kappa}{\sqrt{2}} \text{Tr}[\Lambda dA]$$

$$L^{D=11} = \frac{1}{2\kappa^2} \left[R - \frac{1}{2 \cdot 4!} G_4^2 \right] - i \bar{\psi}_\mu \Gamma^{\mu\nu\rho} \tilde{\nabla}_\nu \psi_\rho + \frac{1}{2\kappa^2 (144)^2} G_4 \wedge G_4 \wedge \hat{C}$$

$$+ \frac{1}{192} [\bar{\psi}_\mu \Gamma^{\mu\nu\rho\sigma\tau\nu} \psi_\nu + 12 \bar{\psi}^\nu \Gamma^{\rho\sigma} \psi^\tau] (G + \hat{G})_{\nu\rho\sigma\tau}$$

$$\tilde{\omega}_{\mu,ab} = \omega_{\mu,ab} + \frac{i\kappa^2}{4} [-\bar{\psi}^\nu \Gamma_{\nu\mu ab\rho} \psi^\rho + 2(\bar{\psi}_\mu \Gamma_b \psi_a - \bar{\psi}_\mu \Gamma_a \psi_b + \bar{\psi}_b \Gamma_\mu \psi_a)]$$

$$G_{\mu\nu\rho\sigma} = \partial_\mu \hat{C}_{\nu\rho\sigma} - \partial_\nu \hat{C}_{\rho\sigma\mu} + \partial_\rho \hat{C}_{\sigma\mu\nu} - \partial_\sigma \hat{C}_{\mu\nu\rho}$$



$$\tilde{G}_{\mu\nu\rho\sigma} = G_{\mu\nu\rho\sigma} - 6\kappa^2 \bar{\psi}_{[\mu} \Gamma_{\nu\rho} \psi_{\sigma]}$$

$$G_{\mu\nu} = \begin{pmatrix} g_{\mu\nu} + e^{2\sigma} A_\mu A_\nu & e^{2\sigma} A_\mu \\ e^{2\sigma} A_\mu & e^{2\sigma} \end{pmatrix}$$

$$C_{\mu\nu\rho} = \hat{C}_{\mu\nu\rho} - (\hat{C}_{\nu\rho,11} A_\mu + \text{cyclic}) , B_{\mu\nu} = \hat{C}_{\mu\nu,11}$$

$$S^{IIA} = \frac{1}{2\kappa^2} \int d^{10}x \sqrt{g} e^\sigma \left[R - \frac{1}{2 \cdot 4!} \hat{G}^2 - \frac{1}{2 \cdot 3!} e^{-2\sigma} H^2 - \frac{1}{4} e^{2\sigma} F^2 \right] + \frac{1}{2\kappa^2 (48)^2} \int B \wedge G \wedge G$$

$$F_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu , H_{\mu\nu\rho} = \partial_\mu B_{\nu\rho} + \text{cíclico}$$

$$\hat{G}_{\mu\nu\rho\sigma} = G_{\mu\nu\rho\sigma} + (F_{\mu\nu} B_{\rho\sigma} + 5 \text{ permutaciones}).$$

$$\tilde{S}_{10} = \frac{1}{2\kappa^2} \int d^{10}x \sqrt{g} e^{-\Phi} \left[\left(R + (\nabla\Phi)^2 - \frac{1}{12} H^2 \right) - \frac{1}{2 \cdot 4!} \hat{G}^2 - \frac{1}{4} F^2 \right] + \frac{1}{2\kappa^2 (48)^2} \int B \wedge G \wedge G$$

$$S=\chi+ie^{-\phi/2}$$

$$S \rightarrow \frac{aS+b}{cS+d}, \begin{pmatrix} B_{\mu\nu}^N \\ B_{\mu\nu}^R \end{pmatrix} \rightarrow \begin{pmatrix} d & -c \\ -b & a \end{pmatrix} \begin{pmatrix} B_{\mu\nu}^N \\ B_{\mu\nu}^R \end{pmatrix}$$

$$S^{IIB} = \frac{1}{2\kappa^2} \int d^{10}x \sqrt{-\det g} \left[R - \frac{1}{2} \frac{\partial S \partial \bar{S}}{S_2^2} - \frac{1}{12} \frac{|H^R + S H^N|^2}{S_2} \right]$$

$$J^\mu = \frac{\delta \Gamma^{\rm eff}}{\delta A_\mu}, T^{\mu\nu} = \frac{1}{\sqrt{-g}} \frac{\delta \Gamma^{\rm eff}}{\delta g_{\mu\nu}}$$

$$\delta_\Lambda \Gamma^{\rm eff} = \text{Tr} \int D_\mu \Lambda \frac{\delta \Gamma^{\rm eff}}{\delta A_\mu} = \text{Tr} \int \Lambda D_\mu \frac{\delta \Gamma^{\rm eff}}{\delta A_\mu} = \int \text{Tr} [\Lambda D_\mu J^\mu]$$

$$\delta_{diff} \Gamma^{\rm eff} = \int (\nabla^\mu \epsilon^\nu + \nabla^\nu \epsilon^\mu) \frac{\delta \Gamma^{\rm eff}}{\delta g_{\mu\nu}} = \int \epsilon^\mu \nabla_\nu T^{\mu\nu}$$

$$\delta\Gamma|_{\rm gauge} \sim \int d^{10}x [c_1 {\rm Tr}[\Lambda F_0^5] + c_2 {\rm Tr}[\Lambda F_0] {\rm Tr}[F_0^4] + c_3 {\rm Tr}[\Lambda F_0] ({\rm Tr}[F_0^2])^2]$$

$$\delta\Gamma|_{\rm grav} \sim \int d^{10}x [d_1 {\rm Tr}[\Theta R_0^5] + d_2 {\rm Tr}[\Theta R_0] {\rm Tr}[R_0^4] + d_3 {\rm Tr}[\Theta R_0] ({\rm Tr}[R_0^2])^2]$$

$$\delta\Gamma|_{\rm mixed} \sim \int d^{10}x [e_1 {\rm Tr}[\Lambda F_0] {\rm Tr}[R_0^4] + e_2 {\rm Tr}[\Theta R_0] {\rm Tr}[F_0^4] \\ + e_3 {\rm Tr}[\Theta R_0] ({\rm Tr}[F_0^2])^2 + e_4 {\rm Tr}[\Lambda F_0] ({\rm Tr}[R_0])^2]$$

- Wess-Zumino:

$$\delta_{\Lambda_1} G(\Lambda_2) - \delta_{\Lambda_2} G(\Lambda_1) = G([\Lambda_1, \Lambda_2])$$



$$\delta F = [F, \Lambda], \delta R = [R, \Theta]$$

$$d\text{Tr}[R^m] = d\text{Tr}[F^m] = 0$$

$$I^{D+2}(R, F) = d\Omega^{D+1}(\omega, A)$$

$$\delta_\Lambda \Omega^{D+1}(\omega, A) = d\Omega^D(\omega, A, \Lambda)$$

- Mecanismo de Green-Schwarz:

$$\begin{aligned} \delta\Gamma|_{\text{reduc}} \sim \int d^{10}x & (\text{Tr}[\Lambda F_0] + \text{Tr}[\Theta R_0]) (a_1 \text{Tr}[F_0^4] \\ & + a_2 \text{Tr}[R_0^4] + a_3 (\text{Tr}[F_0^2])^2 + a_4 (\text{Tr}[R_0^2])^2 + a_5 \text{Tr}[F_0^2] \text{Tr}[R_0^2]) \end{aligned}$$

$$\hat{H} = dB + \Omega^{CS}(A) + \Omega^{CS}(\omega)$$

$$\delta_\Lambda \Omega^{CS}(A) = d\text{Tr}[\Lambda dA], \delta_\Theta \Omega^{CS}(\omega) = d\text{Tr}[\Theta d\omega]$$

$$\delta B = -\text{Tr}[\Lambda F_0 + \Theta R_0]$$

$$\begin{aligned} \Gamma_{\text{counter}} \sim \int d^{10}x & B (a_1 \text{Tr}[F_0^4] + a_2 \text{Tr}[R_0^4] + a_3 (\text{Tr}[F_0^2])^2 + \\ & + a_4 (\text{Tr}[R_0^2])^2 + a_5 \text{Tr}[F_0^2] \text{Tr}[R_0^2]) \end{aligned}$$

$$F_{\mu_1 \dots \mu_{D/2}} = \pm \frac{i}{(D/2)!} \epsilon_{\mu_1 \dots \mu_D} F^{\mu_{D/2+1} \dots \mu_D}$$

$$R_0 = \begin{pmatrix} 0 & x_1 & 0 & 0 & \dots & \\ -x_1 & 0 & 0 & 0 & \dots & \\ 0 & 0 & 0 & x_2 & \dots & \\ 0 & 0 & -x_2 & 0 & \dots & \\ \dots & \dots & \dots & \dots & \dots & \\ \dots & & & & 0 & x_{D/2} \\ \dots & & & & -x_{D/2} & 0 \end{pmatrix}$$

$$\hat{I}_{1/2}(R) = \prod_{i=1}^{D/2} \left(\frac{x_i/2}{\sinh(x_i/2)} \right)$$

$$I_{3/2}(R) = \hat{I}_{1/2}(R) \left(-1 + 2 \sum_{i=1}^{D/2} \cosh(x_i) \right)$$

$$I_A(R) = -\frac{1}{8} \prod_{i=1}^{D/2} \left(\frac{x_i}{\tanh(x_i)} \right)$$

$$I_{1/2}(R, F) = \text{Tr}[e^{iF}] \hat{I}_{1/2}(R)$$



$$\begin{aligned}
I_{1/2}(R, F)|_{12-\text{form}} &= -\frac{\text{Tr}[F^6]}{720} + \frac{\text{Tr}[F^4]\text{Tr}[R^2]}{24 \cdot 48} + \\
&\quad -\frac{\text{Tr}[F^2]}{256} \left(\frac{\text{Tr}[R^4]}{45} + \frac{(\text{Tr}[R^2])^2}{36} \right) + \\
&\quad + \frac{n}{64} \left(\frac{\text{Tr}[R^6]}{5670} + \frac{\text{Tr}[R^2]\text{Tr}[R^4]}{4320} + \frac{(\text{Tr}[R^2])^3}{10368} \right) \\
I_{3/2}(R)|_{12-\text{form}} &= -\frac{495}{64} \left(\frac{\text{Tr}[R^6]}{5670} + \frac{\text{Tr}[R^2]\text{Tr}[R^4]}{4320} + \frac{(\text{Tr}[R^2])^3}{10368} \right) + \\
&\quad + \frac{\text{Tr}[R^2]}{384} \left(\text{Tr}[R^4] + \frac{(\text{Tr}[R^2])^2}{4} \right)
\end{aligned}$$

$$I_A(R)|_{12-\text{form}} = \hat{I}_{1/2}(R)|_{12-\text{form}} - I_{3/2}(R)|_{12-\text{form}}$$

$$2I^{N=1} = I_{3/2}(R) - I_{1/2}(R) + I_{1/2}(R, F)$$

$$\begin{aligned}
96I^{\text{total}} &= -\frac{\text{Tr}[F^6]}{15} + \frac{\text{Tr}[R^2]\text{Tr}[F^4]}{24} + \frac{\text{Tr}[R^2]\text{Tr}[R^4]}{8} + \frac{(\text{Tr}[R^2])^3}{32} - \\
&\quad - \frac{\text{Tr}[F^2]}{960} (4\text{Tr}[R^4] + 5(\text{Tr}[R^2])^2)
\end{aligned}$$

$$\text{Tr}[F^6] = \frac{1}{48} \text{Tr}[F^2]\text{Tr}[F^4] - \frac{1}{14400} (\text{Tr}[F^2])^3.$$

$$96I^{\text{total}} = \left(\text{Tr}[R^2] - \frac{1}{30} \text{Tr}[F^2] \right) X_8$$

$$X_8 = \frac{\text{Tr}[F^4]}{24} - \frac{(\text{Tr}[F^2])^2}{720} - \frac{\text{Tr}[F^2]\text{Tr}[R^2]}{240} + \frac{\text{Tr}[R^4]}{8} + \frac{(\text{Tr}[R^2])^2}{32},$$

$$\begin{aligned}
\text{Tr}[F^6] &= (N-32)\text{tr}[F^6] + 15\text{tr}[F^2]\text{tr}[F^4] \\
\text{Tr}[F^4] &= (N-8)\text{tr}[F^4] + 3(\text{tr}[F^2])^2, \text{Tr}[F^2] = (N-2)\text{tr}[F^2]
\end{aligned}$$

$$\text{Tr}[F^6] = \frac{1}{7200} (\text{Tr}[F^2])^3, \text{Tr}[F^4] = \frac{1}{100} (\text{Tr}[F^2])^2$$

$$d\hat{H} = \text{tr}[R^2] - \frac{1}{30} \text{Tr}[F^2]$$

$$\int \text{tr}[R^2] = \frac{1}{30} \int \text{Tr}[F^2]$$

$$\begin{aligned}
\text{tr}_S[F^6] &= 16\text{tr}[F^6] - 15\text{tr}[F^2]\text{tr}[F^4] + \frac{15}{4} (\text{tr}[F^2])^3 \\
\text{tr}_S[F^4] &= -8\text{tr}[F^4] + 6(\text{tr}[F^2])^2, \text{tr}_S[F^2] = 16\text{tr}[F^2]
\end{aligned}$$

Compactificación y rompimiento de supersimetría. Métrica de Kaluza-Klein.

$$Z_D^{\text{heterotic}} = \frac{\Gamma_{10-D, 10-D}(G, B) \bar{\Gamma}_H}{\tau_2^{\frac{D-2}{2}} \eta^8 \bar{\eta}^8} \frac{1}{2} \sum_{a,b=0}^1 (-1)^{a+b+ab} \frac{\vartheta^4 \begin{bmatrix} a \\ b \end{bmatrix}}{\eta^4},$$



$$V_{\text{Higgs}} \sim f^a{}_{bc} f^a{}_{b'c'} G^{\alpha\gamma} G^{\beta\delta} Y^b_\alpha Y^c_\beta Y^{b'}_{\gamma} Y^{c'}_{\delta}$$

$$[m^2]^{ab} \sim G^{\alpha\beta} f_d^{ca} f_d^{cb'} Y^d_\alpha Y^{a'}_\beta$$

$$Z_D^{\text{heterotic}} = \frac{\Gamma_{10-D,26-D}(G,B,Y)}{\tau_2^{\frac{D-2}{2}} \eta^8 \bar{\eta}^8} \frac{1}{2} \sum_{a,b=0}^1 (-1)^{a+b+ab} \frac{\vartheta^4 \left[\begin{smallmatrix} a \\ b \end{smallmatrix} \right]}{\eta^4},$$

$$\alpha'^8 S_{10-d}^{\text{heterotic}} = \int d^{10}x \sqrt{-\det G_{10}} e^{-\Phi} \left[R + (\nabla \Phi)^2 - \frac{1}{12} \hat{H}^2 - \frac{1}{4} \text{Tr}[F^2] \right] + \mathcal{O}(\alpha')$$

$$F^I_{\mu\nu} = \partial_\mu A^I_\nu - \partial_\nu A^I_\mu$$

$$\hat{H}_{\mu\nu\rho} = \partial_\mu B_{\nu\rho} - \frac{1}{2} \sum_I A^I_\mu F^I_{\nu\rho} + \text{cyclic}$$

$$S_D^{\text{heterotic}} = \int d^D x \sqrt{-\det G} e^{-\Phi} \left[R + \partial^\mu \Phi \partial_\mu \Phi - \frac{1}{12} \hat{H}^{\mu\nu\rho} \hat{H}_{\mu\nu\rho} - \right. \\ \left. - \frac{1}{4} (\hat{M}^{-1})_{ij} F^i_{\mu\nu} F^{j\mu\nu} + \frac{1}{8} \text{Tr}(\partial_\mu \hat{M} \partial^\mu \hat{M}^{-1}) \right]$$

$$\hat{H}_{\mu\nu\rho} = \partial_\mu B_{\nu\rho} - \frac{1}{2} L_{ij} A^i_\mu F^j_{\nu\rho} + \text{cyclic}$$

$$\hat{M} \rightarrow \Omega \hat{M} \Omega^T, A_\mu \rightarrow \Omega \cdot A_\mu$$

$$S_D^{\text{heterotic}} = \int d^D x \sqrt{-\det G_E} \left[R - \frac{1}{D-2} \partial^\mu \Phi \partial_\mu \Phi \right. \\ \left. - \frac{e^{-\frac{4\Phi}{D-2}}}{12} \hat{H}^{\mu\nu\rho} \hat{H}_{\mu\nu\rho} - \frac{e^{-\frac{2\Phi}{D-2}}}{4} (\hat{M}^{-1})_{ij} F^i_{\mu\nu} F^{j\mu\nu} + \frac{1}{8} \text{Tr}(\partial_\mu \hat{M} \partial^\mu \hat{M}^{-1}) \right]$$

$$e^{-2\phi} \hat{H}_{\mu\nu\rho} = \frac{\epsilon^{\sigma}_{\mu\nu\rho}}{\sqrt{-\det g_E}} \nabla_\sigma a$$

$$\frac{\epsilon^{\mu\nu\rho\sigma}}{\sqrt{-\det g_E}} \partial_\mu \hat{H}_{\nu\rho\sigma} = -L_{ij} F^i_{\mu\nu} \tilde{F}^{j,\mu\nu},$$

$$\tilde{F}^{\mu\nu} = \frac{1}{2} \frac{\epsilon^{\mu\nu\rho\sigma}}{\sqrt{-\det g_E}} F_{\rho\sigma},$$

$$\nabla^\mu e^{2\phi} \nabla_\mu a = -\frac{1}{4} F^i_{\mu\nu} \tilde{F}^{j,\mu\nu}$$

$$\tilde{S}_{D=4}^{\text{heterotic}} = \int d^4x \sqrt{-\det g_E} \left[R - \frac{1}{2} \partial^\mu \phi \partial_\mu \phi - \frac{1}{2} e^{2\phi} \partial^\mu a \partial_\mu a \pm \frac{1}{4} e^{-\phi} (M^{-1})_{ij} F_{\mu\nu}^i F^{j,\mu\nu} \right. \\ \left. + \frac{1}{4} a L_{ij} F_{\mu\nu}^i \tilde{F}^{j,\mu\nu} + \frac{1}{8} \text{Tr}(\partial_\mu M \partial^\mu M^{-1}) \right]$$

$$S = a + i e^{-\phi}$$

$$\tilde{S}_{D=4}^{\text{heterotic}} = \int d^4x \sqrt{-\det g_E} \left[R - \frac{1}{2} \frac{\partial^\mu S \partial_\mu \bar{S}}{\text{Im} S^2} - \frac{1}{4} \text{Im} S (M^{-1})_{ij} F_{\mu\nu}^i F^{j,\mu\nu} + \frac{1}{4} \text{Re} S L_{ij} F_{\mu\nu}^i \tilde{F}^{j,\mu\nu} \right. \\ \left. + \frac{1}{8} \text{Tr}(\partial_\mu M \partial^\mu M^{-1}) \right]$$

$$\delta\psi_\mu = \nabla_\mu \epsilon + \frac{\sqrt{2}}{32} e^{2\Phi} (\gamma_\mu \gamma_5 \otimes H) \epsilon,$$

$$\delta\psi_m = \nabla_m \epsilon + \frac{\sqrt{2}}{32} e^{2\Phi} (\gamma_m H - 12 H_m) \epsilon,$$

$$\delta\lambda = \sqrt{2} (\gamma^m \nabla_m \Phi) \epsilon + \frac{1}{8} e^{2\Phi} H \epsilon,$$

$$\delta\chi^a = -\frac{1}{4} e^\Phi F_{m,n}^a \gamma^{mn} \epsilon,$$

$$H = H_{mnr} \gamma^{mnr}, H_m = H_{mnr} \gamma^{nr}$$

$$\Gamma^\mu = \gamma^\mu \otimes \mathbf{1}_6, \Gamma^m = \gamma^5 \otimes \gamma^m$$

$$\gamma^5 = \frac{i}{4!} \epsilon_{\mu\nu\rho\sigma} \gamma^{\mu\nu\rho\sigma}, \gamma = \frac{i}{6!} \sqrt{\det g} \epsilon_{mnrpq} \gamma^{mnrpq}$$

$$\nabla_m \xi = 0$$

$$F_{mn}^a \gamma^{mn} \xi = 0$$

$$R_{[mn}^{rs} R_{pq]rs} = \frac{1}{30} F_{[mn}^a F_{pq]}^a$$

$$G_{MN} \sim h_{\mu\nu}(x) \otimes \phi(y) + A_\mu(x) \otimes f_m(y) + \Phi(x) \otimes h_{mn}(y)$$

$$\Box_y \phi(y) = 0$$

- Lichnerowicz:

$$-\Box h_{mn} + 2R_{mnrs} h^{rs} = 0, \nabla^m h_{mn} = g^{mn} h_{mn} = 0.$$

$$\Box f_{mn} - R_{mnrs} f^{rs} = \Box f_{mn} + 2R_{mrsn} f^{rs} = 0, \\ \nabla_m A_{np} + \nabla_p A_{mn} + \nabla_n A_{pm} = 0, \nabla^m A_{mn} = 0$$

$$h_{mn} = A_m^p S_{pm} + A_n^p S_{pm}$$

$$B_{MN} \sim B_{\mu\nu}(x) \otimes \phi(y) + B_\mu(x) \otimes f_m(y) + \Phi(x) \otimes B_{mn}(y)$$



$$\begin{aligned}
S_{K3}^{IIA} = & \int d^6x \sqrt{-\det G_6} e^{-\Phi} \left[R + \nabla^\mu \Phi \nabla_\mu \Phi - \frac{1}{12} H^{\mu\nu\rho} H_{\mu\nu\rho} \right. \\
& + \frac{1}{8} \text{Tr}(\partial_\mu \hat{M} \partial^\mu \hat{M}^{-1}) + \frac{1}{4} \int d^6x \sqrt{-\det G} (\hat{M}^{-1})_{IJ} F_{\mu\nu}^I F^{J\mu\nu} \Big] \\
& + \frac{1}{16} \int d^6x \epsilon^{\mu\nu\rho\sigma\tau\nu} B_{\mu\nu} F_{\rho\sigma}^I \hat{L}_{IJ} F_{\tau\nu}^J
\end{aligned}$$

Espacio – tiempo cuántico supersimétrico.

$$\begin{aligned}
\{Q_\alpha^I, Q_\beta^J\} &= \epsilon_{\alpha\beta} Z^{IJ} \\
\{\bar{Q}_\alpha^I, Q_\beta^J\} &= \epsilon_{\dot{\alpha}\dot{\beta}} \bar{Z}^{IJ} \\
\{Q_\alpha^I, \bar{Q}_{\dot{\alpha}}^J\} &= \delta^{IJ} \sigma_{\alpha\dot{\alpha}}^\mu P_\mu
\end{aligned}$$

$$Q_\alpha^I = \frac{1}{2\pi i} \oint dz e^{-\phi/2} S_\alpha \Sigma^I, \bar{Q}_{\dot{\alpha}}^I = \frac{1}{2\pi i} \oint dz e^{-\phi/2} C_{\dot{\alpha}} \bar{\Sigma}^I$$

$$:e^{q_1\phi(z)}::e^{q_2\phi(w)}:= (z-w)^{-q_1q_2}:e^{(q_1+q_2)\phi(w)}: + \dots,$$

$$S_\alpha(z)C_{\dot{\alpha}}(w) = \sigma_{\alpha\dot{\alpha}}^\mu \psi^\mu(w) + \mathcal{O}(z-w),$$

$$S_\alpha(z)S_\beta(w) = \frac{\epsilon_{\alpha\beta}}{\sqrt{z-w}} + \mathcal{O}(\sqrt{z-w}),$$

$$C_{\dot{\alpha}}(z)C_{\dot{\beta}}(w) = \frac{\epsilon_{\dot{\alpha}\dot{\beta}}}{\sqrt{z-w}} + \mathcal{O}(\sqrt{z-w}),$$

$$\Sigma^I(z)\bar{\Sigma}^J(w) = \frac{\delta^{IJ}}{(z-w)^{3/4}} + (z-w)^{1/4} J^{IJ}(w) + \dots$$

$$\Sigma^I(z)\Sigma^J(w) = (z-w)^{-1/4} \Psi^{IJ}(w) + \dots$$

$$\bar{\Sigma}^I(z)\bar{\Sigma}^J(w) = (z-w)^{-1/4} \bar{\Psi}^{IJ}(w) + \dots$$

$$G^{\text{int}}(z)\Sigma^I(w) \sim (z-w)^{-1/2}, G^{\text{int}}(z)\bar{\Sigma}^I(w) \sim (z-w)^{-1/2}$$

$$\lambda^I(z) = \Sigma^I(z)e^{iX/2}, \bar{\lambda}^I(z) = \bar{\Sigma}^I(z)e^{-iX/2}$$

$$\lambda^I(z)\bar{\lambda}^J(w) = \frac{\delta^{IJ}}{z-w} + \hat{J}^{IJ} + \mathcal{O}(z-w),$$

$$\lambda^I(z)\lambda^J(w) = e^{iX}\Psi^{IJ} + \mathcal{O}(z-w),$$

$$\bar{\lambda}^I(z)\bar{\lambda}^J(w) = e^{-iX}\bar{\Psi}^{IJ} + \mathcal{O}(z-w),$$

$$\hat{J}^{II}(z)\hat{J}^{JJ}(w) = \frac{\delta^{IJ}}{(z-w)^2} + \text{regular}$$

$$J^{II}(z)J^{JJ}(w) = \frac{\delta^{IJ} - 3/4}{(z-w)^2} + \text{regular}$$

$$J = 2J^{11}, J(z)J(w) = \frac{3}{(z-w)^2} + \text{regular}$$



$$\langle J(z_1)\Sigma(z_2)\bar{\Sigma}(z_3)\rangle=\frac{3}{2}\frac{z_{23}^{1/4}}{z_{12}z_{13}}$$

$$J=i\sqrt{3}\partial\Phi,\Sigma=e^{i\sqrt{3}\Phi/2}W^+,\bar{\Sigma}=e^{-i\sqrt{3}\Phi/2}W^-$$

$$G^{int}=\sum_{q\geq 0}e^{iq\Phi}T^{(q)}+e^{-q\Phi}T^{(-q)}$$

$$J(z)G^{\pm}(w)=\pm\frac{G^{\pm}(w)}{(z-w)}+\cdots$$

$$|h,q;\bar{h}\rangle:|0,0;0\rangle,|0,0;1\rangle^I,|1/2,\pm1;1\rangle^i$$

$$J^s(z)J^s(w)=\frac{1}{(z-w)^2}+\cdots$$

$$J^3(z)J^3(w)=\frac{1/2}{(z-w)^2}+\cdots$$

$$J^s(z)J^3(w)=\cdots$$

$$J^s=i\partial\phi\,,J^3=\frac{i}{\sqrt{2}}\partial\chi$$

$$\Sigma^1=\exp\left[\frac{i}{2}\phi+\frac{i}{\sqrt{2}}\chi\right],\Sigma^2=\exp\left[\frac{i}{2}\phi-\frac{i}{\sqrt{2}}\chi\right]$$

$$\bar{\Sigma}^1=\exp\left[-\frac{i}{2}\phi-\frac{i}{\sqrt{2}}\chi\right],\bar{\Sigma}^2=\exp\left[-\frac{i}{2}\phi+\frac{i}{\sqrt{2}}\chi\right].$$

$$G^{\rm int}=G_{(2)}+G_{(4)},G_{(2)}=G_{(2)}^++G_{(2)}^-$$

$$G_{(4)}=G_{(4)}^++G_{(4)}^-,$$

$$J^s(z)G_{(2)}^{\pm}(w)=\pm\frac{G_{(2)}^{\pm}(w)}{z-w}+\cdots,$$

$$J^3(z)G_{(4)}^{\pm}(w)=\pm\frac{1}{2}\frac{G_{(4)}^{\pm}(w)}{z-w}+\cdots,$$

$$J^s(z)G_{(4)}^{\pm}(w)=\text{finite},J^3(z)G_{(2)}^{\pm}(w)=\text{finite}$$

$$G_{(2)}^{\pm}=e^{\pm i\phi}Z^{\pm}.$$

$$\psi^0=\frac{1}{\sqrt{2}}(\psi^3+i\psi^4),\psi^1=\frac{1}{\sqrt{2}}(\psi^5+i\psi^6)$$

$$\psi^2=\frac{1}{\sqrt{2}}(\psi^7+i\psi^8),\psi^3=\frac{1}{\sqrt{2}}(\psi^9+i\psi^{10})$$

$$\langle \psi^I(z)\bar{\psi}^J(w)\rangle=\frac{\delta^{IJ}}{z-w}\,,\langle \psi^I(z)\psi^J(w)\rangle=\langle \bar{\psi}^I(z)\bar{\psi}^J(w)\rangle=0$$

$$J^I(z)=i\partial_z\phi^I(z)\,,\langle\phi^I(z)\phi^J(w)\rangle=-\delta^{IJ}\log{(z-w)}$$

$$\psi^I=:e^{i\phi^I}:,\bar{\psi}^I=:e^{-i\phi^I}:$$



$$V(\epsilon_I) =: \exp \left[\frac{i}{2} \sum_{I=0}^3 \epsilon_I \phi^I \right]:$$

$$\psi^I \rightarrow e^{2\pi i \theta^I} \psi^I, \bar{\psi}^I \rightarrow e^{-2\pi i \theta^I} \bar{\psi}^I$$

$$\phi^I \rightarrow \phi^I + 2\pi \theta^I$$

$$V^{\pm,\epsilon}=\bar{\partial}X^{\pm}V_S(\epsilon)e^{ip\cdot X},X^{\pm}=\frac{1}{\sqrt{2}}(X^3\pm iX^4)$$

$$G^{\rm int}=\sum_{i=5}^{10}\psi^i\partial X^i$$

$$\phi^2\rightarrow\phi^2+\pi\,,\phi^3\rightarrow\phi^3-\pi$$

$$[{\bf 120}]\rightarrow [{\bf 3},{\bf 1},{\bf 1}]\oplus [{\bf 1},{\bf 3},{\bf 1}]\oplus [{\bf 1},{\bf 1},{\bf 66}]\oplus [{\bf 2},{\bf 1},{\bf 12}]\oplus [{\bf 1},{\bf 2},{\bf 12}]$$

$$\in \mathrm{SU}(2)\times \mathrm{SU}(2)\times \mathrm{O}(12)[{\bf 128}]\rightarrow [{\bf 2},{\bf 1},{\bf 32}]\oplus [{\bf 1},\overline{\bf 2},{\bf 32}]$$

$$\in \mathrm{SU}(2)\times \mathrm{SU}(2)\times \mathrm{O}(12)$$

$$\mathrm{E}_8\ni[{\bf 248}]\rightarrow [{\bf 1},{\bf 133}]\oplus [{\bf 3},{\bf 1}]\oplus [{\bf 2},{\bf 56}]\in \mathrm{SU}(2)\times \mathrm{E}_7$$

$$\frac{1}{2}\sum_{a,b=0}^1(-1)^{a+b+ab}\frac{\vartheta^2\left[\begin{smallmatrix}a\\b\end{smallmatrix}\right]\vartheta\left[\begin{smallmatrix}a+h\\b+g\end{smallmatrix}\right]\vartheta\left[\begin{smallmatrix}a-h\\b-g\end{smallmatrix}\right]}{\eta^4}$$

$$Z_{(4,4)}\left[\begin{smallmatrix}0\\0\end{smallmatrix}\right]=\frac{\Gamma_{4,4}}{\eta^4\bar\eta^4},Z_{(4,4)}\left[\begin{smallmatrix}h\\g\end{smallmatrix}\right]=2^4\frac{\eta^2\bar\eta^2}{\vartheta^2\left[\begin{smallmatrix}1-h\\1-g\end{smallmatrix}\right]\bar\vartheta^2\left[\begin{smallmatrix}1-h\\1-g\end{smallmatrix}\right]},(h,g)\neq(0,0)$$

$$\frac{1}{2}\sum_{\gamma,\delta=0}^1\frac{\bar{\vartheta}\left[\begin{smallmatrix}\gamma+h\\\delta+g\end{smallmatrix}\right]\bar{\vartheta}\left[\begin{smallmatrix}\gamma-h\\\delta-g\end{smallmatrix}\right]\bar{\vartheta}^6\left[\begin{smallmatrix}\gamma\\\delta\end{smallmatrix}\right]}{\bar{\eta}^8}$$

$$\begin{aligned} Z_{N=2}^{\rm heterotic} &= \frac{1}{2} \sum_{h,g=0}^1 \frac{\Gamma_{2,2} \bar{\Gamma}_{E_8} Z_{(4,4)} \left[\begin{smallmatrix} h \\ g \end{smallmatrix} \right]}{\tau_2 \eta^4 \bar{\eta}^{12}} \frac{1}{2} \sum_{\gamma,\delta=0}^1 \frac{\bar{\vartheta} \left[\begin{smallmatrix} \gamma+h \\ \delta+g \end{smallmatrix} \right] \bar{\vartheta} \left[\begin{smallmatrix} \gamma-h \\ \delta-g \end{smallmatrix} \right] \bar{\vartheta}^6 \left[\begin{smallmatrix} \gamma \\ \delta \end{smallmatrix} \right]}{\bar{\eta}^8} \\ &\times \frac{1}{2} \sum_{a,b=0}^1 (-1)^{a+b+ab} \frac{\vartheta^2 \left[\begin{smallmatrix} a \\ b \end{smallmatrix} \right] \vartheta \left[\begin{smallmatrix} a+h \\ b+g \end{smallmatrix} \right] \vartheta \left[\begin{smallmatrix} a-h \\ b-g \end{smallmatrix} \right]}{\eta^4} \end{aligned}$$

$$J^a=-\frac{i}{2}\Big[\psi^0\psi^a+\frac{1}{2}\epsilon^{abc}\psi^b\psi^c\Big],\tilde{J}^a=-\frac{i}{2}\Big[\psi^0\psi^a-\frac{1}{2}\epsilon^{abc}\psi^b\psi^c\Big].$$

$$V_{\alpha\beta}^{\pm}=\pm i(\delta_{\alpha\beta}\psi^0\pm i\sigma^a_{\alpha\beta}\psi^a)$$



$$\begin{aligned}
V_{\alpha\gamma}^+(z)V_{\gamma\beta}^+(w) &= V_{\alpha\gamma}^-(z)V_{\gamma\beta}^-(w) = \frac{\delta_{\alpha\beta}}{z-w} - 2\sigma_{\alpha\beta}^a(J^a(w) - \tilde{J}^a(w)) + \mathcal{O}(z-w) \\
V_{\alpha\gamma}^+(z)V_{\gamma\beta}^-(w) &= \frac{3\delta_{\alpha\beta}}{z-w} + 4\sigma_{\alpha\beta}^a\tilde{J}^a(w) + \mathcal{O}(z-w) \\
V_{\alpha\gamma}^-(z)V_{\gamma\beta}^+(w) &= \frac{3\delta_{\alpha\beta}}{z-w} - 4\sigma_{\alpha\beta}^aJ^a(w) + \mathcal{O}(z-w)
\end{aligned}$$

$$Z_{N=2}^{\text{heterotic}} = \frac{1}{2} \sum_{h,g=0}^1 \frac{\Gamma_{2,18}(\epsilon) \begin{bmatrix} h \\ g \end{bmatrix} Z_{(4,4)} \begin{bmatrix} h \\ g \end{bmatrix}}{\tau_2 \eta^4 \bar{\eta}^{20}} \frac{1}{2} \sum_{a,b=0}^1 (-1)^{a+b+ab} \frac{\vartheta^2 \begin{bmatrix} a \\ b \end{bmatrix} \vartheta \begin{bmatrix} a+h \\ b+g \end{bmatrix} \vartheta \begin{bmatrix} a-h \\ b-g \end{bmatrix}}{\eta^4}$$

$$\begin{aligned}
G &= \frac{T_2 - \frac{W_2^I W_2^I}{2U_2}}{U_2} \begin{pmatrix} 1 & U_1 \\ U_1 & |U|^2 \end{pmatrix} \\
B &= \left(T_1 - \frac{W_1^I W_2^I}{2U_2} \right) \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix}
\end{aligned}$$

- Kähler:

$$f=S\left(TU-\frac{1}{2}W^IW^I\right), K=-\log\left(S_2\right)-\log\left[U_2T_2-\frac{1}{2}W_2^IW_2^I\right]$$

$$\begin{aligned}
\tau_2 B_2 = \tau_2 \langle \lambda^2 \rangle &= \Gamma_{2,18} \begin{bmatrix} 0 \\ 1 \end{bmatrix} \frac{\bar{\vartheta}_3^2 \bar{\vartheta}_4^2}{\bar{\eta}^{24}} - \Gamma_{2,18} \begin{bmatrix} 1 \\ 0 \end{bmatrix} \frac{\bar{\vartheta}_2^2 \bar{\vartheta}_3^2}{\bar{\eta}^{24}} - \Gamma_{2,18} \begin{bmatrix} 1 \\ 1 \end{bmatrix} \frac{\bar{\vartheta}_2^2 \bar{\vartheta}_4^2}{\bar{\eta}^{24}} \\
&= \frac{\Gamma_{2,18} \begin{bmatrix} 0 \\ 0 \end{bmatrix} + \Gamma_{2,18} \begin{bmatrix} 0 \\ 1 \end{bmatrix}}{2} \bar{F}_1 - \frac{\Gamma_{2,18} \begin{bmatrix} 0 \\ 0 \end{bmatrix} - \Gamma_{2,18} \begin{bmatrix} 0 \\ 1 \end{bmatrix}}{2} \bar{F}_1 + \frac{\Gamma_{2,18} \begin{bmatrix} 1 \\ 0 \end{bmatrix} + \Gamma_{2,18} \begin{bmatrix} 1 \\ 0 \end{bmatrix}}{2} \bar{F}_+ \\
&\quad - \frac{\Gamma_{2,18} \begin{bmatrix} 1 \\ 0 \end{bmatrix} - \Gamma_{2,18} \begin{bmatrix} 1 \\ 0 \end{bmatrix}}{2} \bar{F}_- \\
\bar{F}_1 &= \frac{\bar{\vartheta}_3^2 \bar{\vartheta}_4^2}{\bar{\eta}^{24}}, \bar{F}_{\pm} = \frac{\bar{\vartheta}_2^2 (\bar{\vartheta}_3^2 \pm \bar{\vartheta}_4^2)}{\bar{\eta}^{24}}
\end{aligned}$$

$$\tau\rightarrow\tau+1:~B_2\rightarrow B_2, \tau\rightarrow-\frac{1}{\tau}:~B_2\rightarrow\tau^2B_2$$

$$\begin{aligned}
F_1 &= \frac{1}{q} + \sum_{n=0}^{\infty} d_1(n) q^n = \frac{1}{q} + 16 + 156q + \mathcal{O}(q^2) \\
F_+ &= \frac{8}{q^{3/4}} + q^{1/4} \sum_{n=0}^{\infty} d_+(n) q^n = \frac{8}{q^{3/4}} + 8q^{1/4} (30 + 481q + \mathcal{O}(q^2)) \\
F_- &= \frac{32}{q^{1/4}} + q^{3/4} \sum_{n=0}^{\infty} d_-(n) q^n = \frac{32}{q^{1/4}} + 32q^{3/4} (26 + 375q + \mathcal{O}(q^2))
\end{aligned}$$

$$M^2=\frac{\left| -m_1U+m_2+Tn_1+\left(TU-\frac{1}{2}\vec{W}^2\right)n_2+\vec{W}\cdot\vec{Q}\right|^2}{4S_2\left(T_2U_2-\frac{1}{2}\mathrm{Im}\vec{W}^2\right)}$$



$$\rho = \vec{m} \cdot \vec{\epsilon}_R + \vec{n} \cdot \epsilon_L - \vec{Q} \cdot \vec{\zeta}$$

$$s = \vec{m} \cdot \vec{n} - \frac{1}{2} \vec{Q} \cdot \vec{Q}$$

$$M^2 = \left| \left(m_1 + \frac{1}{2} \epsilon_L^1 \right) U - \left(m_2 + \frac{1}{2} \epsilon_L^2 \right) - T \left(n_1 + \frac{1}{2} \epsilon_R^1 \right) - \left(TU - \frac{1}{2} \vec{W}^2 \right) \left(n_2 + \frac{1}{2} \epsilon_R^2 \right) \right.$$

$$+ \\ - \vec{W} \cdot \left(\vec{Q} + \frac{1}{2} \vec{\zeta} \right) \Big|^2 / 4 S_2 \left(T_2 U_2 - \frac{1}{2} \text{Im} \vec{W}^2 \right) (12.4.36)$$

$$s' = \left(\vec{m} + \frac{\vec{\epsilon}_L}{2}\right) \cdot \left(\vec{n} + \frac{\vec{\epsilon}_R}{2}\right) - \frac{1}{2} \left(\vec{Q} + \frac{\vec{\zeta}}{2}\right) \cdot \left(\vec{Q} + \frac{\vec{\zeta}}{2}\right),$$

$$m_L^2 = \frac{1}{4} \left(\frac{1}{R} + nR \right)^2, m_R^2 = \frac{1}{4} \left(\frac{1}{R} - nR \right)^2$$

$$Z_{Z_2 \times Z_2}^{N=1} = \frac{1}{\tau_2 \eta^2 \bar{\eta}^2} \frac{1}{4} \sum_{h_1, g_1=0}^1 \sum_{h_2, g_2=0}^1 \frac{1}{2} \sum_{\alpha, \beta=0}^1 (-)^{\alpha+\beta+\alpha\beta} \\ \times \frac{\vartheta \left[\begin{smallmatrix} \alpha \\ \beta \end{smallmatrix} \right]}{\eta} \frac{\vartheta \left[\begin{smallmatrix} \alpha + h_1 \\ \beta + g_1 \end{smallmatrix} \right]}{\eta} \frac{\vartheta \left[\begin{smallmatrix} \alpha + h_2 \\ \beta + g_2 \end{smallmatrix} \right]}{\eta} \frac{\vartheta \left[\begin{smallmatrix} \alpha - h_1 - h_2 \\ \beta - g_1 - g_2 \end{smallmatrix} \right]}{\eta} \frac{\bar{\Gamma}_8}{\bar{\eta}^8} Z_{2,2}^1 \left[\begin{smallmatrix} h_1 \\ g_1 \end{smallmatrix} \right] \times Z_{2,2}^2 \left[\begin{smallmatrix} h_2 \\ g_2 \end{smallmatrix} \right] Z_{2,2}^3 \left[\begin{smallmatrix} h_1 + h_2 \\ g_1 + g_2 \end{smallmatrix} \right] \\ \times \frac{1}{2} \sum_{\bar{\alpha}, \bar{\beta}=0}^1 \frac{\bar{\vartheta} \left[\begin{smallmatrix} \bar{\alpha} \\ \bar{\beta} \end{smallmatrix} \right]^5}{\bar{\eta}^5} \frac{\bar{\vartheta} \left[\begin{smallmatrix} \bar{\alpha} + h_1 \\ \bar{\beta} + g_1 \end{smallmatrix} \right]}{\bar{\eta}} \frac{\bar{\vartheta} \left[\begin{smallmatrix} \bar{\alpha} + h_2 \\ \bar{\beta} + g_2 \end{smallmatrix} \right]}{\bar{\eta}} \frac{\bar{\vartheta} \left[\begin{smallmatrix} \bar{\alpha} - h_1 - h_2 \\ \bar{\beta} - g_1 - g_2 \end{smallmatrix} \right]}{\bar{\eta}}$$

$$Z_{6-d}^{II-\lambda} = \frac{1}{2} \sum_{h,g=0}^1 \frac{Z_{(4,4)} \left[\begin{smallmatrix} h \\ g \end{smallmatrix} \right]}{\tau_2^2 \eta^4 \bar{\eta}^4} \times \frac{1}{2} \sum_{a,b=0}^1 (-1)^{a+b+ab} \frac{\vartheta^2 \left[\begin{smallmatrix} a \\ b \end{smallmatrix} \right] \vartheta \left[\begin{smallmatrix} a+h \\ b+g \end{smallmatrix} \right] \vartheta \left[\begin{smallmatrix} a-h \\ b-g \end{smallmatrix} \right]}{\eta^4} \times \\ \frac{1}{2} \sum_{\bar{a},\bar{b}=0}^1 (-1)^{\bar{a}+\bar{b}+\lambda \bar{a} \bar{b}} \frac{\bar{\vartheta}^2 \left[\begin{smallmatrix} \bar{a} \\ \bar{b} \end{smallmatrix} \right] \bar{\vartheta} \left[\begin{smallmatrix} \bar{a}+h \\ \bar{b}+g \end{smallmatrix} \right] \bar{\vartheta} \left[\begin{smallmatrix} \bar{a}-h \\ \bar{b}-g \end{smallmatrix} \right]}{\bar{\eta}^4}$$

$$\chi = \frac{1}{g} [\chi(M) - \chi(F)] + \chi(N).$$

$$\left.\frac{1}{g_i^2}\right|_{\rm tree}=\frac{k_i}{g_{\rm spacetime\; dimensions}^2}=k_iS_2$$

$$V_G^{\mu,a} \sim (\partial X^\mu + i(p\cdot\psi)\psi^\mu) \bar{f}^a e^{ip\cdot X}$$



$$\int d^4x \frac{1}{g^2(T_i)} F_{\mu\nu}^a F^{a,\mu\nu}$$

$$V_{\text{modulus}}^{IJ}=(\partial X^I+i(p\cdot\psi)\psi^I)\bar{\partial}X^J e^{ip\cdot X}$$

$$I_{1-\text{loop}}=\int \left\langle V^{a,\mu}(p_1,z)V^{b,\nu}(p_2,w)V_{\text{modulus}}^{IJ}(p_3,0)\right\rangle \sim \delta^{ab}(p_1\cdot p_2\eta^{\mu\nu}-p_1^\mu p_2^\nu)F^{IJ}(T,U)+\mathcal{O}(p^4)$$

$$F^{IJ}=\int_{\mathcal{F}}\frac{d^2\tau}{\tau_2^2}\int\frac{d^2z}{\tau_2}\int d^2w\langle\psi(z)\psi(w)\rangle^2\langle\bar{J}^a(\bar{z})\bar{J}^b(\bar{w})\rangle\langle\partial X^I(0)\bar{\partial}X^J(0)\rangle$$

$$S\left[\begin{smallmatrix} a \\ b \end{smallmatrix} \right](z)=\langle \psi(z)\psi(0)\rangle|_b^a=\frac{\vartheta\left[\begin{smallmatrix} a \\ b \end{smallmatrix} \right](z)\vartheta_1'(0)}{\vartheta_1(z)\vartheta\left[\begin{smallmatrix} a \\ b \end{smallmatrix} \right](0)}=\frac{1}{z}+\cdots$$

$$S^2\left[\begin{smallmatrix} a \\ b \end{smallmatrix} \right](z)=\mathcal{P}(z)+4\pi i\partial_\tau\log\frac{\theta\left[\begin{smallmatrix} a \\ b \end{smallmatrix} \right](\tau)}{\eta(\tau)}$$

$$\langle\langle\psi(z)\psi(0)\rangle\rangle=\frac{1}{2}\sum_{(a,b)\neq(1,1)}(-1)^{a+b+ab}\frac{\vartheta^2\left[\begin{smallmatrix} a \\ b \end{smallmatrix} \right]\vartheta\left[\begin{smallmatrix} a+h \\ b+g \end{smallmatrix} \right]\vartheta\left[\begin{smallmatrix} a-h \\ b-g \end{smallmatrix} \right]}{\eta^4}S^2\left[\begin{smallmatrix} a \\ b \end{smallmatrix} \right](z)$$

$$=4\pi^2\eta^2\vartheta\left[\begin{smallmatrix} 1+h \\ 1+g \end{smallmatrix} \right]\vartheta\left[\begin{smallmatrix} 1-h \\ 1-g \end{smallmatrix} \right]$$

$$\langle \bar{J}^a(\bar{z})\bar{J}^b(0)\rangle=\frac{k\delta^{ab}}{4\pi^2}\,\bar{\partial}_{\bar{z}}^2\log\,\bar{\vartheta}_1(\bar{z})+\text{Tr}[J_0^aJ_0^b]=\delta^{ab}\left(\frac{k}{4\pi^2}\,\bar{\partial}_{\bar{z}}^2\log\,\bar{\vartheta}_1(\bar{z})+\text{Tr}[Q^2]\right)$$

$$\begin{aligned}\langle\partial X^I(0)\bar{\partial}X^J(0)\rangle&=\frac{\sqrt{\det G}}{(\sqrt{\tau_2}\eta\bar{\eta})^2}\sum_{\vec{m},\vec{n}}(m^I+n^I\tau)(m^J+n^J\bar{\tau})\times\\&\times\exp\left[-\frac{\pi(G_{KL}+B_{KL})}{\tau_2}(m_K+n_K\tau)(m_L+n_L\bar{\tau})\right]\end{aligned}$$

$$V_{T_i}=v_{IJ}(T_i)\partial X^I\bar{\partial}X^J$$

$$v(T)=-\frac{i}{2U_2}\begin{pmatrix}1&U\\ \bar{U}&|U|^2\end{pmatrix},\,v(U)=\frac{iT_2}{U_2^2}\begin{pmatrix}1&\bar{U}\\ \bar{U}&U^2\end{pmatrix}$$

$$v(\bar{T})=\overline{v(T)},v(\bar{U})=\overline{v(U)}$$

$$\langle V_{T_i}\rangle=-\frac{\tau_2}{2\pi}\partial_{T_i}\frac{\Gamma_{2,2}}{\eta^2\bar{\eta}^2}$$

$$\left.\frac{\partial}{\partial T_i}\frac{16\pi^2}{g_i^2}\right|_{1-\text{loop}}\sim\frac{\partial}{\partial T_i}\int_{\mathcal{F}}\frac{d^2\tau}{\tau_2^2}\frac{\tau_2\Gamma_{2,2}}{\bar{\eta}^4}\text{Tr}_R^{int}\left[(-1)^F\left(Q_i^2-\frac{k_i}{4\pi\tau_2}\right)\right]+\text{constant}$$

$$\langle X(z,\bar{z})X(0)\rangle=-\log|\vartheta_1(z)|^2+2\pi\frac{\text{Im}z^2}{\tau_2}$$



$$\int \frac{d^2z}{\tau_2} (S^2[\begin{smallmatrix} a \\ b \end{smallmatrix}](z) - \langle X \partial X \rangle^2) \left(\frac{k}{4\pi^2} \bar{\partial}^2 \log \bar{\vartheta}_1(\bar{z}) + \text{Tr}[Q^2] \right) \\ = 4\pi i \partial_\tau \log \frac{\vartheta[\begin{smallmatrix} a \\ b \end{smallmatrix}]}{\eta} \left(\text{Tr}[Q^2] - \frac{k}{4\pi\tau_2} \right)$$

$$Z_2^I = \frac{16\pi^2}{g_I^2} \Big|_{1-\text{loop}} = \frac{1}{4\pi^2} \int_{\mathcal{F}} \frac{d^2\tau}{\tau_2} \frac{1}{\eta^2 \bar{\eta}^2} \sum_{\text{even}} 4\pi i \partial_\tau \left(\frac{\vartheta[\begin{smallmatrix} a \\ b \end{smallmatrix}]}{\eta} \right) \text{Tr}_{\text{int}} \left[Q_I^2 - \frac{k_I}{4\pi\tau_2} \right] \begin{bmatrix} a \\ b \end{bmatrix}$$

$$\frac{16\pi^2}{g_I^2} \Big|_{1-\text{loop}}^{\text{IR}} = \int_{\mathcal{F}} \frac{d^2\tau}{\tau_2} \text{Str} Q_I^2 \left(\frac{1}{12} - s^2 \right)$$

$$\frac{16\pi^2}{g_I^2} \Big|_{1-\text{loop}}^{\text{IR}} = b_I \log (\mu^2 \alpha') + \text{finite}$$

$$b_I = \text{Str} Q_I^2 \left(\frac{1}{12} - s^2 \right) \Big|_{\text{massless}}$$

$$V_{\text{grav}} = \epsilon_{\mu\nu} (\partial X^\mu + i p \cdot \psi \psi^\mu) \bar{\partial} X^\nu$$

$$\int \frac{d^2z}{\tau_2} \langle X \bar{\partial}^2 X \rangle = \int \frac{d^2z}{\tau_2} \left(\bar{\partial}^2 \log \bar{\vartheta}_1(\bar{z}) + \frac{\pi}{\tau_2} \right) = 0.$$

$$ds^2 = G_{\mu\nu} dx^\mu dx^\nu = (dX^0)^2 + \frac{N}{4} (d\alpha^2 + d\beta^2 + d\gamma^2 + 2\sin(\beta/\sqrt{\alpha'}) d\alpha d\gamma)$$

$$B_{\mu\nu} dX^\mu \wedge dX^\nu = \frac{N}{2} \cos(\beta/\sqrt{\alpha'}) d\alpha \wedge d\gamma, \Phi = \frac{X^0 \sqrt{\alpha'}}{\sqrt{N+2}}$$

$$L_0 = -\frac{1}{2\alpha'} + E^2 + \frac{1}{4\alpha'(N+2)} + \frac{j(j+1)}{\alpha'(N+2)} + \cdots$$

$$\mu^2 = \frac{M_{\text{spacetime dimensions}}^2}{2(N+2)}, M_{\text{spacetime dimensions}} = \frac{1}{\sqrt{\alpha'}}.$$

$$Z(\mu) = \Gamma(\mu/M_{\text{spacetime dimensions}})Z(0),$$

$$\Gamma(\mu/M_{\text{spacetime dimensions}}) = 4\sqrt{x} \frac{\partial}{\partial x} [\rho(x) - \rho(x/4)] \Big|_{x=N+2} \\ \rho(x) = \sqrt{x} \sum_{m,n \in \mathbb{Z}} \exp \left[-\frac{\pi x}{\tau_2} |m + n\tau|^2 \right]$$

$$\delta I = \int d^2z (A_\mu^a(X) \partial X^\mu + F_{\mu\nu}^a \psi^\mu \psi^\nu) \bar{f}^a$$

$$\delta I = \int d^2z B^a (J^3 + i\psi^1\psi^2) \bar{f}^a$$



$$k_I \frac{16\pi^2}{g_{\text{spacetime dimensions}}^2} + Z_2^I(\mu/M_{\text{spacetime dimensions}})$$

$$\frac{16\pi^2}{g_{I\text{ bare}}^2} + b_I(4\pi)^\epsilon \int_0^\infty \frac{dt}{t^{1-\epsilon}} \Gamma_{\text{EFT}}\left(\frac{\mu}{\sqrt{\pi}M_{\text{spacetime dimensions}}}, t\right)$$

$$\frac{16\pi^2}{g_{I\text{ bare}}^2} = \frac{16\pi^2}{g_I^2(\mu)} - b_I(4\pi)^\epsilon \int_0^\infty \frac{dt}{t^{1-\epsilon}} e^{-t\mu^2/M^2}$$

$$\left. \frac{16\pi^2}{g_I^2(\mu)} \right|_{\overline{DR}} = k_I \frac{16\pi^2}{g_{\text{spacetime dimensions}}^2} + Z_2^I(\mu/M_{\text{spacetime dimensions}}) - b_I(2\gamma + 2),$$

$$\int_{\mathcal{F}} \frac{d^2\tau}{\tau_2} \Gamma(\mu/M_{\text{spacetime dimensions}}) = \log \frac{M_{\text{spacetime dimensions}}^2}{\mu^2} + \log \frac{2e^{\gamma+3}}{\pi\sqrt{27}} + \mathcal{O}\left(\frac{\mu}{M_{\text{spacetime dimensions}}}\right)$$

$$\left. \frac{16\pi^2}{g_I^2(\mu)} \right|_{\overline{DR}} = k_I \frac{16\pi^2}{g_{\text{spacetime dimensions}}^2} + b_I \log \frac{M_{\text{spacetime dimensions}}^2}{\mu^2} + b_I \log \frac{2e^{1-\gamma}}{\pi\sqrt{27}} + \Delta_I,$$

$$\Delta_I = \int_{\mathcal{F}} \frac{d^2\tau}{\tau_2} \left[\frac{1}{|\eta|^4} \sum_{\text{even}} \frac{i}{\pi} \partial_\tau \left(\frac{\vartheta \left[\begin{smallmatrix} a \\ b \end{smallmatrix} \right]}{\eta} \right) \text{Tr}_{\text{int}} \left[Q_I^2 - \frac{k_I}{4\pi\tau_2} \right] \left[\begin{smallmatrix} a \\ b \end{smallmatrix} \right] - b_I \right].$$

Umbrales gravitacionales.

$$\Delta_{\text{grav}} = \int_{\mathcal{F}} \frac{d^2\tau}{\tau_2} \left[\frac{1}{|\eta|^4} \sum_{\text{even}} \frac{i}{\pi} \partial_\tau \left(\frac{\vartheta \left[\begin{smallmatrix} a \\ b \end{smallmatrix} \right]}{\eta} \right) \frac{\hat{E}_2}{12} C^{\text{int}} \left[\begin{smallmatrix} a \\ b \end{smallmatrix} \right] - b_{\text{grav}} \right]$$

$$\zeta_{U(1)} = \epsilon_{\mu\nu}^1 \epsilon_\rho^2 \int \frac{\delta^2 z}{\tau_2} \left\langle (\partial x^\mu + i p_1 \cdot \psi \psi^\mu) \bar{\partial} X^\nu e^{i p_1 \cdot X} \Big|_z \times \right.$$

$$\left. \times \psi^\rho \bar{J} e^{i p_2 \cdot X} \Big|_0 \oint dw (\psi^\sigma \partial X^\sigma + G^{\text{int}}) \Big|_w \right\rangle$$

$$\zeta_{U(1)} = \epsilon_{\mu\nu}^1 \epsilon_\rho^2 \epsilon^{a\mu\rho\sigma} p_1^a \langle \partial X^\sigma \bar{\partial} X^\nu(\bar{z}) \rangle \langle \bar{J} \rangle + \mathcal{O}(p^2)$$

$$\zeta_{U(1)} \sim \int_{\mathcal{F}} \frac{d^2\tau}{\tau_2^2} \frac{1}{\bar{\eta}^2} \text{Tr}[(-1)^F Q]_R$$

$$\zeta_{U(1)} \sim \sum_{i, \text{ massless}} q^i$$

$$S = \int \sqrt{\det G} \left[-\frac{1}{12} e^{-2\phi} H_{\mu\nu\rho} H^{\mu\nu\rho} + \zeta B \wedge F \right]$$

$$\tilde{S} = \int \sqrt{\det G} e^{2\phi} (\partial_\mu a + \zeta A_\mu)^2$$

$$V_D \sim \zeta e^\phi \left(e^{-\phi} + \sum_i q^i h_i |c_i|^2 \right)^2$$



$$\frac{i}{2\pi} \frac{1}{2} \sum_{\text{even}} (-1)^{a+b+ab} \partial_{\tau} \left(\frac{\vartheta \left[\begin{smallmatrix} a \\ b \end{smallmatrix} \right]}{\eta} \right) \frac{\vartheta \left[\begin{smallmatrix} a \\ b \end{smallmatrix} \right] \vartheta \left[\begin{smallmatrix} a+h \\ b+g \end{smallmatrix} \right] \vartheta \left[\begin{smallmatrix} a-h \\ b-g \end{smallmatrix} \right] Z_{4,4} \left[\begin{smallmatrix} h \\ g \end{smallmatrix} \right]}{\eta^3 |\eta|^4} = 4 \frac{\eta^2}{\bar{\vartheta} \left[\begin{smallmatrix} 1+h \\ 1+g \end{smallmatrix} \right] \bar{\vartheta} \left[\begin{smallmatrix} 1-h \\ 1-g \end{smallmatrix} \right]},$$

$$\bar{\chi}_0^{\text{E}_8}(v_i) = \frac{1}{2} \sum_{a,b=0}^1 \frac{\prod_{i=1}^8 \bar{\vartheta} \left[\begin{smallmatrix} a \\ b \end{smallmatrix} \right] (v_i)}{\bar{\eta}^8}$$

$$\left[\frac{1}{(2\pi i)^2} \partial_{v_1}^2 - \frac{1}{4\pi\tau_2} \right] \bar{\chi}_0^{\text{E}_8}(v_i) \Big|_{v_i=0} = \frac{1}{12} (\hat{\bar{E}}_2 \bar{E}_4 - \bar{E}_6)$$

$$\frac{1}{2} \sum_{(h,g) \neq (0,0)} \sum_{a,b=0}^1 \frac{\bar{\vartheta} \left[\begin{smallmatrix} a \\ b \end{smallmatrix} \right] \sqrt{6} \left[\begin{smallmatrix} a+h \\ b+g \end{smallmatrix} \right] \bar{\vartheta} \left[\begin{smallmatrix} a-h \\ b-g \end{smallmatrix} \right]}{\bar{\vartheta} \left[\begin{smallmatrix} 1+h \\ 1+g \end{smallmatrix} \right] \bar{\vartheta} \left[\begin{smallmatrix} 1-h \\ 1-g \end{smallmatrix} \right]} = -\frac{1}{4} \frac{\bar{E}_6}{\bar{\eta}^6}$$

$$\Delta_{\text{E}_8} = \int_{\mathcal{F}} \frac{d^2\tau}{\tau_2} \left[-\frac{1}{12} \Gamma_{2,2} \frac{\hat{\bar{E}}_2 \bar{E}_4 \bar{E}_6 - \bar{E}_6^2}{\bar{\eta}^{24}} + 60 \right]$$

$$\left[\text{Tr} Q_{\text{E}_7}^2 - \frac{1}{4\pi\tau_2} \right] = \left[\frac{1}{(2\pi i)^2} \partial_v^2 - \frac{1}{4\pi\tau_2} \right] \frac{1}{2} \sum_{a,b} \frac{\bar{\vartheta} \left[\begin{smallmatrix} a \\ b \end{smallmatrix} \right] (v) \bar{\vartheta}^5 \left[\begin{smallmatrix} a \\ b \end{smallmatrix} \right] \bar{\vartheta} \left[\begin{smallmatrix} a+h \\ b+g \end{smallmatrix} \right] \bar{\vartheta} \left[\begin{smallmatrix} a-h \\ b-g \end{smallmatrix} \right]}{\bar{\eta}^8} \Big|_{v=0}$$

$$\Delta_{\text{E}_7} = \int_{\mathcal{F}} \frac{d^2\tau}{\tau_2} \left[-\frac{1}{12} \Gamma_{2,2} \frac{\hat{\bar{E}}_2 \bar{E}_4 \bar{E}_6 - \bar{E}_4^3}{\bar{\eta}^{24}} - 84 \right]$$

$$\Delta_{\text{E}_8} - \Delta_{\text{E}_7} = -144\Delta, \Delta = \int_{\mathcal{F}} \frac{d^2\tau}{\tau_2} (\Gamma_{2,2} - 1)$$

$$\Delta = -\log\left[4\pi^2 T_2 U_2 |\eta(T)\eta(u)|^4\right]$$

$$\lim_{T_2 \rightarrow \infty} \Delta = \frac{\pi}{3} T_2 + \mathcal{O}(\log T_2)$$

$$\Delta_{\text{E}_8}^{N=1} = \int_{\mathcal{F}} \frac{d^2\tau}{\tau_2} \left[-\frac{1}{12} \sum_{i=1}^3 \Gamma_{2,2}(T_i, U_i) \frac{\hat{\bar{E}}_2 \bar{E}_4 \bar{E}_6 - \bar{E}_6^2}{\bar{\eta}^{24}} + \frac{3}{2} 60 \right]$$

$$\Delta_{\text{E}_6}^{N=1} = \int_{\mathcal{F}} \frac{d^2\tau}{\tau_2} \left[-\frac{1}{12} \sum_{i=1}^3 \Gamma_{2,2}(T_i, U_i) \frac{\hat{\bar{E}}_2 \bar{E}_4 \bar{E}_6 - \bar{E}_4^3}{\bar{\eta}^{24}} - \frac{3}{2} 84 \right].$$

$$Z_{N=4 \rightarrow N=2} = \frac{1}{2} \sum_{h,g=0}^1 \frac{1}{\tau_2 |\eta|^4} \frac{\Gamma_{2,2} \left[\begin{smallmatrix} h \\ g \end{smallmatrix} \right] \bar{\Gamma}_{\text{E}_8}}{|\eta|^4} Z_{(4,4)} \left[\begin{smallmatrix} h \\ g \end{smallmatrix} \right] \frac{1}{2} \sum_{\gamma, \delta=0}^1 \frac{\bar{\vartheta} \left[\begin{smallmatrix} \gamma+h \\ \delta+g \end{smallmatrix} \right] \bar{\vartheta} \left[\begin{smallmatrix} \gamma-h \\ \delta-g \end{smallmatrix} \right] \bar{\vartheta}^6 \left[\begin{smallmatrix} \gamma \\ \delta \end{smallmatrix} \right]}{\bar{\eta}^8}$$

$$\times \frac{1}{2} \sum_{a,b=0}^1 (-1)^{a+b+ab} \frac{\vartheta^2 \left[\begin{smallmatrix} a \\ b \end{smallmatrix} \right] \vartheta \left[\begin{smallmatrix} a+h \\ b+g \end{smallmatrix} \right] \vartheta \left[\begin{smallmatrix} a-h \\ b-g \end{smallmatrix} \right]}{\eta^4}$$



$$m_{3/2}^2 = \frac{|U|^2}{T_2 U_2}$$

$$b_{\mathrm{E}_8} = -60 \, , b_{\mathrm{E}_7} = -12 \, , b_{\mathrm{SU}(2)} = 52$$

$$\Delta_I = b_I \Delta + \left(\frac{\tilde{b}_I}{3} - b_I \right) \delta - k_I Y$$

$$\Delta = \int_{\mathcal{F}} \frac{d^2\tau}{\tau_2} \left[\sum'_{h,g} \Gamma_{2,2} \left[\begin{smallmatrix} h \\ g \end{smallmatrix} \right] - 1 \right] = -\log \left[\frac{\pi^2}{4} |\vartheta_4(T)|^4 |\vartheta_2(U)|^4 T_2 U_2 \right]$$

$$\delta = \int_{\mathcal{F}} \frac{d^2\tau}{\tau_2} \sum'_{h,g} \Gamma_{2,2} \left[\begin{smallmatrix} h \\ g \end{smallmatrix} \right] \bar{\sigma} \left[\begin{smallmatrix} h \\ g \end{smallmatrix} \right]$$

$$Y = \int_{\mathcal{F}} \frac{d^2\tau}{\tau_2} \sum'_{h,g} \Gamma_{2,2} \left[\frac{1}{12} \frac{\hat{\tilde{E}}_2}{\bar{\eta}^{24}} \bar{\Omega} \left[\begin{smallmatrix} h \\ g \end{smallmatrix} \right] + \bar{\rho} \left[\begin{smallmatrix} h \\ g \end{smallmatrix} \right] + 40 \bar{\sigma} \left[\begin{smallmatrix} h \\ g \end{smallmatrix} \right] \right]$$

$$\Omega \left[\begin{smallmatrix} 0 \\ 1 \end{smallmatrix} \right] = \frac{1}{2} E_4 \vartheta_3^4 \vartheta_4^4 (\vartheta_3^4 + \vartheta_4^4),$$

$$\sigma \left[\begin{smallmatrix} h \\ g \end{smallmatrix} \right] = -\frac{1}{4} \frac{\vartheta^{12} \left[\begin{smallmatrix} h \\ g \end{smallmatrix} \right]}{\eta^{12}},$$

$$\rho \left[\begin{smallmatrix} 0 \\ 1 \end{smallmatrix} \right] = f(1-x) \, , \rho \left[\begin{smallmatrix} 1 \\ 0 \end{smallmatrix} \right] = f(x) \, , \rho \left[\begin{smallmatrix} 1 \\ 1 \end{smallmatrix} \right] = f(x/(x-1)),$$

$$f(x) = \frac{4(8-49x+66x^2-49x^3+8x^4)}{3x(1-x)^2}$$

$$\Delta_I = \int_{\mathcal{F}} \frac{d^2\tau}{\tau_2} \left[\frac{\Gamma_{2,2}}{\bar{\eta}^{24}} \left(\mathrm{Tr}[Q_I^2] - \frac{k_I}{4\pi\tau_2} \right) \bar{\Omega} - b_I \right]$$

$$\Delta_{\mathrm{grav}} = \int_{\mathcal{F}} \frac{d^2\tau}{\tau_2} \left[\frac{\Gamma_{2,2}}{\bar{\eta}^{24}} \frac{\hat{\tilde{E}}_2}{12} \bar{\Omega} - b_{\mathrm{grav}} \right]$$

$$\bar{\Omega} = \xi \bar{E}_4 \bar{E}_6$$

$$\Delta_I = \int_{\mathcal{F}} \frac{d^2\tau}{\tau_2} \left[\Gamma_{2,2} \left(\frac{\xi k_I}{12} \frac{\hat{\tilde{E}}_2}{\bar{\eta}^{24}} \bar{E}_4 \bar{E}_6 + A_I \bar{J} + B_I \right) - b_I \right]$$

$$\Delta_{\mathrm{grav}} = \xi \int_{\mathcal{F}} \frac{d^2\tau}{\tau_2} \left[\Gamma_{2,2} \frac{\hat{\tilde{E}}_2}{12 \bar{\eta}^{24}} \bar{E}_4 \bar{E}_6 - b_{\mathrm{grav}} \right]$$

$$A_I = -\frac{\xi k_I}{12}$$

$$744A_I + B_I - b_I + k_ib_{\mathrm{grav}} = 0$$



$$b_{\text{grav}} = \frac{22 - N_V + N_H}{12}$$

$$\Delta_I = b_I \Delta - k_I Y$$

$$\begin{aligned} \Delta &= \int_{\mathcal{F}} \frac{d^2 \tau}{\tau_2} [\Gamma_{2,2}(T, U) - 1] \\ &= -\log(4\pi^2 |\eta(T)|^4 |\eta(U)|^4 \text{Im} T \text{Im} U) \\ Y &= \frac{1}{12} \int_{\mathcal{F}} \frac{d^2 \tau}{\tau_2} \Gamma_{2,2}(T, U) \left[\frac{\hat{\bar{E}}_2 \bar{E}_4 \bar{E}_6}{\bar{\eta}^{24}} - \bar{J} + 1008 \right] \end{aligned}$$

$$\Delta_{\text{grav}} = - \int_{\mathcal{F}} \frac{d^2 \tau}{\tau_2} \left[\Gamma_{2,2} \frac{\hat{\bar{E}}_2 \bar{E}_4 \bar{E}_6}{12 \bar{\eta}^{24}} - 22 \right]$$

$$\frac{16\pi^2}{g_I^2(\mu)} = k_I \frac{16\pi^2}{g_{\text{renorm}}^2} + b_I \log \frac{M_s^2}{\mu^2} + \hat{\Delta}_I$$

$$g_{\text{renorm}}^2 = \frac{g_{\text{string}}^2}{1 - \frac{Y}{16\pi^2} g_{\text{string}}^2}$$

$$M_{\text{string}} = \frac{M_P g_{\text{renorm}}}{\sqrt{1 + \frac{Y}{16\pi^2} g_{\text{renorm}}^2}}$$

$$\frac{1}{g_I^2} = \frac{k_I}{g_U^2}$$

$$\frac{16\pi^2}{g_I^2(\mu)} = k_I \frac{16\pi^2}{g_U^2} + b_I \log \frac{M_U^2}{\mu^2}$$

$$g_U = g_{\text{renorm}} = \frac{g_{\text{spacetime dimensions}}}{\sqrt{1 - \frac{g_{\text{spacetime dimensions}}^2 Y}{16\pi^2}}}$$

$$M_U^2 = \frac{2e^{1-\gamma}}{\pi\sqrt{27}} e^{\Delta} M_P^2 g_{\text{spacetime dimensions}}^2 = \frac{2e^{1-\gamma}}{\pi\sqrt{27}} e^{\Delta} M_P^2 \frac{g_U}{\sqrt{1 + \frac{g_U^2 Y}{16\pi^2}}}$$

$$Z_{D=9}^{0(32)} = \frac{1}{(\sqrt{\tau_2} \eta \bar{\eta})^7} \frac{\Gamma_{1,17}(R, Y^I)}{\eta \bar{\eta}^{17}} \frac{1}{2} \sum_{a,b=0}^1 (-1)^{a+b+ab} \frac{\vartheta^4 \left[\begin{smallmatrix} a \\ b \end{smallmatrix} \right]}{\eta^4},$$



$$\begin{aligned}
\Gamma_{1,17}(R) &= R \sum_{m,n \in \mathbb{Z}} \exp \left[-\frac{\pi R^2}{\tau_2} |m + \tau n|^2 \right] \frac{1}{2} \sum_{a,b} \bar{\vartheta}^8 \left[\begin{smallmatrix} a \\ b \end{smallmatrix} \right] \bar{\vartheta}^8 \left[\begin{smallmatrix} a+n \\ b+m \end{smallmatrix} \right] \\
&= \frac{1}{2} \sum_{h,g=0}^1 \Gamma_{1,1}(2R) \left[\begin{smallmatrix} h \\ g \end{smallmatrix} \right] \frac{1}{2} \sum_{a,b} \bar{\vartheta}^8 \left[\begin{smallmatrix} a \\ b \end{smallmatrix} \right] \bar{\vartheta}^8 \left[\begin{smallmatrix} a+h \\ b+g \end{smallmatrix} \right] \\
\Gamma_{1,1}(R) \left[\begin{smallmatrix} h \\ g \end{smallmatrix} \right] &= R \sum_{m,n \in \mathbb{Z}} \exp \left[-\frac{\pi R^2}{\tau_2} \left| \left(m + \frac{g}{2} \right) + \tau \left(n + \frac{h}{2} \right) \right|^2 \right] \\
&= \frac{1}{R} \sum_{m,n \in \mathbb{Z}} (-1)^{mh+ng} \exp \left[-\frac{\pi}{\tau_2 R^2} |m + \tau n|^2 \right].
\end{aligned}$$

$$\partial X^9 \rightarrow \partial X^9, \psi^9 \rightarrow \psi^9, \bar{\partial} X^9 \rightarrow -\bar{\partial} X^9, \bar{\psi}^9 \rightarrow -\bar{\psi}^9.$$

Tensores antisimétricos y supermembranas.

$$A_p \equiv A_{\mu_1 \mu_2 \dots \mu_p} dx^{\mu_1} \wedge \dots \wedge dx^{\mu_p}$$

$$A_p \rightarrow A_p + d\Lambda_{p-1},$$

$$F_{p+1} = dA_p$$

$$d^*F_{p+1} = 0$$

$$\exp \left[iQ_p \int_{\text{world-volume}} A_{p+1} \right] = \exp \left[iQ_p \int A_{\mu_0 \dots \mu_p} dx^{\mu_0} \wedge \dots \wedge dx^{\mu_p} \right]$$

$$d\tilde{A}_{D-p-3} = \tilde{F}_{D-p-2} = {}^*F_{p+2} = {}^*dA_{p+1}$$

$$\Phi = \int_{S^{D-p-2}} {}^*F_{p+2} = \int_{S^{D-p-3}} \tilde{A}_{D-p-3}$$

$$\Phi \tilde{Q}_{D-p-4} = Q_p \tilde{Q}_{D-p-4} = 2\pi N, n \in \mathbb{Z}$$

$$M_{m,n}^2 = \frac{|m+n\tau|^2}{\tau_2}$$

$$M_{m_0,n_0} = \sum_{i=1}^N \sqrt{M_i^2 + \vec{p}_i^2} \geq \sum_{i=1}^N M_i$$

$$M_{m_0,n_0} \geq \sum_{i=1}^N M_{m_i,n_i}$$

$$\|v_0\| \leq \sum_{i=1}^N \|v_i\|$$



N - Dimensiones y p - supermembranas.

$$S^{\text{het}} = \int d^{10}x \sqrt{G} e^{-\Phi} \left[R + (\nabla\Phi)^2 - \frac{1}{12} \hat{H}^2 - \frac{1}{4} F^2 \right]$$

$$S^I = \int d^{10}x \sqrt{G} \left[e^{-\Phi} (R + (\nabla\Phi)^2) - \frac{1}{4} e^{-\Phi/2} F^2 - \frac{1}{12} \hat{H}^2 \right]$$

$$S_E^{\text{het}} = \int d^{10}x \sqrt{g} \left[R - \frac{1}{8} (\nabla\Phi)^2 - \frac{1}{4} e^{-\Phi/4} F^2 - \frac{1}{12} e^{-\Phi/2} \hat{H}^2 \right]$$

$$S_E^I = \int d^{10}x \sqrt{g} \left[R - \frac{1}{8} (\nabla\Phi)^2 - \frac{1}{4} e^{\Phi/4} F^2 - \frac{1}{12} e^{\Phi/2} \hat{H}^2 \right]$$

$$\text{NN NS sector } \psi + \bar{\psi}|_{\sigma=0} = \psi - \bar{\psi}|_{\sigma=\pi} = 0$$

$$\text{NN R sector } \psi - \bar{\psi}|_{\sigma=0} = \psi - \bar{\psi}|_{\sigma=\pi} = 0.$$

$$\text{DDNS sector } \psi - \bar{\psi}|_{\sigma=0} = \psi + \bar{\psi}|_{\sigma=\pi} = 0$$

$$\text{DDR sector } \psi + \bar{\psi}|_{\sigma=0} = \psi + \bar{\psi}|_{\sigma=\pi} = 0$$

$$X^I(\sigma, \tau) = x^I + w^I \sigma + 2 \sum_{n \neq 0} \frac{a_n^I}{n} e^{in\tau} \sin(n\sigma),$$

$$X^\mu(\sigma, \tau) = x^\mu + p^\mu \tau - 2i \sum_{n \neq 0} \frac{a_n^\mu}{n} e^{in\tau} \cos(n\sigma).$$

$$\psi^I(\sigma, \tau) = \sum_{n \in \mathbb{Z}} b_{n+1/2}^I e^{i(n+1/2)(\sigma+\tau)}, \psi^\mu(\sigma, \tau) = \sum_{n \in \mathbb{Z}} b_{n+1/2}^\mu e^{i(n+1/2)(\sigma+\tau)},$$

$$\psi^I(\sigma, \tau) = \sum_{n \in \mathbb{Z}} b_n^I e^{in(\sigma+\tau)}, \psi^\mu(\sigma, \tau) = \sum_{n \in \mathbb{Z}} b_n^\mu e^{in(\sigma+\tau)}.$$

$$\bar{b}_{n+1/2}^I = b_{n+1/2}^I, \bar{b}_n^I = -b_n^I$$

$$\bar{b}_{n+1/2}^\mu = -b_{n+1/2}^\mu, \bar{b}_n^\mu = b_n^\mu.$$

$$\Omega b_{-1/2}^\mu |0\rangle = \bar{b}_{-1/2}^\mu |0\rangle = -b_{-1/2}^\mu |0\rangle,$$

$$\Omega b_{-1/2}^I |0\rangle = \bar{b}_{-1/2}^I |0\rangle = b_{-1/2}^I |0\rangle.$$

$$\Gamma^{11}|R\rangle = |R\rangle$$

$$\Omega|R\rangle = -\Gamma^2 \dots \Gamma^9 |R\rangle = |R\rangle$$

$$\Gamma^0 \Gamma^1 |R\rangle = -|R\rangle$$

$$\partial_\tau X^I|_{\sigma=0} = 0, \partial_\sigma X^I|_{\sigma=\pi} = 0,$$

$$\text{DNNS sector } \psi + \bar{\psi}|_{\sigma=0} = \psi + \bar{\psi}|_{\sigma=\pi} = 0,$$

$$\text{DNR sector } \psi - \bar{\psi}|_{\sigma=0} = \psi + \bar{\psi}|_{\sigma=\pi} = 0,$$



$$b_{-1/2}^{\mu}|p;i,j\rangle,b_{-1/2}^l|p;i,j\rangle$$

$$\tilde{S}^{IIA}=\frac{1}{2\kappa_{10}^2}\Big[\int\,d^{10}x\sqrt{g}e^{-\Phi}\Big[\Big(R+(\nabla\Phi)^2-\frac{1}{12}H^2\Big)-\frac{1}{2\cdot4!}\hat{G}^2-\frac{1}{4}F^2\Big]+\frac{1}{(48)^2}\int\,B\wedge G\wedge G\Big]$$

$$\{Q^1_{\alpha},Q^2_{\alpha}\}=\delta_{\alpha\dot{\alpha}}W$$

$$M\geq \frac{c_0}{\lambda}|W|$$

$$M=\frac{c}{\lambda}|n| \, , n\in Z.$$

$$R=\lambda^{2/3}$$

$$e^{\phi/2}=\lambda+\frac{c}{r^8}\,,\chi=\chi_0+i\frac{c}{\lambda(\lambda r^8+c)}\,,$$

$$ds^2=A(r)^{-3/4}[-(dx^0)^2+(dx^1)^2]+A(r)^{1/4}dy\cdot dy$$

$$S=\chi_0+i\frac{e^{-\phi_0/2}}{\sqrt{A(r)}}$$

$$B^1=0\,,B_{01}^2=\frac{1}{\sqrt{\Delta}A(r)}$$

$$A(r)=1+\frac{Q\sqrt{\Delta}}{3r^6}\,,Q=\frac{3\kappa^2T}{\pi^4}\,,\Delta=e^{\phi_0/2}[\chi_0^2+e^{-\phi_0}].$$

$$\tilde{T}=T\sqrt{\Delta}.$$

$$T_{p,q}=T\frac{|p+qS|}{\sqrt{S_2}}$$

$$M_B^2=\frac{m^2}{R_B^2}+\left(2\pi R_B nT_{p,q}\right)^2+4\pi T_{p,q}(N_L+N_R)$$

$$M_B^2|_{\rm BPS}=\left(\frac{m}{R_B}+2\pi R_B nT_{p,q}\right)^2$$

$$M_B^2|_{BPS}=\left(\frac{m}{R_B}+2\pi R_B T\frac{|n_1+n_2S|}{\sqrt{S_2}}\right)^2$$

$$M_{11}^2=(m(2\pi R_{11})^2\tau_2T_{11})^2+\frac{|n_1+n_2\tau|^2}{R_{11}^2\tau_2^2}+\cdots$$

$$S=\tau,\frac{1}{R_B^2}=TT_{11}A_{11}^{3/2},\beta=2\pi R_{11}\frac{\sqrt{\tau_2}T_{11}}{T}$$

$$\begin{aligned}\mathcal{A} &= 2V_{p+1} \int \frac{d^{p+1}k}{(2\pi)^{p+1}} \int_0^\infty \frac{dt}{2t} e^{-2\pi\alpha' t k^2 - t \frac{|Y|^2}{2\pi\alpha'}} \frac{1}{\eta^{12}(it)} \frac{1}{2} \sum_{a,b} (-1)^{a+b+ab} g^4 \begin{bmatrix} a \\ b \end{bmatrix} (it) \\ &= 2V_{p+1} \int_0^\infty \frac{dt}{2t} (8\pi^2 \alpha' t)^{-\frac{p+1}{2}} e^{-t \frac{|Y|^2}{2\pi\alpha'}} \frac{1}{\eta^{12}(it)} \frac{1}{2} \sum_{a,b=0}^1 (-1)^{a+b+ab} g^4 \begin{bmatrix} a \\ b \end{bmatrix} (it)\end{aligned}$$

$$\begin{aligned}\mathcal{A}|_{\substack{\text{spacetime dimensions} \\ \text{supermassive.}}} &= 8(1-1)V_{p+1} \int_0^\infty \frac{dt}{t} (8\pi^2 \alpha' t)^{-\frac{p+1}{2}} t^4 e^{-t \frac{|Y|^2}{2\pi\alpha'}} \\ &= 2\pi(1-1)V_{p+1} (4\pi^2 \alpha')^{3-p} G_{9-p}(|Y|)\end{aligned}$$

$$G_d(|Y|) = \frac{1}{4\pi^{d/2}} \int_0^\infty \frac{dt}{t^{(4-d)/2}} e^{-t|Y|^2}$$

$$S = \frac{\alpha_p}{2} \int F_{p+2}^* F_{p+2} + iT_p \int_{\text{branes}} A_{p+1}$$

$$\mathcal{A}|_{\text{field theory}} = \frac{(iT_p)^2}{\alpha_p} V_{p+1} G_{9-p}(|Y|)$$

$$\frac{T_p^2}{\alpha_p} = 2\pi(4\pi^2 \alpha')^{3-p}$$

$$\frac{T_p T_{6-p}}{\alpha_p} = 2\pi n$$

$$2\kappa_{10}^2 = (2\pi)^7 \alpha'^4$$

$$S_p = -T_p \int_{W_{p+1}} d^{p+1} \xi e^{-\Phi/2} \sqrt{\det \hat{G}} - iT_p \int A_{p+1}$$

$$\hat{G}_{\alpha\beta} = G_{\mu\nu} \frac{\partial X^\mu}{\partial \xi^\alpha} \frac{\partial X^\nu}{\partial \xi^\beta}$$

$$\int A_{p+1} = \frac{1}{(p+1)!} \int d^{p+1} \xi A_{\mu_1 \dots \mu_{p+1}} \frac{\partial X^{\mu_1}}{\partial \xi^{\alpha_1}} \dots \frac{\partial X^{\mu_{p+1}}}{\partial \xi^{\alpha_{p+1}}} \epsilon^{\alpha_1 \dots \alpha_{p+1}}$$

$$2\kappa_{10}^2 T_p T_{6-p} = 2\pi n$$

$$T_p = \frac{1}{(2\pi)^p (\alpha')^{(p+1)/2}}$$

$$2\kappa_{11}^2 = (2\pi)^8 (\alpha')^{9/2}$$

$$T_2^M = T_2 = \frac{1}{(2\pi)^2 (\alpha')^{3/2}}$$

$$2\kappa_{11}^2 T_2^M T_5^M = 2\pi \rightarrow T_5^M = \frac{1}{(2\pi)^5 (\alpha')^3}$$

$$2\pi\sqrt{\alpha'} T_5^M = T_4$$



$$S_{BI} = \int d^{10}x e^{-\Phi/2} \sqrt{\det(\delta_{\mu\nu} + 2\pi\alpha' F_{\mu\nu})}$$

$$S_B = \frac{i}{2\pi\alpha'} \int_{M_2} d^2\xi \epsilon^{\alpha\beta} B_{\mu\nu} \partial_\alpha x^\mu \partial_\beta x^\nu - \frac{i}{2} \int_{B_1} ds A_\mu \partial_s x^\mu$$

$$\delta S_B = \frac{i}{\pi\alpha'} \int_{B_1} ds \Lambda_\mu \partial_s x^\mu$$

$$\mathcal{F}_{\mu\nu} = 2\pi\alpha' F_{\mu\nu} - B_{\mu\nu} = 2\pi\alpha' (\partial_\mu A_\nu - \partial_\nu A_\mu) - B_{\mu\nu}$$

$$S_p = -T_p \int_{W_{p+1}} d^{p+1}\xi e^{-\Phi/2} \sqrt{\det(\hat{G} + \mathcal{F})} - iT_p \int A_{p+1}$$

$$S_p = -T_p \int_{W_{p+1}} d^{p+1}\xi e^{-\frac{\Phi}{2}} \sqrt{\det(\hat{G} + \mathcal{F})} \pm iT_p \int A$$

$$S_1 = -\frac{1}{2\pi\alpha'} \left[\int d^2\xi \frac{|S|}{\sqrt{S_2}} \sqrt{\det(\hat{G} + \mathcal{F})} + i \int \left(B^N + \frac{iS_1}{2\pi} \mathcal{F} \right) \right]$$

$$e^{-\Phi/2} \rightarrow \frac{\sqrt{\alpha'}}{R} e^{-\Phi/2}$$

$$T_{p+1} = \frac{T_p}{2\pi\sqrt{\alpha'}}$$

$$S_p^N = -T_p \text{Str} \int_{W_{p+1}} d^{p+1}\xi e^{-\Phi/2} (F_{\mu\nu}^2 + 2F_{\mu l}^2 + F_{IJ}^2)$$

$$F_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu + [A_\mu, A_\nu]$$

$$F_{\mu l} = \partial_\mu X^l + [A_\mu, X^l], F_{IJ} = [X^I, X^J]$$

$$S_{D=6}^{\text{het}} = \int d^6x \sqrt{-G} \left[R - \frac{1}{4} \partial^\mu \Phi \partial_\mu \Phi - \frac{e^{-\Phi}}{12} \hat{H}^{\mu\nu\rho} \hat{H}_{\mu\nu\rho} - \frac{e^{-\frac{\Phi}{2}}}{4} (\hat{M}^{-1})_{ij} F_{\mu\nu}^i F^{j\mu\nu} + \frac{1}{8} \text{Tr}(\partial_\mu \hat{M} \partial^\mu \hat{M}^{-1}) \right]$$

$$S_{D=6}^{IIA} = \int d^6x \sqrt{-G} \left[R - \frac{1}{4} \partial^\mu \Phi \partial_\mu \Phi \right.$$

$$\left. - \frac{1}{12} e^{-\Phi} H^{\mu\nu\rho} H_{\mu\nu\rho} - \frac{1}{4} e^{\Phi/2} (\hat{M}^{-1})_{ij} F_{\mu\nu}^i F^{j\mu\nu} + \frac{1}{8} \text{Tr}(\partial_\mu \hat{M} \partial^\mu \hat{M}^{-1}) \right]$$

$$+ \frac{1}{16} \int d^6x \epsilon^{\mu\nu\rho\sigma\tau\epsilon} B_{\mu\nu} F_{\rho\sigma}^i \hat{L}_{ij} F_{\tau\epsilon}^j$$

$$\Phi' = -\Phi, G'_{\mu\nu} = G_{\mu\nu}, \hat{M}' = \hat{M}, A'^i_\mu = A^i_\mu$$

$$e^{-\Phi} \hat{H}_{\mu\nu\rho} = \frac{1}{6} \frac{\epsilon_{\mu\nu\rho\sigma\tau\epsilon}}{\sqrt{-G}} H'_{\sigma\tau\epsilon}$$



$$\phi = \Phi - \frac{1}{2} \log [\det G_{\alpha\beta}]$$

$$S_{D=4}^{\text{het}} = \int d^4x \sqrt{-g} e^{-\phi} [R + L_B + L_{\text{gauge}} + L_{\text{scalar}}]$$

$$\mathcal{L}_{g+\phi} = R + \partial^\mu \phi \partial_\mu \phi,$$

$$\mathcal{L}_B = -\frac{1}{12} H^{\mu\nu\rho} H_{\mu\nu\rho},$$

$$\begin{aligned} H_{\mu\nu\rho} &= \partial_\mu B_{\nu\rho} - \frac{1}{2} [B_{\mu\alpha} F_{\nu\rho}^{A,\alpha} + A_\mu^\alpha F_{\alpha,\nu\rho}^B + \hat{L}_{ij} A_\mu^i F_{\nu\rho}^j] + \text{cyclic} \\ &\equiv \partial_\mu B_{\nu\rho} - \frac{1}{2} L_{IJ} A_\mu^I F_{\nu\rho}^J + \text{cyclic} \end{aligned}$$

$$L = \begin{pmatrix} 0 & 0 & 1 & 0 & \vec{0} \\ 0 & 0 & 0 & 1 & \vec{0} \\ 1 & 0 & 0 & 0 & \vec{0} \\ 0 & 1 & 0 & 0 & \vec{0} \\ \vec{0} & \vec{0} & \vec{0} & \vec{0} & \hat{L} \end{pmatrix}$$

$$C_{\alpha\beta} = \epsilon_{\alpha\beta} B - \frac{1}{2} \hat{L}_{ij} Y_\alpha^i Y_\beta^j$$

$$\begin{aligned} \mathcal{L}_{\text{gauge}} &= -\frac{1}{4} \left\{ \left[(\hat{M}^{-1})_{ij} + \hat{L}_{ki} \hat{L}_{lj} Y_\alpha^k G^{\alpha\beta} Y_\beta^l \right] F_{\mu\nu}^i F^{j,\mu\nu} + G^{\alpha\beta} F_{\alpha,\mu\nu}^B F_{B,\beta}^{\mu\nu} \right. \\ &\quad + \left[G_{\alpha\beta} + C_{\gamma\alpha} G^{\gamma\delta} C_{\delta\beta} + Y_\alpha^i (\hat{M}^{-1})_{ij} Y_\beta^j \right] F_{\mu\nu}^{A,\alpha} F_A^{\beta,\mu\nu} - 2 G^{\alpha\gamma} C_{\gamma\beta} F_{\alpha,\mu\nu}^B F^{A,\beta,\mu\nu} \\ &\quad \left. - 2 \hat{L}_{ij} Y_\alpha^i G^{\alpha\beta} F_{\mu\nu}^j F_\beta^{B,\mu\nu} + 2 \left(Y_\alpha^i (\hat{M}^{-1})_{ij} + C_{\gamma\alpha} G^{\gamma\beta} \hat{L}_{ij} Y_\beta^i \right) F_{\mu\nu}^{a,A} F^{j,\mu\nu} \right\} \\ &\equiv -\frac{1}{4} (M^{-1})_{IJ} F_{\mu\nu}^I F^{J,\mu\nu} \end{aligned}$$

$$\begin{aligned} \mathcal{L}_{\text{scalar}} &= \partial_\mu \phi \partial^\mu \phi + \frac{1}{8} \text{Tr} [\partial_\mu \hat{M} \partial^\mu \hat{M}^{-1}] + \frac{1}{2} G^{\alpha\beta} (\hat{M}^{-1})_{ij} \partial_\mu Y_\alpha^i \partial^\mu Y_\beta^j + \frac{1}{4} \partial_\mu G_{\alpha\beta} \partial^\mu G^{\alpha\beta} \\ &\quad + \frac{1}{2 \det G} \left[\partial_\mu B + \epsilon^{\alpha\beta} \hat{L}_{ij} Y_\alpha^i \partial_\mu Y_\beta^j \right] \left[\partial^\mu B + \epsilon^{\alpha\beta} \hat{L}_{ij} Y_\alpha^i \partial^\mu Y_\beta^j \right] \\ &= \partial_\mu \phi \partial^\mu \phi + \frac{1}{8} \text{Tr} [\partial_\mu M \partial^\mu M^{-1}] \end{aligned}$$

$$g_{\mu\nu} \rightarrow e^{-\phi} g_{\mu\nu}$$

$$\begin{aligned} S_{D=4}^{\text{het,E}} &= \int d^4x \sqrt{-g} \left[R - \frac{1}{2} \partial^\mu \phi \partial_\mu \phi \right. \\ &\quad \left. - \frac{1}{12} e^{-2\phi} H^{\mu\nu\rho} H_{\mu\nu\rho} - \frac{1}{4} e^{-\phi} (M^{-1})_{IJ} F_{\mu\nu}^I F^{J,\mu\nu} + \frac{1}{8} \text{Tr} (\partial_\mu M \partial^\mu M^{-1}) \right] \end{aligned}$$



$$e^{-2\phi} H_{\mu\nu\rho} = \frac{\epsilon_{\mu\nu\rho}^\sigma}{\sqrt{-g}} \partial_\sigma a$$

$$\begin{aligned} \tilde{S}_{D=4}^{\text{het}} = \int d^4x \sqrt{-g} \left[R - \frac{1}{2} \partial^\mu \phi \partial_\mu \phi - \frac{1}{2} e^{2\phi} \partial^\mu a \partial_\mu a - \frac{1}{4} e^{-\phi} (M^{-1})_{IJ} F_{\mu\nu}^I F^{J,\mu\nu} \right. \\ \left. + \frac{1}{4} a L_{IJ} F_{\mu\nu}^I \tilde{F}^{J,\mu\nu} + \frac{1}{8} \text{Tr}(\partial_\mu M \partial^\mu M^{-1}) \right] \end{aligned}$$

$$\tilde{F}^{\mu\nu} = \frac{1}{2} \frac{\epsilon^{\mu\nu\rho\sigma}}{\sqrt{-g}} F_{\rho\sigma}$$

$$S = a + i e^{-\phi}$$

$$\tilde{S}_{D=4}^{\text{het}} = \int d^4x \sqrt{-g} \left[R - \frac{1}{2} \frac{\partial^\mu S \partial_\mu \bar{S}}{\text{Im} S^2} - \frac{1}{4} \text{Im} S (M^{-1})_{IJ} F_{\mu\nu}^I F^{J,\mu\nu} + \frac{1}{4} \text{Re} S L_{IJ} F_{\mu\nu}^I \tilde{F}^{J,\mu\nu} + \frac{1}{8} \text{Tr}(\partial_\mu M \partial^\mu M^{-1}) \right]$$

$$e^{-2\phi} H_{\mu\nu\rho} = \frac{\epsilon_{\mu\nu\rho}^\sigma}{\sqrt{-g}} \left[\partial_\sigma a + \frac{1}{2} \hat{L}_{ij} Y_\alpha^i \delta_\sigma Y_\beta^j \epsilon^{\alpha\beta} \right]$$

$$\tilde{S}_{D=4}^{IIA} = \int d^4x \sqrt{-g} [R + \mathcal{L}_{\text{gauge}}^{\text{even}} + \mathcal{L}_{\text{gauge}}^{\text{odd}} + \mathcal{L}_{\text{scalar}}]$$

$$\begin{aligned} \mathcal{L}_{\text{gauge}}^{\text{even}} = -\frac{1}{4} \int d^4x \sqrt{-g} \left[e^{-\phi} G^{\alpha\beta} (F_{\alpha,\mu\nu}^B - B_{\alpha\gamma} F_{\mu\nu}^{A,\gamma}) (F_\beta^{B,\mu\nu} - B_{\alpha\delta} F_A^{\delta,\mu\nu}) \right. \\ \left. + e^{-\phi} G_{\alpha\beta} F_{\mu\nu}^{A,\alpha} F_A^{\beta,\mu\nu} \sqrt{\det G_{\alpha\beta}} (\hat{M}^{-1})_{ij} (F_{\mu\nu}^i + Y_\alpha^i F_{\mu\nu}^{A,\alpha}) (F^{j,\mu\nu} + Y_\beta^j F_A^{\beta,\mu\nu}) \right] \end{aligned}$$

$$\begin{aligned} \mathcal{L}_{\text{gauge}}^{\text{odd}} = \frac{1}{2} \int d^4x \epsilon^{\mu\nu\rho\sigma} \left[\frac{1}{4} a F_{\alpha,\mu\nu}^B F_{\rho\sigma}^{A,\alpha} \right. \\ \left. + \frac{1}{2} \epsilon^{\alpha\beta} \hat{L}_{ij} Y_\beta^i F_{\alpha,\mu\nu}^B \left(F_{\rho\sigma}^j + \frac{1}{2} Y_\gamma^j F_{\rho\sigma}^{A,\gamma} \right) - \frac{1}{8} \epsilon^{\alpha\beta} \hat{L}_{ij} B_{\alpha\beta} (F_{\mu\nu}^i + Y_\gamma^i F_{\mu\nu}^{A,\gamma}) (F_{\rho\sigma}^j + Y_\delta^j F_{\rho\sigma}^{A,\delta}) \right] \end{aligned}$$

$$\begin{aligned} \mathcal{L}_{\text{scalar}} = -\frac{1}{2} (\partial\phi)^2 + \frac{1}{4} \partial^\mu G_{\alpha\beta} \partial_\mu G^{\alpha\beta} - \frac{1}{2 \det G} \partial_\mu B \partial^\mu B + \frac{1}{8} \text{Tr}[\partial_\mu \hat{M} \partial^\mu \hat{M}^{-1}] \\ - \frac{1}{2} e^{2\phi} \left(\partial_\mu a + \frac{1}{2} \hat{L}_{ij} \epsilon^{\alpha\beta} Y_\alpha^i \partial^\mu Y_\beta^j \right)^2 - \frac{1}{2} e^\phi \sqrt{\det G_{\alpha\beta}} \hat{M}^{-1} \Big|_{ij} G^{\alpha\beta} \partial_\mu Y_\alpha^i \partial^\mu Y_\beta^j \end{aligned}$$

$$e^{-\phi} = \sqrt{\det G'_{\alpha\beta}}, e^{-\phi'} = \sqrt{\det G_{\alpha\beta}}$$

$$\frac{G_{\alpha\beta}}{\sqrt{\det G_{\alpha\beta}}} = \frac{G'_{\alpha\beta}}{\sqrt{\det G'_{\alpha\beta}}}, A_\mu'^\alpha = A_\mu^\alpha$$

$$g_{\mu\nu} = g'_{\mu\nu} \text{ Einstein frame}$$

$$\hat{M}' = \hat{M}, A_\mu^i = A_\mu'^i, Y_\alpha^i = Y_\alpha'^i$$



$$A=B', A'=B$$

$$\frac{1}{2}\frac{\epsilon_{\mu\nu}^{\rho\sigma}}{\sqrt{-g}}\epsilon^{\alpha\beta}F_{\beta,\rho\sigma}^{B'}=e^{-\phi}G^{\alpha\beta}[F_{\beta,\mu\nu}^B-C_{\beta\gamma}F_{\mu\nu}^{A,\gamma}-\hat{L}_{ij}Y_{\beta}^iF_{\mu\nu}^j]-\frac{1}{2}a\frac{\epsilon_{\mu\nu}^{\rho\sigma}}{\sqrt{-g}}F_{\rho\sigma}^{A,\alpha}$$

$$W^i=W_1^i+iW_2^i=-Y_2^i+UY_1^i\\ G_{\alpha\beta}=\frac{T_2-\frac{\sum_i\left(W_2^i\right)^2}{2U_2}}{U_2}\begin{pmatrix}1&U_1\\U_1&|U|^2\end{pmatrix},B=T_1-\frac{\sum_iW_1^iW_2^i}{2U_2}$$

$$\mathcal{L}_{\text{scalar}}^{\text{het}}=-\frac{1}{2}\partial_{z^i}\partial_{\bar{z}^j}K(z_k,\bar{z}_k)\partial_{\mu}z^i\partial^{\mu}\bar{z}^j$$

$$K=\log\left[S_2\left(T_2U_2-\frac{1}{2}\sum_i\left(W_2^i\right)^2\right)\right]$$

$$G_{\alpha\beta}=\frac{T_2}{U_2}\begin{pmatrix}1&U_1\\U_1&|U|^2\end{pmatrix},B=T_1$$

$$S=a-\frac{\sum_i W_1^iW_2^i}{2U_2}+i\left(e^{-\phi}-\frac{\sum_i\left(W_2^i\right)^2}{2U_2}\right)$$

$$S'=T,T'=S,U=U',W^i=W'^i$$

$$\begin{pmatrix} m_1 \\ m_2 \\ n_1 \\ n_2 \\ q^i \end{pmatrix} \rightarrow \begin{pmatrix} m_1 \\ m_2 \\ \tilde{n}_2 \\ -\tilde{n}_1 \\ q^i \end{pmatrix}, \begin{pmatrix} \tilde{m}_1 \\ \tilde{m}_2 \\ \tilde{n}_1 \\ \tilde{n}_2 \\ \tilde{q}^i \end{pmatrix} \rightarrow \begin{pmatrix} \tilde{m}_1 \\ \tilde{m}_2 \\ -n_2 \\ n_1 \\ \tilde{q}^i \end{pmatrix}$$

$$\vartheta\left[\begin{smallmatrix}a\\b\end{smallmatrix}\right](v\mid\tau)=\sum_{n\in Z}q^{\frac{1}{2}\left(n-\frac{a}{2}\right)^2}e^{2\pi i\left(v-\frac{b}{2}\right)\left(n-\frac{a}{2}\right)}$$

$$\vartheta\left[\begin{smallmatrix}a+2\\b\end{smallmatrix}\right](v\mid\tau)=\vartheta\left[\begin{smallmatrix}a\\b\end{smallmatrix}\right](v\mid\tau),\vartheta\left[\begin{smallmatrix}a\\b+2\end{smallmatrix}\right](v\mid\tau)=e^{i\pi a}\vartheta\left[\begin{smallmatrix}a\\b\end{smallmatrix}\right](v\mid\tau),\\ \vartheta\left[\begin{smallmatrix}-a\\-b\end{smallmatrix}\right](v\mid\tau)=\vartheta\left[\begin{smallmatrix}a\\b\end{smallmatrix}\right](-v\mid\tau),\vartheta\left[\begin{smallmatrix}a\\b\end{smallmatrix}\right](-v\mid\tau)=e^{i\pi ab}\vartheta\left[\begin{smallmatrix}a\\b\end{smallmatrix}\right](v\mid\tau)\;(a,b\in Z).$$

$$\vartheta\left[\begin{smallmatrix}a\\b\end{smallmatrix}\right](v\mid\tau+1)=e^{-\frac{i\pi}{4}a(a-2)}\vartheta\left[\begin{smallmatrix}a\\a+b-1\end{smallmatrix}\right](v\mid\tau),\\ \vartheta\left[\begin{smallmatrix}a\\b\end{smallmatrix}\right]\left(\frac{v}{\tau}\middle|\!-\frac{1}{\tau}\right)=\sqrt{-i\tau}e^{\frac{i\pi}{2}ab+i\pi\frac{v^2}{\tau}}\vartheta\left[\begin{smallmatrix}b\\-a\end{smallmatrix}\right](v\mid\tau).$$



$$\vartheta_1(v|\tau) = 2q^{\frac{1}{8}} \sin[\pi v] \prod_{n=1}^{\infty} (1-q^n)(1-q^n e^{2\pi i v})(1-q^n e^{-2\pi i v})$$

$$\vartheta_2(v|\tau) = 2q^{\frac{1}{8}} \cos[\pi v] \prod_{n=1}^{\infty} (1-q^n)(1+q^n e^{2\pi i v})(1+q^n e^{-2\pi i v})$$

$$\vartheta_3(v|\tau) = \prod_{n=1}^{\infty} (1-q^n)(1+q^{n-1/2} e^{2\pi i v})(1+q^{n-1/2} e^{-2\pi i v})$$

$$\vartheta_4(v|\tau) = \prod_{n=1}^{\infty} (1-q^n)(1-q^{n-1/2} e^{2\pi i v})(1-q^{n-1/2} e^{-2\pi i v})$$

$$\eta(\tau) = q^{\frac{1}{24}} \prod_{n=1}^{\infty} (1-q^n).$$

$$\left. \frac{\partial}{\partial v} \vartheta_1(v) \right|_{v=0} \equiv \vartheta_1' = 2\pi \eta^3(\tau)$$

$$\eta\left(-\frac{1}{\tau}\right) = \sqrt{-i\tau} \eta(\tau)$$

$$\vartheta\left[\begin{smallmatrix} a \\ b \end{smallmatrix}\right]\left(v+\frac{\epsilon_1}{2}\tau+\frac{\epsilon_2}{2}\middle|\tau\right)=e^{-\frac{i\pi\tau}{4}\epsilon_1^2-\frac{i\pi\epsilon_1}{2}(2v-b)-\frac{i\pi}{2}\epsilon_1\epsilon_2}\vartheta\left[\begin{smallmatrix} a-\epsilon_1 \\ b-\epsilon_2 \end{smallmatrix}\right](v|\tau)$$

$$\vartheta_2(0|\tau)\vartheta_3(0|\tau)\vartheta_4(0|\tau)=2\eta^3$$

$$\vartheta_2^4(v|\tau)-\vartheta_1^4(v|\tau)=\vartheta_3^4(v|\tau)-\vartheta_4^4(v|\tau)$$

$$\vartheta_2(2\tau)=\frac{1}{\sqrt{2}}\sqrt{\vartheta_3^2(\tau)-\vartheta_4^2(\tau)},\vartheta_3(2\tau)=\frac{1}{\sqrt{2}}\sqrt{\vartheta_3^2(\tau)+\vartheta_4^2(\tau)}$$

$$\vartheta_4(2\tau)=\sqrt{\vartheta_3(\tau)\vartheta_4(\tau)},\eta(2\tau)=\sqrt{\frac{\vartheta_2(\tau)\eta(\tau)}{2}}$$

$$\frac{1}{2}\sum_{a,b=0}^1(-1)^{a+b+ab}\prod_{i=1}^4\vartheta\left[\begin{smallmatrix} a \\ b \end{smallmatrix}\right](v_i)=-\prod_{i=1}^4\vartheta_1(v_i'),$$

$$v_1'=\frac{1}{2}(-v_1+v_2+v_3+v_4),\quad v_2'=\frac{1}{2}(v_1-v_2+v_3+v_4)$$

$$v_3'=\frac{1}{2}(v_1+v_2-v_3+v_4),\quad v_4'=\frac{1}{2}(v_1+v_2+v_3-v_4)$$

$$\frac{1}{2}\sum_{a,b=0}^1(-1)^{a+b+ab}\prod_{i=1}^4\vartheta\left[\begin{smallmatrix} a+h_i \\ b+g_i \end{smallmatrix}\right](v_i)=-\prod_{i=1}^4\vartheta\left[\begin{smallmatrix} 1-h_i \\ 1-g_i \end{smallmatrix}\right](v_i')$$

$$\frac{1}{2}\sum_{a,b=0}^1(-1)^{a+b}\prod_{i=1}^4\vartheta\left[\begin{smallmatrix} a \\ b \end{smallmatrix}\right](v_i)=-\prod_{i=1}^4\vartheta_1(v_i')+\prod_{i=1}^4\vartheta_1(v_i)$$



$$\frac{1}{2} \sum_{a,b=0}^1 (-1)^{a+b} \prod_{i=1}^4 \vartheta \left[\begin{smallmatrix} a+h_i \\ b+g_i \end{smallmatrix} \right] (v_i) = - \prod_{i=1}^4 \vartheta \left[\begin{smallmatrix} 1-h_i \\ 1-g_i \end{smallmatrix} \right] (v'_i) + \prod_{i=1}^4 \vartheta \left[\begin{smallmatrix} 1+h_i \\ 1+g_i \end{smallmatrix} \right] (v_i)$$

$$\left[\frac{1}{(2\pi i)^2}\frac{\partial^2}{\partial v^2}-\frac{1}{i\pi}\frac{\partial}{\partial \tau}\right]\vartheta\left[\begin{smallmatrix} a \\ b \end{smallmatrix}\right](v\mid \tau)=0$$

$$\frac{1}{4\pi i}\frac{\vartheta_2''}{\vartheta_2}=\partial_\tau\log\vartheta_2=\frac{i\pi}{12}(E_2+\vartheta_3^4+\vartheta_4^4)$$

$$\frac{1}{4\pi i}\frac{\vartheta_3''}{\vartheta_3}=\partial_\tau\log\vartheta_3=\frac{i\pi}{12}(E_2+\vartheta_2^4-\vartheta_4^4)$$

$$\frac{1}{4\pi i}\frac{\vartheta_4''}{\vartheta_4}=\partial_\tau\log\vartheta_4=\frac{i\pi}{12}(E_2-\vartheta_2^4-\vartheta_3^4)$$

$$\mathcal{P}(z)=4\pi i\partial_\tau\log\eta(\tau)-\partial_z^2\log\vartheta_1(z)=\frac{1}{z^2}+\mathcal{O}(z^2)$$

$$\mathcal{P}(-z)=\mathcal{P}(z)\,,\mathcal{P}(z+1)=\mathcal{P}(z+\tau)=\mathcal{P}(z)$$

$$\mathcal{P}(z,\tau+1)=\mathcal{P}(z,\tau)\,,\mathcal{P}\left(\frac{z}{\tau},-\frac{1}{\tau}\right)=\tau^2\mathcal{P}(z,\tau)$$

$$\int \frac{d^2z}{\tau_2} \mathcal{P}(z,\tau) = 4\pi i \partial_\tau \log \left(\sqrt{\tau_2} \eta \right)$$

$$\int \frac{d^2z}{\tau_2} |\mathcal{P}(z,\tau)|^2 = \left| 4\pi i \partial_\tau \log \left(\sqrt{\tau_2} \eta \right) \right|^2$$

$$\int \frac{d^2z}{\tau_2} \overline{\mathcal{P}}(\bar{z},\bar{\tau}) \left[\partial_z \log \vartheta_1(z) + 2\pi i \frac{Imz}{\tau_2} \right]^2 = 4\pi i \partial_\tau \log \left(\eta \sqrt{\tau_2} \right)$$

$$\int \frac{d^2z}{\tau_2} \partial_z^2 \log \vartheta_1(z) = -\frac{\pi}{\tau_2}$$

$$\tilde{f}(k)\equiv\frac{1}{2\pi}\int_{-\infty}^{+\infty}f(x)e^{ikx}dx$$

$$\sum_{n\in\mathbb{Z}}f(2\pi n)=\sum_{n\in\mathbb{Z}}\tilde{f}(n).$$

$$\sum_{n\in\mathbb{Z}}e^{-\pi an^2+\pi bn}=\frac{1}{\sqrt{a}}\sum_{n\in\mathbb{Z}}e^{-\frac{\pi}{a}\left(n+i\frac{b}{2}\right)^2},$$

$$\sum_{n\in\mathbb{Z}}ne^{-\pi an^2+\pi bn}=-\frac{i}{\sqrt{a}}\sum_{n\in\mathbb{Z}}\frac{\left(n+i\frac{b}{2}\right)}{a}e^{-\frac{\pi}{a}\left(n+i\frac{b}{2}\right)^2},$$

$$\sum_{n\in\mathbb{Z}}n^2e^{-\pi an^2+\pi bn}=\frac{1}{\sqrt{a}}\sum_{n\in\mathbb{Z}}\left[\frac{1}{2\pi a}-\frac{\left(n+i\frac{b}{2}\right)^2}{a^2}\right]e^{-\frac{\pi}{a}\left(n+i\frac{b}{2}\right)^2}.$$

$$\sum_{m_i\in\mathbb{Z}}e^{-\pi m_im_jA_{ij}+\pi B_im_i}=(\det A)^{-\frac{1}{2}}\sum_{m_i\in\mathbb{Z}}e^{-\pi(m_k+iB_k/2)(A^{-1})_{kl}(m_l+iB_l/2)}$$



$$S_{p,q} = \frac{1}{4\pi} \int d^2\sigma \sqrt{\det g} g^{ab} G_{\alpha\beta} \partial_a X^\alpha \partial_b X^\beta + \frac{1}{4\pi} \int d^2\sigma \epsilon^{ab} B_{\alpha\beta} \partial_a X^\alpha \partial_b X^\beta \\ + \frac{1}{4\pi} \int d^2\sigma \sqrt{\det g} \sum_I \psi^I (\bar{\nabla} + Y_\alpha^I (\bar{\nabla} X^\alpha)) \bar{\psi}^I$$

$$Z_{p,p+16}(G,B,Y) = \frac{\sqrt{\det G}}{\tau_2^{p/2} \eta^p \bar{\eta}^{p+16}} \\ \times \sum_{m^\alpha, n^\alpha \in \mathbb{Z}} \exp \left[-\frac{\pi}{\tau_2} (G+B)_{\alpha\beta} (m^\alpha + \tau n^\alpha) (m^\beta + \bar{\tau} n^\beta) \right] \\ \times \frac{1}{2} \sum_{a,b=0}^1 \prod_{l=1}^{16} e^{i\pi(m^\alpha Y_\alpha^l Y_\beta^l n^\beta - b n^\alpha Y_\alpha^l)} \bar{\vartheta} \left[\begin{matrix} a - 2n^\alpha Y_\alpha^l \\ b - 2m^\beta Y_\beta^l \end{matrix} \right] \\ = \frac{\sqrt{\det G}}{\tau_2^{p/2} \eta^p \bar{\eta}^{p+16}} \sum_{m^\alpha, n^\alpha \in \mathbb{Z}} \exp \left[-\frac{\pi}{\tau_2} (G+B)_{\alpha\beta} (m^\alpha + \tau n^\alpha) (m^\beta + \bar{\tau} n^\beta) \right] \\ \times \exp \left[-i\pi \sum_I n^\alpha (m^\beta + \bar{\tau} n^\beta) Y_\alpha^I Y_\beta^I \right] \frac{1}{2} \sum_{a,b=0}^1 \prod_{l=1}^{16} \bar{\vartheta} \left[\begin{matrix} a \\ b \end{matrix} \right] (Y_\gamma^I (m^\gamma + \bar{\tau} n^\gamma) | \bar{\tau})$$

$$\tau \rightarrow \tau + 1, Z_{p,p+16} \rightarrow e^{4\pi i/3} Z_{p,p+16}$$

$$\Gamma_{p,p+16}(G,B,Y) = \sum_{m_\alpha,n_\alpha,Q_I} q^{P_L^2/2} \bar{q}^{P_R^2/2}$$

$$M = \begin{pmatrix} G^{-1} & G^{-1}C & G^{-1}Y^t \\ C^t G^{-1} & G + C^t G^{-1}C + Y^t Y & C^t G^{-1}Y^t + Y^t \\ YG^{-1} & YG^{-1}C + Y & \mathbf{1}_{16} + YG^{-1}Y^t \end{pmatrix}$$

$$C_{\alpha\beta} = B_{\alpha\beta} - \frac{1}{2} Y_\alpha^I Y_\beta^I$$

$$L = \begin{pmatrix} 0 & \mathbf{1}_p & 0 \\ \mathbf{1}_p & 0 & 0 \\ 0 & 0 & \mathbf{1}_{16} \end{pmatrix}$$

$$M^T L M = M L M = L, M^{-1} = L M L$$

$$\frac{1}{2} P_L^2 = \frac{1}{4} (m^\alpha, n_\alpha, Q_I) \cdot (M - L) \cdot \begin{pmatrix} m^\alpha \\ n_\alpha \\ Q_I \end{pmatrix} \\ \frac{1}{2} P_R^2 = \frac{1}{4} (m^\alpha, n_\alpha, Q_I) \cdot (M + L) \cdot \begin{pmatrix} m^\alpha \\ n_\alpha \\ Q_I \end{pmatrix}.$$

$$\frac{1}{2} P_R^2 - \frac{1}{2} P_L^2 = m^\alpha n_\alpha - \frac{1}{2} Q_I Q_I$$

$$\Gamma_{p,p+16}(G,B,Y=0) = \Gamma_{p,p}(G,B) \bar{\Gamma}_{0(32)/Z_2}$$

$$\Gamma_{p,p+16}(G,B,Y=\tilde{Y}) = \Gamma_{p,p}(G',B') \bar{\Gamma}_{E_8 \times E_8}$$



$$\begin{pmatrix} m^\alpha \\ n_\alpha \\ Q_I \end{pmatrix} \rightarrow \Omega \cdot \begin{pmatrix} m^\alpha \\ n_\alpha \\ Q_I \end{pmatrix}, M \rightarrow \Omega M \Omega^T$$

$$Z_{d,d+16}^N(\epsilon) \begin{bmatrix} h \\ g \end{bmatrix} = \frac{\Gamma_{p,p+16}(\epsilon) \begin{bmatrix} h \\ g \end{bmatrix}}{\eta^p \bar{\eta}^{p+16}} = \frac{\sum_{\lambda \in L + \epsilon \frac{h}{N}} e^{\frac{2\pi i g \epsilon \cdot \lambda}{N}} q^{p_L^2/2} \bar{q}^{p_R^2/2}}{\eta^p \bar{\eta}^{p+16}}$$

$$\begin{aligned} Z^N(-\epsilon) \begin{bmatrix} h \\ g \end{bmatrix} &= Z^N(\epsilon) \begin{bmatrix} h \\ g \end{bmatrix}, Z^N(\epsilon) \begin{bmatrix} -h \\ -g \end{bmatrix} = Z^N(\epsilon) \begin{bmatrix} h \\ g \end{bmatrix} \\ Z^N(\epsilon) \begin{bmatrix} h+1 \\ g \end{bmatrix} &= \exp \left[-\frac{i\pi g \epsilon^2}{N} \right] Z^N(\epsilon) \begin{bmatrix} h \\ g \end{bmatrix}, Z^N(\epsilon) \begin{bmatrix} h \\ g+1 \end{bmatrix} = Z^N(\epsilon) \begin{bmatrix} h \\ g \end{bmatrix} \\ Z^N(\epsilon + N\epsilon') \begin{bmatrix} h \\ g \end{bmatrix} &= \exp \left[\frac{2\pi i g h \epsilon \cdot \epsilon'}{N} \right] Z^N(\epsilon) \begin{bmatrix} h \\ g \end{bmatrix} \end{aligned}$$

$$\begin{aligned} \tau \rightarrow \tau + 1 : Z^N(\epsilon) \begin{bmatrix} h \\ g \end{bmatrix} &\rightarrow \exp \left[\frac{4\pi i}{3} + \frac{i\pi h^2 \epsilon^2}{N^2} \right] Z^N(\epsilon) \begin{bmatrix} h \\ h+g \end{bmatrix} \\ \tau \rightarrow -\frac{1}{\tau} : Z^N(\epsilon) \begin{bmatrix} h \\ g \end{bmatrix} &\rightarrow \exp \left[-\frac{2\pi i h g \epsilon^2}{N^2} \right] Z^N(\epsilon) \begin{bmatrix} g \\ -h \end{bmatrix} \end{aligned}$$

$$Z^N(\epsilon, \Omega M \Omega^T) \begin{bmatrix} h \\ g \end{bmatrix} = Z^N(\Omega \cdot \epsilon, M) \begin{bmatrix} h \\ g \end{bmatrix}$$

$$\hat{e}_{\hat{\mu}}^{\hat{r}} = \begin{pmatrix} e_{\mu}^r & A_{\mu}^{\beta} E_{\beta}^{\alpha} \\ 0 & E_{\alpha}^{\alpha} \end{pmatrix}, \hat{e}_{\hat{r}}^{\hat{\mu}} = \begin{pmatrix} e_r^{\mu} & -e_r^{\nu} A_{\nu}^{\alpha} \\ 0 & E_{\alpha}^{\alpha} \end{pmatrix}$$

$$\hat{G}_{\hat{\mu}\hat{\nu}} = \begin{pmatrix} g_{\mu\nu} + A_{\mu}^{\alpha} G_{\alpha\beta} A_{\nu}^{\beta} & G_{\alpha\beta} A_{\mu}^{\beta} \\ G_{\alpha\beta} A_{\nu}^{\beta} & G_{\alpha\beta} \end{pmatrix}, \hat{G}^{\hat{\mu}\hat{\nu}} = \begin{pmatrix} g^{\mu\nu} & -A^{\mu\alpha} \\ -A^{\nu\alpha} & G^{\alpha\beta} + A_{\rho}^{\alpha} A^{\beta,\rho} \end{pmatrix}.$$

$$\alpha'^{D-2} S_D^{\text{heterotic}} = \int d^D x \sqrt{-\text{det} g} e^{-\phi} \left[R + \partial_{\mu} \phi \partial^{\mu} \phi + \frac{1}{4} \partial_{\mu} G_{\alpha\beta} \partial^{\mu} G^{\alpha\beta} + \right. \\ \left. - \frac{1}{4} G_{\alpha\beta} F_{\mu\nu}^{A\alpha} F_A^{\beta,\mu\nu} \right]$$

$$\begin{aligned} \phi &= \hat{\Phi} - \frac{1}{2} \log \left(\text{det} G_{\alpha\beta} \right) \\ F_{\mu\nu}^{A\alpha} &= \partial_{\mu} A_{\nu}^{\alpha} - \partial_{\nu} A_{\mu}^{\alpha} \end{aligned}$$

$$\begin{aligned} &-\frac{1}{12} \int d^{10}x \sqrt{-\text{det} \hat{G}} e^{-\hat{\Phi}} \hat{H}^{\hat{\mu}\hat{\nu}\hat{\rho}} \hat{H}_{\hat{\mu}\hat{\nu}\hat{\rho}} \\ &= - \int d^D x \sqrt{-\text{det} g} e^{-\phi} \left[\frac{1}{4} H_{\mu\alpha\beta} H^{\mu\alpha\beta} + \frac{1}{4} H_{\mu\nu\alpha} H^{\mu\nu\alpha} + \frac{1}{12} H_{\mu\nu\rho} H^{\mu\nu\rho} \right] \end{aligned}$$

$$H_{\mu\alpha\beta} = e_{\mu}^r \hat{e}_{\hat{r}}^{\hat{\mu}} \hat{H}_{\hat{\mu}\alpha\beta} = \hat{H}_{\mu\alpha\beta}$$

$$\begin{aligned} H_{\mu\nu\alpha} &= e_{\mu}^r e_{\nu}^s \hat{e}_{\hat{r}}^{\hat{\mu}} \hat{e}_{\hat{s}}^{\hat{\nu}} \hat{H}_{\hat{\mu}\hat{\nu}\alpha} = \hat{H}_{\mu\nu\alpha} - A_{\mu}^{\beta} \hat{H}_{\nu\alpha\beta} + A_{\nu}^{\beta} \hat{H}_{\mu\alpha\beta} \\ H_{\mu\nu\rho} &= e_{\mu}^r e_{\nu}^s e_t^t \hat{e}_{\hat{r}}^{\hat{\mu}} \hat{e}_{\hat{s}}^{\hat{\nu}} \hat{e}_{\hat{t}}^{\hat{\rho}} \hat{H}_{\hat{\mu}\hat{\nu}\hat{\rho}} = \hat{H}_{\mu\nu\rho} + \left[-A_{\mu}^{\alpha} \hat{H}_{\alpha\nu\rho} + A_{\mu}^{\alpha} A_{\nu}^{\beta} \hat{H}_{\alpha\beta\rho} + \text{cyclic} \right] \end{aligned}$$

$$\int d^{10}x \sqrt{-\det \hat{G}} e^{-\hat{\Phi}} \sum_{I=1}^{16} \hat{F}_{\hat{\mu}\hat{\nu}}^I F^{I,\hat{\mu}\hat{\nu}} = \int d^Dx \sqrt{-\det g} e^{-\phi} \sum_{I=1}^{16} [\tilde{F}_{\mu\nu}^I \tilde{F}^{I,\mu\nu} + 2\tilde{F}_{\mu\alpha}^I \tilde{F}^{I,\mu\alpha}]$$

$$Y_\alpha^I = \hat{A}_\alpha^I, A_\mu^I = \hat{A}_\mu^I - Y_\alpha^I A_\mu^\alpha, \tilde{F}_{\mu\nu}^I = F_{\mu\nu}^I + Y_\alpha^I F_{\mu\nu}^{A,\alpha}$$

$$\tilde{F}_{\mu\alpha}^I = \partial_\mu Y_\alpha^I, F_{\mu\nu}^I = \partial_\mu A_\nu^I - \partial_\nu A_\mu^I$$

$$\hat{H}_{\mu\alpha\beta} = \partial_\mu \hat{B}_{\alpha\beta} + \frac{1}{2} \sum_I [Y_\alpha^I \partial_\mu Y_\beta^I - Y_\beta^I \partial_\mu Y_\alpha^I]$$

$$C_{\alpha\beta} \equiv \hat{B}_{\alpha\beta} - \frac{1}{2} \sum_I Y_\alpha^I Y_\beta^I$$

$$H_{\mu\alpha\beta} = \partial_\mu C_{\alpha\beta} + \sum_I Y_\alpha^I \partial_\mu Y_\beta^I$$

$$\hat{H}_{\mu\nu\alpha} = \partial_\mu \hat{B}_{\nu\alpha} - \partial_\nu \hat{B}_{\mu\alpha} + \frac{1}{2} \sum_I [\hat{A}_\nu^I \partial_\mu Y_\alpha^I - \hat{A}_\mu^I \partial_\nu Y_\alpha^I - Y_\alpha^I \hat{F}_{\mu\nu}^I]$$

$$B_{\mu,\alpha} \equiv \hat{B}_{\mu\alpha} + B_{\alpha\beta} A_\mu^\beta + \frac{1}{2} \sum_I Y_\alpha^I A_\mu^I$$

$$F_{\alpha,\mu\nu}^B = \partial_\mu B_{\alpha,\nu} - \partial_\nu B_{\alpha,\mu}$$

$$H_{\mu\nu\alpha} = F_{\alpha\mu\nu}^B - C_{\alpha\beta} F_{\mu\nu}^{A,\beta} - \sum_I Y_\alpha^I F_{\mu\nu}^I$$

$$B_{\mu\nu} = \hat{B}_{\mu\nu} + \frac{1}{2} \left[A_\mu^\alpha B_{\nu\alpha} + \sum_I A_\mu^I A_\nu^\alpha Y_\alpha^I - (\mu \leftrightarrow \nu) \right] - A_\mu^\alpha A_\nu^\beta B_{\alpha\beta}$$

$$H_{\mu\nu\rho} = \partial_\mu B_{\nu\rho} - \frac{1}{2} \left[B_{\mu\alpha} F_{\nu\rho}^{A,\alpha} + A_\mu^\alpha F_{\alpha,\nu\rho}^B + \sum_I A_\mu^I F_{\nu\rho}^I \right] + \text{cyclic}$$

$$\equiv \partial_\mu B_{\nu\rho} - \frac{1}{2} L_{ij} A_\mu^i F_{\nu\rho}^j + \text{cyclic}$$

$$S_D^{\text{heterotic}} = \int d^Dx \sqrt{-\det g} e^{-\phi} \left[R + \partial^\mu \phi \partial_\mu \phi - \frac{1}{12} \tilde{H}^{\mu\nu\rho} \tilde{H}_{\mu\nu\rho} - \right.$$

$$\left. - \frac{1}{4} (M^{-1})_{ij} F_{\mu\nu}^i F^{j\mu\nu} + \frac{1}{8} \text{Tr}(\partial_\mu M \partial^\mu M^{-1}) \right]$$

$$S_C = -\frac{1}{2 \cdot 4!} \int d^d x \sqrt{-G} \hat{F}^2$$

$$\hat{F}_{\mu\nu\rho\sigma} = \partial_\mu \hat{C}_{\nu\rho\sigma} - \partial_\sigma \hat{C}_{\mu\nu\rho} + \partial_\rho \hat{C}_{\sigma\mu\nu} - \partial_\nu \hat{C}_{\rho\sigma\mu}$$

$$C_{\alpha\beta\gamma} = \hat{C}_{\alpha\beta\gamma}, C_{\mu\alpha\beta} = \hat{C}_{\mu\alpha\beta} - C_{\alpha\beta\gamma} A_\mu^\gamma$$

$$C_{\mu\nu\alpha} = \hat{C}_{\mu\nu\alpha} + \hat{C}_{\mu\alpha\beta} A_\nu^\beta - \hat{C}_{\nu\alpha\beta} A_\mu^\beta + C_{\alpha\beta\gamma} A_\mu^\alpha A_\nu^\gamma$$

$$C_{\mu\nu\rho} = \hat{C}_{\mu\nu\rho} + \left(-\hat{C}_{\nu\rho\alpha} A_\mu^\alpha + \hat{C}_{\alpha\beta\rho} A_\mu^\alpha A_\nu^\beta + \text{cyclic} \right) - C_{\alpha\beta\gamma} A_\mu^\alpha A_\nu^\beta A_\rho^\gamma$$



$$S_C = -\frac{1}{2 \cdot 4!} \int d^D x \sqrt{-g} \sqrt{\det G_{\alpha\beta}} [F_{\mu\nu\rho\sigma} F^{\mu\nu\rho\sigma} + 4F_{\mu\nu\rho\alpha} F^{\mu\nu\rho\alpha} + 6F_{\mu\nu\alpha\beta} F^{\mu\nu\alpha\beta} + 4F_{\mu\alpha\beta\gamma} F^{\mu\alpha\beta\gamma}]$$

$$\begin{aligned} F_{\mu\alpha\beta\gamma} &= \partial_\mu C_{\alpha\beta\gamma}, F_{\mu\nu\alpha\beta} = \partial_\mu C_{\nu\alpha\beta} - \partial_\nu C_{\mu\alpha\beta} + C_{\alpha\beta\gamma} F_{\mu\nu}^\gamma \\ F_{\mu\nu\rho\alpha} &= \partial_\mu C_{\nu\rho\alpha} + C_{\mu\alpha\beta} F_{\nu\rho}^\beta + \text{cyclic} \\ F_{\mu\nu\rho\sigma} &= (\partial_\mu C_{\nu\rho\sigma} + 3 \text{ perm.}) + (C_{\rho\sigma\alpha} F_{\mu\nu}^\alpha + 5 \text{ perm.}) \end{aligned}$$

$$\mathcal{L}_{N=1} = -\frac{1}{2\kappa^2}R + G_{i\bar{j}}D_\mu\phi^iD^\mu\bar{\phi}^{\bar{j}} + V(\phi,\bar{\phi}) + \sum_a\frac{1}{4g_a^2}[F_{\mu\nu}F^{\mu\nu}]_a + \frac{\theta_a}{4}[F_{\mu\nu}\tilde{F}^{\mu\nu}]_a$$

$$G_{i\bar{j}}=\partial_i\partial_{\bar{j}}K(\phi,\bar{\phi})$$

$$\frac{1}{g_a^2} = \text{Re} f_a(\phi) \, , \, \theta_a = -\text{Im} f_a(\phi)$$

$$V(\phi,\bar{\phi})=e^{\kappa^2K}(D_iWG^{i\bar{i}}\bar{D}_{\bar{i}}\bar{W}-3\kappa^2|W|^2)$$

$$D_iW=\frac{\partial W}{\partial \phi^i}+\kappa^2\frac{\partial K}{\partial \phi^i}W$$

$$K\rightarrow K+\Lambda(\phi)+\bar{\Lambda}(\bar{\phi})\,,W\rightarrow We^{-\Lambda}\,,f_a\rightarrow f_a$$

$$K=-\log\left[i(\bar{Z}^IF_I-Z^I\bar{F}_I)\right]$$

$$K=-\log\left[2\left(f(T^i)+\bar{f}(\bar{T}^i)\right)-\left(T^i-\bar{T}^i\right)(f_i-\bar{f}_i)\right]$$

$$R_{ij\bar{k}\bar{l}}=G_{ij}G_{k\bar{l}}+G_{i\bar{l}}G_{k\bar{j}}-e^{-2K}W_{ikm}G^{m\bar{m}}\bar{W}_{\bar{m}\bar{j}\bar{l}}$$

$$\mathcal{L}^{\text{vectors}}=-\frac{1}{4}\Xi_{IJ}F_{\mu\nu}^IF^{J,\mu\nu}-\frac{\theta_{IJ}}{4}F_{\mu\nu}^I\tilde{F}^{J,\mu\nu}$$

$$\Xi_{IJ}=\frac{i}{4}\left[N_{IJ}-\bar{N}_{IJ}\right],\theta_{IJ}=\frac{1}{4}\left[N_{IJ}+\bar{N}_{IJ}\right]$$

$$N_{IJ}=\bar{F}_{IJ}+2i\frac{\mathrm{Im}F_{IK}\mathrm{Im}F_{JL}Z^KZ^L}{\mathrm{Im}F_{MN}Z^MZ^N}$$

$$M_{BPS}^2=\frac{|e_I Z^I+q^I F_I|^2}{\mathrm{Im}(Z^I \bar{F}_I)}$$

$$M_{BPS}^2=\frac{1}{4\mathrm{Im}S}(\alpha^t+S\beta^t)M_+(\alpha+\bar{S}\beta)+\frac{1}{2}\sqrt{(\alpha^tM_+\alpha)(\beta^tM_+\beta)-(\alpha^tM_+\beta)^2},$$

$$\left\{Q^I_\alpha,Q^J_\beta\right\}=\epsilon_{\alpha\beta}Z^{IJ}\,,\left\{\bar{Q}^I_{\dot{\alpha}},Q^J_{\dot{\beta}}\right\}=\epsilon_{\dot{\alpha}\dot{\beta}}\bar{Z}^{IJ}\,,\left\{Q^I_\alpha,\bar{Q}^J_{\dot{\alpha}}\right\}=\delta^{IJ}2\sigma^\mu_{\alpha\dot{\alpha}}P_\mu$$

$$\left\{Q^I_\alpha,\bar{Q}^J_{\dot{\alpha}}\right\}=2M\delta_{\alpha\dot{\alpha}}\delta^{IJ}\,,\left\{Q^I_\alpha,Q^J_\beta\right\}=\left\{\bar{Q}^I_{\dot{\alpha}},\bar{Q}^J_{\dot{\beta}}\right\}=0$$



$$A_{\alpha}^I = \frac{1}{\sqrt{2M}} Q_{\alpha}^I, A_{\alpha}^{\dagger I} = \frac{1}{\sqrt{2M}} \bar{Q}_{\dot{\alpha}}^I$$

$$\{Q_{\alpha}^I, \bar{Q}_{\dot{\alpha}}^J\} = 2 \begin{pmatrix} 2E & 0 \\ 0 & 0 \end{pmatrix} \delta^{IJ}$$

$$\{Q_{\alpha}^{am}, \bar{Q}_{\dot{\alpha}}^{bn}\} = 2M\delta^{\alpha\dot{\alpha}}\delta^{ab}\delta^{mn}, \{Q_{\alpha}^{am}, Q_{\beta}^{bn}\} = Z_n\epsilon^{\alpha\beta}\epsilon^{ab}\delta^{mn}$$

$$A_{\alpha}^m = \frac{1}{\sqrt{2}}[Q_{\alpha}^{1m} + \epsilon_{\alpha\beta}Q_{\beta}^{2m}], B_{\alpha}^m = \frac{1}{\sqrt{2}}[Q_{\alpha}^{1m} - \epsilon_{\alpha\beta}Q_{\beta}^{2m}],$$

$$\{A_{\alpha}^m, A_{\beta}^n\} = \{A_{\alpha}^m, B_{\beta}^n\} = \{B_{\alpha}^m, B_{\beta}^n\} = 0, \\ \{A_{\alpha}^m, A_{\beta}^{\dagger n}\} = \delta_{\alpha\beta}\delta^{mn}(2M + Z_n), \{B_{\alpha}^m, B_{\beta}^{\dagger n}\} = \delta_{\alpha\beta}\delta^{mn}(2M - Z_n)$$

$$M \geq \max\left[\frac{Z_n}{2}\right]$$

$$B_{2n}(R) = \mathrm{Tr}_R[(-1)^{2\lambda}\lambda^{2n}]$$

$$Z_R(y) = \mathrm{stry}^{2\lambda}$$

$$Z_{[j]} = \begin{cases} (-)^{2j} \left(\frac{y^{2j+1} - y^{-2j-1}}{y - 1/y} \right) & \text{massive} \\ (-)^{2j} (y^{2j} + y^{-2j}) & \text{supermassive} \end{cases}.$$

$$Z_{r\otimes\tilde{r}}=Z_rZ_{\tilde{r}}$$

$$B_n(R)=\left(y^2\frac{d}{dy^2}\right)^nZ_R(y)\Big|_{y=1}$$

$$Z_m(y) = Z_{[j]}(y)(1-y)^m(1-1/y)^m$$

$$L_j:~[j] \otimes ([1] + 4[1/2] + 5[0])$$

$$S_j:~[j] \otimes (2[1/2] + 4[0])$$

$$M_{\lambda}^0\colon \pm\,(\lambda+1/2)+2(\pm\lambda)+\pm(\lambda-1/2)$$

$$B_0(\text{ any rep }) = 0 \\ B_2(M_{\lambda}^0) = (-1)^{2\lambda+1}, B_2(S_j) = (-1)^{2j+1}D_j, B_2(L_j) = 0.$$

$$L_j:~[j] \otimes (42[0] + 48[1/2] + 27[1] + 8[3/2] + [2])$$

$$s=2 \text{ supermassive long}: 42[0] + 48[1/2] + 27[1] + 8[3/2] + [2]$$

$$I_j\colon [j] \otimes (14[0] + 14[1/2] + 6[1] + [3/2])$$

$$I_{3/2}\colon 14[0] + 14[1/2] + 6[1] + [3/2]$$

$$S_j:~[j] \otimes (5[0] + 4[1/2] + [1]),$$



$$S_1: 5[0] + 4[1/2] + [1]$$

$$M_\lambda^0: [\pm(\lambda + 1)] + 4[\pm(\lambda + 1/2)] + 6[\pm(\lambda)] + 4[\pm(\lambda - 1/2)] + [\pm(\lambda - 1)],$$

$$M_1^0: [\pm 2] + 4[\pm 3/2] + 6[\pm 1] + 4[\pm 1/2] + 2[0]$$

$$L_j \rightarrow 2I_j + I_{j+1/2} + I_{j-1/2}$$

$$I_j \rightarrow 2S_j + S_{j+1/2} + S_{j-1/2}$$

$$S_j \rightarrow \sum_{\lambda=0}^j M_\lambda^0, j - \lambda \in \mathbb{Z}$$

$$B_n(\text{ any rep }) = 0 \text{ for } n = 0, 2.$$

$$B_4(L_j) = B_4(I_j) = 0, B_4(S_j) = (-1)^{2j} \frac{3}{2} D_j$$

$$B_4(M_\lambda^0) = (-1)^{2\lambda} 3, B_4(V^0) = \frac{3}{2}$$

$$B_6(L_j) = 0, B_6(I_j) = (-1)^{2j+1} \frac{45}{4} D_j, B_6(S_j) = (-1)^{2j} \frac{15}{8} D_j^3$$

$$B_6(M_\lambda^0) = (-1)^{2\lambda} \frac{15}{4} (1 + 12\lambda^2), B_6(V^0) = \frac{15}{8}$$

$$B_8(L_j) = (-1)^{2j} \frac{315}{4} D_j, B_8(I_j) = (-1)^{2j+1} \frac{105}{16} D_j (1 + D_j^2)$$

$$B_8(S_j) = (-1)^{2j} \frac{21}{64} D_j (1 + 2D_j^4)$$

$$B_8(M_\lambda^0) = (-1)^{2\lambda} \frac{21}{16} (1 + 80\lambda^2 + 160\lambda^4), B_8(V^0) = \frac{63}{32}$$

$$(\lambda \pm 2) + 8\left(\lambda \pm \frac{3}{2}\right) + 28(\lambda \pm 1) + 56\left(\lambda \pm \frac{1}{2}\right) + 70(\lambda).$$

$$(\pm 2) + 8\left(\pm \frac{3}{2}\right) + 28(\pm 1) + 56\left(\pm \frac{1}{2}\right) + 70(0)$$

$$[j] \otimes ([2] + 8[3/2] + 27[1] + 48[1/2] + 42[0])$$

$$S^j \rightarrow \sum_{\lambda=0}^j M_0^\lambda$$

$$I_1^j: [j] \otimes ([5/2] + 10[2] + 44[3/2] + 110[1] + 165[1/2] + 132[0])$$

$$I_2^j: [j] \otimes ([3] + 12[5/2] + 65[2] + 208[3/2] + 429[1] + 572[1/2] + 429[0])$$

$$I_3^j: [j] \otimes ([7/2] + 14[3] + 90[5/2] + 350[2] + 910[3/2] + 1638[1] + 2002[1/2] + 1430[0])$$



$$[j] \otimes ([4] + 16[7/2] + 119[3] + 544[5/2] + 1700[2] + 3808[3/2] + 6188[1] + 7072[1/2] + 4862[0])$$

$$L^j \rightarrow I_3^{j+\frac{1}{2}} + 2I_3^j + I_3^{j-\frac{1}{2}}$$

$$I_3^j \rightarrow I_2^{j+\frac{1}{2}} + 2I_2^j + I_2^{j-\frac{1}{2}}$$

$$I_2^j \rightarrow I_1^{j+\frac{1}{2}} + 2I_1^j + I_1^{j-\frac{1}{2}}$$

$$I_1^j \rightarrow S^{j+\frac{1}{2}} + 2S^j + S^{j-\frac{1}{2}}$$

$$B_8(M_0^\lambda) = (-1)^{2\lambda} 315$$

$$B_{10}(M_0^\lambda) = (-1)^{2\lambda} \frac{4725}{2} (6\lambda^2 + 1)$$

$$B_{12}(M_0^\lambda) = (-1)^{2\lambda} \frac{10395}{16} (240\lambda^4 + 240\lambda^2 + 19)$$

$$B_{14}(M_0^\lambda) = (-1)^{2\lambda} \frac{45045}{16} (336\lambda^6 + 840\lambda^4 + 399\lambda^2 + 20)$$

$$B_{16}(M_0^\lambda) = (-1)^{2\lambda} \frac{135135}{256} (7680\lambda^8 + 35840\lambda^6 + 42560\lambda^4 + 12800\lambda^2 + 457)$$

$$B_8(S^j) = (-1)^{2j} \cdot \frac{315}{2} D_j,$$

$$B_{10}(S^j) = (-1)^{2j} \cdot \frac{4725}{8} D_j (D_j^2 + 1),$$

$$B_{12}(S^j) = (-1)^{2j} \cdot \frac{10395}{32} D_j (3D_j^4 + 10D_j^2 + 6),$$

$$B_{14}(S^j) = (-1)^{2j} \cdot \frac{45045}{128} D_j (3D_j^6 + 21D_j^4 + 42D_j^2 + 14),$$

$$B_{16}(S^j) = (-1)^{2j} \cdot \frac{45045}{512} D_j (10D_j^8 + 120D_j^6 + 504D_j^4 + 560D_j^2 + 177),$$

$$B_8(I_1^j) = 0$$

$$B_{10}(I_1^j) = (-1)^{2j+1} \cdot \frac{14175}{4} D_j$$

$$B_{12}(I_1^j) = (-1)^{2j+1} \cdot \frac{155925}{16} D_j (2D_j^2 + 3)$$

$$B_{14}(I_1^j) = (-1)^{2j+1} \cdot \frac{2837835}{64} D_j (D_j^2 + 1)(D_j^2 + 4)$$

$$B_{16}(I_1^j) = (-1)^{2j+1} \cdot \frac{2027025}{128} D_j (4D_j^6 + 42D_j^4 + 112D_j^2 + 57)$$

$$B_8(I_3^j) = B_{10}(I_3^j) = B_{12}(I_3^j) = 0$$



$$\begin{aligned}
B_{14}(I_3^j) &= (-1)^{2j+1} \cdot \frac{42567525}{8} D_j, \\
B_{16}(I_3^j) &= (-1)^{2j+1} \cdot \frac{212837625}{8} D_j (2D_j^2 + 5), \\
B_8(L^j) &= B_{10}(L^j) = B_{12}(L^j) = B_{14}(L^j) = 0, \\
B_{16}(L^j) &= (-1)^{2j} \cdot \frac{638512875}{2} D_j.
\end{aligned}$$

$$F_d(-1/\tau) = \tau^d F_d(\tau) F_d(\tau + 1) = F_d(\tau)$$

$$\begin{aligned}
E_2 &= \frac{12}{i\pi} \partial_\tau \log \eta = 1 - 24 \sum_{n=1}^{\infty} \frac{nq^n}{1-q^n} \\
E_4 &= \frac{1}{2} (\vartheta_2^8 + \vartheta_3^8 + \vartheta_4^8) = 1 + 240 \sum_{n=1}^{\infty} \frac{n^3 q^n}{1-q^n} \\
E_6 &= \frac{1}{2} (\vartheta_2^4 + \vartheta_3^4)(\vartheta_3^4 + \vartheta_4^4)(\vartheta_4^4 - \vartheta_2^4) = 1 - 504 \sum_{n=1}^{\infty} \frac{n^5 q^n}{1-q^n}.
\end{aligned}$$

$$\begin{aligned}
H_2 &\equiv \frac{1-E_2}{24} = \sum_{n=1}^{\infty} \frac{nq^n}{1-q^n} \equiv \sum_{n=1}^{\infty} d_2(n)q^n \\
H_4 &\equiv \frac{E_4-1}{240} = \sum_{n=1}^{\infty} \frac{n^3q^n}{1-q^n} \equiv \sum_{n=1}^{\infty} d_4(n)q^n \\
H_6 &\equiv \frac{1-E_6}{504} = \sum_{n=1}^{\infty} \frac{n^5q^n}{1-q^n} \equiv \sum_{n=1}^{\infty} d_6(n)q^n
\end{aligned}$$

$$d_{2k}(N)=\sum_{n|N}n^{2k-1}, k=1,2,3$$

$$\hat{E}_2=E_2-\frac{3}{\pi\tau_2}$$

$$j=\frac{E_4^3}{\eta^{24}}=\frac{1}{q}+744+\cdots, \eta^{24}=\frac{1}{2^6\cdot 3^3}[E_4^3-E_6^2].$$

$$F_{d+2}=\left(\frac{i}{\pi}\partial_\tau+\frac{d/2}{\pi\tau_2}\right)F_d\equiv D_dF_d$$

$$D_{d_1+d_2}(F_{d_1}F_{d_2})=F_{d_2}(D_{d_1}F_{d_1})+F_{d_1}(D_{d_2}F_{d_2})$$

$$D_2\hat{E}_2=\frac{1}{6}E_4-\frac{1}{6}\hat{E}_2^2, D_4E_4=\frac{2}{3}E_6-\frac{2}{3}\hat{E}_2E_4, D_6E_6=E_4^2-\hat{E}_2E_6$$

$$\begin{aligned}\frac{\vartheta_1'''}{\vartheta_1'} &= -\pi^2 E_2, \frac{\vartheta_1^{(5)}}{\vartheta_1'} = -\pi^2 E_2 (4\pi i \partial_\tau \log E_2 - \pi^2 E_2) \\ &\quad - 3 \frac{\vartheta_1^{(5)}}{\vartheta_1'} + 5 \left(\frac{\vartheta_1'''}{\vartheta_1'} \right)^2 = 2\pi^4 E_4 \\ &\quad - 15 \frac{\vartheta_1^{(7)}}{\vartheta_1'} - \frac{350}{3} \left(\frac{\vartheta_1'''}{\vartheta_1'} \right)^3 + 105 \frac{\vartheta_1^{(5)} \vartheta_1'''}{\vartheta_1'^2} = \frac{80\pi^6}{3} E_6 \\ &\quad \frac{1}{2} \sum_{i=2}^4 \frac{\vartheta_i'' \vartheta_i^7}{(2\pi i)^2} = \frac{1}{12} (E_2 E_4 - E_6)\end{aligned}$$

$$\xi(v) = \prod_{n=1}^{\infty} \frac{(1-q^n)^2}{(1-q^n e^{2\pi i v})(1-q^n e^{-2\pi i v})} = \frac{\sin \pi v}{\pi} \frac{\vartheta_1'}{\vartheta_1(v)} \quad \xi(v) = \xi(-v)$$

$$\xi(0) = 1, \xi^{(2)}(0) = -\frac{1}{3} \left(\pi^2 + \frac{\vartheta_1'''}{\vartheta_1'} \right) = -\frac{\pi^2}{3} (1 - E_2)$$

$$\xi^{(4)}(0) = \frac{\pi^4}{5} + \frac{2\pi^2}{3} \frac{\vartheta_1'''}{\vartheta_1'} + \frac{2}{3} \left(\frac{\vartheta_1'''}{\vartheta_1'} \right)^2 - \frac{1}{5} \frac{\vartheta_1^{(5)}}{\vartheta_1'} = \frac{\pi^4}{15} (3 - 10E_2 + 2E_4 + 5E_2^2)$$

$$\begin{aligned}\xi^{(6)}(0) &= -\frac{\pi^6}{7} - \pi^4 \frac{\vartheta_1'''}{\vartheta_1'} - \frac{10\pi^2}{3} \left(\frac{\vartheta_1'''}{\vartheta_1'} \right)^2 + \pi^2 \frac{\vartheta_1^{(5)}}{\vartheta_1'} + \\ &\quad - \frac{10}{3} \left(\frac{\vartheta_1'''}{\vartheta_1'} \right)^3 + 2 \frac{\vartheta_1^{(5)} \vartheta_1'''}{\vartheta_1'^2} - \frac{1}{7} \frac{\vartheta_1^{(7)}}{\vartheta_1'} \\ &= \frac{\pi^6}{63} (-9 + 63E_2 - 105E_2^2 - 42E_4 + 16E_6 + \\ &\quad + 42E_2 E_4 + 35E_2^3)\end{aligned}$$

$$Z(v, \bar{v}) = \text{Str}[q^{L_0} \bar{q}^{\bar{L}_0} e^{2\pi i v \lambda_R - 2\pi i \bar{v} \lambda_L}]$$

$$Z_{D=4}^{\text{heterotic}} = \frac{1}{\tau_2 \eta^2 \bar{\eta}^2} \sum_{a,b=0}^1 (-1)^{a+b+ab} \frac{\vartheta \left[\begin{smallmatrix} a \\ b \end{smallmatrix} \right]}{\eta} C^{\text{int}} \left[\begin{smallmatrix} a \\ b \end{smallmatrix} \right]$$

$$Z_{D=4}^{\text{heterotic}}(v, \bar{v}) = \frac{\xi(v) \bar{\xi}(\bar{v})}{\tau_2 \eta^2 \bar{\eta}^2} \sum_{a,b=0}^1 (-1)^{a+b+ab} \frac{\vartheta \left[\begin{smallmatrix} a \\ b \end{smallmatrix} \right](v)}{\eta} C^{\text{int}} \left[\begin{smallmatrix} a \\ b \end{smallmatrix} \right]$$

$$Z_{D=4}^{\text{heterotic}}(v, \bar{v}) = \frac{\xi(v) \bar{\xi}(\bar{v})}{\tau_2 \eta^2 \bar{\eta}^2} \frac{\vartheta \left[\begin{smallmatrix} 1 \\ 1 \end{smallmatrix} \right](v/2)}{\eta} C^{\text{int}} \left[\begin{smallmatrix} 1 \\ 1 \end{smallmatrix} \right](v/2)$$

$$C^{\text{int}} \left[\begin{smallmatrix} 1 \\ 1 \end{smallmatrix} \right](v) = \text{Tr}_R \left[(-1)^{F^{\text{int}}} e^{2\pi i v J_0} q^{L_0^{\text{int}} - 3/8} \bar{q}^{\bar{L}_0^{\text{int}} - 11/12} \right]$$

$$Q = \frac{1}{2\pi i} \frac{\partial}{\partial v}, \bar{Q} = -\frac{1}{2\pi i} \frac{\partial}{\partial \bar{v}}$$

$$B_{2n} \equiv \text{Str}[\lambda^{2n}] = (Q + \bar{Q})^{2n} Z_{D=4}^{\text{heterotic}}(v, \bar{v})|_{v=\bar{v}=0}$$



$$Z_{N=4}^{\text{heterotic}}(v, \bar{v}) = \frac{\vartheta_1^4(v/2)}{\eta^{12} \bar{\eta}^{24}} \xi(v) \bar{\xi}(\bar{v}) \frac{\Gamma_{6,22}}{\tau_2}$$

$$B_4 = \langle (Q + \bar{Q})^4 \rangle = \langle Q^4 \rangle = \frac{3}{2} \frac{1}{\bar{\eta}^{24}}$$

$$B_6 = \langle (Q + \bar{Q})^6 \rangle = \langle Q^6 + 15Q^4\bar{Q}^2 \rangle = \frac{15}{8} \frac{2 - \bar{E}_2}{\bar{\eta}^{24}}$$

$$Z_{N=8}^{II}(v, \bar{v}) = \text{Str}[q^{L_0} \bar{q}^{\bar{L}_0} e^{2\pi i v \lambda_R - 2\pi i \bar{v} \lambda_L}]$$

$$\begin{aligned} &= \frac{1}{4} \sum_{\alpha, \beta=0}^1 \sum_{\bar{\alpha}, \bar{\beta}=0}^1 (-1)^{\alpha+\beta+\alpha\bar{\beta}} \frac{\vartheta \left[\begin{smallmatrix} \alpha \\ \beta \end{smallmatrix} \right] (v) \vartheta^3 \left[\begin{smallmatrix} \alpha \\ \beta \end{smallmatrix} \right] (0)}{\eta^4} \\ &\times (-1)^{\bar{\alpha}+\bar{\beta}+\bar{\alpha}\bar{\beta}} \frac{\bar{\vartheta} \left[\begin{smallmatrix} \alpha \\ \beta \end{smallmatrix} \right] \bar{\vartheta}^3 \left[\begin{smallmatrix} \bar{\alpha} \\ \bar{\beta} \end{smallmatrix} \right] (0)}{\bar{\eta}^4} \frac{\xi(v) \bar{\xi}(\bar{v})}{\text{Im}\tau |\eta|^4} \frac{\Gamma_{6,6}}{|\eta|^{12}} = \frac{\Gamma_{6,6}}{\text{Im}\tau} \frac{\vartheta_1^4(v/2)}{\eta^{12}} \frac{\bar{\vartheta}_1^4(\bar{v}/2)}{\bar{\eta}^{12}} \xi(v) \bar{\xi}(\bar{v}) \end{aligned}$$

$$B_8 = \langle (Q + \bar{Q})^8 \rangle = 70 \langle Q^4 \bar{Q}^4 \rangle = \frac{315}{2} \frac{\Gamma_{6,6}}{\text{Im}\tau}$$

$$M^2 = \frac{1}{4} p_L^2, \vec{m} \cdot \vec{n} = 0$$

$$\begin{aligned} B_{10} &= \langle (Q + \bar{Q})^{10} \rangle = 210 \langle Q^6 \bar{Q}^4 + Q^4 \bar{Q}^6 \rangle \\ &= -\frac{4725}{8\pi^2} \frac{\Gamma_{6,6}}{\text{Im}\tau} \left(\frac{\vartheta_1'''}{\vartheta_1'} + 3\xi'' + cc \right) = \frac{4725}{4} \frac{\Gamma_{6,6}}{\text{Im}\tau} \end{aligned}$$

$$\begin{aligned} B_{12} &= \langle 495(Q^4 \bar{Q}^8 + Q^8 \bar{Q}^4) + 924Q^6 \bar{Q}^6 \rangle = \left[\frac{10395}{2} + \frac{31185}{64} (E_4 + \bar{E}_4) \right] \frac{\Gamma_{6,6}}{\text{Im}\tau} \\ &= \left[\frac{10395 \cdot 19}{32} + \frac{10395 \cdot 45}{4} \left(\frac{E_4 - 1}{240} + cc \right) \right] \frac{\Gamma_{6,6}}{\text{Im}\tau} \end{aligned}$$

$$I_2^j: \sum_j (-1)^{2j} D_j = d_4(N)$$

$$\begin{aligned} B_{14} &= \langle (Q + \bar{Q})^{14} \rangle = \\ &= \left[\frac{45045}{32} 20 + \frac{14189175}{16} \left(2 \frac{E_4 - 1}{240} + \frac{1 - E_6}{504} + cc \right) \right] \frac{\Gamma_{6,6}}{\text{Im}\tau} \end{aligned}$$

$$I_2^j: \sum_j (-1)^{2j} D_j^3 = d_6(N)$$

$$\begin{aligned} Z^{II} &= \frac{1}{8} \sum_{g,h=0}^1 \sum_{\alpha,\beta=0}^1 \sum_{\bar{\alpha},\bar{\beta}=0}^1 (-1)^{\alpha+\beta+\alpha\bar{\beta}} \frac{\vartheta^2 \left[\begin{smallmatrix} \alpha \\ \beta \end{smallmatrix} \right]}{\eta^2} \frac{\vartheta \left[\begin{smallmatrix} \alpha+h \\ \beta+g \end{smallmatrix} \right]}{\eta} \frac{\vartheta \left[\begin{smallmatrix} \alpha-h \\ \beta-g \end{smallmatrix} \right]}{\eta} \times \\ &(-1)^{\bar{\alpha}+\bar{\beta}+\bar{\alpha}\bar{\beta}} \frac{\bar{\vartheta} \left[\begin{smallmatrix} \bar{\alpha} \\ \bar{\beta} \end{smallmatrix} \right]}{\bar{\eta}^2} \frac{\bar{\vartheta} \left[\begin{smallmatrix} \bar{\alpha}+h \\ \bar{\beta}+g \end{smallmatrix} \right]}{\bar{\eta}} \frac{\bar{\vartheta} \left[\begin{smallmatrix} \bar{\alpha}-h \\ \bar{\beta}-g \end{smallmatrix} \right]}{\bar{\eta}} \frac{1}{\text{Im}\tau |\eta|^4} \frac{\Gamma_{2,2}}{|\eta|^4} Z_{4,4} [h] \end{aligned}$$



$$Z_{4,4} \begin{bmatrix} 0 \\ 0 \end{bmatrix} = \frac{\Gamma_{4,4}}{|\eta|^8}, Z_{4,4} \begin{bmatrix} 0 \\ 1 \end{bmatrix} = 16 \frac{|\eta|^4}{|\vartheta_2|^4} = \frac{|\vartheta_3 \vartheta_4|^4}{|\eta|^8}$$

$$Z_{4,4} \begin{bmatrix} 1 \\ 0 \end{bmatrix} = 16 \frac{|\eta|^4}{|\vartheta_4|^4} = \frac{|\vartheta_2 \vartheta_3|^4}{|\eta|^8}, Z_{4,4} \begin{bmatrix} 1 \\ 1 \end{bmatrix} = 16 \frac{|\eta|^4}{|\vartheta_3|^4} = \frac{|\vartheta_2 \vartheta_4|^4}{|\eta|^8}$$

$$Z^H(v, \bar{v}) = \frac{1}{4} \sum_{\alpha \beta \bar{\alpha} \bar{\beta}} (-1)^{\alpha + \beta + \alpha \beta + \bar{\alpha} + \bar{\beta} + \bar{\alpha} \bar{\beta}} \frac{\vartheta \begin{bmatrix} \alpha \\ \beta \end{bmatrix} (v) \vartheta \begin{bmatrix} \alpha \\ \beta \end{bmatrix} (0)}{\eta^6}$$

$$\times \frac{\bar{\vartheta} \begin{bmatrix} \bar{\alpha} \\ \bar{\beta} \end{bmatrix} (\bar{v}) \bar{\vartheta} \begin{bmatrix} \bar{\alpha} \\ \bar{\beta} \end{bmatrix} (0)}{\bar{\eta}^6} \xi(v) \bar{\xi}(\bar{v}) C \begin{bmatrix} \alpha & \bar{\alpha} \\ \beta & \bar{\beta} \end{bmatrix} \frac{\Gamma_{2,2}}{\tau_2}$$

$$= \frac{\vartheta_1^2(v/2) \bar{\vartheta}_1^2(\bar{v}/2)}{\eta^6 \bar{\eta}^6} \xi(v) \bar{\xi}(\bar{v}) C \begin{bmatrix} 1 & 1 \\ 1 & 1 \end{bmatrix} (v/2, \bar{v}/2) \frac{\Gamma_{2,2}}{\tau_2}$$

$$C \begin{bmatrix} 1 & 1 \\ 1 & 1 \end{bmatrix} (v, 0) = 8 \sum_{i=2}^4 \frac{\vartheta_i^2(v)}{\vartheta_i^2(0)}$$

$$\langle \lambda^4 \rangle = \langle (Q + \bar{Q})^4 \rangle = 6 \langle Q^2 \bar{Q}^2 + Q^2 \bar{Q}^4 \rangle = 36 \frac{\Gamma_{2,2}}{\tau_2}.$$

$$\langle \lambda^6 \rangle = \langle (Q + \bar{Q})^6 \rangle = 15 \langle Q^4 \bar{Q}^2 + Q^2 \bar{Q}^4 \rangle = 90 \frac{\Gamma_{2,2}}{\tau_2},$$

$$\partial_v^2 C \begin{bmatrix} 1 & 1 \\ 1 & 1 \end{bmatrix} (v, 0) \Big|_{v=0} = -16\pi^2 E_2$$

$$L_{\rm gauge} = -\frac{1}{8} {\rm Im} \int d^4x \sqrt{-{\rm det} g} {\bf F}^i_{\mu\nu} N_{ij} {\bf F}^{j,\mu\nu}$$

$${\bf F}_{\mu\nu} = F_{\mu\nu} + i {}^\star F_{\mu\nu}, \quad {}^\star F_{\mu\nu} = \frac{1}{2} \frac{\epsilon_{\mu\nu}{}^{\rho\sigma}}{\sqrt{-g}} F_{\rho\sigma}$$

$$L_{\rm gauge} = -\frac{1}{4} \int d^4x [\sqrt{-g} F_{\mu\nu}^i N_2^{ij} F^{j,\mu\nu} + F_{\mu\nu}^i N_1^{ij} {}^\star F^{j,\mu\nu}]$$

$$\mathbf{G}^i_{\mu\nu} = N_{ij} \mathbf{F}^j_{\mu\nu} = N_1 F - N_2 {}^\star F + i(N_2 F + N_1 {}^\star F)$$

$$\mathrm{Im}\nabla^\mu\begin{pmatrix}\mathbf{G}^i_{\mu\nu}\\\mathbf{F}^i_{\mu\nu}\end{pmatrix}=\begin{pmatrix}0\\0\end{pmatrix}$$

$$\begin{pmatrix} \mathbf{G}'_{\mu\nu} \\ \mathbf{F}'_{\mu\nu} \end{pmatrix} = \begin{pmatrix} A & B \\ C & D \end{pmatrix} \begin{pmatrix} \mathbf{G}_{\mu\nu} \\ \mathbf{F}_{\mu\nu} \end{pmatrix}$$

$$N'=(AN+B)(CN+D)^{-1}$$

$$F' = C(N_1 F - N_2 {}^\star F) + D F, \quad {}^\star F' = C(N_2 F + N_1 {}^\star F) + D {}^\star F$$



$$F' = N_1 F - N_2^* F, {}^* F' = N_2 F + N_1^* F, N' = -\frac{1}{N}.$$

$$\begin{pmatrix} A & B \\ C & D \end{pmatrix} = \begin{pmatrix} \mathbf{1} - e & -e \\ e & \mathbf{1} - e \end{pmatrix}, e = \begin{pmatrix} 1 & 0 & \dots \\ 0 & 0 & \dots \\ . & . & . \end{pmatrix},$$

$$N'_{00} = -\frac{1}{N_{00}}, N'_{0i} = \frac{N_{0i}}{N_{00}}, N'_{i0} = \frac{N_{i0}}{N_{00}}, N'_{ij} = N_{ij} - \frac{N_{i0}N_{0j}}{N_{00}}.$$

$$\begin{pmatrix} A & B \\ C & D \end{pmatrix} = \begin{pmatrix} \mathbf{1} - e_1 & e_2 \\ -e_2 & \mathbf{1} - e_1 \end{pmatrix},$$

$$e_1 = \begin{pmatrix} 1 & 0 & 0 & \dots \\ 0 & 1 & 0 & \dots \\ 0 & 0 & 0 & . \\ . & . & . & . \end{pmatrix}, e_2 = \begin{pmatrix} 0 & 1 & 0 & \dots \\ -1 & 0 & 0 & \dots \\ 0 & 0 & 0 & . \\ . & . & . & . \end{pmatrix}.$$

$$N'_{\alpha\beta} = -\frac{N_{\alpha\beta}}{\det N_{\alpha\beta}} N'_{\alpha i} = -\frac{N_{\alpha\beta} \epsilon^{\beta\gamma} N_{\gamma i}}{\det N_{\alpha\beta}}, N'_{i\alpha} = \frac{N_{i\beta} \epsilon^{\beta\gamma} N_{\alpha\gamma}}{\det N_{\alpha\beta}}$$

$$N'_{ij} = N_{ij} + \frac{N_{i\alpha} \epsilon^{\alpha\beta} N_{\beta\gamma} \epsilon^{\gamma\delta} N_{\delta j}}{\det N_{\alpha\beta}}$$

$$N = S_1 L + i S_2 M^{-1}, S = S_1 + i S_2$$

$$N' = -N^{-1} = -\frac{S_1}{|S|^2} L + i \frac{S_2}{|S|^2} M = -\frac{S_1}{|S|^2} L + i \frac{S_2}{|S|^2} L M^{-1} L$$

$$N' = -LN^{-1}L = -\frac{S_1}{|S|^2} L + i \frac{S_2}{|S|^2} M^{-1}$$

$$\begin{pmatrix} A & B \\ C & D \end{pmatrix} = \begin{pmatrix} a\mathbf{1}_{28} & bL \\ cL & d\mathbf{1}_{28} \end{pmatrix}, ad - bc = 1$$

Supergravedad cuántica perturbativa en campos cuánticos relativistas. Modelo Matemático.

Espacio de Moduli, amplitudes y gauge fixing. Invariante BRST y tensores. Invariancia difeomorfa. Invariancia de gauge. Operadores de Vértice.

$$I = I_X + I_{\text{gh}}$$

$$I_{\text{gh}} = \frac{1}{2\pi} \int d^2\sigma \sqrt{g} b_{ij} D^i c^j$$

$$\{Q_B, b_{ij}\} = T_{ij}$$

$$\delta I = \frac{1}{4\pi} \int d^2\sigma \sqrt{g} \delta g_{ij} T^{ij}$$

$$[Q_B, g_{ij}] = \delta g_{ij}, \{Q_B, \delta g_{ij}\} = 0$$

$$I \rightarrow \hat{I} = I + \frac{1}{4\pi} \int d^2\sigma \sqrt{g} \delta g_{ij} b^{ij}$$



$$F(g \mid \delta g) = \int \mathcal{D}(X, b, c) \exp(-\hat{I}(X, b, c; g, \delta g)) = \int \mathcal{D}(X, b, c) \exp(-I) \exp\left(-\frac{1}{4\pi} \int_{\Sigma} d^2\sigma \sqrt{g} \delta g_{ij} b^{ij}\right)$$

$$[Q_B, F(g \mid \delta g)] = 0$$

$$\exp\left(-\frac{1}{4\pi} \int_{\Sigma} d^2\sigma \sqrt{g} \delta g_{ij} b^{ij}\right) = \sum_{n=0}^{\infty} \frac{1}{n!} \left(-\frac{1}{4\pi} \int_{\Sigma} d^2\sigma \sqrt{g} \delta g_{ij} b^{ij}\right)^n$$

$$Q_B = \int_{\Sigma} d^2\sigma \sum_{i,j=1,2} \sqrt{g} \delta g_{ij} \frac{\delta}{\delta g_{ij}}$$

$$d = \sum_{i=1}^s dx^i \frac{\partial}{\partial x^i},$$

$$F(x^1 \dots | \dots dx^s) = \sum_{i_1 < \dots < i_p} F_{i_1 \dots i_p}(x^1 \dots x^s) dx^{i_1} \dots dx^{i_p}$$

$$g_{ij} \rightarrow g_{ij} + \epsilon (D_i v_j + D_j v_i)$$

$$\delta g_{ij} \rightarrow \delta g_{ij} + \epsilon (D_i v_j + D_j v_i)$$

$$\int_{\Sigma} d^2\sigma \sqrt{g} (D_i v_j + D_j v_i) \frac{\delta}{\delta (\delta g_{ij})} F(g \mid \delta g) = 0$$

$$\int \mathcal{D}(X, b, c) \exp(-\hat{I}) \int_{\Sigma} d^2\sigma \sqrt{g} (D_i v_j + D_j v_i) b^{ij} = 0$$

$$D_i b^{ij} = 0.$$

$$c^i \rightarrow c^i + \epsilon v^i,$$

$$\int \mathcal{D}(X, b, c) \exp(-\hat{I}) \rightarrow \int \mathcal{D}(X, b, c) \exp(-\hat{I}) \left(1 + \frac{\epsilon}{2\pi} \int_{\Sigma} d^2\sigma \sqrt{g} b_{ij} D^i v^j\right)$$

$$b_{ij} = \sum_{\alpha=1}^{6g-6} u_{\alpha} b_{\alpha,ij} + \sum_{\lambda} w_{\lambda} b'_{\lambda ij}.$$

$$\hat{I} = \dots + \frac{1}{4\pi} \sum_{\alpha=1}^{6g-6} \int_{\Sigma} d^2\sigma \sqrt{g} \delta g_{ij} u_{\alpha} b_{\alpha}^{ij}.$$

$$\prod_{\alpha=1}^{6g-6} \int du_{\alpha} \exp\left(\frac{u_{\alpha}}{4\pi} \int_{\Sigma} d^2\sigma \sqrt{g} \delta g_{ij} b_{\alpha}^{ij}\right) = \prod_{\alpha=1}^{6g-6} \frac{1}{4\pi} \int_{\Sigma} d^2\sigma \sqrt{g} \delta g_{ij} b_{\alpha}^{ij}.$$

$$v_{\alpha} = \frac{1}{4\pi} \int_{\Sigma} d^2\sigma \sqrt{g} \delta g_{ij} b_{\alpha}^{ij}$$



$$\prod_{\alpha=1}^{6g-6} \delta(v_{\alpha})$$

$$c^i=\sum_{\lambda}\gamma_{\lambda}c_{\lambda}^i$$

$$\prod_{\lambda}\int D\gamma_{\lambda}\exp\left(m_{\lambda}\gamma_{\lambda}w_{\lambda}\right)=\prod_{\lambda}m_{\lambda}\cdot\prod_{\lambda}\delta(w_{\lambda})$$

$$Z_{\mathbf{g}} = \int_{\mathcal{M}_{\mathbf{g}}} F(g \mid \delta g)$$

$$\delta g_{ij} = \sum_{s=1}^p \frac{\partial g_{ij}}{\partial m_s} \, \mathrm{d} m_s.$$

$$\frac{1}{4\pi}\int_{\Sigma}\mathrm{d}^2\sigma\sqrt{g}\delta g_{ij}b^{ij}=\frac{1}{4\pi}\sum_{s=1}^p\mathrm{d}m_s\int_{\Sigma}\mathrm{d}^2\sigma\frac{\partial(\sqrt{g}g_{ij})}{\partial m_s}b^{ij}$$

$$F(g\mid\delta g)=\int\mathcal{D}(X,b,c)\exp\left(-\hat{I}\right)=\int\mathcal{D}(X,b,c)\exp\left(-I-\frac{1}{4\pi}\sum_{s=1}^p\mathrm{d}m_s\int_{\Sigma}\mathrm{d}^2\sigma\frac{\partial(\sqrt{g}g_{ij})}{\partial m_s}b^{ij}\right)$$

$$\mathrm{d}m_1,\ldots,\mathrm{d}m_{\mathfrak{p}}$$

$$F_{\mathrm{top}}\left(m_1\ldots\mid\cdots\mathrm{d}m_{\mathfrak{p}}\right)=(-1)^{\mathfrak{p}(\mathfrak{p}+1)/2}\,\mathrm{d}m_1\ldots\,\mathrm{d}m_{\mathfrak{p}}\int\mathcal{D}(X,b,c)e^{-I}\prod_{s=1}^p\frac{1}{4\pi}\int_{\Sigma}\mathrm{d}^2\sigma\frac{\partial(\sqrt{g}g_{ij})}{\partial m_s}b^{ij}$$

$$\Psi_s=\frac{1}{4\pi}\int_{\Sigma}\mathrm{d}^2\sigma\frac{\partial(\sqrt{g}g_{ij})}{\partial m_s}b^{ij}.$$

$$(-1)^{\mathfrak{p}(\mathfrak{p}+1)/2}\int\mathcal{D}(m_1\ldots\mid\cdots\mathrm{d}m_{\mathfrak{p}})\mathrm{d}m_1\ldots\,\mathrm{d}m_{\mathfrak{p}}\int\mathcal{D}(X,b,c)\exp\left(-I\right)\Psi_1\ldots\Psi_{\mathfrak{p}}.$$

$$(-1)^{\mathfrak{p}(\mathfrak{p}+1)/2}\int\mathcal{D}(m_1\ldots\mid\cdots\mathrm{d}m_{\mathfrak{p}})\mathrm{d}m_1\ldots\,\mathrm{d}m_{\mathfrak{p}}\int\mathcal{D}(X,b,c)\exp\left(-I\right)\delta(\Psi_1)\ldots\delta(\Psi_{\mathfrak{p}}).$$

$$\int\mathcal{D}(m_1,\ldots,m_{\mathfrak{p}})\int\mathcal{D}(X,b,c)\exp\left(-I\right)\delta(\Psi_1)\ldots\delta(\Psi_{\mathfrak{p}})$$

$$F_{\Omega}(g\mid\delta g)=\int\mathcal{D}(X,b,c)\exp\left(-\hat{I}(X,b,c,g,\delta g)\right)\Omega.$$

$$\mathrm{d}F_{\Omega}+F_{Q_B\Omega}=0$$



$$\Omega=\prod_{s=1}^n\mathcal{V}_s(X,b,c;p_s)$$

$$b_n\mathcal{V}_s=\tilde{b}_n\mathcal{V}_s=0,n\geq 0$$

$$b_n=\frac{1}{2\pi i}\oint\,\,\,\mathrm{d} z z^{n+1}b_{zz},\,\tilde{b}_n=\frac{1}{2\pi i}\oint\,\,\,\mathrm{d} \tilde{z} \tilde{z}^{n+1}b_{\tilde{z}\tilde{z}}$$

$$b_n\mathcal{W}_s=\tilde{b}_n\mathcal{W}_s=0,n\geq 0.$$

$$\delta c^i=c^j\partial_jc^i$$

$$\delta X=c^j\partial_jX$$

$$\delta g_{ij}=D_ic_j+D_jc_i-g_{ij}D_kc^k$$

$$v^i(p_1)=\cdots=v^i(p_{\mathfrak{n}})=0$$

$$\mathcal{M}_{\mathbf{g},\mathbf{n}}=\mathcal{I}/\mathcal{D}_{p_1,\ldots,p_{\mathbf{n}}}$$

$$\int \mathcal{D}(X,b,c)\exp\left(-\hat{I}\right)\prod_{s=1}^n\mathcal{V}_s(X,b,c;p_s)\int_{\Sigma}\mathrm{d}^2\sigma\sqrt{g}(D_iv_j+D_jv_i)b^{ij}=0$$

$$\langle \mathcal{V}_1\ldots \mathcal{V}_{\mathfrak{n}}\rangle_{\mathbf{g}}=\int_{\mathcal{M}_{\mathbf{g},\mathbf{n}}}F_{\mathcal{V}_1,\ldots,\mathcal{V}_{\mathfrak{n}}}$$

$$\int_{\mathcal{M}_{\mathbf{g},\mathbf{n}}}F_{\{Q_B,\mathcal{W}_1\},\mathcal{V}_2,\ldots,\mathcal{V}_{\mathfrak{n}}}=0$$

$$\mathrm{d} F_{\mathcal{W}_1,\mathcal{V}_2,\ldots,\mathcal{V}_{\mathfrak{n}}}+F_{\{Q_B,\mathcal{W}_1\},\mathcal{V}_2,\ldots,\mathcal{V}_{\mathfrak{n}}}=0$$

$$\int_{\mathcal{M}_{\mathbf{g},\mathbf{n}}}F_{\{Q_B,\mathcal{W}_1\},\mathcal{V}_2,\ldots,\mathcal{V}_{\mathfrak{n}}}=-\int_{\mathcal{M}_{\mathbf{g},\mathbf{n}}}\mathrm{d} F_{\mathcal{W}_1,\mathcal{V}_2,\ldots,\mathcal{V}_{\mathfrak{n}}}$$

- Deligne-Mumford – Riemann.

$$z=\mathbf{z},\tilde{z}=\tilde{\mathbf{z}}$$

$$z=\begin{cases}w+\mathbf{z} & \text{if }|w|<\epsilon\\w & \text{if }|w|>1-\epsilon\end{cases}$$

$$z=w+f(|w|)\mathbf{z}$$

$$\tilde{z}=\tilde{w}+f(|w|)\tilde{\mathbf{z}}$$

$$\partial_{\mathbf{z}}g_{IJ}=D_Iv_J+D_Jv_I$$

$$\partial_{\tilde{\mathbf{z}}}g_{IJ}=D_I\tilde{v}_J+D_J\tilde{v}_I,$$

$$(v^w,v^{\tilde{w}})=-\big(\partial_{\mathbf{z}}|_{z,\tilde{z}}w,\partial_{\mathbf{z}}|_{z,\tilde{z}}\tilde{w}\big)$$



$$(v^w, v^{\tilde{w}}) = (1, 0)$$

$$(\tilde{v}^w, \tilde{v}^{\tilde{w}}) = (0, 1)$$

$$\mathrm{d}\mathbf{z}\,\mathrm{d}\tilde{\mathbf{z}}\Psi_{\mathbf{z}}\Psi_{\tilde{\mathbf{z}}}$$

$$\Psi_{\mathbf{z}}=\frac{1}{2\pi}\int_{\Sigma}\mathrm{d}w\,\mathrm{d}\tilde{w}\sqrt{g}b^{ij}D_iv_j,\,\Psi_{\tilde{\mathbf{z}}}=\frac{1}{2\pi}\int_{\Sigma}\mathrm{d}w\,\mathrm{d}\tilde{w}\sqrt{g}b^{ij}D_i\tilde{v}_j$$

$$\Psi_{\mathbf{z}}\Psi_{\tilde{\mathbf{z}}}\tilde{c}V(0)$$

$$\int_{\Sigma} \mathrm{d}\mathbf{z}\,\mathrm{d}\tilde{\mathbf{z}} V(\mathbf{z},\tilde{\mathbf{z}})$$

$$\begin{array}{c} \Sigma \rightarrow \mathcal{M}_{\mathbf{g},\mathbf{n}} \\ \downarrow \pi \\ \mathcal{M}_{\mathbf{g},\mathbf{n}-1} \end{array}$$

$$\mathcal{V}=\delta(\tilde{c})\delta(c)V$$

$$\delta I = \frac{-i}{4\pi} \int_{\Sigma} \mathrm{d}\tilde{z} \, \mathrm{d}z \big(\delta J_{\tilde{z}}^z T_{zz} - \delta J_z^{\tilde{z}} T_{\tilde{z}\tilde{z}} \big)$$

$$\hat{I} = I + \frac{-i}{4\pi} \int_{\Sigma} \mathrm{d}\tilde{z} \, \mathrm{d}z \big(\delta J_{\tilde{z}}^z b_{zz} - \delta J_z^{\tilde{z}} b_{\tilde{z}\tilde{z}} \big)$$

Espacios de Supermoduli - Riemann. Variables holomórficas y antiholomórficas y deformaciones.

$$D_\theta=\frac{\partial}{\partial\theta}+\theta\frac{\partial}{\partial z}.$$

$$\Phi^{[n]}=u+\theta v,$$

$$v_f=f(z)(\partial_\theta-\theta\partial_z)$$

$$V_g=g(z)\partial_z+\frac{g'(z)}{2}\theta\partial_\theta$$

$$G_r=z^{r+1/2}(\partial_\theta-\theta\partial_z),\qquad r\in\mathbb{Z}+1/2$$

$$L_n=-z^{n+1}\partial_z-\frac{1}{2}(n+1)z^n\theta\partial_\theta,\quad n\in\mathbb{Z}$$

$$r\geq-1/2, n\geq-1$$

$$[L_m,L_n]=(m-n)L_{m+n}$$

$$\{G_r,G_s\}=2L_{r+s}$$

$$[L_m,G_r]=\left(\frac{m}{2}-r\right)G_{m+r}$$

$$\delta J_{\tilde{z}}^z=h_{\tilde{z}}^z+\theta\chi_{\tilde{z}}^\theta$$

$$\delta J_{\theta}^{\tilde{z}}=e_{\theta}^{\tilde{z}}+\theta h_z^{\tilde{z}}$$



$$q^{\bar{z}}\partial_{\bar{z}} + \left(q^z\partial_z + \frac{1}{2}D_\theta q^z D_\theta\right) + q^\theta D_\theta$$

$$\delta D_\theta = [D_\theta, q^\theta D_\theta] = (D_\theta q^\theta)D_\theta - 2q^\theta\partial_z$$

$$\delta J_{\bar{z}}^z = \partial_{\bar{z}}q^z$$

$$\delta J_\theta^{\bar{z}} = D_\theta q^{\bar{z}} \quad (3.10)$$

Propagadores, identidades y números fantasma.

$$\delta I = \frac{1}{2\pi} \int_{\Sigma} \mathcal{D}(\bar{z}, z | \theta) (\delta J_{\bar{z}}^z \mathcal{S}_{z\theta} + \delta J_\theta^{\bar{z}} T_{\bar{z}\bar{z}})$$

$$\mathcal{D}(\bar{z}, z | \theta) = -i[d\bar{z}; dz | d\theta]$$

$$\mathcal{S}_{z\theta} = S_{z\theta} + \theta T_{zz}$$

$$\partial_{\bar{z}} \mathcal{S}_{z\theta} = 0 = D_\theta T_{\bar{z}\bar{z}} = \partial_z T_{\bar{z}\bar{z}},$$

$$\mathcal{S}_{z\theta}(z | \theta) = \frac{1}{2} \sum_{r \in \mathbb{Z} + 1/2} z^{-r-3/2} G_r + \theta \sum_{n \in \mathbb{Z}} z^{-n-2} L_n$$

$$T_{\bar{z}\bar{z}} = \sum_{n \in \mathbb{Z}} \bar{z}^{-n-2} \tilde{L}_n.$$

$$C^z = c^z + \theta \gamma^\theta$$

$$B_{z\theta} = \beta_{z\theta} + \theta b_{zz}$$

$$\tilde{B}_{\bar{z}\bar{z}} = \tilde{b}_{\bar{z}\bar{z}} + \theta \tilde{f}_{\bar{z}\bar{z}\theta}, \tilde{C}^{\bar{z}} = \tilde{c}^{\bar{z}} + \theta \tilde{g}_\theta^{\bar{z}}$$

$$I_{\text{gh}} = \frac{1}{2\pi} \int \mathcal{D}(\bar{z}, z | \theta) (B_{z\theta} \partial_{\bar{z}} C^z + \tilde{B}_{\bar{z}\bar{z}} D_\theta \tilde{C}^{\bar{z}})$$

$$\partial_{\bar{z}} B_{z\theta} = 0 = D_\theta \tilde{B}_{\bar{z}\bar{z}}$$

$$[Q_B, B_{z\theta}] = \mathcal{S}_{z\theta}$$

$$\{Q_B, \tilde{B}_{\bar{z}\bar{z}}\} = T_{\bar{z}\bar{z}}$$

$$\{Q_B, \delta J_{\bar{z}}^z\} = [Q_B, \delta J_\theta^{\bar{z}}] = 0$$

$$I \rightarrow \hat{I} = I + \frac{1}{2\pi} \int_{\Sigma} \mathcal{D}(\bar{z}, z | \theta) (\delta J_{\bar{z}}^z B_{z\theta} - \delta J_\theta^{\bar{z}} \tilde{B}_{\bar{z}\bar{z}})$$

$$F(\mathcal{J}, \delta \mathcal{J}) = \int \mathcal{D}(X, B, C, \tilde{B}, \tilde{C}) \exp(-\hat{I})$$

$$[Q_B, F(\mathcal{J}, \delta \mathcal{J})] = 0$$

$$dF(\mathcal{J}, \delta \mathcal{J}) = 0$$



Supermanifolds e integración bosónica.

$$F(x, \lambda dx) = \lambda^S F(x, dx)$$

$$d = \sum_I dx^I \frac{\partial}{\partial x^I}, \quad d^2 = 0$$

$$\int \mathcal{D}(x, dx) F(x, dx)$$

$$F(x, dx) = f(t^1 \dots | \dots \theta^n) dt^1 \dots dt^m \delta(d\theta^1) \dots \delta(d\theta^n)$$

$$F(x, dx) = f(t^1 \dots | \dots \theta^n) \delta(dt^1) \dots \delta(d\theta^n) \\ = f(t^1 \dots | \dots \theta^n) \delta^{m|n}(dt^1 \dots | \dots d\theta^n)$$

$$\frac{\partial^r}{\partial (d\theta)^r} \delta(d\theta)$$

$$F(x, dx) = \theta^2 dt \delta'(d\theta^1)$$

$$I_{\beta\gamma} = \frac{1}{\pi} \int_{\Sigma} d^2z \beta \partial_{\bar{z}} \gamma$$

$$\int \mathcal{D}(d\theta) \frac{\partial^r}{\partial (d\theta)^r} \delta(d\theta) = \delta_{r,0}$$

$$\int \mathcal{D}(u,v) \exp \left(- \sum_{i,j} u_i m_{ij} v_j \right) = \frac{1}{\det m}$$

$$\int \mathcal{D}(u,v) \exp \left(- \sum_{i,j} u_i m_{ij} v_j - \sum_k (r_k u_k + s_k v_k) \right) = \frac{\exp(\sum_{i,j} s_i (m^{-1})_{ij} r_j)}{\det m}$$

$$\int \mathcal{D}(u,v) \exp \left(- \sum_{i,j} u_i m_{ij} v_j - \sum_k s_k v_k \right) = \frac{1}{\det m}$$

$$G(v) = \int \mathcal{D}(u) \exp(-(u, mv))$$

$$\int \mathcal{D}(v) A(v) G(v) = \int \mathcal{D}(u, v) A(v) \exp(-(u, mv))$$

$$\int \mathcal{D}(v) \exp(-(s, v)) G(v) = \frac{1}{\det m}$$

$$\int \mathcal{D}(u) \exp(-(u, mv)) = \delta(mv)$$



$$\int \mathcal{D}(v) \delta(mv) = \frac{1}{\det m}$$

$$\tilde{I}_{\delta JB} = \frac{1}{2\pi} \int \mathcal{D}(\tilde{z}, z \mid \theta) \delta J_{\tilde{z}}^z B_{z\theta}$$

$$B=\sum_{\alpha=1\dots|\cdots 2g-2}u_{\alpha}B_{\alpha}+\sum_{\lambda}w_{\lambda}B'_{\lambda}.$$

$$C=\sum_{\lambda}\gamma_{\lambda}C_{\lambda},$$

$$\int \, \mathrm{d} \gamma \mathrm{exp}\left(-m w \gamma\right) = \delta(m w)$$

$$\tilde{I}_{\delta JB} = \sum_{\alpha=1\dots|\cdots 2g-2} u_{\alpha} \Psi_{\alpha},$$

$$\Psi_{\alpha} = \frac{1}{2\pi} \int_{\Sigma} \mathcal{D}(\tilde{z}, z \mid \theta) \delta J_{\tilde{z}}^z \, B_{z\theta\alpha}$$

$$\prod_{\alpha=1\dots|\cdots 2g-2}\int \, \mathrm{d} u_{\alpha} \mathrm{exp}\left(u_{\alpha}\Psi_{\alpha}\right)=\prod_{\alpha=1\dots|\cdots 2g-2}\delta(\Psi_{\alpha})=\delta^{3g-3|2g-2}(\Psi_1'\dots|\dots\Psi_{2g-2}'')$$

$$\tilde{I}_{\delta J\tilde{B}} = \frac{1}{2\pi} \int_{\Sigma} \mathcal{D}(\tilde{z}, z \mid \theta) \delta J_{\tilde{\theta}}^{\tilde{z}} \tilde{B}_{\tilde{z}\tilde{z}}$$

$$\tilde{I}_{\delta J\tilde{B}} = \sum_{\alpha=1}^{3g-3} \tilde{u}_{\alpha} \tilde{\Psi}_{\alpha}$$

$$\tilde{\Psi}_{\alpha} = \frac{1}{2\pi} \int_{\Sigma} [\mathrm{d}\tilde{z}; \mathrm{d}z \mid \mathrm{d}\theta] \delta J_{\tilde{\theta}}^{\tilde{z}} \, \tilde{B}_{\tilde{z}\tilde{z}}$$

$$\prod_{\alpha=1}^{3g-3}\int \, \mathrm{d} \tilde{u}_{\alpha} \mathrm{exp}\left(\tilde{u}_{\alpha}\tilde{\Psi}_{\alpha}\right)=\delta^{3g-3}(\tilde{\Psi}_1,\dots,\tilde{\Psi}_{3g-3})$$

$$\delta^{3g-3}(\tilde{\Psi}_1\ldots\tilde{\Psi}_{3g-3})\delta^{3g-3|2g-2}(\Psi_1'\ldots|\ldots\Psi_{2g-2}'')$$

$$\delta J_{\tilde{z}}^z \rightarrow \delta J_{\tilde{z}}^z + \partial_{\tilde{z}} q^z, \delta J_{\tilde{\theta}}^{\tilde{z}} \rightarrow \delta J_{\tilde{\theta}}^{\tilde{z}} + D_{\theta} q^{\tilde{z}}$$

$$m'=f(m\mid \eta^1,\eta^2)\\ \eta'^i=\psi^i(m\mid \eta^1,\eta^2), i=1,2$$

$$m'=f_0(m)+\eta^1\eta^2f_2(m)$$

$$\tilde{m}' = \tilde{f}(\tilde{m})$$

$$Z_{\mathrm{g}} = \int_{\Gamma} \mathcal{D}(\mathcal{J}, \delta \mathcal{J}) F(\mathcal{J}, \delta \mathcal{J})$$



$$\begin{aligned}
\delta J_{\bar{z}}^z &= \sum_{\alpha=1 \dots | \dots 2g-2} \frac{\partial J_{\bar{z}}^z}{\partial \mathbf{m}_{\alpha}} d\mathbf{m}_{\alpha} \\
\delta J_{\bar{\theta}}^{\bar{z}} &= \sum_{\beta=1 \dots 3g-3} \frac{\partial J_{\bar{\theta}}^{\bar{z}}}{\partial \tilde{m}_{\beta}} d\tilde{m}_{\beta} \\
I_{d\mathbf{m}, d\tilde{m}} &= \sum_{\alpha=1 \dots | \dots 2g-2} d\mathbf{m}_{\alpha} B^{(\alpha)} + \sum_{\beta=1 \dots 3g-3} d\tilde{m}_{\beta} \tilde{B}^{(\beta)} \\
B^{(\alpha)} &= \frac{1}{2\pi} \int_{\Sigma} \mathcal{D}(\tilde{z}, z \mid \theta) \frac{\partial J_{\bar{z}}^z}{\partial \mathbf{m}_{\alpha}} B_{z\theta} \\
\tilde{B}^{(\beta)} &= \frac{1}{2\pi} \int_{\Sigma} \mathcal{D}(\tilde{z}, z \mid \theta) \frac{\partial J_{\bar{\theta}}^{\bar{z}}}{\partial \tilde{m}_{\beta}} \tilde{B}_{\bar{z}\bar{z}} \\
\prod_{\beta=1 \dots 3g-3} \int \mathcal{D}(d\tilde{m}_{\beta}) \exp(-d\tilde{m}_{\beta} \tilde{B}^{(\beta)}) \cdot \prod_{\alpha=1 \dots | \dots 2g-2} \mathcal{D}(d\mathbf{m}_{\alpha}) \exp(-d\mathbf{m}_{\alpha} B^{(\alpha)}) \\
&= \prod_{\beta=1 \dots 3g-3} \delta(\tilde{B}^{(\beta)}) \cdot \prod_{\alpha=1 \dots | \dots 2g-2} \delta(B^{(\alpha)}) \\
&\quad \delta^{3g-3}(\tilde{B}^{(\beta)}) \delta^{3g-3|2g-2}(B^{(\alpha)}) \\
\Lambda(\tilde{m}, \mathbf{m}) &= \int \mathcal{D}(X, B, C, \tilde{B}, \tilde{C}) \exp(-I) \delta^{3g-3}(\tilde{B}^{(\beta)}) \delta^{3g-3|2g-2}(B^{(\alpha)}) \\
\mathcal{D}(x) &= [d\tilde{m}_1 \dots d\tilde{m}_{3g-3}; dm_1 \dots dm_{3g-3} \mid d\eta_1 \dots d\eta_{2g-2}] \\
\Xi(x) &= \mathcal{D}(x) \Lambda(x) \\
\Xi(x) &= \int \mathcal{D}(dx) F(J, \delta J) \\
Z_g &= \int_{\Gamma} \Xi(x)
\end{aligned}$$

Sistema de Coordenadas – Moduli y transformaciones de gauge.

$$\begin{aligned}
\Xi &= [d\tilde{m}; dm \mid d\eta^1, d\eta^2] (Y_0(\tilde{m}, m) + Y_2(\tilde{m}, m) \eta^1 \eta^2) \\
m' &= m + a(m) \eta^1 \eta^2 \\
\eta'^1 &= \eta^1 \\
\eta'^2 &= \eta^2 \\
\Xi &= [d\tilde{m}; dm' \mid d\eta'^1, d\eta'^2] (Y'_0(\tilde{m}, m') + Y'_2(\tilde{m}, m') \eta'^1 \eta'^2) \\
Y'_2(\tilde{m}, m') &= Y_2(\tilde{m}, m') - \partial_{m'} (a(m') Y_0(\tilde{m}; m')) \\
\delta J_{\bar{z}}^z &= h_{\bar{z}}^z + \theta \chi_{\bar{z}}^{\theta} \\
\chi_{\bar{z}}^{\theta} &\rightarrow \chi_{\bar{z}}^{\theta} + \partial_{\bar{z}} y^{\theta}
\end{aligned}$$



$$\chi_{\bar{z}}^{\theta} = \sum_{\sigma=1}^{2g-2} \eta_{\sigma} \chi_{\bar{z}}^{(\sigma)\theta}$$

$$\partial_{\bar{z}} y^{\theta} = \sum_{\sigma} e_{\sigma} \chi_{\bar{z}}^{(\sigma)\theta}$$

$$\chi_{\bar{z}}^{(\sigma)\theta} \rightarrow \chi_{\bar{z}}^{(\sigma)\theta} + \partial_{\bar{z}} y^{(\sigma)\theta}$$

$$\chi \rightarrow \chi + \partial_{\bar{z}} y,$$

$$y = \sum_{\sigma=1}^{2g-2} \eta_{\sigma} y^{(\sigma)}$$

$$h_{\bar{z}}^z \rightarrow h_{\bar{z}}^z + y^{\theta} \chi_{\bar{z}}^{\theta}.$$

$$h_{\bar{z}}^z \rightarrow h_{\bar{z}}^z + \sum_{\sigma, \sigma'} \eta_{\sigma} \eta_{\sigma'} y^{(\sigma)\theta} \chi_{\bar{z}}^{(\sigma')\theta}$$

$$I_{\eta} = \sum_{\sigma=1}^{2g-2} \frac{\eta_{\sigma}}{2\pi} \int_{\Sigma_{\text{red}}} d^2 z \chi_{\bar{z}}^{(\sigma)\theta} S_{z\theta}$$

$$I_{d\eta} = \sum_{\sigma=1}^{2g-2} \frac{d\eta_{\sigma}}{2\pi} \int_{\Sigma_{\text{red}}} d^2 z \chi_{\bar{z}}^{(\sigma)\theta} \beta_{z\theta}$$

$$\exp \left(-\frac{1}{2\pi} \int_{\Sigma_{\text{red}}} d^2 z \chi_{\bar{z}}^{(\sigma)\theta} (\eta_{\sigma} S_{z\theta} + d\eta_{\sigma} \beta_{z\theta}) \right)$$

$$\delta \left(\int_{\Sigma_{\text{red}}} d^2 z \chi_{\bar{z}}^{(\sigma)\theta} \beta_{z\theta} \right) \cdot \int_{\Sigma_{\text{red}}} d^2 z \chi_{\bar{z}}^{(\sigma)\theta} S_{z\theta}$$

$$\delta \left(\int_{\Sigma_{\text{red}}} d^2 z \chi_{\bar{z}}^{(\sigma)\theta} \beta_{z\theta} \right) \cdot \delta \left(\int_{\Sigma_{\text{red}}} d^2 z \chi_{\bar{z}}^{(\sigma)\theta} S_{z\theta} \right)$$

$$\chi_{\bar{z}}^{(\sigma)\theta} \rightarrow \lambda \chi_{\bar{z}}^{(\sigma)\theta}, \eta_{\sigma} \rightarrow \lambda^{-1} \eta_{\sigma}, \lambda \in \mathbb{C}^*$$

$$\begin{pmatrix} \chi_{\bar{z}}^{(1)} \\ \vdots \\ \chi_{\bar{z}}^{(s)\theta} \end{pmatrix} \rightarrow M \begin{pmatrix} \chi_{\bar{z}}^{(1)} \\ \vdots \\ \chi_{\bar{z}}^{(s)\theta} \end{pmatrix}$$

$$(\eta_1 \dots \eta_s) \rightarrow (\eta_1 \dots \eta_s) M^{-1}, (d\eta_1 \dots d\eta_s) \rightarrow (d\eta_1 \dots d\eta_s) M^{-1}$$

$$\prod_{\sigma=1}^s \left(\delta \left(\int_{\Sigma_{\text{red}}} d^2 z \chi_{\bar{z}}^{(\sigma)\theta} \beta_{z\theta} \right) \cdot \int_{\Sigma_{\text{red}}} d^2 z \chi_{\bar{z}}^{(\sigma)\theta} S_{z\theta} \right)$$



Operador Picture-Changing.

$$\chi_{\tilde{z}}^{(\sigma)\theta} = \delta_{p_{\sigma}}.$$

$$y(p)=\delta(\beta(p))S_{z\theta}(p)$$

$$\chi_{\tilde{z}}^{(\sigma)\theta} = \delta_{p_{\sigma}}.$$

$$\prod_{\sigma=1}^{2g-2}y(p_{\sigma})$$

$$\partial_{\tilde{z}}y^{\theta}=\sum_{\sigma=1}^{2g-2}e_{\sigma}\delta_{p_{\sigma}}$$

$$H^0\left(\Sigma_{\mathrm{red}},T^{1/2}(\Sigma_{\sigma}p_{\sigma})\right)\neq 0$$

$$\chi_{\tilde{z}}^{(1)\theta}=\delta^2(z),\chi_{\tilde{z}}^{(2)\theta}=\delta^2(z-\epsilon)$$

$$\hat{\chi}_{\tilde{z}}^{(2)\theta}=-\frac{1}{\epsilon}\Big(\chi_{\tilde{z}}^{(2)\theta}-\chi_{\tilde{z}}^{(1)\theta}\Big)=-\frac{1}{\epsilon}(\delta^2(z-\epsilon)-\delta^2(z)).$$

$$\hat{\chi}_{\tilde{z}}^{(2)\theta}=\partial_z\delta^2(z).$$

$$\delta\big(\partial_z\beta(z_0)\big)\partial_zS_{z\theta}(z_0).$$

$$\delta\left(\int_{\Sigma_{\mathrm{red}}}\mathrm{d}^2z\chi_{\tilde{z}}^{(\sigma)\theta}\beta_{z\theta}\right)$$

$$\partial_z\delta(\beta(z))=\partial_z\beta(z)\cdot\delta'(\beta(z))$$

$$\delta^{(k)}\left(\int_{\Sigma_{\mathrm{red}}}\mathrm{d}^2z\chi_{\tilde{z}}^{(\sigma)\theta}\beta_{z\theta}\right)$$

$$r-s=2g-2$$

$$\delta^{3g-3|2g-2}(B^{(\alpha)})=\prod_{\alpha'=1}^{3g-3}\delta(b^{(\alpha')})\prod_{\alpha''=1}^{2g-2}\delta(\beta^{(\alpha'')})$$

Métrica Neveu-Schwarz.

$$b_n\mathcal{V}=\tilde{b}_n\mathcal{V}=\beta_r\mathcal{V}=0, n,r\geq 0$$

$$\beta_r=\frac{1}{2\pi i}\oint\,\,\,\mathrm{d}zz^{r+1/2}\beta_{z\theta}$$

$$\begin{array}{l} L_n\mathcal{V}=0,\quad n\geq 0\\ G_r\mathcal{V}=0,\quad r\geq 1/2, \end{array}$$



$$\tilde{L}_n\mathcal{V}=0, n\geq 0.$$

$$\delta(\tilde{b})\delta(b)\delta(\beta)$$

$$\mathcal{V}=\tilde{c}c\delta(\gamma)V,$$

$$L_n^XV=\frac{1}{2}\delta_{n,0}V, n\geq 0$$

$$G_r^XV=0, r>0$$

$$\tilde{L}_n^XV=\delta_{n,0}V, n\geq 0$$

$$Q_B^*=\sum_{r\in\mathbb{Z}+1/2}\gamma_rG_{-r}^X$$

$$\gamma(z)=\sum_{r\in\mathbb{Z}+1/2}z^{-r+1/2}\gamma_r$$

$$F_{\mathcal{V}_1,\ldots,\mathcal{V}_n}(\mathcal{J},\delta\mathcal{J})=\int\mathcal{D}(X,B,C,\tilde{B},\tilde{C})\exp{(-\hat{I})}\prod_{i=1}^n\mathcal{V}_i(p_i)$$

$$F_{\{Q_B,\mathcal{W}_1\},\mathcal{V}_2,\ldots,\mathcal{V}_n}+\mathrm{d}F_{\mathcal{W}_1,\mathcal{V}_2,\ldots,\mathcal{V}_n}=0$$

$$\tilde{b}_n\mathcal{V}=b_n\mathcal{V}=\beta_r\mathcal{V}=0, n,r\geq 0.$$

$$\tilde{b}_n\mathcal{W}=b_n\mathcal{W}=\beta_r\mathcal{W}=0, n,r\geq 0.$$

$$\hat{I}\rightarrow \hat{I}+\frac{1}{2\pi}\int_{\Sigma}\mathcal{D}(\check{z},z\mid \theta)\partial_{\check{z}}q^zB_{z\theta}$$

$$\langle \mathcal{V}_1\ldots \mathcal{V}_n\rangle =\int_{\Gamma}F_{\mathcal{V}_1,\ldots,\mathcal{V}_n}(\mathcal{J},\delta\mathcal{J})$$

$$\int_{\Gamma}F_{\{Q_B,\mathcal{W}_1\},\mathcal{V}_2,\ldots,\mathcal{V}_n}=-\int_{\Gamma}\mathrm{d}F_{\mathcal{W}_1,\mathcal{V}_2,\ldots,\mathcal{V}_n}=0$$

$$K_{\mathcal{V}}=\int_{\Sigma}\mathcal{D}(\check{z},z\mid \theta)V(\check{z};z\mid \theta)$$

$$\left(\frac{1}{2\pi}\int_{\Sigma}\mathrm{d}^2zg^z\partial_{\check{z}}b_{zz}\right)\cdot c^z(0)=g^z(0)$$

$$\delta\left(\frac{1}{2\pi}\int_{\Sigma}\mathrm{d}^2zg^z\partial_{\check{z}}b_{zz}\right)\cdot\delta(c^z(0))=g^z(0)$$

$$\delta\left(\frac{1}{2\pi}\int_{\Sigma}\mathrm{d}^2zf^{\theta}\partial_{\check{z}}\beta_{z\theta}\right)\cdot\delta(\gamma^{\theta}(0))=\frac{1}{f^{\theta}(0)}$$

$$\delta(\tilde{B}^{(\tilde{\mathbf{z}})})\delta(B^{(\mathbf{z})})\delta(B^{(\boldsymbol{\theta})})$$

$$[\mathrm{d}\tilde{\mathbf{z}};\mathrm{d}\mathbf{z}\mid \mathrm{d}\boldsymbol{\theta}]\delta(\tilde{B}^{(\tilde{\mathbf{z}})})\delta(B^{(\mathbf{z})})\delta(B^{(\boldsymbol{\theta})})\tilde{c}c\delta(\gamma)V$$



$$[\mathrm{d}\tilde{\mathbf{z}}; \mathrm{d}\mathbf{z} \mid \mathrm{d}\boldsymbol{\theta}]V(\tilde{\mathbf{z}}; \mathbf{z} \mid \boldsymbol{\theta}).$$

$$\int_{\Sigma} [\mathrm{d}\tilde{\mathbf{z}}; \mathrm{d}\mathbf{z} \mid \mathrm{d}\boldsymbol{\theta}]V = \left\{ \int_{\Sigma} [\mathrm{d}\tilde{\mathbf{z}}; \mathrm{d}\mathbf{z} \mid \mathrm{d}\boldsymbol{\theta}]D_{\boldsymbol{\theta}}W \right. \\ \left. \int_{\Sigma} [\mathrm{d}\tilde{\mathbf{z}}; \mathrm{d}\mathbf{z} \mid \mathrm{d}\boldsymbol{\theta}]\partial_{\tilde{\mathbf{z}}}W' \right.$$

$$\left\langle \prod_{i=1}^3 \tilde{c} V_i(\tilde{z}_i,z_i) \prod_{j=4}^n \int \mathrm{d}^2 z_j V(\tilde{z}_j,z_j) \right\rangle$$

$$\int \mathrm{d}\theta_3 \tilde{c} c V_3(\tilde{z}_3; z_3 \mid \theta_3) = \tilde{c} c D_{\theta} V_3(\tilde{z}_3; z_3 \mid 0)$$

$$\langle \tilde{c} c \delta(\gamma) V_1(\tilde{z}_1,z_1 \mid 0) \tilde{c} c \delta(\gamma) V_2(\tilde{z}_2,z_2 \mid 0) \tilde{c} c D_{\theta} V_3(\tilde{z}_3,z_3 \mid 0) \rangle$$

$$\left\langle \tilde{c} c \delta(\gamma) V_1(\tilde{z}_1,z_1 \mid 0) \tilde{c} c \delta(\gamma) V_2(\tilde{z}_2,z_2 \mid 0) \tilde{c} c D_{\theta} V_3(\tilde{z}_3,z_3 \mid 0) \prod_{j=4}^n \int \left[\mathrm{d}\tilde{z}_j; \mathrm{d}z_j \mid \mathrm{d}\theta_j\right] V_j(\tilde{z}_j; z_j \mid \theta_j) \right\rangle$$

$$\chi_{\tilde{z}}^{\theta} = \sum_{\sigma=1}^{2g-2+n} \eta_{\sigma} \chi_{\tilde{z}}^{(\sigma)\theta}$$

$$H^0\left(\Sigma_{\mathrm{red}}, T^{1/2}\otimes \mathcal{O}\left(\sum_{\sigma=1}^{2\,\mathrm{g}-2+\mathrm{n}}p_{\sigma}-\sum_{i=1}^{\mathrm{n}}q_i\right)\right)\neq 0$$

$$\left\langle \prod_{i=1}^3 \tilde{c} c \mathcal{V}^{(s_i)}(\tilde{z}_i,z_i) \prod_{j=4}^n \int \mathrm{d}^2 z_j \hat{V}_j^{(s_j)}(\tilde{z}_j,z_j) \right\rangle$$

$$\sum_i s_i = -2$$

Métrica Ramond.

$$D_{\theta}^* = \partial_{\theta} + \theta z \partial_z.$$

$$\tau|\zeta=(\log z)|\theta$$

$$(D_{\theta}^*)^2 = z\partial_z,$$

$$v_f=f(z)(\partial_{\theta}-\theta z\partial_z)$$

$$V_g = z \left(g(z) \partial_z + \frac{g'(z)}{2} \theta \partial_{\theta} \right)$$



$$G_r = z^r (\partial_\theta - \theta z \partial_z)$$

$$L_n = -z^{n+1} \partial_z - \frac{n z^n}{2} \theta \partial_\theta$$

$$[L_m, L_n] = (m-n)L_{m+n}$$

$$\{G_r, G_s\} = 2L_{r+s}$$

$$[L_m, G_r] = \left(\frac{m}{2} - r\right) G_{m+r}$$

$$G_0|_{z=0} = \partial_\theta$$

$$\theta \rightarrow \theta + \alpha$$

$$G_r \mathcal{V} = L_n \mathcal{V} = 0, r, n \geq 0$$

$$\mathcal{V} = \tilde{c} c \Theta V,$$

$$L_n^X V = \frac{5}{8} \delta_{n,0} V, n \geq 0$$

$$G_r^X V = 0, r \geq 0$$

$$\beta(z) = \sum_r z^{-r-3/2} \beta_r, \gamma(z) = \sum_r z^{-r+1/2} \gamma_r$$

$$Q_B^* = \sum_{r \in \mathbb{Z}} \gamma_{-r} G_r^X$$

$$\gamma_r \Theta = 0, r > 0$$

$$\beta_r \Theta = 0, r \geq 0$$

$$\dim \mathfrak{M}_{\mathfrak{g}, \mathfrak{n}_{\mathrm{NS}}, \mathfrak{n}_{\mathrm{R}}} = 3 \, \mathfrak{g} - 3 + \mathfrak{n}_{\mathrm{NS}} + \mathfrak{n}_{\mathrm{R}} \left| 2 \, \mathfrak{g} - 2 + \mathfrak{n}_{\mathrm{NS}} + \frac{1}{2} \mathfrak{n}_{\mathrm{R}} \right.$$

$$\begin{array}{l} \gamma_r \Theta = 0, \quad r \geq 0 \\ \beta_r \Theta = 0, \quad r > 0. \end{array}$$

$$D_\theta^* = \partial_\theta + \theta w(z) \partial_z$$

$$w(z)=\prod_{i=1}^{\mathfrak{n}_{\mathrm{R}}}(z-z_i)$$

$$\nu_f = f(z)(\partial_\theta - \theta w(z) \partial_z)$$

$$V_g = w(z) \left(g(z) \partial_z + \frac{g'(z)}{2} \theta \partial_\theta \right)$$

$$D_\theta^* = \partial_\theta + \theta z (z-a) \partial_z$$

$$\nu_s = z^s (\partial_\theta - \theta z (z-a) \partial_z)$$

$$V_n = z(z-a) \left(z^n \partial_z + \frac{n z^{n-1}}{2} \theta \partial_\theta \right)$$



$$c\Theta(a)c\Theta(0)\sim c\partial c\delta(\gamma),a\rightarrow 0$$

$$\Theta(a)\Theta(0)\sim \delta(\gamma),a\rightarrow 0$$

$$I_X=\frac{1}{2\pi i}\int_{\Sigma}[\mathrm{d}\tilde{\tau};\mathrm{d}\tau\mid\mathrm{d}\zeta]G_{IJ}(X^K)\partial_{\tilde{\tau}}X^ID_{\tilde{\zeta}}X^J$$

$$\begin{array}{l} \tilde{\tau}=\log \tilde{z} \\ \tau=\log z \\ \zeta=\theta \end{array}$$

$$I_X=\frac{1}{2\pi i}\int_{\Sigma}[\mathrm{d}\tilde{z};\mathrm{d}z\mid\mathrm{d}\theta]\frac{1}{z}G_{IJ}(X^K)\partial_{\tilde{z}}X^ID_{\theta}^*X^J$$

$$I_X=\frac{1}{2\pi}\int_{\Sigma_{\rm red}}\mathrm{d}^2z\left(G_{IJ}\partial_{\bar{z}}x^I\partial_zx^J+\frac{1}{z}G_{IJ}\psi^I\frac{D}{D\bar{z}}\psi^J\right)$$

$$\tau \cong \tau + 2\pi i$$

$$I_\psi=\frac{1}{2\pi}\int_{\Sigma_{\rm red}}\mathrm{d}^2\tau G_{IJ}(x^K)\psi^I\frac{D}{D\tilde{\tau}}\psi^J$$

$$I_{BC}=\frac{1}{2\pi i}\int_{\Sigma}[\mathrm{d}\tilde{z};\mathrm{d}z\mid\mathrm{d}\theta]B\partial_{\tilde{z}}C$$

$$[dz\mid\mathrm{d}\theta]=[dz^*\mid\mathrm{d}\theta^*]\mathrm{Ber}^{-1}\begin{pmatrix}\partial_zz^*&\partial_z\theta^*\\ \partial_\theta z^*&\partial_\theta\theta^*\end{pmatrix}$$

$$C=C^*\cdot\mathrm{Ber}^{-2}\begin{pmatrix}\partial_zz^*&\partial_z\theta^*\\ \partial_\theta z^*&\partial_\theta\theta^*\end{pmatrix}$$

$$B=B^*\cdot\mathrm{Ber}^3\begin{pmatrix}\partial_zz^*&\partial_z\theta^*\\ \partial_\theta z^*&\partial_\theta\theta^*\end{pmatrix}$$

$$\mathcal{R}^2\cong T\otimes\mathcal{O}(-q_1-\cdots-q_{\mathfrak{n}_{\mathsf{R}}})$$

$$\mathcal{R}^2\cong T\otimes\mathcal{O}(-q)$$

$$\mathcal{R}^{-2}\cong K\otimes\mathcal{O}(q)$$

$$C(z\mid\theta)=\hat{c}(z)+\theta\hat{\gamma}(z)$$

$$B=\hat{\beta}+\theta\hat{b}$$

$$I_{BC}=\frac{1}{2\pi}\int_{\Sigma_{\rm red}}\mathrm{d}^2z\big(\hat{b}\partial_{\bar{z}}\hat{c}+\hat{\beta}\partial_{\bar{z}}\hat{\gamma}\big)$$

$$\gamma\sim z^{1/2},\beta\sim z^{-1/2},z\rightarrow 0$$

$$\gamma(z)=\sum_{r\in\mathbb{Z}}z^{-r+1/2}\gamma_r,\beta(z)=\sum_{r\in\mathbb{Z}}z^{-r-3/2}\beta_r$$

$$\begin{aligned}\gamma_r\Theta_{-1/2} &= 0, \quad r > 0 \\ \beta_r\Theta_{-1/2} &= 0, \quad r \geq 0\end{aligned}$$

$$\beta_{z\theta'}(z)\gamma^{\theta'}(w)\sim-\frac{1}{z-w}$$

$$\langle \beta_{z\theta'}(z)\gamma^{\theta'}(w)\rangle_{\Theta_{-1/2}}=-\frac{1}{z-w}\sqrt{\frac{w}{z}}.$$

$$J_{\beta\gamma}(w)=\lim_{z\rightarrow w}\Big(-\beta(z)\gamma(w)-\frac{1}{z-w}\Big),$$

$$\langle J_{\beta\gamma}(w)\rangle_{\Theta_{-1/2}}=-\frac{1}{2w},$$

$$T_{\beta\gamma} =: \partial_z \beta_{z\theta'} \cdot \gamma^{\theta'} : - \frac{3}{2} \partial_z (:\beta_{z\theta'} \gamma^{\theta'}:)$$

$$\langle T_{\beta\gamma}(w)\rangle_{\Theta_{-1/2}}=\frac{3}{8w^2},$$

$$\langle \beta_{z\theta'}(z)\gamma^{\theta'}(w)\rangle_{\Theta_{-t}}=-\frac{1}{z-w}\Big(\frac{w}{z}\Big)^t$$

$$\begin{aligned}\langle J_{\beta\gamma}(w)\rangle_{\Theta_{-t}} &= -\frac{t}{w} \\ \langle T_{\beta\gamma}(w)\rangle_{\Theta_{-t}} &= -\frac{t(t-2)}{2w^2}\end{aligned}$$

$$b_n\mathcal{V}_1=\beta_r\mathcal{V}_1=0, n,r\geq 0,$$

$$F_{\{Q_B,\mathcal{W}_1\},\mathcal{V}_2,\ldots,\mathcal{V}_s}+\mathrm{d}F_{\mathcal{W}_1,\mathcal{V}_2,\ldots,\mathcal{V}_s}=0$$

$$\int_{\Gamma} F_{\{Q_B,\mathcal{W}_1\},\mathcal{V}_2,\ldots,\mathcal{V}_s}+\int_{\Gamma} \mathrm{d}F_{\mathcal{W}_1,\mathcal{V}_2,\ldots,\mathcal{V}_s}=0$$

$$\begin{array}{c} \Sigma \rightarrow \mathfrak{M}_{\mathbf{g},n_{\mathrm{NS}},n_{\mathrm{R}}} \\ \downarrow \pi \\ \mathfrak{M}_{\mathbf{g},n_{\mathrm{NS}}-1,n_{\mathrm{R}}}. \end{array}$$

$$\mathcal{R}^2\cong T\otimes \mathcal{O}(-q_1-\cdots-q_{n_{\mathrm{R}}})$$

$$\deg \mathcal{R}=1-\mathrm{g}-\frac{1}{2}n_{\mathrm{R}}$$

$$\deg \hat{\mathcal{R}}=1-\mathrm{g}-n_{\mathrm{NS}}-\frac{1}{2}n_{\mathrm{R}}$$

$$\chi_z^\theta = \sum_{\sigma=1}^{\Delta} \eta_\sigma \delta_{r_\sigma}$$



$$H^0\left(\Sigma,\hat{\mathcal{R}}\otimes\mathcal{O}\left(\sum_{\sigma=1}^{\Delta}r_{\sigma}\right)\right)\neq 0$$

$$\langle \tilde{c} c \delta(\gamma) V_1(\tilde{z}_1; z_1 \mid 0) \tilde{c} c \Theta_{-1/2} V_2'(\tilde{z}_2; z_2) \tilde{c} c \Theta_{-1/2} V_3'(\tilde{z}_3; z_3) \rangle.$$

$$\left\langle \prod_{i=1}^3 \tilde{c} c \Theta_{-1/2} V_i'(\tilde{z}_i; z_i) \int \mathrm{d}^2 z_4 \Theta_{-1/2} V_4'(\tilde{z}_4; z_4) \right\rangle$$

Dualidad isotr3pica.

$$\omega(\mathcal{V},\mathcal{W})=\langle \mathcal{V}(0)\mathcal{W}(1)\rangle.$$

$$\omega(Q_B\mathcal{U},\mathcal{V})+(-1)^{|\mathcal{U}|}\omega(\mathcal{U},Q_B\mathcal{V})=0.$$

$$\langle c\partial cU(0)cV(1)\rangle\neq 0.$$

$$\omega\colon \mathcal{H}_n^*\cong \mathcal{H}_{3-n}.$$

$$\omega\colon \mathcal{H}_{n,k}^*\cong \mathcal{H}_{1-n;-2-k}$$

$$\langle c\partial c\delta(\gamma)U(0)c\delta(\gamma)V(1)\rangle\neq 0$$

$$\omega\colon \mathcal{H}_{n,-3/2}^*\cong \mathcal{H}_{1-n;-1/2}$$

$$D_\theta^*=\partial_\theta+z(z-1)\theta\partial_z$$

$$v=\partial_\theta-z(z-1)\theta\partial_z$$

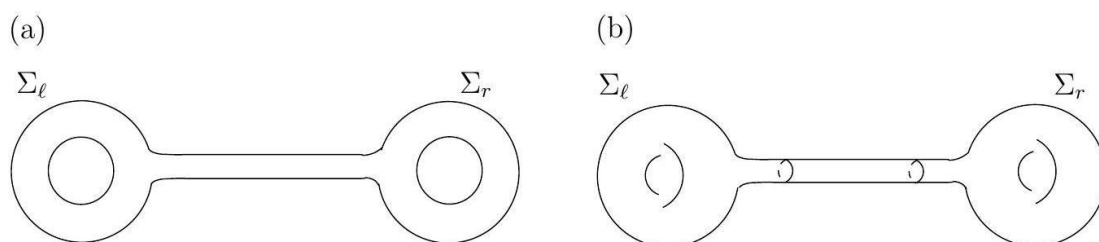
Supermasa y deformaciones dimensionales – GSO – Propagador de una Partícula Supermasiva – NS - Ramond en un espacio de Fock.

$$\left\langle c\partial c\Theta_{-1/2}u^\alpha\Sigma_\alpha(p)\cdot c\Theta_{-3/2}v_\beta\Sigma^\beta(q)\right\rangle=u^\alpha v_\alpha(2\pi)^{10}\delta^{10}(p+q).$$

$$p^I\Gamma_{I\alpha\beta}u^\beta=0$$

$$v_\beta\cong v_\beta+p^I\Gamma_{I\beta\gamma}w^\gamma$$

$$\frac{1}{p^2+m^2}=\int_0^\infty\mathrm{d}s\exp\left(-s(p^2+m^2)\right)$$



$$xy=q$$

$$x=e^{\varrho},y=qe^{-\varrho}$$

$$\varrho = u + i\varphi, u, \varphi \in \mathbb{R}$$

$$0\leq u\leq s, 0\leq \varphi\leq \pi$$

$$\mathrm{d}s^2=\mathrm{d}u^2+\mathrm{d}\varphi^2$$

$$L_0=-\frac{1}{2\pi}\int_0^\pi\mathrm{d}\varphi T_{uu}$$

$$b_0=-\frac{1}{2\pi}\int_0^\pi\mathrm{d}\varphi b_{uu}$$

$$z=-i\varrho=\varphi-iu$$

$$T_{zz}(u,\varphi)=T_{\bar{z}}(u,-\varphi)=T_{zz}(u,\varphi+2\pi)\\ b_{zz}(u,\varphi)=b_{\bar{z}\bar{z}}(u,-\varphi)=b_{zz}(u,\varphi+2\pi)$$

$$L_0=\frac{1}{2\pi}\int_0^{2\pi}\mathrm{d}\varphi T_{zz}$$

$$b_0=\frac{1}{2\pi}\int_0^{2\pi}\mathrm{d}\varphi b_{zz}$$

$$0\leq t\leq 1, 0\leq \varphi\leq \pi$$

$$\mathrm{d}s^2=s^2\,\mathrm{d}t^2+\mathrm{d}\varphi^2$$

$$\begin{aligned}\Psi_s&=\frac{1}{4\pi}\int_S\mathrm{d}t\,\mathrm{d}\varphi\frac{\partial(\sqrt{g}g_{ij})}{\mathrm{d}s}b^{ij}=\frac{1}{2\pi}\int_0^1\mathrm{d}t\int_0^\pi\mathrm{d}\varphi s^2b^{tt}\\&=\frac{1}{2\pi}\int_0^1\mathrm{d}t\int_0^\pi\mathrm{d}\varphi b_{uu}=\int_0^1\mathrm{d}tb_0=b_0\end{aligned}$$

$$\int_0^\infty \mathrm{d}s b_0 \exp\left(-sL_0\right)=\frac{b_0}{L_0}$$

$$b_0\int_0^1\frac{\mathrm{d}q}{q}q^{L_0}$$

$$L_0=\frac{\alpha'}{4}p^2+N,$$

$$\frac{4}{\alpha'}\frac{b_0}{p^2+M^2}$$

$$Z_{\Sigma;s}=\langle\psi_\ell|\frac{b_0}{L_0}|\psi_r\rangle.$$

$$\tilde{\theta} = \begin{cases} \theta & \text{if } \varphi = 0 \\ -\theta & \text{if } \varphi = \pi \end{cases}$$

$$\varphi \rightarrow \varphi + 2\pi, \theta \rightarrow -\theta$$

$$\chi_z^{\tilde{\theta}} = \begin{cases} \chi_z^{\theta} & \text{if } \varphi = 0 \\ -\chi_z^{\theta} & \text{if } \varphi = \pi \end{cases}$$

$$\chi_z^{\theta}(u,\varphi+2\pi)=-\chi_z^{\theta}(u,\varphi)$$

$$\chi_z^{\theta} \rightarrow \chi_z^{\theta} + \partial_z y^{\theta},$$

$$\frac{b_0}{L_0}\Pi_{\rm GSO}.$$

$$\int_0^\infty \mathrm{d} s \mathrm{Tr} \mathcal{O} b_0 \exp \left(-s L_0 \right) \Pi_{\rm GSO}$$

$$\begin{array}{l} \beta_r|q\rangle=0, \quad r>-q-3/2 \\ \gamma_r|q\rangle=0, \quad r\geq q+3/2 \end{array}$$

$$\beta_r|q\rangle=\gamma_r|q\rangle=0, r>0$$

$$\det \tilde{\partial}_{\beta \gamma} = \mathrm{Tr} (-1)^F \exp \left(-s L_{0; \beta^* \gamma^*} \right) = q^f \prod_{r=\frac{1}{2}, \frac{3}{2}, \dots}^{\frac{1}{2}, \frac{3}{2}, \dots} (1 - q^r)$$

$$\frac{1}{\det \tilde{\partial}_{\beta \gamma}} = q^{-f} \prod_{r=\frac{1}{2}, \frac{3}{2}, \dots}^{\frac{1}{2}, \frac{3}{2}, \dots} \frac{1}{1 - q^r}.$$

$$\mathrm{Tr} \mathcal{O} q^{L_0}$$

$$\tilde{\theta} = \begin{cases} \theta & \text{if } \varphi = 0 \\ \theta & \text{if } \varphi = \pi \end{cases}$$

$$G_0=\frac{1}{2\pi}\int_0^{2\pi}\mathrm{d}\varphi S_{z\theta}(0,\varphi)$$

$$\beta_0=\frac{1}{2\pi}\int_0^{2\pi}\mathrm{d}\varphi\beta_{z\theta}(u,\varphi)$$

$$\chi_z^{\theta}(u,\varphi+2\pi)=\chi_z^{\theta}(u,\varphi).$$

$$\chi_z^{\theta}=\eta,$$

$$\chi_z^{\theta} \rightarrow \chi_z^{\theta} + \partial_z y^{\theta}$$

- Deligne-Mumford:



$$\exp\left(-\frac{1}{2\pi}\int_0^{2\pi}d\varphi\int_0^sdu(\eta S_{z\theta}+d\eta\beta_{z\theta})\right)=\exp\left(-s(\eta G_0+d\eta\beta_0)\right)$$

$$G_0\delta(\beta_0)$$

$$\frac{b_0\delta(\beta_0)\Pi_{\rm GSO}G_0}{L_0}$$

$$\frac{b_0\delta(\beta_0)\Pi_{\rm GSO}}{G_0}$$

$$\mathrm{d}s^2=\mathrm{d}u^2+\mathrm{d}\varphi^2, 0\leq u\leq s, 0\leq \varphi\leq 2\pi.$$

$$L_0=-\frac{1}{2\pi}\int_0^{2\pi}\mathrm{d}\varphi T_{zz}$$

$$\tilde{L}_0=-\frac{1}{2\pi}\int_0^{2\pi}\mathrm{d}\varphi T_{\bar{z}\bar{z}}$$

$$b_0=-\frac{1}{2\pi}\int_0^{2\pi}\mathrm{d}\varphi b_{zz}$$

$$\tilde{b}_0=-\frac{1}{2\pi}\int_0^{2\pi}\mathrm{d}\varphi b_{\bar{z}\bar{z}}$$

$$\hat{\varphi}=\varphi-\alpha f(u),$$

$$\mathrm{d}s^2=\mathrm{d}u^2+\mathrm{d}(\hat{\varphi}+\alpha f(u))^2.$$

$$\Psi_\alpha=\frac{1}{4\pi}\int_0^{2\pi}\mathrm{d}\varphi\int_0^s\mathrm{d}u\frac{\partial(\sqrt{g}g_{ij})}{\partial\alpha}b^{ij}$$

$$\frac{\partial g_{ij}}{\partial \alpha}=D_iv_j+D_jv_i$$

$$v=f(u)\frac{\partial}{\partial\varphi},$$

$$\Psi_\alpha=\frac{1}{4\pi}\int_0^{2\pi}\mathrm{d}\varphi\int_0^s\mathrm{d}u\sqrt{g}b^{ij}(D_iv_j+D_jv_i)$$

$$\Psi_\alpha=-\frac{1}{2\pi}\int_0^{2\pi}\mathrm{d}\varphi b_{u\varphi}(\varphi,s)=b_0-\tilde{b}_0$$

$$\exp\left(-s(L_0+\tilde{L}_0)\right)\exp\left(-i\alpha(L_0-\tilde{L}_0)\right)$$

$$2\Psi_s\Psi_\alpha\int_0^\infty\mathrm{d}s\int_0^{2\pi}\mathrm{d}\alpha\exp\left(-s(L_0+\tilde{L}_0)\right)\exp\left(-i\alpha(L_0-\tilde{L}_0)\right)$$

$$=4\pi\tilde{b}_0b_0\delta_{L_0-\tilde{L}_0}\int_0^\infty\mathrm{d}s\exp\left(-s(L_0+\tilde{L}_0)\right)=\frac{2\pi\tilde{b}_0b_0\delta_{L_0-\tilde{L}_0}}{L_0}$$



$$\begin{aligned}
& \tilde{b}_0 b_0 \int_{|q| \leq 1} \frac{d^2 q}{|q|^2} q^{L_0} \tilde{q}^{\tilde{L}_0} \\
& \frac{2\pi \tilde{b}_0 b_0 \delta_{L_0 - \tilde{L}_0} \Pi_{\text{GSO}}}{L_0} \\
& \frac{2\pi \tilde{b}_0 b_0 \delta(\beta_0) \delta_{L_0 - \tilde{L}_0} G_0 \Pi_{\text{GSO}}}{L_0} = \frac{2\pi \tilde{b}_0 b_0 \delta(\beta_0) \delta_{L_0 - \tilde{L}_0} \Pi_{\text{GSO}}}{G_0} \\
& \frac{2\pi \tilde{b}_0 b_0 \delta_{L_0 - \tilde{L}_0} \Pi_{\text{GSO}} \tilde{\Pi}_{\text{GSO}}}{L_0} \\
& \frac{2\pi \tilde{b}_0 b_0 \delta_{L_0 - \tilde{L}_0} \delta(\beta_0) G_0 \Pi_{\text{GSO}} \tilde{\Pi}_{\text{GSO}}}{L_0} \\
& \frac{2\pi \tilde{b}_0 b_0 \delta_{L_0 - \tilde{L}_0} \delta(\tilde{\beta}_0) \tilde{G}_0 \Pi_{\text{GSO}} \tilde{\Pi}_{\text{GSO}}}{L_0} \\
& \frac{2\pi \tilde{b}_0 b_0 \delta(\beta_0) G_0 \delta(\tilde{\beta}_0) \tilde{G}_0 \delta_{L_0 - \tilde{L}_0} \Pi_{\text{GSO}} \tilde{\Pi}_{\text{GSO}}}{L_0} \\
& \langle \phi_{\alpha\beta}(p) \phi_{\alpha'\beta'}(-p) \rangle = \frac{(\Gamma \cdot p)_{\alpha\alpha'} (\Gamma \cdot p)_{\beta\beta'}}{p^2} + \text{constant}.
\end{aligned}$$

Métrica de Feynman – singularidad y propagador de Lorentz – BRST – OPE.

$$\begin{aligned}
\frac{1}{p^2 + m^2} &= \int_0^\infty dt \exp(-t(p^2 + m^2)) \\
\frac{1}{p^2 + m^2 - i\epsilon} &= i \int_0^\infty d\tau \exp(-i\tau(p^2 + m^2) - \epsilon\tau)
\end{aligned}$$

$$\mathcal{D}_{\text{nonsep}} \cong \widehat{\mathcal{M}}_{g-1, n+2}.$$

$$\mathcal{D}_{\text{sep}} \cong \widehat{\mathcal{M}}_{g_1, n_1+1} \times \widehat{\mathcal{M}}_{g_2, n_2+1},$$

$$\mathbf{g}_1 + \mathbf{g}_2 = \mathbf{g}, \mathbf{n}_1 + \mathbf{n}_2 = \mathbf{n}.$$

$$(x-a)(y-b)=q.$$

$$\tilde{b}_0 b_0 \int_{|q| \leq 1} \frac{d^2 q}{|q|^2} q^{L_0} \bar{q}^{\tilde{L}_0}$$

$$\int \frac{d^d p}{(2\pi)^d} \frac{1}{p^2}$$



$$\begin{aligned}xy&=-\varepsilon^2\\y\theta&=\varepsilon\psi\\x\psi&=-\varepsilon\theta\\ \theta\psi&=0.\end{aligned}$$

$$q_{\mathrm{NS}}=-\varepsilon^2$$

$$D_\theta^*=\frac{\partial}{\partial\theta}+\theta x\frac{\partial}{\partial x}, D_\psi^*=\frac{\partial}{\partial\psi}+\psi y\frac{\partial}{\partial y}.$$

$$\begin{array}{l}xy=q_{\mathrm{R}}\\ \theta=\pm\sqrt{-1}\psi.\end{array}$$

$$\mathcal{D}_{\mathrm{nonsep}}\cong \widehat{\mathfrak{M}}_{\mathbf{g},\mathbf{n}_{\mathrm{NS}}+2,\mathbf{n}_{\mathrm{R}}}$$

$$\mathcal{D}_{\mathrm{sep}}\cong \widehat{\mathfrak{M}}_{\mathbf{g}_1,\mathbf{n}_{\mathrm{NS},1}+1,\mathbf{n}_{\mathrm{R},1}}\times \widehat{\mathfrak{M}}_{\mathbf{g}_2,\mathbf{n}_{\mathrm{NS},2}+1,\mathbf{n}_{\mathrm{R},2}}$$

$$\Pi\colon \mathcal{D}_{\mathrm{sep}}\rightarrow \widehat{\mathfrak{M}}_{\mathbf{g}_1,\mathbf{n}_{\mathrm{NS},1},\mathbf{n}_{\mathrm{R},1}+1}\times \widehat{\mathfrak{M}}_{\mathbf{g}_2,\mathbf{n}_{\mathrm{NS},2},\mathbf{n}_{\mathrm{R},2}+1}$$

$$G_0=\frac{\partial}{\partial\psi}-\psi y\frac{\partial}{\partial y}$$

$$\mathcal{D}_{\mathrm{sep}}\cong \widehat{\mathfrak{M}}_{\mathbf{g}_1,\mathbf{n}_{\mathrm{NS},1}+1,\mathbf{n}_{\mathrm{R},1}}\times \widehat{\mathfrak{M}}_{\mathbf{g}_2,\mathbf{n}_{\mathrm{NS},2}+1,\mathbf{n}_{\mathrm{R},2}'}$$

$$\mathrm{R}=b_0\Pi_0$$

$$\{Q_B,\mathrm{R}\}=\{Q_B,b_0\}\Pi_0=L_0\Pi_0=0$$

$$\oplus_{n\in\mathbb{Z}}\mathcal{H}_n\otimes\mathcal{H}_{n+1}^*$$

$$\mathrm{R}'\in\oplus_{n\in\mathbb{Z}}\mathcal{H}_n\otimes\mathcal{H}_{2-n}$$

$$\mathrm{R}'\in (\mathcal{H}\otimes\mathcal{H})_2$$

$$\mathrm{R}'=\sum_i cU_i\otimes cU^i+Q_B\mathcal{X}$$

$$\mathcal{X}=\sum_j \mathcal{S}_j\otimes \mathcal{T}_j, \mathcal{S}_j, \mathcal{T}_j\in \mathcal{H}$$

$$\mathrm{d}F_\Omega + F_{Q_B\Omega} = 0.$$

$$\mathrm{d}F_{\nu_1\cdots\nu_n}\mathcal{X}+F_{\nu_1\cdots\nu_n}Q_B\mathcal{X}=0$$

$$\int_{\mathcal{M}_{\mathbf{g}-1,\mathbf{n}+2}}F_{\nu_1\cdots\nu_n}Q_B\mathcal{X}=-\int_{\mathcal{M}_{\mathbf{g}-1,\mathbf{n}+2}}\mathrm{d}F_{\nu_1\cdots\nu_n}\mathcal{X}$$

$$\widetilde{\mathrm{R}}'=\sum_i\left(c\partial cU_i\otimes cU^i+cU_i\otimes c\partial cU^i\right)+Q_B\mathcal{Y}$$



$$(c\partial c\partial^2c\otimes 1+1\otimes c\partial c\partial^2c)$$

Cohomología de Ramond.

$$R'=\sum_i c\delta(\gamma)U_i\otimes c\delta(\gamma)U^i+Q_B\mathcal{X}$$

$$R=b_0\delta(\beta_0)G_0\Pi_0$$

$$R'\in (\mathcal{H}\otimes \mathcal{H})_{1;-1/2\otimes -1/2}$$

$$R'=\sum_i c\Theta_{-1/2}\Phi_i\otimes c\Theta_{-1/2}\Phi_i'+Q_B\mathcal{X}$$

$$\sum_{\alpha\beta} (p\cdot\Gamma)^{\alpha\beta} c\Theta_{-1/2}\Sigma_{\alpha}(p)\otimes c\Theta_{-1/2}\Sigma_{\beta}(-p)$$

$$R'=\sum_i \tilde{c}cU_i\otimes \tilde{c}cU^i+\{Q_B,\mathcal{X}\}$$

$$\widetilde{R}'=\sum_i\left(\tilde{c}c(\tilde{\partial}\tilde{c}+\partial c)U_i\otimes\tilde{c}cU^i+\tilde{c}cU_i\otimes\tilde{c}c(\tilde{\partial}\tilde{c}+\partial c)U^i\right)+\{Q_B,\mathcal{X}\}$$

$$\sum_I\left(\tilde{c}\tilde{\partial}\tilde{c}\tilde{\partial}^2\tilde{c}\cdot c\partial X^I\otimes c\partial X_I+c\partial X_I\otimes\tilde{c}\tilde{\partial}\tilde{c}\tilde{\partial}^2\tilde{c}\cdot c\partial X^I\right)+z\leftrightarrow\tilde{z}$$

$$(x-a)(y-b)=q$$

$$(x,a)\rightarrow (\lambda x,\lambda a), (y,b)\rightarrow (\tilde{\lambda}y,\tilde{\lambda}b), q\rightarrow \lambda\tilde{\lambda}q.$$

$$\Omega\sim\mathrm{d}a\,\mathrm{d}b\,\frac{\mathrm{d}q}{q^2}$$

$$\Omega\sim c(a)\otimes c(b)\frac{\mathrm{d}q}{q^2}$$

$$\mathcal{V}_\ell(a)\otimes \mathcal{V}_r(b)\mathrm{d}q q^{L_0-1}$$

$$\begin{array}{l} (x-a-\alpha\theta)(y-b-\beta\psi)=-\varepsilon^2\\ (y-b-\beta\psi)(\theta-\alpha)=\varepsilon(\psi-\beta)\\ (x-a-\alpha\theta)(\psi-\beta)=-\varepsilon(\theta-\alpha)\\ (\theta-\alpha)(\psi-\beta)=0 \end{array}$$

$$\begin{array}{l} (x,a,\alpha)\!\!\rightarrow\!\! (\lambda x,\lambda a,\lambda^{1/2}\alpha)\\ (y,b,\beta)\!\!\rightarrow\!\! (\tilde{\lambda}y,\tilde{\lambda}b,\tilde{\lambda}^{1/2}\beta)\\ \varepsilon\rightarrow (\lambda\tilde{\lambda})^{1/2}\varepsilon \end{array}$$

$$\Omega\sim [\mathrm{d}a\mid \mathrm{d}\alpha]\otimes [\mathrm{d}b\mid \mathrm{d}\beta]\frac{\mathrm{d}\varepsilon}{\varepsilon^2}$$



$$\Omega \sim c\delta(\gamma) \otimes c\delta(\gamma) \frac{d\varepsilon}{\varepsilon^2}$$

$$\mathcal{V}_\ell(a\mid\alpha)\otimes\mathcal{V}_r(b\mid\beta)d\varepsilon\varepsilon^{2L_0-1}$$

$$d\varepsilon\varepsilon^{2L_0-1}\sim dq_{\rm NS}q_{\rm NS}^{L_0-1}$$

$$(\hat{x}-\hat{a})(\hat{y}-\hat{b})=\hat{q}$$

$$\hat{q}=q\frac{\partial \hat{a}}{\partial a}\frac{\partial \hat{b}}{\partial b}$$

Compactificación Deligne-Mumford.

$$\overline{\tilde{q}}=q_{\rm NS}$$

$$\overline{\tilde{q}}=q_{\rm NS}\Big(1+\mathcal{O}(\eta_i\eta_j)\Big),$$

$$\overline{\tilde{q}}=q_{\rm NS}+\mathcal{O}(\eta_i\eta_j).$$

$$\Xi=[\mathrm{d}\tilde{q};\mathrm{d}q_{\mathrm{NS}}\mid \mathrm{d}\eta_1,\mathrm{d}\eta_2]\tilde{q}^{-1}.$$

$$\int \Xi \rightarrow \int \Xi + \int [\mathrm{d}\tilde{q};\mathrm{d}q_{\mathrm{NS}}\mid \mathrm{d}\eta_1,\mathrm{d}\eta_2]\eta_1\eta_2\frac{\partial}{\partial q_{\mathrm{NS}}}\frac{1}{\tilde{q}}.$$

$$\tilde{q}\rightarrow e^{\tilde{\varphi}}\tilde{q},q_{\rm NS}\rightarrow e^{\varphi}q_{\rm NS}$$

$$\begin{array}{l} \tilde{q}=\tilde{u}_1-\tilde{u}_2\\ q_{\rm NS}=u_1-u_2-\zeta_1\zeta_2. \end{array}$$

$$q_{\rm NS}(1+\mathcal{O}(\eta^2))=t^1+it^2$$

$$\tilde{q}_{\rm NS}(1+\mathcal{O}(\tilde{\eta}^2))=t^1-it^2$$

$$t^1=q_{\rm NS}(1+\mathcal{O}(\eta^2))$$

Anomalías BRST para las partículas sin masa. Supersimetrías de gauge.

$$\int_{\Gamma} F_{\{Q_B,\mathcal{W}_1\},\mathcal{V}_2,\ldots,\mathcal{V}_n}=-\int_{\partial\Gamma} F_{\mathcal{W}_1,\mathcal{V}_2,\ldots,\mathcal{V}_n}$$

$$b_0\int_0^\infty\mathrm{d}s\exp\left(-sL_0\right)$$

$$2\pi(b_0-\tilde{b}_0)\delta_{L_0-\tilde{L}_0}\exp\left(-s(L_0+\tilde{L}_0)\right).$$

$$\sum_i\left(c\partial cU_i\otimes cU^i+cU_i\otimes c\partial cU^i\right)$$

$$\mathcal{O}=\sum_i a_i \mathcal{Y}^i$$



$$\sum_i a_i \langle y^i \mathcal{V}_n \rangle_{\mathrm{gr}}.$$

$$\mathcal{Q}_B(\mathcal{W}_1)=\{Q_B,\mathcal{W}_1\}-g_{\mathrm{st}}^2\mathfrak{g}_\ell\mathcal{O}(\mathcal{W}_1),$$

$$\mathcal{Q}_B(\mathcal{V}_n)=\{Q_B,\mathcal{V}_n\}+\sum_i y^i c\partial c U_i,$$

Rompimiento de simetría de gauge - BRST.

$$\sigma \rightarrow \sigma - g_{\mathrm{st}}^2 \mathfrak{g}_\ell \lambda$$

$$\int \left(F_{IJ} F^{IJ} + (\partial_I \sigma + A_I)^2 \right)$$

$$\lambda \rightarrow \lambda - g_{\mathrm{st}}^{2g_\ell} \zeta.$$

$$\mathcal{W}_1=\varepsilon_I\cdot(c^z\partial_zX^I-\tilde{c}\tilde{z}\partial_{\tilde{z}}X^I)\exp{(ip\cdot X)},p^2=\varepsilon\cdot p=0,$$

$$W=\varepsilon_I\,\mathrm{d}X^I\exp{(ip\cdot X)},p^2=\varepsilon\cdot p=0.$$

$$\mathcal{O}(\mathcal{W})=\mathcal{W}_{\mathrm{open}}, \mathcal{W}_{\mathrm{open}}=cW_{\mathrm{open}}$$

$$\int \mathrm{d}^Dx \left(\frac{1}{g_{\mathrm{st}}^2} H_{IJK}^2 + \frac{1}{g_{\mathrm{st}}} \left(\partial_I A_J - \partial_J A_I + B_{IJ} \right)^2 \right), H = \mathrm{d}B$$

$$\int_{\mathbb{R}^4} \mathrm{d}^4x e^{-2\phi} ((\mathrm{d}A)^2 + (\mathrm{d}B)^2)$$

$$\int_{\mathbb{R}^4} B \wedge \mathrm{d}A$$

$$xy=q,$$

$$\mathcal{D}=\widehat{\mathcal{M}}_\ell\times\widehat{\mathcal{M}}_r$$

$$q\rightarrow e^{f_\ell+f_r}q$$

$$\mathcal{N}=\mathcal{L}_\ell\otimes\mathcal{L}_r$$

$$\mathcal{A}_g(p_1,\zeta_1;\ldots;p_n,\zeta_n)=\int_{\widehat{\mathcal{M}}}F_{p_1,\zeta_1;\ldots;p_n,\zeta_n}(g\mid\delta g)$$

$$\mathcal{A}_{g,\,\mathrm{sing}}=\sum_{\alpha=1}^s\int\frac{\mathrm{d}^2q}{\bar{q}q}\int_{\widehat{\mathcal{M}}_\ell}\mathcal{G}_{\ell,\alpha}\int_{\widehat{\mathcal{M}}_r}\mathcal{G}_{r,\alpha}.$$

$$\int_{\widehat{\mathcal{M}}_r} \mathcal{G}_{r,\alpha} = 0, \alpha = 1, \ldots, s.$$



$$\int_{\epsilon < |q| < \Lambda} \frac{d^2 q}{\bar{q} q}$$

$$\mathcal{A}_{\mathbf{g},\epsilon}=\int_{\widehat{\mathcal{M}}_{\epsilon}}F(g\mid\delta g),$$

$$\mathcal{A}_{\mathbf{g}}\rightarrow\mathcal{A}_{\mathbf{g}}-4\pi\sum_{\alpha}\int_{\widehat{\mathcal{M}}_{\ell}\times\widehat{\mathcal{M}}_r}h\mathcal{G}_{\ell,\alpha}\mathcal{G}_{r,\alpha}$$

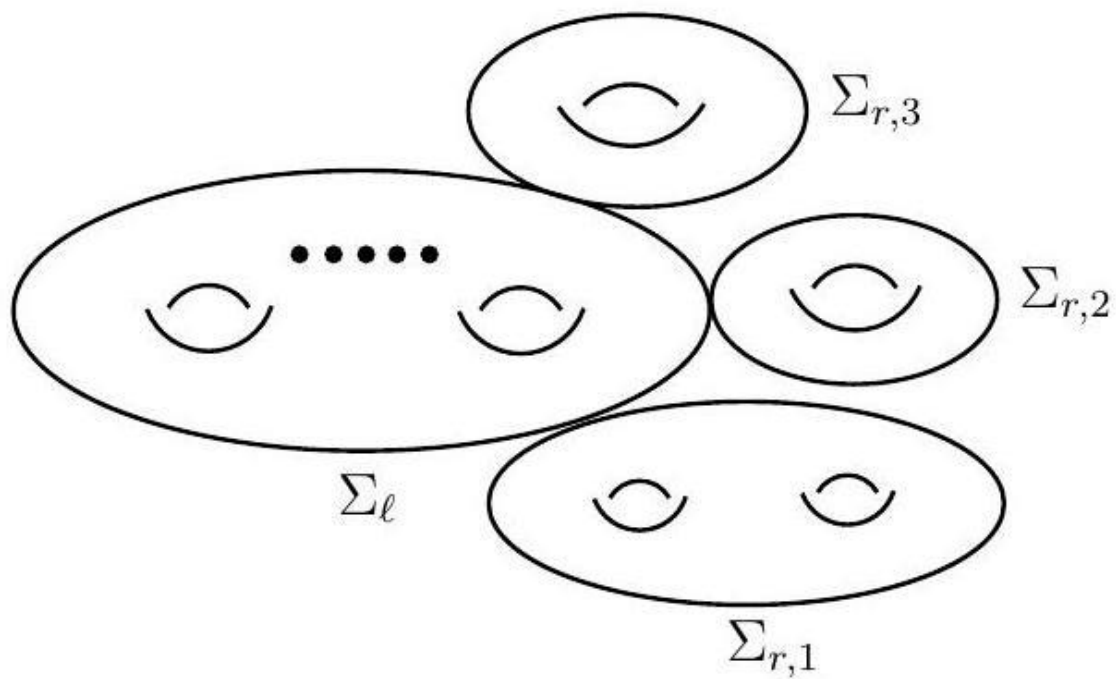
$$\Delta_{g_r}\phi_\alpha=-\frac{1}{4\pi}\int_{\widehat{\mathcal{M}}_r}h_r\mathcal{G}_{r,\alpha}$$

$$\mathcal{A}_{\mathbf{g}}\rightarrow\mathcal{A}_{\mathbf{g}}+\sum_{\alpha}\Delta_{\mathbf{g}_r}\phi_{\alpha}\int_{\widehat{\mathcal{M}}_{\ell}}\mathcal{G}_{\ell,\alpha}$$

$$\mathcal{A}_{\mathbf{g}}\rightarrow\mathcal{A}_{\mathbf{g}}+\sum_{\alpha}\Delta_{\mathbf{g}_r}\phi_{\alpha}\frac{\partial}{\partial\phi_{\alpha}}\mathcal{A}_{\mathbf{g}_{\ell}}$$

$$\mathcal{A}_{\mathbf{g}}\rightarrow\mathcal{A}_{\mathbf{g}}+\sum_{\mathbf{g}_{\ell}+\mathbf{g}_r=\mathbf{g}}\sum_{\alpha}\Delta_{\mathbf{g}_r}\phi_{\alpha}\frac{\partial}{\partial\phi_{\alpha}}\mathcal{A}_{\mathbf{g}_{\ell}}$$

$$\mathcal{A}\rightarrow\mathcal{A}+\sum_{\alpha}\Delta\phi_{\alpha}\frac{\partial}{\partial\phi_{\alpha}}\mathcal{A}.$$



$$\mathcal{A}\rightarrow\exp\left(\sum_{\alpha}\Delta\phi_{\alpha}\frac{\partial}{\partial\phi_{\alpha}}\right)\mathcal{A}.$$

$$\mathcal{K} = \sum_{\alpha} \Delta\phi_{\alpha} \frac{\partial}{\partial\phi_{\alpha}}$$

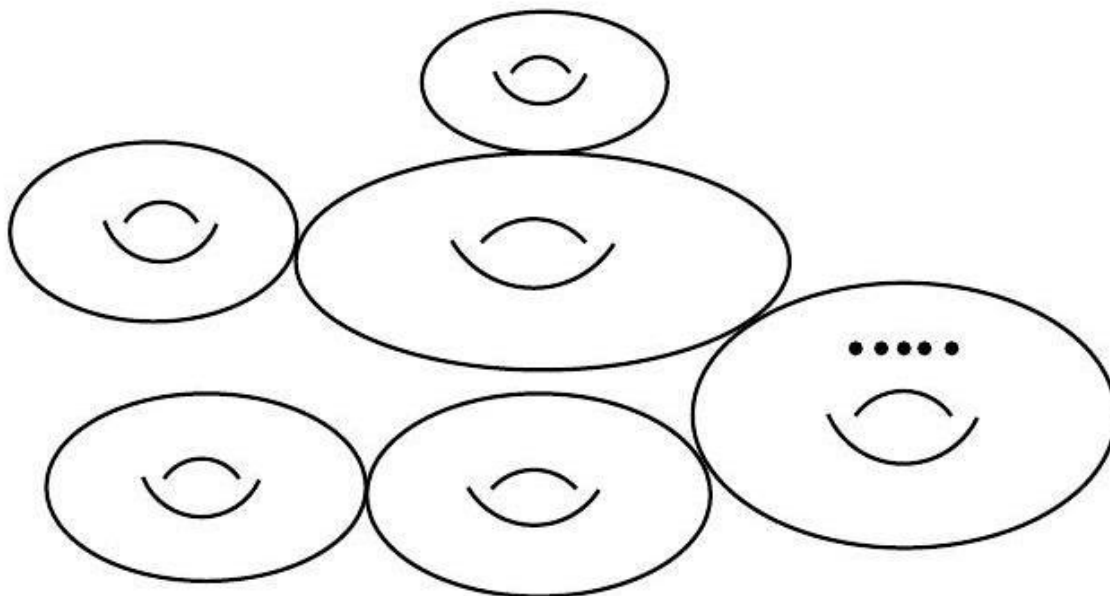


Figura 1. Superespacios producidos por supermembranas.

$$F_0 = \sum_{\alpha} \frac{d^2q}{\bar{q}q} \wedge \mathcal{G}_{\ell,\alpha} \wedge \mathcal{G}_{r,\alpha}$$

$$\mathcal{A}_{g,\epsilon} = \int_{\widehat{\mathcal{M}}_{\epsilon}} F - \int_{\partial\widehat{\mathcal{M}}_{\epsilon}} \Lambda$$

$$\chi_{r,\alpha} = \frac{d^2q}{\bar{q}q} \mathcal{G}_{r,\alpha}$$

$$\chi_{r,\alpha} = d\lambda_{r,\alpha}$$

$$\Lambda = \sum_{\alpha} \mathcal{G}_{\ell,\alpha} \wedge \lambda_{r,\alpha}$$

$$\Delta\phi_{\alpha} = - \int_{\mathcal{X}_{\epsilon}} \Delta\lambda_{r,\alpha}$$

$$\Lambda \rightarrow \Lambda + \sum_{\alpha} \mathcal{G}_{\ell,\alpha} \wedge \Delta\lambda_{r,\alpha}$$

$$\mathcal{A}_{\mathbf{g}} \rightarrow \mathcal{A}_{\mathbf{g}} + \sum_{\alpha} \Delta\phi_{\alpha} \int_{\widehat{\mathcal{M}}_{\ell}} \mathcal{G}_{\ell,\alpha}.$$

$$\lambda_{r,\alpha}^{(0)} = \frac{d^2q}{\bar{q}q} \beta_{r,\alpha}.$$

$$\lambda_{r,\alpha}=\frac{\mathrm{d}^2q}{\bar{q}q}\beta_{r,\alpha}+\frac{\mathrm{d}q}{q}\gamma_{r,\alpha}+\frac{\mathrm{d}\bar{q}}{\bar{q}}\tilde{\gamma}_{r,\alpha},$$

Renormalización de la función de onda.

$$\mathcal{A}_{\mathrm{g},\mathrm{sing}}=\sum_{\alpha=1}^s\int\frac{\mathrm{d}^2q}{\bar{q}q}\int_{\widehat{\mathcal{M}}_\ell}\mathcal{G}_{\ell,\alpha}\int_{\widehat{\mathcal{M}}_r}\mathcal{G}_{r,\alpha}$$

$$\int_{\widehat{\mathcal{M}}_\ell}\mathcal{G}_{\ell,\alpha}=0$$

$$xy=q$$

$$\mathcal{B}\cong \widehat{\mathcal{M}}_\ell\times \widehat{\mathcal{M}}_r$$

$$\sum_i\left(c\partial cU_i\otimes cU^i+cU_i\otimes c\partial cU^i\right)$$

$$\mathcal{O}(\mathcal{W}_1)=\sum_i a_i \mathcal{V}_i$$

$$\int_{\widehat{\mathcal{M}}_{\mathrm{g},\mathrm{n}}}F_{\{Q_B,\mathcal{W}_1\},\mathcal{V}_2,\ldots,\mathcal{V}_{\mathrm{n}}}=-\int_{\partial\widehat{\mathcal{M}}_{\mathrm{g},\mathrm{n}}}F_{\mathcal{W}_1,\mathcal{V}_2,\ldots,\mathcal{V}_{\mathrm{n}}}$$

$$\frac{\mathrm{d}q}{q}\sum_i cU_i\otimes cU^i$$

$$\frac{\mathrm{d}q}{q}\sum\; cU_i\otimes cU^i+\sum_i\; (c\partial cU_i\otimes cU^i+cU_i\otimes c\partial cU^i)$$

$$v(x)\partial_x+\frac{\partial v(x)}{\partial x}\Big|_{x=0}q\partial_q, \, v(0)=0$$

$$\mathcal{X}_i=\frac{\mathrm{d}q}{q}cU_i+c\partial cU_i$$

$$a_i=\int_{\mathcal{M}^*_\ell}F_{\mathcal{W}_1,\ldots,\mathcal{X}_i}$$

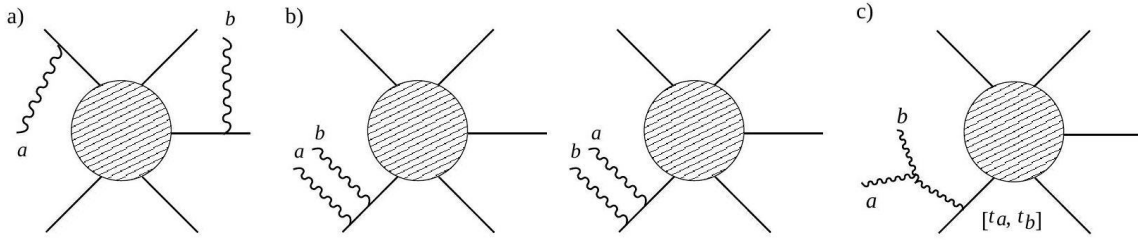
$$\frac{\mathrm{d}(|q|^2)}{|q|^2}\sum_i\;\tilde{c}cU_i\otimes\tilde{c}cU^i+\sum_i\;(\tilde{c}c(\tilde{\partial}\tilde{c}+\partial c)U_i\otimes\tilde{c}cU^i+\tilde{c}cU_i\otimes\tilde{c}c(\tilde{\partial}\tilde{c}+\partial c)U^i)$$

Supersimetrías de espacio – tiempo.

$$\frac{i}{p_i^2-m_i^2}$$

$$-ie_i\varepsilon\cdot(p_i+p_i'),p_i'=p_i+k$$





$$-ie_i \varepsilon \cdot (p_i + p'_i) \frac{i}{(p'_i)^2 - m_i^2} = \frac{e_i \varepsilon \cdot p_i}{k \cdot p_i},$$

$$\mathcal{A} \sim \sum_{i=1}^n \frac{e_i \varepsilon \cdot p_i}{k \cdot p_i} \mathcal{A}'.$$

$$\sum_i e_i = 0,$$

$$-it_{a,i} \varepsilon \cdot (p + p')$$

$$\mathcal{A} \sim \sum_{i=1}^n \frac{t_{a,i} \varepsilon \cdot p_i}{k \cdot p_i} \mathcal{A}'$$

$$\sum_i t_{a,i} \mathcal{A}' = 0.$$

$$\sum_{i \neq i'} \frac{t_{a,i} \varepsilon_a \cdot p_i}{k_a \cdot p_i} \frac{t_{b,i'} \varepsilon_b \cdot p_{i'}}{k_b \cdot p_{i'}} \mathcal{A}'$$

$$\sum_{i \neq i'} t_{a,i} \frac{t_{b,i'} \varepsilon_b \cdot p_{i'}}{k_b \cdot p_{i'}} \mathcal{A}' = - \sum_i \frac{t_{b,i} \varepsilon_b \cdot p_i}{k_b \cdot p_i} t_{a,i} \mathcal{A}'$$

$$\left(\sum_i \frac{t_{a,i} \varepsilon_a \cdot p_i}{k_a \cdot p_i} \frac{t_{b,i} \varepsilon_b \cdot p_i}{(k_a + k_b) \cdot p_i} + a \leftrightarrow b \right) \mathcal{A}'.$$

$$\left(\sum_i \frac{t_{b,i} \varepsilon_b \cdot p_i}{k_b \cdot p_i} t_{a,i} + \sum_i [t_{a,i}, t_{b,i}] \frac{\varepsilon_b \cdot p_i}{(k_a + k_b) \cdot p_i} \right) \mathcal{A}'.$$

Gravedad y supergravedad. Métrica de Ward.

$$\sum_{i=1}^n Q_{\alpha,i} \mathcal{A}' = 0$$

$$\mathcal{V} = \{Q_B, \mathcal{W}\} = \tilde{c} c i k_j \varepsilon_l \tilde{\partial} X^j \partial X^l e^{ik \cdot X}$$

$$\mathcal{V} = \tilde{c} c \tilde{\partial} (\varepsilon_l \partial X^l e^{ik \cdot X})$$

$$0 = \{Q_B, \mathcal{W}\}$$



$$0 = \tilde{\partial} J, J = \varepsilon_I \partial X^I.$$

$$J^I = \partial X^I.$$

$$\frac{1}{2\pi\alpha'} \oint \gamma J^I \cdot \nu = p^I \nu$$

$$0 = \left\langle \int_{\Sigma'} dJ \cdot \nu_1 \dots \nu_n \right\rangle = \sum_{i=1}^n \langle \nu_1 \dots \nu_{i-1} (\oint \gamma_i J \cdot \nu_i) \nu_{i+1} \dots \nu_n \rangle.$$

$$0 = \left(\sum_{i=1}^n p_i \right) \langle \nu_1 \dots \nu_n \rangle.$$

$$\sum_{i=1}^n p_i = 0.$$

Spacetime Supersymmetry. Métrica de Ramond – BRST – Stokes – Ward. Conservación de energía – momentum – stress.

$$\mathcal{W}(k,u) = c\Theta_{-1/2}u^\alpha\Sigma_\alpha e^{ik\cdot X}$$

$$(\gamma \cdot k)_{\alpha\beta} u^\beta = 0$$

$$\mathcal{V}(k,u) = \{Q_B, \mathcal{W}(k,u)\} = i\tilde{c}ck \cdot \tilde{\partial} X \Theta_{-1/2} u^\alpha \Sigma_\alpha e^{ik\cdot X}$$

$$\mathcal{S}_\alpha = c\Theta_{-1/2}\Sigma_\alpha$$

$$dF_{\mathcal{S}_\alpha} \nu_1 \dots \nu_n = 0$$

$$0 = \int_{\Gamma} dF_{\mathcal{S}_\alpha} \nu_1 \dots \nu_n = \int_{\partial\Gamma} F_{\mathcal{S}_\alpha} \nu_1 \dots \nu_n$$

$$\sum_{i=1}^n \langle \nu_1 \dots \nu_{i-1} Q_\alpha(\nu_i) \nu_{i+1} \dots \nu_n \rangle = 0$$

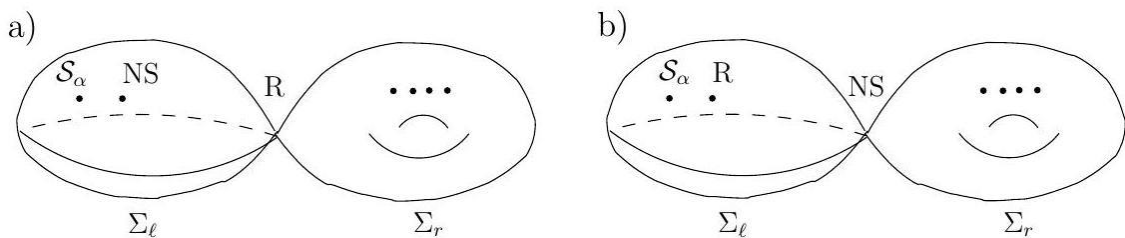
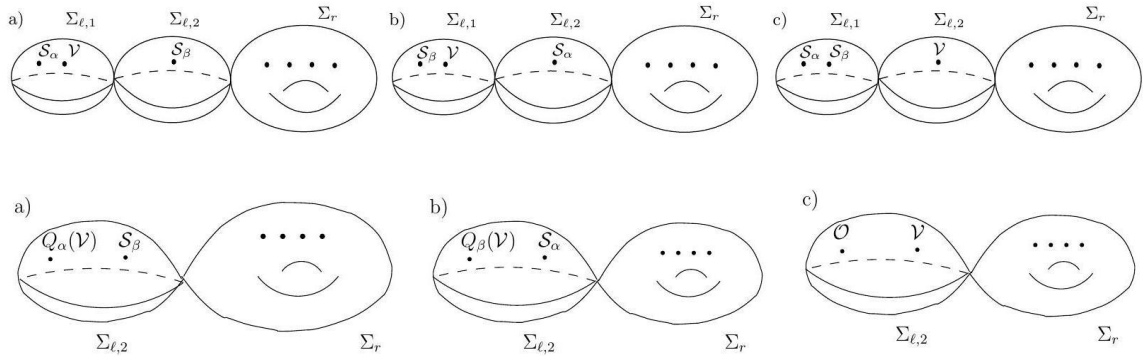


Figura 2. Supermembranas.

$$\mathcal{S}_\alpha = c\hat{S}_\alpha, \hat{S}_\alpha = \Theta_{-1/2}\Sigma_\alpha$$

$$\frac{1}{2\pi i}\oint_{|z|=\epsilon}dz\hat{S}_\alpha\cdot\mathcal{V}$$

$$0=\int_{\widehat{\mathcal{M}}_{\mathbf{g},\mathbf{n}+1}}\mathrm{d}F_{\mathcal{P}^I\mathcal{V}_1\dots\mathcal{V}_\mathbf{n}}=\int_{\partial\widehat{\mathcal{M}}_{\mathbf{g},\mathbf{n}+1}}F_{\mathcal{P}^I\mathcal{V}_1\dots\mathcal{V}_\mathbf{n}}$$



$$0=\int_B\mathrm{d}F(\mathcal{J},\delta\mathcal{J})=\int_{\partial B}F(\mathcal{J},\delta\mathcal{J})$$

$$Q_\alpha(Q_\beta(\mathcal{V})) + Q_\beta(Q_\alpha(\mathcal{V})) + \mathcal{O}(\mathcal{V}) = 0$$

$$\mathcal{O}=-\{Q_\alpha,Q_\beta\}$$

$$\tilde{c}\tilde{\partial}\tilde{c}\tilde{\partial}^2\tilde{c}c\delta(\gamma)D_\theta X^I\otimes c\delta(\gamma)D_\theta X_I+\cdots$$

$$\mathcal{O}_{\mathcal{S}_\alpha\mathcal{S}_\beta}=\frac{1}{2\pi i}\oint_{|z|=\epsilon}dz\hat{S}_\alpha(z)\mathcal{S}_\beta(0)$$

$$\mathcal{V}_\phi=\sum_\alpha Q_\alpha(\mathcal{V}_{\psi_\alpha})$$

$$\mathcal{V}_\phi=\tilde{c}c\delta(\gamma)\tilde{\partial}X_IDX^I$$

$$\mathcal{V}_{\psi_\alpha}=\Gamma_I^{\alpha\beta}\tilde{c}c\tilde{\partial}X^I\Theta_{-1/2}\Sigma_\beta$$

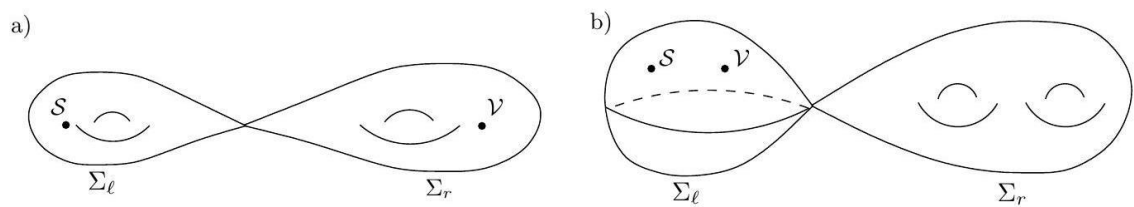


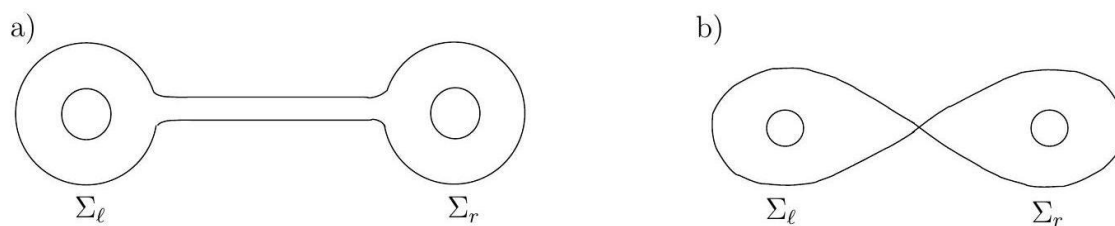
Figura 3. Supermembranas.

$$\sum_\alpha \langle s_\alpha \mathcal{V}_{\psi_\alpha} \rangle$$

$$0 = \int_{\Gamma} dF_{\mathcal{SV}} = \int_{\partial\Gamma} F_{\mathcal{SV}}$$

Supermembranas. Métrica de Ward – Chan – Paton. Función Delta.

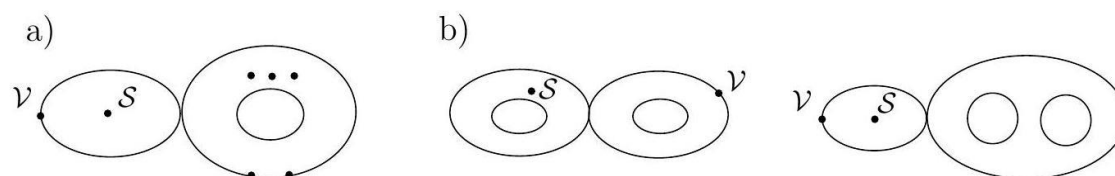
$$\mathcal{V}_{\text{NS,NS}} = \sum_{\alpha} \{Q'_{\alpha}, \mathcal{V}_{\text{R,NS}}^{\alpha}\} = \sum_{\beta} \{Q''_{\beta}, \mathcal{V}_{\text{NS,R}}^{\beta}\}$$



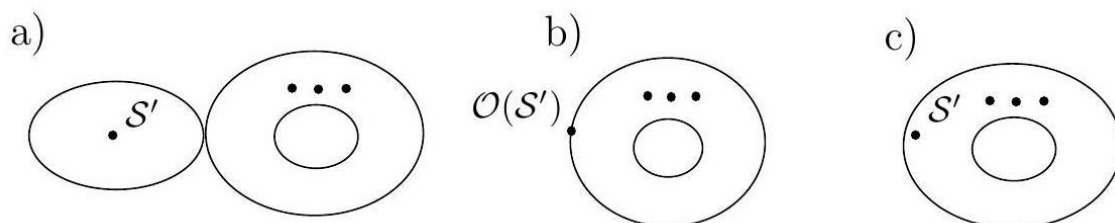
$$q \rightarrow e^f q$$

$$q \rightarrow e^{f_{\ell} + f_r} q$$

$$\frac{q_{\Sigma_{\ell}, \mathbb{D}}}{q_{\Sigma_{\ell}, \mathbb{RP}^2}} \rightarrow e^{\kappa} \frac{q_{\Sigma_{\ell}, \mathbb{D}}}{q_{\Sigma_{\ell}, \mathbb{RP}^2}}, \kappa \in \mathbb{R}$$



$$0 = \int_{\partial\tilde{\Gamma}} F_{S_{\alpha}} \mathcal{V}_1 \dots \mathcal{V}_n$$



$$\mathcal{O}(\mathcal{S}'_{\alpha}) + \mathcal{O}(\mathcal{S}''_{\alpha}) = 0$$

$$\mathcal{S}_{\alpha} = \mathcal{S}'_{\alpha} + \phi_{\text{D}}(\mathcal{S}'_{\alpha})$$

$$\langle \mathcal{V}_{\text{NS-NS}} \rangle + \langle \mathcal{V}_{\text{R-R}} \rangle = 0,$$

$$\mathcal{V}_{\text{NS-NS}} = \sum_{\alpha} \{Q'_{\alpha}, \mathcal{V}_{\text{R-NS}}^{\alpha}\}$$

$$\mathcal{V}_{\text{NS-NS}} + \mathcal{V}_{\text{R-R}} = \sum_{\alpha} \{Q'_{\alpha} + Q''_{\alpha}, \mathcal{V}_{\text{R-NS}}^{\alpha}\}.$$

$$\mathcal{V}_{\text{NS-NS}} + \mathcal{V}_{\text{R-R}} = \sum_{\alpha} \{Q'_{\alpha} + Q''_{\alpha}, \mathcal{V}_{\text{R-NS}}^{\alpha} + \mathcal{V}_{\text{NS-R}}^{\alpha}\}$$

$$\begin{aligned} \mathcal{Y}^I &= c\delta(\gamma)D_{\theta}X^I\exp\left(ik\cdot X\right) \\ \mathcal{Z}_{\alpha} &= c\Theta_{-1/2}\Sigma_{\alpha}\exp\left(ik\cdot X\right) \end{aligned}$$

$$\begin{aligned} \tilde{\mathcal{Y}}^I &= \tilde{c}\delta(\tilde{\gamma})D_{\tilde{\theta}}X^I\exp\left(ik\cdot X\right) \\ \tilde{\mathcal{Z}}_{\alpha} &= \tilde{c}\tilde{\Theta}_{-1/2}\tilde{\Sigma}_{\alpha}\exp\left(ik\cdot X\right) \end{aligned}$$

$$\begin{aligned} \mathcal{Z}_*^{\alpha} &= c\Theta_{-3/2}\Sigma^{\alpha}\exp\left(ik\cdot X\right) \\ \tilde{\mathcal{Z}}_*^{\alpha} &= \tilde{c}\tilde{\Theta}_{-3/2}\tilde{\Sigma}^{\alpha}\exp\left(ik\cdot \tilde{X}\right) \end{aligned}$$

$$y\cdot \mathcal{Z}_*^{\alpha}=(\Gamma\cdot k)^{\alpha\beta}\mathcal{Z}_{\beta}+\{Q_B,\cdot\}$$

$$\hat{\mathcal{S}}_{\alpha}(z)\mathcal{Z}_{\beta}(w)\sim\frac{1}{z-w}\Gamma_{\alpha\beta}^I\mathcal{Y}_I$$

$$\{Q_{\alpha},\mathcal{Z}_{\beta}\}=\Gamma_{\alpha\beta}^I\mathcal{Y}_I$$

$$\hat{\mathcal{S}}_{\alpha}(z)\mathcal{Y}_I(w)\sim\frac{1}{z-w}\Gamma_{I\alpha\beta}\mathcal{Z}_*^{\beta}$$

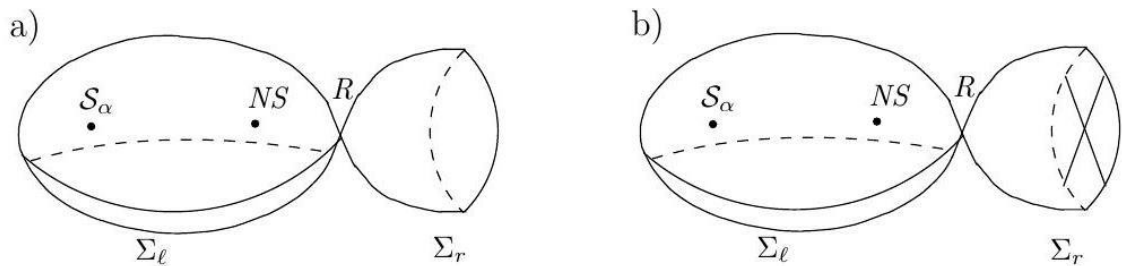
$$\{Q_{\alpha},\mathcal{Y}_I\}=(\Gamma_I\Gamma\cdot k)_{\alpha}^{\;\;\beta}\mathcal{Z}_{\beta}$$

$$\mathcal{S}_{\alpha}(z)\mathcal{Y}_I(w)\sim\Gamma_{I\alpha\beta}\partial c\mathcal{Z}_*^{\beta}$$

$$\tilde{v}=\partial_{\tilde{\theta}}-\tilde{\theta}\tilde{z}\partial_{\tilde{z}}$$

$$\tilde{v}^2=-\tilde{z}\partial_{\tilde{z}}$$

$$\tilde{\theta}\rightarrow\tilde{\theta}+\alpha,$$



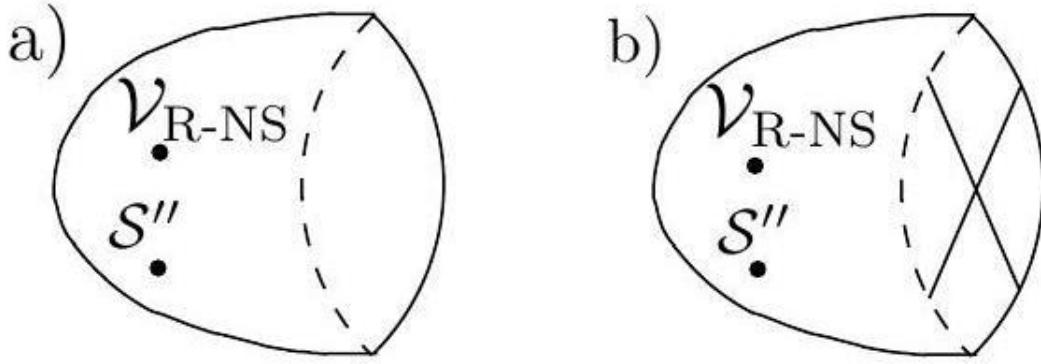


Figura 4. Supermembranas.

$$0 = \int_{\partial \hat{\Gamma}} F_{\hat{s}\alpha} \mathcal{V}_{\text{NS/R}}^{\alpha}$$

$$\langle \mathcal{V}_{\text{NS-NS}} \rangle_{\Sigma} = 0$$

$$\langle \mathcal{V}_{\text{NS-NS}} \rangle_{\Sigma} + \langle V_{\text{R-R}} \rangle_{\Sigma} = 0.$$

$$\mathcal{V}_{\text{R-R}} = \tilde{\partial} \tilde{c} \tilde{Z}_{*}^{\alpha} Z_{\alpha} + \tilde{Z}_{\alpha} \partial c Z_{*}^{\alpha}$$

$$\langle \mathcal{V}_{\text{NS-NS}} \rangle_{\text{D}} + \langle \mathcal{V}_{\text{NS-NS}} \rangle_{\mathbb{RP}^2} = 0.$$

$$I_1 \sim \int [dm_1 \dots | \dots d\eta_s] \mathcal{G}_1(m_1 \dots; q_{\Sigma_{\ell}, \text{D}} | \dots \eta_s) \frac{dq_{\Sigma_{\ell}, \text{D}}}{q_{\Sigma_{\ell}, \text{D}}}$$

$$I_2 \sim \int [dm_1 \dots | \dots d\eta_s] \mathcal{G}_2(m_1 \dots; q_{\Sigma_{\ell}, \mathbb{RP}^2} | \dots \eta_s) \frac{dq_{\Sigma_{\ell}, \mathbb{RP}^2}}{q_{\Sigma_{\ell}, \mathbb{RP}^2}}$$

$$I_{1,\log} = \int_{\mathcal{M}_{\ell}} [dm_1 \dots | \dots d\eta_s] \mathcal{G}_1(m_1 \dots; 0 | \dots \eta_s)$$

$$I_{2,\log} = \int_{\mathcal{M}_{\ell}} [dm_1 \dots | \dots d\eta_s] \mathcal{G}_2(m_1 \dots; 0 | \dots \eta_s)$$

$$\mathcal{G}_1(m_1 \dots; 0 | \dots \eta_s) = \mathcal{G}_0(m_1 \dots | \dots \eta_s) \langle \mathcal{V}_{\text{NS-NS}} \rangle_{\text{D}}$$

$$\mathcal{G}_2(m_1 \dots; 0 | \dots \eta_s) = \mathcal{G}_0(m_1 \dots | \dots \eta_s) \langle \mathcal{V}_{\text{NS-NS}} \rangle_{\mathbb{RP}^2}$$

$$\mathcal{A}_{\mathcal{V}_1 \dots \mathcal{V}_n; \mathcal{V}_{\text{NS-NS}}} = \int_{\mathcal{M}_{\ell}} [dm_1 \dots | \dots d\eta_s] \mathcal{G}_0$$

$$I_{1,\log} + I_{2,\log} = 0$$

$$I_1 \rightarrow I_1 - \int_{\mathcal{M}_{\ell}} [dm_1 \dots | \dots d\eta_s] h \mathcal{G}_1(m_1 \dots; 0 | \dots \eta_s)$$

$$= I_1 - \langle \mathcal{V}_{\text{NS-NS}} \rangle_{\text{D}} \int_{\mathcal{M}_{\ell}} [dm_1 \dots | \dots d\eta_s] h \mathcal{G}_0(m_1, \dots | \dots \eta_s)$$

$$I_2 \rightarrow I_2 - \langle \mathcal{V}_{\text{NS-NS}} \rangle_{\mathbb{RP}^2} \int_{\mathcal{M}_\ell} [dm_1 \dots | \dots d\eta_s] h \mathcal{G}_0(m_1, \dots | \dots \eta_s)$$

$$I_2 \rightarrow I_2 + \kappa \langle \mathcal{V}_{\text{NS-NS}} \rangle_{\mathbb{RP}^2} \int_{\mathcal{M}_\ell} [dm_1 \dots | \dots d\eta_s] \mathcal{G}_0(m_1, \dots | \dots \eta_s)$$

$$\mathcal{A}_{\mathcal{V}_1 \dots \mathcal{V}_n} \rightarrow \mathcal{A}_{\mathcal{V}_1 \dots \mathcal{V}_n} + \kappa \langle \mathcal{V}_{\text{NS-NS}} \rangle_{\mathbb{RP}^2} \mathcal{A}_{\mathcal{V}_1 \dots \mathcal{V}_n; \mathcal{V}_{\text{NS-NS}}}$$

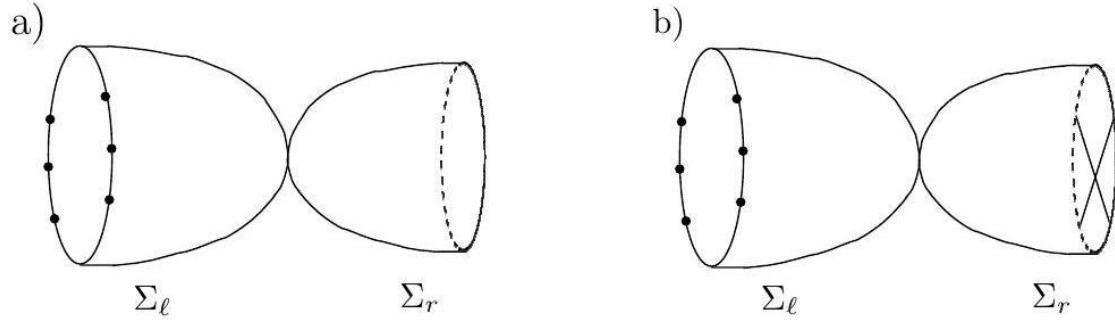


Figura 5. Supermembranas.

$$\langle \delta(\gamma(z_1)) \delta(\beta(z_2)) \rangle \neq 0$$

$$\langle \delta'(\gamma(z_1)) \delta(\beta(z_2)) \rangle = 0$$

$$\langle \delta'(\gamma(z_1)) \delta(\beta(z_2)) \gamma(z_3) \rangle \neq 0$$

$$\langle b(z_1) c(z_2) \rangle \neq 0.$$

$$\langle bc \delta(\beta) \delta(\gamma) \rangle \neq 0$$

$$\langle bc \delta(\beta) \delta'(\gamma) \gamma \rangle \neq 0$$

$$\prod_{i=2}^n \oint_{\partial_\ell \Sigma} [dz | d\theta] \varepsilon_i^{(i)} D_\theta X^i \exp(ik^{(i)} \cdot X)$$

$$\mathcal{V}_1 = c \delta(\gamma) \varepsilon^{(1)} \cdot D_\theta X \exp(ik^{(1)} \cdot X)$$

$$(b_0 + \tilde{b}_0) \int_{s_0}^{\infty} ds \exp(-s(L_0 + \tilde{L}_0)) \delta(\beta_0 + \tilde{\beta}_0) (G_0 + \tilde{G}_0)$$

$$c \delta(\gamma) \varepsilon^{(1)} \cdot D_\theta X \exp(ik^{(1)} \cdot X) (b_0 + \tilde{b}_0) \int_{s_0}^{\infty} ds \exp(-s(L_0 + \tilde{L}_0)) \delta(\beta_0 + \tilde{\beta}_0) (G_0 + \tilde{G}_0)$$

$$G_0 = G_0^X + G_0^{\text{gh}}$$

$$\tilde{G}_0 = \tilde{G}_0^X + \tilde{G}_0^{\text{gh}}$$

$$\mathcal{W}_1 = c \delta'(\gamma) \exp(ik^{(1)} \cdot X)$$

$$c\delta'(\gamma)\exp\left(ik^{(1)}\cdot X\right)\exp\left(-s(L_0+\tilde{L}_0)\right)\delta(\beta_0+\tilde{\beta}_0)(G_0+\tilde{G}_0)$$

$$I_{\beta\gamma}=\frac{1}{2\pi}\int_{\Sigma}d^2z\beta\partial_{\bar{z}}\gamma$$

$$I_{\beta^*\gamma^*}=\frac{1}{2\pi}\int_{\Sigma}d^2z\beta^*\partial_{\bar{z}}\gamma^*$$

$$\int \mathcal{D}\beta^*\mathcal{D}\gamma^*\exp\left(-I_{\beta^*\gamma^*}\right)=\det M$$

$$\int \mathcal{D}\beta\mathcal{D}\gamma\exp\left(-I_{\beta\gamma}\right)=\frac{1}{\det M}$$

$$\langle \mathcal{O} \rangle_N = \frac{\langle \mathcal{O} \rangle}{\langle 1 \rangle}$$

$$\int \mathrm{d}^n x \exp\left(-\frac{1}{2}(x,Nx)\right)=\frac{1}{\sqrt{\det N}}$$

$$N=\left(\begin{array}{cc}0&M\\M^t&0\end{array}\right)$$

$$\sqrt{\det N}=\det M,$$

$$\tilde{N}=\left(\begin{array}{cc}0&M\\-M^t&0\end{array}\right)$$

$$\langle \gamma(u)\beta(w)\rangle_N=\langle \gamma^*(u)\beta^*(w)\rangle_N=S(u,w),$$

$$\frac{\partial_{\bar{z}}}{2\pi}S(z,z')=\delta^2(z,z')$$

$$\langle \gamma^*(u_1)\gamma^*(u_2)...\gamma^*(u_s)\beta^*(w_s)\beta^*(w_{s-1})...\beta^*(w_1)\rangle_N=\sum_{\pi}(-1)^{\pi}\prod_{i=1}^sS(u_i,w_{\pi(i)}),$$

$$\left\langle \prod_{i=1}^s\gamma(u_i)\prod_{j=1}^s\beta(w_j)\right\rangle_N=\sum_{\pi}\prod_{i=1}^sS(u_i,w_{\pi(i)}).$$

$$\langle \delta(\gamma^*(u))\delta(\beta^*(w))\rangle=(\det M)S(u,w)$$

$$\int \mathrm{d}\tau^*\,\mathrm{d}\sigma^*\exp\left(-\sigma^*\gamma^*(u)-\beta^*(w)\tau^*\right)=\gamma^*(u)\beta^*(w)=\delta(\gamma^*(u))\delta(\beta^*(w))$$

$$\langle \delta(\gamma^*(u))\delta(\beta^*(w))\rangle=\int \mathcal{D}\beta^*\mathcal{D}\gamma^*\,\mathrm{d}\tau^*\,\mathrm{d}\sigma^*\exp\left(-I_{\beta^*\gamma^*}-\sigma^*\gamma^*(u)-\beta^*(w)\tau^*\right)$$

$$\hat{I}_{\hat{\beta}^*\hat{\gamma}^*}=I_{\beta^*\gamma^*}+\sigma^*\gamma^*(u)+\beta^*(w)\tau^*$$

$$\hat{\beta}^*=(\beta^*(z)\sigma^*),\hat{\gamma}^*=\begin{pmatrix}\gamma^*(z)\\ \tau^*\end{pmatrix}$$



$$\langle \delta(\gamma^*(u))\delta(\beta^*(w)) \rangle = \det \widehat{M}$$

$$\det \widehat{M} = (\det M)S(u,w)$$

$$\delta(\gamma(u))\delta(\beta(w)) = \int \mathrm{d}\tau \, \mathrm{d}\sigma \exp \left(-\sigma \gamma(u) - \beta(w)\tau \right)$$

$$\hat{I}_{\hat{\beta}\hat{\gamma}} = I_{\beta\gamma} + \sigma\gamma(u) + \beta(w)\tau,$$

$$\langle \delta(\gamma(u))\delta(\beta(w)) \rangle = \int \mathcal{D}\hat{\beta}\mathcal{D}\hat{\gamma} \exp \left(-\hat{I}_{\hat{\beta}\hat{\gamma}} \right) = \frac{1}{\det \widehat{M}}$$

$$\langle \delta(\gamma(u))\delta(\beta(w)) \rangle_N = \frac{1}{S(u,w)}$$

$$\delta(\gamma(u))\delta(\beta(w)) \sim u-w$$

$$\gamma^*(u)=0$$

$$\frac{\partial_{\bar{z}}\gamma^*}{2\pi}+\delta^2(z,w)\tau^*=0$$

$$\int \mathcal{D}\beta^*\mathcal{D}\gamma^*\prod_{i=1}^s\mathrm{d}\tau_i^*\mathrm{d}\sigma_i^*\exp\left(-I_{\beta^*\gamma^*}-\sum_i\sigma_i^*\gamma^*(u_i)-\sum_j\beta^*(w_j)\tau_j^*\right)$$

$$\left\langle \prod_{i=1}^s \delta(\gamma(u_i)) \prod_{j=1}^s \delta(\beta(w_j)) \right\rangle_N = \frac{1}{\sum_{\pi} (-1)^{\pi} \prod_{i=1}^s S(u_i, w_{\pi(i)})}$$

$$\delta(\gamma(u_1))\delta(\gamma(u_2))\sim -\frac{1}{u_1-u_2}\delta(\gamma)\delta(\partial\gamma)(u_2)$$

$$\Theta_{-t}=\delta(\gamma)\delta(\partial\gamma)\ldots\delta(\partial^{t-1}\gamma)$$

$$\delta(\gamma(u_1))\delta(\beta)\delta(\partial\beta)(u_2)\sim (u_1-u_2)^2\delta(\beta(u_2))$$

$$\langle \delta(\gamma^*(u))\delta(\beta^*(w))\gamma^*(u')\beta^*(w') \rangle = \det M \big(S(u,w)S(u',w') - S(u,w')S(u',w) \big)$$

$$\langle \delta(\gamma^*(u))\delta(\beta^*(w))\gamma^*(u')\beta^*(w') \rangle = \int \mathcal{D}\hat{\beta}^*\mathcal{D}\hat{\gamma}^* \exp \left(-\hat{I}_{\hat{\beta}^*\hat{\gamma}^*} \right) \gamma^*(u')\beta^*(w')$$

$$\langle \delta(\gamma^*(u))\delta(\beta^*(w))\gamma^*(u')\beta^*(w') \rangle = (\det \widehat{M})\hat{S}(u',w'),$$

$$\hat{S}(u',w') = \frac{1}{S(u,w)} \big(S(u,w)S(u',w') - S(u,w')S(u',w) \big).$$

$$\langle \delta(\gamma(u))\delta(\beta(w))\gamma(u')\beta(w') \rangle = \frac{1}{\det \widehat{M}} \hat{S}(u',w').$$

$$\langle \delta(\gamma(u))\delta(\beta(w))\gamma(u')\beta(w') \rangle_N = \frac{1}{S(u,w)^2} \big(S(u,w)S(u',w') - S(u,w')S(u',w) \big).$$



$$\gamma(u')\delta(\gamma(u)) \sim (u' - u)\partial\gamma \cdot \delta(\gamma)(u).$$

$$\beta(w')\delta(\gamma(u)) \sim \frac{1}{w' - u} \delta'(\gamma(u)).$$

$$\gamma(u')\delta'(\gamma(u)) \sim -\delta(\gamma(u))$$

$$\begin{aligned}\gamma(u')\delta(\beta(w)) &\sim \frac{1}{u' - w} \delta'(\beta(w)) \\ \beta(w')\delta(\beta(w)) &\sim (w' - w)\partial\beta \cdot \delta(\beta)(w)\end{aligned}$$

$$\langle \gamma(u)\beta(w)\gamma(u')\beta(w') \rangle = S(u,w)S(u',w') + S(u,w')S(u',w)$$

$$\delta(\gamma(u))\delta(\gamma(0)) = \delta\left(\gamma(0) + u\gamma'(0) + \frac{u^2}{2}\gamma''(0) + \dots\right)\delta(\gamma(0))$$

$$\delta\left(\gamma(0) + u\gamma'(0) + \frac{u^2}{2}\gamma''(0) + \dots\right)\delta(\gamma(0)) = \delta\left(u\gamma'(0) + \frac{u^2}{2}\gamma''(0) + \dots\right)\delta(\gamma(0)) = \frac{1}{u}\delta\left(\gamma'(0) + \frac{u}{2}\gamma''(0) + \dots\right)\delta(\gamma(0))$$

$$\delta\left(\gamma'(0) + \frac{u}{2}\gamma''(0) + \dots\right) = \delta(\gamma'(0)) + \frac{u}{2}\gamma''(0)\delta'(\gamma'(0)) + \dots$$

$$\delta(\gamma(u))\delta(\gamma(0)) \sim \frac{1}{u}\delta(\gamma')\delta(\gamma)(0) + \frac{1}{2}\gamma''\delta'(\gamma')\delta(\gamma)(0) + \dots$$

$$\gamma_f = \int_{\Sigma} f(z)\gamma(z), \beta_g = \int_{\Sigma} g(z)\beta(z)$$

$$\langle \delta(\gamma_f^*)\delta(\beta_g^*) \rangle_N = \langle \gamma_f^*\beta_g^* \rangle_N = \int_{\Sigma \times \Sigma} f(z)S(z,z')g(z')$$

$$\langle \delta(\gamma_f)\delta(\beta_g) \rangle_N = \frac{1}{\int_{\Sigma \times \Sigma} f(z)S(z,z')g(z')}$$

$$\langle \delta(\gamma(u))\delta(\partial\beta(w)) \rangle_N = \frac{1}{\partial_w S(u,w)}$$

$$H^1(\Sigma,K\otimes\mathcal{L}^{-1})=0$$

$$M^t\mathbf{y}_i(z)=0, i=1,\ldots,t$$

$$\langle \beta^*(w_1) \ldots \beta^*(w_t) \rangle = \det' M \det N.$$

$$\langle \delta(\beta^*(w_1)) \ldots \delta(\beta^*(w_t)) \rangle = \det' M \det N.$$

$$\langle \delta(\beta(w_1)) \ldots \delta(\beta(w_t)) \rangle = \frac{1}{\det' M} \frac{1}{\det N}$$

$$\delta(\beta^*(w_1)) \ldots \delta(\beta^*(w_t)) = \int \mathrm{d}\tau_1^* \ldots \mathrm{d}\tau_t^* \exp\left(-\sum_i \beta(w_i)\tau_i^*\right)$$



$$\hat{I}_{\beta^* \hat{\gamma}^*} = I_{\beta^* \gamma^*} + \sum_{i=1}^t \beta^*(w_i) \tau_i^*$$

$$\langle \delta(\beta^*(w_1)) \dots \delta(\beta^*(w_t)) \rangle = \int \mathcal{D}\beta^* \mathcal{D}\hat{\gamma}^* \exp(-\hat{I}_{\beta^* \hat{\gamma}^*})$$

$$\hat{I}_{\beta \hat{\gamma}} = I_{\beta \gamma} + \sum_{i=1}^t \beta(w_i) \tau_i,$$

$$\langle \delta(\beta(w_1)) \dots \delta(\beta(w_t)) \rangle = \int \mathcal{D}\beta \mathcal{D}\hat{\gamma} \exp(-\hat{I}_{\beta \hat{\gamma}})$$

$$\langle \gamma^*(u) \beta^*(w_1) \dots \beta^*(w_{t+1}) \rangle.$$

$$S(u, w) \rightarrow S(u, w) + \sum_i h_i(u) y_i(w)$$

$$N_{(t+1)} = \begin{pmatrix} y_1(w_1) & y_1(w_2) & \dots & y_1(w_{t+1}) \\ y_2(w_1) & y_2(w_2) & \dots & y_2(w_{t+1}) \\ & & \ddots & \\ y_t(w_1) & y_t(w_2) & \dots & y_t(w_{t+1}) \\ S(u, w_1) & S(u, w_2) & \dots & S(u, w_{t+1}) \end{pmatrix}$$

$$\langle \gamma^*(u) \beta^*(w_1) \dots \beta^*(w_{t+1}) \rangle = \det' M \det N_{(t+1)}$$

$$\langle \delta(\gamma^*(u)) \delta(\beta^*(w_1)) \dots \delta(\beta^*(w_{t+1})) \rangle = \det' M \det N_{(t+1)}$$

$$\langle \delta(\gamma(u)) \delta(\beta(w_1)) \dots \delta(\beta(w_{t+1})) \rangle = \frac{1}{\det' M} \frac{1}{\det N_{(t+1)}}$$

$$\frac{\langle \gamma^*(u) \delta(\beta^*(w_1)) \dots \delta(\beta^*(w_t)) \beta^*(w_{t+1}) \rangle}{\langle \delta(\beta^*(w_1)) \dots \delta(\beta^*(w_t)) \rangle} = \frac{\det N_{(t+1)}}{\det N_{(t)}}$$

$$\frac{\langle \gamma(u) \delta(\beta(w_1)) \dots \delta(\beta(w_t)) \beta(w_{t+1}) \rangle}{\langle \delta(\beta(w_1)) \dots \delta(\beta(w_t)) \rangle} = \frac{\det N_{(t+1)}}{\det N_{(t)}}$$

$$\langle \gamma(u) \delta(\beta(w_1)) \dots \delta(\beta(w_t)) \beta(w_{t+1}) \rangle = \frac{\det N_{(t+1)}}{\det' M (\det N_{(t)})^2}$$

$$N_{(t+s)} = \begin{pmatrix} y_1(w_1) & y_1(w_2) & \dots & y_1(w_{t+s}) \\ y_2(w_1) & y_2(w_2) & \dots & y_2(w_{t+s}) \\ & & \ddots & \\ y_t(w_1) & y_t(w_2) & \dots & y_t(w_{t+s}) \\ S(u_1, w_1) & S(u_1, w_2) & \dots & S(u_1, w_{t+s}) \\ & & \ddots & \\ S(u_s, w_1) & S(u_s, w_2) & \dots & S(u_s, w_{t+s}) \end{pmatrix}$$

$$\langle \gamma^*(u_1) \dots \gamma^*(u_s) \beta^*(u_1) \dots \beta^*(u_{t+s}) \rangle = \det' M \det N_{(t+s)}$$



$$\langle \delta(\gamma(u_1)) \dots \delta(\gamma^*(u_s)) \delta(\beta^*(u_1)) \dots \delta(\beta^*(u_{t+s})) \rangle = \frac{1}{\det' M} \frac{1}{\det N_{(t+s)}}$$

$$\langle \delta(\gamma(u_1)) \delta(\gamma(u_2)) \dots \delta(\gamma(u_s)) \gamma(u_{s+1}) \delta(\beta(w_1)) \dots \delta(\beta(w_{t+s})) \beta(w_{t+s+1}) \rangle = \frac{\det N_{(t+s+1)}}{\det' M (\det N_{(t+s)})^2}$$

$$\frac{\langle \delta(\gamma(u_1)) \delta(\gamma(u_2)) \dots \delta(\gamma(u_s)) \gamma(u_{s+1}) \delta(\beta(w_1)) \dots \delta(\beta(w_{t+s})) \beta(w_{t+s+1}) \rangle}{\langle \delta(\gamma(u_1)) \delta(\gamma(u_2)) \dots \delta(\gamma(u_s)) \delta(\beta(w_1)) \dots \delta(\beta(w_{t+s})) \rangle} = \frac{\det N_{(t+s+1)}}{\det N_{(t+s)}}$$

$$\mathcal{L}^2 \cong K^{-1} \otimes_{i=1}^{n_{\mathrm{R}}} \mathcal{O}(-p_i)$$

$$\Theta_{-t}=\delta(\gamma)\delta(\partial\gamma)\ldots\delta(\partial^{t-1}\gamma)$$

$$\mathcal{L}^2 \cong K^{-1} \otimes \mathcal{O}(-p)$$

$$\hat{\gamma}(z)(z^{-1}\partial_z)^{1/2}$$

$$s=\frac{1}{z^{1/2}}s^\circ$$

$$\gamma^\circ(z)(\partial_z)^{1/2}.$$

$$\langle \gamma(u) \beta(w) \rangle_{N,\delta} = \frac{\det N_{(t+s+1)}}{\det N_{(t+s)}}$$

$$\tilde{\mathcal{L}} = \mathcal{L} \otimes_{i=1}^s \mathcal{O}(-u_i) \otimes_{j=1}^{t+s} \mathcal{O}(w_j)$$

$$V_a=\sum_{i=1}^{n-m}v_{a,i}\frac{\partial}{\partial f_i}$$

$$\mathbf{i}_{V_a} = \sum_{i=1}^{n-m} v_{a,i} \frac{\partial}{\partial \operatorname{d} f_i}.$$

$$V=\sum_{n=1}^{\infty}L_{-n}U_n$$

$$b_n\mathcal{W}=0, n\geq 0,$$

$$V=L_{-1}\Phi_0,$$

$$\mathcal{W}=\Phi_0.$$

$$V=\left(L_{-2}+\frac{3}{2}L_{-1}^2\right)\Phi_{-1},$$

$$\mathcal{W}=bc\Phi_{-1}+\frac{3}{2}L_{-1}\Phi_{-1}.$$

$$V=L_{-1}\Phi_0+\left(L_{-2}+\frac{3}{2}L_{-1}^2\right)\Phi_{-1}$$



$$\mathcal{W} = \Phi_0 + \left(bc + \frac{3}{2} L_{-1} \right) \Phi_{-1}$$

$$Q_B = Q_{>} + Q_0$$

$$Q_B \mathcal{W} = \mathcal{V},$$

$$N_*^{\text{lc}} = \sum_{m \geq 1} \frac{1}{m} \alpha_{-m}^+ \alpha_m^-$$

$$Q_{>} = Q_{>,1} + Q_{>,0} + Q_{>,-1}$$

$$Q_{>,1} = -(2\alpha')^{1/2} k^+ \sum_{m \geq 1} \alpha_m^- c_{-m}$$

$$R = \frac{1}{(2\alpha')^{1/2} k^+} \sum_{m \geq 1} \alpha_{-m}^+ b_m$$

$$S = \{Q_{>,1}, R\} = \sum_{m=1}^{\infty} (m c_{-m} b_m - \alpha_{-m}^+ \alpha_m^-)$$

$$V = \sum_{n>0} L_{-n}^X W_n + \sum_{r>0} G_{-r}^X \Lambda_r$$

$$V = G_{-1/2}^X \Phi_0$$

$$\mathcal{W} = c \delta'(\gamma) \Phi_0$$

$$\beta_r|-1\rangle = \gamma_r|-1\rangle = 0, r>0$$

$$\mathcal{W} = -c_1 \beta_{-1/2} |-1\rangle \otimes \Phi_0$$

$$V = (G_{-3/2}^X + 2G_{-1/2}^X L_{-1}^X) \Phi_{-1}$$

$$\mathcal{W} = (\delta(\gamma) G_{-1/2}^X - c\beta\delta(\gamma) + c\delta'(\gamma) L_{-1}^X) \Phi_{-1}$$

$$\mathcal{W} = (G_{-1/2}^X - c_1 \beta_{-3/2} - c_1 \beta_{-1/2} L_{-1}^X) |-1\rangle \otimes \Phi_{-1}$$

$$V = \left(L_{-1}^X - \frac{1}{2} G_{-1}^X G_0^X \right) \Phi$$

$$\begin{aligned} \beta_n \Theta_{-1/2} &= 0, & n \geq 0 \\ \gamma_n \Theta_{-1/2} &= 0, & n > 0 \end{aligned}$$

$$\mathcal{W} = \left(1 - G_0^{\text{gh}} G_0^X \right) \Theta_{-1/2} \Phi$$

$$\mathcal{W} = \left(1 + \frac{1}{2} c_1 \beta_{-1} G_0^X \right) \Theta_{-1/2} \Phi$$

$$Q_B = Q_{>} + Q_0,$$



$$N_*^{\text{lc}} = \sum_{m \geq 1} \frac{1}{m} \alpha_{-m}^+ \alpha_m^- - \sum_{r \geq 1/2} \psi_{-r}^+ \psi_r^-$$

$$Q_{>} = Q_{>,1} + Q_{>,0} + Q_{>,-1}$$

$$Q_{>,1} = -(2\alpha')^{1/2} k^+ \sum_{m \geq 1} \alpha_m^- c_{-m} + (2\alpha')^{1/2} k^+ \sum_{r \geq 1/2} \gamma_{-r} \psi_r^-$$

$$\prod_{s \geq 1} \beta_{-s}^{n_s} \prod_{r \geq 0} \gamma_{-r}^{m_r} \Theta_{-1/2}$$

$$N_*^{\text{lc}} = \sum_{m \geq 1} \frac{1}{m} \alpha_{-m}^+ \alpha_m^- - \sum_{r \geq 0} \psi_{-r}^+ \psi_r^-$$

$$Q_{>,1} = -(2\alpha')^{1/2} k^+ \sum_{m \geq 1} \alpha_m^- c_{-m} + (2\alpha')^{1/2} k^+ \sum_{r \geq 0} \gamma_{-r} \psi_r^-$$

$$\begin{array}{l} z \cong z+1 \\ \theta \cong -\theta \end{array}$$

$$\begin{array}{l} z \cong z+\tau \\ \theta \cong \theta \end{array}$$

$$\tilde{z} \cong \tilde{z} + 1 \cong \tilde{z} + \tilde{\tau}$$

$$q_{\rm NS} = u_1 - u_2 - \zeta_1 \zeta_2$$

$$\begin{array}{l} u \cong u+1 \\ \zeta_1 \cong -\zeta_1 \\ \zeta_2 \cong \zeta_2 \\ \tilde{u} \cong \tilde{u}+1 \end{array}$$

$$\begin{array}{l} u \cong u+\tau \\ \zeta_1 \cong \zeta_1 \\ \zeta_2 \cong \zeta_2 \\ \tilde{u} \cong \tilde{u}+\bar{\tau} \end{array}$$

$$\tilde{q}=\tilde{u}$$

$$q_{\rm NS} = u - \zeta_1 \zeta_2$$

$$\overline{\tilde{u}} = u + \zeta_1 \zeta_2 h(\tilde{u}, u)$$

$$h(\tilde{u}+1,u+1)=-h(\tilde{u},u), h(\tilde{u}+\bar{\tau},u+\tau)=h(\tilde{u},u)$$

$$h(0,0)=-1$$

$$h_\lambda=\lambda h_1+(1-\lambda)h_2, 0\leq \lambda\leq 1$$

$$\Omega=[\mathrm{d}\tilde{u};\mathrm{d}u\mid \mathrm{d}\zeta_1\ \mathrm{d}\zeta_2]P(\tilde{u}),$$

$$P(\tilde{u}+1)=-P(\tilde{u}), P(\tilde{u}+\bar{\tau})=P(\tilde{u})$$



$$P(\tilde{u}) = \sum_{n,m \in \mathbb{Z}} \frac{(-1)^n}{\tilde{u} + n + m\bar{\tau}}$$

$$I = \int_{\Gamma} \Omega$$

$$u = \bar{\tilde{u}} - \zeta_1 \zeta_2 h(\tilde{u}, \bar{\tilde{u}})$$

$$[d\tilde{u}; du \mid d\zeta_1 \, d\zeta_2] = \left(1 - \zeta_1 \zeta_2 \frac{\partial h}{\partial \bar{\tilde{u}}}\right) [d\tilde{u} \, d\bar{\tilde{u}} \mid d\zeta_1 \, d\zeta_2]$$

$$I = \int_{\Gamma} [d\tilde{u} \, d\bar{\tilde{u}} \mid d\zeta_1 \, d\zeta_2] \left(1 - \zeta_1 \zeta_2 \frac{\partial h}{\partial \bar{\tilde{u}}}\right) P(\tilde{u})$$

$$I = - \int_{\Sigma_{\text{red}}} d\tilde{u} \wedge d\bar{\tilde{u}} \frac{\partial h}{\partial \bar{\tilde{u}}} P(\tilde{u})$$

$$I=4\pi i$$

$$\begin{aligned} z &\rightarrow -z \\ \theta &\rightarrow \pm\sqrt{-1}\theta. \end{aligned}$$

$$\begin{aligned} u &\rightarrow -u \\ \zeta_i &\rightarrow \pm\sqrt{-1}\zeta_i, i=1,2 \end{aligned}$$

$$\begin{aligned} u &\rightarrow u \\ \zeta_i &\rightarrow -\zeta_i, i=1,2 \end{aligned}$$

CONCLUSIONES.

Se concluye entonces, que la supergravedad cuántica, es un estado de la gravedad propiamente dicha, en la que, el tejido temporo – espacial, no conserva su geometría inicial. La deformación y por ende, la torsión del espacio – tiempo cuántico, no es local, de tal suerte, que la producción de multiespacios es inminente. La supergravedad cuántica, ocurre en circunstancias extremadamente hostiles de la materia, en las que, a escala microscópica, una partícula estrella u oscura, según sea el caso, se aniquila o colapsa, provocando, inicialmente, un agujero negro masivo o supermasivo, según la emisión de radiación que se tenga por implícita, y simultáneamente, la formación de dimensiones en $\mathbb{R}^n$. En este punto, es indispensable anotar, que una partícula, sea cual sea, puede interactuar en distintas dimensiones, al mismo tiempo. En un escenario de supergravedad cuántica, la materia y la energía alcanzan densidades exponenciales, lo que las funde, transformándolas finalmente, en materia y energía oscuras, yuxtapuestas en un tejido espacio – tiempo pluridimensional y no local.



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