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GRAVEDAD CUÁNTICA RELATIVISTA

RELATIVISTIC QUANTUM GRAVITY

Manuel Ignacio Albuja Bustamante
Investigador independiente



Gravedad cuántica relativista

Manuel Ignacio Albuja Bustamante¹

ignaciomanuelaljubustamante@gmail.com

Investigador independiente

RESUMEN

La teoría cuántica de campos relativistas o curvos, esbozada por este autor en trabajos anteriores, define la gravedad, como una característica fenomenológica de las llamadas “partículas supermasivas”. Este efecto gravitacional cuántico, puede ser ordinario o extraordinario, según sea el caso. Cabe precisar, que el efecto gravitacional ordinario, es el resultado de la deformación del espacio – tiempo cuántico, sin causar pliegues de campo, es decir, distintas capas geométricas y en D – dimensiones, conservándose la misma dimensión en $\mathbb{R}^4$, en tanto que, el efecto gravitacional extraordinario, es el resultado de la deformación con ruptura de simetría del espacio – tiempo cuántico, a razón de la existencia de un agujero negro cuántico en dimensión D , es decir, en multidimensiones, por lo que, el tejido del espacio – tiempo se hipergeometrizo, en una superficie de Riemann y más concretamente, en un espacio superplanckiano. Sin embargo, en este artículo, nos ocuparemos específicamente del efecto gravitacional ordinario, el mismo que puede ser endógeno o exógeno, según sea el caso, es decir, causado por la propia partícula supermasiva a propósito de su masa o por la interacción de una partícula con el gravitón, esto es, por la permeabilización de un espacio – tiempo cuántico por un campo gravitónico de gauge y local.

Palabras clave: gravedad cuántica, espacio de Hilbert – Einstein, Dimensión en $\mathbb{R}^4$, efecto gravitacional ordinario.

¹ Autor principal

Correspondencia: ignaciomanuelaljubustamante@gmail.com



Relativistic quantum gravity

ABSTRACT

The quantum theory of relativistic or curved fields, outlined by this author in previous works, defines gravity as a phenomenological characteristic of the so-called "supermassive particles". This quantum gravitational effect can be ordinary or extraordinary, as the case may be. It should be noted that the ordinary gravitational effect is the result of the deformation of quantum space-time, without causing field folds, that is, different geometric and D-dimensional layers, preserving the same dimension in $\mathbb{R}^4$, while the extraordinary gravitational effect is the result of the deformation with symmetry breaking of quantum space-time, due to the existence of a quantum black hole in dimension D, that is, in multidimensions, so that the fabric of space-time is hypergeometrized, on a Riemann surface and more specifically, in a superplanckian space. However, in this article, we will deal specifically with the ordinary gravitational effect, which can be endogenous or exogenous, as the case may be, that is, caused by the supermassive particle itself in terms of its mass or by the interaction of a particle with the graviton, that is, by the permeabilization of a quantum space-time by a gauge and local gravitonic field.

Keywords: quantum gravity, Hilbert-Einstein space, Dimension in, ordinary gravitational effect. $\mathbb{R}^4$

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INTRODUCCIÓN

En este trabajo, nos proponemos demostrar el comportamiento de un espacio – tiempo cuántico por efecto gravitacional ordinario. Para estos efectos, consideramos el propagador de una partícula supermasiva, el mismo que se desplaza en una superficie de Riemann y por ende, en un espacio de Hilbert – Einstein, en el que, la partícula supermasiva, a propósito de la gravedad causada, deforma el espacio – tiempo cuántico, provocando un rompimiento de simetría de gauge susceptible de reparación, afectando por ende, las trayectorias inicial y final de las partículas repercutidas, su momento angular y las líneas geométricas de colisión y aniquilación. Asimismo, se considera la deformación del espacio – tiempo cuántico, cuando la energía – stress – momentum de una partícula supermasiva, curva el espacio – tiempo cuántico, esto a propósito de la brecha de masa, que supera el estado de vacío, es decir, cuando el salto de energía es superior a cero. Nos ocuparemos también de las hiperpartículas, esto es, de aquellas partículas que aunque tengan o no tengan masa, deforman el espacio – tiempo cuántico, específicamente cuando superan la velocidad de la luz, liberando cargas de energía infinitas. En este trabajo, no nos ocuparemos de la partícula oscura ni de la partícula estrella respectivamente, pues aquellas, son propias de la supergravedad cuántica.

En este trabajo, se intenta extrapolar las ecuaciones de campo de Einstein a un espacio – tiempo cuántico específico, en dimensión $\mathbb{R}^4$.

RESULTADOS Y DISCUSIÓN.

En este apartado, pasamos a diseñar el modelo matemático de gravedad cuántica relativo a un espacio – tiempo cuántico relativista o curvo.



Cálculos preliminares de un campo de gauge y sus transformaciones e invariancias, respecto de un espacio planckiano.

$$X^\mu(\tau, \sigma) \sim x^\mu + p^\mu \tau + \frac{i}{\sqrt{2}} \sum_{n \in \mathbb{Z}^*} \frac{1}{n} (\alpha_n^\mu e^{-in(\tau-\sigma)} + \bar{\alpha}_n^\mu e^{-in(\tau+\sigma)}),$$

$$[x^\mu, p^\nu] = i\eta^{\mu\nu}$$

$$[\alpha_m^\mu, \alpha_n^\nu] = m\eta^{\mu\nu} \delta_{m+n,0}$$

$$H_{\text{closed}} = -\frac{m^2}{2} + N + \bar{N} - 2$$

$$H_{\text{open}} = -m^2 + N - 1$$

$$N = \sum_{n \in \mathbb{N}} n N_n, \quad N_n = \frac{1}{n} \alpha_{-n} \cdot \alpha_n,$$

$$\bar{N} = \sum_{n \in \mathbb{N}} n \bar{N}_n, \quad \bar{N}_n = \frac{1}{n} \bar{\alpha}_{-n} \cdot \bar{\alpha}_n.$$

$$H|\psi\rangle = 0$$

$$(N - \bar{N})|\psi\rangle = 0$$

$$p^\mu |k\rangle = k^\mu |k\rangle, \forall n > 0: \alpha_n^\mu |k\rangle = 0$$

$$|\psi\rangle = \prod_{n>0} \prod_{\mu=0}^{D-1} (\alpha_{-n}^\mu)^{N_{n,\mu}} |k\rangle,$$

$$\text{closed} : m^2 = -4, \text{ open} : m^2 = -1$$

$$\alpha_{-1}^\mu |k\rangle$$

$$m^2 = 0$$

$$|A\rangle = \int d^D k A_\mu(k) \alpha_{-1}^\mu |k\rangle$$

$$k^2 A_\mu = 0$$

$$k^\mu A_\mu = 0$$

$$A_\mu \longrightarrow A_\mu + k_\mu \lambda$$

$$\alpha_{-1}^\mu \bar{\alpha}_{-1}^\nu |k\rangle$$

$$m^2 = 0$$



$$\begin{aligned} & \left(\alpha_{-1}^{\mu} \bar{\alpha}_{-1}^{\nu} + \alpha_{-1}^{\nu} \bar{\alpha}_{-1}^{\mu} - \frac{1}{D} \eta^{\mu\nu} \alpha_{-1} \cdot \bar{\alpha}_{-1} \right) |p\rangle \\ & (\alpha_{-1}^{\mu} \bar{\alpha}_{-1}^{\nu} - \alpha_{-1}^{\nu} \bar{\alpha}_{-1}^{\mu}) |p\rangle, \frac{1}{D} \eta_{\mu\nu} \alpha_{-1}^{\mu} \bar{\alpha}_{-1}^{\nu} |p\rangle \end{aligned}$$

$$| \text{ boson } \rangle = Q | \text{ fermion } \rangle.$$

$$D(N=0)=26, D(N=1)=10$$

Superficies de Riemann y amplitudes en base a la función de Green.

$$\chi=2-2g-b$$

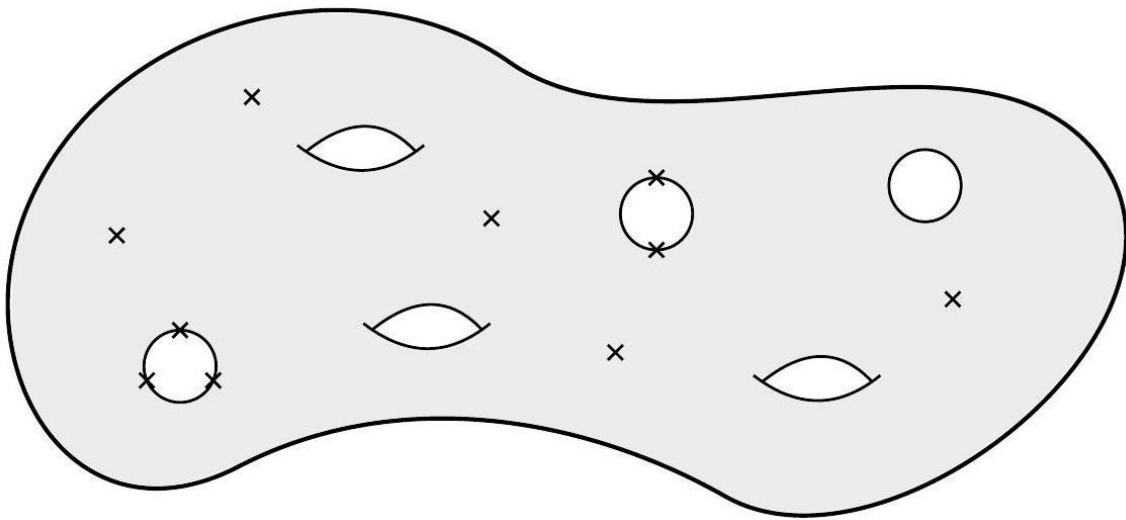


Figura 1. Modelamiento de curvaturas locales.

$$A_n(k_1,\ldots,k_n)=\sum_{g\geq 0}g_s^{n-2+2g}A_{g,n}$$

$$A_{g,n}=\int \prod_{i=1}^n \mathrm{d}^2z_i \int \mathrm{d}g_{ab} \, \mathrm{d}\Psi \mathrm{e}^{-S_{\mathrm{cft}}[g_{ab},\Psi]} \prod_{i=1}^n V_i(k_i,z_i)$$

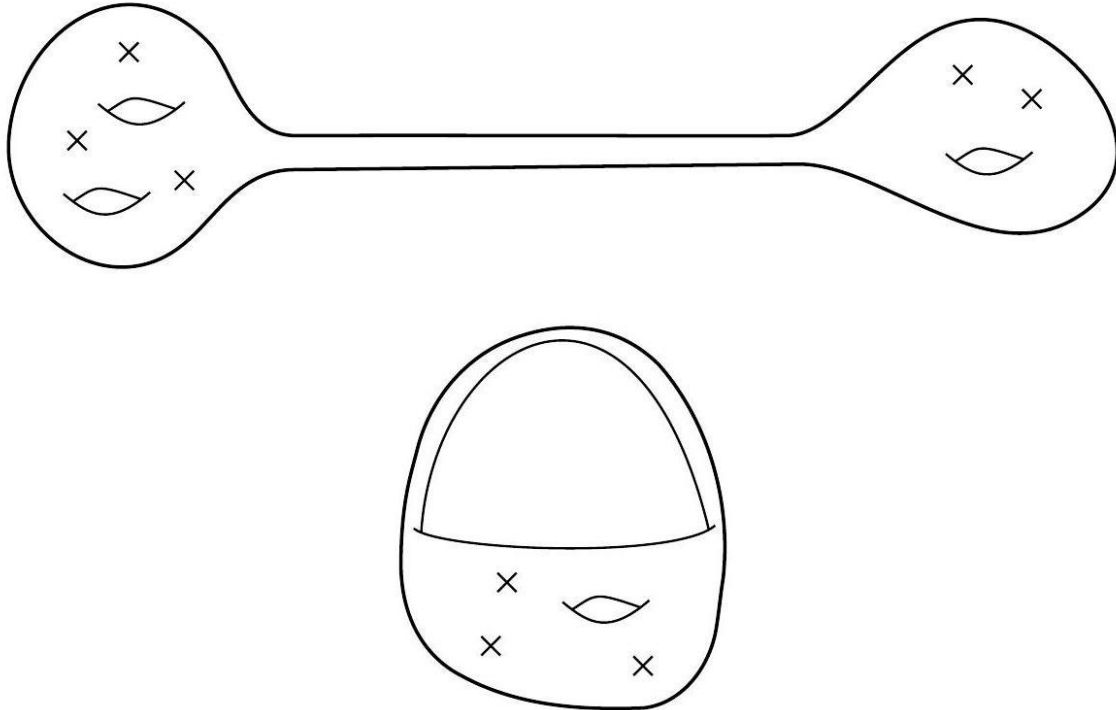
$$\dim_{\mathbb{R}}\mathcal{M}_{g,n}=6g-6+2n.$$

$$A_{g,n}=\int_{\mathcal{M}_{g,n}}\prod_{\lambda=1}^{6g-6+2n}\mathrm{d}t_{\lambda}F(t)$$

$$F(t)=\left\langle \prod_{i=1}^n V_i \times \text{ghosts} \times \text{super-ghosts} \right\rangle_{\Sigma_{g,n}}$$

Divergencias de Feynman – acción de los osciladores y propagadores de campo.

$$\frac{1}{k^2 + m^2} = \int_0^\infty ds e^{-s(k^2 + m^2)}$$



Figuras 2 y 3. Curvatura bilateral y curvatura local. La gravedad cuántica ocupa la segunda demostración, en tanto que la supergravedad cuántica, ocupa la primera demostración.

$$S_0 = \frac{1}{2} K_\Psi(\Psi, \Psi) + \frac{1}{2} K_\Phi(\Phi, \Phi)$$

$$S_{\text{int}} = \sum_{m,n} \mathcal{V}_{m,n}(\Phi^m, \Psi^n)$$

$$|\Psi\rangle = \sum_n \int \frac{d^D k}{(2\pi)^D} \psi_\alpha(k) |k, \alpha\rangle$$

$$\psi_\alpha = \{T, G_{\mu\nu}, B_{\mu\nu}, \Phi, \dots\}.$$

$$\sigma \in [0, 2\pi), \sigma \sim \sigma + 2\pi.$$

$$S_{\text{NG}}[X^\mu] = \frac{1}{2\pi\alpha'} \int d^2\sigma \sqrt{\det G_{\mu\nu}(X)} \frac{\partial X^\mu}{\partial \sigma^a} \frac{\partial X^\nu}{\partial \sigma^b}$$

$$S_{\text{P}}[g, X^\mu] = \frac{1}{4\pi\alpha'} \int d^2\sigma \sqrt{g} g^{ab} G_{\mu\nu}(X) \frac{\partial X^\mu}{\partial \sigma^a} \frac{\partial X^\nu}{\partial \sigma^b}$$

$$\chi_g := \chi(\Sigma_g) := 2 - 2g = \frac{1}{4\pi} \int_{\Sigma_g} d^2\sigma \sqrt{g} R$$

$$\sigma'^a = f^a(\sigma^b), g'(\sigma') = f^*g(\sigma), \Psi'(\sigma') = f^*\Psi(\sigma)$$

$$g'_{ab}(\sigma') = \frac{\partial \sigma^c}{\partial \sigma'^a} \frac{\partial \sigma^d}{\partial \sigma'^b} g_{cd}(\sigma), X'^\mu(\sigma') = X^\mu(\sigma)$$

$$\delta_\xi \sigma^a = \xi^a, \delta_\xi \Psi = \mathcal{L}_\xi \Psi, \delta_\xi g_{ab} = \mathcal{L}_\xi g_{ab}$$

$$\mathcal{L}_\xi g_{ab} = \xi^c \partial_c g_{ab} + g_{ac} \partial_b \xi^c + g_{bc} \partial_a \xi^c = \nabla_a \xi_b + \nabla_b \xi_a.$$

$$\Gamma_g := \pi_0 \left(\text{Diff}(\Sigma_g) \right) = \frac{\text{Diff}(\Sigma_g)}{\text{Diff}_0(\Sigma_g)}$$

$$g'_{ab}(\sigma) = e^{2\omega(\sigma)} g_{ab}(\sigma), \Psi'(\sigma) = \Psi(\sigma)$$

$$\delta_\omega g_{ab} = 2\omega g_{ab}, \delta_\omega \Psi = 0$$

$$\text{Conf}(\Sigma_g) := \frac{\text{Met}(\Sigma_g)}{\text{Weyl}(\Sigma_g)}$$

$$G := \text{Diff}(\Sigma_g) \ltimes \text{Weyl}(\Sigma_g).$$

$$G_0 := \text{Diff}_0(\Sigma_g) \ltimes \text{Weyl}(\Sigma_g).$$

$$g' = f^*(e^{2\omega}g) = e^{2f^*\omega}f^*g$$

$$g_{ab}(\sigma) = e^{2\phi(\sigma)} \hat{g}_{ab}(\sigma)$$

$$\hat{g}_{ab} = \delta_{ab}, \phi$$

$$\hat{g}_{ab} = \delta_{ab}, \phi = 0$$

$$S_m[f^*g, f^*\Psi] = S_m[g, \Psi], S_m[e^{2\omega}g, \Psi] = S_m[g, \Psi]$$

$$X^\mu(\tau, \sigma) \sim X^\mu(\tau, \sigma + 2\pi)$$

$$c_m = D + c_{\perp}.$$

$$\mathbb{M}_{\text{matter CFT parameters}}\colon D, c_m.$$

$$T_{m,ab}\colon=-\frac{4\pi}{\sqrt{g}}\frac{\delta S_m}{\delta g^{ab}}$$

$$\nabla^a T_{m,ab}=0\,\,(\text{on-shell}).$$

$$g^{ab}T_{m,ab}=0\,\,(\text{off-shell})$$

$$P^a\colon=\int\,\mathrm{d}\sigma T_m^{0a}$$

$$g^{ab}T_{m,ab}\propto R$$

$$Z_g\colon=\int\frac{\mathrm{d}_g g_{ab}}{\Omega_{\text{gauge}}\left[g\right]}Z_m[g],Z_m[g]\colon=\int\,\mathrm{d}_g\Psi\mathrm{e}^{-S_m[g,\Psi]}$$

$$(\delta\Phi_1,\delta\Phi_2)_g\colon=\int\,\mathrm{d}^2\sigma\sqrt{g}\gamma_g(\delta\Phi_1,\delta\Phi_2),|\delta\Phi|_g^2\colon=(\delta\Phi,\delta\Phi)_g$$

$$\int\,\mathrm{d}_g\delta\Phi\mathrm{e}^{-\frac{1}{2}(\delta\Phi,\delta\Phi)_g}=\frac{1}{\sqrt{\mathrm{det}\gamma_g}}$$

$$\int\,\mathrm{d}\Phi\sqrt{\mathrm{det}\gamma_g}$$

$$\int\,\mathrm{d}_g\delta\Phi\mathrm{e}^{-\frac{1}{2}(\delta\Phi,\delta\Phi)_g}=1$$

$$\Phi(\sigma)\rightarrow\Phi'(\sigma)=\Phi(\sigma)+\varepsilon(\sigma)$$

$$(\delta f,\delta f)_g\coloneqq\int\,\mathrm{d}^2\sigma\sqrt{g}\delta f^2$$

$$(\delta V^a,\delta V^a)_g\coloneqq\int\,\mathrm{d}^2\sigma\sqrt{g}g_{ab}\delta V^a\delta V^b$$

$$(\delta T_{ab},\delta T_{ab})_g\coloneqq\int\,\mathrm{d}^2\sigma\sqrt{g}G^{abcd}\delta T_{ab}\delta T_{cd}$$

$$G^{abcd}\colon=G_{\perp}^{abcd}+ug^{ab}g^{cd},G_{\perp}^{abcd}\colon=g^{ac}g^{bd}+g^{ad}g^{bc}-g^{ab}g^{cd}$$

$$G^{abcd}T_{cd}=G_{\perp}^{abcd}T_{cd}=2T_{ab},G^{abcd}(\Lambda g_{cd})=2u(\Lambda g_{ab})$$

$$\delta g_{ab}=g_{ab}\delta\Lambda+\delta g_{ab}^{\perp},\delta\Lambda=\frac{1}{2}g^{ab}\delta g_{ab},g^{ab}\delta g_{ab}^{\perp}=0$$

$$|\delta g_{ab}|_g^2=4u|\delta\Lambda|_g^2+|\delta g_{\mu\nu}^{\perp}|_g^2$$

$$u>0$$

$$\mathrm{d}_g g_{ab} = \mathrm{d}_g \Lambda \, \mathrm{d}_g g_{ab}^{\perp}$$



$$\begin{aligned}
G^{abcd}\delta g_{ab}\delta g_{cd} &= (G_{\perp}^{abcd} + u g^{ab} g^{cd})(g_{ab}\delta\Lambda + \delta g_{ab}^{\perp})(g_{cd}\delta\Lambda + \delta g_{cd}^{\perp}) \\
&= (2u g^{cd}\delta\Lambda + G_{\perp}^{abcd}\delta g_{ab}^{\perp})(g_{cd}\delta\Lambda + \delta g_{cd}^{\perp}) \\
&= 4u(\delta\Lambda)^2 + G_{\perp}^{abcd}\delta g_{ab}^{\perp}\delta g_{cd}^{\perp} \\
&= 4u\delta\Lambda^2 + 2g^{ac}g^{bd}\delta g_{ab}^{\perp}\delta g_{cd}^{\perp}
\end{aligned}$$

$$G^{abcd}=g^{ac}g^{bd}+cg^{ab}g^{cd}$$

$$(\delta X^{\mu},\delta X^{\mu})_g=\int\,\mathrm{d}^2\sigma\sqrt{g}G_{\mu\nu}(X)\delta X^{\mu}\delta X^{\nu}$$

Reparación de gauge Faddeev-Popov en superficies de Riemann – Roch. Acción de Moduli.

Métrica de David-Distler-Kawai y acción de Liouville - Faddeev-Popov.

$$(f^*g,f^*\Psi)\sim (g,\Psi), (e^{2\omega}g,\Psi)\sim (g,\Psi)$$

$$\mathrm{d}(\text{ fields }) = \text{Jacobian} \times \mathrm{d}(\text{ gauge }) \times \mathrm{d}(\text{ physical }).$$

$$Z=\int_{\mathbb{R}^2}\mathrm{d}x\,\mathrm{d}y\mathrm{e}^{-(x-y)^2}$$

$$r=x-y, y=a$$

$$Z=\int_{\mathbb{R}}\mathrm{d}a\int_0^\infty\mathrm{e}^{-r^2}=\frac{\sqrt{\pi}}{2}\mathrm{Vol}(\mathbb{R})$$

$$\mathcal{M}_g\!:=\!\frac{\mathrm{Met}(\Sigma_g)}{G}$$

$$\mathcal{T}_g\!:=\!\frac{\mathrm{Met}(\Sigma_g)}{G_0}$$

$$\mathcal{M}_g=\frac{\mathcal{T}_g}{\Gamma_g}$$

$$M_g\!:=\!\dim_{\mathbb{R}}\mathcal{M}_g=\dim_{\mathbb{R}}\mathcal{T}_g=\begin{cases}0&g=0\\2&g=1\\6g-6&g\geq2\end{cases}$$

$$\int_{\mathcal{M}_g} \mathrm{d}^{\mathbf{M}_g} t = \frac{1}{\Omega_{\Gamma_g}} \int_{\mathcal{T}_g} \mathrm{d}^{\mathbf{M}_g} t$$

$$\Omega_G\!:=\!\int_G\,\mathrm{d}g=\int_G\,\mathrm{d}(hg),$$

$$\alpha=\alpha^iT_i.$$

$$\Omega_G=\int\,\mathrm{d}\alpha\!:=\!\int\,\prod_i\,\mathrm{d}\alpha^i$$

$$g_{ab}=\hat{g}_{ab}^{(f,\phi)}(t)\!:=\!e^{2f^*\phi}f^*\hat{g}_{ab}(t)=f^*(e^{2\phi}\hat{g}_{ab}(t))$$



$$g_{ab}(\sigma) = \hat{g}_{ab}^{(f,\phi)}(\sigma; t) := e^{2\phi(\sigma)} \hat{g}'_{ab}(\sigma; t), \hat{g}'_{ab}(\sigma; t) = \frac{\partial \sigma'^c}{\partial \sigma^a} \frac{\partial \sigma'^d}{\partial \sigma^b} \hat{g}_{cd}(\sigma'; t)$$

$$\delta g_{ab} = 2\phi g_{ab} + \nabla_a \xi_b + \nabla_b \xi_a + \delta t_i \partial_i g_{ab}$$

$$\partial_i := \frac{\partial}{\partial t_i}.$$

$$\Omega_{\mathrm{Diff}_0}[g] := \int \mathrm{d}_g \xi$$

$$\Omega_{\mathrm{Weyl}}[g] := \int \mathrm{d}_g \phi$$

$$\Omega_{\mathrm{Diff}}[g] = \Omega_{\mathrm{Diff}_0}[g] \Omega_{\Gamma_g}$$

$$\Omega_{\mathrm{Diff}_0}[g] := \Omega_{\mathrm{Diff}_0}[e^{2\phi}\hat{g}] = \Omega_{\mathrm{Diff}_0}[\hat{g}], \Omega_{\mathrm{Weyl}}[g] := \Omega_{\mathrm{Weyl}}[\mathcal{L}_\xi\hat{g}] = \Omega_{\mathrm{Weyl}}[\hat{g}]$$

$$\Omega_{\mathrm{gauge}}[g] := \Omega_{\mathrm{gauge}}[\hat{g}]$$

$$\Omega_{\mathrm{Diff}_0}[e^{2\phi}\hat{g}] = \int \mathrm{d}_{e^{2\phi}\mathcal{L}_\xi\hat{g}} \xi = \int \mathrm{d}_{e^{2\phi}\hat{g}} \xi = \int \mathrm{d}_{\hat{g}} \xi = \Omega_{\mathrm{Diff}_0}[\hat{g}]$$

$$\Omega_{\mathrm{Weyl}}[\mathcal{L}_\xi\hat{g}] = \int \mathrm{d}_{e^{2\phi}\mathcal{L}_\xi\hat{g}} \phi = \int \mathrm{d}_{e^{2\phi}\hat{g}} \phi = \Omega_{\mathrm{Weyl}}[\hat{g}]$$

$$|\delta\phi|^2 = \int \mathrm{d}^2\sigma \sqrt{g} \delta\phi^2 = \int \mathrm{d}^2\sigma \sqrt{\hat{g}} e^{2\phi} \delta\phi^2$$

$$F_{ab} := g_{ab} - \hat{g}_{ab}^{(f,\phi)}(t),$$

$$\delta g_{ab} = 2\tilde{\Lambda} g_{ab} + (P_1 \xi)_{ab} + \delta t_i \mu_{iab}$$

$$(P_1 \xi)_{ab} = \nabla_a \xi_b + \nabla_b \xi_a - g_{ab} \nabla_c \xi^c,$$

$$\mu_{iab} = \partial_i g_{ab} - \frac{1}{2} g_{ab} g^{cd} \partial_i g_{cd},$$

$$\tilde{\Lambda} = \Lambda + \frac{1}{2} \delta t_i g^{ab} \partial_i g_{ab}, \Lambda = \phi + \frac{1}{2} \nabla_c \xi^c.$$

$$Z_g = \int \mathrm{d}^{\mathbf{M}_g} t \, \mathrm{d}_g \tilde{\Lambda} \, \mathrm{d}_g (P_1 \xi) \Omega_{\mathrm{gauge}}[g]^{-1} Z_m[g]$$

$$(P_1\xi,\tilde{\Lambda})\longrightarrow(\xi,\phi)$$

$$\mathrm{d}_g(P_1\xi)\mathrm{d}_g\tilde{\Lambda}\stackrel{?}{=}\mathrm{d}_g\xi\,\mathrm{d}_g\phi\Delta_{\mathrm{FP}}[g]$$

$$\Delta_{\mathrm{FP}}[g] = \det \frac{\partial (P_1 \xi, \tilde{\Lambda})}{\partial (\xi, \phi)} = \det \begin{pmatrix} P_1 & 0 \\ \star & 1 \end{pmatrix} = \det P_1$$

$$(T,P_1v)_g=(P_1^{\dagger}T,v)_g,$$

$$(P_1^{\dagger}T)_a=-2\nabla^bT_{ab}$$

$$\dim \ker P_1^\dagger - \dim \ker P_1 = -3\chi_g = 6g - 6$$

Deformaciones de Teichmüller.

$$(\delta g, P_1 \xi)_g = 0 \Rightarrow (P_1^\dagger \delta g, \xi)_g = 0$$

$$P_1^\dagger \delta g = 0$$

$$\delta g \in \ker P_1^\dagger$$

$$\ker P_1^\dagger = \text{Span}\{\phi_i\}, i = 1, \dots, \dim \ker P_1^\dagger$$

$$\dim_{\mathbb{R}} \ker P_1^\dagger = M_g = \begin{cases} 0 & g = 0 \\ 2 & g = 1 \\ 6g - 6 & g > 1 \end{cases}$$

$$\Pi := P_1 \frac{1}{P_1^\dagger P_1} P_1^\dagger$$

$$\delta t_i \mu_i = \delta t_i (1 - \Pi) \mu_i + \delta t_i \Pi \mu_i = \delta t_i (1 - \Pi) \mu_i + \delta t_i P_1 \zeta_i.$$

$$\zeta_i := \frac{1}{P_1^\dagger P_1} P_1^\dagger \mu_i.$$

$$(1 - \Pi) \mu_i = \phi_j (M^{-1})_{jk} (\phi_k, \mu_i)_g$$

$$M_{ij} := (\phi_i, \phi_j)_g.$$

$$\delta g_{ab} = (P_1 \tilde{\xi})_{ab} + 2\tilde{\Lambda} g_{ab} + Q_{iab} \delta t_i.$$

$$\tilde{\xi} = \xi + \zeta_i \delta t_i, Q_{iab} = \phi_{jab} (M^{-1})_{jk} (\phi_k, \mu_i)_g.$$

$$|\delta g|_g^2 = |\delta \tilde{\Lambda}|_g^2 + |P_1 \tilde{\xi}|_g^2 + |Q_i \delta t_i|_g^2$$

$$d_g g_{ab} = d_g \tilde{\Lambda} d_g (P_1 \tilde{\xi}) d_g (Q_i \delta t_i).$$

$$(\tilde{\xi}, \tilde{\Lambda}, Q_i \delta t_i) \rightarrow (\xi, \Lambda, \delta t_i),$$

$$d_g \tilde{\Lambda} d_g (P_1 \tilde{\xi}) d_g (Q_i \delta t_i) = d^{M_g} t d_g \Lambda d_g (P_1 \xi) \frac{\det(\phi_i, \mu_j)_g}{\sqrt{\det(\phi_i, \phi_j)_g}},$$

$$Z_g = \int_{\mathcal{T}_g} d^{M_g} t \frac{1}{\Omega_{\text{gauge}}[\hat{g}]} \int d_g \Lambda d_g (P_1 \xi) \frac{\det(\phi_i, \mu_j)_g}{\sqrt{\det(\phi_i, \phi_j)_g}} Z_m[g]$$



Vectores Conformal Killing.

$$\int d_g \xi = \Omega_{\text{Diff}_0}[\hat{g}]$$

$$\xi^{(0)} \in \mathcal{K}_g := \ker P_1$$

$$(P_1 \xi^{(0)})_{ab} = \nabla_a \xi_b^{(0)} + \nabla_b \xi_a^{(0)} - g_{ab} \nabla_c \xi^{(0)c} = 0$$

$$K_g := \dim_{\mathbb{R}} \mathcal{K}_g = \dim_{\mathbb{R}} \ker P_1 = \begin{cases} 6 & g = 0 \\ 2 & g = 1 \\ 0 & g > 1 \end{cases}$$

$$g = 0: \mathcal{K}_0 = \text{SL}(2, \mathbb{C}), g = 1: \mathcal{K}_1 = \text{U}(1) \times \text{U}(1).$$

$$\xi = \xi^{(0)} + \xi'$$

$$(\xi^{(0)}, \xi')_g = 0$$

$$(P_1 \xi, \Lambda) \rightarrow (\xi', \phi).$$

$$d_g \Lambda d_g (P_1 \xi) = d_g \phi d_g \xi' \Delta_{\text{FP}}[g]$$

$$\Delta_{\text{FP}}[g] = \det' \frac{\partial(P_1 \xi, \Lambda)}{\partial(\xi', \phi)} = \det' P_1 = \sqrt{\det' P_1 P_1^\dagger}$$

$$Z_g = \int_{\mathcal{T}_g} d^{\text{M}_g} t \Omega_{\text{gauge}}[\hat{g}]^{-1} \int d_g \phi d_g \xi' \frac{\det(\phi_i, \mu_j)_g}{\sqrt{\det(\phi_i, \phi_j)_g}} \Delta_{\text{FP}}[g] Z_m[g]$$

$$\Delta_{\text{FP}}[g] = \det' \frac{\partial(P_1 \xi, \Lambda)}{\partial(\xi', \phi)} = \det' \begin{pmatrix} P_1 & 0 \\ \frac{1}{2} \nabla & 1 \end{pmatrix} = \det' P_1$$

$$\sqrt{\det' P_1^\dagger P_1} = \det' P_1$$

$$\begin{aligned} 1 &= \int d_g \delta \Lambda d_g (P_1 \delta \xi) e^{-|\delta \Lambda|_g^2 - |P_1 \delta \xi'|_g^2} \\ &= \Delta_{\text{FP}}[g] \int d_g \delta \phi d_g \delta \xi' e^{-|\delta \phi + \frac{1}{2} \nabla_c \delta \xi^c|_g^2 - |P_1 \delta \xi'|_g^2} \\ &= \Delta_{\text{FP}}[g] \int d_g \delta \phi d_g \delta \xi' e^{-|\delta \phi|_g^2 - (\delta \xi', P_1^\dagger P_1 \delta \xi')_g} \\ &= \Delta_{\text{FP}}[g] (\det' P_1^\dagger P_1)^{-1/2} \end{aligned}$$

$$\Omega'_{\text{Diff}_0}[g] := \Omega'_{\text{Diff}_0}[\hat{g}] = \int d_g \xi'$$

$$\ker P_1 = \text{Span}\{\psi_i\}, i = 1, \dots, K_g$$



$$\xi' \rightarrow \xi$$

$$d_g \xi' = \frac{1}{\sqrt{\det(\psi_i, \psi_j)_g}} \frac{d_g \xi}{\Omega_{\text{ckv}}[g]},$$

$$\Omega_{\text{Diff}_0}[g] = \sqrt{\det(\psi_i, \psi_j)_g} \Omega_{\text{ckv}}[g] \Omega'_{\text{Diff}_0}[g]$$

$$\xi^{(0)} = \alpha_i \psi_i$$

$$\xi \rightarrow (\xi', \alpha_i).$$

$$\begin{aligned} 1 &= \int d\xi e^{-|\xi|^2_g} = J \int d\xi^{(0)} d\xi' e^{-|\xi'|^2_g - |\xi^{(0)}|^2_g} \\ &= J \int \prod_i d\alpha_i e^{-\alpha_i \alpha_j (\psi_i, \psi_j)_g} \int d\xi' e^{-|\xi'|^2_g} \\ &= J \left(\det(\psi_i, \psi_j)_g \right)^{-1/2} \end{aligned}$$

$$d\xi = \sqrt{\det(\psi_i, \psi_j)_g} d\xi' \prod_i d\alpha_i$$

$$\Omega_{\text{ckv}}[g] = \int \prod_i d\alpha_i$$

$$Z_g = \int_{\mathcal{T}_g} d^{\text{M}_g} t \Omega_{\text{gauge}}[\hat{g}]^{-1} \int d_g \phi d_g \xi \frac{\det(\phi_i, \mu_j)_g}{\sqrt{\det(\phi_i, \phi_j)_g}} \frac{\Omega_{\text{ckv}}[g]^{-1}}{\sqrt{\det(\psi_i, \psi_j)_g}} \Delta_{\text{FP}}[g] Z_m[g]$$

$$(f^*\hat{g}, f^*\phi, f^*\Psi) \rightarrow (\hat{g}, \phi, \Psi)$$

$$Z_g = \int_{\mathcal{T}_g} d^{\text{M}_g} t \frac{\Omega_{\text{Diff}_0}[\hat{g}]}{\Omega_{\text{gauge}}[\hat{g}]} \int d_g \phi \frac{\det(\phi_i, \mu_j)_g}{\sqrt{\det(\phi_i, \phi_j)_g}} \frac{\Omega_{\text{ckv}}[g]^{-1}}{\sqrt{\det(\psi_i, \psi_j)_g}} \Delta_{\text{FP}}[g] Z_m[g]$$

$$g_{ab} := g_{ab}^{(\phi)} = e^{2\phi} \hat{g}_{ab}$$

$$Z_g = \frac{1}{\Omega_{\Gamma_g}} \int_{\mathcal{T}_g} d^{\text{M}_g} t \frac{\Omega_{\text{Diff}}[\hat{g}]}{\Omega_{\text{gauge}}[\hat{g}]} \int d_g \phi \frac{\det(\phi_i, \mu_j)_g}{\sqrt{\det(\phi_i, \phi_j)_g}} \frac{\Omega_{\text{ckv}}[g]^{-1}}{\sqrt{\det(\psi_i, \psi_j)_g}} \Delta_{\text{FP}}[g] Z_m[g]$$

$$Z_g = \int_{\mathcal{M}_g} d^{\text{M}_g} t \frac{\Omega_{\text{Diff}}[\hat{g}]}{\Omega_{\text{gauge}}[\hat{g}]} \int d_g \phi \frac{\det(\phi_i, \mu_j)_g}{\sqrt{\det(\phi_i, \phi_j)_g}} \frac{\Omega_{\text{ckv}}[g]^{-1}}{\sqrt{\det(\psi_i, \psi_j)_g}} \Delta_{\text{FP}}[g] Z_m[g]$$

Transformaciones de Weyl – difeomorfismos gravitacionales y anomalías cuánticas – nodos tensoriales de campo – parámetros de Ricci.

$$\frac{\Delta_{\text{FP}}[e^{2\phi}\hat{g}]}{\sqrt{\det(\phi_i, \phi_j)_{e^{2\phi}\hat{g}}}} = e^{\frac{c_{\text{gh}}}{6}S_L[\hat{g}, \phi]} \frac{\Delta_{\text{FP}}[\hat{g}]}{\sqrt{\det(\hat{\phi}_i, \hat{\phi}_j)_{\hat{g}}}}$$

$$Z_m[e^{2\phi}\hat{g}] = e^{\frac{c_m}{6}S_L[\hat{g}, \phi]} Z_m[\hat{g}]$$

$$S_L[\hat{g}, \phi] := \frac{1}{4\pi} \int d^2\sigma \sqrt{\hat{g}} (\hat{g}^{ab} \partial_a \phi \partial_b \phi + \hat{R} \phi)$$

$$c_{\text{gh}} = -26$$

$$\det(\phi_i, \mu_j)_{e^{2\phi}\hat{g}} = \det(\hat{\phi}_i, \hat{\mu}_j)_{\hat{g}}, \det(\psi_i, \psi_j)_{e^{2\phi}\hat{g}} = \det(\psi_i, \psi_j)_{\hat{g}},$$

$$\Omega_{\text{ckv}}[e^{2\phi}\hat{g}] = \Omega_{\text{ckv}}[\hat{g}].$$

$$\langle g^{\mu\nu} T_{\mu\nu} \rangle = \frac{c}{12} R$$

$$Z_g = \int_{\mathcal{M}_g} d^{M_g} t \frac{\Omega_{\text{Diff}}[\hat{g}]}{\Omega_{\text{gauge}}[\hat{g}]} \frac{\det(\phi_i, \hat{\mu}_j)_{\hat{g}}}{\sqrt{\det(\phi_i, \phi_j)_{\hat{g}}}} \frac{\Omega_{\text{ckv}}[\hat{g}]^{-1}}{\sqrt{\det(\psi_i, \psi_j)_{\hat{g}}}} \Delta_{\text{FP}}[\hat{g}] Z_m[\hat{g}] \int d_g \phi e^{-\frac{c_L}{6} S_L[\hat{g}, \phi]}$$

$$c_L := 26 - c_m$$

$$c_L = 0 \Rightarrow c_m = 26$$

$$c_{\perp} = 26 - D$$

$$\int d_g \phi = \Omega_{\text{Weyl}}[\hat{g}]$$

$$\Omega_{\text{gauge}}[\hat{g}] = \Omega_{\text{Diff}}[\hat{g}] \times \Omega_{\text{Weyl}}[\hat{g}]$$

$$Z_g = \int_{\mathcal{M}_g} d^{M_g} t \frac{\det(\phi_i, \hat{\mu}_j)_{\hat{g}}}{\sqrt{\det(\phi_i, \phi_j)_{\hat{g}}}} \frac{\Omega_{\text{ckv}}[\hat{g}]^{-1}}{\sqrt{\det(\psi_i, \psi_j)_{\hat{g}}}} \Delta_{\text{FP}}[\hat{g}] Z_m[\hat{g}]$$

Ambigüedades, ultralocalidad y constante cosmológica – Métrica Weil-Petersson.

$$S_{\mu}[g] = \int d^2\sigma \sqrt{g}$$

$$\sqrt{\det \gamma_g} = e^{-\mu_{\gamma} S_{\mu}[g]},$$

$$\Omega_{\Phi} = \lim_{\lambda \rightarrow 0} \int d_g \Phi e^{-\lambda(\Phi, \Phi)_g}$$



$$\int d_g \Phi e^{-\lambda(\Phi, \Phi)_g} = e^{-\mu(\lambda) S_\mu[g]}$$

$$\Omega_\Phi = \int d_g \Phi = e^{-\mu(0) S_\mu[g]}$$

$$\lim_{\epsilon \rightarrow 0} \frac{1}{\epsilon} \int d^2 \sigma \sqrt{g}$$

$$\begin{aligned} Z_g &= \int_{\mathcal{M}_g} d^{\mathbf{M}_g} t \frac{\det(\phi_i, \hat{\mu}_j)_{\hat{g}}}{\sqrt{\det(\phi_i, \phi_j)_{\hat{g}}}} \frac{\Omega_{\text{ckv}}[\hat{g}]^{-1}}{\sqrt{\det(\psi_i, \psi_j)_{\hat{g}}}} \Delta_{\text{FP}}[\hat{g}] Z_m[\hat{g}] \\ &= \int_{\mathcal{M}_g} d^{\mathbf{M}_g} t \frac{\det(\phi_i, \hat{\mu}_j)_{\hat{g}}^2}{\det(\phi_i, \phi_j)_{\hat{g}}} \frac{\det' \hat{P}_1^\dagger \hat{P}_1}{\det(\psi_i, \psi_j)_{\hat{g}}} \frac{Z_m[\hat{g}]}{\Omega_{\text{ckv}}[\hat{g}]} \end{aligned}$$

$$d(\text{WP}) = \int_{\mathcal{M}_g} d^{\mathbf{M}_g} t \frac{\det(\phi_i, \hat{\mu}_j)_{\hat{g}}}{\sqrt{\det(\phi_i, \phi_j)_{\hat{g}}}}$$

$$\sigma'^a = \hat{f}^a(\sigma^b), \hat{g}'(\sigma') = f^* \hat{g}(\sigma), \phi'(\sigma') = f^* \phi(\sigma), \Psi'(\sigma') = f^* \Psi(\sigma)$$

$$g_{ab} = f^*(\mathrm{e}^{2\phi} \hat{g}_{ab}(t))$$

$$g'_{ab}(\sigma) = \mathrm{e}^{2\omega(\sigma)} g_{ab}(\sigma), \phi'(\sigma) = \phi(\sigma) - \omega(\sigma), \Psi'(\sigma) = \Psi(\sigma)$$

$$g'_{ab} = F^* g_{ab}, g'_{ab} = f'^*(\mathrm{e}^{2\phi'} \hat{g}'_{ab}), g_{ab} = f^*(\mathrm{e}^{2\phi} \hat{g}_{ab})$$

$$\hat{F} = f'^{-1} \circ F \circ f, \phi' = \hat{F}^*(\phi - \omega), \hat{g}'_{ab} = \hat{F}^*(\mathrm{e}^{2\omega} \hat{g}_{ab})$$

$$g'_{ab} = \tilde{f}^*(\mathrm{e}^{2\phi} \hat{g}_{ab})$$

$$\begin{aligned} g'_{ab} &= F^* g_{ab} = F^* \left(f^* (\mathrm{e}^{2\phi} \hat{g}_{ab}) \right) = F^* \left(f^* (\mathrm{e}^{2(\phi-\omega)} \mathrm{e}^{2\omega} \hat{g}_{ab}) \right) \\ &= f'^*(\mathrm{e}^{2\phi'} \hat{g}'_{ab}) \end{aligned}$$

Campos fantasmas. Métrica de Grassmann y Métrica de Faddeev-Popov.

$$\Delta_{\text{FP}}[g] = \int d'_g b \, d'_g c e^{-S_{\text{gh}}[g,b,c]}$$

$$\begin{aligned} S_{\text{gh}}[g,b,c] &:= \frac{1}{4\pi} \int d^2 \sigma \sqrt{g} g^{ab} g^{cd} b_{ac} (P_1 c)_{bd} \\ &= \frac{1}{4\pi} \int d^2 \sigma \sqrt{g} g^{ab} (b_{ac} \nabla_b c^c + b_{bc} \nabla_a c^c - b_{ab} \nabla_c c^c) \end{aligned}$$

$$(P_1 c)_{ab} = \nabla_a c_b + \nabla_b c_a - g_{ab} \nabla_c c^c = 0, (P_1^\dagger b)_a = -2 \nabla^b b_{ab} = 0$$

$$T_{ab}^{\text{gh}} = -b_{ac} \nabla_b c^c - b_{bc} \nabla_a c^c + c^c \nabla_c b_{ab} + g_{ab} b_{cd} \nabla^c c^d$$



$$g^{ab}T_{ab}^{\text{gh}} = 0$$

$$S_{\text{gh}}[e^{2\omega}g, b, c] = S_{\text{gh}}[g, b, c]$$

$$N_{\text{gh}}(b) = -1, N_{\text{gh}}(c) = 1.$$

$$N_{\text{gh}}(\Psi) = 0$$

$$Z_g = \int_{\mathcal{M}_g} d^{M_g} t \frac{\det(\phi_i, \hat{\mu}_j)_{\hat{g}}}{\sqrt{\det(\phi_i, \phi_j)_{\hat{g}}}} \frac{\Omega_{\text{ckv}}[\hat{g}]^{-1}}{\sqrt{\det(\psi_i, \psi_j)_{\hat{g}}}} \int d_{\hat{g}} \Psi d'_{\hat{g}} b d'_{\hat{g}} c e^{-S_m[\hat{g}, \Psi] - S_{\text{gh}}[\hat{g}, b, c]}$$

$$\langle 0 | c_{-1} \bar{c}_{-1} c_0 \bar{c}_0 c_1 \bar{c}_1 | 0 \rangle = 1$$

$$A_{0,2}(k,k')=\frac{C_{S^2}}{\text{Vol}\mathcal{K}_{0,2}}\langle \mathcal{V}_k(\infty,\infty)c_0\bar{c}_0\mathcal{V}_{k'}(0,0)\rangle_{S^2}$$

Weyl ghost.

$$F_{ab}^{\perp}=\sqrt{g}g_{ab}-\sqrt{\hat{g}}\hat{g}_{ab}=0$$

$$S'_{\text{gh}}[g,b,c,c_w]=\frac{1}{4\pi}\int\,\,\,\text{d}^2\sigma\sqrt{g}g^{ab}(b_{ac}\nabla_b c^c+b_{bc}\nabla_a c^c+2b_{ab}c_w)$$

$$\nabla_a c_b + \nabla_b c_a + 2g_{ab}c_w = 0, \nabla^a b_{ab} = 0, g^{ab}b_{ab} = 0$$

$$c_w=-\frac{1}{2}\nabla_a c^a$$

$$T_{ab}'^{\text{gh}}=-(b_{ac}\nabla_b c^c+b_{bc}\nabla_a c^c+2b_{ab}c_w)-\nabla_c(b_{ab}c^c)+\frac{1}{2}g_{ab}g^{cd}(b_{ce}\nabla_d c^e+b_{de}\nabla_c c^e+2b_{cd}c_w)$$

$$g^{ab}T_{ab}'^{\text{gh}}=-g^{ab}\nabla_c(b_{ab}c^c)$$

$$\begin{aligned} g^{ab}(b_{ac}\delta\nabla_b c^c+b_{bc}\delta\nabla_a c^c)&=2g^{ab}b_{ac}\delta\nabla_b c^c=2g^{ab}b_{ac}\delta\Gamma_{bd}^c c^d\\ &=g^{ab}b_{ac}g^{ce}(\nabla_b\delta g_{de}+\nabla_d\delta g_{be}-\nabla_e\delta g_{bd})c^d\\ &=b^{ab}(\nabla_a\delta g_{bc}+\nabla_c\delta g_{ab}-\nabla_b\delta g_{ac})c^c\\ &=b^{ab}\nabla_c\delta g_{ab}c^c \end{aligned}$$

$$\int \, \mathrm{d}_g c_w \mathrm{e}^{-(c_w, g^{ab} b_{ab})_g} = \delta(g^{ab} b_{ab})$$

Modo Cero de los campos fantasmas.

$$Z_g=\int_{\mathcal{M}_g} \mathrm{d}^{M_g}t\frac{\Omega_{\text{ckv}}[\hat{g}]^{-1}}{\sqrt{\det(\psi_i,\psi_j)_{\hat{g}}}}\int \, \mathrm{d}_{\hat{g}}\Psi \, \mathrm{d}_{\hat{g}}b \, \mathrm{d}_{\hat{g}}c \prod_{i=1}^{M_g}(b,\hat{\mu}_i)_{\hat{g}}\mathrm{e}^{-S_m[\hat{g},\Psi]-S_{\text{gh}}[\hat{g},b,c]}$$

$$(b,\hat{\mu}_i)_{\hat{g}}=\int \, \mathrm{d}^2\sigma\sqrt{\hat{g}}G_{\perp}^{abcd}b_{ab}\hat{\mu}_{i,cd}=\int \, \mathrm{d}^2\sigma\sqrt{\hat{g}}g^{ac}g^{bd}b_{ab}\hat{\mu}_{i,cd}.$$



$$b = b_0 + b', b_0 = b_{0i} \phi_i$$

$$1 = \int d_{\hat{g}} b e^{-|b|_{\hat{g}}^2} = J \int d_{\hat{g}} b' \prod_i db_{0i} e^{-|b'|_{\hat{g}}^2 - |b_{0i} \phi_i|^2} = J \sqrt{\det(\phi_i, \phi_j)}.$$

$$\int d^{M_g} b_{0i} \prod_j b_0(\sigma_j^0) = \int d^{M_g} b_{0i} \prod_j [b_{0i} \phi_i(\sigma_j^0)] = \det \phi_i(\sigma_j^0).$$

$$\frac{d_{\hat{g}} b'}{\sqrt{\det(\phi_i, \phi_j)_{\hat{g}}}} = \frac{d_{\hat{g}} b}{\det \phi_i(\sigma_j^0)} \prod_{j=1}^{M_g} b(\sigma_j^0).$$

$$d_{\hat{g}} b' \frac{\det(\phi_i, \hat{\mu}_j)_{\hat{g}}}{\sqrt{\det(\phi_i, \phi_j)_{\hat{g}}}} = d_{\hat{g}} b \prod_{j=1}^{M_g} (b, \hat{\mu}_j)_{\hat{g}}$$

$$\prod_{j=1}^{M_g} b(\sigma_j^0) = \prod_{j=1}^{M_g} [b_{0i} \phi_i(\sigma_j^0)] = \det \phi_i(\sigma_j^0) \prod_{j=1}^{M_g} b_{0i}$$

$$\det(\phi_i, \hat{\mu}_j)_{\hat{g}} \prod_{j=1}^{M_g} b_{0i} = \prod_{j=1}^{M_g} [b_{0i}(\phi_i, \hat{\mu}_j)_{\hat{g}}] = \prod_{j=1}^{M_g} (b_{0i} \phi_i, \hat{\mu}_j)_{\hat{g}} = \prod_{j=1}^{M_g} (b, \hat{\mu}_j)_{\hat{g}}$$

$$Z_g = \int_{\mathcal{M}_g} d^{M_g} t \frac{\Omega_{\text{ckv}}[\hat{g}]^{-1}}{\det \psi_i(\sigma_j^0)} \int d_{\hat{g}} \Psi d_{\hat{g}} b d_{\hat{g}} c \prod_{j=1}^{K_g^c} \frac{\epsilon_{ab}}{2} c^a(\sigma_j^0) c^b(\sigma_j^0)$$

$$\times \prod_{i=1}^{M_g} (\hat{\mu}_i, b)_{\hat{g}} e^{-S_m[\hat{g}, \Psi] - S_{\text{gh}}[\hat{g}, b, c]}$$

$$c \longrightarrow c + c_0, P_1 c_0 = 0.$$

$$b \longrightarrow b + b_0, P_1^\dagger b_0 = 0.$$

$$\Phi_0 = \ln g_s$$

$$g_s^{2g-2} = e^{-\Phi_0 \chi_g} = \exp \left(-\frac{\Phi_0}{4\pi} \int d^2 \sigma \sqrt{g} R \right) = e^{-\Phi_0 S_{\text{EH}}[g]}$$

Amplitudes en un espacio de Moduli – Operadores de vértice. Métrica perturbativa de Green.

$$Z_g = \int \frac{d_g g_{ab}}{\Omega_{\text{gauge}}[g]} Z_m[g], Z_m[g] = \int d_g \Psi e^{-S_m[g, \Psi]}$$

$$V_\alpha(k_i) := \int d^2 \sigma \sqrt{g(\sigma)} V_\alpha(k; \sigma).$$

$$S_{\text{EH}}[g] := \frac{1}{4\pi} \int d^2 \sigma \sqrt{g} R + \frac{1}{2\pi} \oint dsk = \chi_{g,n}$$



$$\chi_{g,n} := \chi(\Sigma_{g,n}) = 2 - 2g - n$$

$$g_s^{-\chi_{g,n}} = e^{-\Phi_0 S_{\text{EH}}[g]}, \Phi_0 := \ln g_s.$$

$$A_{g,n}(\{k_i\})_{\{\alpha_i\}} := \int \frac{d_g g_{ab}}{\Omega_{\text{gauge}}[g]} d_g \Psi e^{-S_m[g,\Psi] - \Phi_0 S_{\text{EH}}[g]} \prod_{i=1}^n \left(\int d^2 \sigma_i \sqrt{g(\sigma_i)} V_{\alpha_i}(k_i; \sigma_i) \right)$$

$$A_{g,n}(\{k_i\})_{\{\alpha_i\}} := A_{g,n}(k_1, \dots, k_n)_{\alpha_1, \dots, \alpha_n} := A_{g,n}(V_{\alpha_1}(k_1), \dots, V_{\alpha_n}(k_n)).$$

$$A_n(k_1, \dots, k_n)_{\alpha_1, \dots, \alpha_n} = \sum_{g=0}^{\infty} A_{g,n}(k_1, \dots, k_n)_{\alpha_1, \dots, \alpha_n}$$

$$A_{g,n}(\{k_i\})_{\{\alpha_i\}} = \int \frac{d_g g_{ab}}{\Omega_{\text{gauge}}[g]} e^{-\Phi_0 S_{\text{EH}}[g]} \int \prod_{i=1}^n d^2 \sigma_i \sqrt{g} \left\langle \prod_{i=1}^n V_{\alpha_i}(k_i; \sigma_i) \right\rangle_{m,g}$$

$$H=0.$$

$$S=1+\mathrm{i}T,$$

$$A_n(k_1,\ldots,k_n)=G_n(k_1,\ldots,k_n)\prod_{i=1}^n(k_i^2+m_i^2).$$

$$G_n(k_1,\ldots,k_n)\propto\delta^{(D)}(k_1+\cdots+k_n).$$

$$T_2=G_2(k,k')(k^2+m^2)^2\sim (k^2+m^2)\delta^{(D)}(k+k')\stackrel{k^2\rightarrow -m^2}{\rightarrow}0$$

$$G_2(k,k')=\frac{\delta^{(D)}(k+k')}{k^2+m^2}$$

$$A_2(k,k')=2k^0(2\pi)^{D-1}\delta^{(D-1)}(\boldsymbol{k}-\boldsymbol{k}').$$

$$\left[a(\boldsymbol{k}),a^\dagger(\boldsymbol{k}')\right]=2k^0(2\pi)^{D-1}\delta^{(D-1)}(\boldsymbol{k}-\boldsymbol{k}').$$

$$\phi(\alpha)=\int\,\mathrm{d}^{D-1}\boldsymbol{k}\alpha(\boldsymbol{k})^*a^\dagger(\boldsymbol{k}).$$

Reparaciones de Gauge. Métrica de Faddeev-Popov.

$$\delta_\xi V_{\alpha_i}(k_i)=\delta_\xi\int\,\mathrm{d}^2\sigma\sqrt{g}V_{\alpha_i}(k_i;\sigma)=0$$

$$\delta_\omega V_{\alpha_i}(k_i)=\delta_\omega\int\,\mathrm{d}^2\sigma\sqrt{g}V_{\alpha_i}(k_i;\sigma)=0$$



$$\begin{aligned}
A_{g,n}(\{k_i\})_{\{\alpha_i\}} &= g_s^{-\chi_{g,n}} \int_{\mathcal{M}_g} d^{\mathbf{M}_g} t \frac{\det(\phi_i, \hat{\mu}_j)_{\hat{g}}}{\sqrt{\det(\phi_i, \phi_j)_{\hat{g}}}} \frac{\Omega_{\text{ckv}}[\hat{g}]^{-1}}{\sqrt{\det(\psi_i, \psi_j)_{\hat{g}}}} \\
&\quad \times \int \prod_{i=1}^n d^2 \sigma_i \sqrt{\hat{g}} \left\langle \prod_{i=1}^n \hat{V}_{\alpha_i}(k_i; \sigma_i) \right\rangle_{m, \hat{g}} \\
A_{g,n}(\{k_i\})_{\{\alpha_i\}} &= g_s^{-\chi_{g,n}} \int_{\mathcal{M}_g} d^{\mathbf{M}_g} t \frac{\Omega_{\text{ckv}}[\hat{g}]^{-1}}{\sqrt{\det(\psi_i, \psi_j)_{\hat{g}}}} \int d_{\hat{g}} b d'_{\hat{g}} c \prod_{i=1}^{\mathbf{M}_g} (b, \hat{\mu}_i)_{\hat{g}} e^{-S_{\text{gh}}[\hat{g}, b, c]} \\
&\quad \times \int \prod_{i=1}^n d^2 \sigma_i \sqrt{\hat{g}} \left\langle \prod_{i=1}^n \hat{V}_{\alpha_i}(k_i; \sigma_i) \right\rangle_{m, \hat{g}} \\
A_{g,n} &= g_s^{-\chi_{g,n}} \int_{\mathcal{M}_g} d^{\mathbf{M}_g} t \frac{\Omega_{\text{ckv}}[\hat{g}]^{-1}}{\det \psi_i(\sigma_j^0)} \int d_{\hat{g}} b d_{\hat{g}} c \prod_{j=1}^{\mathbf{K}_g^c} \frac{\epsilon_{ab}}{2} c^a(\sigma_j^0) c^b(\sigma_j^0) \prod_{i=1}^{\mathbf{M}_g} (\hat{\mu}_i, b)_{\hat{g}} e^{-S_{\text{gh}}[\hat{g}, b, c]} \\
&\quad \times \int \prod_{i=1}^n d^2 \sigma_i \sqrt{\hat{g}} \left\langle \prod_{i=1}^n \hat{V}_{\alpha_i}(k_i; \sigma_i) \right\rangle \\
1 &= \Delta(\sigma_j^0) \int d\xi \prod_{j=1}^{\mathbf{K}_g^c} \delta^{(2)}(\sigma_j - \sigma_j^{0(\xi)}), \sigma_j^{0(\xi)} = \sigma_j^0 + \delta_{\xi} \sigma_j^0, \delta_{\xi} \sigma_j^0 = \xi(\sigma_j^0) \\
\Delta(\sigma_j^0) &= \det \psi_i(\sigma_j^0) \\
\xi(\sigma_j^0) &= \alpha_i \psi_i(\sigma_j^0) \\
1 &= \int \prod_{j=1}^{\mathbf{K}_g^c} d^2 \delta \sigma_j e^{-\sum_j (\delta \sigma_j, \delta \sigma_j)} = \Delta \int \prod_{j=1}^{\mathbf{K}_g} d\alpha_i e^{-\sum_{j,i,i'} (\alpha_i \psi_i(\sigma_j), \alpha_{i'} \psi_{i'}(\sigma_j))} \\
&= \Delta \left(\det \psi_i(\sigma_j) \right)^{-1}
\end{aligned}$$

$$\chi_{g,n} = 2 - 2g - n < 0.$$

$$\begin{aligned}
A_{g,n}(\{k_i\})_{\{\alpha_i\}} &= g_s^{-\chi_{g,n}} \int_{\mathcal{M}_g} d^{\mathbf{M}_g} t \int d_{\hat{g}} b d_{\hat{g}} c \prod_{j=1}^{\mathbf{K}_g^c} \frac{\epsilon_{ab}}{2} c^a(\sigma_j^0) c^b(\sigma_j^0) \prod_{i=1}^{\mathbf{M}_g} (\hat{\mu}_i, b)_{\hat{g}} e^{-S_{\text{gh}}[\hat{g}, b, c]} \\
&\quad \times \int \prod_{i=\mathbf{K}_g^c+1}^n d^2 \sigma_i \sqrt{\hat{g}} \left\langle \prod_{j=1}^{\mathbf{K}_g^c} \hat{V}_{\alpha_j}(k_j; \sigma_j^0) \prod_{i=\mathbf{K}_g^c+1}^n \hat{V}_{\alpha_i}(k_i; \sigma_i) \right\rangle
\end{aligned}$$



$$A_{g,n}(\{k_i\})_{\{\alpha_i\}} = g_s^{-\chi_{g,n}} \int_{\mathcal{M}_g} d^{M_g} t \int_{i=K_g^c+1}^n d^2 \sigma_i \sqrt{\hat{g}} \left\langle \prod_{j=1}^{K_g^c} \frac{\epsilon_{ab}}{2} c^a(\sigma_j^0) c^b(\sigma_j^0) \prod_{i=1}^{M_g} (\hat{\mu}_i, b)_{\hat{g}} \right\rangle_{\text{gh}, \hat{g}} \\ \times \left\langle \prod_{j=1}^{K_g^c} \hat{V}_{\alpha_j}(k_j; \sigma_j^0) \prod_{i=K_g^c+1}^n \hat{V}_{\alpha_i}(k_i; \sigma_i) \right\rangle_{m, \hat{g}}$$

$$A_{g,n}(\{k_i\})_{\{\alpha_i\}} = g_s^{-\chi_{g,n}} \int_{\mathcal{M}_g} d^{M_g} t \int_{i=K_g^c+1}^n d^2 \sigma_i \sqrt{\hat{g}} \left\langle \prod_{i=1}^{M_g} \hat{B}_i \prod_{j=1}^{K_g^c} \hat{V}_{\alpha_j}(k_j; \sigma_j^0) \prod_{i=K_g^c+1}^n \hat{V}_{\alpha_i}(k_i; \sigma_i) \right\rangle$$

$$\hat{V}_{\alpha_j}(k_j; \sigma_j^0) := \frac{\epsilon_{ab}}{2} c^a(\sigma_j^0) c^b(\sigma_j^0) \hat{V}_{\alpha_j}(k_j; \sigma_j^0), \hat{B}_i := (\hat{\mu}_i, b)_{\hat{g}}$$

$$A_{g,n}(\{k_i\})_{\{\alpha_i\}} = g_s^{-\chi_{g,n}} \int_{\mathcal{M}_g \times \mathbb{C}^{n-K_g^c}} \left\langle \bigwedge_{i=1}^{M_g} \hat{B}_i dt_i \prod_{j=1}^{K_g^c} \hat{V}_{\alpha_i}(k_i; \sigma_j^0) \prod_{i=K_g^c+1}^n \hat{V}_{\alpha_i}(k_i; \sigma_i) d^2 \sigma_i \sqrt{\hat{g}} \right\rangle.$$

$$\langle V_k(z,\bar{z})V_{k'}(z',\bar{z}')\rangle_{S^2}=\frac{\mathrm{i}(2\pi)^D\delta^{(D)}(k+k')}{|z-z'|^4}$$

$$A_{0,2}(k,k')=\frac{C_{S_2}}{\mathrm{Vol}\mathcal{K}_{0,0}}\int\mathrm{d}^2z\,\mathrm{d}^2z'\langle V_k(z,\bar{z})V_{k'}(z',\bar{z}')\rangle_{S^2}$$

$$A_{0,2}(k,k')=\frac{C_{S^2}}{\mathrm{Vol}\mathcal{K}_{0,2}}\langle V_k(\infty,\infty)V_{k'}(0,0)\rangle_{S^2}$$

$$A_{0,2}(k,k')=(2\pi)^{D-1}\delta^{(D-1)}(\boldsymbol{k}+\boldsymbol{k}')\frac{C_{S_2}2\pi\mathrm{i}\delta(0)}{\mathrm{Vol}\mathcal{K}_{0,2}}$$

$$A_{0,2}(k,k')=\frac{C_{S^2}}{\mathrm{Vol}\mathcal{K}_{0,2}}\langle V_k(\infty,\infty)V_{k'}(0,0)\rangle_{S^2}$$

$$\mathrm{Vol}\mathcal{K}_{0,2}=\int\frac{\mathrm{d}^2z}{|z|^2}=2\int_0^{2\pi}\mathrm{d}\sigma\int_0^\infty\frac{\mathrm{d}r}{r}$$

$$\mathrm{Vol}\mathcal{K}_{0,2}=4\pi\int_0^\infty\frac{\mathrm{d}r}{r}=4\pi\int_{-\infty}^\infty\mathrm{d}\tau=4\pi\lim_{\varepsilon\rightarrow 0}\int_{-\infty}^\infty\mathrm{d}\tau\mathrm{e}^{\mathrm{i}\varepsilon\tau}=4\pi\times 2\pi\lim_{\varepsilon\rightarrow 0}\delta(\varepsilon)$$

$$\mathrm{Vol}_\varepsilon\mathcal{K}_{0,2}=8\pi^2\delta(\varepsilon)$$

$$\mathrm{Vol}_{M,E}\mathcal{K}_{0,2}=8\pi^2\mathrm{i}\delta(E)$$

$$X^0(z,\bar{z})=x^0+\frac{\mathrm{i}}{2}\alpha'k^0\ln\,|z|^2=x^0+\mathrm{i}\alpha'k^0\tau,$$

$$X_M^0=x_M^0+\alpha'k_M^0t$$

$$\text{Vol}_M \mathcal{K}_{0,2} \rightarrow \frac{8\pi^2 i \delta(0)}{\alpha' k_M^0} = \frac{C_{S_2} 2\pi i \delta(0)}{2k_M^0}$$

$$A_{0,2}(k,k')=2k^0(2\pi)^{D-1}\delta^{(D-1)}(\boldsymbol{k}+\boldsymbol{k}')$$

$$Z_0 \sim \frac{\delta^{(D)}(0)}{\text{Vol}\mathcal{K}_0}$$

Cuántización y simetría BRST – Transformaciones BRST. Parámetros de Grassmann.

$$Z_g=\int_{\mathcal{M}_g}\frac{\mathrm{d}^{M_g}t}{\Omega_{\text{ckv}}[g]}\int\mathrm{d}_g g_{ab}\,\mathrm{d}_g\Psi\,\mathrm{d}_g b\,\mathrm{d}_g'c\delta(\sqrt{g}g_{ab}-\sqrt{\hat{g}}\hat{g}_{ab})\prod_{i=1}^{M_g}(\phi_i,b)_g\mathrm{e}^{-S_m[g,\Psi]-S_{\text{gh}}[g,b,c]}$$

$$Z_g=\int_{\mathcal{M}_g}\frac{\mathrm{d}^{M_g}t}{\Omega_{\text{ckv}}[g]}\int\mathrm{d}_g g_{ab}\,\mathrm{d}_g B^{ab}\,\mathrm{d}_g\Psi\,\mathrm{d}_g b\,\mathrm{d}_g'c\prod_{i=1}^{M_g}(\phi_i,b)_g\mathrm{e}^{-S_m[g,\Psi]-S_{\text{gf}}[g,\hat{g},B]-S_{\text{gh}}[g,b,c]}$$

$$S_{\text{gf}}[g,\hat{g},B]=-\frac{\text{i}}{4\pi}\int\,\mathrm{d}^2\sigma B^{ab}(\sqrt{g}g_{ab}-\sqrt{\hat{g}}\hat{g}_{ab})$$

$$\delta_{\epsilon}g_{ab}=\text{i}\epsilon\mathcal{L}_cg_{ab},\delta_{\epsilon}\Psi=\text{i}\epsilon\mathcal{L}_c\Psi,\\ \delta_{\epsilon}c^a=\text{i}\epsilon\mathcal{L}_cc^a,\delta_{\epsilon}b_{ab}=\epsilon B_{ab},\delta_{\epsilon}B_{ab}=0,$$

$$\delta_{\epsilon}g_{ab}=\text{i}\epsilon\mathcal{L}_cg_{ab}+\text{i}\epsilon g_{ab}c_w,\delta_{\epsilon}c_w=\text{i}\epsilon\mathcal{L}_cc_w$$

$$B_{ab}=\text{i}T_{ab}\!:=\!\text{i}\left(T_{ab}^m+T_{ab}^{\text{gh}}\right),$$

$$\delta_{\epsilon}\Psi=\text{i}\epsilon\mathcal{L}_c\Psi,\delta_{\epsilon}c^a=\text{i}\epsilon\mathcal{L}_cc^a,\delta_{\epsilon}b_{ab}=\text{i}\epsilon T_{ab}$$

$$\delta_{\epsilon}c^a=\epsilon c^b\partial_b c^a$$

$$Q_B=\int\,\mathrm{d}\sigma j_B^0$$

$$Q_B^2=0$$

$$N_{\text{gh}}(Q_B)=1$$

$$\delta_{\epsilon}\Psi=\text{i}[\epsilon Q_B,\Psi]_{\pm}$$

$$T_{ab}=[Q_B,b_{ab}]$$

$$|\psi\rangle\in\mathcal{H}(Q_B)\!:=\frac{\ker Q_B}{\text{Im}Q_B},$$

$$Q_B|\psi\rangle=0,\nexists|\chi\rangle\!:\!|\psi\rangle=Q_B|\chi\rangle.$$

$$|\psi\rangle\sim|\psi\rangle+Q_B|\Lambda\rangle.$$

$$\int\,\mathrm{d}\sigma b_{ab}|\psi\rangle=0$$



$$b^+ := \int d\sigma b_{00}, b^- := \int d\sigma b_{01}$$

$$\mathcal{H}^-(Q_B) = \mathcal{H}(Q_B) \cap \ker b^-, \mathcal{H}^0(Q_B) = \mathcal{H}^-(Q_B) \cap \ker b^+.$$

$$P_\sigma|\psi\rangle=0$$

$$P_\sigma|\psi\rangle=\int d\sigma T_{01}|\psi\rangle=\int d\sigma\{Q_B,b_{01}\}|\psi\rangle=Q_B\int d\sigma b_{01}|\psi\rangle,$$

$$b^-|\psi\rangle=0$$

$$\mathcal{H}^-:=\mathcal{H}_\downarrow\oplus\mathcal{H}_\uparrow,\mathcal{H}_\downarrow:=\mathcal{H}^0:=\mathcal{H}^-\cap\ker b^+.$$

$$b^+|\downarrow\rangle=0,b^+|\uparrow\rangle=|\downarrow\rangle$$

$$Q_B|\psi_\downarrow\rangle=H|\psi_\uparrow\rangle,Q_B|\psi_\uparrow\rangle=0$$

$$\begin{array}{l} \text{Im}Q_B=\{|\psi_\uparrow\rangle\in\mathcal{H}_\uparrow|H|\psi_\uparrow\rangle\neq0\}\\ \text{ker}Q_B=\{|\psi_\uparrow\rangle\in\mathcal{H}_\uparrow\}\cup\{|\psi_\downarrow\rangle\in\mathcal{H}_\downarrow|H|\psi_\downarrow\rangle=0\} \end{array}$$

$$H|\psi\rangle=0$$

$$b^+|\psi\rangle=0$$

$$Q_B\int d\sigma b_{00}|\psi\rangle=0$$

Coordenadas geométricas complejas de un campo de gauge. Métrica de Liouville y simetrías tensoriales.

$$ds^2=g_{ab}\,d\sigma^a\,d\sigma^b=e^{2\phi(\tau,\sigma)}(d\tau^2+d\sigma^2),$$

$$g_{ab}=e^{2\phi}\delta_{ab},\hat{g}_{ab}=\delta_{ab}.$$

$$\begin{array}{l} z=\tau+i\sigma,\quad \bar{z}=\tau-i\sigma\\ \tau=\frac{z+\bar{z}}{2},\quad \sigma=\frac{z-\bar{z}}{2i} \end{array}$$

$$ds^2=2g_{z\bar{z}}\,dz\,d\bar{z}=e^{2\phi(z,\bar{z})}|dz|^2.$$

$$\begin{array}{l} g_{z\bar{z}}=\frac{e^{2\phi}}{2},g_{zz}=g_{\bar{z}\bar{z}}=0\\ g^{z\bar{z}}=2e^{-2\phi},g^{zz}=g^{\bar{z}\bar{z}}=0 \end{array}$$

$$\hat{g}_{z\bar{z}}=\frac{1}{2},\hat{g}^{z\bar{z}}=2$$

$$w=w(z),\bar{w}=\bar{w}(\bar{z}).$$

$$e^{2\phi(z,\bar{z})}=\left|\frac{\partial w}{\partial z}\right|^2e^{2\phi(w,\bar{w})}$$



$$ds^2 = e^{2\phi(w, \bar{w})} |dw|^2$$

$$d^2\sigma:=d\tau\,d\sigma=\frac{1}{2}\,d^2z,\,d^2z:=dz\,d\bar{z}$$

$$\delta^{(2)}(z)=\frac{1}{2}\delta^{(2)}(\sigma)$$

$$\int\,d^2z\delta^{(2)}(z)=\int\,d^2\sigma\delta^{(2)}(\sigma)=1$$

$$\partial_z=\frac{1}{2}(\partial_\tau-i\partial_\sigma),\partial_{\bar{z}}=\frac{1}{2}(\partial_\tau+i\partial_\sigma)\\ dz=d\tau+i\,d\sigma,d\bar{z}=d\tau-i\,d\sigma$$

$$\epsilon_{01}=\epsilon^{01}=1\\ \epsilon_{z\bar{z}}=\frac{i}{2},\epsilon^{z\bar{z}}=-2i$$

$$V^z=V^0+iV^1,V^{\bar{z}}=V^0-iV^1$$

$$V=V^0\partial_0+V^1\partial_1=V^z\partial_z+V^{\bar{z}}\partial_{\bar{z}}$$

$$V^w=\frac{\partial w}{\partial z}V^z,V^{\bar{w}}=\frac{\partial \bar{w}}{\partial \bar{z}}V^{\bar{z}}$$

$$T\Sigma_g\simeq T\Sigma_g^+\oplus T\Sigma_g^-\\ V^z\partial_z\in T\Sigma_g^+,V^{\bar{z}}\partial_{\bar{z}}\in T\Sigma_g^-$$

$$\omega_z=\frac{1}{2}(\omega_0-i\omega_1),\omega_{\bar{z}}=\frac{1}{2}(\omega_0+i\omega_1)$$

$$\omega=\omega_0\,d\sigma^0+\omega_1\,d\sigma^1=\omega_z\,dz+\omega_{\bar{z}}\,d\bar{z}$$

$$T^*\Sigma_g\simeq\Omega^{1,0}(\Sigma_g)\oplus\Omega^{0,1}(\Sigma_g)\\ \omega_z\,dz\in\Omega^{1,0}(\Sigma_g),\omega_{\bar{z}}\,d\bar{z}\in\Omega^{0,1}(\Sigma_g)$$

$$V_z=g_{z\bar{z}}V^{\bar{z}},V_{\bar{z}}=g_{z\bar{z}}V^z$$

$$T^{\overbrace{z'\cdots z}^{q_++p_-}}\underbrace{z\cdots z}_{p_++q_-}=(g^{z\bar{z}})^{p_-}(g_{z\bar{z}})^{q_-}T^{\overbrace{z''z\bar{z}\cdots\bar{z}}^{q_+}\overbrace{\phantom{z''z\bar{z}\cdots\bar{z}}}^{q_-}}\underbrace{z\cdots z}_{p_+}\underbrace{\bar{z}\cdots\bar{z}}_{p_-}.$$

$$T^{\overbrace{w\cdots w}^q}\underbrace{w\cdots w}_p=\left(\frac{\partial w}{\partial z}\right)^nT^{\overbrace{z\cdots z}^q}\underbrace{z\cdots z}_p,n:=q-p.$$

$$T^{zz}=2(T^{00}+iT^{01})\in\mathcal{T}^2,T^{\bar{z}\bar{z}}=2(T^{00}-iT^{01})\in\mathcal{T}^{-2},T^z{}_z=0$$

$$T_{zz}=g_{z\bar{z}}g_{z\bar{z}}T^{\bar{z}\bar{z}}=\frac{1}{2}(T^{00}-iT^{01})$$

$$T^{zz}=\left(\frac{\partial z}{\partial \tau}\right)^2T^{00}+\left(\frac{\partial z}{\partial \sigma}\right)^2T^{11}+2\frac{\partial z}{\partial \tau}\frac{\partial z}{\partial \sigma}T^{01}=T^{00}-T^{11}+2i T^{01}$$



$$\int d^2z(\partial_z v^z + \partial_{\bar{z}} v^{\bar{z}}) = -i\oint (dz v^{\bar{z}} - d\bar{z} v^z) = -2i\oint_{\partial R} (v_z dz - v_{\bar{z}} d\bar{z})$$

$$\psi_i(z,\bar{z})=\psi_i^z\partial_z+\psi_i^{\bar{z}}\partial_{\bar{z}},\phi_i(z,\bar{z})=\phi_{i,zz}(\mathrm{d}z)^2+\phi_{i,\bar{z}}(\mathrm{d}\bar{z})^2$$

$$(P_1\xi)_{zz}=2\nabla_z\xi_z=\partial_z\xi^{\bar{z}},(P_1\xi)_{\bar{z}\bar{z}}=2\nabla_{\bar{z}}\xi_{\bar{z}}=\partial_{\bar{z}}\xi^z\\ (P_1^\dagger T)_z=-2\nabla^zT_{zz}=-4\partial_{\bar{z}}T_{z\bar{z}},(P_1^\dagger T)_{\bar{z}}=-2\nabla^{\bar{z}}T_{\bar{z}\bar{z}}=-4\partial_zT_{\bar{z}\bar{z}}$$

$$\psi^z=\psi^z(z),\psi^{\bar{z}}=\psi^{\bar{z}}(\bar{z}),\phi_{zz}=\phi_{zz}(z),\phi_{\bar{z}\bar{z}}=\phi_{\bar{z}\bar{z}}(\bar{z})$$

$$\ker P_1=\mathrm{Span}\{\psi_K(z)\}\oplus\mathrm{Span}\{\bar{\psi}_K(\bar{z})\},K=1,\ldots,K_g^c$$

$$\ker P_1^\dagger=\mathrm{Span}\{\phi_I(z)\}\oplus\mathrm{Span}\{\bar{\phi}_I(\bar{z})\},I=1,\ldots,M_g^c$$

$$m_I=t_{2I-1}+it_{2I},\bar{m}_I=t_{2I-1}-it_{2I},I=1,\ldots,M_g^c$$

$$\mathrm{d}^{\mathsf{M}_g}t=\mathrm{d}^{2\mathsf{M}_g^c}m$$

$$(T_1,T_2)=2\int d^2\sigma\sqrt{\hat{g}}\hat{g}^{ac}g^{bd}T_{1,ab}T_{2,cd}=4\int d^2z(T_{1,zz}T_{2,\bar{z}\bar{z}}+T_{1,\bar{z}\bar{z}}T_{2,zz})\\ (\xi_1,\xi_2)=\int d^2\sigma\sqrt{\hat{g}}\hat{g}_{ab}\xi^a\xi^b=\frac{1}{4}\int d^2z(\xi_1^z\xi_2^{\bar{z}}+\xi_1^{\bar{z}}\xi_2^z)$$

$$\mu_{izz}=\partial_i\bar{g}_{zz},\mu_{i\bar{z}\bar{z}}=\partial_i\bar{g}_{\bar{z}\bar{z}}$$

$$Z_g=\int_{\mathcal{M}_g} \mathrm{d}^{2\mathsf{M}_g^c}m\frac{|\det(\phi_I,\mu_J)|^2}{|\det(\phi_I,\bar{\phi}_J)|}\frac{\det'P_1^\dagger P_1}{|\det(\psi_I,\bar{\psi}_J)|}\frac{Z_m[\delta]}{\Omega_{\mathrm{ckv}}[\delta]}$$

$$c\!:=\!c^z,\bar{c}\!:=\!c^{\bar{z}},b\!:=\!b_{zz},\bar{b}\!:=\!b_{\bar{z}\bar{z}}$$

$$S_{\mathrm{gh}}[g,b,c]=\frac{1}{2\pi}\int d^2z(b\partial_{\bar{z}}c+\bar{b}\partial_z\bar{c})$$

$$\partial_zc=0,\partial_zb=0,\partial_{\bar{z}}\bar{c}=0,\partial_{\bar{z}}\bar{b}=0$$

$$\bigwedge_{i=1}^{\mathsf{M}_g} B_i\,\mathrm{d}t_i=\bigwedge_{I=1}^{\mathsf{M}_g^c} B_I\bar{B}_I\,\mathrm{d}m_I\wedge\bar{m}_I, B_I\!:=\!(\mu_I,b)$$

$$Z_g=\int_{\mathcal{M}_g} \mathrm{d}^{2\mathsf{M}_g^c}m\frac{\Omega_{\mathrm{ckv}}[\delta]^{-1}}{|\det\psi_I(z_j^0)|^2}\int \mathrm{d}(b,\bar{b})\mathrm{d}(c,\bar{c})\prod_{j=1}^{K_g^c} c(z_j^0)\bar{c}(\bar{z}_j^0)\prod_{I=1}^{\mathsf{M}_g^c} |(\mu_I,b)|^2\mathrm{e}^{-S_{\mathrm{gh}}[b,c]}Z_m[\delta].$$

Simetría conforme en dimensión $\mathbb{R}^4$. Ecuaciones de Killing y métrica Beltrami-Laplace.

$$x^\mu\longrightarrow x'^\mu=x'^\mu(x)$$

$$g_{\mu\nu}(x)\longrightarrow g'_{\mu\nu}(x')=\frac{\partial x^\rho}{\partial x'^\mu}\frac{\partial x^\sigma}{\partial x'^\nu}g_{\rho\sigma}(x)=\Omega(x')^2g_{\mu\nu}(x')$$

$$\frac{u\cdot v}{|u||v|}=\frac{u'\cdot v'}{|u'||v'|}$$



$$\Omega:=\mathrm{e}^{\omega}.$$

$$\delta x^\mu=\xi^\mu$$

$$\delta g_{\mu\nu}=\mathcal{L}_\xi g_{\mu\nu}=\nabla_\mu\xi_\nu+\nabla_\nu\xi_\mu=\frac{2}{d}g_{\mu\nu}\nabla_\rho\xi^\rho$$

$$\Omega^2=1+\frac{2}{d}\nabla_\rho\xi^\rho$$

$$\mathrm{ISO}(\mathcal{M})\subset \mathrm{CISO}(\mathcal{M})$$

$$\delta g_{\mu\nu}=\mathcal{L}_\xi g_{\mu\nu}=\nabla_\mu\xi_\nu+\nabla_\nu\xi_\mu=0$$

$$\eta=\mathrm{diag}(\underbrace{-1,\ldots,-1}_q,\underbrace{1,\ldots,1}_p)$$

$$(\eta_{\mu\nu}\Delta+(D-2)\partial_\mu\partial_\nu)\partial\cdot\epsilon=0$$

$$\begin{array}{l} \xi^\mu=a^\mu \\ \xi^\mu=\omega^\mu{}_\nu x^\nu \\ \xi^\mu=\lambda x^\mu \\ \xi^\mu=b^\mu x^2-2b\cdot xx^\mu \end{array}$$

$$\mathrm{ISO}(\mathbb{R}^{p,q})=\mathrm{SO}(p,q), \mathrm{CISO}(\mathbb{R}^{p,q})=\mathrm{SO}(p+1,q+1)$$

$$\dim\mathrm{SO}(p+1,q+1)=\frac{1}{2}(p+q+2)(p+q+1)$$

Planos complejos.

Métrica de Riemann.

$$\overline{\mathbb{C}}=\mathbb{C}\cup\{\infty\}.$$

$$\lim_{r\rightarrow\infty}r\mathrm{e}^{\mathrm{i}\theta}:\!=\infty$$

$$z=\mathrm{e}^{\mathrm{i}\phi}\cot\frac{\theta}{2}$$

$$\begin{array}{l} z=x+\mathrm{i}y,\quad \bar{z}=x-\mathrm{i}y \\ x=\frac{z+\bar{z}}{2},\quad y=\frac{z-\bar{z}}{2\mathrm{i}} \end{array}$$

$$\mathrm{d}s^2=\mathrm{d}x^2+\mathrm{d}y^2=\mathrm{d}z\,\mathrm{d}\bar{z}$$

$$\partial\!:=\partial_z=\frac{1}{2}(\partial_x-\mathrm{i}\partial_y),\bar{\partial}\!:=\partial_{\bar{z}}=\frac{1}{2}(\partial_x+\mathrm{i}\partial_y)$$

$$\partial\phi(z_1)\partial\phi(z_2)\!:=\partial_{z_1}\partial_{z_2}\phi(z_1)\phi(z_2).$$



$$w = \frac{1}{z}$$

Métrica de Lorentz.

$$t \in \mathbb{R}, \sigma \in [0, L), \sigma \sim \sigma + L,$$

$$ds^2 = -dt^2 + d\sigma^2 = -d\sigma^+ d\sigma^-,$$

$$d\sigma^\pm = dt \pm d\sigma$$

$$\tau = it,$$

$$ds^2 = d\tau^2 + d\sigma^2.$$

$$w = \tau + i\sigma, \bar{w} = \tau - i\sigma$$

$$ds^2 = dw d\bar{w}.$$

$$w = i(t + \sigma) = i\sigma^+, \bar{w} = i(t - \sigma) = i\sigma^-.$$

$$z = e^{2\pi w/L}, \bar{z} = e^{2\pi \bar{w}/L},$$

$$ds^2 = \left(\frac{L}{2\pi}\right)^2 \frac{dz d\bar{z}}{|z|^2}.$$

Álgebra de Witt – Métrica de Cauchy-Riemann- Killing – Laurent – Möbius. Definiciones cuánticas de un CFT y sus operadores y propagadores. Métrica de Virasoro. Funciones de Correlación.

$$ds^2 = dz d\bar{z},$$

$$z \rightarrow z' = f(z), \bar{z} \rightarrow \bar{z}' = \bar{f}(\bar{z})$$

$$ds^2 = dz' d\bar{z}' = \left| \frac{df}{dz} \right|^2 dz d\bar{z}.$$

$$\delta z = v(z), \delta \bar{z} = \bar{v}(\bar{z}),$$

$$\bar{\partial}v = 0, \partial\bar{v} = 0.$$

$$v(z) = \sum_{n \in \mathbb{Z}} v_n z^{n+1}, \bar{v}(\bar{z}) = \sum_{n \in \mathbb{Z}} \bar{v}_n \bar{z}^{n+1},$$

$$\ell_n = -z^{n+1} \partial_z, \bar{\ell}_n = -\bar{z}^{n+1} \partial_{\bar{z}}, n \in \mathbb{Z}.$$

$$[\ell_m, \ell_n] = (m-n)\ell_{m+n}, [\bar{\ell}_m, \bar{\ell}_n] = (m-n)\bar{\ell}_{m+n}, [\ell_m, \bar{\ell}_n] = 0.$$

$$\delta g_{z\bar{z}} = \partial v + \bar{\partial} \bar{v}, \delta g_{zz} = \delta g_{\bar{z}\bar{z}} = 0$$

$$\lim_{|z| \rightarrow 0} v(z) < \infty \implies \forall n < -1: v_n = 0$$



$$v(1/w) = \frac{dz}{dw} \sum_n v_n w^{-n-1}$$

$$\lim_{|z|\rightarrow\infty} v(z) = \lim_{|w|\rightarrow 0} \frac{dz}{dw} v(1/w) = - \lim_{|w|\rightarrow 0} \frac{v(1/w)}{w^2} < \infty \implies \forall n > 1: v_n = 0$$

$$\{\ell_{-1},\ell_0,\ell_1\}\cup\{\bar{\ell}_{-1},\bar{\ell}_0,\bar{\ell}_1\}$$

$$\ell_{-1}=-\partial_z, \ell_0=-z\partial_z, \ell_1=-z^2\partial_z$$

$$[\ell_0,\ell_{\pm 1}]=\mp \ell_{\pm 1}, [\ell_1,\ell_{-1}]=2\ell_0$$

$$\mathrm{PSL}(2,\mathbb{C})\colon=\mathrm{SL}(2,\mathbb{C})/\mathbb{Z}_2\sim\mathrm{SO}(3,1)$$

$$\mathcal{K}_0=\mathrm{PSL}(2,\mathbb{C})$$

$$g=\begin{pmatrix}a&b\\c&d\end{pmatrix}, a,b,c,d\in\mathbb{C}, \det g=ad-bc=1$$

$$K_0:=\dim\mathrm{SL}(2,\mathbb{C})=6.$$

$$f_g(z)=\frac{az+b}{cz+d}$$

$$v(z)=\beta+2\alpha z+\gamma z^2, \bar{v}(\bar{z})=\bar{\beta}+2\bar{\alpha}\bar{z}+\bar{\gamma}\bar{z}^2$$

$$a=1+\alpha, b=\beta, c=-\gamma, d=1-\alpha$$

$$\text{translation:}\quad f_g(z)=z+a,\qquad a\in\mathbb{C},$$

$$\text{rotation:}\quad f_g(z)=\zeta z,\qquad |\zeta|=1,$$

$$\text{dilatation:}\quad f_g(z)=\lambda z,\qquad \lambda\in\mathbb{R},$$

$$\text{SCT:}\quad f_g(z)=\frac{z}{cz+1},\qquad c\in\mathbb{C}.$$

$$\text{inversion: } I^+(z)\colon=I(z)\colon=\frac{1}{z}$$

$$I^-(z)\colon=-I(z)=I(-z)=-\frac{1}{z}$$

$$g_{\infty,0,1}(z)=\frac{1}{1-z}$$

$$\forall f \text{ meromorphic : } \mathcal{O}(z,\bar{z})=\left(\frac{\mathrm{d}f}{\mathrm{d}z}\right)^h\left(\frac{\mathrm{d}\bar{f}}{\mathrm{d}\bar{z}}\right)^{\bar{h}}\mathcal{O}'(f(z),\bar{f}(\bar{z})),$$



$$\forall f \in \text{PSL}(2, \mathbb{C}): \mathcal{O}(z, \bar{z}) = \left(\frac{df}{dz} \right)^h \left(\frac{d\bar{f}}{d\bar{z}} \right)^{\bar{h}} \mathcal{O}'(f(z), \bar{f}(\bar{z}))$$

$$\Delta := h + \bar{h}, s := h - \bar{h}.$$

$$\mathcal{O}(z,\bar{z})\mathrm{d}z^h\,\mathrm{d}\bar{z}^{\bar{h}}$$

$$f\circ \mathcal{O}(z,\bar{z}):=f'(z)^h\bar{f}'(\bar{z})^{\bar{h}}\mathcal{O}'(f(z),\bar{f}(\bar{z}))$$

$$\delta z=v(z),\delta\bar{z}=\bar{v}(\bar{z})$$

$$\delta \mathcal{O}(z,\bar{z})=(h\partial v+v\partial)\mathcal{O}(z,\bar{z})+(\bar{h}\bar{\partial}\bar{v}+\bar{v}\bar{\partial})\mathcal{O}(z,\bar{z})$$

$$\nabla^\nu T_{\mu\nu}=0, g^{\mu\nu}T_{\mu\nu}=0$$

$$g^{\mu\nu}T_{\mu\nu}=4T_{z\bar{z}}=T_{xx}+T_{yy}=0$$

$$T_{z\bar{z}}=0$$

$$\partial_z T_{\bar{z}\bar{z}}=0, \partial_{\bar{z}} T_{zz}=0$$

$$T(z)\!:=\!T_{zz}(z),\bar{T}(\bar{z})\!:=\!T_{\bar{z}\bar{z}}(\bar{z})$$

$$J_v(z)\!:=\!J_v^{\bar{z}}(z)=-T(z)v(z),\bar{J}_v(\bar{z})\!:=\!J_v^z(\bar{z})=-\bar{T}(\bar{z})\bar{v}(\bar{z})$$

$$Z=\int\,\mathrm{d}\Psi\mathrm{e}^{-S[\Psi]}$$

$$[L_m,L_n]=(m-n)L_{m+n}+\frac{c}{12}m(m-1)(m+1)\delta_{m+n}$$

$$[\bar{L}_m,\bar{L}_n]=(m-n)\bar{L}_{m+n}+\frac{\bar{c}}{12}m(m-1)(m+1)\delta_{m+n}$$

$$[L_m,\bar{L}_n]=0,[c,L_m]=0,[\bar{c},\bar{L}_m]=0$$

$$\left\langle \prod_{i=1}^n \mathcal{O}_i(z_i,\bar{z}_i) \right\rangle = \int \mathrm{d}\Psi \mathrm{e}^{-S[\Psi]} \prod_{i=1}^n \mathcal{O}_i(z_i,\bar{z}_i),$$

$$\left\langle \prod_{i=1}^n \mathcal{O}_i(z_i,\bar{z}_i) \right\rangle = \prod_{i=1}^n \left(\frac{\mathrm{d}f}{\mathrm{d}z}(z_i) \right)^{h_i} \left(\frac{\mathrm{d}\bar{f}}{\mathrm{d}\bar{z}}(\bar{z}_i) \right)^{\bar{h}_i} \times \left\langle \prod_{i=1}^n \mathcal{O}_i\left(f(z_i),\bar{f}(\bar{z}_i)\right) \right\rangle.$$

$$\delta \left\langle \prod_{i=1}^n \mathcal{O}_i(z_i,\bar{z}_i) \right\rangle = \sum_{i=1}^n \left(h_i \partial_i v(z_i) + v(z_i) \partial_i + \text{c.c.} \right) \left\langle \prod_{i=1}^n \mathcal{O}_i(z_i,\bar{z}_i) \right\rangle = 0,$$



$$\begin{aligned}
\langle \mathcal{O}_i(z_i, \bar{z}_i) \rangle &= \delta_{h_i, 0} \delta_{\bar{h}_i, 0}, \\
\langle \mathcal{O}_i(z_i, \bar{z}_i) \mathcal{O}_j(z_j, \bar{z}_j) \rangle &= \delta_{h_i, h_j} \delta_{\bar{h}_i, \bar{h}_j} \frac{g_{ij}}{z_{ij}^{2h_i} \bar{z}_{ij}^{2\bar{h}_i}}, \\
\langle \mathcal{O}_i(z_i, \bar{z}_i) \mathcal{O}_j(z_j, \bar{z}_j) \mathcal{O}_k(z_k, \bar{z}_k) \rangle &= \frac{C_{ijk}}{z_{ij}^{h_i+h_j-h_k} z_{jk}^{h_j+h_k-h_i} z_{ki}^{h_i+h_k-h_j}} \\
&\times \frac{1}{\bar{z}_{ij}^{\bar{h}_i+\bar{h}_j-\bar{h}_k} \bar{z}_{jk}^{\bar{h}_j+\bar{h}_k-\bar{h}_i} \bar{z}_{ki}^{\bar{h}_i+\bar{h}_k-\bar{h}_j}}, \\
z_{ij} &= z_i - z_j.
\end{aligned}$$

$$\left\langle \prod_{i=1}^4 \mathcal{O}_i(z_i, \bar{z}_i) \right\rangle = f(x, \bar{x}) \prod_{i < j} \frac{1}{z_{ij}^{(h_i+h_j)-h/3}} \times \text{c.c.}$$

$$h := \sum_{i=1}^4 h_i, \bar{h} := \sum_{i=1}^4 \bar{h}_i.$$

$$x := \frac{z_{12} z_{34}}{z_{13} z_{24}}.$$

Operadores formales y cuantización radial. Ordenadores radiales y conmutadores. Productos de expansión.

$$z = e^{\tau + i\sigma} = x + iy.$$

$$\tau \longrightarrow \tau + T$$

$$z \longrightarrow e^T z.$$

$$H = \frac{2\pi}{L} (L_0 + \bar{L}_0)$$

$$\text{on-shell state: } h + \bar{h} = 0.$$

$$R(A(z)B(w)) = \begin{cases} A(z)B(w) & |z| > |w| \\ (-1)^F B(w)A(z) & |w| > |z| \end{cases}$$

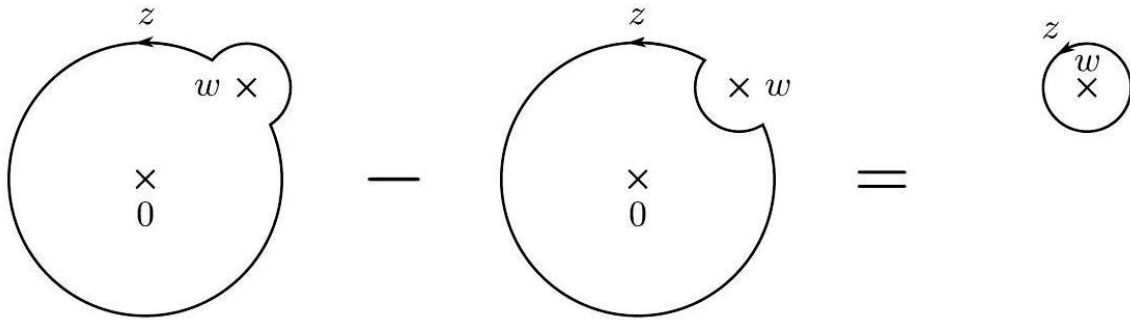
$$[A(z), B(w)]_{\pm, |z|=|w|} = \lim_{\delta \rightarrow 0} (A(z)B(w)|_{|z|=|w|+\delta} \pm B(w)A(z)|_{|z|=|w|-\delta})$$

$$A = \oint_{c_0} \frac{dz}{2\pi i} a(z), B = \oint_{c_0} \frac{dz}{2\pi i} b(z)$$

$$[A, B]_{\pm} = \oint_{c_0} \frac{dw}{2\pi i} \oint_{c_w} \frac{dz}{2\pi i} a(z)b(w)$$

$$[A, b(w)]_{\pm} = \oint_{c_w} \frac{dz}{2\pi i} a(z)b(w)$$





$$\partial_{\mu} j^{\mu} = \partial j^z + \bar{\partial} j^{\bar{z}} = 2(\partial j_{\bar{z}} + \bar{\partial} j_z) = 0$$

$$Q = \frac{1}{2\pi i} \oint_{C_0} (j_z dz - j_{\bar{z}} d\bar{z})$$

$$j(z) := j_z(z), \bar{j}(\bar{z}) := j_{\bar{z}}(\bar{z})$$

$$Q = Q_L + Q_R, Q_L := \frac{1}{2\pi i} \oint_{C_0} j(z) dz, Q_R := -\frac{1}{2\pi i} \oint_{C_0} \bar{j}(\bar{z}) d\bar{z}$$

$$\delta_{\epsilon} \mathcal{O}(z, \bar{z}) = -[\epsilon Q, \mathcal{O}(z, \bar{z})] = -\epsilon \oint_{C_z} \frac{dw}{2\pi i} j(w) \mathcal{O}(z, \bar{z}) + \epsilon \oint_{C_{\bar{z}}} \frac{d\bar{w}}{2\pi i} \bar{j}(\bar{w}) \mathcal{O}(z, \bar{z})$$

$$Q=\frac{1}{2\pi}\int\, \mathrm{d}\sigma j^0$$

$$Q=\frac{1}{2\pi}\int\, (\,\mathrm{d}\sigma j^0-\mathrm{d}\tau j^1)= -\frac{1}{2\pi}\int\, \epsilon_{\mu\nu}j^{\mu}\mathrm{d}x^{\nu}$$

$$Q=-\frac{1}{2\pi}\oint\,\epsilon_{z\bar{z}}(j^z\,\mathrm{d}\bar{z}-j^{\bar{z}}\,\mathrm{d}z)=-\frac{\mathrm{i}}{4\pi}\oint\,(j^z\,\mathrm{d}\bar{z}-j^{\bar{z}}\,\mathrm{d}z)=-\frac{1}{2\pi\mathrm{i}}\oint\,(j_z\,\mathrm{d}z-j_{\bar{z}}\,\mathrm{d}\bar{z})$$

$$\mathcal{O}_i(z_i,\bar{z}_i)\mathcal{O}_j(z_j,\bar{z}_j)=\sum_k\frac{c_{ij}^k}{z_{ij}^{h_i+h_j-h_k}\bar{z}_{ij}^{\bar{h}_i+\bar{h}_j-\bar{h}_k}}\mathcal{O}_k(z_j,\bar{z}_j),$$

$$C_{ijk}=g_{k\ell}c_{ij}^{\ell}.$$

Identidad OPE.

$$\phi(z)1=\sum_{n\in\mathbb{N}}\frac{(z-w)^n}{n!}\partial^n\phi(w)$$

$$A(z)B(w) \coloneqq \sum_{n=-\infty}^N \frac{\{AB\}_n(z)}{(z-w)^n}$$

$$A(z)B(w) \sim \sum_{n=1}^N \frac{\{AB\}_n(z)}{(z-w)^n} =: A(z)B(w)$$



$$\phi_i(z_i)\phi_j(z_j) \sim \sum_k \theta(h_i + h_j - h_k) \frac{c_{ij}^k}{(z-w)^{h_i+h_j-h_k}} \phi_k(w)$$

$$T(z)\phi(w) \sim \frac{h\phi(w)}{(z-w)^2} + \frac{\partial\phi(w)}{z-w}$$

$$\begin{aligned}\delta\phi(z) &= \oint_{c_z} \frac{dw}{2\pi i} v(w)T(w)\phi(z) \sim \oint_{c_z} \frac{dw}{2\pi i} v(w) \left(\frac{h\phi(z)}{(w-z)^2} + \frac{\partial\phi(z)}{w-z} \right) \\ &= h\partial v(z)\phi(z) + v(z)\partial\phi(z)\end{aligned}$$

$$T(z)T(w) \sim \frac{c/2}{(z-w)^4} + \frac{2T(w)}{(z-w)^2} + \frac{\partial T(w)}{z-w}$$

$$\delta T = 2\partial v T + v\partial T + \frac{c}{12}\partial^3 v$$

$$T'(w) = \left(\frac{dz}{dw}\right)^{-2} \left(T(z) - \frac{c}{12}S(w,z)\right) = \left(\frac{dz}{dw}\right)^{-2} T(z) + \frac{c}{12}S(z,w)$$

$$S(w,z) = \frac{w^{(3)}}{w'} - \frac{3}{2}\left(\frac{w''}{w'}\right)^2$$

$$S(u,z) = S(w,z) + \left(\frac{dw}{dz}\right)^2 S(u,w)$$

$$\begin{aligned}\delta T(z) &= \oint_{c_z} \frac{dw}{2\pi i} v(w)T(w)T(z) \sim \oint_{c_z} \frac{dw}{2\pi i} v(w) \left(\frac{c/2}{(z-w)^4} + \frac{2T(w)}{(z-w)^2} + \frac{\partial T(w)}{z-w} \right) \\ &= \frac{c}{2 \times 3!} \partial^3 v(z) + 2\partial v(z)T(z) + v(z)\partial T(z)\end{aligned}$$

Conjugaciones BPZ y Modo de Expansi3n.

$$t^\dagger = -i\tau^\dagger = t \Rightarrow \tau^\dagger = -\tau.$$

$$z \stackrel{\tau \rightarrow -\tau}{\rightarrow} e^{-\tau+i\sigma} = \frac{1}{z^*} = I(\bar{z}),$$

$$\mathcal{O}(z,\bar{z})^\dagger := (\bar{I} \circ \mathcal{O}(z,\bar{z}))^\dagger,$$

$$\mathcal{O}(z,\bar{z})^\dagger = \left[\frac{1}{\bar{z}^{2h} z^{2\bar{h}}} \mathcal{O}\left(\frac{1}{\bar{z}}, \frac{1}{z}\right) \right]^\dagger = \frac{1}{z^{2h} \bar{z}^{2\bar{h}}} \mathcal{O}^\dagger\left(\frac{1}{z}, \frac{1}{\bar{z}}\right)$$

$$(\lambda \mathcal{O}_1 \cdots \mathcal{O}_n)^\dagger = \lambda^* \mathcal{O}_n^\dagger \cdots \mathcal{O}_1^\dagger, \lambda \in \mathbb{C},$$

$$\mathcal{O}(z,\bar{z})^t := I^\pm \circ \mathcal{O}(z,\bar{z}) = \frac{(\mp 1)^{h+\bar{h}}}{z^{2h} \bar{z}^{2\bar{h}}} \mathcal{O}\left(\pm \frac{1}{z}, \pm \frac{1}{\bar{z}}\right)$$

$$(\lambda \mathcal{O}_1 \cdots \mathcal{O}_n)^t = \lambda \mathcal{O}_1^t \cdots \mathcal{O}_n^t, \lambda \in \mathbb{C}$$

$$1^\dagger = 1^t = 1$$



$$\mathcal{O}(z, \bar{z}) = \sum_{m,n} \frac{\mathcal{O}_{m,n}}{z^{m+h} \bar{z}^{n+\bar{h}}}$$

$$m+h \in \mathbb{Z} + \nu, n+\bar{h} \in \mathbb{Z} + \bar{\nu}, \nu, \bar{\nu} = \begin{cases} 0 & \text{periodic} \\ 1/2 & \text{anti-periodic} \end{cases}$$

$$\mathcal{O}(e^{2\pi i} z, \bar{z}) = e^{2\pi i \nu} \mathcal{O}(z, \bar{z}), \mathcal{O}(z, e^{2\pi i} \bar{z}) = e^{2\pi i \bar{\nu}} \mathcal{O}(z, \bar{z})$$

$$\nu, \bar{\nu} = \begin{cases} 0 & \text{NS} \\ 1/2 & \text{R} \end{cases}$$

$$\nu, \bar{\nu} = \begin{cases} 0 & \text{untwisted} \\ 1/2 & \text{twisted} \end{cases}$$

$$(\mathcal{O}^\dagger)_{-m,-n} = (\mathcal{O}_{m,n})^\dagger.$$

$$\mathcal{O}^\dagger = \mathcal{O} \Rightarrow (\mathcal{O}_{m,n})^\dagger = \mathcal{O}_{-m,-n}.$$

$$\phi(z) = \sum_{n \in \mathbb{Z} + h + \nu} \frac{\phi_n}{z^{n+h}}.$$

$$\phi_n = \oint_{C_0} \frac{dz}{2\pi i} z^{n+h-1} \phi(z)$$

$$\phi^\dagger = \phi \Rightarrow (\phi_n)^\dagger = \phi_{-n}.$$

$$\phi_n^t = (I^\pm \circ \phi)_n = (-1)^h (\pm 1)^n \phi_{-n}.$$

$$\begin{aligned} \phi_n^t &= (I^\pm \circ \phi)_n = \oint \frac{dz}{2\pi i} z^{n+h-1} I^\pm \circ \phi(z) \\ &= \oint \frac{dz}{2\pi i} z^{n+h-1} \left(\mp \frac{1}{z^2} \right)^h \phi\left(\pm \frac{1}{z}\right) \\ &= (\mp 1)^h \oint \frac{dz}{2\pi i} z^{n-h-1} \phi\left(\pm \frac{1}{z}\right) \\ &= (\mp 1)^h \oint \frac{dw}{2\pi i} \left(\pm \frac{1}{w}\right)^{n-h} w^{-1} \phi(w) \\ &= (\mp 1)^h (\pm 1)^{n-h} \oint \frac{dw}{2\pi i} w^{-n+h-1} \phi(w) \end{aligned}$$

$$\frac{dz}{z} = \mp \frac{dw}{w^2 z} = -\frac{dw}{w}$$

$$T(z) = \sum_{n \in \mathbb{Z}} \frac{L_n}{z^{n+2}}, L_n = \oint \frac{dz}{2\pi i} T(z) z^{n+1}$$

$$L_n^\dagger = L_{-n}.$$

$$[L_m, \phi_n] = (m(h-1) - n) \phi_{m+n}$$

$$[L_m, \phi(z)] = z^m (z\partial + (n+1)h) \phi(z).$$



$$[L_0, L_{-n}] = nL_{-n}, [L_0, \phi_{-n}] = n\phi_{-n}$$

$$L_n^\pm = L_n \pm \bar{L}_n$$

$$j(z) = \sum_n \frac{j_n}{z^{n+1}}$$

$$Q_L = j_0$$

Espacios de Hilbert y correspondencia operador – estado.

$$L_0|0\rangle = 0, L_{\pm 1}|0\rangle = 0$$

$$\langle \mathcal{O} \rangle := \langle 0 | \mathcal{O} | 0 \rangle.$$

$$|\mathcal{O}\rangle = \lim_{z, \bar{z} \rightarrow 0} \mathcal{O}(z, \bar{z})|0\rangle = \mathcal{O}(0,0)|0\rangle$$

$$|\phi\rangle = \lim_{z \rightarrow 0} \phi(z)|0\rangle = \phi(0)|0\rangle.$$

$$\forall n \geq -h+1: \phi_n|0\rangle = 0$$

$$|\phi\rangle = \phi_{-h}|0\rangle = \oint \frac{dz}{2\pi i} \frac{\phi(z)}{z} |0\rangle.$$

$$\phi(z)|0\rangle = e^{zL_{-1}}\phi(0)e^{-zL_{-1}}|0\rangle = e^{zL_{-1}}|\phi\rangle.$$

$$\forall n \geq -1: L_n|0\rangle = 0$$

$$\forall |\phi\rangle \in \mathcal{H}: \langle \Omega | L_0 | \Omega \rangle \leq \langle \phi | L_0 | \phi \rangle.$$

$$L_0|\Omega\rangle:=a_\Omega|\Omega\rangle,$$

$$\langle 0| = |0\rangle^\dagger = |0\rangle^t$$

$$\langle 0|L_0 = 0, \langle 0|L_{\pm 1} = 0.$$

$$\begin{aligned} \langle \mathcal{O}^\dagger | &= \lim_{w, \bar{w} \rightarrow 0} \langle 0 | \mathcal{O}(w, \bar{w})^\dagger = \lim_{w, \bar{w} \rightarrow 0} \frac{1}{w^{2h} \bar{w}^{2\bar{h}}} \langle 0 | \mathcal{O} \left(\frac{1}{\bar{w}}, \frac{1}{w} \right)^\dagger \\ &= \lim_{z, \bar{z} \rightarrow \infty} z^{2h} \bar{z}^{2\bar{h}} \langle 0 | \mathcal{O}^\dagger(z, \bar{z}) \\ &= \langle 0 | I \circ \mathcal{O}^\dagger(0,0) \end{aligned}$$

$$\begin{aligned} \langle \phi^\dagger | &= \lim_{w \rightarrow 0} \langle 0 | \phi(w)^\dagger = \lim_{w \rightarrow 0} \frac{1}{w^{2h}} \langle 0 | \phi^\dagger \left(\frac{1}{w} \right) \\ &= \lim_{z \rightarrow \infty} z^{2h} \langle 0 | \phi^\dagger(z) \\ &= \langle 0 | I \circ \phi^\dagger(0). \end{aligned}$$

$$\langle \phi^\dagger | = \langle 0 | (\phi^\dagger)_h.$$



$$\begin{aligned}\langle \phi | &:= \lim_{w \rightarrow 0} \langle 0 | \phi(w)^t \\ &= (\pm 1)^h \lim_{z \rightarrow \infty} z^{2h} \langle 0 | \phi(z) \\ &= \langle 0 | I^\pm \circ \phi(0)\end{aligned}$$

$$\langle \phi | = (\pm 1)^h \langle 0 | \phi_h.$$

$$\langle \phi^\dagger | = (\pm 1)^h \langle \phi |.$$

$$\forall n \leq h-1: \langle 0 | \phi_n = 0,$$

$$\forall n \leq 1: \langle 0 | L_n = 0.$$

$$\langle 0 | T(z) | 0 \rangle = 0.$$

$$\begin{aligned}\langle \phi_i^\dagger | \phi_j \rangle &= \langle 0 | \bar{I} \circ \phi_j(0) \phi_i(0) | 0 \rangle = \lim_{\substack{z \rightarrow \infty \\ w \rightarrow 0}} z^{2h_i} \langle 0 | \phi_i^\dagger(z) \phi_j(w) | 0 \rangle, \\ \langle \phi_i | \phi_j \rangle &= \langle 0 | I \circ \phi_j(0) \phi_i(0) | 0 \rangle = (\pm 1)^{h_i} \lim_{\substack{z \rightarrow \infty \\ w \rightarrow 0}} z^{2h_i} \langle 0 | \phi_i(z) \phi_j(w) | 0 \rangle.\end{aligned}$$

$$\langle \phi_i | \phi_j \rangle = \langle I \circ \phi_i(0) \phi_j(0) \rangle, \langle \phi_i^\dagger | \phi_j \rangle = \langle I \circ \phi_i^\dagger(0) \phi_j(0) \rangle.$$

$$\langle \phi_i | \phi_j(z) | \phi_k \rangle = (\pm 1)^{h_i} \lim_{w \rightarrow \infty} w^{2h_i} \langle \phi_i(w) \phi_j(z) \phi_k(0) \rangle.$$

$$\langle \phi_i^c | \phi_j \rangle = \delta_{ij}$$

Módulos Verma.

$$L_0 |\phi\rangle = h |\phi\rangle, \forall n \geq 1: L_n |\phi\rangle = 0.$$

$$|\phi_{\{n_i\}}\rangle := \prod_i L_{-n_i} |\phi\rangle,$$

$$L_0 = h + \sum_i n_i.$$

$$\langle 0 | : \mathcal{O} : | 0 \rangle = 0.$$

$$\langle \Omega | \star \mathcal{O}_\star^\star | \Omega \rangle = 0.$$

$$: A(z) B(w) : \stackrel{?}{=} A(z) B(w) - \langle A(z) B(w) \rangle.$$

$$: A(z) B(z) : \stackrel{?}{=} \lim_{w \rightarrow z} (A(z) B(w) - \langle A(z) B(w) \rangle).$$

$$: A(z) B(w) : := A(z) B(w) - A(z) B(w) = \sum_{n \in \mathbb{N}} (z-w)^n \{AB\}_{-n}(z)$$

$$: AB(z) : := A(z) B(z) := \lim_{w \rightarrow z} : A(z) B(w) : = \{AB\}_0(z)$$

$$: AB(z) : = \oint_{c_z} \frac{dw}{2\pi i} \frac{A(z) B(w)}{z-w}$$



$$AB(z): = AB(z):.$$

$$:AB(z): = \sum_m \frac{:AB:_m}{z^{m+h_A+h_B}}$$

$$:AB:_m = \sum_{n \leq -h_A} A_n B_{m-n} + \sum_{n > -h_A} B_{m-n} A_n$$

$$:AB(z): \neq BA(z):, :A(BC)(z): \neq (AB)C(z):.$$

$$A_1(z):A_2A_3(w): = A_1(z)A_2A_3(w): + \sqrt{A_1(z):A_2A_3(w):} \\ A_1(z):A_2A_3(w): = A_1(z)A_2(w):A_3(w): + A_1(z)A_3(w):A_2(w):.$$

$$A(z):B(w)^n: = nA(z)B(w):B(w)^{n-1}:$$

$$A(z):e^{B(w)}: = A(z)B(w):e^{B(w)}:$$

$$:e^{A(z)}:e^{B(w)}: = \exp(A(z)B(w)):e^{A(z)}e^{B(w)}:.$$

$$\prod_{i=1}^n :e^{A_i}: = \exp\left(\sum_{i=1}^n A_i\right): \exp \sum_{i<j} \langle A_i A_j \rangle,$$

$$\left\langle \prod_{i=1}^n :e^{A_i}: \right\rangle = \exp \sum_{i<j} \langle A_i A_j \rangle.$$

$$A(z):e^{B(w)}: = A(z) \sum_n \frac{1}{n!} :B(w)^n: = A(z)B(w) \sum_n \frac{1}{(n-1)!} :B(w)^{n-1}:.$$

$$:e^{A(z)}:e^{B(w)}: = \sum_{m,n} \frac{1}{m!n!} :A(z)^m::B(w)^n: \\ = \sum_{m,n,k} \frac{k!}{m!n!} \binom{m}{k} \binom{n}{k} (A(z)B(w))^k :A(z)^{m-k}::B(w)^{n-k}: \\ = \sum_{m,n,k} \frac{1}{k!(m-k)!(n-k)!} (A(z)B(w))^k :A(z)^{m-k}::B(w)^{n-k}:.$$

$${}^{\star}AB(z){}^{\star} = \sum_m \frac{{}^{\star}AB(z){}^{\star}_m}{z^{m+h_A+h_B}},$$

$${}^{\star}AB{}^{\star}_m = \sum_{n \leq 0} A_n B_{m-n} + \sum_{n > 0} B_{m-n} A_n.$$

$$:AB:_m = {}^{\star}AB{}^{\star}_m + \sum_{n=0}^{h_A-1} [B_{m+n}, A_{-n}].$$



$$\begin{aligned}
:AB:_m &= \sum_{n \leq -h_A} A_n B_{m-n} + \sum_{n > -h_A} B_{m-n} A_n \\
&= \sum_{n \geq h_A} A_{-n} B_{m+n} + \sum_{n > 0} B_{m-n} A_n + \sum_{n=0}^{h_A-1} B_{m+n} A_{-n} \\
&= \sum_{n \geq 0} A_{-n} B_{m+n} + \sum_{n > 0} B_{m-n} A_n + \sum_{n=0}^{h_A-1} [B_{m+n}, A_{-n}] \\
&= {}^*AB^*_m + \sum_{n=0}^{h_A-1} [B_{m+n}, A_{-n}].
\end{aligned}$$

$$\phi(z) = \left(\frac{L}{2\pi}\right)^h z^{-h} \phi_{\text{cyl}}(w)$$

$$\phi_{\text{cyl}} = \left(\frac{2\pi}{L}\right)^h \sum_{n \in \mathbb{Z}} \phi_n e^{-\frac{2\pi}{L}w} = \left(\frac{2\pi}{L}\right)^h \sum_{n \in \mathbb{Z}} \frac{\phi_n}{z^n}.$$

$$T_{\text{cyl}}(w) = \left(\frac{2\pi}{L}\right)^2 \left(T(z)z^2 - \frac{c}{24}\right)$$

$$(L_0)_{\text{cyl}} = L_0 - \frac{c}{24}$$

$$H = (L_0)_{\text{cyl}} + (\bar{L}_0)_{\text{cyl}} = L_0 + \bar{L}_0 - \frac{c + \bar{c}}{24}$$

Sistemas CFT – Escalar libre – Acción covariante.

$$S = \frac{\epsilon}{4\pi\ell^2} \int d^2x \sqrt{g} g^{\mu\nu} \partial_\mu X \partial_\nu X$$

$$\epsilon := \begin{cases} +1 & \text{spacelike} \\ -1 & \text{timelike} \end{cases}, \sqrt{\epsilon} := \begin{cases} +1 & \text{spacelike} \\ i & \text{timelike} \end{cases}$$

$$X(\tau,\sigma) \sim X(\tau,\sigma+2\pi).$$

$$T_{\mu\nu} = -\frac{\epsilon}{\ell^2} \left[\partial_\mu X \partial_\nu X - \frac{1}{2} g_{\mu\nu} (\partial X)^2 \right],$$

$$T^\mu_\mu = 0.$$

$$\Delta X = 0,$$

$$0 = \int dX \frac{\delta}{\delta X(\sigma)} \left(e^{-S[X]} X(\sigma') \right)$$

$$\langle \partial^2 X(\sigma) X(\sigma') \rangle = -2\pi\epsilon\ell^2 \delta^{(2)}(\sigma - \sigma').$$

$$\langle X(\sigma) X(\sigma') \rangle = -\frac{\epsilon\ell^2}{2} \ln |\sigma - \sigma'|^2.$$



$$\langle X(\sigma)X(\sigma') \rangle = G(r), r = |\sigma - \sigma'|.$$

$$\Delta G(r) = \frac{1}{r} \partial_r (r G'(r)).$$

$$-2\pi\epsilon\ell^2 = 2\pi \int_0^r dr' r' \times \frac{1}{r'} \partial_{r'} (r' G'(r')) = 2\pi r G'(r).$$

$$G'(r) = -\epsilon\ell^2 \ln r$$

$$\ln r = \frac{1}{2} \ln r^2 = \frac{1}{2} \ln |\sigma - \sigma'|^2.$$

$$X \longrightarrow X + a, a \in \mathbb{R}.$$

$$J^\mu := 2\pi i \epsilon \frac{\partial \mathcal{L}}{\partial (\partial_\mu X)} = \frac{i}{\ell^2} g^{\mu\nu} \partial_\nu X, \nabla_\mu J^\mu = 0.$$

$$p = \frac{1}{2\pi} \int d\sigma J^0 = \frac{i}{2\pi\ell^2} \int d\sigma \partial^0 X.$$

$$\tilde{J}^\mu := -i\epsilon^{\mu\nu} J_\nu = \frac{1}{\ell^2} \epsilon^{\mu\nu} \partial_\nu X$$

$$\nabla_\mu \tilde{J}^\mu \propto \epsilon^{\mu\nu} [\nabla_\mu, \nabla_\nu] X = 0$$

$$w = \frac{1}{2\pi} \int d\sigma \tilde{J}^0 = \frac{1}{2\pi\ell^2} \int_0^{2\pi} d\sigma \partial_1 X = \frac{1}{2\pi\ell^2} (X(\tau, 2\pi) - X(\tau, 0))$$

$$J^\mu_a = \frac{i}{2\pi\ell^2} \eta_{ab} \partial^\mu X^b$$

Acción en un plano complejo.

$$S = \frac{\epsilon}{2\pi\ell^2} \int dz d\bar{z} \partial_z X \partial_{\bar{z}} X$$

$$\partial_z \partial_{\bar{z}} X = 0.$$

$$X(z, \bar{z}) = X_L(z) + X_R(\bar{z})$$

$$X(z) := X_L(z), X(\bar{z}) := X_R(\bar{z})$$

$$J := J_z = \frac{i}{\ell^2} \partial_z X, \bar{J} := J_{\bar{z}} = \frac{i}{\ell^2} \partial_{\bar{z}} X$$

$$\bar{\partial} J = 0, \partial \bar{J} = 0$$

$$p = p_L + p_R, p_L = \frac{1}{2\pi i} \oint dz J, p_R = -\frac{1}{2\pi i} \oint d\bar{z} \bar{J}$$

$$\tilde{J}_z = \frac{i}{\ell^2} \partial_z X = J, \tilde{J}_{\bar{z}} = -\frac{i}{\ell^2} \partial_{\bar{z}} X = -\bar{J}$$



$$w = p_L - p_R$$

$$p_L = \frac{p+w}{2}, p_R = \frac{p-w}{2}$$

$$p^2 + w^2 = p_L^2 + p_R^2, 2pw = p_L^2 - p_R^2$$

$$T:=T_{zz}=-\frac{\epsilon}{\ell^2}\partial_zX\partial_zX,\bar{T}:=T_{\bar{z}\bar{z}}=-\frac{\epsilon}{\ell^2}\partial_{\bar{z}}X\partial_{\bar{z}}X,T_{z\bar{z}}=0$$

$$V_k(z,\bar{z})\mathrel{\mathop:}=e^{i\epsilon kX(z,\bar{z})}.$$

$$V_{k_L,k_R}(z,\bar{z})\mathrel{\mathop:}=e^{2i\epsilon(k_LX(z)+k_RX(\bar{z}))}.$$

OPE y Modos de Expansi3n.

$$X(z)X(w)\sim -\frac{\epsilon\ell^2}{2}\ln(z-w).$$

$$\partial X(z)X(w)\sim -\frac{\epsilon\ell^2}{2}\frac{1}{z-w},$$

$$\partial X(z)\partial X(w)\sim -\frac{\epsilon\ell^2}{2}\frac{1}{(z-w)^2}$$

$$T(z)\partial X(w)\sim \frac{\partial X(w)}{(z-w)^2}+\frac{\partial(\partial X(w))}{z-w}.$$

$$T(z)T(w)\sim \frac{1}{2}\frac{1}{(z-w)^4}+\frac{2T(w)}{(z-w)^2}+\frac{\partial T(w)}{z-w}$$

$$c=1.$$

$$h=n$$

$$T(z)\partial^nX(w)\sim \cdots+\frac{n\partial^nX(w)}{(z-w)^2}+\frac{\partial(\partial^nX(w))}{z-w}$$

$$T(z)\partial^2X(w)\sim \frac{2\partial X(w)}{(z-w)^3}+\frac{2\partial^2X}{(z-w)^2}+\frac{\partial(\partial^2X(w))}{z-w}.$$

$$J(z)V_k(w,\bar{w})\sim \frac{\ell^2k}{2}\frac{V_k(w,\bar{w})}{z-w}.$$

$$T(z)V_k(w,\bar{w})\sim \frac{h_kV_k(w,\bar{w})}{(z-w)^2}+\frac{\partial V_k(w,\bar{w})}{z-w}$$

$$(h_k,\bar{h}_k)=\left(\frac{\epsilon\ell^2k^2}{4},\frac{\epsilon\ell^2k^2}{4}\right),\Delta_k=\frac{\epsilon\ell^2k^2}{2},s_k=0$$

$$V_k(z,\bar{z})V_{k'}(w,,\bar{w})\sim \frac{V_{k+k'}(w,\bar{w})}{(z-w)^{-\epsilon k k'\ell^2/2}}$$



$$T(z)\partial X(w) = -\frac{\epsilon}{\ell^2} : \partial X(z) \partial X(z) : \partial X(w) \sim -\frac{2\epsilon}{\ell^2} : \partial X(z) \partial X(z) : \partial X(w) \sim \frac{\partial X(z)}{(z-w)^2}$$

$$\begin{aligned} T(z)\partial X(w) &= \frac{1}{\ell^4} : \partial X(z) \partial X(z) : : \partial X(w) \partial X(w) : \\ &\sim \frac{1}{\ell^4} [: \partial X(z) \partial X(z) : : \partial X(w) \partial X(w) : + : \partial X(z) \partial X(z) : : \partial X(w) \partial X(w) : \\ &\quad + : \partial X(z) \partial X(z) : : \partial X(w) \partial X(w) : + \text{perms}] \\ &\sim 2 \times \frac{1}{4} \frac{1}{(z-w)^4} - 4 \times \frac{1}{2\ell^2} \frac{1}{(z-w)^2} : \partial X(z) \partial X(w) : \\ &\sim \frac{1}{2} \frac{1}{(z-w)^4} - \frac{2}{\ell^2} \frac{1}{(z-w)^2} (: \partial X(w) \partial X(w) : + (z-w) : \partial^2 X(w) \partial X(w) :). \end{aligned}$$

$$\begin{aligned} T(z)\partial^n X(w) &\sim \partial_w^{n-1} \frac{\partial X(z)}{(z-w)^2} \\ &\sim n! \frac{\partial X(z)}{(z-w)^{n+1}} \\ &\sim \frac{n!}{(z-w)^{n+1}} \left(\dots + \frac{1}{(n-1)!} (z-w)^{n-1} \partial^{n-1} (\partial X(w)) \right. \\ &\quad \left. + \frac{1}{n!} (z-w)^n \partial^n (\partial X(w)) \right) \end{aligned}$$

$$\partial X(z) V_k(w, \bar{w}) \sim i\epsilon k \partial X(z) X(w) V_k(w, \bar{w}) \sim i\epsilon k \left(-\frac{\epsilon \ell^2}{2} \frac{1}{z-w} \right) V_k(w, \bar{w})$$

$$\begin{aligned} T(z) V_k(w, \bar{w}) &\sim -\frac{\epsilon}{\ell^2} : \partial X(z) \partial X(z) : : e^{i\epsilon k X(w, \bar{w})} : \\ &\sim \frac{i\epsilon k}{2} \frac{1}{z-w} \partial X(z) : e^{i\epsilon k X(w, \bar{w})} : - \frac{\epsilon}{\ell^2} \partial X(z) : \partial X(z) e^{i\epsilon k X(w, \bar{w})} : \\ &\sim \frac{i\epsilon k}{2} \frac{1}{z-w} (: \partial X(z) e^{i\epsilon k X(w, \bar{w})} : + \partial X(z) : e^{i\epsilon k X(w, \bar{w})} :) \\ &\quad + \frac{i\epsilon k}{2} \frac{\partial X(z) e^{i\epsilon k X(w, \bar{w})}}{z-w} \\ &\sim \frac{\epsilon k^2 \ell^2}{4} \frac{V_k(w, \bar{w})}{(z-w)^2} + i\epsilon k \frac{: \partial X(w) e^{i\epsilon k X(w, \bar{w})} :}{z-w}. \end{aligned}$$

$$\begin{aligned} V_k(z, \bar{z}) V_{k'}(w, \bar{w}) &\sim \exp(-kk' X(z, \bar{z}) X(w, \bar{w})) : e^{i\epsilon k X(z, \bar{z})} e^{i\epsilon k' X(w, \bar{w})} : \\ &\sim (z-w)^{\epsilon k k' \ell^2 / 2} V_{k+k'}(w, \bar{w}) \end{aligned}$$

$$\partial X = -i \sqrt{\frac{\ell^2}{2}} \sum_{n \in \mathbb{Z}} \alpha_n z^{-n-1}, \bar{\partial} X = -i \sqrt{\frac{\ell^2}{2}} \sum_{n \in \mathbb{Z}} \bar{\alpha}_n \bar{z}^{-n-1}$$

$$\alpha_n = i\oint \frac{dz}{2\pi i} z^{n-1} \partial X(z), \bar{\alpha}_n = i\oint \frac{d\bar{z}}{2\pi i} \bar{z}^{n-1} \bar{\partial} X(z).$$

$$X(z) = \frac{x_L}{2} - i \sqrt{\frac{\ell^2}{2}} \alpha_0 \ln z + i \sqrt{\frac{\ell^2}{2}} \sum_{n \neq 0} \frac{\alpha_n}{n} z^{-n}$$



$$X(\bar{z}) = \frac{x_R}{2} - i \sqrt{\frac{\ell^2}{2}} \bar{\alpha}_0 \ln \bar{z} + i \sqrt{\frac{\ell^2}{2}} \sum_{n \neq 0} \frac{\bar{\alpha}_n}{n} \bar{z}^{-n}$$

$$p_L = \frac{\alpha_0}{\sqrt{2\ell^2}}, p_R = \frac{\bar{\alpha}_0}{\sqrt{2\ell^2}}$$

$$X(z) = \frac{x_L}{2} - i \ell^2 p_L \ln z + i \sqrt{\frac{\ell^2}{2}} \sum_{n \neq 0} \frac{\alpha_n}{n} z^{-n},$$

$$p = \frac{1}{\sqrt{2\ell^2}}(\alpha_0 + \bar{\alpha}_0), w = \frac{1}{\sqrt{2\ell^2}}(\alpha_0 - \bar{\alpha}_0)$$

$$\alpha_0 = \sqrt{\frac{\ell^2}{2}}(p + w), \bar{\alpha}_0 = \sqrt{\frac{\ell^2}{2}}(p - w)$$

$$x_L = x + q, x_R = x - q$$

$$x = \frac{1}{2}(x_L + x_R), q = \frac{1}{2}(x_L - x_R)$$

$$X(z, \bar{z}) = x - i \frac{\ell^2}{2} \left(p \ln |z|^2 + w \ln \frac{z}{\bar{z}} \right) + i \sqrt{\frac{\ell^2}{2}} \sum_{n \neq 0} \frac{1}{n} (\alpha_n z^{-n} + \bar{\alpha}_n \bar{z}^{-n})$$

$$X(\tau, \sigma) = x - i \ell^2 p \tau + \ell^2 w \sigma + \dots$$

$$\begin{aligned} p_L &= \frac{1}{2\pi i} \oint dz J = \frac{i}{\ell^2} \frac{1}{2\pi i} \oint dz \partial X = \frac{i}{\ell^2} \frac{1}{2\pi i} \oint dz \partial X \\ &= \frac{1}{\sqrt{2\ell^2}} \frac{1}{2\pi i} \oint dz \sum_n \alpha_n z^{-n-1} = \frac{1}{\sqrt{2\ell^2}} \alpha_0. \end{aligned}$$

$$X(\tau, \sigma + 2\pi) \sim X(\tau, \sigma)$$

$$X(e^{2\pi i} z, e^{-2\pi i} \bar{z}) \sim X(z, \bar{z}).$$

$$X(e^{2\pi i} z, e^{-2\pi i} \bar{z}) = X(z, \bar{z}) - i \sqrt{\frac{\ell^2}{2}} (\alpha_0 - \bar{\alpha}_0)$$

$$\alpha_0 = \bar{\alpha}_0 \implies p_L = p_R = \frac{p}{2}, w = 0.$$

$$N_n = \frac{\epsilon}{n} \alpha_{-n} \alpha_n, \bar{N}_n = \frac{\epsilon}{n} \bar{\alpha}_{-n} \bar{\alpha}_n.$$

$$N = \sum_{n>0} n N_n.$$



$$L_m = \frac{\epsilon}{2} \sum_n : \alpha_n \alpha_{m-n}$$

$$m \neq 0: L_m = \frac{\epsilon}{2} \sum_{n \neq 0, m} : \alpha_n \alpha_{m-n} : + \epsilon \alpha_0 \alpha_m,$$

$$L_0 = \frac{\epsilon}{2} \sum_n : \alpha_n \alpha_{-n} : = N + \frac{\epsilon}{2} \alpha_0^2 = N + \epsilon \ell^2 p_L^2$$

$$\hat{L}_0 := N$$

$$\bar{L}_0 = \bar{N} + \epsilon \ell^2 p_R^2, \hat{\bar{L}}_0 := \bar{N}$$

$$L_0^+ = N + \bar{N} + \epsilon \ell^2 (p_L^2 + p_R^2) = N + \bar{N} + \frac{\epsilon \ell^2}{2} (p^2 + w^2)$$

$$L_0^- = N - \bar{N} + \epsilon \ell^2 (p_L^2 - p_R^2) = N - \bar{N} + \epsilon \ell^2 w p$$

Conmutadores.

$$[\alpha_m, \alpha_n] = \epsilon m \delta_{m+n,0}, [\bar{\alpha}_m, \bar{\alpha}_n] = \epsilon m \delta_{m+n,0}, [\alpha_m, \bar{\alpha}_n] = 0$$

$$[p, w] = [p, p] = [w, w] = 0, [p, \alpha_n] = [p, \bar{\alpha}_n] = [w, \alpha_n] = [w, \bar{\alpha}_n] = 0$$

$$[x_L, p_L] = i\epsilon, [x_R, p_R] = i\epsilon$$

$$[x, p] = [q, w] = i\epsilon, [x, w] = [q, p] = 0$$

$$[L_m, \alpha_n] = -n \alpha_{m+n}$$

$$[L_0, \alpha_{-n}] = n \alpha_{-n}$$

$$[N_m, \alpha_{-n}] = \alpha_{-m} \delta_{m,n}$$

$$|k\rangle := \lim_{z, \bar{z} \rightarrow 0} V_k(z, \bar{z}) |0\rangle = e^{i\epsilon k x} |0\rangle,$$

$$p|k\rangle = k|k\rangle.$$

$$\forall n > 0: \alpha_n |k\rangle = 0,$$

$$N_n |k\rangle = 0.$$

$$L_0^+ |k\rangle = 2\epsilon \ell^2 k^2 |k\rangle, L_0^- |k\rangle = 0,$$

$$\mathcal{F}(k) = \text{Span}\{|k; \{N_n\}\rangle\},$$

$$|k; \{N_n\}\rangle := \prod_{n \geq 1} \frac{(\alpha_{-n})^{N_n}}{\sqrt{n^{N_n} N_n!}} |k\rangle, N_n \in \mathbb{N}^*$$

$$\mathcal{H} = \int_{\mathbb{R}} dk \mathcal{F}(k)$$



$$\begin{aligned}\lim_{z, \bar{z} \rightarrow 0} e^{i\epsilon k X(z, \bar{z})} |0\rangle &= \lim_{z, \bar{z} \rightarrow 0} \exp i\epsilon k \left[x - i \frac{\ell^2}{2} p \ln |z|^2 + i \sqrt{\frac{\ell^2}{2}} \sum_{n \neq 0} \frac{1}{n} (\alpha_n z^{-n} + \bar{\alpha}_n \bar{z}^{-n}) \right] |0\rangle \\ &= \lim_{z, \bar{z} \rightarrow 0} \exp \left[i\epsilon k x - \epsilon k \sqrt{\frac{\ell^2}{2}} \sum_{n \neq 0} \frac{1}{n} (\alpha_n z^{-n} + \bar{\alpha}_n \bar{z}^{-n}) \right] |0\rangle.\end{aligned}$$

$$\begin{aligned}p|k\rangle &= \frac{1}{\ell^2} \frac{1}{2\pi i} \oint (dz i \partial X(z) + d\bar{z} i \bar{\partial} X(\bar{z})) V_k(0,0) |0\rangle \\ &= \frac{1}{\ell^2} \frac{1}{2\pi i} \oint \left(\frac{dz}{z} \frac{\ell^2 k}{2} + \frac{d\bar{z}}{\bar{z}} \frac{\ell^2 k}{2} \right) V_k(0,0) |0\rangle \\ &= k V_k(0,0) |0\rangle\end{aligned}$$

$$x^\dagger = x \, p^\dagger = p, \alpha_n^\dagger = \alpha_{-n}.$$

$$L_n^\dagger = L_{-n}$$

$$\langle k| = |k\rangle^\dagger = \langle 0| e^{-i\epsilon k x}, \langle k| p = \langle k| k.$$

$$\alpha_n^t = -(\pm 1)^n \alpha_{-n}$$

$$p^t = -p, \langle -k| = |k\rangle^t.$$

$$\langle k \mid k' \rangle = 2\pi \delta(k - k')$$

$$\langle k^c| = \frac{1}{2\pi} \langle k|.$$

$$|k\rangle^\dagger = -|k\rangle^t$$

Sistema Ghost.

Acción covariante y holomórfica o antiholomórfica.

$$S = \frac{1}{4\pi} \int d^2x \sqrt{g} g^{\mu\nu} b_{\mu\mu_1\cdots\mu_{\lambda-1}} \nabla_\nu c^{\mu_1\cdots\mu_{\lambda-1}}$$

$$b_{\mu_1\cdots\mu_n} \longrightarrow e^{-i\theta} b_{\mu_1\cdots\mu_n}, c^{\mu_1\cdots\mu_{n-1}} \longrightarrow e^{i\theta} c^{\mu_1\cdots\mu_{n-1}}.$$

$$b(z) \mathop{:}= b_{z\dots z}(z), \bar{b}(\bar{z}) \mathop{:}= b_{\bar{z}\dots\bar{z}}(\bar{z}), c(z) \mathop{:}= c^{z\dots z}(\bar{z}), \bar{c}(\bar{z}) \mathop{:}= c^{\bar{z}\dots\bar{z}}(z)$$

$$S=\frac{1}{2\pi}\int d^2z(b\bar{\partial}c+\bar{b}\partial\bar{c})$$

$$\partial \bar{b} = 0, \bar{\partial} b = 0, \partial \bar{c} = 0, \bar{\partial} c = 0$$

$$h(b)=\lambda, h(c)=1-\lambda, h(\bar{b})=\lambda, h(\bar{c})=1-\lambda,$$



$$\begin{aligned}
T &= -\lambda: b\partial c: + (1-\lambda): \partial bc: \\
&= -\lambda: \partial(bc): +: \partial bc: \\
&= (1-\lambda): \partial(bc): -: b\partial c: .
\end{aligned}$$

$$b(z)c(w) = -\epsilon c(w)b(z), b(z)b(w) = -\epsilon b(w)b(z), c(z)c(w) = -\epsilon c(w)c(z)$$

$$\epsilon = \begin{cases} +1 & \text{anticommuting} \\ -1 & \text{commuting} \end{cases}$$

$$\delta b = -ib, \delta c = ic, \delta \bar{b} = -i\bar{b}, \delta \bar{c} = i\bar{c}$$

$$j(z) = -: b(z)c(z):, \bar{j}(\bar{z}) = -: \bar{b}(\bar{z})\bar{c}(\bar{z}):$$

$$N_{\text{gh}} = N_{\text{gh},L} + N_{\text{gh},R}, N_{\text{gh},L} = \oint \frac{dz}{2\pi i} j(z), N_{\text{gh},R} = -\oint \frac{d\bar{z}}{2\pi i} \bar{j}(\bar{z})$$

$$N_{\text{gh}}(c) = 1, N_{\text{gh}}(b) = -1, N_{\text{gh}}(\bar{c}) = 1, N_{\text{gh}}(\bar{b}) = -1$$

$$\int d'b\,d'c\,\frac{\delta}{\delta b(z)}\big[b(w)\mathrm{e}^{-S[b,c]}\big]=0$$

$$\delta^{(2)}(z-w)+\frac{1}{2\pi}\langle b(w)\bar{\partial}c(z)\rangle=0$$

$$\langle c(z)b(w)\rangle=\frac{1}{z-w}$$

$$T(z)b(w) \sim \lambda \frac{b(w)}{(z-w)^2} + \frac{\partial b(w)}{z-w}$$

$$T(z)c(w) \sim (1-\lambda) \frac{c(w)}{(z-w)^2} + \frac{\partial c(w)}{z-w}$$

$$T(z)T(w) \sim \frac{c_\lambda/2}{(z-w)^4} + \frac{2T(w)}{(z-w)^2} + \frac{\partial T(w)}{z-w}$$

$$c_\lambda=2\epsilon(-1+6\lambda-6\lambda^2)=-2\epsilon(1+6\lambda(\lambda-1))$$

$$q_\lambda=\epsilon(1-2\lambda)$$

$$c_\lambda=\epsilon(1-3q_\lambda^2)$$

$$j(z)b(w) \sim -\frac{b(w)}{z-w}$$

$$j(z)c(w) \sim \frac{c(w)}{z-w}$$

$$j(z)\mathcal{O}(w) \sim N_{\text{gh}}(\mathcal{O})\frac{\mathcal{O}(w)}{z-w}$$

$$j(z)j(w) \sim \frac{\epsilon}{(z-w)^2}$$



$$T(z)j(w) \sim \frac{q_\lambda}{(z-w)^3} + \frac{j(w)}{(z-w)^2} + \frac{\partial j(w)}{z-w}.$$

$$j(z) = \frac{dw}{dz} j'(w) + \frac{q_\lambda}{2} \frac{d}{dz} \ln \frac{dw}{dz} = \frac{dw}{dz} j'(w) + \frac{q_\lambda}{2} \frac{\partial_z^2 w}{\partial_z w}.$$

$$j(z) = \frac{dw}{dz} \left(j^{\text{cyl}}(w) - \frac{q_\lambda}{2} \right),$$

$$N_{\text{gh}} = N_{\text{gh}}^{\text{cyl}} - q_\lambda, N_{\text{gh},L} = N_{\text{gh},L}^{\text{cyl}} - \frac{q_\lambda}{2}, N_{\text{gh},R} = N_{\text{gh},R}^{\text{cyl}} - \frac{q_\lambda}{2}.$$

$$N^c - N^b = -\frac{\epsilon q_\lambda}{2} \chi_g = (1-2\lambda)(g-1)$$

$$\begin{aligned} T(z)b(w) &= (-\lambda: b(z)\partial c(z): + (1-\lambda): \partial b(z)c(z):)b(w) \\ &\sim -\lambda: b(z)\partial c(z): b(w) + (1-\lambda): \partial b(z)c(z): b(w) \\ &\sim -\lambda b(z)\partial_z \frac{1}{z-w} + (1-\lambda)\partial b(z)\frac{1}{z-w} \\ &\sim \lambda(b(w) + (z-w)\partial b(w))\frac{1}{(z-w)^2} + (1-\lambda)\frac{\partial b(w)}{z-w} \end{aligned}$$

$$\begin{aligned} T(z)c(w) &= (-\lambda: b(z)\partial c(z): + (1-\lambda): \partial b(z)c(z):)c(w) \\ &\sim \epsilon\lambda: \partial c(z)b(z): c(w) - \epsilon(1-\lambda): c(z)\partial b(z): c(w) \\ &\sim \lambda\frac{\partial c(z)}{z-w} - (1-\lambda)c(z)\partial_z \frac{1}{z-w} \\ &\sim \lambda\frac{\partial c(w)}{z-w} + (1-\lambda)(c(w) + (z-w)\partial c(w))\frac{1}{(z-w)^2} \\ &\sim (1-\lambda)\frac{c(w)}{(z-w)^2} + \frac{\partial c(w)}{(z-w)^2} \end{aligned}$$

$$j(z)b(w) = -: b(z)c(z): b(w) \sim -: b(z)c(z): b(w) \sim -\frac{b(z)}{z-w} \sim -\frac{b(w)}{z-w}.$$

$$j(z)c(w) = -: b(z)c(z): c(w) \sim \epsilon: c(z)b(z): c(w) \sim \frac{c(z)}{z-w} \sim \frac{c(w)}{z-w}.$$

$$\begin{aligned} j(z)j(w) &=: b(z)c(z)::b(w)c(w): \\ &\sim: b(z)c(z)::b(w)c(w): +: b(z)c(z)::b(w)c(w): +: b(z)c(z)::b(w)c(w): \\ &\sim \frac{\epsilon}{(z-w)^2} + \frac{\epsilon: c(z)b(w):}{z-w} + \frac{: b(z)c(w):}{z-w} \sim \frac{\epsilon}{(z-w)^2}. \end{aligned}$$

$$b(z) = \sum_{n \in \mathbb{Z} + \lambda + \nu} \frac{b_n}{z^{n+\lambda}}, c(z) = \sum_{n \in \mathbb{Z} + \lambda + \nu} \frac{c_n}{z^{n+1-\lambda}},$$

$$b_n = \oint \frac{dz}{2\pi i} z^{n+\lambda-1} b(z), c_n = \oint \frac{dz}{2\pi i} z^{n-\lambda} c(z).$$

$$b \rightarrow -b, c \rightarrow -c.$$

$$N_n^b =: b_{-n}c_n:, N_n^c = \epsilon: c_{-n}b_n:$$



$$N^b = \sum_{n>0} n N_n^b, N^c = \sum_{n>0} n N_n^c$$

$$L_m = \sum_n (n - (1 - \lambda)m) : b_{m-n} c_n : = \sum_n (\lambda m - n) : b_n c_{m-n} :.$$

$$L_0 = - \sum_n n : b_n c_{-n} : = \sum_n n : b_{-n} c_n :.$$

$$j_m = - \sum_n : b_{m-n} c_n : = - \sum_n : b_n c_{m-n} :.$$

$$N_{\text{gh},L} = j_0 = - \sum_n : b_{-n} c_n :.$$

$$b_n^\pm = b_n \pm \bar{b}_n, c_n^\pm = \frac{1}{2}(c_n \pm \bar{c}_n).$$

$$b_n^- b_n^+ = 2b_n \bar{b}_n, c_n^- c_n^+ = \frac{1}{2} c_n \bar{c}_n$$

$$\begin{aligned} T &= -\lambda : b \partial c : + (1 - \lambda) : \partial b c : \\ &= \sum_{m,n} \left(\lambda : b_m c_n : \frac{n+1-\lambda}{z^{m+\lambda} z^{m+2-\lambda}} - (1-\lambda) : b_m c_n : \frac{m+\lambda}{z^{m+\lambda+1} z^{m+1-\lambda}} \right) \\ &= \sum_{m,n} (\lambda(n+1-\lambda) - (1-\lambda)(m+\lambda)) \frac{: b_m c_n :}{z^{m-n+2}} \\ &= \sum_{m,n} (\lambda(n+1-\lambda) - (1-\lambda)(m-n+\lambda)) \frac{: b_{m-n} c_n :}{z^{m+2}} \\ &= \sum_{m,n} (n-m+\lambda m) \frac{: b_{m-n} c_n :}{z^{m+2}} = \sum_m \frac{L_m}{z^{m+2}}. \end{aligned}$$

$$j = -: bc : = \sum_{m,n} \frac{: b_m c_n :}{z^{m+\lambda} z^{n+1-\lambda}} = \sum_{m,n} \frac{: b_{m-n} c_n :}{z^{m+1}} = \sum_m \frac{j_m}{z^{m+1}}.$$

$$[b_m, c_n]_\epsilon = \delta_{m+n,0}, [b_m, b_n]_\epsilon = 0, [c_m, c_n]_\epsilon = 0$$

$$[b_m^+, c_n^+]_\epsilon = \delta_{m+n}, [b_m^-, c_n^-]_\epsilon = \delta_{m+n}$$

$$[N_m^b, b_{-n}] = b_{-n} \delta_{m,n}, [N_m^c, c_{-n}] = c_{-n} \delta_{m,n}$$

$$[L_m, b_n] = (m(\lambda - 1) - n) b_{m+n}, [L_m, c_n] = -(m\lambda + n) c_{m+n}$$

$$[L_0, b_0] = 0, [L_0, c_0] = 0$$

$$[j_m, j_n] = m \delta_{m+n,0}$$

$$[L_m, j_n] = -n j_{m+n} + \frac{q_\lambda}{2} m(m+1) \delta_{m+n,0}$$



$$\begin{aligned}
[N_{\text{gh}}, b(w)] &= -b(w), [N_{\text{gh}}, c(w)] = c(w) \\
[b_m, c_n]_\epsilon &= \epsilon \oint_{c_0} \frac{dw}{2\pi i} w^{-1} \oint_{c_w} \frac{dz}{2\pi i} z^{-1} w^{n+\lambda} z^{m-\lambda+1} b(z) c(w) \\
&\sim \epsilon \oint_{c_0} \frac{dw}{2\pi i} w^{-1} \oint_{c_w} \frac{dz}{2\pi i} z^{-1} w^{n+\lambda} z^{m-\lambda+1} \frac{\epsilon}{z-w} \\
&= \oint_{c_0} \frac{dw}{2\pi i} w^{m+n-1} = \delta_{m+n,0}
\end{aligned}$$

$$[N_{\text{gh}}, b(w)] = \oint \frac{dz}{2\pi i} j(z) b(w) \sim -\oint \frac{dz}{2\pi i} \frac{b(w)}{z-w} = -b(w)$$

$$\forall n > -\lambda: b_n|0\rangle = 0, \forall n > \lambda - 1: c_n|0\rangle = 0$$

$$|n_1, \dots, n_{\lambda-1}\rangle = c_1^{n_1} \dots c_{\lambda-1}^{n_{\lambda-1}}|0\rangle$$

$$L_0|n_1, \dots, n_{\lambda-1}\rangle = -\left(\sum_{j=1}^{\lambda-1} j n_j\right)|n_1, \dots, n_{\lambda-1}\rangle$$

Estado de energía – momentum. Métrica Grassmann.

$$\{| \downarrow \rangle, | \uparrow \rangle\},$$

$$| \downarrow \rangle := c_1 \dots c_{\lambda-1}|0\rangle, | \uparrow \rangle := c_0 c_1 \dots c_{\lambda-1}|0\rangle.$$

$$|\Omega\rangle = \omega_\downarrow | \downarrow \rangle + \omega_\uparrow | \uparrow \rangle, \omega_\downarrow, \omega_\uparrow \in \mathbb{C}.$$

$$b_0 | \uparrow \rangle = | \downarrow \rangle, c_0 | \downarrow \rangle = | \uparrow \rangle, b_0 | \downarrow \rangle = 0, c_0 | \uparrow \rangle = 0.$$

$$\forall n > 0: b_n | \downarrow \rangle = b_n | \uparrow \rangle = 0, c_n | \downarrow \rangle = b_n | \downarrow \rangle = 0.$$

$$|0\rangle = b_{1-\lambda} \dots b_{-1} | \downarrow \rangle = b_{1-\lambda} \dots b_{-1} b_0 | \uparrow \rangle.$$

$$L_0 | \downarrow \rangle = a_\lambda | \downarrow \rangle, L_0 | \uparrow \rangle = a_\lambda | \uparrow \rangle,$$

$$a_\lambda = -\sum_{n=1}^{\lambda-1} n = -\frac{\lambda(\lambda-1)}{2} = \frac{c_\lambda}{24} + \frac{2}{24}.$$

$$\{| \downarrow \downarrow \rangle, | \uparrow \downarrow \rangle, | \downarrow \uparrow \rangle, | \uparrow \uparrow \rangle\},$$

$$| \downarrow \downarrow \rangle := c_1 \bar{c}_1 \dots c_{\lambda-1} \bar{c}_{\lambda-1} |0\rangle$$

$$| \uparrow \downarrow \rangle := c_0 | \downarrow \downarrow \rangle, | \downarrow \uparrow \rangle := \bar{c}_0 | \downarrow \downarrow \rangle, | \uparrow \uparrow \rangle := c_0 \bar{c}_0 | \downarrow \downarrow \rangle.$$

$$c_0 | \downarrow \downarrow \rangle = | \uparrow \downarrow \rangle, \bar{c}_0 | \downarrow \downarrow \rangle = | \downarrow \uparrow \rangle, c_0 | \downarrow \uparrow \rangle = -\bar{c}_0 | \uparrow \downarrow \rangle = | \uparrow \uparrow \rangle$$

$$b_0 | \uparrow \uparrow \rangle = | \downarrow \uparrow \rangle, \bar{b}_0 | \uparrow \uparrow \rangle = -| \uparrow \downarrow \rangle, b_0 | \uparrow \downarrow \rangle = \bar{b}_0 | \downarrow \uparrow \rangle = | \downarrow \downarrow \rangle,$$



$$\begin{aligned} b_0|\downarrow\downarrow\rangle &= \bar{b}_0|\downarrow\downarrow\rangle = 0, & c_0|\uparrow\downarrow\rangle &= \bar{b}_0|\uparrow\downarrow\rangle = 0, \\ b_0|\downarrow\uparrow\rangle &= \bar{c}_0|\downarrow\uparrow\rangle = 0, & c_0|\uparrow\uparrow\rangle &= \bar{c}_0|\uparrow\uparrow\rangle = 0. \end{aligned}$$

$$\{|\downarrow\downarrow\rangle, |+\rangle, |-\rangle, |\uparrow\uparrow\rangle\},$$

$$\begin{aligned} c_0^\pm|\downarrow\downarrow\rangle &= \frac{1}{2}|\pm\rangle, & c_0^\mp|\pm\rangle &= \pm|\uparrow\uparrow\rangle, \\ b_0^\pm|\pm\rangle &= \pm 2|\downarrow\downarrow\rangle, & b_0^\mp|\uparrow\uparrow\rangle &= \pm|\pm\rangle. \end{aligned}$$

$$\begin{aligned} b_0^+|\downarrow\downarrow\rangle &= b_0^-|\downarrow\downarrow\rangle = 0, & c_0^-|-\rangle &= b_0^+|-\rangle = 0, \\ c_0^+|+\rangle &= b_0^-|+\rangle = 0 & c_0^+|\uparrow\uparrow\rangle &= c_0^-|\uparrow\uparrow\rangle = 0, \end{aligned}$$

$$c_0^-c_0^+|\downarrow\downarrow\rangle = \frac{1}{2}|\uparrow\uparrow\rangle, b_0^+b_0^-|\uparrow\uparrow\rangle = 2|\downarrow\downarrow\rangle.$$

$$\begin{aligned} 2c_0^+|\pm\rangle &= (c_0 + \bar{c}_0)|\uparrow\downarrow\rangle \pm (c_0 + \bar{c}_0)|\downarrow\uparrow\rangle = \bar{c}_0|\uparrow\downarrow\rangle \pm c_0|\downarrow\uparrow\rangle = (-1 \pm 1)|\uparrow\uparrow\rangle \\ b_0^+|\pm\rangle &= (b_0 + \bar{b}_0)|\uparrow\downarrow\rangle \pm (b_0 + \bar{b}_0)|\downarrow\uparrow\rangle = b_0|\uparrow\downarrow\rangle \pm \bar{b}_0|\downarrow\uparrow\rangle = (1 \pm 1)|\downarrow\downarrow\rangle \\ 2c_0^\pm|\downarrow\downarrow\rangle &= (c_0 \pm \bar{c}_0)|\downarrow\downarrow\rangle = c_0|\downarrow\downarrow\rangle \pm \bar{c}_0|\downarrow\downarrow\rangle = |\uparrow\downarrow\rangle \pm |\downarrow\uparrow\rangle = |\pm\rangle \\ b_0^\pm|\uparrow\uparrow\rangle &= (b_0 \pm \bar{b}_0)|\uparrow\uparrow\rangle = b_0|\uparrow\uparrow\rangle \pm \bar{b}_0|\uparrow\uparrow\rangle = |\downarrow\uparrow\rangle \mp |\uparrow\downarrow\rangle = \mp|\mp\rangle \end{aligned}$$

$$L_0 = \sum_n n^* b_{-n} c_n^* + a_\lambda = N^b + N^c + a_\lambda$$

$$\hat{L}_0 = N^b + N^c.$$

$$L_m = \sum_n (n - (1 - \lambda)m)^* b_{m-n} c_n^* + a_\lambda \delta_{m,0}$$

$$\begin{aligned} N_{\text{gh},L} &= j_0 = \sum_n {}^* b_{-n} c_n^* - \left(\frac{q_\lambda}{2} + \frac{1}{2} \right) \\ &= \sum_{n>0} (N_n^c - N_n^b) + \frac{1}{2} (N_0^c - N_0^b) - \frac{q_\lambda}{2} \end{aligned}$$

$$j_m = \sum_n {}^* b_{m-n} c_n^* - \left(\frac{q_\lambda}{2} + \frac{1}{2} \right) \delta_{m,0}.$$

$$\hat{N}_{\text{gh},L} := \sum_{n>0} (N_n^c - N_n^b)$$

$$\begin{aligned} j_0|\downarrow\rangle &= (\lambda - 1)|\downarrow\rangle = \left(-\frac{q_\lambda}{2} - \frac{1}{2}\right)|\downarrow\rangle, \\ j_0|\uparrow\rangle &= \lambda|\uparrow\rangle = \left(-\frac{q_\lambda}{2} + \frac{1}{2}\right)|\uparrow\rangle. \end{aligned}$$

$$j_0|0\rangle = 0$$

$$j_0^{\text{cyl}}|\downarrow\rangle = -\frac{1}{2}|\downarrow\rangle, j_0^{\text{cyl}}|\uparrow\rangle = \frac{1}{2}|\uparrow\rangle.$$



$$\begin{aligned}
L_0 &= - \sum_n n: b_n c_{-n}: = - \sum_{n \leq -\lambda} n b_n c_{-n} + \epsilon \sum_{n > -\lambda} n c_{-n} b_n \\
&= \sum_{n \geq \lambda} n b_{-n} c_n + \epsilon \sum_{n > -\lambda} n c_{-n} b_n \\
&= \sum_{n \geq \lambda} n b_{-n} c_n + \epsilon \sum_{n > 0} n c_{-n} b_n + \epsilon \sum_{n=-\lambda+1}^0 n c_{-n} b_n \\
&= \sum_{n \geq \lambda} n b_{-n} c_n + \epsilon \sum_{n > 0} n c_{-n} b_n + \epsilon \sum_{n=0}^{\lambda-1} n b_{-n} c_n + a_\lambda \\
&= \sum_{n > 0} n b_{-n} c_n + \epsilon \sum_{n > 0} n c_{-n} b_n + a_\lambda \\
&= \sum_n n \star b_{-n} c_n^\star + a_\lambda \\
\sum_{n=-\lambda+1}^0 c_{-n} b_n &= - \sum_{n=0}^{\lambda-1} n c_n b_{-n} = - \sum_{n=0}^{\lambda-1} n (-\epsilon b_{-n} c_n + 1) = \epsilon \sum_{n=0}^{\lambda-1} n b_{-n} c_n + a_\lambda. \\
j_0 &= - \sum_n : b_{-n} c_n: = - \sum_{n \geq \lambda} b_{-n} c_n + \epsilon \sum_{n > -\lambda} c_{-n} b_n \\
&= - \sum_{n \geq \lambda} b_{-n} c_n + \epsilon \sum_{n > 0} c_{-n} b_n + \epsilon \sum_{n=1}^{\lambda-1} c_n b_{-n} + \epsilon c_0 b_0 \\
&= - \sum_{n \geq \lambda} b_{-n} c_n + \epsilon \sum_{n > 0} c_{-n} b_n - \sum_{n=1}^{\lambda-1} b_{-n} c_n + \epsilon(\lambda - 1) + \epsilon c_0 b_0 \\
&= - \sum_{n > 0} b_{-n} c_n + \epsilon \sum_{n > 0} c_{-n} b_n + \epsilon(\lambda - 1) + \epsilon c_0 b_0. \\
\epsilon(\lambda - 1) &= -\frac{q_\lambda}{2} - \frac{\epsilon}{2}. \\
\epsilon c_0 b_0 + \epsilon(\lambda - 1) &= \frac{\epsilon}{2} c_0 b_0 + \frac{1}{2} (-b_0 c_0 + \epsilon) + \epsilon(\lambda - 1) \\
&= \frac{1}{2} (\epsilon c_0 b_0 - b_0 c_0) + \epsilon \left(\lambda - \frac{1}{2} \right).
\end{aligned}$$

Estructura del espacio de Hilbert. Métrica Grassmann odd.

$$\begin{aligned}
\mathcal{H}_{\text{gh}} &= \mathcal{H}_{\text{gh},0} \oplus c_0 \mathcal{H}_{\text{gh},0}, \mathcal{H}_{\text{gh},0} := \mathcal{H}_{\text{gh}} \cap \ker b_0 \\
\mathcal{H}_{\text{gh},0} &= \text{Span}\{|\downarrow; \{N_n^b\}; \{N_n^c\}\rangle\}, \\
|\downarrow; \{N_n^b\}; \{N_n^c\}\rangle &= \prod_{n \geq 1} (b_{-n})^{N_n^b} (c_{-n})^{N_n^c} |\downarrow\rangle, N_n^b, N_n^c \in \mathbb{N}^* \\
\psi &= \psi_\downarrow + \psi_\uparrow, \psi_\downarrow \in \mathcal{H}_{\text{gh},0}, \psi_\uparrow \in c_0 \mathcal{H}_{\text{gh},0} \\
\mathcal{H}_{\text{gh}} &= \mathcal{H}_{\text{gh},0} \oplus c_0 \mathcal{H}_{\text{gh},0} \oplus \bar{c}_0 \mathcal{H}_{\text{gh},0} \oplus c_0 \bar{c}_0 \mathcal{H}_{\text{gh},0}
\end{aligned}$$



$$\mathcal{H}_{\text{gh},0} := \mathcal{H}_{\text{gh}} \cap \ker b_0 \cap \ker \bar{b}_0.$$

$$|\downarrow\downarrow; \{N_n^b\}; \{N_n^c\}; \{\bar{N}_n^b\}; \{\bar{N}_n^c\}\rangle = \prod_{n \geq 1} (b_{-n})^{N_n^b} (\bar{b}_{-n})^{\bar{N}_n^b} (c_{-n})^{N_n^c} (\bar{c}_{-n})^{\bar{N}_n^c} |\downarrow\downarrow\rangle, \\ N_n^b, \bar{N}_n^b, N_n^c, \bar{N}_n^c \in \mathbb{N}^*.$$

$$\psi = \psi_{\downarrow\downarrow} + \psi_{\uparrow\downarrow} + \psi_{\downarrow\uparrow} + \psi_{\uparrow\uparrow},$$

$$\mathcal{H}_{\text{gh}} = \mathcal{H}_{\text{gh},0} \oplus c_0^+ \mathcal{H}_{\text{gh},0} \oplus c_0^- \mathcal{H}_{\text{gh},0} \oplus c_0^- c_0^+ \mathcal{H}_{\text{gh},0}$$

$$\mathcal{H}_{\text{gh},0} := \mathcal{H}_{\text{gh}} \cap \ker b_0^- \cap \ker b_0^+$$

$$\mathcal{H}_{\text{gh},\pm} := \mathcal{H}_{\text{gh}} \cap \ker b_0^\pm = \mathcal{H}_{\text{gh},0} \oplus c_0^\mp \mathcal{H}_{\text{gh},0},$$

$$\mathcal{H}_{\text{gh}} = \mathcal{H}_{\text{gh},\pm} \oplus c_0^\pm \mathcal{H}_{\text{gh},\pm}.$$

$$b_n^\dagger = \epsilon b_{-n}, c_n^\dagger = c_{-n}.$$

$$b_n^t = (-1)^\lambda b_{-n}, c_n^t = (-1)^{1-\lambda} c_{-n},$$

$$|\downarrow\rangle^\dagger = \langle 0|c_{1-\lambda} \cdots c_{-1}, |\uparrow\rangle^\dagger = \langle 0|c_{1-\lambda} \cdots c_{-1}c_0.$$

$$\langle \downarrow| := |\downarrow\rangle^t = (-1)^{(1-\lambda)^2} \langle 0|c_{-1} \cdots c_{1-\lambda} \\ \langle \uparrow| := |\uparrow\rangle^t = (-1)^{\lambda(1-\lambda)} \langle 0|c_0 c_{-1} \cdots c_{1-\lambda}.$$

$$\langle \downarrow| = (-1)^{a_\lambda + (1-\lambda)(2-\lambda)} |\downarrow\rangle^\dagger, \langle \uparrow| = (-1)^{a_\lambda} |\uparrow\rangle^\dagger,$$

$$\langle \downarrow| = (-1)^{(1-\lambda)^2 + \frac{1}{2}(2-\lambda)(1-\lambda)} |\downarrow\rangle^\dagger = (-1)^{-a_\lambda + (1-\lambda)(2-\lambda)} |\downarrow\rangle^\dagger$$

$$\sum_{i=1}^{\lambda-2} i = \frac{1}{2}(2-\lambda)(1-\lambda) = -a_\lambda + 1 - \lambda.$$

$$\langle \uparrow| = (-1)^{\lambda(1-\lambda) - \frac{1}{2}\lambda(1-\lambda)} |\uparrow\rangle^\dagger = (-1)^{\frac{1}{2}\lambda(1-\lambda)} |\uparrow\rangle^\dagger.$$

$$\langle \uparrow|b_0 = \langle \downarrow|, \langle \downarrow|c_0 = \langle \uparrow|, \langle \downarrow|b_0 = 0, \langle \uparrow|c_0 = 0.$$

$$|\downarrow\rangle = \begin{pmatrix} 0 \\ 1 \end{pmatrix}, |\uparrow\rangle = \begin{pmatrix} 1 \\ 0 \end{pmatrix},$$

$$b_0 = \begin{pmatrix} 0 & 0 \\ 1 & 0 \end{pmatrix}, c_0 = \begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix}$$

$$\langle \downarrow\downarrow\rangle = \langle \uparrow\uparrow\rangle = 0$$

$$\langle \uparrow\downarrow\rangle = \langle \downarrow|c_0|\downarrow\rangle = \langle 0|c_{1-\lambda} \cdots c_{-1}c_0c_1 \cdots c_{\lambda-1}|0\rangle = 1.$$

$$\langle 0^c| = \langle 0|c_{1-\lambda} \cdots c_{-1}c_0c_1 \cdots c_{\lambda-1}, \langle \downarrow^c| = \langle \uparrow|.$$



Cuantización BRST. Reparametrización e invariancia.

$$c = c_m + c_{\text{gh}} = c_m - 26, T(z) = T^m(z) + T^{\text{gh}}(z), \mathcal{H} = \mathcal{H}_m \otimes \mathcal{H}_{\text{gh}}.$$

$$Q_B |\psi\rangle = 0$$

$$|\psi\rangle \sim |\psi\rangle + Q_B |\Lambda\rangle.$$

$$\begin{aligned} j_B(z) &=: c(z) \left(T^m(z) + \frac{1}{2} T^{\text{gh}}(z) \right) : + \kappa \partial^2 c(z) \\ &= c(z) T^m(z) + : b(z) c(z) \partial c(z) : + \kappa \partial^2 c(z) \end{aligned}$$

$$Q_B = Q_{B,L} + Q_{B,R}, Q_{B,L} = \oint \frac{dz}{2\pi i} j_B(z), Q_{B,R} = \oint \frac{d\bar{z}}{2\pi i} \bar{j}_B(\bar{z}).$$

$$T(z)j_B(w) \sim \left(\frac{c_m}{2} - 4 - 6\kappa \right) \frac{c(w)}{(z-w)^4} + (3 - 2\kappa) \frac{\partial c(w)}{(z-w)^3} + \frac{j_B(w)}{(z-w)^2} + \frac{\partial j_B(w)}{z-w}.$$

$$c_m = 26, \kappa = \frac{3}{2}.$$

$$T(z)j_B(w) \sim \frac{j_B(w)}{(z-w)^2} + \frac{\partial j_B(w)}{z-w}.$$

$$\begin{aligned} j_B(z)b(w) &\sim \frac{2\kappa}{(z-w)^3} + \frac{j(w)}{(z-w)^2} + \frac{T(w)}{z-w} \\ j_B(z)c(w) &\sim \frac{: c(w) \partial c(w) :}{z-w} \end{aligned}$$

$$j_B(z)\phi(w) \sim h \frac{c(w)\phi(w)}{(z-w)^2} + \frac{: h \partial c(w) \phi(w) + c(w) \partial \phi(w) :}{z-w}$$

$$j_B(z)j(w) \sim \frac{2\kappa + 1}{(z-w)^3} - \frac{2\partial c(w)}{(z-w)^2} - \frac{j_B(w)}{z-w}$$

$$j_B(z)j_B(w) \sim -\frac{c_m - 18 : c(w) \partial c(w) :}{2 (z-w)^3} - \frac{c_m - 18 : c(w) \partial^2 c(w) :}{4 (z-w)^2} - \frac{c_m - 26 : c(w) \partial^3 c(w) :}{12 (z-w)}.$$

$$\begin{aligned} Q_B &= \sum_m : c_m \left(L_{-m}^m + \frac{1}{2} L_{-m}^{\text{gh}} \right) : \\ &= \sum_m c_{-m} L_m^m + \frac{1}{2} \sum_{m,n} (n-m) : c_{-m} c_{-n} b_{m+n} : \\ Q_B &= \sum_m {}^* c_m \left(L_{-m}^m + \frac{1}{2} L_{-m}^{\text{gh}} \right) {}^* - \frac{c_0}{2} \\ &= \sum_n c_m L_{-m}^m + \frac{1}{2} \sum_{m,n} (n-m) {}^* c_{-m} c_{-n} b_{m+n} {}^* - c_0, \end{aligned}$$

$$Q_B = c_0 L_0 - b_0 M + \hat{Q}_B$$



$$\hat{Q}_B = \sum_{m \neq 0} c_{-m} L_m^m - \frac{1}{2} \sum_{\substack{m, n \neq 0 \\ m+n \neq 0}} (m-n)^* c_{-m} c_{-n} b_{m+n}^*,$$

$$M = \sum_{m \neq 0} m c_{-m} c_m$$

$$[L_0, M] = [\hat{Q}_B, M] = [\hat{Q}_B, L_0] = 0, \hat{Q}_B^2 = L_0 M$$

$$\{Q_B, b(z)\} = T(z)$$

$$\{Q_B, c(z)\} = c(z) \partial c(z)$$

$$[Q_B, \phi(z)] = h \partial c(z) \phi(z) + c(z) \partial \phi(z)$$

$$[Q_B, \phi(z)] = \partial(c(z) \phi(z))$$

$$\{Q_B, c(z) \phi(z)\} = (1-h) c(z) \partial c(z) \phi(z)$$

$$[Q_B, j(z)] = -j_B(z)$$

$$[N_{\text{gh}}, Q_B] = Q_B$$

$$\{Q_B, Q_B\} = 0$$

$$[Q_B, T(z)] = 0$$

$$c_m = 26.$$

$$L_n = \{Q_B, b_n\}$$

$$[Q_B, L_n] = 0$$

Cohomología BRST para sectores holomórficos y antiholomórficos.

$$k_{\parallel}^2 = \epsilon_0 (k^0)^2 + (k^1)^2.$$

$$\mathcal{H}:=\mathcal{H}_{\parallel}\otimes\mathcal{H}_{\perp},\mathcal{H}_{\parallel}:=\int\mathrm{d}k^0\mathcal{F}_0(k^0)\otimes\int\mathrm{d}k^1\mathcal{F}_1(k^1)\otimes\mathcal{H}_{\text{gh}}$$

$$|\psi\rangle=|\psi_{\parallel}\rangle\otimes|\psi_{\perp}\rangle,$$

$$|\psi_{\parallel}\rangle=c_0^{N_0^c}\prod_{m>0}(\alpha_{-m}^0)^{N_m^0}(\alpha_{-m}^1)^{N_m^1}(b_{-m})^{N_m^b}(c_{-m})^{N_m^c}|k^0,k^1,\downarrow\rangle$$

$$|k^0,k^1,\downarrow\rangle:=|k^0\rangle\otimes|k^1\rangle\otimes|\downarrow\rangle, N_m^0, N_m^1\in\mathbb{N}, N_m^b, N_m^c=0,1.$$

$$\mathcal{H}_0=\mathcal{H}\cap\ker b_0$$

$$\mathcal{H}=\mathcal{H}_0\oplus c_0\mathcal{H}_0.$$

$$L_0=L_0^m+L_0^{\text{gh}}=(L_0^m-1)+N^b+N^c,$$

$$L_0=\left(L_0^\perp-m_{\parallel,L}^2\ell^2-1\right)+\hat{L}_0^\parallel$$



$$m_{\parallel,L}^2 = -p_{\parallel,L}^2, \hat{L}_0^{\parallel} = N^0 + N^1 + N^b + N^c \in \mathbb{N}.$$

$$\text{on-shell: } L_0|\psi\rangle = 0$$

$$\mathcal{H}_{\text{abs}}(Q_B) := \{|\psi\rangle \in \mathcal{H} | Q_B|\psi\rangle = 0, \nexists |\chi\rangle \in \mathcal{H} | |\psi\rangle = Q_B|\chi\rangle\}.$$

$$\{Q,\Delta\}=1$$

$$|\psi\rangle = \{Q_B,\Delta\}|\psi\rangle = Q_B(\Delta|\psi\rangle)$$

$$\Delta\!:=\frac{b_0}{L_0}$$

$$L_0=\{Q_B,b_0\}$$

$$Q_B|\psi\rangle=0, L_0|\psi\rangle\neq 0$$

$$|\psi\rangle=Q_B\left(\frac{b_0}{L_0}|\psi\rangle\right)$$

$$\mathcal{H}_{\text{abs}}(Q_B)\subset \ker L_0$$

$$|\psi\rangle=\frac{L_0}{L_0}|\psi\rangle=\frac{1}{L_0}\{Q_B,b_0\}|\psi\rangle=\frac{1}{L_0}Q_B(b_0|\psi\rangle)$$

$$|\psi\rangle\in\ker L_0\colon P_0|\psi\rangle=|\psi\rangle, |\psi\rangle\in(\ker L_0)^\perp\colon P_0|\psi\rangle=0.$$

$$\{Q_B,\Delta(1-P_0)\}=(1-P_0).$$

$$\mathcal{H}_0:=\mathcal{H}\cap\ker b_0=\mathcal{H}_m\otimes\mathcal{H}_{\text{gh},0}$$

$$|\psi\rangle\in\mathcal{H}_0\implies b_0|\psi\rangle=0.$$

$$b_0|\psi\rangle=Q_B|\psi\rangle=0, L_0|\psi\rangle\neq 0\implies |\psi\rangle=0.$$

$$\mathcal{H}_{\text{rel}}(Q_B) := \mathcal{H}_0(Q_B) = \{|\psi\rangle \in \mathcal{H}_0 | Q_B|\psi\rangle = 0, \nexists |\chi\rangle \in \mathcal{H} | |\psi\rangle = Q_B|\chi\rangle\}.$$

$$Q_B|\psi\rangle=0, L_0|\psi\rangle=0, b_0|\psi\rangle=0$$

$$Q_B=c_0L_0-b_0M+\hat{Q}_B,\hat{Q}_B^2=L_0M$$

$$|\psi\rangle\in\mathcal{H}_0\cap\ker L_0\implies Q_B|\psi\rangle=\hat{Q}_B|\psi\rangle,\hat{Q}_B^2|\psi\rangle=0$$

$$\mathcal{H}_0(Q_B)=\mathcal{H}_0(\hat{Q}_B).$$

$$X_L^\pm=\frac{1}{\sqrt{2}}\bigg(X_L^0\pm\frac{\mathrm{i}}{\sqrt{\epsilon_0}}X_L^1\bigg)$$

$$\alpha_n^\pm = \frac{1}{\sqrt{2}} \left(\alpha_n^0 \pm \frac{i}{\sqrt{\epsilon_0}} \alpha_n^1 \right), n \neq 0,$$

$$x_L^\pm = \frac{1}{\sqrt{2}} \left(x_L^0 \pm \frac{i}{\sqrt{\epsilon_0}} x_L^1 \right), p_L^\pm = \frac{1}{\sqrt{2}} \left(p_L^0 \pm \frac{i}{\sqrt{\epsilon_0}} p_L^1 \right),$$

$$[\alpha_m^+, \alpha_n^-] = \epsilon_0 m \delta_{m+n,0}, [x_L^\pm, p_L^\mp] = i\epsilon_0.$$

$$2p_L^+ p_L^- = (p_L^0)^2 + \epsilon_0 (p_L^1)^2 = \epsilon_0 p_{\parallel,L}^2,$$

$$x^+ p^- + x^- p^+ = x^0 p^0 + \epsilon_0 x^1 p^1,$$

$$\sum_n \alpha_n^+ \alpha_{m-n}^- = \frac{1}{2} \sum_n (\alpha_n^0 \alpha_{m-n}^0 + \epsilon_0 \alpha_n^1 \alpha_{m-n}^1).$$

$$N_n^\pm = \frac{\epsilon_0}{n} \alpha_{-n}^\pm \alpha_n^\mp, N^\pm = \sum_{n>0} n N_n^\pm.$$

$$N^+ + N^- = N^0 + N^1.$$

$$L_0 = (L_0^\perp - m_{\parallel,L}^2 \ell^2 - 1) + \hat{L}_0^\parallel$$

$$m_{\parallel,L}^2 = -2\epsilon_0 p_L^+ p_L^-, \hat{L}_0^\parallel = N^+ + N^- + N^b + N^c$$

$$L_m^0 + L_m^1 = \epsilon_0 \sum_n : \alpha_n^+ \alpha_{m-n}^- : = \epsilon_0 \sum_{n \neq 0, m} : \alpha_n^+ \alpha_{m-n}^- : + \epsilon_0 (\alpha_0^- \alpha_m^+ + \alpha_m^+ \alpha_0^-).$$

$$[\alpha_m^+, \alpha_n^\pm] = \frac{1}{2} \left[\left(\alpha_m^0 + \frac{i}{\sqrt{\epsilon_0}} \alpha_m^1 \right), \left(\alpha_n^0 \pm \frac{i}{\sqrt{\epsilon_0}} \alpha_n^1 \right) \right]$$

$$= \frac{1}{2} ([\alpha_m^0, \alpha_n^0] \mp \frac{1}{\epsilon_0} [\alpha_m^1, \alpha_n^1]) = \frac{\epsilon_0}{2} m \delta_{m+n,0} (1 \mp 1),$$

$$[x_L^-, p_L^\pm] = \frac{1}{2} \left[\left(x_L^0 - \frac{i}{\sqrt{\epsilon_0}} x_L^1 \right), \left(p_L^0 \pm \frac{i}{\sqrt{\epsilon_0}} p_L^1 \right) \right]$$

$$= \frac{1}{2} ([x_L^0, p_L^0] \pm \epsilon_0 [x_L^1, p_L^1]) = \frac{\epsilon_0}{2} (1 \pm 1)$$

$$\sum_n \alpha_n^+ \alpha_{m-n}^- = \frac{1}{2} \sum_n \left(\alpha_n^0 + \frac{i}{\sqrt{\epsilon_0}} \alpha_n^1 \right) \left(\alpha_{m-n}^0 - \frac{i}{\sqrt{\epsilon_0}} \alpha_{m-n}^1 \right)$$

$$= \frac{1}{2} \sum_n \left(\alpha_n^0 \alpha_{m-n}^0 + \epsilon_0 \alpha_n^1 \alpha_{m-n}^1 + \frac{i}{\sqrt{\epsilon_0}} (\alpha_{m-n}^0 \alpha_n^1 - \alpha_n^0 \alpha_{m-n}^1) \right)$$

$$N^0 + N^1 = \sum_n n (N_n^0 + N_n^1) = \sum_n n (N_n^+ + N_n^-) = N^+ + N^-.$$

$$\hat{Q}_B = \sum_{m \neq 0} c_{-m} \left(L_m^\perp + \epsilon_0 \sum_n \alpha_n^+ \alpha_{m-n}^- \right) + \frac{1}{2} \sum_{m,n} (n-m) : c_{-m} c_{-n} b_{m+n} :$$

$$\text{deg} = N^+ - N^- + \hat{N}^c - \hat{N}^b$$



$$\forall m \neq 0: \deg(\alpha_m^+) = \deg(c_m) = 1, \deg(\alpha_m^-) = \deg(b_m) = -1$$

$$\hat{Q}_B = Q_0 + Q_1 + Q_2, \deg(Q_j) = j$$

$$Q_1 = \sum_{m \neq 0} c_{-m} L_m^\perp + \sum_{\substack{m, n \neq 0 \\ m+n \neq 0}} {}^* c_{-m} \left(\epsilon_0 \alpha_n^+ \alpha_{m-n}^- + \frac{1}{2} (m-n) c_{-m} b_{m+n} \right) {}^*$$

$$Q_0 = \sum_{n \neq 0} \alpha_0^+ c_{-n} \alpha_n^-, Q_2 = \sum_{n \neq 0} \alpha_0^- c_{-n} \alpha_n^+$$

$$Q_0^2 = Q_2^2 = 0, \{Q_0, Q_1\} = \{Q_1, Q_2\} = 0, Q_1^2 + \{Q_0, Q_2\} = 0$$

$$\mathcal{H}_0(\hat{Q}_B) \simeq \mathcal{H}_0(Q_0)$$

$$\hat{\Delta} := \frac{B}{\hat{L}_0^\parallel}, B := \epsilon_0 \sum_{n \neq 0} \frac{1}{\alpha_0^+} \alpha_{-n}^+ b_n.$$

$$\hat{L}_0^\parallel = \{Q_0, B\} \Rightarrow \{Q_0, \hat{\Delta}\} = 1.$$

$$\hat{L}_0^\parallel |\psi\rangle = 0, \Rightarrow N^\pm |\psi\rangle = N^c |\psi\rangle = N^b |\psi\rangle = 0,$$

$$0 = Q_0 \hat{L}_0^\parallel |\psi\rangle = \hat{L}_0^\parallel Q_0 |\psi\rangle.$$

$$Q_0 |\psi\rangle = 0$$

$$L_0 = L_0^\perp - m_{\parallel,L}^2 \ell^2 - 1 = 0$$

$$\{Q_0, Q_1\} |\psi_0\rangle = 0 \Rightarrow Q_0(Q_1 |\psi_0\rangle) = 0.$$

$$Q_1 |\psi_0\rangle =: -Q_0 |\psi_1\rangle \Rightarrow |\psi_1\rangle = -\frac{B}{\hat{L}_0^\parallel} Q_1 |\psi_0\rangle.$$

$$Q_1 |\psi_0\rangle = \left\{ Q_0, \frac{B}{\hat{L}_0^\parallel} \right\} Q_1 |\psi_0\rangle = Q_0 \left(\frac{B}{\hat{L}_0^\parallel} Q_1 |\psi_0\rangle \right).$$

$$\{Q_0, Q_1\} |\psi_1\rangle = Q_0(Q_1 |\psi_1\rangle + Q_2 |\psi_0\rangle).$$

$$Q_1 |\psi_1\rangle + Q_2 |\psi_0\rangle = Q_0 |\psi_2\rangle, |\psi_2\rangle = -\frac{B}{\hat{L}_0^\parallel} (Q_1 |\psi_1\rangle + Q_2 |\psi_0\rangle)$$

$$\{Q_0, Q_1\} |\psi_1\rangle = Q_0 Q_1 |\psi_1\rangle - Q_1^2 |\psi_0\rangle = Q_0 Q_1 |\psi_1\rangle + \{Q_0, Q_2\} |\psi_0\rangle.$$

$$|\psi_{k+1}\rangle = -\frac{B}{\hat{L}_0^\parallel} (Q_1 |\psi_k\rangle + Q_2 |\psi_{k-1}\rangle)$$

$$|\psi\rangle = \sum_{k \in \mathbb{N}} |\psi_k\rangle.$$

$$\hat{Q}_B |\psi\rangle = 0.$$



$$N_{\text{gh}}(\psi) = N_{\text{gh}}(\psi_0) = 1$$

$$Q_1|\psi_0\rangle = Q_2|\psi_0\rangle = 0$$

$$\begin{aligned}\hat{Q}_B|\psi\rangle &= \sum_{k\in\mathbb{N}} \hat{Q}_B|\psi_k\rangle \\ &= Q_0|\psi_0\rangle + \underbrace{Q_1|\psi_0\rangle + Q_0|\psi_1\rangle}_{=0} + \underbrace{Q_2|\psi_0\rangle + Q_1|\psi_1\rangle + Q_0|\psi_2\rangle}_{=0} + \cdots \\ &= 0\end{aligned}$$

$$\mathcal{H}_{\text{abs}}(Q_B) = \mathcal{H}_{\text{rel}}(Q_B) \oplus c_0 \mathcal{H}_{\text{rel}}(Q_B)$$

$$\begin{aligned}|\psi\rangle &= |k^0,k^1,\downarrow\rangle\otimes|\psi_\perp\rangle, \quad |\psi_\perp\rangle\in\mathcal{H}_\perp \\ (L_0^\perp-m_{\parallel,L}^2\ell^2-1)|\psi\rangle &= 0, \quad p_{L,\parallel}^2=-m_{\parallel,L}^2\ell^2\end{aligned}$$

$$(L_0^m-1)|\psi\rangle=0, \forall n>0: L_n^m|\psi\rangle=0.$$

$$Q_B=c_0L_0-b_0M+\hat{Q}_B+\bar{c}_0\bar{L}_0-\bar{b}_0\bar{M}+\hat{\bar{Q}}_B$$

$$Q_B=c_0^+L_0^+-b_0^+M^++c_0^-L_0^--b_0^-M^-+\hat{Q}_B^+,$$

$$L_0^+=\left(L_0^{\perp+}-\frac{m_{\parallel}^2\ell^2}{2}\right)+\hat{L}_0^{\parallel+}, L_0^-=L_0^{\perp-}+\hat{L}_0^{\parallel-}$$

$$M^{\pm}:=\frac{1}{2}(M\pm\bar{M}).$$

$$L_0^+|\psi\rangle=L_0^-|\psi\rangle=0.$$

$$\hat{L}_0^{\parallel\pm}=N^0\pm\bar{N}^0+N^1\pm\bar{N}^1+N^b\pm\bar{N}^b+N^c\pm\bar{N}^c=0.$$

$$(L_0^m+\bar{L}_0^m-2)|\psi\rangle=0, (L_0^m-\bar{L}_0^m)|\psi\rangle=0$$

$$\forall n>0: L_n^m|\psi\rangle=\bar{L}_n^m|\psi\rangle=0$$

$$|\downarrow\downarrow\rangle=c(0)\bar{c}(0)|0\rangle=c_1\bar{c}_1|0\rangle.$$

$$\hat{L}_0|\psi_{\ell,\bar{\ell}}\rangle=\ell|\psi_{\ell,\bar{\ell}}\rangle,\hat{\bar{L}}_0|\psi_{\ell,\bar{\ell}}\rangle=\bar{\ell}|\psi_{\ell,\bar{\ell}}\rangle.$$

$$k^2=-m^2, m^2:=\frac{2}{\ell^2}(N+\bar{N}-2),$$

$$m^2\ell^2=-4<0.$$

$$\mathcal{V}(k,z,\bar{z})=c(z)\bar{c}(\bar{z})\mathrm{e}^{\mathrm{i}k\cdot X(z,\bar{z})}.$$



Campo gravitónico de gauge y campos cuánticos gravitacionales e interacciones locales y no locales de las partículas supermasivas en simetrías axiales.

Expansión de campo.

$$\Psi[X(\sigma), c(\sigma)] := \langle X(\sigma), c(\sigma) | \Psi \rangle$$

$$|\Psi\rangle = \sum_{\alpha} \int \frac{d^D k}{(2\pi)^D} \psi_{\alpha}(k) |\phi_{\alpha}(k)\rangle$$

$$\psi_{\alpha}(x) = \int \frac{d^D k}{(2\pi)^D} e^{ik \cdot x} \psi_{\alpha}(k)$$

$$|\Psi\rangle = \sum_r \psi_r |\phi_r\rangle$$

Campo escalar.

$$\phi(x) = \int \frac{d^D k}{(2\pi)^D} \phi(k) e^{ik \cdot x}$$

$$|\phi\rangle = \int \frac{d^D k}{(2\pi)^D} \phi(k) |k\rangle, \phi(k) = \langle k | \phi \rangle$$

$$\phi(x) = \langle x | \phi \rangle = \int \frac{d^D k}{(2\pi)^D} \langle x | k \rangle \langle k | \phi \rangle, \langle x | k \rangle = e^{ik \cdot x}$$

Warm-up.

$$(-\Delta + m^2)\phi(x) = 0.$$

$$k^2 = -m^2.$$

$$\phi(x) = \int dk \phi(k) e^{ikx}$$

$$\phi(x) = \langle x | \phi \rangle, \phi(k) = \langle k | \phi \rangle,$$

$$|\phi\rangle = \int dx \phi(x) |x\rangle = \int dk \phi(k) |k\rangle$$

$$K(x, x') := \langle x | K | x' \rangle = \delta(x - x')(-\Delta_x + m^2),$$

$$\int dx' K(x, x') \phi(x') = 0 \Leftrightarrow K|\phi\rangle = 0.$$

$$S = \frac{1}{2} \int dx \phi(x) (-\Delta + m^2) \phi(x) = \frac{1}{2} \int dx dx' \phi(x) K(x, x') \phi(x')$$

$$S = \frac{1}{2} \langle \phi | K | \phi \rangle$$



$$Q_B|\psi\rangle = 0$$

$$\Phi \in \mathcal{H}$$

$$Q_B|\Phi\rangle = 0$$

$$\langle A, B \rangle := \langle A \mid B \rangle,$$

$$S = \frac{1}{2} \langle \Phi, Q_B \Phi \rangle = \frac{1}{2} \langle \Phi | Q_B | \Phi \rangle.$$

$$\langle A, B \rangle = (-1)^{|A||B|} \langle B, A \rangle, \langle Q_B A, B \rangle = -(-1)^{|A|} \langle A, Q_B B \rangle,$$

$$N_{\mathrm{gh}}(\Phi) = 1$$

$$|\Phi| = 1$$

$$|\Phi\rangle^\ddagger = |\Phi\rangle^t$$

$$\begin{aligned} \langle \Phi, Q_B \Phi \rangle &= (-1)^{|\Phi|(|Q_B \Phi|)} \langle Q_B \Phi, \Phi \rangle = (-1)^{|\Phi|(1+|\Phi|)} \langle Q_B \Phi, \Phi \rangle \\ &= \langle Q_B \Phi, \Phi \rangle = -(-1)^{|\Phi|} \langle \Phi, Q_B \Phi \rangle \end{aligned}$$

$$|\Phi\rangle = |\Phi_{\downarrow}\rangle + c_0|\widetilde{\Phi}_{\downarrow}\rangle,$$

$$\Phi_{\downarrow}, \widetilde{\Phi}_{\downarrow} \in \mathcal{H}_0 \implies b_0|\Phi_{\downarrow}\rangle = b_0|\widetilde{\Phi}_{\downarrow}\rangle = 0.$$

$$N_{\mathrm{gh}}(\Phi_{\downarrow}) = 1, N_{\mathrm{gh}}(\widetilde{\Phi}_{\downarrow}) = 0.$$

$$Q_B = c_0 L_0 - b_0 M + \hat{Q}_B$$

$$S = \frac{1}{2} \langle \Phi_{\downarrow} | c_0 L_0 | \Phi_{\downarrow} \rangle + \frac{1}{2} \langle \widetilde{\Phi}_{\downarrow} | c_0 M | \widetilde{\Phi}_{\downarrow} \rangle + \langle \widetilde{\Phi}_{\downarrow} | c_0 \hat{Q}_B | \Phi_{\downarrow} \rangle.$$

$$0 = -M|\widetilde{\Phi}_{\downarrow}\rangle + \hat{Q}_B|\Phi_{\downarrow}\rangle, 0 = c_0 L_0|\Phi_{\downarrow}\rangle + c_0 \hat{Q}_B|\widetilde{\Phi}_{\downarrow}\rangle.$$

$$|\Phi\rangle = |\Phi_{\downarrow}\rangle + |\Phi_{\uparrow}\rangle, |\Phi_{\downarrow}\rangle = \Pi_s|\Phi\rangle, |\Phi_{\uparrow}\rangle = \bar{\Pi}_s|\Phi\rangle.$$

$$\Pi_s Q_B |\Phi\rangle = -b_0 M |\Phi_{\uparrow}\rangle + \hat{Q}_B |\Phi_{\downarrow}\rangle, \bar{\Pi}_s Q_B |\Phi\rangle = c_0 L_0 |\Phi_{\downarrow}\rangle + \hat{Q}_B |\Phi_{\uparrow}\rangle,$$

$$[\Pi_s, \hat{Q}_B] = [\Pi_s, M] = [\Pi_s, L_0] = 0$$

$$\begin{aligned} S &= \frac{1}{2} \langle \Phi, Q_B \Phi \rangle \\ &= \frac{1}{2} \langle \Pi_s \Phi + \bar{\Pi}_s \Phi, Q_B \Phi \rangle \\ &= \frac{1}{2} \langle \Pi_s \Phi, \bar{\Pi}_s Q_B \Phi \rangle + \frac{1}{2} \langle \bar{\Pi}_s \Phi, \Pi_s Q_B \Phi \rangle \\ &= \frac{1}{2} \langle \Phi_{\downarrow}, c_0 L_0 \Phi_{\downarrow} + \hat{Q}_B \Phi_{\uparrow} \rangle + \frac{1}{2} \langle \Phi_{\uparrow}, -b_0 M \Phi_{\uparrow} + \hat{Q}_B \Phi_{\downarrow} \rangle \\ &= \frac{1}{2} \langle \Phi_{\downarrow}, c_0 L_0 \Phi_{\downarrow} \rangle + \frac{1}{2} \langle \Phi_{\downarrow}, \hat{Q}_B \Phi_{\uparrow} \rangle - \frac{1}{2} \langle \Phi_{\uparrow}, b_0 M \Phi_{\uparrow} \rangle + \frac{1}{2} \langle \Phi_{\uparrow}, \hat{Q}_B \Phi_{\downarrow} \rangle. \end{aligned}$$



Invariancia de Gauge.

$$|\phi\rangle \sim |\psi\rangle + Q_B|\lambda\rangle.$$

$$|\Phi\rangle \longrightarrow |\Phi'\rangle = |\Phi\rangle + \delta_\Lambda|\Phi\rangle, \delta_\Lambda|\Phi\rangle = Q_B|\Lambda\rangle \quad N_{\text{gh}}(\Lambda) = 0$$

Gauge de Siegel y singularidad.

$$b_0|\Phi\rangle = 0$$

$$|\Lambda\rangle = -\Delta|\Phi\rangle, \Delta = \frac{b_0}{L_0}$$

$$0 = \{Q_B, b_0\}|\psi\rangle = L_0|\psi\rangle$$

$$b_0|\Phi'\rangle = b_0|\Phi\rangle + b_0Q_B|\Lambda\rangle = 0$$

$$b_0|\Phi\rangle = b_0\{Q_B, \Delta\}|\Phi\rangle = b_0Q_B\Delta|\Phi\rangle,$$

$$b_0Q_B(\Delta|\Phi\rangle + |\Lambda\rangle) = 0$$

$$A'_\mu = A_\mu + \partial_\mu\lambda$$

$$\partial^\mu A'_\mu = 0 \implies \Delta\lambda = -\partial^\mu A_\mu$$

$$\lambda = -\frac{k^\mu}{k^2}A_\mu$$

$$|\tilde{\Phi}_\downarrow\rangle = 0 \implies |\Phi\rangle = |\Phi_\downarrow\rangle.$$

$$S = \frac{1}{2}\langle\Phi|c_0L_0|\Phi\rangle,$$

$$L_0|\Phi\rangle = 0$$

$$\hat{Q}_B|\Phi\rangle = 0$$

$$\begin{aligned} S &= \frac{1}{2}\langle\Phi|Q_B\{c_0, b_0\}|\Phi\rangle = \frac{1}{2}\langle\Phi|Q_Bb_0c_0|\Phi\rangle \\ &= \frac{1}{2}\langle\Phi|\{b_0, Q_B\}c_0|\Phi\rangle - \frac{1}{2}\langle\Phi|b_0Q_Bc_0|\Phi\rangle \\ &= \frac{1}{2}\langle\Phi|c_0L_0|\Phi\rangle \end{aligned}$$

Campo expansivo, paridad y número fantasma.

$$n_r := N_{\text{gh}}(\phi_r), |\phi_r| = n_r \bmod 2.$$

$$\langle\phi_r^c | \phi_s\rangle = \delta_{rs}$$

$$|\phi_r\rangle = |\phi_{\downarrow,r}\rangle + |\phi_{\uparrow,r}\rangle, b_0|\phi_{\downarrow,r}\rangle = c_0|\phi_{\uparrow,r}\rangle = 0$$

$$|\psi_\uparrow\rangle = c_0|\tilde{\psi}_\downarrow\rangle, b_0|\tilde{\psi}_\downarrow\rangle = 0, N_{\text{gh}}(\psi_\uparrow) = N_{\text{gh}}(\tilde{\psi}_\downarrow) + 1.$$



$$|\Phi\rangle = \sum_r \psi_r |\phi_r\rangle$$

$$\forall r: |\Phi| = |\psi_r| |\phi_r|.$$

$$G(\psi_r) = 1 - n_r$$

$$\Phi = \sum_{n \in \mathbb{Z}} \Phi_n, N_{\text{gh}}(\Phi_n) = n$$

$$\Phi = \Phi_+ + \Phi_-, \Phi_+ = \sum_{n>1} \Phi_n, \Phi_- = \sum_{n \leq 1} \Phi_n$$

$$\forall n_r \neq n: \psi_r = 0$$

$$|\Phi_n\rangle = \delta(N_{\text{gh}} - n) |\Psi\rangle = \sum_r \delta(n_r - n) \psi_r |\phi_r\rangle$$

$$\begin{aligned} |\Phi\rangle &= |\Phi_{\downarrow}\rangle + |\Phi_{\uparrow}\rangle = |\Phi_{\downarrow}\rangle + c_0 |\tilde{\Phi}_{\downarrow}\rangle, \\ |\Phi_{\uparrow}\rangle &= c_0 |\tilde{\Phi}_{\downarrow}\rangle, |\tilde{\Phi}_{\downarrow}\rangle = b_0 |\Phi_{\uparrow}\rangle, \end{aligned}$$

$$b_0 |\Phi_{\downarrow}\rangle = 0, c_0 |\Phi_{\uparrow}\rangle = 0, b_0 |\tilde{\Phi}_{\downarrow}\rangle = 0$$

$$|\Phi_{\downarrow}\rangle = \sum_r \psi_{\downarrow,r} |\phi_{\downarrow,r}\rangle, |\Phi_{\uparrow}\rangle = \sum_r \psi_{\uparrow,r} |\phi_{\uparrow,r}\rangle.$$

$$Z = \int d\Phi_{\text{cl}} e^{-S[\Phi_{\text{cl}}]} = \int d\Phi_{\text{cl}} e^{-\frac{1}{2} \langle \Phi_{\text{cl}} | Q_B | \Phi_{\text{cl}} \rangle}$$

$$Z = \int \prod_s d\psi_s e^{-S[\{\psi_r\}]}$$

Transformaciones de gauge bajo la métrica de Faddeev-Popov.

$$F(\Phi_{\text{cl}}) := b_0 |\Phi_{\text{cl}}\rangle = 0.$$

$$\delta F = b_0 Q_B |\Lambda_{\text{cl}}\rangle,$$

$$\det \frac{\delta F}{\delta \Lambda_{\text{cl}}} = \det b_0 Q_B.$$

$$\det b_0 Q_B = \int dB' dC e^{-S_{\text{FP}}}, S_{\text{FP}} = -\langle B' | b_0 Q_B | C \rangle.$$

$$N_{\text{gh}}(B') = 3, N_{\text{gh}}(C) = 0.$$

$$|B'\rangle = \delta(N_{\text{gh}} - 3) \sum_r b'_r |\phi_r\rangle, |C\rangle = \delta(N_{\text{gh}}) \sum_r c_r |\phi_r\rangle,$$

$$|b_r| = |c_r| = 1.$$



$$|B'\rangle = 0, |C\rangle = 1.$$

$$\begin{aligned}\delta|C\rangle &= Q_B|\Lambda_{-1}\rangle, & N_{\text{gh}}(\Lambda_{-1}) &= -1 \\ \delta|B'\rangle &= b_0|\Lambda'\rangle, & N_{\text{gh}}(\Lambda') &= 4\end{aligned}$$

$$|\Lambda\rangle=|\Lambda_0\rangle+Q_B|\Lambda_{-1}\rangle,$$

$$|\Phi'_{\text{cl}}\rangle\longrightarrow|\Phi_{\text{cl}}\rangle+Q_B|\Lambda_0\rangle$$

$$\begin{aligned}|B'\rangle &= |B_{\downarrow}'\rangle + c_0|B\rangle, |B\rangle\colon=|\tilde{B}_{\downarrow}'\rangle, \\ |\Lambda'\rangle &= |\Lambda_{\downarrow}'\rangle + c_0|\tilde{\Lambda}_{\downarrow}'\rangle.\end{aligned}$$

$$\delta|B_{\downarrow}'\rangle=|\tilde{\Lambda}_{\downarrow}'\rangle,\delta|B\rangle=0$$

$$|F'\rangle=c_0|B'\rangle=0\implies|B_{\downarrow}'\rangle=0.$$

$$\det\frac{\delta F'}{\delta\Lambda'}=\det c_0b_0=\det c_0\det b_0=\frac{1}{2}\det\{b_0,c_0\}=\frac{1}{2}$$

$$\Delta_{\text{FP}}=\int\text{d}B\,\text{d}C\,\text{e}^{-S_{\text{FP}}[B,C]}, S_{\text{FP}}=\langle B|Q_B|C\rangle$$

$$b_0|B\rangle=0, N_{\text{gh}}(B)=2, |B|=1$$

$$S_{\text{FP}}=\frac{1}{2}(\langle B|Q_B|C\rangle+\langle C|Q_B|B\rangle).$$

$$|B\rangle=|B_{\downarrow}\rangle+c_0|\tilde{B}_{\downarrow}\rangle$$

$$\delta|B\rangle=Q_B|\Lambda_1\rangle$$

$$|\Phi\rangle=\sum_n|\Phi_n\rangle$$

$$b_0|\Phi\rangle=0\implies b_0|\Phi_n\rangle=0.$$

$$|\beta\rangle=\sum_{n\in\mathbb{Z}}|\beta_n\rangle.$$

$$Z=\int\mathrm{d}\Phi\,\mathrm{d}\beta\mathrm{e}^{-S[\Phi,\beta]}$$

$$\begin{aligned}S[\Phi,\beta]&=\frac{1}{2}\langle\Phi|Q_B|\Phi\rangle+\langle\beta|b_0|\Phi\rangle\\&=\sum_{n\in\mathbb{Z}}\left(\frac{1}{2}\langle\Phi_{2-n}|Q_B|\Phi_n\rangle+\langle\beta_{4-n}|b_0|\Phi_n\rangle\right)\end{aligned}$$

$$\delta|\Phi\rangle=Q_B|\Lambda\rangle$$

$$|\Lambda\rangle=\sum_{n\in\mathbb{Z}}|\Lambda_n\rangle.$$

Acción de espacio – tiempo curvo.

$$|\Phi\rangle = \frac{1}{\sqrt{\alpha'}} \int \frac{d^D k}{(2\pi)^D} \left(T(k) + A_\mu(k) \alpha_{-1}^\mu + i \sqrt{\frac{\alpha'}{2}} B(k) b_{-1} c_0 + \dots \right) |k, \downarrow\rangle$$

$$(\alpha' k^2 - 1)T(k) = 0, k^2 A_\mu(k) + i k_\mu B(k) = 0, \\ k^\mu A_\mu(k) + i B(k) = 0.$$

$$k^2 A_\mu(k) - k_\mu k \cdot A(k) = 0$$

$$(\alpha' \Delta + 1)T = 0, B = \partial^\mu A_\mu, \Delta A_\mu = \partial_\mu B$$

$$Q_B = c_0 L_0 - b_0 M + \hat{Q}_B \\ M \sim 2c_{-1} c_1, \hat{Q}_B \sim c_1 L_{-1}^m + c_{-1} L_1^m, \\ L_1^m \sim \alpha_0 \cdot \alpha_1, L_{-1}^m \sim \alpha_0 \cdot \alpha_{-1}.$$

$$Q_B |\Phi\rangle = \frac{1}{\sqrt{\alpha'}} \int \frac{d^D k}{(2\pi)^D} \left(T(k) c_0 L_0 |k, \downarrow\rangle + A_\mu(k) (c_0 L_0 + \eta_{\nu\rho} c_{-1} \alpha_1^\nu \alpha_0^\rho) \alpha_{-1}^\mu |k, \downarrow\rangle \right. \\ \left. + i \sqrt{\frac{\alpha'}{2}} B(k) (-2b_0 c_{-1} c_1 + \eta_{\nu\rho} c_1 \alpha_{-1}^\nu \alpha_0^\rho) b_{-1} c_0 |k, \downarrow\rangle \right)$$

$$= \frac{1}{\sqrt{\alpha'}} \int \frac{d^D k}{(2\pi)^D} \left(T(k) (\alpha' k^2 - 1) c_0 |k, \downarrow\rangle \right. \\ \left. + A_\mu(k) (\alpha' k^2 c_0 \alpha_{-1}^\mu + \sqrt{2\alpha'} \eta_{\nu\rho} \eta^{\mu\nu} k^\rho c_{-1}) |k, \downarrow\rangle \right. \\ \left. + i \sqrt{\frac{\alpha'}{2}} B(k) (2c_{-1} + \sqrt{2\alpha'} \eta_{\nu\rho} k^\rho \alpha_{-1}^\nu c_0) |k, \downarrow\rangle \right)$$

$$= \frac{1}{\sqrt{\alpha'}} \int \frac{d^D k}{(2\pi)^D} \left(T(k) (\alpha' k^2 - 1) c_0 |k, \downarrow\rangle \right. \\ \left. + \alpha' (A_\mu(k) k^2 + i k^\mu B(k) c_0 \alpha_{-1}^\mu) |k, \downarrow\rangle \right. \\ \left. + \sqrt{2\alpha'} (k^\mu A_\mu(k) + i B(k) c_{-1}) |k, \downarrow\rangle \right)$$

$$|\Lambda\rangle = \frac{i}{\sqrt{2\alpha'}} \int \frac{d^D k}{(2\pi)^D} (\lambda(k) b_{-1} |k, \downarrow\rangle + \dots).$$

$$Q_B |\Lambda\rangle = \frac{i}{\sqrt{\alpha'}} \int \frac{d^D k}{(2\pi)^D} \lambda(k) \left(-\sqrt{\frac{\alpha'}{2}} k^2 b_{-1} c_0 + k_\mu \alpha_{-1}^\mu \right) |k, \downarrow\rangle,$$

$$\delta A_\mu = -i k_\mu \lambda, \delta B = k^2 \lambda.$$

$$|T\rangle = \int \frac{d^D k}{(2\pi)^D} T(k) c_1 |k, 0\rangle,$$



$$\langle T| = \int \frac{d^D k}{(2\pi)^D} T(k) \langle -k, 0 | c_{-1}$$

$$\langle T^\dagger| = \int \frac{d^D k}{(2\pi)^D} T(k)^* \langle k, 0 | c_{-1}$$

$$T(k)^* = T(-k)$$

$$S[T] = \frac{1}{2} \int \frac{d^D k}{(2\pi)^D} T(-k) \left(k^2 - \frac{1}{\alpha'} \right) T(k)$$

$$S[A] = \frac{1}{2} \int \frac{d^D k}{(2\pi)^D} A_\mu(-k) k^2 A^\mu(k)$$

$$\frac{L_0^+}{2} = \frac{1}{2} (k^2 + m^2)$$

$$\begin{aligned} \langle T | c_0 L_0 | T \rangle &= \frac{1}{\alpha'} \int \frac{d^D k}{(2\pi)^D} \frac{d^D k'}{(2\pi)^D} T(k) T(k') \langle -k', 0 | c_{-1} c_0 L_0 c_1 | k, 0 \rangle \\ &= \frac{1}{\alpha'} \int \frac{d^D k}{(2\pi)^D} \frac{d^D k'}{(2\pi)^D} T(k) T(k') (\alpha' k^2 - 1) \langle -k', 0 | c_{-1} c_0 c_1 | k, 0 \rangle \\ &= \frac{1}{\alpha'} \int \frac{d^D k}{(2\pi)^D} d^D k' T(k) T(k') (\alpha' k^2 - 1) \delta^{(D)}(k + k') \end{aligned}$$

$$Q_B|\Psi\rangle=0$$

$$S=\frac{1}{2}\langle\Psi,Q_B\Psi\rangle.$$

$$N_{\mathrm{gh}}(\Psi)=2$$

$$N_{\mathrm{gh}}(\langle\cdot,\cdot\rangle)=1.$$

$$\langle A,B\rangle=\langle A|c_0^-|B\rangle.$$

$$S=\frac{1}{2}\langle\Psi|c_0^-Q_B|\Psi\rangle.$$

$$\mathcal{H}=\mathcal{H}^-\oplus c_0^-\mathcal{H}^-, \mathcal{H}^-:=\mathcal{H}\cap \ker b_0^-,$$

$$|\Psi\rangle=|\Psi_{-}\rangle+c_0^{-}|\tilde{\Psi}_{-}\rangle, \Psi_{-},\tilde{\Psi}_{-}\in\mathcal{H}^{-}$$

$$c_0^-|\Psi\rangle=c_0^-|\Psi_{-}\rangle.$$

$$b_0^-|\Psi\rangle=0$$

$$L_0^-|\Psi\rangle=0,$$

$$\Psi\in\mathcal{H}^-\cap\ker L_0^-.$$

$$|\Psi\rangle\rightarrow|\Psi'\rangle=|\Psi\rangle+\delta_\Lambda|\Psi\rangle,\delta_\Lambda|\Psi\rangle=Q_B|\Lambda\rangle,$$

$$N_{\mathrm{gh}}(\Lambda)=1, L_0^-|\Lambda\rangle=0, b_0^-|\Lambda\rangle=0$$



$$b_0^+|\Psi\rangle=0.$$

$$S=\frac{1}{2}\langle\Psi|c_0^-c_0^+L_0^+|\Psi\rangle=\frac{1}{4}\langle\Psi|c_0\bar{c}_0L_0^+|\Psi\rangle.$$

$$L_0^+|\Psi\rangle=0.$$

$$c_0^-Q_B=(c_0-\bar{c}_0)(c_0L_0+\bar{c}_0\bar{L}_0)=c_0\bar{c}_0(L_0+\bar{L}_0).$$

Off-shell.

$$A_{0,3}=\left\langle \prod_{i=1}^3\mathcal{V}_i(z_i)\right\rangle_{S^2}\propto (z_1-z_2)^{h_3-h_1-h_2}\times \text{perms}\times \text{c.c.}$$

$$z\longrightarrow f_g(z)=\frac{az+b}{cz+d}\in\mathrm{SL}(2,\mathbb{C})$$

$$z=f_i(w_i), z_i=f_i(0)$$

$$f\circ \mathcal{V}(w)=f'(w)^h\overline{f'(w)}^{\bar{h}}\mathcal{V}(f(w))$$

$$\begin{aligned} A_{0,3} &= \left\langle \prod_{i=1}^3 f_i \circ \mathcal{V}_i(0) \right\rangle_{S^2} = \left(\prod_{i=1}^3 f_i'(0)^{h_i} \overline{f_i'(0)}^{\bar{h}_i} \right) \left\langle \prod_{i=1}^3 \mathcal{V}_i(f_i(0)) \right\rangle_{S^2} \\ &\propto \left(\prod_{i=1}^3 f_i'(0)^{h_i} \overline{f_i'(0)}^{\bar{h}_i} \right) (f_1(0) - f_2(0))^{h_3 - h_1 - h_2} \times \text{perms} \times \text{c.c.} \end{aligned}$$

$$f_i\rightarrow \frac{af_i+b}{cf_i+d}$$

$$f_i'\rightarrow \frac{f_i'}{(cf_i+d)^2}, f_i-f_j\rightarrow \frac{f_i-f_j}{(cf_i+d)(cf_j+d)}$$

$$\mathcal{V}_{0,3}(\mathcal{V}_1,\mathcal{V}_2,\mathcal{V}_3):= \begin{array}{c} \mathcal{V}_1 \\ \diagdown \\ \\ \diagup \\ \mathcal{V}_2 \end{array} \begin{array}{c} \\ \diagup \\ \\ \diagdown \end{array} \mathcal{V}_3 \quad = A_{0,3}(\mathcal{V}_1,\mathcal{V}_2,\mathcal{V}_3).$$

$$A_{0,4} = \int d^2 z_4 \left\langle \prod_{i=1}^3 c \bar{c} V_i(z_i) V_4(z_4) \right\rangle_{S^2}$$

$$z_4 \longrightarrow z_1, z_2, z_3$$

$$A_{0,4} \propto \prod_{\substack{i,j=1 \\ i < j}}^3 |z_i - z_j|^{2+k_i \cdot k_j} \int d^2 z_4 \prod_{i=1}^3 |z_4 - z_i|^{k_i \cdot k_4}$$

$$\mathcal{F}_{0,4}^{(s)} = \int \frac{d^2 q}{|q|^2} \langle c \bar{c} V_1(z_1) c \bar{c} V_2(z_2) c \bar{c} V_3(0) | q y_4 \Big|^2 V_4(q y_4) \rangle$$

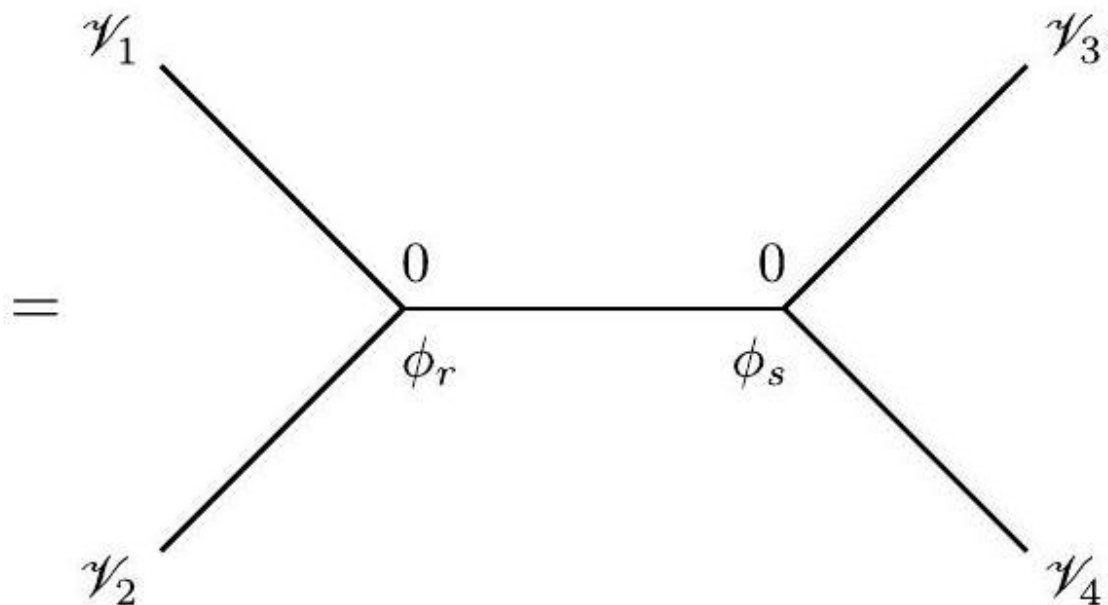
$$\mathcal{F}_{0,4}^{(s)} = - \int \frac{d^2 q}{|q|^2} \langle c \bar{c} V_1(z_1) c \bar{c} V_2(z_2) \oint_{|w|=|q|^{1/2}} dw w(w) \oint_{|w|=|q|^{1/2}} d\bar{w} \overline{b(w)} c \bar{c} V_4(q y_4) c \bar{c} V_3(0) \rangle$$

$$\mathcal{F}_{0,4}^{(s)} = - \int \frac{d^2 q}{|q|^2} \langle c \bar{c} V_1(z_1) c \bar{c} V_2(z_2) \oint dw w b(w) \oint d\bar{w} \bar{w} \overline{b(w)} q^{L_0} \bar{q}^{\bar{L}_0} c \bar{c} V_4(y_4) c \bar{c} V_3(0) \rangle$$

$$\mathcal{F}_{0,4}^{(s)} = \langle c \bar{c} V_1(z_1) c \bar{c} V_2(z_2) \phi_r(0) \rangle \langle c \bar{c} V_3(z_3) c \bar{c} V_4(y_4) \phi_s(0) \rangle \int \frac{d^2 q}{|q|^2} \langle \phi_r^c q^{L_0} \bar{q}^{\bar{L}_0} b_0 \bar{b}_0 \phi_s^c \rangle$$

$$\Delta(\phi_r^c, \phi_s^c) := \langle \phi_r^c | \Delta | \phi_s^c \rangle := - \int \frac{d^2 q}{|q|^2} \langle \phi_r^c q^{L_0} \bar{q}^{\bar{L}_0} b_0 \bar{b}_0 \phi_s^c \rangle$$

$$\mathcal{F}_{0,4}^{(s)} = \mathcal{V}_{0,3}(c \bar{c} V_1(z_1), c \bar{c} V_2(z_2), \phi_r(0)) \times \Delta(\phi_r^c, \phi_s^c) \times \mathcal{V}_{0,3}(c \bar{c} V_3(z_3), c \bar{c} V_4(y_4), \phi_s(0))$$



$$q = e^{-s+i\theta}, s \in \mathbb{R}_+, \theta \in [0, 2\pi)$$

$$\int \frac{d^2 q}{|q|^2} q^{L_0} \bar{q}^{\bar{L}_0} = 2 \int_0^\infty ds \int_0^{2\pi} d\theta e^{-s(L_0 + \bar{L}_0)} e^{i\theta(L_0 - \bar{L}_0)} = \frac{2}{L_0 + \bar{L}_0} \delta_{L_0, \bar{L}_0}$$

$$\Delta = -\frac{2b_0 \bar{b}_0}{L_0 + \bar{L}_0} \delta_{L_0, \bar{L}_0} = \frac{b_0^+}{L_0^+} b_0^- \delta_{L_0^-, 0}$$

$$L_0 |\phi_\alpha(k)\rangle = \bar{L}_0 |\phi_\alpha(k)\rangle = \frac{\alpha'}{4} (k^2 + m_\alpha^2) |\phi_\alpha(k)\rangle$$

$$\Delta_{\alpha\beta}(k) = \int \frac{d^2 q}{|q|^2} \langle \phi_\alpha^c(k) q^{L_0} \bar{q}^{\bar{L}_0} b_0 \bar{b}_0 \phi_\beta^c(-k) \rangle = \frac{M_{\alpha\beta}(k)}{k^2 + m_\alpha^2}$$

$$M_{\alpha\beta}(k) := \frac{2}{\alpha'} \langle \phi_\alpha^c(k) | b_0^+ b_0^- | \phi_\beta^c(-k) \rangle.$$

$$\mathcal{F}_{0,4}^{(t)} = \begin{array}{c} \begin{array}{cc} \mathcal{V}_1 & \mathcal{V}_3 \\ & \diagdown \quad \diagup \\ & 0 \\ & \diagup \quad \diagdown \\ & 0 \\ \mathcal{V}_2 & \mathcal{V}_4 \end{array} \end{array} \qquad \mathcal{F}_{0,4}^{(u)} = \begin{array}{c} \begin{array}{cc} \mathcal{V}_1 & \mathcal{V}_3 \\ & \diagdown \quad \diagup \\ & \text{---} 0 \text{---} \\ & \diagup \quad \diagdown \\ & \text{---} 0 \text{---} \\ \mathcal{V}_2 & \mathcal{V}_4 \end{array} \end{array}$$

$$\mathcal{V}_{0,4} = \begin{array}{c} \begin{array}{cc} \mathcal{V}_1 & \mathcal{V}_3 \\ & \diagdown \quad \diagup \\ & 0 \\ & \diagup \quad \diagdown \\ \mathcal{V}_2 & \mathcal{V}_4 \end{array} \end{array}$$

$$A_{0,4} = \mathcal{F}_{0,4}^{(s)} + \mathcal{F}_{0,4}^{(t)} + \mathcal{F}_{0,4}^{(u)} + \mathcal{V}_{0,4}.$$

Estados Off-shell.

$$\mathcal{H} = \mathcal{H}_m \otimes \mathcal{H}_{\text{gh}}.$$

$$\mathcal{H} = \text{Span}\{|\phi_r\rangle\}.$$

$$n_r := N_{\text{gh}}(\phi_r) \in \mathbb{Z}$$

$$|\phi_r| := N_{\text{gh}}(\phi_r) \bmod 2 = \begin{cases} 0 & N_{\text{gh}}(\phi_r) \text{ even} \\ 1 & N_{\text{gh}}(\phi_r) \text{ odd} \end{cases}$$

$$\langle \phi_r^c | \phi_s \rangle = \delta_{rs}$$

$$n_r^c := N_{\text{gh}}(\phi_r^c)$$

$$n_r^c + n_r = 6$$

$$\langle \phi_r | \phi_s \rangle = 0$$

$$\langle \phi_r | \phi_s^c \rangle = (-1)^{|\phi_r|} \delta_{rs}.$$

$$1 = \sum_r |\phi_r\rangle \langle \phi_r^c| = \sum_r (-1)^{|\phi_r|} |\phi_r^c\rangle \langle \phi_r|.$$

$$\mathcal{H} = \mathcal{H}_{\pm} \oplus c_0^{\pm} \mathcal{H}_{\pm}$$

$$\mathcal{H}_{\pm} := \mathcal{H} \cap \ker b_0^{\pm} = \mathcal{H}_0 \oplus c_0^{\mp} \mathcal{H}_0, \mathcal{H}_0 := \mathcal{H} \cap \ker b_0^- \cap \ker b_0^+.$$

$$L_0^- |\phi\rangle = 0, b_0^- |\phi\rangle = 0$$

$$\mathcal{H} \sim \mathcal{H}_{\downarrow\downarrow} \oplus \mathcal{H}_{\downarrow\uparrow} \oplus \mathcal{H}_{\uparrow\downarrow} \oplus \mathcal{H}_{\uparrow\uparrow},$$

$$\mathcal{H}_{\downarrow\uparrow} \sim \mathcal{H}_0, \mathcal{H}_{\uparrow\downarrow} \sim \bar{c}_0 \mathcal{H}_{\downarrow\downarrow}, \mathcal{H}_{\uparrow\uparrow} \sim c_0 \mathcal{H}_{\downarrow\downarrow}, \mathcal{H}_{\downarrow\downarrow} \sim c_0 \bar{c}_0 \mathcal{H}_{\downarrow\downarrow}$$

$$\phi_r = \phi_{\downarrow\downarrow,r} + \phi_{\downarrow\uparrow,r} + \phi_{\uparrow\downarrow,r} + \phi_{\uparrow\uparrow,r}$$

$$\begin{aligned} b_0 |\phi_{\downarrow\downarrow,r}\rangle &= \bar{b}_0 |\phi_{\downarrow\downarrow,r}\rangle = 0, & b_0 |\phi_{\downarrow\uparrow,r}\rangle &= \bar{c}_0 |\phi_{\downarrow\uparrow,r}\rangle = 0, \\ c_0 |\phi_{\uparrow\downarrow,r}\rangle &= \bar{b}_0 |\phi_{\uparrow\downarrow,r}\rangle = 0, & c_0 |\phi_{\uparrow\uparrow,r}\rangle &= \bar{c}_0 |\phi_{\uparrow\uparrow,r}\rangle = 0. \end{aligned}$$

$$|\phi_{\downarrow\uparrow,r}\rangle = \bar{c}_0 |\phi_{\downarrow\downarrow,r}\rangle |\phi_{\uparrow\downarrow,r}\rangle = c_0 |\phi_{\downarrow\downarrow,r}\rangle |\phi_{\uparrow\uparrow,r}\rangle = c_0 \bar{c}_0 |\phi_{\downarrow\downarrow,r}\rangle.$$

$$\phi_r^c = \phi_{\downarrow\downarrow,r}^c + \phi_{\downarrow\uparrow,r}^c + \phi_{\uparrow\downarrow,r}^c + \phi_{\uparrow\uparrow,r}^c$$

$$\begin{aligned} \langle \phi_{\downarrow\downarrow,r}^c | c_0 &= \langle \phi_{\downarrow\downarrow,r}^c | \bar{c}_0 = 0, & \langle \phi_{\downarrow\uparrow,r}^c | c_0 &= \langle \phi_{\downarrow\uparrow,r}^c | \bar{b}_0 = 0, \\ \langle \phi_{\uparrow\downarrow,r}^c | b_0 &= \langle \phi_{\uparrow\downarrow,r}^c | \bar{c}_0 = 0, & \langle \phi_{\uparrow\uparrow,r}^c | b_0 &= \langle \phi_{\uparrow\uparrow,r}^c | \bar{b}_0 = 0. \end{aligned}$$

$$\langle \phi_{\downarrow\uparrow,r}^c | = \langle \phi_{\downarrow\downarrow,r}^c | \bar{b}_0 \langle \phi_{\uparrow\downarrow,r}^c | = \langle \phi_{\downarrow\downarrow,r}^c | b_0 \langle \phi_{\uparrow\uparrow,r}^c | = \langle \phi_{\downarrow\downarrow,r}^c | \bar{b}_0 b_0.$$

$$\langle \phi_{x,r}^c | \phi_{y,s} \rangle = \delta_{xy} \delta_{rs}$$

$$b_0^- |\phi_r\rangle = 0 \Rightarrow b_0 |\phi_r\rangle = \bar{b}_0 |\phi_r\rangle$$



$$\phi_{\uparrow\downarrow,r} + \phi_{\uparrow\uparrow,r} = \phi_{\downarrow\downarrow,r} + \phi_{\downarrow\uparrow,r}$$

$$\phi_r = 2(\phi_{\downarrow\downarrow,r} + \phi_{\downarrow\uparrow,r})$$

Amplitudes Off-shell.

$$\chi_{g,n} := \chi(\Sigma_{g,n}) = 2 - 2g - n.$$

$$M_{g,n} := \dim_{\mathbb{R}} \mathcal{M}_{g,n} = 6g - 6 + 2n, \text{ for } \begin{cases} g \geq 2, \\ g = 1, n \geq 1, \\ g = 0, n \geq 3. \end{cases}$$

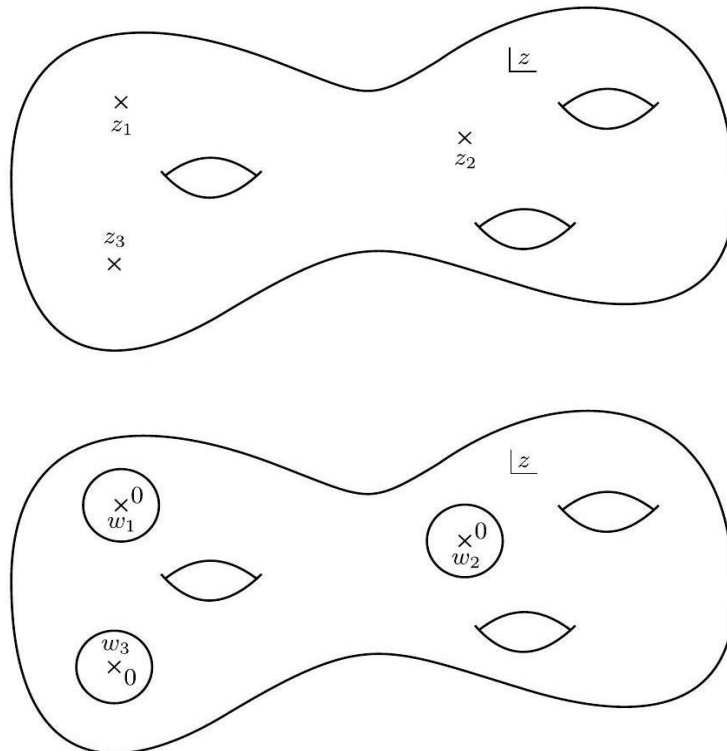
$$A_{g,n}(\mathcal{V}_1, \dots, \mathcal{V}_n) = \int_{\mathcal{M}_{g,n}} \omega_{\mathbb{M}_{g,n}}^{g,n}(\mathcal{V}_1, \dots, \mathcal{V}_n)$$

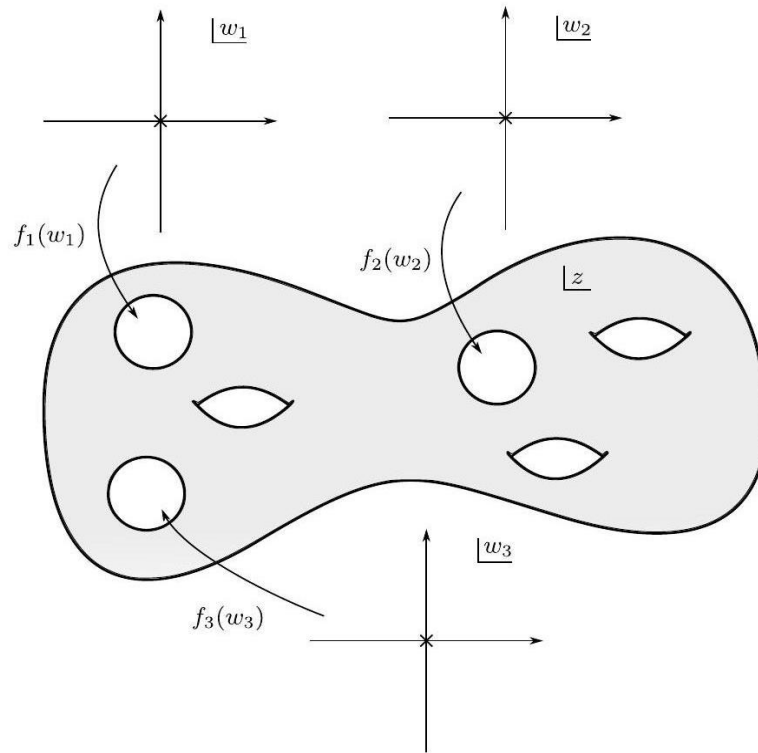
$$\omega_{\mathbb{M}_{g,n}}^{g,n}(\mathcal{V}_1, \dots, \mathcal{V}_n) = \left\langle \text{ghosts} \times \prod_{i=1}^n \mathcal{V}_i \right\rangle_{\Sigma_{g,n}} \times \bigwedge_{\lambda=1}^{M_{g,n}} dt_{\lambda}$$

Coordenadas locales.

$$z = f_i(w_i), z_i = f_i(0)$$

$$A_{g,n}(\mathcal{V}_1, \dots, \mathcal{V}_n)_{\mathcal{S}_{g,n}} = \int_{\mathcal{S}_{g,n}} \omega_{\mathbb{M}_{g,n}}^{g,n}(\mathcal{V}_1, \dots, \mathcal{V}_n) \Big|_{\mathcal{S}_{g,n}}$$





Figuras 4, 5 y 6. Curvatura local provocada por una partícula supermasiva.

$$\forall \mathcal{S}_{g,n}: A_{g,n}(\mathcal{V}_1, \dots, \mathcal{V}_n)_{\mathcal{S}_{g,n}} = A_{g,n}(\mathcal{V}_1, \dots, \mathcal{V}_n) \text{ (on-shell)}.$$

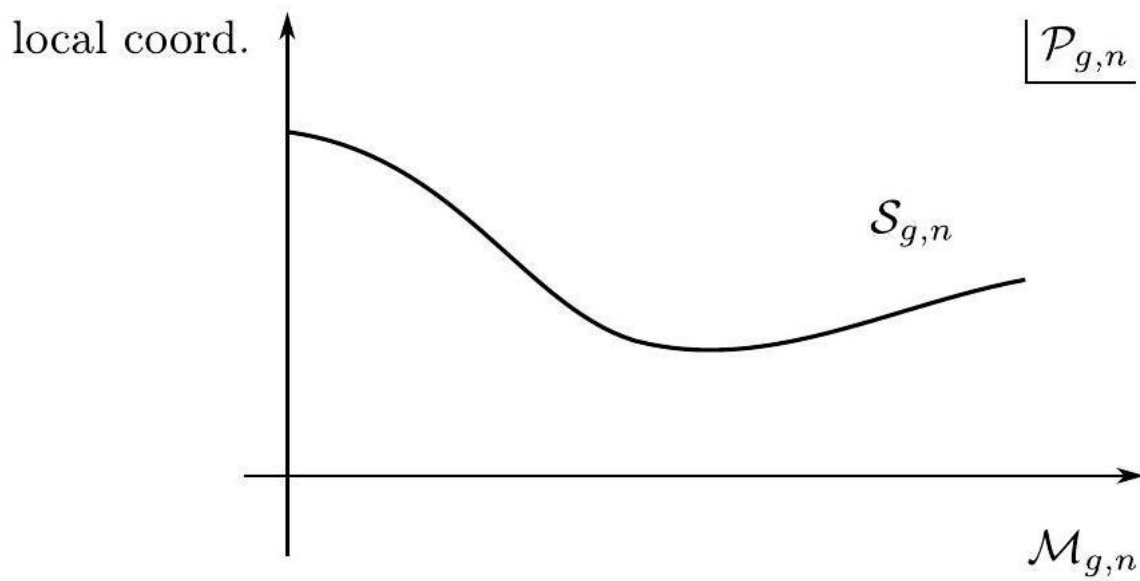


Figura 7. Fluctuaciones en el sistema de coordenadas de una partícula supermasiva al deformar o curvar el espacio – tiempo cuántico.

Geometría en espacios de Moduli y superficies de Riemann.

$$\text{disks} = n.$$

$$\backslash \# \text{spheres} = 2g - 2 + n.$$

$$\text{circles} = \frac{n + 3(2g - 2 + n)}{2} = 3g - 3 + 2n$$

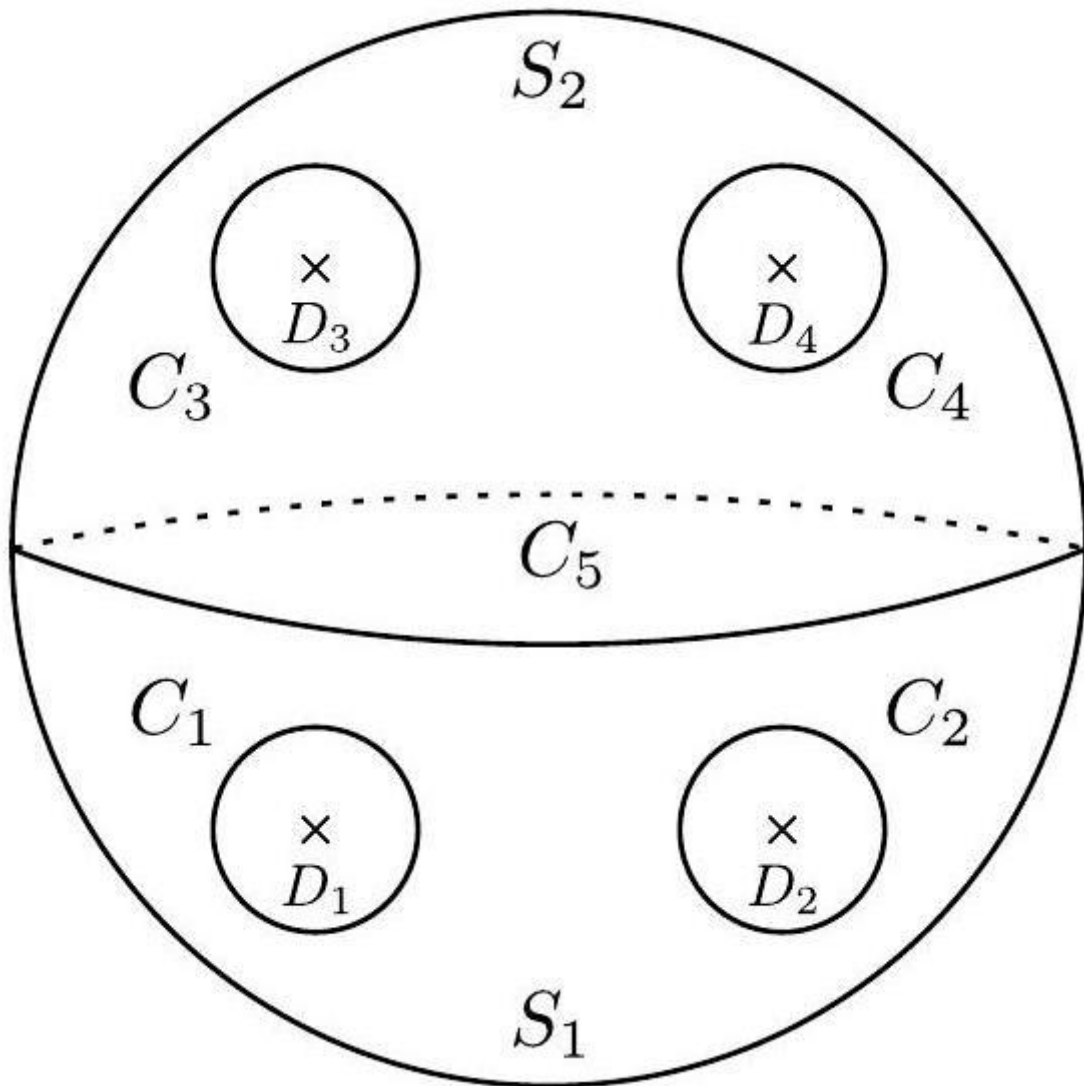
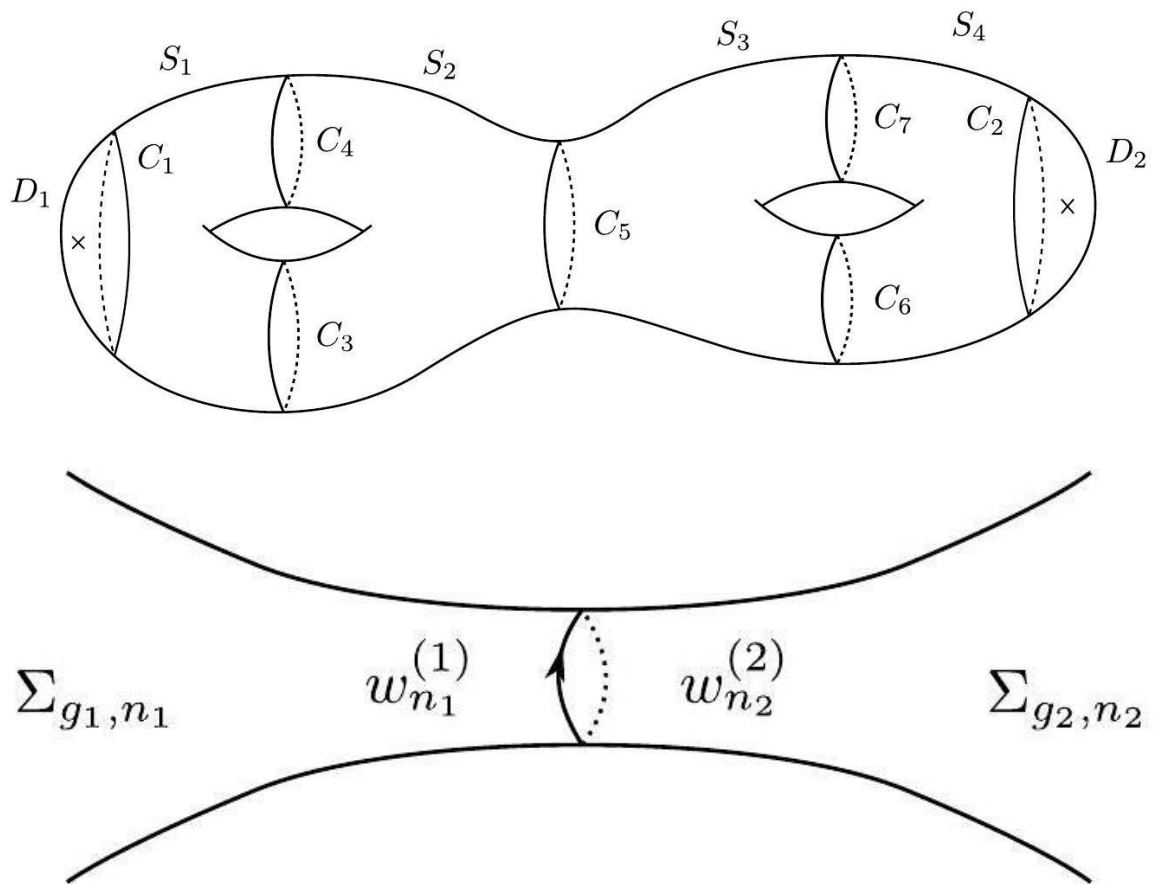


Figura 8. Propagadores de una partícula reperiutada en relación al propagador de la partícula supermasiva.



Figuras 9 y 10. Entrelazamiento cuántico por curvatura local en dimensión $\mathbb{R}^4$, es decir, en la misma dimensión.

$$C_{\Lambda(ab)} := S_a \cap S_b, C_{i(a)} := S_a \cap D_i, \{C_\alpha\} = \{C_\Lambda, C_i\}$$

$$\Lambda = 1, \dots, M_{g,n}^c, M_{g,n}^c = 3g - 3 + n$$

$$\text{on } C_{\Lambda(ab)}: z_a = F_{ab}(z_b)$$

$$\text{on } C_{i(a)}: z_a = f_{ai}(w_i)$$

$$w_i \rightarrow \tilde{w}_i = e^{i\alpha_i} w_i$$

$$z = f_i(w_i), z = \tilde{f}_i(\tilde{w}_i)$$

$$f_i(w_i) = \tilde{f}_i(e^{i\alpha_i} w_i)$$

$$\hat{\mathcal{P}}_{g,n} = \mathcal{P}_{g,n}/U(1)^n$$

$$\text{on } C_{\Lambda(ab)}: z_a - z_{a,m} = \frac{q_\Lambda}{z_b - z_{b,n}}$$

$$\text{on } C_{i(a)}: z_a - z_{a,m} = w_i + \sum_{N=1}^{\infty} p_{i,N} w_i^N$$

$$\{x_s\} = \{q_\Lambda, p_{i,N}\}$$

$$z_a = F_{ab}(z_b; x_s), z_a = f_{ai}(w_i; x_s)$$

$$\text{on } C_\alpha\colon \sigma_\alpha = F_\alpha(\tau_\alpha; x_s)$$

Espacio tangencial.

$$\delta x_s = \epsilon V_s$$

$$\epsilon V_s \partial_s f = f(x_s + \epsilon V_s) - f(x_s)$$

$$F_\alpha \longrightarrow F_\alpha + \epsilon \delta F_\alpha$$

$$\sigma_\alpha = F_\alpha(\tau_\alpha)$$

$$\sigma'_\alpha = F_\alpha(\tau_\alpha) + \epsilon \delta F_\alpha(\tau_\alpha) = \sigma_\alpha + \epsilon \delta F_\alpha(\tau_\alpha) = \sigma_\alpha + \epsilon \delta F_\alpha(F_\alpha^{-1}(\sigma_\alpha))$$

$$\sigma'_\alpha = \sigma_\alpha + \epsilon v^{(\alpha)}(\sigma_\alpha), v^{(\alpha)} = \delta F_\alpha \circ F_\alpha^{-1}$$

$$V^{(\alpha)} \sim (v^{(\alpha)}, C_{(\alpha)}).$$

$$V \sim (v, C), C \subseteq \bigcup_\alpha C_\alpha$$

$$v|_{C_\alpha} = v^{(\alpha)}$$

$$x_s \longrightarrow x_s + \epsilon \delta x_s$$

$$C_\alpha\colon F_\alpha \longrightarrow F_\alpha + \epsilon \delta F_\alpha, \delta F_\alpha = \frac{\partial F_\alpha}{\partial x_s} \delta x_s$$

$$\sigma'_\alpha = \sigma_\alpha + \epsilon v_s^{(\alpha)}(\sigma_\alpha) \delta x_s, v_s^{(\alpha)}(\sigma_\alpha) = \frac{\partial F_\alpha}{\partial x_s}(F_\alpha^{-1}(\sigma_\alpha))$$

Plumbing fixture.

$$D_q^{(1)} = \left\{ \left| w_{n_1}^{(1)} \right| \leq |q|^{1/2} \right\}, D_q^{(2)} = \left\{ \left| w_{n_2}^{(2)} \right| \leq |q|^{1/2} \right\}$$

$$\Sigma_{g,n} = \Sigma_{g_1,n_1}$$

$$\Sigma_{g_2,n_2}, \begin{cases} g = g_1 + g_2 \\ n = n_1 + n_2 - 2 \end{cases}$$

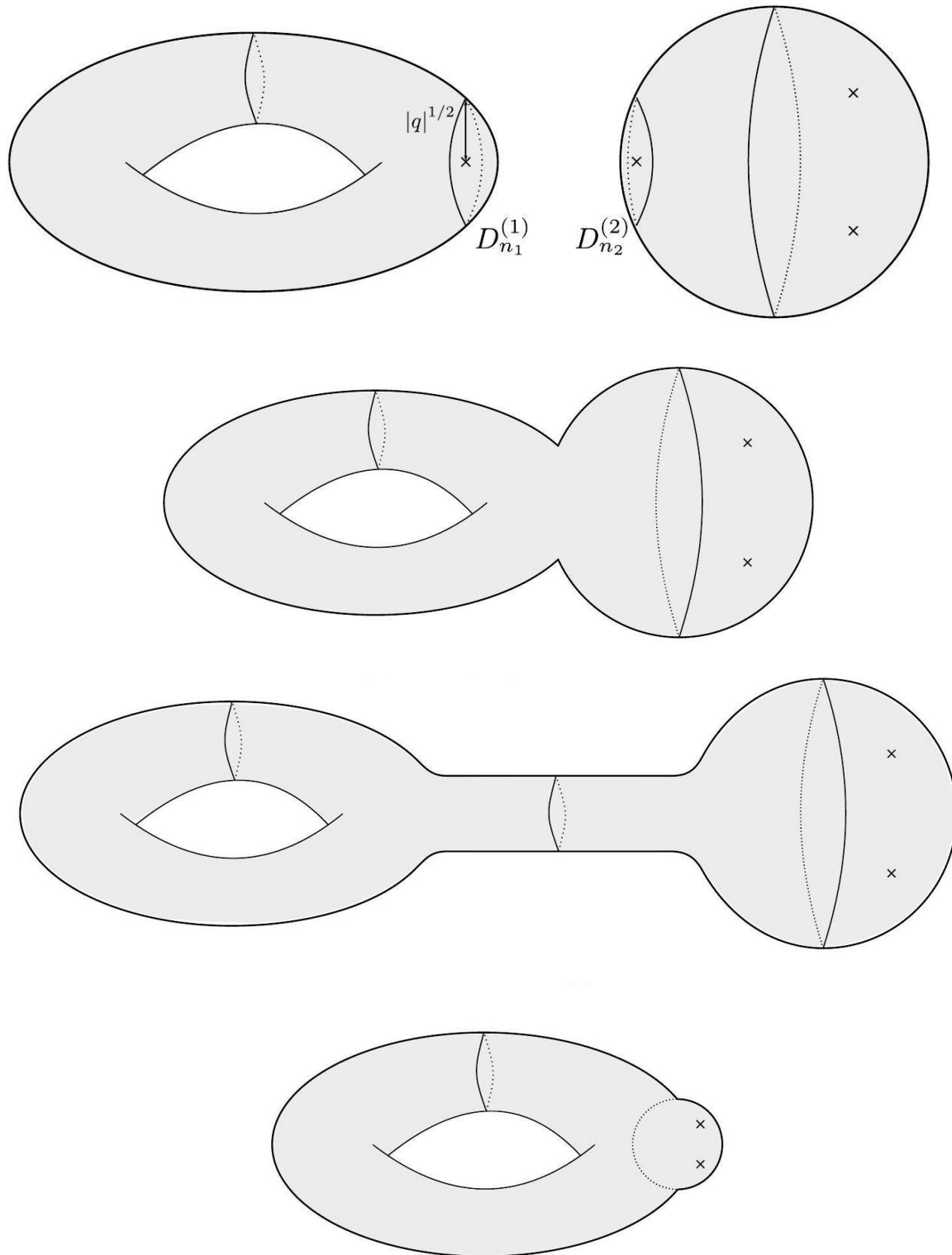
$$w_{n_1}^{(1)}w_{n_2}^{(2)}=q, |q|\leq 1,$$

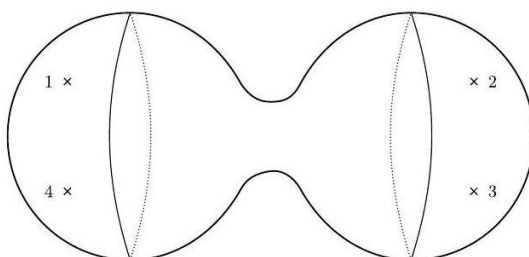
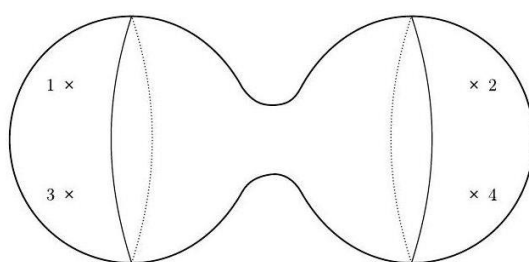
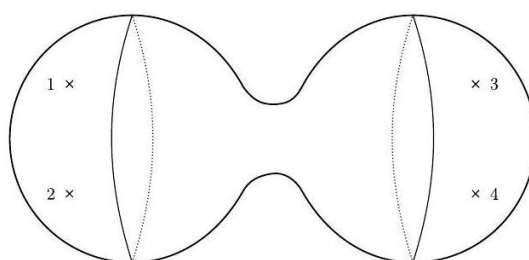
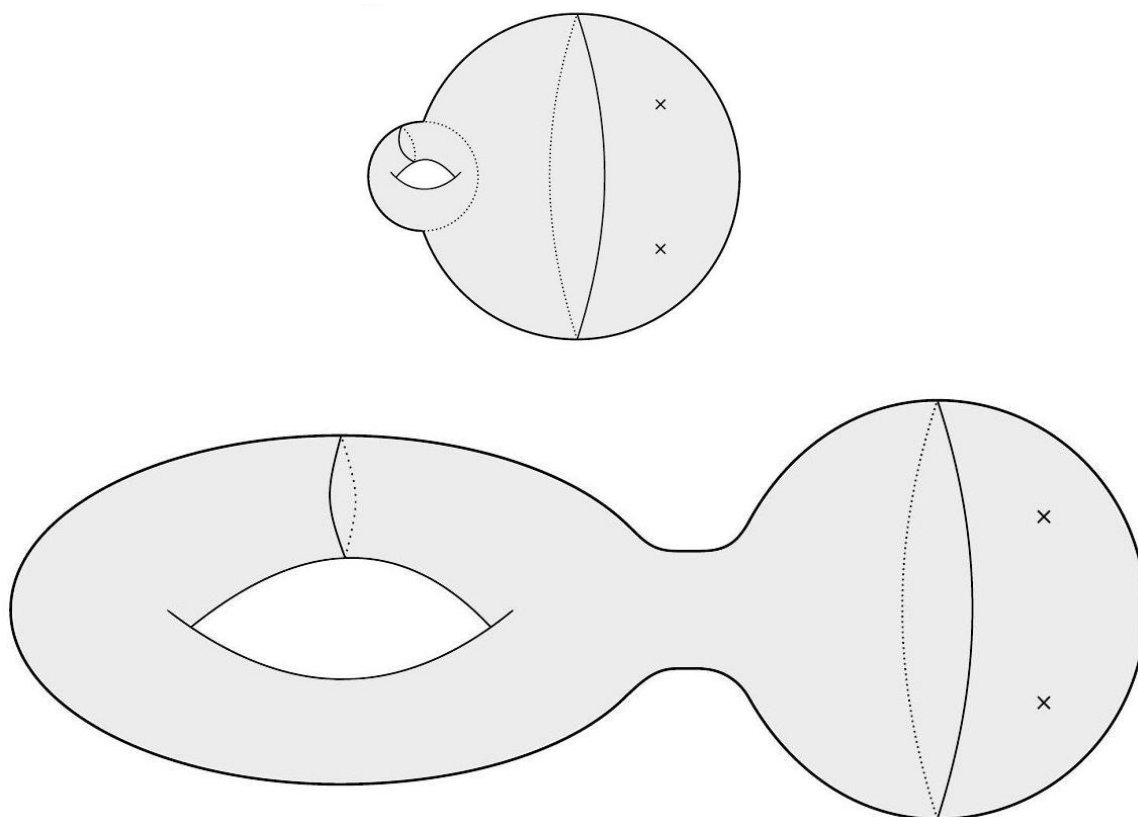
$$q = \mathrm{e}^{-s+\mathrm{i}\theta}, s \in \mathbb{R}_+, \theta \in [0,2\pi).$$

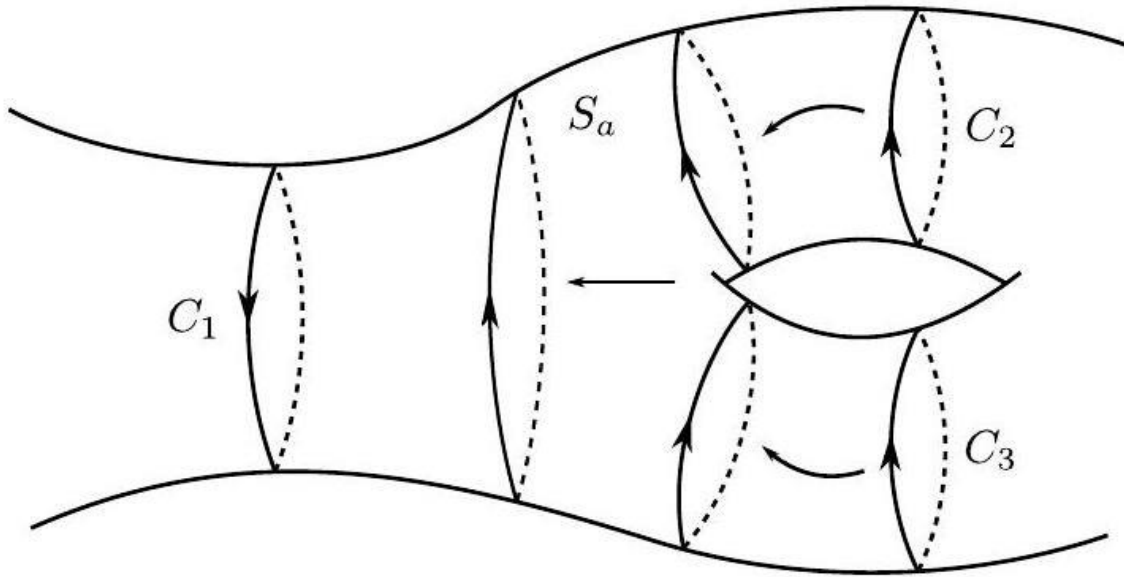


$$|w_{n_2}^{(2)}| = \frac{|q|}{|w_{n_1}^{(1)}|} > |q|^{1/2}$$

$$z_a^{(1)} = f_{an_1}^{(1)}(w_{n_1}^{(1)}), z_b^{(2)} = f_{bn_2}^{(2)}(w_{n_2}^{(2)})$$







Figuras 11, 12, 13, 14, 15, 16, 17, 18, 19 y 20. Puente de Einstein – Rosen a escala cuántica por curvatura local en dimensión $\mathbb{R}^4$, es decir, en la misma dimensión.

$$z_a^{(1)} = f_{an_1}^{(1)}(w_{n_1}^{(1)}) = f_{an_1}^{(1)}\left(\frac{q}{w_{n_2}^{(2)}}\right) = f_{an_1}^{(1)}\left(\frac{q}{f_{bn_2}^{(2)-1}(z_b^{(2)})}\right)$$

$$z_a = F_{ab}(z_b), F_{ab} = f_{an_1}^{(1)} \circ (q \cdot I) \circ f_{bn_2}^{(2)-1},$$

$$M_{g_1, n_1} + M_{g_2, n_2} = 6g_1 - 6 + 2n_1 + 6g_2 - 6 + 2n_2 = M_{g, n} - 2.$$

$$\Sigma_{g, n} = \Sigma_{g_1, n_1}, \begin{cases} g = g_1 + 1 \\ n = n_1 - 2 \end{cases}$$

$$w_{n_1-1}^{(1)} w_{n_1}^{(1)} = q$$

$$M_{g_1, n_1} = M_{g, n} - 2.$$

$$\mathcal{M}_{g_1, n_1} \mathcal{M}_{g_2, n_2} \subset \mathcal{M}_{g, n}$$

$$\mathcal{M}_{g_1, n_1} \subset \mathcal{M}_{g, n}$$

$$\mathcal{F}_{g, n} := \mathcal{M}_{g-1, n+2} \bigcup \left(\bigcup_{\substack{n_1+n_2=n+2 \\ g_1+g_2=g}} \mathcal{M}_{g_1, n_1} \# \mathcal{M}_{g_2, n_2} \right)$$

$$\mathcal{V}_{g, n} := \mathcal{M}_{g, n} - \mathcal{F}_{g, n}$$

$$\mathcal{F}_{g, n}^{1\text{PR}} := \bigcup_{\substack{n_1+n_2=n+2 \\ g_1+g_2=g}} \mathcal{M}_{g_1, n_1} \mathcal{M}_{g_2, n_2}$$

$$\mathcal{V}_{g, n}^{1\text{PI}} := \mathcal{M}_{g, n} - \mathcal{F}_{g, n}^{1\text{PR}}$$

$$\mathcal{V}_{g,n}^{1\text{PI}} = \mathcal{V}_{g,n} U \left(\bigcup_{g'} \# \mathcal{M}_{g-g', n+g'} \right)$$

$$\mathcal{V}_{0,3} = \mathcal{M}_{0,3}, \mathcal{F}_{0,3} = \emptyset.$$

$$\mathcal{F}_{0,4} = \mathcal{V}_{0,3} \mathcal{V}_{0,3}$$

$$\begin{aligned} \mathcal{F}_{0,5} &= \mathcal{M}_{0,4} \mathcal{M}_{0,3} \\ \mathcal{F}_{1,1} &= \mathcal{V}_{0,3} \end{aligned}$$

$$\begin{aligned} \mathcal{V}_{0,3} &= \mathcal{V}_{0,3} \mathcal{V}_{0,3} \mathcal{V}_{0,3} + \mathcal{V}_{0,4} \mathcal{V}_{0,3} \\ \mathcal{F}_{0,4} &= \mathcal{V}_{0,3} + \mathcal{V}_{0,4} \end{aligned}$$

$$r(\Sigma_{g_1,n_1} \# \Sigma_{0,3}) = r(\Sigma_{g_1,n_1}) + 1, r(\# \Sigma_{g_1,n_1}) = r(\Sigma_{g_1,n_1}) + 1.$$

$$r(\Sigma_{g,n}) = 3g + n - 2 \in \mathbb{N}^*$$

$$r(\Sigma_{0,3}) = 1$$

$$r(\Sigma_{g_1,n_1} \# \Sigma_{g_2,n_2}) = r(\Sigma_{g_1,n_1}) + r(\Sigma_{g_2,n_2}).$$

Stubs.

$$q = \mathrm{e}^{-s+\mathrm{i}\theta}, s \in [s_0, \infty), \theta \in [0, 2\pi), s_0 \geq 0.$$

$$s_0 < s'_0\colon \mathcal{F}_{g,n}(s'_0) \subset \mathcal{F}_{g,n}(s_0) \; \mathcal{V}_{g,n}(s_0) \subset \mathcal{V}_{g,n}(s'_0)$$

$$w_1=\lambda\tilde{w}_1, w_2=\lambda\tilde{w}_2$$

$$\tilde{w}_1\tilde{w}_2=\mathrm{e}^{-\tilde{s}+\mathrm{i}\tilde{\theta}}$$

$$\tilde{s}=s+2\ln\ |\lambda|, \tilde{\theta}=\theta+\mathrm{i}\ln\ \frac{\lambda}{\overline{\lambda}}$$

$$\tilde{s}\in[s_0,\infty), s_0\colon=2\ln\ |\lambda|.$$

Amplitudes y estados de superficie Off-shell.

$$\omega_{i_1\cdots i_p}:=\omega_p\left(\partial_{s_1},\ldots,\partial_{s_p}\right),$$

$$\omega_{i_1i_2\cdots i_p}=-\omega_{i_2i_1\cdots i_p},$$

$$\omega_p(V^{(1)},\ldots,V^{(p)})=\omega_p\left(V_{s_1}^{(1)}\partial_{s_1},\ldots,V_{s_p}^{(p)}\partial_{s_p}\right)=\omega_{i_1\cdots i_p}V_{s_1}^{(1)}\cdots V_{s_p}^{(p)},$$

$$\omega_p(\mathcal{V}_1,\ldots,\mathcal{V}_n)\colon=\omega_p(\otimes_i\mathcal{V}_i).$$

$$\omega_0 = (2\pi \mathrm{i})^{-\mathrm{M}_{g,n}^c} \left\langle \prod_{i=1}^n f_i \circ \mathcal{V}_i(0) \right\rangle_{\Sigma_{g,n}}$$

$$B(V)\colon=\oint_c\frac{\mathrm{d}z}{2\pi\mathrm{i}}b(z)v(z)+\oint_c\frac{\mathrm{d}\bar{z}}{2\pi\mathrm{i}}\bar{b}(\bar{z})\bar{v}(\bar{z})$$



$$B(V) := \sum_{\alpha} \oint_{c_{\alpha}} \frac{dz}{2\pi i} b(z) v(z) + \text{c.c.}$$

$$T(V) := \oint_c \frac{dz}{2\pi i} T(z) v(z) + \oint_c \frac{d\bar{z}}{2\pi i} \bar{T}(\bar{z}) \bar{v}(\bar{z})$$

$$T(V) = \{Q_B, B(V)\}$$

$$B = B_S \, dx_S, B_S := B(\partial_S) \\ B_S = \sum_{\alpha} \oint_{c_{\alpha}} \frac{d\sigma_{\alpha}}{2\pi i} b(\sigma_{\alpha}) \frac{\partial F_{\alpha}}{\partial x_S} (F_{\alpha}^{-1}(\sigma_{\alpha})) + \sum_{\alpha} \oint_{c_{\alpha}} \frac{d\bar{\sigma}_{\alpha}}{2\pi i} \bar{b}(\bar{\sigma}_{\alpha}) \frac{\partial \bar{F}_{\alpha}}{\partial x_S} (\bar{F}_{\alpha}^{-1}(\bar{\sigma}_{\alpha}))$$

$$\omega_p(V^{(1)},\ldots,V^{(p)})(\mathcal{V}_1,\ldots,\mathcal{V}_n) := (2\pi i)^{-M_{g,n}^c} \left\langle B(V^{(1)}) \cdots B(V^{(p)}) \prod_{i=1}^n \mathcal{V}_i \right\rangle_{\Sigma_{g,n}}$$

$$\omega_p = \omega_{p,s_1 \cdots s_p} \, dx_{s_1} \wedge \cdots \wedge dx_{s_p} \\ = (2\pi i)^{-M_{g,n}^c} \left\langle B_{s_1} \, dx_{s_1} \wedge \cdots \wedge B_{s_p} \, dx_{s_p} \prod_{i=1}^n \mathcal{V}_i \right\rangle_{\Sigma_{g,n}}$$

$$x_s = x_s(t_1,\ldots,t_q)$$

$$\forall p \leq q: \omega_p|_{\mathcal{S}} = (2\pi i)^{-M_{g,n}^c} \left\langle B_{r_1} \frac{\partial x_{s_1}}{\partial t_{r_1}} dt_{r_1} \wedge \cdots \wedge B_{r_p} \frac{\partial x_{s_p}}{\partial t_{r_p}} dt_{r_p} \prod_{i=1}^n \mathcal{V}_i \right\rangle_{\Sigma_{g,n}}, \\ \forall p > q: \omega_p \mid \mathcal{S} = 0.$$

$$B_r := \frac{\partial x_s}{\partial t_r} B_s.$$

$$A_{g,n}(\mathcal{V}_1,\ldots,\mathcal{V}_n)_{\mathcal{S}_{g,n}} := \int_{\mathcal{S}_{g,n}} \omega_{\mathbb{M}_{g,n}}^{g,n}(\mathcal{V}_1,\ldots,\mathcal{V}_n) \Big|_{\mathcal{S}_{g,n}} \\ \omega_{\mathbb{M}_{g,n}}^{g,n}(\mathcal{V}_1,\ldots,\mathcal{V}_n) \Big|_{\mathcal{S}_{g,n}} = (2\pi i)^{-M_{g,n}^c} \left\langle \bigwedge_{\lambda=1}^{M_{g,n}} B_s \frac{\partial x_s}{\partial t_{\lambda}} dt_{\lambda} \prod_{i=1}^n f_i \circ \mathcal{V}_i(0) \right\rangle$$

$$A_n(\mathcal{V}_1,\ldots,\mathcal{V}_n) := \sum_{g \geq 0} A_{g,n}(\mathcal{V}_1,\ldots,\mathcal{V}_n).$$

$$\mathcal{R}_{g,n}(\mathcal{V}_1,\ldots,\mathcal{V}_n) := \int_{\mathcal{R}_{g,n}} \omega_{\mathbb{M}_{g,n}}^{g,n}(\mathcal{V}_1,\ldots,\mathcal{V}_n),$$

$$\mathcal{R}_n := \sum_{g \geq 0} \mathcal{R}_{g,n},$$

$$\mathcal{R}_n(\mathcal{V}_1,\ldots,\mathcal{V}_n) := \sum_{g \geq 0} \mathcal{R}_{g,n}(\mathcal{V}_1,\ldots,\mathcal{V}_n) = \sum_{g \geq 0} \int_{\mathcal{R}_{g,n}} \omega_{\mathbb{M}_{g,n}}^{g,n}(\mathcal{V}_1,\ldots,\mathcal{V}_n).$$



$$\begin{aligned}\langle \Sigma_{g,n} | B_{s_1} \cdots B_{s_p} | \otimes_i \mathcal{V}_i \rangle &:= \omega_{s_1 \cdots s_p}(\mathcal{V}_1, \dots, \mathcal{V}_n), \\ \langle \omega_p^{g,n} | \otimes_i \mathcal{V}_i \rangle &:= \omega_p(\mathcal{V}_1, \dots, \mathcal{V}_n), \\ \langle A_{g,n} | \otimes_i \mathcal{V}_i \rangle &:= A_{g,n}(\mathcal{V}_1, \dots, \mathcal{V}_n).\end{aligned}$$

$$\langle \mathcal{R}_{g,n} | \otimes_i \mathcal{V}_i \rangle := \mathcal{R}_{g,n}(\mathcal{V}_1, \dots, \mathcal{V}_n).$$

$$\langle \omega_p^{g,n} | = \langle \Sigma_{g,n} | B_{s_1} \, dx^{s_1} \cdots B_{s_p} \, dx^{s_p}, \langle A_{g,n} | = \int_{\mathcal{M}_{g,n}} \langle \omega_p^{g,n} |.$$

$$\begin{array}{lll} C_1\colon w_1 = z_1 - y_1, & C_3\colon w_3 = z_2 - y_3, & C_5\colon z_1 = z_2 \\ C_2\colon w_2 = z_1 - y_2, & C_4\colon w_4 = z_2 - y_4. & \end{array}$$

$$y_4\longrightarrow y_4+\delta y_4, \bar{y}_4\longrightarrow \bar{y}_4+\delta \bar{y}_4$$

$$z_2'=z_2+\delta y_4, \bar{z}_2'=\bar{z}_2+\delta \bar{y}_4.$$

$$v=1, \bar{v}=0,$$

$$v=0, \bar{v}=1.$$

$$B(\partial_{y_4})=\oint_{c_4}\,dz_2b(z_2)(+1), B(\partial_{\bar{y}_4})=\oint_{c_4}\,d\bar{z}_2\bar{b}(\bar{z}_2)(+1)$$

$$\begin{aligned}\omega_2(\partial_{y_4},\partial_{\bar{y}_4}) &= \frac{1}{2\pi i}\left\langle B(\partial_{y_4})B(\partial_{\bar{y}_4})\prod_{i=1}^4\mathcal{V}_i\right\rangle_{\Sigma_{0,4}} \\ &= \frac{1}{2\pi i}\left\langle \oint_{c_4}\,dz_2b(z_2)\oint_{c_4}\,d\bar{z}_2\bar{b}(\bar{z}_2)\prod_{i=1}^4\mathcal{V}_i\right\rangle_{\Sigma_{0,4}}\end{aligned}$$

$$\omega_2(\partial_{y_4},\partial_{\bar{y}_4})=\frac{1}{2\pi i}\left\langle \prod_{i=1}^3 c\bar{c}V_i(y_i,\bar{y}_i)\oint_{c_4}\,dz_2b(z_2)\oint_{c_4}\,d\bar{z}_2\bar{b}(\bar{z}_2)\bar{c}(\bar{y}_4)c(y_4)V_4(y_4,\bar{y}_4)\right\rangle_{\Sigma_{0,4}}$$

$$\oint_{c_4}\,dz_2b(z_2)c(y_4)\sim \oint_{c_4}\,dz_2\frac{1}{z_2-y_4}$$

$$A_{0,4}=\frac{1}{2\pi i}\int\,dy_4\wedge\,d\bar{y}_4\left\langle \prod_{i=1}^3 c\bar{c}V_i(y_i)V_4(y_4)\right\rangle_{\Sigma_{0,4}}$$

$$z_a\longrightarrow z_a+\phi(z_a)$$

$$C_i\colon v^{(i)}=\phi|_{C_i}.$$

$$B(V)=\sum_{i=1}^3\oint_{c_i}\,dz_ab(z_a)\phi(z_a)+\text{c.c.}$$

$$w_i\longrightarrow (1+\mathrm{i}\alpha_i)w_i, \bar{w}_i\longrightarrow (1-\mathrm{i}\alpha_i)\bar{w}_i,$$



$$v = \mathrm{i} w_i, \bar{v} = -\mathrm{i} \bar{w}_i.$$

$$B(V) = \mathrm{i} \oint_{\mathcal{C}_i} \mathrm{d} w_i w_i b(w_i) - \mathrm{i} \oint_{\mathcal{C}_i} \mathrm{d} \bar{w}_i \bar{w}_i \bar{b}(\bar{w}_i),$$

$$B(V)\mathcal{V}_i(0) = \mathrm{i} \oint_{\mathcal{C}_i} \mathrm{d} w_i w_i b(w_i) \mathcal{V}_i(0) - \mathrm{i} \oint_{\mathcal{C}_i} \mathrm{d} \bar{w}_i \bar{w}_i \bar{b}(\bar{w}_i) \mathcal{V}_i(0),$$

$$\mathrm{i}(b_0-\bar{b}_0)|\mathcal{V}_i\rangle,$$

$$b_0^-|\mathcal{V}_i\rangle=0$$

$$\mathcal{V}_i\in\mathcal{H}^-,$$

$$w_i\longrightarrow f(w_i), f(0)=0$$

$$f(w_i)=\sum_{m\geq 0}p_mw_i^{m+1}$$

$$v_m=w_i^{m+1}, \bar{v}_m=0$$

$$B(\partial_{p_m})=\oint_{\mathcal{C}_i}\mathrm{d} w_i b(w_i)w_i^{m+1}$$

$$\forall m\geq 0: b_m|\mathcal{V}_i\rangle=0, \bar{b}_m|\mathcal{V}_i\rangle=0$$

$$\omega_p\left(\sum_i Q_B^{(i)}\otimes_i \mathcal{V}_i\right)=(-1)^p\,\mathrm{d}\omega_{p-1}(\otimes \mathcal{V}_i)$$

$$Q_B^{(i)}=1_{i-1}\otimes Q_B\otimes 1_{n-i}$$

$$Q_B\mathcal{V}_i(z,\bar{z})=\frac{1}{2\pi\mathrm{i}}\oint\,\mathrm{d} w j_B(w)\mathcal{V}_i(z,\bar{z})+\,\text{c.c.}$$

$$\begin{aligned}\omega_p\left(\sum_i Q_B^{(i)}\otimes_i \mathcal{V}_i\right) &= \omega_p(Q_B\mathcal{V}_1,\mathcal{V}_2,\ldots,\mathcal{V}_n)+(-1)^{|\mathcal{V}_1|}\omega_p(\mathcal{V}_1,Q_B\mathcal{V}_2,\ldots,\mathcal{V}_n) \\ &+ \cdots + (-1)^{|\mathcal{V}_1|+\cdots+|\mathcal{V}_{n-1}|}\omega_p(\mathcal{V}_1,\mathcal{V}_2,\ldots,Q_B\mathcal{V}_n)\end{aligned}$$

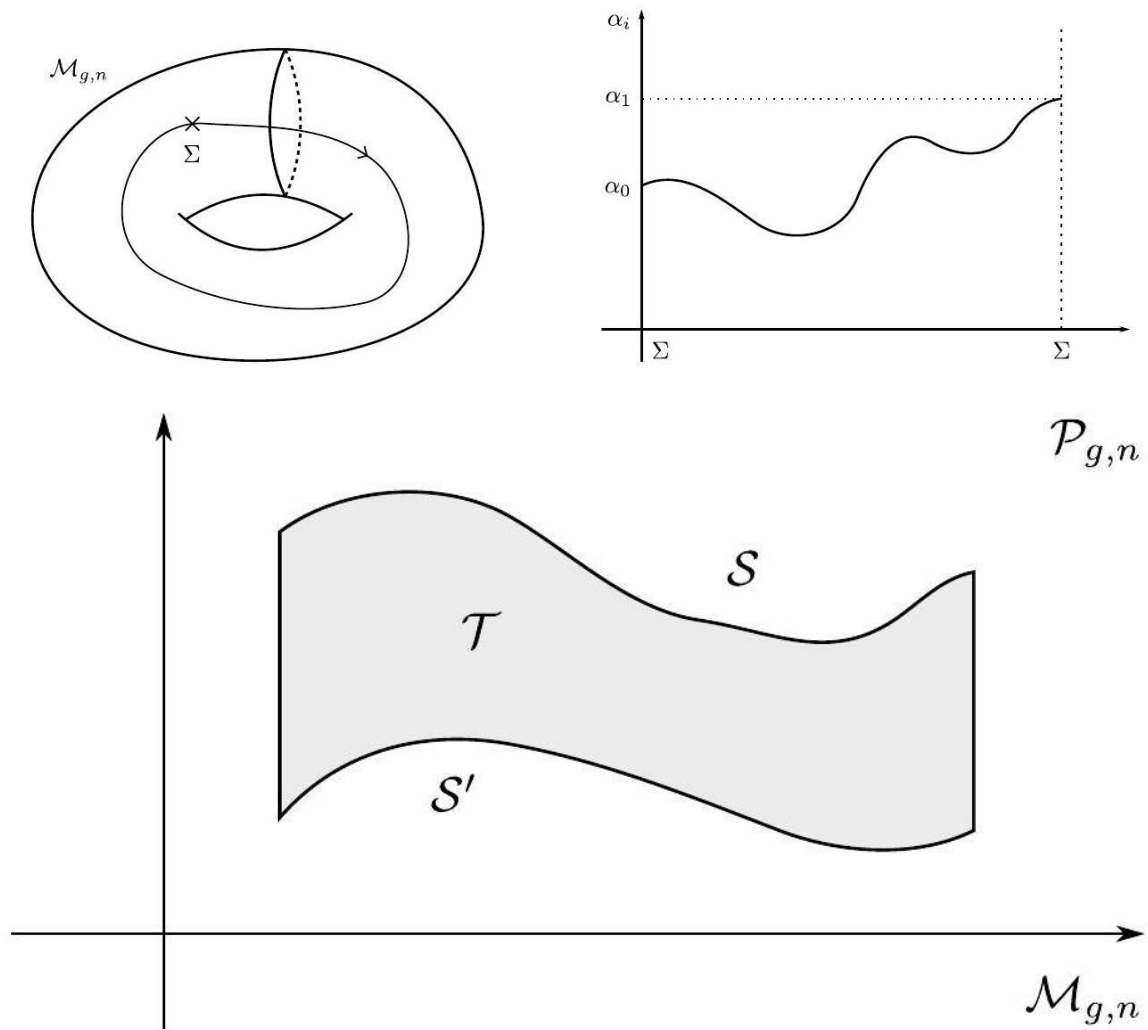
$$T_s=\{Q_B,B_s\}$$

$$\mathrm{d} x_s\{Q_B,B_s\}=\mathrm{d} x_sT_s=\mathrm{d} x_s\partial_s=\mathrm{d}$$

$$N_{\mathrm{gh}}\left(\omega_p(\mathcal{V}_1,\ldots,\mathcal{V}_n)\right)=\sum_{i=1}^nN_{\mathrm{gh}}(\mathcal{V}_i)-p=6-6g$$

$$N_{\mathrm{gh}}\left(\omega_{\mathrm{M}_{g,n}}\right)=6-6g\Rightarrow\sum_{i=1}^nN_{\mathrm{gh}}(\mathcal{V}_i)=2n$$





Figuras 21 y 22. Deformación local del espacio – tiempo cuántico y fluctuaciones de materia y energía de la partícula supermasiva.

$$L_0^- |\mathcal{V}_i\rangle = 0$$

$$b_0^- |\mathcal{V}_i\rangle = 0$$

$$\mathcal{V}_i \in \mathcal{H}^- \cap \ker L_0^-,$$

$$\mathcal{V}_i(0) \rightarrow (e^{i\alpha_i})^h (e^{-i\alpha_i})^{\bar{h}} \mathcal{V}_i(0)$$

$$|\mathcal{V}_i\rangle \rightarrow e^{i\alpha_i(L_0 - \bar{L}_0)} |\mathcal{V}_i\rangle$$

$$\forall i: Q_B |\mathcal{V}_i\rangle = 0 \Rightarrow d\omega_{p-1}(\mathcal{V}_1, \dots, \mathcal{V}_n) = 0.$$

$$\int_{\mathcal{S}} \omega_{\mathcal{M}_{g,n}} - \int_{\mathcal{S}'} \omega_{\mathcal{M}_{g,n}} = \int_{\partial \mathcal{T}} \omega_{\mathcal{M}_{g,n-1}} = \int_{\mathcal{T}} d\omega_{\mathcal{M}_{g,n-1}} = 0$$

$$|\mathcal{V}_1\rangle = Q_B |\Lambda\rangle, Q_B |\mathcal{V}_i\rangle = 0$$

$$\omega_{M_{g,n}}(Q_B \Lambda, \mathcal{V}_2, \dots, \mathcal{V}_n) = d\omega_{M_{g,n-1}}(\Lambda, \mathcal{V}_2, \dots, \mathcal{V}_n),$$

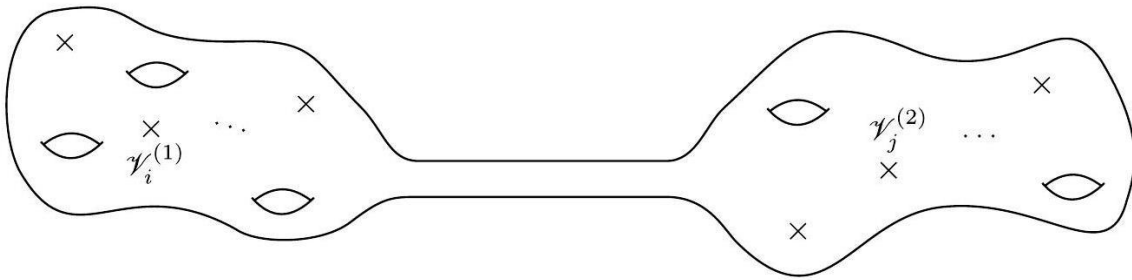
$$\int_S \omega_{M_{g,n}}(Q_B \Lambda, \mathcal{V}_2, \dots, \mathcal{V}_n) = \int_S d\omega_{M_{g,n-1}}(\Lambda, \mathcal{V}_2, \dots, \mathcal{V}_n) = \int_{\partial S} \omega_{M_{g,n-1}}(\Lambda, \mathcal{V}_2, \dots, \mathcal{V}_n)$$

$$\int_S \omega_{M_{g,n}}(Q_B \Lambda, \mathcal{V}_2, \dots, \mathcal{V}_n) = 0$$

Factorización de amplitudes y diagramas de Feynman.

$$w_{n_1}^{(1)} w_{n_2}^{(2)} = q.$$

$$A_{g,n} = \int_{S_{g,n}} \omega_{M_{g,n}}^{g,n}(\mathcal{V}_1^{(1)}, \dots, \mathcal{V}_{n_1-1}^{(1)}, \mathcal{V}_1^{(2)}, \dots, \mathcal{V}_{n_2-1}^{(2)}).$$



$$w_{n_1}'^{(1)} = w_{n_1}^{(1)} + \frac{w_{n_1}^{(1)}}{q} \delta q$$

$$v_q = \frac{w_{n_1}^{(1)}}{q}, B_q = \frac{1}{q} \oint_{c_q} dw_{n_1}^{(1)} b(w_{n_1}^{(1)}) w_{n_1}^{(1)}$$

$$w_{n_1}'^{(1)} w_{n_2}^{(2)} = q + \delta q$$

$$w_{n_1}^{(1)} = \frac{w_{n_1}^{(1)}}{q} (q + \delta q) = w_{n_1}^{(1)} + \frac{q}{w_{n_1}^{(1)}} \delta q$$

$$\begin{aligned} & \omega_{M_{g,n}} \left(V_1^{(1)}, \dots, V_{M_{g_1,n_1}}^{(1)}, \partial_q, \partial_{\bar{q}}, V_1^{(2)}, \dots, V_{M_{g_2,n_2}}^{(2)} \right) \\ &= (2\pi i)^{-M_{g,n}^c} \left\langle \prod_{\lambda=1}^{M_{g_1,n_1}} B(V_\lambda^{(1)}) B(\partial_q) B(\partial_{\bar{q}}) \prod_{\kappa=1}^{M_{g_2,n_2}} B(V_\kappa^{(2)}) \prod_{i=1}^{n_1-1} \mathcal{V}_i^{(1)} \prod_{j=1}^{n_2-1} \mathcal{V}_j^{(2)} \right\rangle_{\Sigma_{g,n}} \end{aligned}$$

$$\left\langle \Sigma_{n_1} \mid \mathcal{V}_{n_1}^{(1)} \right\rangle := \omega_{M_{g_1,n_1}}(\mathcal{V}_1^{(1)}, \dots, \mathcal{V}_{n_1}^{(1)}) = (2\pi i)^{-M_{g_1,n_1}^c} \left\langle \prod_{\lambda=1}^{M_{g_1,n_1}} B(V_\lambda^{(1)}) \prod_{i=1}^{n_1-1} \mathcal{V}_i^{(1)} \right\rangle_{\Sigma_{g_1,n_1}}$$

$$\left\langle \Sigma_{n_2} \mid \mathcal{V}_{n_2}^{(2)} \right\rangle := \omega_{M_{g_2,n_2}}(\mathcal{V}_1^{(2)}, \dots, \mathcal{V}_{n_2}^{(2)}) = (2\pi i)^{-M_{g_2,n_2}^c} \left\langle \prod_{\lambda=1}^{M_{g_2,n_2}} B(V_\lambda^{(2)}) \prod_{j=1}^{n_2-1} \mathcal{V}_j^{(2)} \right\rangle_{\Sigma_{g_2,n_2}}$$

$$\langle \Sigma_{n_1} | = \langle 0 | I \circ \Sigma_{n_1}(0), \langle \Sigma_{n_2} | = \langle 0 | I \circ \Sigma_{n_2}(0),$$

$$\left\langle \Sigma_{n_1} \mid \mathcal{V}_{n_1}^{(1)} \right\rangle = \left\langle I \circ \Sigma_{n_1}(0) \mathcal{V}_{n_1}(0) \right\rangle_{w_{n_1}^{(1)}}, \left\langle \Sigma_{n_2} \mid \mathcal{V}_{n_2}^{(2)} \right\rangle = \left\langle I \circ \Sigma_{n_2}(0) \mathcal{V}_{n_2}(0) \right\rangle_{w_{n_2}^{(2)}}$$

$$w_1=I\left(w_{n_1}^{(1)}\right), w_2=I\left(w_{n_2}^{(2)}\right)$$

$$w_{n_1}^{(1)}=\frac{q}{w_{n_2}^{(2)}}=\frac{q}{I(w_2)}=qw_2:=f(w_2)$$

$$\omega_{\mathbf{M}_{g,n}}=\frac{1}{2\pi\mathrm{i}}\left\langle I\circ\Sigma_{n_1}(0)B_qB_{\bar{q}}f\circ\Sigma_{n_2}(0)\right\rangle_{w_{n_1}^{(1)}}=\frac{1}{2\pi\mathrm{i}}\left\langle \Sigma_{n_1}\mid B_qB_{\bar{q}}q^{L_0}\bar{q}^{\bar{L}_0}\mid\Sigma_2\right\rangle,$$

$$\left\langle \Sigma_{n_1} \mid B_q B_{\bar{q}} \mid \mathcal{V}_{n_1}^{(1)} \right\rangle = \frac{1}{q \bar{q}} \left\langle \Sigma_{n_1} \mid b_0 \bar{b}_0 \mid \mathcal{V}_{n_1}^{(1)} \right\rangle$$

$$B_q\mathcal{V}_{n_1}^{(1)}(z,\bar{z})=\frac{1}{q}\oint_{c_q}\mathrm{d}w_{n_1}^{(1)}b\left(w_{n_1}^{(1)}\right)w_{n_1}^{(1)}\mathcal{V}_{n_1}^{(1)}(z,\bar{z})\longrightarrow\frac{1}{q}b_0\left|\mathcal{V}_{n_1}^{(1)}\right\rangle$$

$$\omega_{\mathbf{M}_{g,n}}=\frac{1}{2\pi\mathrm{i}}\frac{1}{q\bar{q}}\left\langle \Sigma_{n_1}\mid b_0\bar{b}_0q^{L_0}\bar{q}^{\bar{L}_0}\mid\Sigma_{n_2}\right\rangle.$$

$$\mathcal{F}_{g,n}\left(\mathcal{V}_i^{(1)}\mid\mathcal{V}_j^{(2)}\right):=\frac{1}{2\pi\mathrm{i}}\int\bigwedge_{\lambda=1}^{M_{g_1,n_1}}\mathrm{d}t_{\lambda}^{(1)}\bigwedge_{\kappa=1}^{M_{g_2,n_2}}\mathrm{d}t_{\kappa}^{(2)}\wedge\frac{\mathrm{d}q}{q}\wedge\frac{\mathrm{d}\bar{q}}{\bar{q}}\left\langle \Sigma_{n_1}\mid b_0\bar{b}_0q^{L_0}\bar{q}^{\bar{L}_0}\mid\Sigma_{n_2}\right\rangle.$$

$$\mathcal{F}_{g,n}\left(\mathcal{V}_i^{(1)}\mid\mathcal{V}_j^{(2)}\right)$$

$$=\frac{1}{2\pi\mathrm{i}}\int\frac{\mathrm{d}^Dk}{(2\pi)^D}\frac{\mathrm{d}^Dk'}{(2\pi)^D}(-1)^{|\phi_{\alpha}|}\times\int\frac{\mathrm{d}q}{q}$$

$$\wedge\frac{\mathrm{d}\bar{q}}{\bar{q}}\left\langle \phi_{\alpha}(k)^c\mid b_0\bar{b}_0q^{L_0}\bar{q}^{\bar{L}_0}\mid\phi_{\beta}(k')^c\right\rangle$$

$$\times\int\bigwedge_{\lambda=1}^{M_{g_1,n_1}}\mathrm{d}t_{\lambda}^{(1)}\left\langle \Sigma_{n_1}\mid\phi_{\alpha}(k)\right\rangle\bigwedge_{\kappa=1}^{M_{g_2,n_2}}\mathrm{d}t_{\kappa}^{(2)}\left\langle \phi_{\beta}(k')\mid\Sigma_{n_2}\right\rangle$$

$$A_{g_1,n_1}\left(\mathcal{V}_1^{(1)},\ldots,\mathcal{V}_{n_1-1}^{(1)},\phi_{\alpha}(k)\right)=\int_{\mathcal{S}_{g_1,n_1}}\omega_{\mathbf{M}_{g_1,n_1}}\left(\mathcal{V}_1^{(1)},\ldots,\mathcal{V}_{n_1-1}^{(1)},\phi_{\alpha}(k)\right)$$

$$=\int_{\mathcal{S}_{g_1,n_1}}\bigwedge_{\lambda=1}^{M_{g_1,n_1}}\mathrm{d}t_{\lambda}^{(1)}\left\langle \Sigma_{n_1}\mid\phi_{\alpha}(k)\right\rangle$$



$$A_{g_2, n_2} \left(\mathcal{V}_1^{(2)}, \dots, \mathcal{V}_{n_2-1}^{(2)}, \phi_\beta(k') \right) = \int_{\mathcal{S}_{g_2, n_2}} \omega_{M_{g_2, n_2}} \left(\mathcal{V}_1^{(2)}, \dots, \mathcal{V}_{n_2-2}^{(2)}, \phi_\beta(k') \right) \\ = \int_{\mathcal{S}_{g_2, n_2}} \bigwedge_{\lambda=1}^{M_{g_2, n_2}} dt_\lambda^{(2)} \langle \Sigma_{n_2} \mid \phi_\beta(k') \rangle$$

$$\Delta_{\alpha\beta}(k,k')\!:=\!\Delta(\phi_\alpha(k)^c,\phi_\beta(k')^c)\!:=\!\frac{1}{2\pi i}\int\frac{dq}{q}\wedge\frac{d\bar{q}}{\bar{q}}\langle\phi_\alpha(k)^c|b_0\bar{b}_0q^{L_0}\bar{q}^{\bar{L}_0}|\phi_\beta(k')^c\rangle$$

$$\mathcal{F}_{g,n}\left(\mathcal{V}_i^{(1)}\mid\mathcal{V}_j^{(2)}\right)=\int\frac{\mathrm{d}^Dk}{(2\pi)^D}\frac{\mathrm{d}^Dk'}{(2\pi)^D}A_{g_1,n_1}\left(\mathcal{V}_1^{(1)},\dots,\mathcal{V}_{n_1-1}^{(1)},\phi_\alpha(k)\right)\Delta_{\alpha\beta}(k,k') \\ \times A_{g_2,n_2}\left(\mathcal{V}_1^{(2)},\dots,\mathcal{V}_{n_2-1}^{(2)},\phi_\beta(k')\right)$$

$$A_{g_1,n_1}\sim\delta^{(D)}\left(k_1^{(1)}+\dots+k_{n_1-1}^{(1)}+k\right),A_{g_2,n_2}\sim\delta^{(D)}\left(k_1^{(2)}+\dots+k_{n_2-1}^{(2)}+k'\right).$$

$$\mathcal{F}_{g,n}\sim\delta^{(D)}\left(k_1^{(1)}+\dots+k_{n_1-1}^{(1)}+k_1^{(2)}+\dots+k_{n_2-1}^{(2)}\right)$$

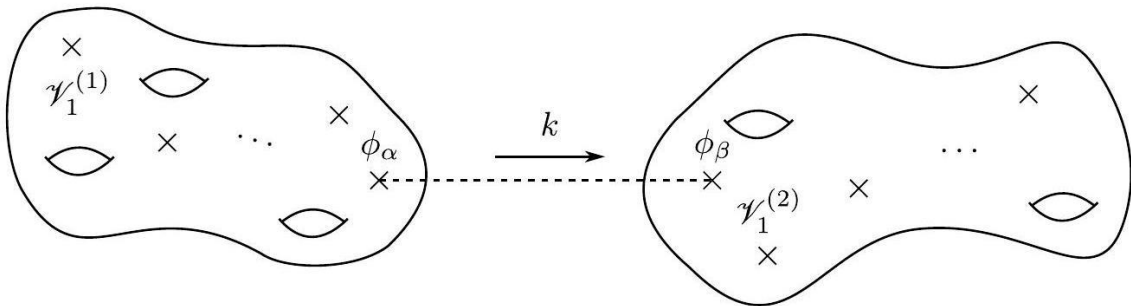
$$N_{\mathrm{gh}}(\phi_\alpha)=2n_1-\sum_{i=1}^{n_1-1}N_{\mathrm{gh}}\left(\mathcal{V}_i^{(1)}\right),N_{\mathrm{gh}}(\phi_\beta)=2n_2-\sum_{j=1}^{n_2-1}N_{\mathrm{gh}}\left(\mathcal{V}_j^{(2)}\right)$$

$$N_{\mathrm{gh}}(\phi_\alpha)+N_{\mathrm{gh}}(\phi_\beta)=4$$

$$N_{\mathrm{gh}}(\phi_\alpha)=N_{\mathrm{gh}}(\phi_\beta)=2$$

$$w_{n_1-1}w_{n_1}=q$$

$$A_{g,n}=\int_{\mathcal{S}_{g,n}}\omega_{M_{g,n}}^{g,n}\left(\mathcal{V}_1^{(1)},\dots,\mathcal{V}_{n_1-2}^{(1)}\right)$$



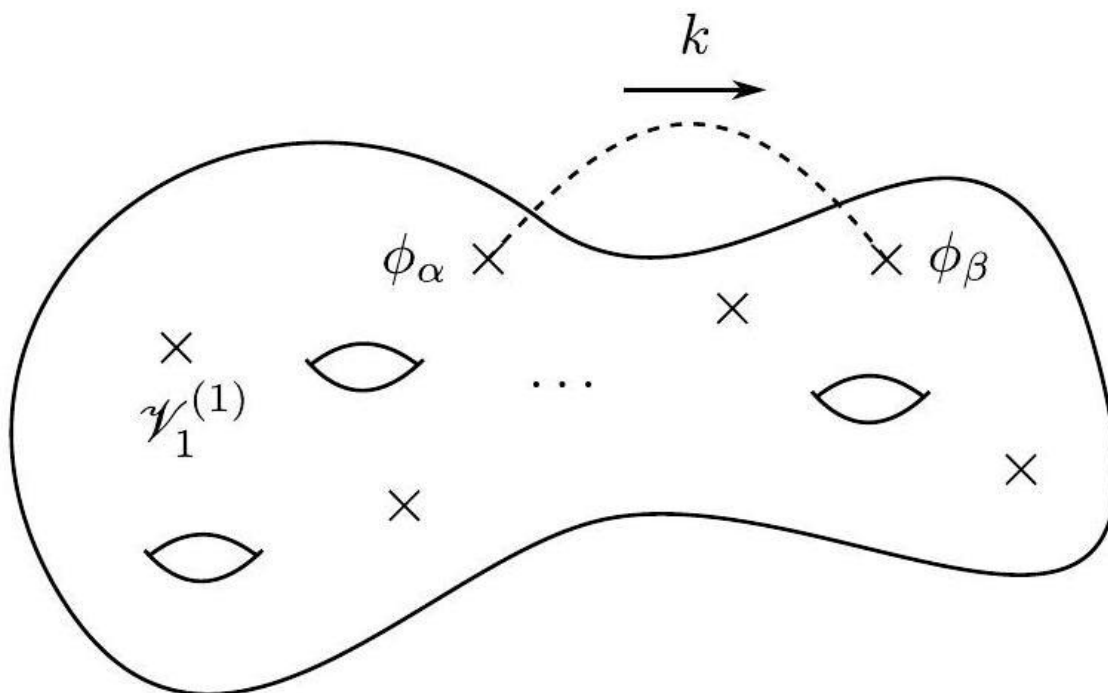
$$\omega_{M_{g,n}}\left(V_1^{(1)},\dots,V_{M_{g_1,n_1}}^{(1)},\partial_q,\partial_{\bar{q}}\right)=(2\pi i)^{-M_{g,n}}\left\langle\prod_{\lambda=1}^{M_{g_1,n_1}}B\left(V_{\lambda}^{(1)}\right)B(\partial_q)B(\partial_{\bar{q}})\prod_{i=1}^{n_1-2}\mathcal{V}_i^{(1)}\right\rangle_{\Sigma_{g,n}}$$

$$\left\langle \Sigma_{n_1-1,n_1} \mid \mathcal{V}_{n_1-1}^{(1)} \otimes \mathcal{V}_{n_1}^{(1)} \right\rangle := \omega_{M_{g_1,n_1}} \left(\mathcal{V}_1^{(1)}, \dots, \mathcal{V}_{n_1}^{(1)} \right)$$

$$\mathcal{F}_{g,n}\left(\mathcal{V}_i^{(1)}\mid\right)=\int\frac{\mathrm{d}^Dk}{(2\pi)^D}\frac{\mathrm{d}^Dk'}{(2\pi)^D}A_{g_1,n_1}\left(\mathcal{V}_1^{(1)},\ldots,\mathcal{V}_{n_1-2}^{(1)},\phi_\alpha(k),\phi_\beta(k')\right)\Delta_{\alpha\beta}(k,k')$$

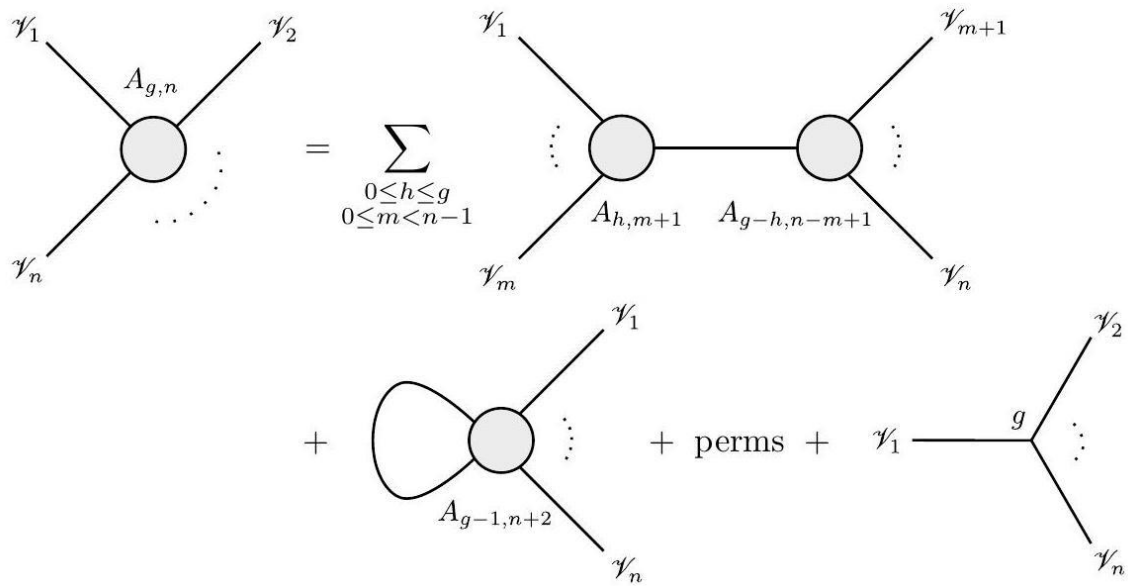
$$N_{\mathrm{gh}}(\phi_\alpha)+N_{\mathrm{gh}}(\phi_\beta)=2n_1-\sum_{i=1}^{n_1-2}N_{\mathrm{gh}}\left(\mathcal{V}_i^{(1)}\right)=4$$

$$A_{g_1,n_1}\left(\mathcal{V}_1^{(1)},\ldots,\mathcal{V}_{n_1-2}^{(1)},\phi_\alpha(k),\phi_\beta(-k)\right)\sim\delta^{(D)}\left(k_1^{(1)}+\cdots+k_{n_1-2}^{(1)}\right)$$



$$r(A_{g,n})\colon=r(\Sigma_{g,n})=3g+n-2$$

$$\mathcal{V}_{g,n}(\mathcal{V}_1,\ldots,\mathcal{V}_n)\colon=\mathcal{V}_1\stackrel{g}{\rightarrow}/\mathcal{V}_n^2\colon=\int_{\mathcal{R}_{g,n}}\omega_{\mathbf{M}_{g,n}}^{g,n}(\mathcal{V}_1,\ldots,\mathcal{V}_n),$$



$$\mathcal{V}_{g,n}(\mathcal{V}_1, \dots, \mathcal{V}_n) := \mathcal{V}_{g,n}(\otimes_i \mathcal{V}_i) := \{\mathcal{V}_1, \dots, \mathcal{V}_n\}_g.$$

Escalar QFT.

Propagador.

$$\Delta = \frac{1}{2\pi i} \int \frac{dq}{q} \wedge \frac{d\bar{q}}{\bar{q}} b_0 \bar{b}_0 q^{L_0} \bar{q}^{\bar{L}_0}$$

$$q = e^{-s+i\theta}, s \in \mathbb{R}_+, \theta \in [0, 2\pi)$$

$$\frac{dq}{q} \wedge \frac{d\bar{q}}{\bar{q}} = -2i ds \wedge d\theta$$

$$\Delta = \frac{1}{2\pi} b_0^+ b_0^- \int_0^\infty ds e^{-sL_0^+} \int_0^{2\pi} d\theta e^{i\theta L_0^-}$$

$$\int_0^\infty ds e^{-sL_0^+} = \frac{1}{L_0^+}, \int_0^{2\pi} d\theta e^{i\theta L_0^-} = 2\pi \delta_{L_0^-, 0}$$

$$\Delta = \frac{b_0^+}{L_0^+} b_0^- \delta_{L_0^-, 0}$$

$$\Delta = -2b_0 \bar{b}_0 \frac{1}{L_0 + \bar{L}_0} \delta_{L_0, \bar{L}_0}$$

$$L_0^+ |\phi_\alpha(k)\rangle = \frac{\alpha'}{2} (k^2 + m_\alpha^2) |\phi_\alpha(k)\rangle, L_0^- |\phi_\alpha(k)\rangle = 0$$

$$\Delta_{\alpha\beta}(k, k') := \langle \phi_\alpha(k)^c | \Delta | \phi_\beta(k')^c \rangle := (2\pi)^D \delta^{(D)}(k + k') \Delta_{\alpha\beta}(k),$$

$$\Delta_{\alpha\beta}(k) := \frac{M_{\alpha\beta}(k)}{k^2 + m_\alpha^2}, M_{\alpha\beta}(k) := \frac{2}{\alpha'} \langle \phi_\alpha^c(k) | b_0^+ b_0^- | \phi_\beta^c(-k) \rangle,$$

$$b_0^+ | \phi_\alpha^c \rangle \neq 0, b_0^- | \phi_\alpha^c \rangle \neq 0$$

$$c_0^+ | \phi_\alpha^c \rangle = 0, c_0^- | \phi_\alpha^c \rangle = 0$$

$$| \phi_\alpha^c \rangle = | \phi_1 \rangle + b_0^\pm | \phi_2 \rangle, c_0^\pm | \phi_1 \rangle = c_0^\pm | \phi_2 \rangle = 0$$

$$c_0^\pm | \phi_\alpha^c \rangle = 0 \Rightarrow | \phi_2 \rangle = 0$$

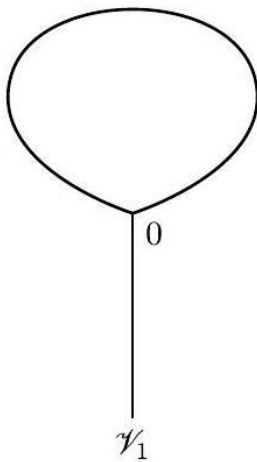
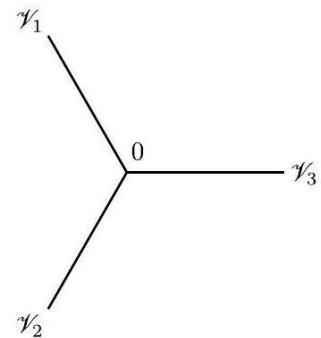
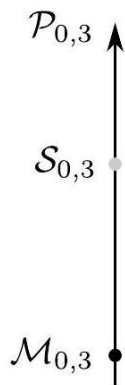
$$b_0^+ | \phi_\alpha \rangle = 0, b_0^- | \phi_\alpha \rangle = 0$$

$$b_0^+ | \mathcal{V}_i \rangle = 0$$

$$\Delta^{-1} = c_0^+ c_0^- L_0^+ \delta_{L_0^-, 0}$$

Vértices fundamentales.

$$\mathcal{V}_{0,3}(\mathcal{V}_1, \mathcal{V}_2, \mathcal{V}_3) := A_{0,3}(\mathcal{V}_1, \mathcal{V}_2, \mathcal{V}_3) = \omega_0^{0,3}(\mathcal{V}_1, \mathcal{V}_2, \mathcal{V}_3).$$



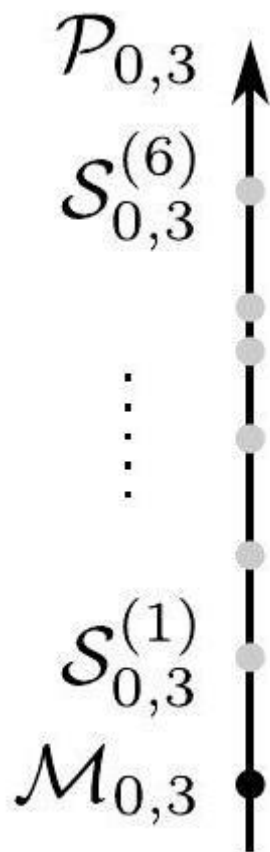
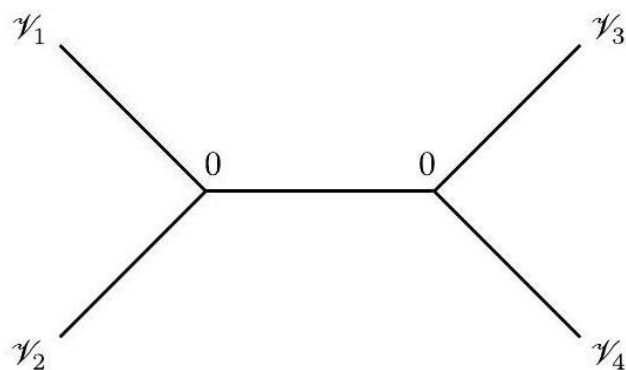


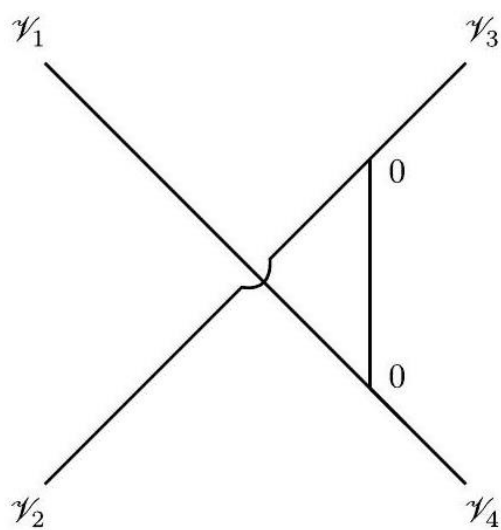
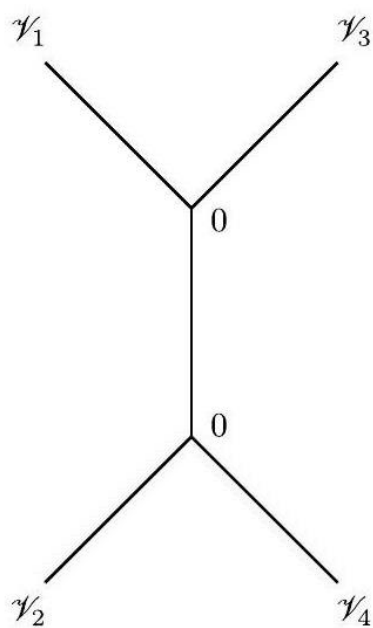
Figura 23, 23A y 23B. Operadores de vértice.

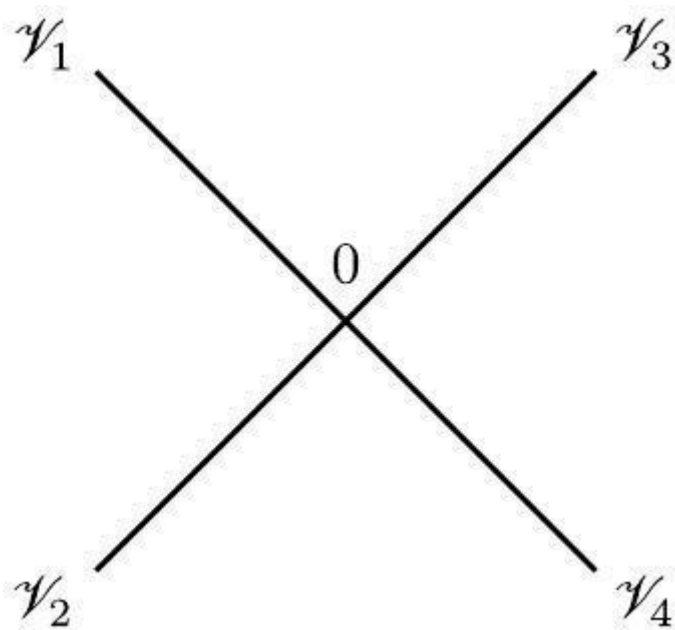
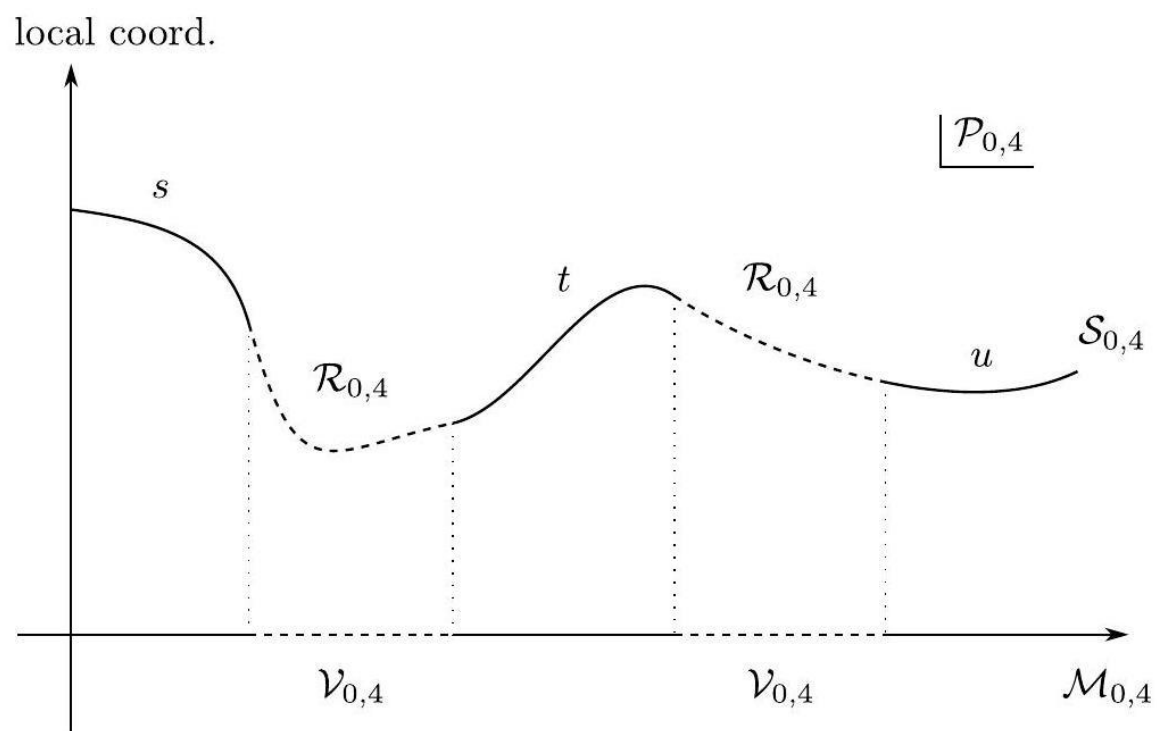
$$\mathcal{V}_{0,4}(\mathcal{V}_1, \mathcal{V}_2, \mathcal{V}_3, \mathcal{V}_4) := \int_{\mathcal{R}_{0,4}} \omega_2^{0,4}(\mathcal{V}_1, \dots, \mathcal{V}_4)$$

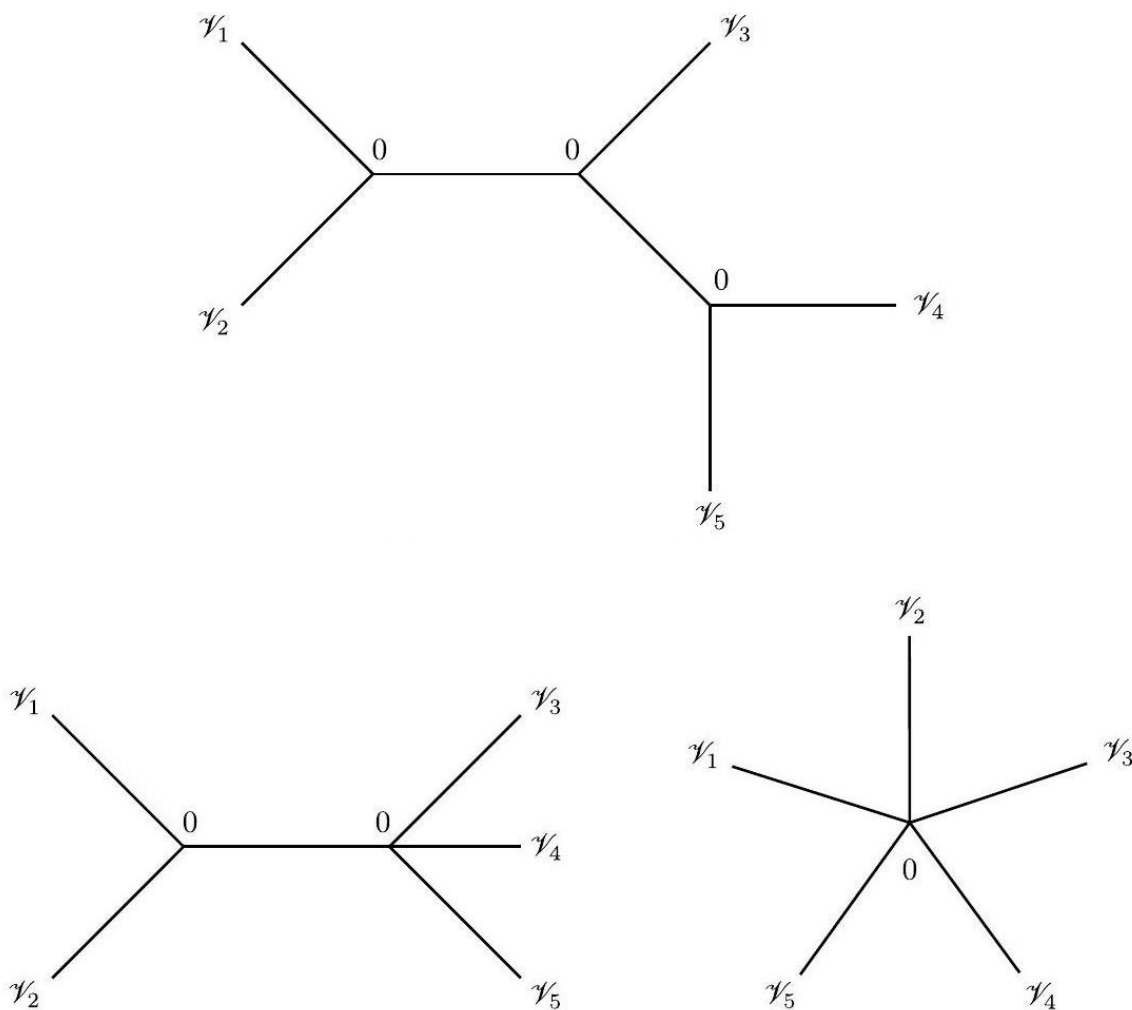
$$A_{0,4} = \mathcal{F}_{0,4}^{(s)} + \mathcal{F}_{0,4}^{(t)} + \mathcal{F}_{0,4}^{(u)} + \mathcal{V}_{0,4}$$

$$\mathcal{V}_{0,5}(\mathcal{V}_1, \dots, \mathcal{V}_5) := \int_{\mathcal{R}_{0,5}} \omega_4^{0,5}(\mathcal{V}_1, \dots, \mathcal{V}_5)$$









Figuras 24, 25, 26, 27, 28, 29 y 30. Diagramas de Feynman de una partícula supermasiva.

$$\mathcal{V}_{0,2} := \Delta^{-1} \Delta \Delta^{-1} = \Delta^{-1}$$

$$\mathcal{V}_{0,2}(\mathcal{V}_1, \mathcal{V}_2) := \langle \mathcal{V}_1 | c_0^+ c_0^- L_0^+ \delta_{L_0^-, 0} | \mathcal{V}_2 \rangle.$$

$$\mathcal{V}_n(\mathcal{V}_1, \dots, \mathcal{V}_n) := \sum_{g \geq 0} (\hbar g_s^2)^g \mathcal{V}_{g,n}(\mathcal{V}_1, \dots, \mathcal{V}_n)$$

$$q = e^{-s+i\theta}, s \in [s_0, \infty), \theta \in [0, 2\pi).$$

$$\int_{s_0}^{\infty} ds e^{-sL_0^+} = \frac{e^{-s_0L_0^+}}{L_0^+}$$

$$\Delta(s_0) = b_0^+ \frac{e^{-s_0L_0^+}}{L_0^+} b_0^- \delta_{L_0^-, 0}.$$

$$\Delta_{\alpha\beta}(k)\!:=\!\frac{\mathrm{e}^{\frac{\alpha' s_0}{2}(k^2+m_\alpha^2)}}{k^2+m_\alpha^2}M_{\alpha\beta}(k)$$

$$\mathcal{V}_{g,n}^{1\mathrm{PI}}(\mathcal{V}_1,\ldots,\mathcal{V}_n)\!:=\!\mathcal{V}_1\,\widehat{\mathcal{V}}_n^{g/1\mathrm{PI}}\!:=\!\int_{\mathcal{R}_{g,n}^{1\mathrm{PI}}}\omega_{\mathbf{M}_{g,n}}^{g,n}(\mathcal{V}_1,\ldots,\mathcal{V}_n),$$

$$\langle \mathcal{V}^{g,n} \mid \otimes_i \mathcal{V}_i \rangle \!:=\! \mathcal{V}_{g,n}(\otimes_i \mathcal{V}_i).$$

$$\mathcal{V}_{g,n+1}(\mathcal{V}_0,\mathcal{V}_1,\ldots,\mathcal{V}_n)\!:=\!\langle \mathcal{V}_0|c_0^-|\ell_{g,n}(\mathcal{V}_1,\ldots,\mathcal{V}_n)\rangle.$$

$$\ell_{g,n}(\mathcal{V}_1,\ldots,\mathcal{V}_n)\!:=\![\mathcal{V}_1,\ldots,\mathcal{V}_n]_g$$

$$\ell_{g,0}\!:=\![\cdot]_g\in\mathcal{H}.$$

$$N_{\mathrm{gh}}\left(\ell_{g,n}(\mathcal{V}_1,\ldots,\mathcal{V}_n)\right)=3-2n+\sum_{i=1}^nN_{\mathrm{gh}}(\mathcal{V}_i)=3+\sum_{i=1}^n\left(N_{\mathrm{gh}}(\mathcal{V}_i)-2\right),$$

$$|\ell_{g,n}(\mathcal{V}_1,\ldots,\mathcal{V}_n)|=1+\sum_{i=1}^n|\mathcal{V}_i|\bmod 2,$$

$$0=\sum_{\substack{g_1,g_2\geq 0\\g_1+g_2=g}}\sum_{\substack{n_1,n_2\geq 0\\n_1+n_2=n}}\frac{n!}{n_1!n_2!}\mathcal{V}_{g_1,n_1+1}\left(\Psi^{n_1},\ell_{g_2,n_2}(\Psi^{n_2})\right)+(-1)^{|\phi_s|}\mathcal{V}_{g-1,n+2}(\phi_s,b_0^-\phi_s^c,\Psi^n)$$

$$\mathcal{V}_{g,n}(\mathcal{V}_1,\ldots,\mathcal{V}_n)=\frac{1}{N}\sum_{a=1}^N\int_{\mathcal{R}_{g,n}^{(a)}}\omega_{\mathbf{M}_{g,n}}^{g,n}(\mathcal{V}_1,\ldots,\mathcal{V}_n)$$

$$\mathcal{V}_{0,3}(\mathcal{V}_1,\mathcal{V}_2,\mathcal{V}_3)=\mathcal{V}_{0,3}(\mathcal{V}_3,\mathcal{V}_1,\mathcal{V}_2)+\cdots$$

$$\mathcal{V}_{0,3}(\mathcal{V}_1,\mathcal{V}_2,\mathcal{V}_3)=\omega_0^{0,3}(\mathcal{V}_1,\mathcal{V}_2,\mathcal{V}_3)\big|_{\mathcal{S}_{0,3}}=\langle f_1\circ\mathcal{V}_1(0)f_2\circ\mathcal{V}_2(0)f_3\circ\mathcal{V}_3(0)\rangle,$$

$$\mathcal{V}_{0,3}(\mathcal{V}_3,\mathcal{V}_1,\mathcal{V}_2)=\langle f_1\circ\mathcal{V}_3(0)f_2\circ\mathcal{V}_1(0)f_3\circ\mathcal{V}_2(0)\rangle\neq\mathcal{V}_{0,3}(\mathcal{V}_1,\mathcal{V}_2,\mathcal{V}_3),$$

$$g(z_1)=z_2,g(z_2)=z_3,g(z_3)=z_1$$

$$\mathcal{V}_{0,3}(\mathcal{V}_3,\mathcal{V}_1,\mathcal{V}_2)=\langle g\circ f_1\circ\mathcal{V}_3(0)g\circ f_2\circ\mathcal{V}_1(0)g\circ f_3\circ\mathcal{V}_2(0)\rangle.$$

$$g\circ f_1=f_2,g\circ f_2=f_3,g\circ f_3=f_1$$

$$\mathcal{V}_{0,3}(\mathcal{V}_1,\mathcal{V}_2,\mathcal{V}_3)=\frac{1}{6}\sum_{a=1}^6\omega_0^{0,3}(\mathcal{V}_1,\mathcal{V}_2,\mathcal{V}_3)\bigg|_{\mathcal{S}_{0,3}^{(a)}}$$

Espacio de Hilbert en CFT.

$$|\Psi\rangle = \sum_r \psi_r |\phi_r\rangle.$$

$$|\Psi\rangle = \sum_r (\psi_{\downarrow\downarrow,r} |\phi_{\downarrow\downarrow,r}\rangle + \psi_{\downarrow\uparrow,r} |\phi_{\downarrow\uparrow,r}\rangle + \psi_{\uparrow\downarrow,r} |\phi_{\uparrow\downarrow,r}\rangle + \psi_{\uparrow\uparrow,r} |\phi_{\uparrow\uparrow,r}\rangle),$$

$$\begin{aligned} b_0 |\phi_{\downarrow\downarrow,r}\rangle &= \bar{b}_0 |\phi_{\downarrow\downarrow,r}\rangle = 0, & b_0 |\phi_{\downarrow\uparrow,r}\rangle &= \bar{c}_0 |\phi_{\downarrow\uparrow,r}\rangle = 0, \\ c_0 |\phi_{\uparrow\downarrow,r}\rangle &= \bar{b}_0 |\phi_{\uparrow\downarrow,r}\rangle = 0, & c_0 |\phi_{\uparrow\uparrow,r}\rangle &= \bar{c}_0 |\phi_{\uparrow\uparrow,r}\rangle = 0. \end{aligned}$$

$$\langle \phi_r^c | \phi_s \rangle = \delta_{rs}$$

$$\begin{aligned} \langle \phi_{\downarrow\downarrow,r}^c | c_0 &= \langle \phi_{\downarrow\downarrow,r}^c | \bar{c}_0 = 0, & \langle \phi_{\downarrow\uparrow,r}^c | c_0 &= \langle \phi_{\downarrow\uparrow,r}^c | \bar{b}_0 = 0 \\ \langle \phi_{\uparrow\downarrow,r}^c | b_0 &= \langle \phi_{\uparrow\downarrow,r}^c | \bar{c}_0 = 0, & \langle \phi_{\uparrow\uparrow,r}^c | b_0 &= \langle \phi_{\uparrow\uparrow,r}^c | \bar{b}_0 = 0, \\ \langle \phi_{x,r}^c | \phi_{y,s} \rangle &= \delta_{xy} \delta_{rs}, \end{aligned}$$

$$G(\psi_r) = 2 - n_r$$

$$n_r = N_{\text{gh}}(\phi_r), n_r^c = N_{\text{gh}}(\phi_r^c) = 6 - n_r$$

Gauge fixed y transformaciones de gauge bajo propagador.

$$\Delta = b_0^+ b_0^- \frac{1}{L_0^+} \delta_{L_0^-,0}, \Delta_{rs} = \langle \phi_r^c | b_0^+ b_0^- \frac{1}{L_0^+} \delta_{L_0^-,0} | \phi_s^c \rangle.$$

$$S_{0,2} = \frac{1}{2} \langle \Psi | K | \Psi \rangle = \frac{1}{2} \psi_r K_{rs} \psi_s$$

$$K = c_0^- c_0^+ L_0^+ \delta_{L_0^-,0} K_{rs} = \langle \phi_r | c_0^- c_0^+ L_0^+ \delta_{L_0^-,0} | \phi_s \rangle$$

$$K = \frac{1}{2} c_0 \bar{c}_0 L_0^+ \delta_{L_0^-,0}$$

$$\ker K|_{\mathcal{H}} \neq \emptyset$$

$$K_{rs} = \frac{1}{2} \begin{pmatrix} \langle \phi_{\downarrow\downarrow,r} | \\ \langle \phi_{\downarrow\uparrow,r} | \\ \langle \phi_{\uparrow\downarrow,r} | \\ \langle \phi_{\uparrow\uparrow,r} | \end{pmatrix}^t \begin{pmatrix} c_0 \bar{c}_0 L_0^+ & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix} \begin{pmatrix} | \phi_{\downarrow\downarrow,s} \rangle \\ | \phi_{\downarrow\uparrow,s} \rangle \\ | \phi_{\uparrow\downarrow,s} \rangle \\ | \phi_{\uparrow\uparrow,s} \rangle \end{pmatrix}.$$

$$L_0^- |\Psi\rangle = 0, b_0^- |\Psi\rangle = 0, b_0^+ |\Psi\rangle = 0$$

$$|\Psi\rangle = \sum_r \psi_{\downarrow\downarrow,r} |\phi_{\downarrow\downarrow,r}\rangle.$$



$$\begin{aligned}
K_{rs}\Delta_{st} &= \langle \phi_r | c_0^- c_0^+ L_0^+ \delta_{L_0^-,0} | \phi_s \rangle \langle \phi_s^c | b_0^+ b_0^- \frac{1}{L_0^+} \delta_{L_0^-,0} | \phi_t^c \rangle \\
&= \langle \phi_r | c_0^- c_0^+ L_0^+ \delta_{L_0^-,0} b_0^+ b_0^- \frac{1}{L_0^+} \delta_{L_0^-,0} | \phi_t^c \rangle \\
&= \langle \phi_r | \{c_0^-, b_0^-\} \{c_0^+, b_0^+\} | \phi_t^c \rangle \\
&= \langle \phi_r | \phi_t^c \rangle = \delta_{rt}.
\end{aligned}$$

$$1 = \sum_r |\phi_r\rangle \langle \phi_r^c| = \sum_r |\phi_{\downarrow\downarrow,r}\rangle \langle \phi_{\downarrow\downarrow,r}^c|.$$

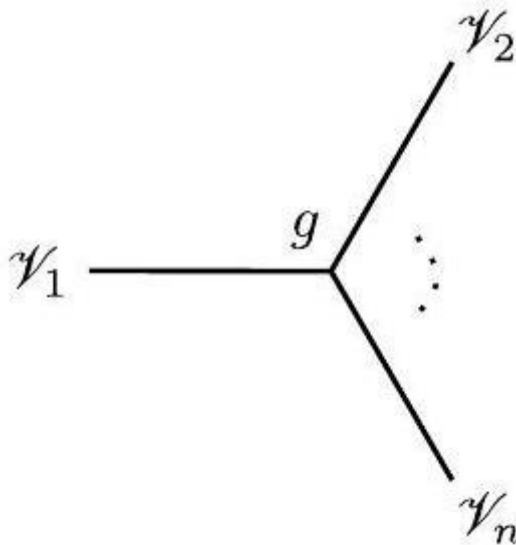
$$S_{0,2} = \frac{1}{2} \mathcal{V}_{0,2}(\Psi^2) = \frac{1}{2} \langle \Psi | c_0^- c_0^+ L_0^+ \delta_{L_0^-,0} | \Psi \rangle.$$

$$S=\int\,\mathrm{d}^Dx\left(\frac{1}{2}\phi(x)(-\partial^2+m^2)\phi(x)+\frac{\lambda}{n!}\phi(x)^n\right)$$

$$\phi_k(x) = \mathrm{e}^{\mathrm{i}k\cdot x}$$

$$\begin{aligned}
V_n(k_1,\ldots,k_n) &= \frac{\lambda}{n!} \int \,\mathrm{d}^Dx n! \prod_{i=1}^n \phi_{k_i}(x) = \lambda \int \,\mathrm{d}^Dx \mathrm{e}^{\mathrm{i}(k_1+\cdots+k_n)x} \\
&= \lambda (2\pi)^D \delta^{(D)}(k_1+\cdots+k_n)
\end{aligned}$$

$$\mathcal{V}_{g,n}(\mathcal{V}_1,\ldots,\mathcal{V}_n)=\int_{\mathcal{R}_{g,n}}\omega_{\mathbf{M}_{g,n}}^{g,n}(\mathcal{V}_1,\ldots,\mathcal{V}_n)=$$



$$S_{g,n}=\hbar^g\frac{g_s^{2g-2+n}}{n!}\mathcal{V}_{g,n}(\Psi^n)$$

$$\mathcal{V}_{g,m}(\mathcal{V}_1,\ldots,\mathcal{V}_{m-1},\phi_r)\langle \phi_r^c|b_0^+b_0^-\frac{1}{L_0^+}|\phi_s^c\rangle\mathcal{V}_{g',n}(\mathcal{W}_1,\ldots,\mathcal{W}_{n-1},\phi_s)$$

$$= \mathcal{V}_{g,m}(\mathcal{V}_1,\ldots,\mathcal{V}_{m-1},\phi_{\downarrow\downarrow,r})\langle \phi_{\downarrow\downarrow,r}^c|b_0^+b_0^-\frac{1}{L_0^+}|\phi_{\downarrow\downarrow,s}^c\rangle\mathcal{V}_{g',n}(\mathcal{W}_1,\ldots,\mathcal{W}_{n-1},\phi_{\downarrow\downarrow,s})$$

$$S = \sum_{g,n \geq 0} \hbar^g \frac{g_s^{2g-2+n}}{n!} \mathcal{V}_{g,n}(\Psi^n) := \frac{1}{2} \langle \Psi | c_0^- c_0^+ L_0^+ \delta_{L_0^-,0} | \Psi \rangle + \sum'_{g,n \geq 0} \hbar^g \frac{g_s^{2g-2+n}}{n!} \mathcal{V}_{g,n}(\Psi^n)$$

$$\mathcal{V}_{g,n}(\Psi^n) := \langle \Psi | c_0^- | \ell_{g,n-1}(\Psi^{n-1}) \rangle$$

$$S = \sum_{g,n \geq 0} \hbar^g \frac{g_s^{2g-2+n}}{n!} \langle \Psi | c_0^- | \ell_{g,n-1}(\Psi^{n-1}) \rangle$$

$$\ell_{0,1}(\Psi_{\text{cl}}) = c_0^+ L_0^+ | \Psi_{\text{cl}} \rangle$$

$$S = \sum_{\substack{g,n \geq 0 \\ \chi_{g,n} \leq 0}} \hbar^g \frac{g_s^{2g-2+n}}{n!} \mathcal{V}_{g,n}(\Psi^n)$$

$$S_{\text{cl}} = \frac{1}{2} \langle \Psi_{\text{cl}} | c_0^- c_0^+ L_0^+ | \Psi_{\text{cl}} \rangle + \sum_{n \geq 3} \frac{g_s^n}{n!} \mathcal{V}_{0,n}(\Psi_{\text{cl}}^n)$$

$$S = \sum_{g,n \geq 0} \hbar^g g_s^{2g-2} \frac{1}{n!} \mathcal{V}_{g,n}(\Psi^n) := \frac{1}{2g_s^2} \langle \Psi | c_0^- c_0^+ L_0^+ \delta_{L_0^-,0} | \Psi \rangle + \frac{1}{g_s^2} \sum'_{g,n \geq 0} \frac{(\hbar g_s^2)^g}{n!} \mathcal{V}_{g,n}(\Psi^n)$$

$$\frac{S}{\hbar} = \sum_{g,n \geq 0} (\hbar g_s^2)^{g-1} \frac{1}{n!} \mathcal{V}_{g,n}(\Psi^n)$$

$$S \rightarrow \alpha S \implies g_s^2 \rightarrow \frac{g_s^2}{\alpha}$$

$$L_0^+ | \Psi \rangle = 0$$

$$\Psi_{\text{cl}} \in \mathcal{H}^- \cap \ker L_0^-, N_{\text{gh}}(\Psi_{\text{cl}}) = 2$$

$$S_{0,2} = \frac{1}{2} \langle \Psi | c_0^- Q_B | \Psi \rangle.$$

$$S_{\text{cl}} = \frac{1}{2} \langle \Psi_{\text{cl}} | c_0^- Q_B | \Psi_{\text{cl}} \rangle + \frac{1}{g_s^2} \sum_{n \geq 3} \frac{g_s^n}{n!} \mathcal{V}_{0,n}(\Psi_{\text{cl}}^n)$$

$$\mathcal{V}_{0,2}(\Psi_{\text{cl}}^2) := \langle \Psi_{\text{cl}} | c_0^- Q_B | \Psi_{\text{cl}} \rangle$$

$$S_{\text{cl}} = \frac{1}{g_s^2} \sum_{n \geq 2} \frac{g_s^n}{n!} \mathcal{V}_{0,n}(\Psi_{\text{cl}}^n) = \frac{1}{g_s^2} \sum_{n \geq 2} \frac{g_s^n}{n!} \langle \Psi_{\text{cl}} | c_0^- | \ell_{0,n-1}(\Psi_{\text{cl}}^{n-1}) \rangle,$$

$$\ell_{0,1}(\Psi_{\text{cl}}) = Q_B | \Psi_{\text{cl}} \rangle.$$

$$\mathcal{F}_{\text{cl}}(\Psi_{\text{cl}}) := \sum_{n \geq 1} \frac{g_s^{n-1}}{n!} \ell_{0,n}(\Psi_{\text{cl}}^n) = Q_B |\Psi_{\text{cl}}\rangle + \sum_{n \geq 2} \frac{g_s^{n-1}}{n!} \ell_{0,n}(\Psi_{\text{cl}}^n) = 0$$

$$\delta S_{\text{cl}} = \frac{1}{g_s^2} \sum_{n \geq 2} \frac{g_s^n}{n!} n \{ \delta \Psi_{\text{cl}}, \Psi_{\text{cl}}^{n-1} \}_0 = \frac{1}{g_s^2} \sum_{n \geq 2} \frac{g_s^n}{(n-1)!} \langle \delta \Psi_{\text{cl}} | c_0^- \rangle \ell_{0,n-1}(\Psi_{\text{cl}}^{n-1}).$$

$$\delta_{\Lambda} S_{\text{cl}} = 0$$

$$\delta_{\Lambda} \Psi_{\text{cl}} = \sum_{n \geq 0} \frac{g_s^n}{n!} \ell_{0,n+1}(\Psi_{\text{cl}}^n, \Lambda) = Q_B |\Lambda\rangle + \sum_{n \geq 1} \frac{g_s^n}{n!} \ell_{0,n+1}(\Psi_{\text{cl}}^n, \Lambda)$$

$$[\delta_{\Lambda_2}, \delta_{\Lambda_1}] \Psi_{\text{cl}} = \delta_{\Lambda(\Lambda_1, \Lambda_2, \Psi_{\text{cl}})} |\Psi_{\text{cl}}\rangle + \sum_{n \geq 0} \frac{g_s^{n+2}}{n!} \ell_{0,n+3}(\Psi_{\text{cl}}^n, \Lambda_2, \Lambda_1, \mathcal{F}_{\text{cl}}(\Psi_{\text{cl}})),$$

$$\Lambda(\Lambda_1, \Lambda_2, \Psi_{\text{cl}}) = \sum_{n \geq 0} \frac{g_s^{n+1}}{n!} \ell_{0,n+2}(\Lambda_1, \Lambda_2, \Psi_{\text{cl}}^n) = g_s \ell_{0,2}(\Lambda_1, \Lambda_2) + \sum_{n \geq 1} \frac{g_s^{n+1}}{n!} \ell_{0,n+2}(\Lambda_1, \Lambda_2, \Psi_{\text{cl}}^n)$$

$$\forall n \geq 4: \mathcal{V}_{0,4}(\mathcal{V}_1, \dots, \mathcal{V}_n) = 0, \ell_{g,n-1}(\mathcal{V}_1, \dots, \mathcal{V}_{n-1}) = 0, (\text{cubic theory}),$$

$$\begin{aligned} \delta_{\Lambda} S_{\text{cl}} &= \sum_{n \geq 2} \frac{g_s^{n-2}}{n!} n \mathcal{V}_{0,n}(\delta \Psi_{\text{cl}}, \Psi_{\text{cl}}^{n-1}) = \sum_{m,n \geq 0} \frac{g_s^{m+n-1}}{m! n!} \mathcal{V}_{0,n+1}(\ell_{0,m+1}(\Psi_{\text{cl}}^m, \Lambda), \Psi_{\text{cl}}^n) \\ &= \sum_{m \geq 0} \sum_{n=0}^m \frac{g_s^{m-1}}{(m-n)! n!} \langle \ell_{m-n+1}(\Psi_{\text{cl}}^{m-n}, \Lambda) | c_0^- \rangle \ell_{0,n}(\Psi_{\text{cl}}^n) \\ &= \langle \ell_{0,n}(\Psi_{\text{cl}}^n) | c_0^- \rangle \ell_{0,m-n+1}(\Psi_{\text{cl}}^{m-n}, \Lambda) \\ &= \mathcal{V}_{0,m-n+2}(\ell_{0,n}(\Psi_{\text{cl}}^n), \Psi_{\text{cl}}^{m-n}, \Lambda) \\ &= -\mathcal{V}_{0,m-n+2}(\Lambda, \ell_{0,n}(\Psi_{\text{cl}}^n), \Psi_{\text{cl}}^{m-n}) \\ &= \langle \Lambda | c_0^- \rangle \ell_{0,m-n+1}(\ell_{0,n}(\Psi_{\text{cl}}^n), \Psi_{\text{cl}}^{m-n}). \end{aligned}$$

$$0 = \sum_{n=0}^m \frac{m!}{(m-n)! n!} \ell_{0,m-n+1}(\ell_{0,n}(\Psi_{\text{cl}}^n), \Psi_{\text{cl}}^{m-n})$$

Formalismo BV.

$$\begin{aligned} S &= \frac{1}{g_s^2} \sum_{g \geq 0} \hbar^g g_s^{2g} \sum_{n \geq 0} \frac{g_s^n}{n!} \mathcal{V}_{g,n}(\Psi^n) \\ &= \frac{1}{2} \langle \Psi | c_0^- Q_B | \Psi \rangle + \sum_{g,n \geq 0} \frac{\hbar^g g_s^{2g-2+n}}{n!} \mathcal{V}_{g,n}(\Psi^n) \\ &= \frac{1}{g_s^2} \sum_{g,n \geq 0} \frac{\hbar^g g_s^{2g-2+n}}{n!} \langle \Psi | c_0^- \rangle \ell_{g,n-1}(\Psi^{n-1}) \end{aligned}$$

$$\Psi \in \mathcal{H}^- \cap \ker L_0^-.$$

$$(S, S) - 2\hbar \Delta S = 0$$



$$|\Psi\rangle = \sum_r \psi_r |\phi_r\rangle$$

$$\Psi = \Psi_+ + \Psi_-$$

$$\Psi_- = \sum_r \sum_{n_r \leq 2} |\phi_r\rangle \psi_r^r, \Psi_+ = \sum_r \sum_{n_r^c > 2} b_0^- |\phi_r^c\rangle \psi_r^*$$

$$G(\psi^r) \geq 0, G(\psi_r^*) < 0$$

$$G(\psi_r^*) = -1 - G(\psi^r)$$

$$\begin{aligned} G(\psi_r^*) &= 2 - N_{\text{gh}}(b_0^- \phi_r^c) = 2 + 1 - n_r^c \\ &= 3 - (6 - n_r) = -3 + (2 - G(\psi^r)) = -1 - G(\psi^r) \end{aligned}$$

$$\frac{\partial_R S}{\partial \psi^r} \frac{\partial_L S}{\partial \psi_r^*} + \hbar \frac{\partial_R \partial_L S}{\partial \psi^r \partial \psi_r^*} = 0$$

$$\sum_{\substack{g_1, g_2 \geq 0 \\ g_1 + g_2 = g}} \sum_{\substack{n_1, n_2 \geq 0 \\ n_1 + n_2 = n}} \frac{\partial_R S_{g_1, n_1}}{\partial \psi^r} \frac{\partial_L S_{g_2, n_2}}{\partial \psi_r^*} + \hbar \frac{\partial_R \partial_L S_{g-1, n}}{\partial \psi^r \partial \psi_r^*} = 0$$

$$\begin{aligned} \mathcal{V}_n^{\text{1PI}}(\mathcal{V}_1, \dots, \mathcal{V}_n) &:= \mathcal{V}_1 \rightarrow \sum_{\mathcal{V}_n}^{\mathcal{V}_2} := \sum_{g \geq 0} (\hbar g_s^2)^g \mathcal{V}_{g,n}^{\text{1PI}}(\mathcal{V}_1, \dots, \mathcal{V}_n), \\ \mathcal{V}_{g,n}^{\text{1PI}}(\mathcal{V}_1, \dots, \mathcal{V}_n) &:= \int_{\mathcal{R}_{g,n}^{\text{1PI}}} \omega_{\mathbf{M}_{g,n}}^{g,n}(\mathcal{V}_1, \dots, \mathcal{V}_n), \end{aligned}$$

$$S_{\text{1PI}} = \frac{1}{g_s^2} \sum_{n \geq 0} \frac{g_s^n}{n!} \mathcal{V}_n^{\text{1PI}}(\Psi^n) := \frac{1}{2} \langle \Psi | c_0^- c_0^+ L_0^+ | \Psi \rangle + \frac{1}{g_s^2} \sum_{n \geq 0} \frac{g_s^n}{n!} \mathcal{V}_n^{\text{1PI}}(\Psi^n)$$

$$S_{\text{1PI}} = \frac{1}{2} \langle \Psi | c_0^- Q_B | \Psi \rangle + \frac{1}{g_s^2} \sum_{n \geq 0} \frac{g_s^n}{n!} \mathcal{V}_n^{\text{1PI}}(\Psi^n)$$

Independencia background.

$$\delta S_{\text{cft}} = \frac{\lambda}{2\pi} \int \mathrm{d}^2z \varphi(z, \bar{z})$$

$$S_1[\Psi_1] = \frac{1}{g_s^2} \left(\frac{1}{2} \langle \Psi_1 | c_0^- Q_B | \Psi_1 \rangle + \sum_{n \geq 0} \frac{1}{n!} \mathcal{V}_n^{\text{1PI}}(\Psi_1^n) \right)$$

$$\mathcal{F}_1(\Psi_1) = Q_B|\Psi_1\rangle + \sum_n \frac{1}{n!} \ell_n(\Psi_1^n) = 0$$



Deformación de CFT.

$$S_{\text{cft},2}[\psi_1] = S_{\text{cft},1}[\psi_1] + \frac{\lambda}{2\pi} \int d^2z \varphi(z, \bar{z})$$

$$\begin{aligned} \left\langle \prod_i \mathcal{O}_i(z_i, \bar{z}_i) \right\rangle_2 &= \left\langle \exp \left(-\frac{\lambda}{2\pi} \int d^2z \varphi(z, \bar{z}) \right) \prod_i \mathcal{O}_i(z_i, \bar{z}_i) \right\rangle_1 \\ &\approx \left\langle \prod_i \mathcal{O}_i(z_i, \bar{z}_i) \right\rangle_1 - \frac{\lambda}{2\pi} \int_{\Sigma} d^2z \left\langle \varphi(z, \bar{z}) \prod_i \mathcal{O}_i(z_i, \bar{z}_i) \right\rangle_1 \end{aligned}$$

$$\delta \left\langle \prod_i \mathcal{O}_i(z_i, \bar{z}_i) \right\rangle_1 = -\frac{\lambda}{2\pi} \int_{\Sigma \cup_i D_i} d^2z \left\langle \varphi(z, \bar{z}) \prod_i \mathcal{O}_i(z_i, \bar{z}_i) \right\rangle_1.$$

$$\delta L_n = \lambda \oint_{|z|=1} \frac{d\bar{z}}{2\pi i} z^{n+1} \varphi(z, \bar{z}), \delta \bar{L}_n = \lambda \oint_{|z|=1} \frac{dz}{2\pi i} \bar{z}^{n+1} \varphi(z, \bar{z}).$$

$$\delta Q_B = \lambda \oint_{|z|=1} \frac{d\bar{z}}{2\pi i} c(z) \varphi(z, \bar{z}) + \lambda \oint_{|z|=1} \frac{d\bar{z}}{2\pi i} \bar{c}(\bar{z}) \varphi(z, \bar{z}).$$

$$\{Q_B, \delta Q_B\} = O(\lambda^2)$$

$$\delta L_0^- |\mathcal{O}\rangle = \lambda \oint_{|z|=1} \frac{d\bar{z}}{2\pi i} z \sum_{p,q} z^{p-1} \bar{z}^{q-1} |\mathcal{O}_{p,q}\rangle - \lambda \oint_{|z|=1} \frac{d\bar{z}}{2\pi i} \sum_{p,q} z^{p-1} \bar{z}^{q-1} |\mathcal{O}_{p,q}\rangle,$$

$$\varphi(z, \bar{z}) \mathcal{O}(0,0) = \sum_{p,q} z^{p-1} \bar{z}^{q-1} \mathcal{O}_{p,q}(0,0)$$

$$S_2[\Psi_1] = S_1[\Psi_1] + \delta S_1[\Psi_1]$$

$$\delta S_1[\Psi_1] = \frac{1}{g_s^2} \left(\frac{1}{2} \langle \Psi_1 | c_0^- \delta Q_B | \Psi_1 \rangle + \sum_{n \geq 0} \frac{1}{n!} \delta \mathcal{V}_n(\Psi_1^n) \right)$$

$$\mathcal{F}_2(\Psi_1) = \mathcal{F}_1(\Psi_1) + \lambda \delta \mathcal{F}_1(\Psi_1) = 0$$

$$\lambda \delta \mathcal{F}_1(\Psi_1) = \delta Q_B | \Psi_1 \rangle + \sum_n \frac{1}{n!} \delta \ell_n(\Psi_1^n)$$

$$|\Psi_1\rangle = \lambda |\Psi_0\rangle, |\Psi_0\rangle = c_1 \bar{c}_1(0) |\varphi\rangle$$

$$|\Psi_1\rangle = \lambda |\Psi_0\rangle + |\Psi'\rangle,$$

$$S_1[\Psi_1] = S_1[\Psi_0] + S'[\Psi']$$

$$S'[\Psi'] = \frac{1}{g_s^2} \left(\frac{1}{2} \langle \Psi' | c_0^- Q_B | \Psi' \rangle + \sum_n \frac{1}{n!} (\mathcal{V}_n(\Psi'^n) + \lambda \mathcal{V}_{n+1}(\Psi_0, \Psi'^n)) \right)$$

$$\mathcal{F}'(\Psi') := \mathcal{F}_1(\Psi') + \lambda \delta \mathcal{F}'(\Psi') = 0$$



$$\delta\mathcal{F}'(\Psi') = \sum_n \frac{1}{n!} \ell_{n+1}(\Psi_0, \Psi'^n)$$

$$\begin{aligned} \mathcal{F}_1(\Psi_1) + \lambda \delta\mathcal{F}_1(\Psi_1) &= (1 + \lambda\mathcal{M}(\Psi'))(\mathcal{F}_1(\Psi') + \lambda\delta\mathcal{F}'(\Psi')) \\ |\Psi_1\rangle &= |\Psi'\rangle + \lambda|\delta\Psi'\rangle \end{aligned}$$

$$\left. \frac{d}{d\lambda} \mathcal{F}_1(\Psi' + \lambda\delta\Psi') \right|_{\lambda=0} + \delta\mathcal{F}_1(\Psi_1) - \delta\mathcal{F}'(\Psi') = \mathcal{M}(\Psi')\mathcal{F}_1(\Psi')$$

$$\begin{aligned} &\lambda Q_B |\delta\Psi'\rangle + \lambda \sum_n \frac{1}{n!} \ell_{n+1}(\delta\Psi', \Psi'^n) + \delta Q_B |\Psi'\rangle \\ &+ \sum_n \frac{1}{n!} \delta \ell_n(\Psi'^n) - \lambda \sum_n \frac{1}{n!} \ell_{n+1}(\Psi_0, \Psi'^n) = 0 \end{aligned}$$

$$\Delta := \lambda \langle A | c_0^- Q_B | \delta\Psi' \rangle + \lambda \sum_n \frac{1}{n!} \mathcal{V}_{n+2}(A, \delta\Psi', \Psi'^n) + \langle A | c_0^- \delta Q_B | \Psi' \rangle + \sum_n \frac{1}{n!} \delta \mathcal{V}_{n+1}(A, \Psi'^n)$$

$$- \lambda \sum_n \frac{1}{n!} \mathcal{V}_{n+2}(A, \Psi_0, \Psi'^n)$$

$$\langle A | c_0^- \delta Q_B | B \rangle = \lambda \mathcal{V}'_{0,3}(\Psi_0, B, A), \delta \mathcal{V}_n(\Psi'^n) = \lambda \mathcal{V}'_{n+1}(\Psi_0, \Psi'^n),$$

$$\langle A | c_0^- | \delta\Psi' \rangle = \sum_n \frac{1}{n!} \mathcal{B}_{n+2}(\Psi_0, \Psi'^n, A)$$

$$\mathcal{V}'_n = \sum_{g \geq 0} \mathcal{V}'_{g,n}, \mathcal{B}_n = \sum_{g \geq 0} \mathcal{B}_{g,n}$$

$$\begin{aligned} \Delta &= - \sum_n \frac{1}{n!} \mathcal{B}_{n+2} \left(\Psi_0, \Psi'^n, Q_B A \right) + \sum_{m,n} \frac{1}{m! n!} \mathcal{B}_{n+2}(\Psi_0, \Psi'^m, \ell_{n+1}(A, \Psi'^n)) \\ &+ \sum_n \frac{1}{n!} \mathcal{V}'_{n+2}(A, \Psi_0, \Psi'^n) - \sum_n \frac{1}{n!} \mathcal{V}_{n+2}(A, \Psi_0, \Psi'^n) \end{aligned}$$

$$\begin{aligned} \mathcal{B}_{n+2}(\Psi_0, \Psi'^n, Q_B A) &= \partial \mathcal{B}_{n+2}(\Psi_0, \Psi'^n, A) + n \mathcal{B}_{n+2}(\Psi_0, \Psi'^{n-1}, Q_B \Psi', A) \\ &= \partial \mathcal{B}_{n+2}(\Psi_0, \Psi'^n, A) - \sum_m \frac{n}{m!} \mathcal{B}_{n+2}(\Psi_0, \Psi'^{n-1}, \ell_m(\Psi'^m), A) \end{aligned}$$



$$\begin{aligned}
\Delta &= \sum_n \frac{1}{n!} \partial \mathcal{B}_{n+2}(\Psi_0, \Psi'^n, A) - \sum_{m,n} \frac{1}{m! n!} \mathcal{B}_{n+3}(\Psi_0, \Psi'^n, \ell_m(\Psi'^m), A) \\
&\quad + \sum_{m,n} \frac{1}{m! n!} \mathcal{B}_{n+2}(\Psi_0, \Psi'^m, \ell_{n+1}(A, \Psi'^n)) + \sum_n \frac{1}{n!} \mathcal{V}'_{n+2}(A, \Psi_0, \Psi'^n) \\
&\quad - \sum_n \frac{1}{n!} \mathcal{V}_{n+2}(A, \Psi_0, \Psi'^n) \\
&\quad \partial \mathcal{B}_{n+2}(\Psi_0, \Psi'^n, A) \\
&= -\mathcal{V}'_{n+2}(A, \Psi_0, \Psi'^n) + \mathcal{V}_{n+2}(A, \Psi_0, \Psi'^n) \\
&\quad + \sum_{\substack{m_1, m_2 \\ m_1+m_2=n}} \frac{n!}{m_1! m_2!} \mathcal{B}_{m_1+3}(\Psi_0, \Psi'^{m_1}, \ell_{m_2}(\Psi'^{m_2}), A) \\
&\quad - \sum_{\substack{m_1, m_2 \\ m_1+m_2=n}} \frac{n!}{m_1! m_2!} \mathcal{B}_{m_1+2}(\Psi_0, \Psi'^{m_1}, \ell_{m_2+1}(A, \Psi'^{m_2}))
\end{aligned}$$

Membranas, Simetrías y espacios planckianos en dimensión $\mathbb{R}^4$. Singularidad.

$$\begin{aligned}
\partial \mathcal{B}_{0,3} &= \mathcal{V}_{0,3} - \mathcal{V}'_{0,3} \\
T(z)T(w) &\sim \frac{c/2}{(z-w)^4} + \frac{2T(w)}{(z-w)^2} + \frac{\partial T(w)}{z-w} \\
G(z)G(w) &\sim \frac{2c/3}{(z-w)^3} + \frac{2T(w)}{(z-w)} \\
T(z)G(w) &\sim \frac{3}{2} \frac{G(w)}{(z-w)^2} + \frac{\partial G(w)}{(z-w)} \\
h(\beta) &= \left(\frac{3}{2}, 0\right), h(\gamma) = \left(-\frac{1}{2}, 0\right) \\
\gamma(z)\beta(w) &\sim \frac{1}{z-w}, \beta(z)\gamma(w) \sim -\frac{1}{z-w} \\
T^{\text{gh}} &= -2b\partial c + c\partial b, T^{\beta\gamma} = \frac{3}{2}\beta\partial\gamma + \frac{1}{2}\gamma\partial\beta \\
N_{\text{gh}}(b) &= N_{\text{gh}}(\beta) = -1, N_{\text{gh}}(c) = N_{\text{gh}}(\gamma) = 1 \\
\gamma &= \eta e^\phi, \beta = \partial \xi e^{-\phi} \\
\delta(\gamma) &= e^{-\phi}, \delta(\beta) = e^\phi \\
\int dc_0 &= 0 \Rightarrow \int dc_0 c_0 = 1
\end{aligned}$$



$$\int \mathrm{d}\gamma_0 = \infty \Rightarrow \int \mathrm{d}\gamma_0 \delta(\gamma_0) = 1$$

$$T^{\beta\gamma}=T^{\eta\xi}+T^{\phi},$$

$$T^{\eta\xi}=-\eta\partial\xi, T^{\phi}=-\frac{1}{2}(\partial\phi)^2-\partial^2\phi.$$

$$\xi(z)\eta(w)\sim\frac{1}{z-w},\mathrm{e}^{q_1\phi(z)}\mathrm{e}^{q_2\phi(w)}\sim\frac{\mathrm{e}^{(q_1+q_2)\phi(w)}}{(z-w)^{q_1q_2}},\partial\phi(z)\partial\phi(w)\sim-\frac{1}{(z-w)^2}.$$

$$N_{\mathrm{gh}}(\eta)=1, N_{\mathrm{gh}}(\xi)=-1, N_{\mathrm{gh}}(\phi)=0.$$

$$N_{\mathrm{pic}}\left(\mathrm{e}^{q\phi}\right)=q, N_{\mathrm{pic}}\left(\xi\right)=1, N_{\mathrm{pic}}\left(\eta\right)=-1.$$

$$N_{\mathrm{pic}}=2(g-1)=-\chi_g.$$

$$h(\mathrm{e}^{q\phi})=-\frac{q}{2}(q+2)$$

$$h(\mathrm{e}^{\phi})=\frac{3}{2}, h(\mathrm{e}^{-\phi})=\frac{1}{2}.$$

$$j_B=c\big(T^{\mathfrak{m}}+T^{\beta\gamma}\big)+\gamma G+bc\partial c-\frac{1}{4}\gamma^2b\\ \bar{j}_B=\bar{c}\bar{T}^{\mathfrak{m}}+\bar{b}\bar{c}\bar{\partial}\bar{c}$$

$$\mathcal{X}(z)=\{Q_B,\xi(z)\}=c\partial\xi+\mathrm{e}^{\phi}G-\frac{1}{4}\partial\eta\mathrm{e}^{2\phi}b-\frac{1}{4}\partial(\eta\mathrm{e}^{2\phi}b),$$

$$\mathcal{X}_0=\frac{1}{2\pi\mathrm{i}}\oint-\frac{\mathrm{d}z}{z}\mathcal{X}(z)$$

$$\mathcal{H}_{\mathrm{small}}=\{|\psi\rangle|\eta_0|\psi\rangle=0\}.$$

$$\mathcal{H}_{\mathrm{small}}\,=\,\mathcal{H}_{\mathrm{large}}\,\cap\,\mathrm{ker}\eta_0$$

$$\langle k|c_{-1}\bar{c}_{-1}c_0\bar{c}_0c_1\bar{c}_1\mathrm{e}^{-2\phi(z)}|k'\rangle=(2\pi)^D\delta^{(D)}(k+k')$$

$$\langle k|c_{-1}\bar{c}_{-1}c_0\bar{c}_0c_1\bar{c}_1\mathrm{e}^{-2\phi(z)}\mathrm{e}^{-\bar{\phi}(\bar{w})}|k'\rangle=-(2\pi)^D\delta^{(D)}(k+k')$$

$$\mathcal{H}_T=\mathcal{H}_{\mathrm{NS}}\oplus\mathcal{H}_{\mathrm{R}}$$

$$\widehat{\mathcal{H}}_T=\mathcal{H}_{-1}\oplus\mathcal{H}_{-1/2},\widetilde{\mathcal{H}}_T=\mathcal{H}_{-1}\oplus\mathcal{H}_{-3/2}$$

$$|p\rangle=\mathrm{e}^{p\phi}(0)|0\rangle.$$

$$\forall n\geq -p-\frac{1}{2}:\beta_n|p\rangle=0$$

$$\forall n\geq p+\frac{3}{2}:\gamma_n|p\rangle=0$$

$$\widehat{\mathcal{H}}_T=\mathrm{Span}\{|\phi_r\rangle\},\widetilde{\mathcal{H}}_T=\mathrm{Span}\{|\phi_r^c\rangle\}$$



$$\langle \phi_r^c \mid \phi_s \rangle = \delta_{rs}.$$

$$1 = \sum_r |\phi_r\rangle \langle \phi_r^c|$$

$$1 = \sum_r (-1)^{|\phi_r|} |\phi_r^c\rangle \langle \phi_r|$$

$$\mathcal{G} = \begin{cases} 1 & \text{NS sector ,} \\ \mathcal{X}_0 & \text{R sector .} \end{cases}$$

$$[\mathcal{G},L_0^\pm]=[\mathcal{G},b_0^\pm]=[\mathcal{G},Q_B]=0.$$

$$\mathrm{M}_{g,m,n} \mathop{:}= \dim \mathcal{M}_{g,m,n} = 6g-6+2m+2n$$

$$N_{\mathrm{gh}}=6-6g, N_{\mathrm{pic}}=2g-2$$

$$n_{\mathrm{pco}}:=2g-2+m+\frac{n}{2}$$

$$A_{g,m,n}(\mathcal{V}_i^{\mathrm{NS}},\mathcal{V}_j^{\mathrm{R}})=\int_{\mathcal{S}_{g,m,n}}\Omega_{\mathrm{M}_{g,m,n}}(\mathcal{V}_i^{\mathrm{NS}},\mathcal{V}_j^{\mathrm{R}}),$$

$$\Omega_{\mathrm{M}_{g,m,n}}=(-2\pi\mathrm{i})^{-\mathrm{M}_{g,m,n}^c}\left\langle\bigwedge_{\lambda=1}^{\mathrm{M}_{g,m,n}}\mathcal{B}_{\lambda}\mathrm{d}t_{\lambda}\prod_{A=1}^{n_{\mathrm{pco}}}\mathcal{X}(y_A)\prod_{i=1}^m\mathcal{V}_i^{\mathrm{NS}}\prod_{j=1}^n\mathcal{V}_j^{\mathrm{R}}\right\rangle_{\Sigma_{g,n}}$$

$$\begin{aligned} \mathcal{B}_{\lambda} = & \sum_{\alpha} \oint_{c_{\alpha}} \frac{\mathrm{d}\sigma_{\alpha}}{2\pi\mathrm{i}} b(\sigma_{\alpha}) \frac{\partial F_{\alpha}}{\partial t_{\lambda}}(F_{\alpha}^{-1}(\sigma_{\alpha})) + \sum_{\alpha} \oint_{c_{\alpha}} \frac{\mathrm{d}\bar{\sigma}_{\alpha}}{2\pi\mathrm{i}} \bar{b}(\bar{\sigma}_{\alpha}) \frac{\partial \bar{F}_{\alpha}}{\partial t_{\lambda}}(\bar{F}_{\alpha}^{-1}(\bar{\sigma}_{\alpha})) \\ & - \sum_A \frac{1}{\mathcal{X}(y_A)} \frac{\partial y_A}{\partial t_{\lambda}} \partial \xi(y_A) \\ & \mathcal{X}(y_A) - \partial \xi(y_A) \mathrm{d}y_A \end{aligned}$$

$$n_{\mathrm{pco}}^{(1)}+n_{\mathrm{pco}}^{(2)}=2(g_1+g_2)-2+(m_1+m_2-2)+\frac{n_1+n_2}{2}=n_{\mathrm{pco}}^{(\mathrm{NS})}$$

$$\Delta_{\mathrm{NS}}=b_0^+b_0^-\frac{1}{L_0+\bar{L}_0}\delta(L_0^-)$$

$$n_{\mathrm{pco}}^{(1)}+n_{\mathrm{pco}}^{(2)}=2(g_1+g_2)-2+(m_1+m_2)+\frac{n_1+n_2-2}{2}-1=n_{\mathrm{pco}}^{(\mathrm{R})}-1$$

$$\Delta_{\mathrm{R}}=b_0^+b_0^-\frac{\mathcal{X}_0}{L_0+\bar{L}_0}\delta(L_0^-)$$

$$\mathcal{X}_0=\frac{1}{2\pi\mathrm{i}}\oint\frac{\mathrm{d}w_n^{(1)}}{w_n^{(1)}}\mathcal{X}\left(w_n^{(1)}\right)=\frac{1}{2\pi\mathrm{i}}\oint\frac{\mathrm{d}w_n^{(2)}}{w_n^{(2)}}\mathcal{X}\left(w_n^{(2)}\right)$$



$$\Delta=b_0^+b_0^-\frac{\mathcal{G}}{L_0+\bar{L}_0}\delta(L_0^-)$$

$$\Delta_{\rm NS}\sim \frac{1}{k^2+m^2}, \Delta_{\rm R}\sim \frac{{\rm i}\,\partial+m}{k^2+m^2}$$

$$\mathcal{C}(x_i,y_j,z_q)=\left\langle \prod_{i=1}^{n+1}\xi(x_i)\prod_{j=1}^n\eta(y_j)\prod_{k=1}^m\mathrm{e}^{q_k\phi(z_k)}\right\rangle$$

$$=\frac{\prod_{j'=1}^n\vartheta_{\delta}(-y_{j'}+\sum_i x_i-\sum_j y_j+\sum_k q_kz_k)}{\prod_{i'=1}^{n+1}\vartheta_{\delta}(-x_{i'}+\sum_i x_i-\sum_j y_j+\sum_k q_kz_k)}\times\frac{\prod_{i<i'}E(x_i,x_{i'})\prod_{j<j'}E(y_j,y_{j'})}{\prod_{i,j}E(x_i,y_j)\prod_{k,\ell}E(z_k,z_{\ell})^{q_kq_{\ell}}}$$

$$\sum_kq_k=0.$$

$$E(x,y)=\frac{\vartheta_1(x-y)}{\vartheta_1'(0)}\sim_{x\rightarrow y}x-y$$

$$\vartheta_{\delta}\left(\sum_{i=2}^{n+1}x_i-\sum_{j=1}^ny_j+\sum_{k=1}^mq_kz_k\right)=0$$

$$\langle A^{(p)}|Q=0,$$

$$Q=Q_B\otimes 1^{\otimes n-1}+\cdots+1^{\otimes n-1}\otimes Q_B$$

$$\langle A^{(p)}|\eta=0$$

$$\eta=\eta_0\otimes 1^{\otimes n-1}+\cdots+1^{\otimes n-1}\otimes \eta_0.$$

$$\langle A^{(p)}|=\langle \alpha^{(p)}|Q,$$

$$\langle \alpha^{(p)}|\eta Q=0$$

$$\langle A^{(p-1)}|=\langle \alpha^{(p)}|\eta$$

$$\Psi=\Psi_{-1}+\Psi_{-1/2}$$

$$\eta_0|\Psi\rangle=0$$

$$\Delta=b_0^+b_0^-\frac{\mathcal{G}}{L_0+\bar{L}_0}\delta(L_0^-), \mathcal{G}=\begin{Bmatrix} 1 & \text{NS} \\ \mathcal{X}_0 & \text{R} \end{Bmatrix}$$

$$b_0^-|\Psi\rangle=L_0^-|\Psi\rangle=0, b_0^+|\Psi\rangle=0$$

$$X=G_0\delta(\beta_0)+b_0\delta'(\beta_0), Y=-c_0\delta'(\gamma_0)$$

$$[Q_B,X]=0$$

$$XYX=X$$



$$XY|\Psi_{-1/2}\rangle = |\Psi_{-1/2}\rangle$$

$$Bc_0^-|\Psi\rangle = |\Psi\rangle,$$

$$B = b_0^- \int_0^{2\pi} \frac{d\theta}{2\pi} e^{i\theta L_0^-} = \delta(b_0^-) \delta(L_0^-)$$

$$S_{0,2} = -\frac{1}{2} \langle \Psi_{-1} | c_0^- Q_B | \Psi_{-1} \rangle - \frac{1}{2} \langle \Psi_{-1/2} | c_0^- Y Q_B | \Psi_{-1/2} \rangle.$$

$$\delta|\Psi\rangle = Q_B|\Lambda\rangle$$

$$\Lambda = \Lambda_{-1} + \Lambda_{-1/2}$$

$$\tilde{\Psi} = \tilde{\Psi}_{-1} + \tilde{\Psi}_{-3/2}$$

$$b_0^-|\tilde{\Psi}\rangle = L_0^-|\tilde{\Psi}\rangle = 0, b_0^+|\tilde{\Psi}\rangle = 0$$

$$S_{0,2} = \frac{1}{2} \langle \tilde{\Psi} | c_0^- c_0^+ L_0^+ \mathcal{G} | \tilde{\Psi} \rangle - \langle \tilde{\Psi} | c_0^- c_0^+ L_0^+ | \Psi \rangle.$$

$$K = c_0^- c_0^+ L_0^+ \begin{pmatrix} -\mathcal{G} & 1 \\ 1 & 0 \end{pmatrix}$$

$$\Delta = b_0^- b_0^+ \frac{1}{L_0^+} \begin{pmatrix} 0 & 1 \\ 1 & \mathcal{G} \end{pmatrix}$$

$$\delta|\Psi\rangle = Q_B|\Lambda\rangle, \delta|\tilde{\Psi}\rangle = Q_B|\tilde{\Lambda}\rangle,$$

$$Q_B|\Psi\rangle = 0, Q_B|\tilde{\Psi}\rangle = 0$$

$$Q_B(|\Psi\rangle - \mathcal{G}|\tilde{\Psi}\rangle) = 0, Q_B|\tilde{\Psi}\rangle = |J(\Psi)\rangle,$$

$$Q_B|\Psi\rangle = \mathcal{G}|J(\Psi)\rangle$$

$$S_{0,2} = -\frac{1}{2} \langle \langle \Psi_0, \eta_0 Q_B \Psi_0 \rangle \rangle,$$

$$\delta|\Psi_0\rangle = Q_B|\Lambda_0\rangle + \eta_0|\Omega_1\rangle,$$

$$Q_B\eta_0|\Psi_0\rangle = 0$$

$$\xi_0|\Psi_0\rangle = 0$$

$$|\Psi_0\rangle = \xi_0|\Psi_{-1}\rangle$$

$$Q_B|\Psi_{-1}\rangle = 0$$

Momentum en dimensión $\mathbb{R}^4$ – espacio SFT.

$$|\Psi\rangle = \sum_j \int \frac{d^D k}{(2\pi)^D} \psi_\alpha(k) |\phi_\alpha(k)\rangle$$



$$S = - \int d^D k \psi_\alpha(k) K_{\alpha\beta}(k) \psi_\beta(-k) - \sum_{n \geq 0} \int d^D k_1 \cdots d^D k_n V_{\alpha_1 \dots \alpha_n}^{(n)}(k_1, \dots, k_n) \psi_{\alpha_1}(k_1) \cdots \psi_{\alpha_n}(k_n)$$

$$\alpha - \vec{k} \beta = K_{\alpha\beta}(k)^{-1} = \frac{-i M_{\alpha\beta}}{k^2 + m_\alpha^2} Q_\alpha(k),$$

$$\begin{array}{c} \alpha_2 \\ \nearrow k_2 \\ \vdots \\ \nwarrow k_n \\ \alpha_n \end{array} \quad \begin{array}{c} \alpha_1 \xrightarrow{k_1} \end{array} \quad \begin{array}{c} \vdots \\ \vdots \\ \vdots \end{array} = i V_{\alpha_1 \dots \alpha_n}^{(n)}(k_1, \dots, k_n) := i \mathcal{V}_n(\phi_{\alpha_1}(k_1), \dots, \phi_{\alpha_n}(k_n))$$

$$= i \int dt e^{-g_{ij}^{\{\alpha_k\}}(t) k_i \cdot k_j - \lambda \sum_{\alpha} m_{\alpha}^2} P_{\alpha_1, \dots, \alpha_n}(k_1, \dots, k_n; t),$$

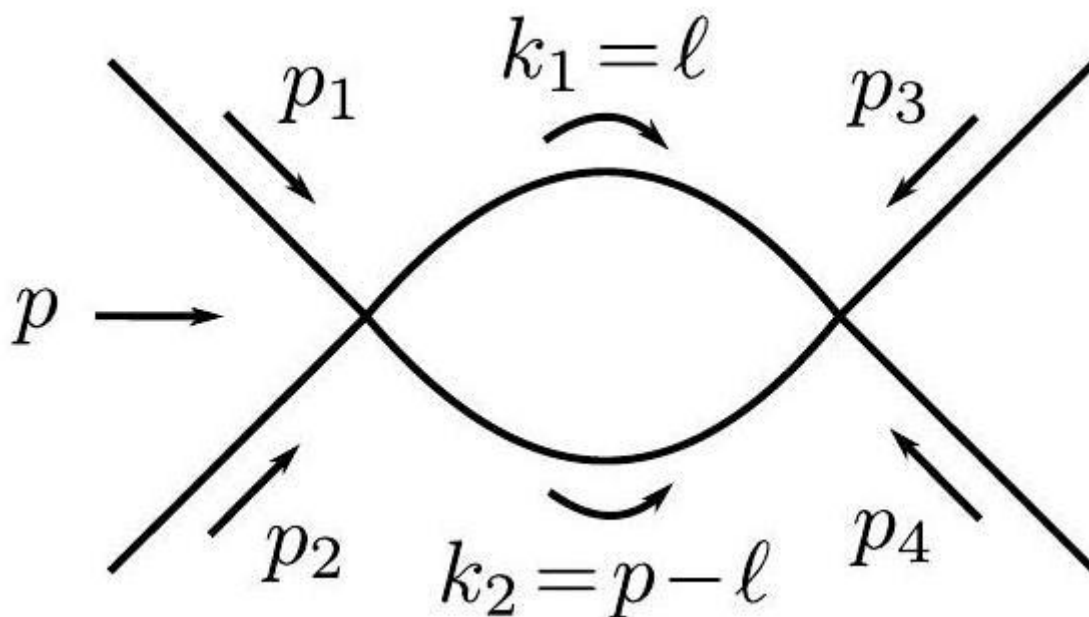
$$\lim_{k^0 \rightarrow \pm i \infty} V^{(n)} = 0, \lim_{k^0 \rightarrow \pm \infty} V^{(n)} = \infty$$

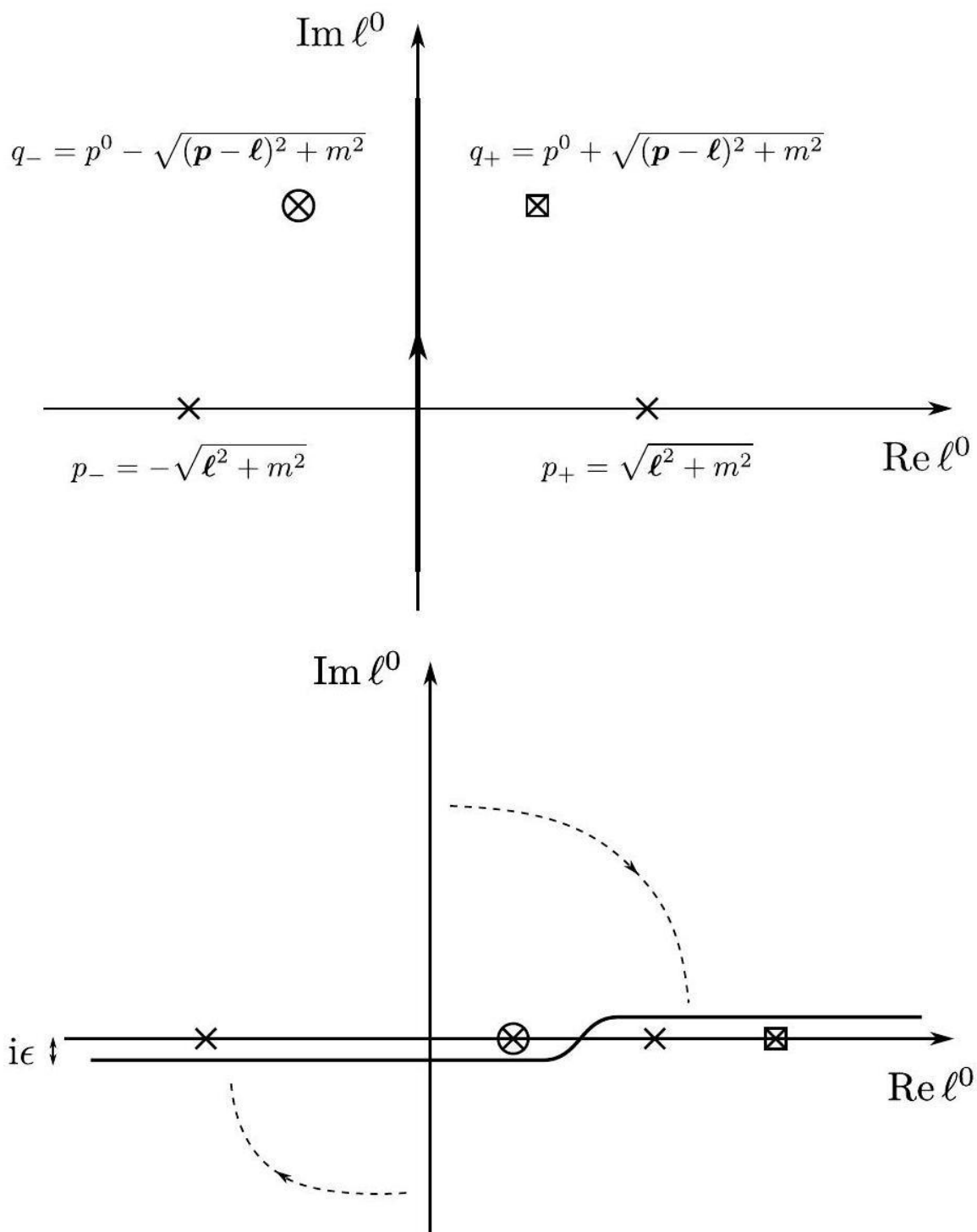
$$F_{g,n}(p_1,\ldots,p_n) \sim \int dT \prod_s d^D \ell_s e^{-G_{rs}(T) \ell_r \cdot \ell_s} - 2 H_{ri}(T) \ell_r \cdot p_i - F_{ij}(T) p_i \cdot p_j$$

$$\times \prod_a \frac{1}{k_a^2+m_a^2} \mathcal{P}(p_i,\ell_r;T)$$

Rotación de Wick.

$$p_{\pm}=\pm\sqrt{\ell^2+m^2},q_{\pm}=p^0\pm\sqrt{(\boldsymbol{p}-\boldsymbol{\ell})^2+m^2}.$$





Figuras 31, 32 y 33. Coordenadas de una partícula supermasiva en aniquilación.

Sistema de Coordenadas.

$$x^\mu = (x^0, x^i), \mu = 0, \dots, D-1 = d \quad i = 1, \dots, d$$

$$\sigma^a = (\sigma^0, \sigma^\alpha), a = 0, \dots, p-1, \alpha = 1, \dots, p.$$

$$\eta_{\mu\nu} = \text{diag}(-1, \underbrace{1, \dots, 1}_d).$$

$$\delta_{\mu\nu} = \text{diag}(\underbrace{1, \dots, 1}_D).$$

$$\epsilon_{01}=-\epsilon^{01}=1$$

$$t=-\mathrm{i}\tau.$$

$$V_M^0=-\mathrm{i}V_E^0, V_{M,0}=\mathrm{i}V_{E,0}.$$

$$x^\pm = x^0 \pm x^1.$$

$$\tau\in\mathbb{R},\sigma\in[0,L),\sigma\sim\sigma+L,$$

$$\mathcal{L}=\frac{1}{2\pi}\int_0^L\mathrm{d}\sigma=\frac{L}{2\pi}$$

$$w=\tau+\mathrm{i}\sigma, \bar{w}=\tau-\mathrm{i}\sigma$$

$$\mathrm{d} s^2=\mathrm{d} \tau^2+\mathrm{d} \sigma^2=\mathrm{d} w\,\mathrm{d} \bar{w}$$

$$w=\mathrm{i}\sigma^+, \bar{w}=\mathrm{i}\sigma^-.$$

$$z=\mathrm{e}^{2\pi w/L}, \bar{z}=\mathrm{e}^{2\pi \bar{w}/L}.$$

$$\epsilon_{z\bar{z}}=\frac{\mathrm{i}}{2}, \epsilon^{z\bar{z}}=-2\mathrm{i}$$

Sistema de Operadores.

$$[A,B]_{:}=[A,B]_{-}=AB-BA,\{A,B\}_{:}=[A,B]_{+}=AB+BA.$$

$$|A|=\begin{cases} +1 & \text{Grassmann odd} \\ 0 & \text{Grassmann even} \end{cases}$$

$$AB=(-1)^{|A||B|}BA.$$

Sistema QFT.

$$p^\mu\colon=(E,p^i)$$

$$(N_L,N_R), N=N_L+N_R$$

$$\delta\phi(x)=\phi'(x)-\phi(x).$$

$$J_a^\mu=\lambda\frac{\partial\mathcal{L}}{\partial(\partial_\mu\phi)}\frac{\delta\phi}{\delta\alpha^a},\nabla_\mu J_a^\mu=0$$

$$Q_a=\frac{1}{\lambda}\oint_{\Sigma}\mathrm{d}^{D-1}x\sqrt{\hbar}J_a^0$$

$$\delta_{\alpha^a}\phi(x)=\mathrm{i}\alpha^a[Q_a,\phi(x)]$$



$$J_a^\mu = i\lambda \frac{\partial \mathcal{L}}{\partial (\partial_\mu \phi)} \frac{\delta \phi}{\delta \alpha^a}$$

$$\delta_{\alpha^a} \phi(x) = -\alpha^a [Q_a, \phi(x)]$$

Espacio curvo y gravedad.

$$\nabla_\mu = \partial_\mu + \Gamma_\mu$$

$$\nabla_\mu A^\nu = \partial_\mu A^\nu + \Gamma_{\mu\rho}^\nu A^\rho$$

$$\Delta = g^{\mu\nu} \nabla_\mu \nabla_\nu = \frac{1}{\sqrt{g}} \nabla_\mu (\sqrt{g} g^{\mu\nu} \nabla_\nu)$$

$$T_{\mu\nu} = -\frac{2\lambda}{\sqrt{g}} \frac{\delta S}{\delta g^{\mu\nu}}$$

Análisis Complejo - Cauchy-Riemann.

$$\oint_{C_z} \frac{dw}{2\pi i} \frac{f(w)}{(w-z)^n} = \frac{f^{(n-1)}(z)}{(n-1)!},$$

$$\bar{\partial} \frac{1}{z} = 2\pi \delta^{(2)}(z)$$

QFT, espacios curvos y gravedad. Función de Green.

$$D_x G(x,y) = \frac{\delta(x-y)}{\sqrt{g}} - P(x,y),$$

$$\nabla_\mu v^\mu = \frac{1}{\sqrt{g}} \partial_\mu (\sqrt{g} v^\mu)$$

$$\delta x^\mu = \xi^\mu,$$

$$\delta g_{\mu\nu} = \mathcal{L}_\xi g_{\mu\nu} = \nabla_\mu \xi_\nu + \nabla_\nu \xi_\mu$$

$$\int_V d^D x \nabla_\mu v^\mu = \oint_{\partial V} d\Sigma_\mu v^\mu, d\Sigma_\mu := \epsilon n_\mu d^{D-1}\Sigma,$$

$$d^{D-1}\Sigma = \sqrt{g} d^{D-1}x, n_\mu = \delta_\mu^0$$

$$Q_S = \frac{1}{\lambda} \int_S d\Sigma_\mu J_a^\mu$$

Dimensiones D. Teorema de Stokes.

$$\int d^2x \partial_\mu v^\mu = \oint \epsilon_{\mu\nu} dx^\nu v^\mu = \oint (v^0 d\sigma - v^1 d\tau)$$



$$\chi_{g;b} := \frac{1}{4\pi} \int d^2\sigma \sqrt{g} R + \frac{1}{2\pi} \oint ds k \\ = 2 - 2g - b$$

$$T_{ab} = -\frac{4\pi}{\sqrt{g}} \frac{\delta S}{\delta g^{ab}}$$

Plano Complejo.

$$z=x+iy, \overline{z}=x-iy$$

$$ds^2=dx^2+dy^2=dz\,d\overline{z}, g_{z\overline{z}}=\frac{1}{2}, g_{zz}=g_{\overline{z}\overline{z}}=0$$

$$\epsilon_{z\overline{z}}=\frac{i}{2}, \epsilon^{z\overline{z}}=-2i$$

$$\partial\colon=\partial_z=\frac{1}{2}(\partial_x-i\partial_y), \bar{\partial}\colon=\partial_{\overline{z}}=\frac{1}{2}(\partial_x+i\partial_y)$$

$$V^z=V^x+iV^y, V^{\overline{z}}=V^x-iV^y$$

$$d^2x=dx\,dy=\frac{1}{2}\,d^2z, \,d^2\overline{z}=d\overline{z}\,d\overline{z}$$

$$\delta(z)=\frac{1}{2}\delta^{(2)}(x), 1=\int d^2z\delta^{(2)}(z)=\int d^2x\delta^{(2)}(x)$$

$$\int_R d^2z(\partial_z v^{\overline{z}}+\partial_{\overline{z}} v^z)=-i\oint_{\partial R}(dzv^{\overline{z}}-d\overline{z}v^z)=-2i\oint_{\partial R}(v_z\,dz-v_{\overline{z}}\,d\overline{z})$$

Propiedades generales.

$$f\circ\phi(z)=\left(\frac{df}{dz}\right)^h\phi(f(z))$$

$$\phi(z)=\sum_n\frac{\phi_n}{z^{n+h}}, \phi_n=\oint_{c_0}\frac{dz}{2\pi i}z^{n+h-1}\phi(z),$$

$$\forall n\geq -h+1: \phi_n|0\rangle=0$$

$$\forall n\leq h-1: \langle 0|\phi_n=0.$$

$$|\phi\rangle\colon=\phi(0)|0\rangle=\phi_{-h}|0\rangle.$$

$$\langle\phi^\dagger| \colon=\langle 0|I\circ\phi^\dagger(0)=\lim_{z\rightarrow\infty}z^{2h}\langle 0|\phi^\dagger(z), \langle\phi| \colon=\langle 0|I^\pm\circ\phi(0)=(\pm 1)^h\lim_{z\rightarrow\infty}z^{2h}\langle 0|\phi(z).$$

$$T(z)=\sum_n\frac{L_n}{z^{n+2}}.$$

$$T(z)\phi(w)\sim \frac{h\phi(w)}{(z-w)^2}+\frac{\partial\phi(w)}{z-w}$$

$$T(z)T(w)\sim \frac{c/2}{(z-w)^4}+\frac{2T(w)}{(z-w)^2}+\frac{\partial T(w)}{z-w}$$



Conjugaciones Hermitianas y BPZ.

Hermitiana.

$$(\lambda A_1 \cdots A_n | 0\rangle)^\dagger = \lambda^* \langle 0 | A_n^\dagger \cdots A_1^\dagger.$$

BPZ.

$$\phi_n^t = (I^\pm \circ \phi)_n = (-1)^h (\pm 1)^n \phi_{-n}$$

$$(\lambda A_1 \cdots A_n | 0\rangle)^t = \lambda \langle 0 | (A_1)^t \cdots (A_n)^t.$$

$$\langle A, B \rangle = (-1)^{|A||B|} \langle B, A \rangle.$$

$$\forall A: \langle A | B \rangle = 0 \Rightarrow |B\rangle = 0.$$

$$\langle \phi_r^c | \phi_s \rangle = \delta_{rs}$$

$$\langle \phi_r | \phi_s^c \rangle = (-1)^{|\phi_r|} \delta_{rs}.$$

Campo escalar.

$$i\partial X^\mu = \sum_n \frac{\alpha_n^\mu}{z^{n+1}}, i\bar{\partial} X^\mu = \sum_n \frac{\bar{\alpha}_n^\mu}{\bar{z}^{n+1}}$$

$$[\alpha_m^\mu, \alpha_n^\nu] = m\delta_{m+n,0}\eta^{\mu\nu}, [\bar{\alpha}_m^\mu, \bar{\alpha}_n^\nu] = m\delta_{m+n,0}\eta^{\mu\nu}, [\alpha_m^\mu, \bar{\alpha}_n^\nu] = 0.$$

$$\alpha_0^\mu = \bar{\alpha}_0^\mu = \sqrt{\frac{\alpha'}{2}} p^\mu.$$

$$[x^\mu, p^\nu] = \eta^{\mu\nu}.$$

$$V_k(z, \bar{z}) = :e^{ik \cdot X(z, \bar{z})}:, h = \bar{h} = \frac{\alpha'^2 k^2}{4}.$$

$$p^\mu |k\rangle = k^\mu |k\rangle, \forall n > 0: \alpha_n^\mu |k\rangle = 0, \bar{\alpha}_n^\mu |k\rangle = 0.$$

$$|k\rangle = V_k(0,0)|0\rangle = e^{ik \cdot x}|0\rangle.$$

$$\langle k | p^\mu = \langle k | k^\mu, \langle k | = |k\rangle^\dagger, \langle -k | = |k\rangle^t.$$

Reparametrizaciones fantasmas.

$$\epsilon = 1, \lambda = 2, c_{\text{gh}} = -26, q_{\text{gh}} = -3, a_{\text{gh}} = -1.$$

$$h(b) = 2, h(c) = -1$$



$$b(z) = \sum_{n \in \mathbb{Z}} \frac{b_n}{z^{n+2}}, c(z) = \sum_{n \in \mathbb{Z}} \frac{c_n}{z^{n-1}},$$

$$b_n = \oint \frac{dz}{2\pi i} z^{n+1} b(z), c_n = \oint \frac{dz}{2\pi i} z^{n-2} c(z).$$

$$\{b_m, c_n\} = \delta_{m+n,0}, \{b_m, b_n\} = 0, \{c_m, c_n\} = 0.$$

$$T = -2 : b \partial c : - : \partial b c : \\ L_m = \sum_n (n+m) : b_{m-n} c_n : = \sum_n (2m-n) : b_n c_{m-n} :$$

$$L_0 = - \sum_n n : b_n c_{-n} : = \sum_n n : b_{-n} c_n :$$

$$[L_m, b_n] = (m-n)b_{m+n}, [L_m, c_n] = -(2m+n)c_{m+n}$$

$$[L_0, b_0] = 0, [L_0, c_0] = 0$$

$$j = - : b c : , N_{\text{gh},L} = \oint \frac{dz}{2\pi i} j(z)$$

$$N_{\text{gh}}(c) = 1, N_{\text{gh}}(b) = -1.$$

$$j_m = - \sum_n : b_{m-n} c_n : = - \sum_n : b_n c_{m-n} : , N_{\text{gh},L} = j_0 = - \sum_n : b_{-n} c_n :$$

$$[j_m, j_n] = m \delta_{m+n,0}, [L_m, j_n] = -n j_{m+n} - \frac{3}{2} m(m+1) \delta_{m+n,0}$$

$$[N_{\text{gh}}, b(w)] = -b(w), [N_{\text{gh}}, c(w)] = c(w)$$

$$N^b = \sum_{n>0} n N_n^b, N^c = \sum_{n>0} n N_n^c, \\ N_n^b = : b_{-n} c_n : , N_n^c = : c_{-n} b_n :$$

$$[N_m^b, b_{-n}] = b_{-n} \delta_{m,n}, [N_m^c, c_{-n}] = c_{-n} \delta_{m,n}$$

$$c(z)b(w) \sim \frac{1}{z-w}, b(z)c(w) \sim \frac{1}{z-w}, b(z)b(w) \sim 0, c(z)c(w) \sim 0,$$

$$T(z)b(w) \sim \frac{2b(w)}{(z-w)^2} + \frac{\partial b(w)}{z-w}, T(z)c(w) \sim \frac{-c(w)}{(z-w)^2} + \frac{\partial c(w)}{z-w}.$$

$$j(z)b(w) \sim -\frac{b(w)}{z-w}, j(z)c(w) \sim \frac{c(w)}{z-w}. j(z)\mathcal{O}(w) \sim N_{\text{gh}}(\mathcal{O}) \frac{\mathcal{O}(w)}{z-w},$$

$$j(z)j(w) \sim \frac{1}{(z-w)^2}.$$

$$T(z)j(w) \sim \frac{-3}{(z-w)^3} + \frac{j(w)}{(z-w)^2} + \frac{\partial j(w)}{z-w}.$$

$$N^c - N^b = 3 - 3g$$



$$N_{\text{gh},L} = N_{\text{gh},L}^{\text{cyl}} + \frac{3}{2}$$

$$\forall n > -2: b_n|0\rangle = 0, \forall n > 1: c_n|0\rangle = 0$$

$$|\downarrow\rangle = c_1|0\rangle, |\uparrow\rangle = c_0c_1|0\rangle.$$

$$L_0|\downarrow\rangle = a_{\text{gh}}|\downarrow\rangle, L_0|\uparrow\rangle = a_{\text{gh}}|\uparrow\rangle, a_{\text{gh}} = -1.$$

$$L_m = \sum_n (n - (1 - \lambda)m)^* b_{m-n} c_n^* + a_{\text{gh}} \delta_{m,0},$$

$$j_m = \sum_n {}^* b_{m-n} c_n^* + \delta_{m,0}.$$

$$L_0 = \sum_n n^* b_{-n} c_{n^*}^* + a_{\text{gh}} = \hat{L}_0 - 1,$$

$$N_{\text{gh},L} = j_0 = \sum_n {}^* b_{-n} c_n^* + 1 = \hat{N}_{\text{gh},L} + \frac{1}{2}(N_0^c - N_0^b) - \frac{3}{2},$$

$$\hat{L}_0 = N^b + N^c, \hat{N}_{\text{gh},L} := \sum_{n>0} (N_n^c - N_n^b).$$

$$N_{\text{gh}}|0\rangle = 0, N_{\text{gh}}|\downarrow\rangle = |\downarrow\rangle, N_{\text{gh}}|\uparrow\rangle = 2|\uparrow\rangle.$$

$$N_{\text{gh}}^{\text{cyl}}|\downarrow\rangle = -\frac{1}{2}|\downarrow\rangle, N_{\text{gh}}^{\text{cyl}}|\uparrow\rangle = \frac{1}{2}|\uparrow\rangle.$$

$$b_n^\dagger = b_{-n}, c_n^\dagger = c_{-n}$$

$$b_n^t = (\pm 1)^n b_{-n}, c_n^t = -(\pm 1)^n c_{-n}$$

$$|\downarrow\rangle^\ddagger = \langle 0|c_{-1}, |\uparrow\rangle^\ddagger = \langle 0|c_{-1}c_0.$$

$$\langle \downarrow | := |\downarrow\rangle^t = \mp \langle 0|c_{-1}, \langle \uparrow | := |\uparrow\rangle^t = \pm \langle 0|c_0c_{-1}$$

$$\langle \downarrow | = \mp |\downarrow\rangle^\ddagger, \langle \uparrow | = \mp |\uparrow\rangle^\ddagger.$$

$$\langle \uparrow | \downarrow \rangle = \langle \downarrow | c_0 | \downarrow \rangle = \langle 0 | c_{-1} c_0 c_1 | 0 \rangle = 1,$$

$$\langle 0^c | = \langle 0 | c_{-1} c_0 c_1$$

$$b_n^\pm = b_n \pm \bar{b}_n, c_n^\pm = \frac{1}{2}(c_n \pm \bar{c}_n)$$

$$\{b_m^+, c_n^+\} = \delta_{m+n}, \{b_m^-, c_n^-\} = \delta_{m+n}$$

$$b_n^- b_n^+ = 2b_n \bar{b}_n, c_n^- c_n^+ = \frac{1}{2}c_n \bar{c}_n$$



Campo gravitónico. Operadores, progagadores, funciones fantasmas, supersimetrías, antisimetría y supermembranas y sistema de coordenadas.

$$(\alpha_n)^t = -(\pm 1)^n \alpha_{-n}, (b_n)^t = (\pm 1)^n b_{-n}, (c_n)^t = -(\pm 1)^n c_{-n}$$

$$L_n^\pm = L_n \pm \bar{L}_n, b_n^\pm = b_n \pm \bar{b}_n, c_n^\pm = \frac{1}{2}(c_n \pm \bar{c}_n)$$

$$\langle A, B \rangle = \langle A | c_0^- | B \rangle,$$

$$\langle A, B \rangle = \langle A \mid B \rangle.$$

$$|k, 0\rangle := |k\rangle \otimes |0\rangle, |k, \downarrow\rangle := |k\rangle \otimes |\downarrow\rangle.$$

$$\langle k, \downarrow | c_0 | k, \downarrow \rangle = \langle k', 0 | c_{-1} c_0 c_1 | k, 0 \rangle = (2\pi)^D \delta^{(D)}(k + k')$$

$$\langle k, \downarrow \downarrow | c_0 \bar{c}_0 | k, \downarrow \downarrow \rangle = \langle k', 0 | c_{-1} \bar{c}_{-1} c_0 \bar{c}_0 c_1 \bar{c}_1 | k, 0 \rangle = (2\pi)^D \delta^{(D)}(k + k')$$

$$|\delta\Phi|^2 = G(\Phi)(\delta\Phi, \delta\Phi),$$

$$|\Phi|^2 = G(\Phi)(\Phi, \Phi),$$

$$\mathrm{d}\Phi\sqrt{\mathrm{det}G(\Phi)}.$$

$$(\delta\Phi_1, \delta\Phi_2) = G(\Phi)(\delta\Phi_1, \delta\Phi_2), (\Phi_1, \Phi_2) = G(\Phi)(\Phi_1, \Phi_2).$$

$$|\delta\Phi_a|^2 = \int \mathrm{d}x \rho(x) \gamma_{ab}(\Phi(x)) \delta\Phi_a(x) \delta\Phi_b(x),$$

$$G_{ab}(x,y)(\Phi) = \delta(x-y) \rho(x) \gamma_{ab}(\Phi(x)).$$

$$\int \mathrm{d}\delta\Phi \mathrm{e}^{-G(\Phi)(\delta\Phi, \delta\Phi)} = \frac{1}{\sqrt{\mathrm{det}G(\Phi)}}$$

$$\int \mathrm{d}\Phi \mathrm{e}^{-G(\Phi)(\Phi, \Phi)}$$

$$\int \mathrm{d}\Phi \sqrt{\mathrm{det}G(\Phi)} F(\Phi)$$

$$S=-\frac{1}{2}\mathrm{tr}\ln \, G(\Phi)$$

$$Z=\int \mathrm{d}\Phi \mathrm{e}^{-S_{cl}(\Phi)}$$

$$G(\delta\Phi,D\delta\Phi)=G(D^\dagger\delta\Phi,\delta\Phi)$$

$$\Phi \rightarrow \Phi + \varepsilon$$



$$\int \mathrm{d}\Phi \mathrm{e}^{-\frac{1}{2}|\Phi+\varepsilon|^2} = \int \mathrm{d}\Phi \mathrm{e}^{-\frac{1}{2}|\Phi|^2}$$

$$\int \mathrm{d}\Phi \mathrm{e}^{-\frac{1}{2}|\Phi+\varepsilon|^2} = \int \mathrm{d}\tilde{\Phi} \det \frac{\delta \Phi}{\delta \tilde{\Phi}} \mathrm{e}^{-\frac{1}{2}|\tilde{\Phi}|^2} = \int \mathrm{d}\tilde{\Phi} \mathrm{e}^{-\frac{1}{2}|\tilde{\Phi}|^2}$$

$$\mathrm{d}\Phi \sqrt{\det G(\Phi)} = \mathrm{d}\tilde{\Phi} \sqrt{\det \tilde{G}(\tilde{\Phi})}, G(\Phi)(\delta\Phi, \delta\Phi) = \tilde{G}(\tilde{\Phi})(\delta\tilde{\Phi}, \delta\tilde{\Phi})$$

$$\mathrm{d}\Phi = J(\Phi, \tilde{\Phi}) \mathrm{d}\tilde{\Phi}, J(\Phi, \tilde{\Phi}) = \left| \det \frac{\partial \Phi}{\partial \tilde{\Phi}} \right| = \sqrt{\frac{\det \tilde{G}(\tilde{\Phi})}{\det G(\Phi)}}$$

$$\int \mathrm{d}\delta\Phi \mathrm{e}^{-|\delta\Phi|^2} = 1$$

$$J(\tilde{\Phi})^{-1} = \int \mathrm{d}\delta\tilde{\Phi} \mathrm{e}^{-\tilde{G}(\delta\tilde{\Phi}, \delta\tilde{\Phi})}$$

$$J=\det\frac{\partial\tilde{x}^\mu}{\partial x^\mu}=\det\frac{\partial\tilde{v}^\mu}{\partial v^\mu}$$

$$\tilde{v}^\mu=v^\nu\frac{\partial\tilde{x}^\mu}{\partial x^\nu}$$

$$D\Phi_0=0$$

$$\Phi=\Phi_0+\Phi',(\Phi_0,\Phi')=0$$

$$Z[D]=\int \mathrm{d}\Phi \sqrt{\det G} \mathrm{e}^{-\frac{1}{2}(\Phi,D\Phi)}=\left(\int \mathrm{d}\Phi_0\right)\int \mathrm{d}\Phi' \mathrm{e}^{-\frac{1}{2}(\Phi',D\Phi')}$$

$$Z[D,J]=\int \mathrm{d}\Phi \sqrt{\det G} \mathrm{e}^{-\frac{1}{2}(\Phi,D\Phi)-(J,\Phi)}$$

$$Z[D,J]=\int \mathrm{d}\Phi_0 \mathrm{e}^{-(J,\Phi_0)}\int \mathrm{d}\Phi' \mathrm{e}^{-\frac{1}{2}(\Phi',D\Phi')-(J,\Phi')}$$

$$\int \mathrm{d}x = \infty$$

$$\int \mathrm{d}\theta = 0$$

$$\int \mathrm{d}\theta \theta = \int \mathrm{d}\theta \delta(\theta) = 1$$

$$\int \mathrm{d}x \delta(x) = 1$$

$$\theta_0(x)=\theta_{0i}\psi_i(x),\ker D=\mathrm{Span}\{\psi_i\}$$



$$d\theta = \frac{1}{\sqrt{\det(\psi_i, \psi_j)}} d\theta' \prod_{i=1}^n d\theta_{0i}$$

$$d\theta \prod_{i=1}^n \theta(x_i) = \frac{\det \psi_i(x_j)}{\sqrt{\det(\psi_i, \psi_j)}} d\theta'$$

$$\begin{aligned} 1 &= \int d\theta e^{-|\theta|^2} = \int d\theta' d\theta_0 e^{-|\theta|^2 - |\theta_0|^2} \\ &= J \int d\theta' \prod_i d\theta_{0i} e^{-|\theta'|^2 - |\theta_{0i} \psi_i|^2} = J \sqrt{\det(\psi_i, \psi_j)} \end{aligned}$$

$$\begin{aligned} \int d\theta_0 \prod_{j=1}^n \theta(x_j) &= \int d\theta_0 \prod_{j=1}^n \theta_0(x_j) = \frac{1}{\sqrt{\det(\psi_i, \psi_j)}} \int d^n \theta_{0i} \prod_{j=1}^n [\theta_{0i} \psi_i(x_j)] \\ &= \frac{\det \psi_i(x_j)}{\sqrt{\det(\psi_i, \psi_j)}} \int \prod_i d\theta_{0i} \theta_{0i} = \frac{\det \psi_i(x_j)}{\sqrt{\det(\psi_i, \psi_j)}} \end{aligned}$$

Cuantización BRST.

$$\delta \phi^i = \epsilon^a \delta_a \phi^i = \epsilon^a R_a^i(\phi)$$

$$\epsilon^a \delta_a S_m = 0$$

$$[\delta_a, \delta_b] = f_{ab}^c \delta_c$$

$$R_a^i(\phi) = (T_a^R)^i_j \phi^j,$$

$$R_{a\mu}^b = \delta_a^b \partial_\mu + f_{ab}^c A_\mu^c$$

$$Z = \Omega_{\text{gauge}}^{-1} \int d\phi^i e^{-S_m}$$

$$F^A(\phi^i) = 0$$

$$S_{\text{gh}} = b_A c^a \delta_a F^A(\phi^i),$$

$$S_{\text{gf}} = -i B_A F^A(\phi^i)$$

$$Z = \int d\phi^i db_A dc^a dB_A e^{-S_{\text{tot}}}$$

$$S_{\text{tot}} = S_m + S_{\text{gf}} + S_{\text{gh}}$$

$$\delta_\epsilon S_{\text{tot}} = 0$$

$$\delta_\epsilon \phi^i = i \epsilon c^a \delta_a \phi^i, \delta_\epsilon c^a = -\frac{i}{2} \epsilon f_{bc}^a c^b c^c, \delta_\epsilon b_A = \epsilon B_A, \delta_\epsilon B_A = 0$$



$$\delta_\epsilon \delta_{\epsilon'} = 0$$

$$\delta_\epsilon \phi^i = i[\epsilon Q_B, \phi^i]$$

$$S_{\text{gf}} + S_{\text{gh}} = \{Q_B, b_A F^A\}$$

$$\delta S = \{Q_B, b_A \delta F^A\}$$

$$[Q_B, \{Q_B, b_A \delta F^A\}] = 0 \implies Q_B^2 = 0$$

$$|\psi\rangle \text{ closed} \iff |\psi\rangle \in \ker Q_B \iff Q_B |\psi\rangle = 0.$$

$$|\psi\rangle \text{ exact} \iff |\psi\rangle \in \text{Im} Q_B \iff \exists |\chi\rangle: |\psi\rangle = Q_B |\chi\rangle.$$

$$|\psi\rangle \in \mathcal{H}(Q_B) \iff |\psi\rangle \in \ker Q_B, \nexists |\chi\rangle: |\psi\rangle = Q_B |\chi\rangle.$$

$$\mathcal{H}(Q_B) = \frac{\ker Q_B}{\text{Im} Q_B}$$

$$|\psi\rangle \simeq |\psi\rangle + Q_B |\chi\rangle.$$

$$\delta_F \langle \psi_f | \psi_i \rangle = \langle \psi_f | \{Q_B, b_A \delta F^A\} | \psi_i \rangle$$

$$Q_B |\psi\rangle = 0$$

$$\langle \psi | Q_B | \chi \rangle = 0$$

$$|\psi\rangle \text{ physical} \iff |\psi\rangle \in \mathcal{H}(Q_B).$$

$$\frac{\delta F^A}{\delta \phi^i} B_A = - \frac{\delta S_m}{\delta \phi^i}$$

$$\delta_\epsilon b_A = -\epsilon \left(\frac{\delta F^A}{\delta \phi^i} \right)^{-1} \frac{\delta S_m}{\delta \phi^i}$$

$$\{Q_B, b_A B_B M^{AB}\} = i B_A M^{AB} B_B$$

$$S[\phi, b, c, B] = S_0[\phi] + Q_B \Psi[\phi, b, c, B]$$

Formalismo Batalin-Vilkovisky.

$$\mathcal{F}_i(\phi) = \frac{\partial S_0}{\partial \phi^i}$$

$$[T_a, T_b] = F_{ab}^c(\phi) T_c + \lambda_{ab}^i \mathcal{F}_i(\phi)$$

$$m_1 = m_0 - \text{rank} R_a^i$$

$$\delta A_p = \text{d}\lambda_{p-1}$$

$$\delta \lambda_{p-1} = \text{d}\lambda_{p-2}$$



$$\delta\lambda_{p-2}=\mathrm{d}\lambda_{p-3}$$

$$\psi^r=\{\phi^i,B_A,b_A,c^a\}$$

$$S[\psi^r,\psi_r^*]=S_0[\phi]+Q_B\psi^r\psi_r^*$$

$$\psi_r^*=\{\phi_i^*,B^{A*},b^{A*},c_a^*\}$$

$$\psi_r^*=\frac{\partial\Psi}{\partial\psi^r}$$

$$\delta\phi^i=\epsilon_0^{a_0}R_{a_0}^i(\phi^i)$$

$$\delta c_0^{a_0}=\epsilon_1^{a_1}R_{a_1}^{a_0}(\phi^i,c^{a_0})$$

$$\delta c_n^{a_n}=\epsilon_{n+1}^{a_{n+1}}R_{a_{n+1}}^{a_n}(\phi^i,c_0^{a_0},\ldots,c_n^{a_n})$$

$$\psi^r=\{c_n^{a_n}\}_{n=-1,\ldots,\ell'}c_{-1}\colon=\phi$$

$$N_{\mathrm{gh}}(\phi^i)=0,N_{\mathrm{gh}}(c_n^{a_n})=n+1$$

$$|c_n|=|\epsilon_n^{a_n}|+n+1$$

$$N_{\mathrm{gh}}(\psi_r^*)=-1-N_{\mathrm{gh}}(\psi^r),|\psi_r^*|=-|\psi^r|.$$

$$\omega=\sum_r\,\mathrm{d}\psi^r\wedge\mathrm{d}\psi_r^*$$

$$(\psi^r,\psi_s^*)=\delta_{rs},(\psi^r,\psi^s)=0,(\psi_r^*,\psi_s^*)=0$$

$$(A,B)=\frac{\partial_RA}{\partial\psi^r}\frac{\partial_LB}{\partial\psi_r^*}-\frac{\partial_RA}{\partial\psi_r^*}\frac{\partial_LB}{\partial\psi^r}$$

$$(A,B)=-(-1)^{(|A|+1)(|B|+1)}(B,A)$$

$$N_{\mathrm{gh}}((A,B))=N_{\mathrm{gh}}(A)+N_{\mathrm{gh}}(B)+1, |(A,B)|=|A|+|B|+1\,\mathrm{mod}2$$

$$(A,BC)=(A,B)C+(-1)^{|B|C}(A,C)B$$

$$N_{\mathrm{gh}}(S)=0, |S|=0$$

$$S[\psi^r,\psi_r^*=0]=S_0[\phi^i],\left.\frac{\partial_L\partial_RS}{\partial c_{n-1,a_{n-1}}^*\partial c_n^{a_n}}\right|_{\psi^*=0}=R_{a_n}^{a_{n-1}}$$

$$\delta_\theta\psi^r=\theta s\psi^r=-\theta(S,\psi_r)=\theta\frac{\partial_RS}{\partial\psi_r^*}$$

$$\delta_\theta\psi_r^*=\theta s\psi_r^*=-\theta(S,\psi_r^*)=-\theta\frac{\partial_RS}{\partial\psi^r}$$

$$\delta_\theta F=\theta sF=-\theta(S,F)$$



$$(S, S) = 0$$

$$(S, (S, F)) = 0$$

$$s^2 = 0$$

$$s\mathcal{O} = 0$$

$$S' = S + (S, \delta F)$$

$$\psi' = \psi - \frac{\delta F}{\delta \psi^*}, \psi'^* = \psi^* + \frac{\delta F}{\delta \psi}$$

$$S'[\psi, \psi^*] = S\left[\psi - \frac{\delta F}{\delta \psi^*}, \psi^* + \frac{\delta F}{\delta \psi}\right]$$

$$S'[\psi, \psi^*] = S[\psi, \psi^*] + (S, \psi) \frac{\delta F}{\delta \psi} + (S, \psi^*) \frac{\delta F}{\delta \psi^*} = S[\psi, \psi^*] - \frac{\partial_R S}{\partial \psi^*} \frac{\delta F}{\delta \psi} + \frac{\partial_R S}{\partial \psi} \frac{\delta F}{\delta \psi^*}$$

$$(\psi'^r, \psi'^*_s) = \delta_{rs}, (S', S') = 0$$

$$G' = G + (\delta F, G)$$

$$S_{\Psi}[\psi^r] = S\left[\psi^r, \frac{\partial \Psi}{\partial \psi^r}\right], \psi_r^* = \frac{\partial \Psi}{\partial \psi^r}$$

$$N_{\text{gh}}(\Psi) = -1, |\Psi| = 1$$

$$|B| = -|\bar{c}|, N_{\text{gh}}(B) = N_{\text{gh}}(\bar{c}) + 1, \\ s\bar{c} = B, sB = 0$$

$$\bar{S} = S[\psi^r, \psi_r^*] - B\bar{c}^*$$

$$(\bar{S}, \bar{S}) = (S, S) = 0$$

$$(B_{0a_0}, \bar{c}_{0a_0}) := (B_{0a_0}^0, \bar{c}_{0a_0}^0)$$

$$|B_0| = |\epsilon_0|, |\bar{c}_0| = -|\epsilon_0|, N_{\text{gh}}(B_0) = 0, N_{\text{gh}}(c_0) = -1.$$

$$(B_{1a_1}^0, \bar{c}_{1a_1}^0), (\bar{B}_1^{1a_1}, c_1^{1a_1})$$

$$\delta_{\theta}\psi^r=\theta s\psi^r=\theta\left.\frac{\partial_R S}{\partial\psi_r^*}\right|_{\psi_r^*=\partial_r\Psi}$$

$$s^2 \propto \text{eom.}$$

Cuántica BV.

$$Z = \int \mathrm{d}\psi^r \, \mathrm{d}\psi_r^* \mathrm{e}^{-W[\psi^r, \psi_r^*]/\hbar}$$

$$\delta_{\theta} F = \theta \sigma F = (W, F) - \hbar \Delta F$$



$$\Delta = \frac{\partial_R}{\partial \psi_r^*} \frac{\partial_L}{\partial \psi^r}$$

$$(W,W)-2\hbar\Delta W=0,$$

$$\Delta \mathrm{e}^{-W/\hbar}=0$$

$$\delta W = \frac{1}{2} (W,W),$$

$$\mathrm{sdet} J \sim 1 + \Delta W$$

$$W=S+\sum_{p\geq 1}\hbar^pW_p$$

$$\sigma \mathcal{O} = 0$$

$$\delta \langle \mathcal{O} \rangle = 0$$

$$Z=\int \mathrm{d}\psi^r \mathrm{e}^{-W_\Psi[\psi^r]}, W_\Psi[\psi^r]=W\left[\psi^r,\frac{\partial\Psi}{\partial\psi^r}\right]$$

$$Z=\int \mathrm{d}\psi^r \mathrm{e}^{-W_\Psi[\psi^r]}\left(\frac{\partial_R S}{\partial \psi_r^*}\right)_{\psi^*=\partial_\psi \Psi}\frac{\partial(\delta\Psi)}{\partial \psi^r}$$

Gravedad cuántica endógena en planos cuánticos relativistas.

Cálculos preliminares.

$$|\dot{H}| \ll H^2,$$

$$\mathrm{d}s^2=-\mathrm{d}t^2+a^2(t)\mathrm{d}\boldsymbol{x}^2.$$

$$M_P\equiv\sqrt{\frac{\hbar c}{G}}=1.2\times10^{19}\mathrm{GeV}/c^2,$$

$$M_{\rm pl}^{-2}\equiv 8\pi G=(2.4\times10^{18}\mathrm{GeV})^{-2}$$

$$\mathcal{R}_{\boldsymbol{k}}=\int\,\mathrm{d}^3x\mathcal{R}(x)e^{i\boldsymbol{k}\cdot\boldsymbol{x}}.$$

$$\langle \mathcal{R}_{\boldsymbol{k}} \mathcal{R}_{\boldsymbol{k}'} \rangle = (2\pi)^3 P_{\mathcal{R}}(k) \delta(\boldsymbol{k} + \boldsymbol{k}').$$

$$\Delta_{\mathcal{R}}^2(k)\equiv \frac{k^3}{2\pi^2}P_{\mathcal{R}}(k).$$

$$\varepsilon \equiv -\frac{\dot{H}}{H^2}, \tilde{\eta} \equiv \frac{\dot{\varepsilon}}{H\varepsilon}$$

$$\epsilon \equiv \frac{M_{\rm pl}^2}{2} \left(\frac{V'}{V}\right)^2, \eta \equiv M_{\rm pl}^2 \frac{V''}{V}$$



$$\ell_s^2 \equiv \alpha', M_s^2 \equiv \frac{1}{\alpha'}$$

$$2\kappa^2 = (2\pi)^7 (\alpha')^4$$

Métrica de Friedmann-Robertson-Walker (FRW).

$$ds^2 = -dt^2 + a^2(t) d\mathbf{x}^2.$$

$$ds^2 = a^2(\tau)[-d\tau^2 + d\mathbf{x}^2],$$

$$\Delta\tau = \int_0^t \frac{dt'}{a(t')} = \int_0^a \frac{d\ln a}{aH}, \text{ where } H \equiv \frac{1}{a} \frac{da}{dt}$$

$$\frac{d}{dt}(aH)^{-1} = -\frac{1}{a}\left[\frac{\dot{H}}{H^2} + 1\right] < 0 \Rightarrow \varepsilon \equiv -\frac{\dot{H}}{H^2} < 1$$

$$a(t) \propto e^{Ht}$$

$$\begin{aligned} 3M_{\text{pl}}^2 H^2 &= \rho \\ 6M_{\text{pl}}^2 (\dot{H} + H^2) &= -(\rho + 3P) \end{aligned}$$

$$2M_{\text{pl}}^2 \dot{H} = -(\rho + P)$$

$$\varepsilon = \frac{3}{2} \left(1 + \frac{P}{\rho} \right)$$

Acción de Goldstone.

$$U(t, \mathbf{x}) \equiv t + \pi(t, \mathbf{x})$$

$$\delta\psi_m(t, \mathbf{x}) \equiv \psi_m(t + \pi(t, \mathbf{x})) - \psi_m(t).$$

$$g_{ij} = a^2(t) \delta_{ij}$$

$$g_{ij} = a^2(t) e^{2\mathcal{R}(t, \mathbf{x})} \delta_{ij}$$

$$\mathcal{R} = -H\pi + \cdots,$$

$$S = \int d^4x \sqrt{-g} \mathcal{L} \left[U, (\partial_\mu U)^2, \square U, \cdots \right]$$

$$S_\pi^{(2)} = \int d^4x \sqrt{-g} \frac{M_{\text{pl}}^2 |\dot{H}|}{c_s^2} \left[\dot{\pi}^2 - \frac{c_s^2}{a^2} (\partial_i \pi)^2 + 3\varepsilon H^2 \pi^2 \right]$$

$$S_{\mathcal{R}}^{(2)} = \frac{1}{2} \int d^4x a^3 y^2(t) \left[\dot{\mathcal{R}}^2 - \frac{c_s^2}{a^2} (\partial_i \mathcal{R})^2 \right]$$

$$y^2 \equiv 2M_{\text{pl}}^2 \frac{\varepsilon}{c_s^2}$$



$$v \equiv y\mathcal{R} = \int d^3k [v_k(t)a_k e^{ik \cdot x} + c.c.]$$

$$\ddot{v}_k + 3H\dot{v}_k + \frac{c_s^2 k^2}{a^2} v_k = 0$$

$$\omega_k(t) \equiv \frac{c_s k}{a(t)}$$

$$\hat{v}_{\mathbf{k}} = v_k(t)\hat{a}_{\mathbf{k}} + h.c.$$

$$|v_k|^2 = \frac{1}{a^3} \frac{1}{2\omega_k}.$$

$$|v_k|^2 = \frac{1}{2} \frac{1}{a_\star^3} \frac{1}{c_s k/a_\star},$$

$$\frac{c_s k}{a_\star} = H.$$

$$|v_k|^2 = \frac{1}{2} \frac{H^2}{(c_s k)^3},$$

Perturbaciones de Curvatura.

$$P_{\mathcal{R}}(k) \equiv |\mathcal{R}_k|^2 = \frac{1}{4} \frac{H^4}{M_{\text{pl}}^2 |\dot{H}| c_s} \frac{1}{k^3}$$

$$\Delta_{\mathcal{R}}^2(k) \equiv \frac{k^3}{2\pi^2} P_{\mathcal{R}}(k) = \frac{1}{8\pi^2} \frac{H^4}{M_{\text{pl}}^2 |\dot{H}| c_s}.$$

$$n_s - 1 \equiv \frac{d \ln \Delta_{\mathcal{R}}^2}{d \ln k} = -2\varepsilon - \tilde{\eta} - \kappa$$

$$\tilde{\eta} \equiv \frac{\dot{\varepsilon}}{H\varepsilon} \quad \text{and} \quad \kappa \equiv \frac{\dot{c}_s}{Hc_s}$$

Ondas gravitacionales.

$$g_{ij} = a^2(t)(\delta_{ij} + 2h_{ij})$$

$$S_h^{(2)} = \frac{1}{2} \int d^4x a^3 y^2 \left[(\dot{h}_{ij})^2 - \frac{1}{a^2} (\partial_k h_{ij})^2 \right]$$

$$y^2 \equiv \frac{1}{4} M_{\text{pl}}^2$$

$$\Delta_h^2(k) \equiv \frac{k^3}{2\pi^2} P_h(k) = \frac{2}{\pi^2} \frac{H^2}{M_{\text{pl}}^2},$$



$$n_t \equiv \frac{d \ln \Delta_h^2}{d \ln k} = -2\varepsilon$$

$$r \equiv \frac{\Delta_h^2}{\Delta_{\mathcal{R}}^2}$$

$$\frac{H}{M_{\rm pl}} = \pi \Delta_{\mathcal{R}}(k_*) \sqrt{\frac{r}{2}},$$

$$H=3\times 10^{-5}\left(\frac{r}{0.1}\right)^{1/2}M_{\rm pl}$$

$$E_{\rm inf} \equiv (3H^2 M_{\rm pl}^2)^{1/4} = 8 \times 10^{-3} \left(\frac{r}{0.1}\right)^{1/4} M_{\rm pl}$$

Anisotropías CMB.

$$e+p\rightarrow H+\gamma$$

$$\Delta T(\boldsymbol{n}) \equiv T(\boldsymbol{n}) - \bar{T}.$$

$$C(\theta) \equiv \left\langle \frac{\Delta T}{\bar{T}}(\boldsymbol{n}) \frac{\Delta T}{\bar{T}}(\boldsymbol{n}') \right\rangle,$$

$$\frac{\Delta T(\boldsymbol{n})}{\bar{T}} = \sum_{\ell=0}^{\infty} \sum_{m=-\ell}^{+\ell} a_{\ell m} Y_{\ell m}(\boldsymbol{n})$$

$$\langle a_{\ell m} a_{\ell' m'}^* \rangle = C_{\ell} \delta_{\ell \ell'} \delta_{m m'}$$

$$C_{\ell}=2\pi\int_{-1}^1\mathrm{d}\cos\theta C(\theta)P_{\ell}(\cos\theta)$$

$$\hat{C}_{\ell} = \frac{1}{2\ell+1} \sum_m |a_{\ell m}|^2$$

$$\mathrm{var}(\hat{C}_{\ell}) \equiv \langle \hat{C}_{\ell} \hat{C}_{\ell} \rangle - \langle \hat{C}_{\ell} \rangle^2 = \frac{2}{2\ell+1} C_{\ell}^2$$

$$C_{\ell} = \int \mathrm{d} \ln k \Delta_{\mathcal{R}}^2(k) T_{\ell}^2(k)$$

Polarización CMB.

$$(Q\pm iU)'(\boldsymbol{n})=e^{\mp 2i\psi}(Q\pm iU)(\boldsymbol{n}).$$

$$(Q\pm iU)(\boldsymbol{n})=\sum_{\ell m}a_{\pm 2,\ell m\pm 2}Y_{\ell m}(\boldsymbol{n})$$



$$E(\mathbf{n}) \equiv a_{E,\ell m} Y_{\ell m}(\mathbf{n}), \quad a_{E,\ell m} \equiv -\frac{a_{2,\ell m} + a_{-2,\ell m}}{2}$$

$$B(\mathbf{n}) \equiv a_{B,\ell m} Y_{\ell m}(\mathbf{n}), \quad a_{B,\ell m} \equiv -\frac{a_{2,\ell m} - a_{-2,\ell m}}{2i}$$

$$\langle a_{X,\ell m} a_{Y,\ell' m'}^* \rangle = C_\ell^{XY} \delta_{\ell\ell'} \delta_{mm'}, X, Y \equiv \{T, E, B\}.$$

Estructura Escalar.

$$P_\delta(z, k) = T_\delta^2(z, k) P_{\mathcal{R}}(k)$$

$$\delta_g(z, \mathbf{x}) = b(z) \delta(z, \mathbf{x})$$

$$D_V(\bar{z}) \equiv \left[(1 + \bar{z})^2 D_A^2(\bar{z}) \frac{c\bar{z}}{H(\bar{z})} \right]^{1/3}$$

Modelo CDM.

$$\Delta_{\mathcal{R}}^2(k) = A_s \left(\frac{k}{k_\star} \right)^{n_s-1}.$$

$$A_s = (2.196_{-0.060}^{+0.051}) \times 10^{-9}.$$

$$n_s = 0.9603 \pm 0.0073$$

Fluctuaciones Escalares.

$$\Delta_{\mathcal{R}}^2(k) = A_s \left(\frac{k}{k_\star} \right)^{n_s-1+\frac{1}{2}\alpha_s \ln(k/k_\star)}$$

Fluctuaciones Tensoriales.

$$r < 0.12 \text{ (95\% limit)}.$$

$$\langle 0 | \hat{\mathcal{R}}_{\mathbf{k}_1} \hat{\mathcal{R}}_{\mathbf{k}_2} | 0 \rangle = (2\pi)^3 P_{\mathcal{R}}(k_1) \delta(\mathbf{k}_1 + \mathbf{k}_2),$$

$$\langle \Omega | \hat{\mathcal{R}}_{\mathbf{k}_1} \cdots \hat{\mathcal{R}}_{\mathbf{k}_n} | \Omega \rangle \propto \int [\mathcal{D}\mathcal{R}] \mathcal{R}_{\mathbf{k}_1} \cdots \mathcal{R}_{\mathbf{k}_n} e^{iS[\mathcal{R}]},$$

$$\langle \Omega | \hat{\mathcal{R}}_{\mathbf{k}_1} \hat{\mathcal{R}}_{\mathbf{k}_2} \hat{\mathcal{R}}_{\mathbf{k}_3} | \Omega \rangle = (2\pi)^3 B_{\mathcal{R}}(k_1, k_2, k_3) \delta(\mathbf{k}_1 + \mathbf{k}_2 + \mathbf{k}_3).$$

$$f_{\text{NL}} \equiv \frac{5}{18} \frac{B_{\mathcal{R}}(k, k, k)}{P_{\mathcal{R}}^2(k)}$$

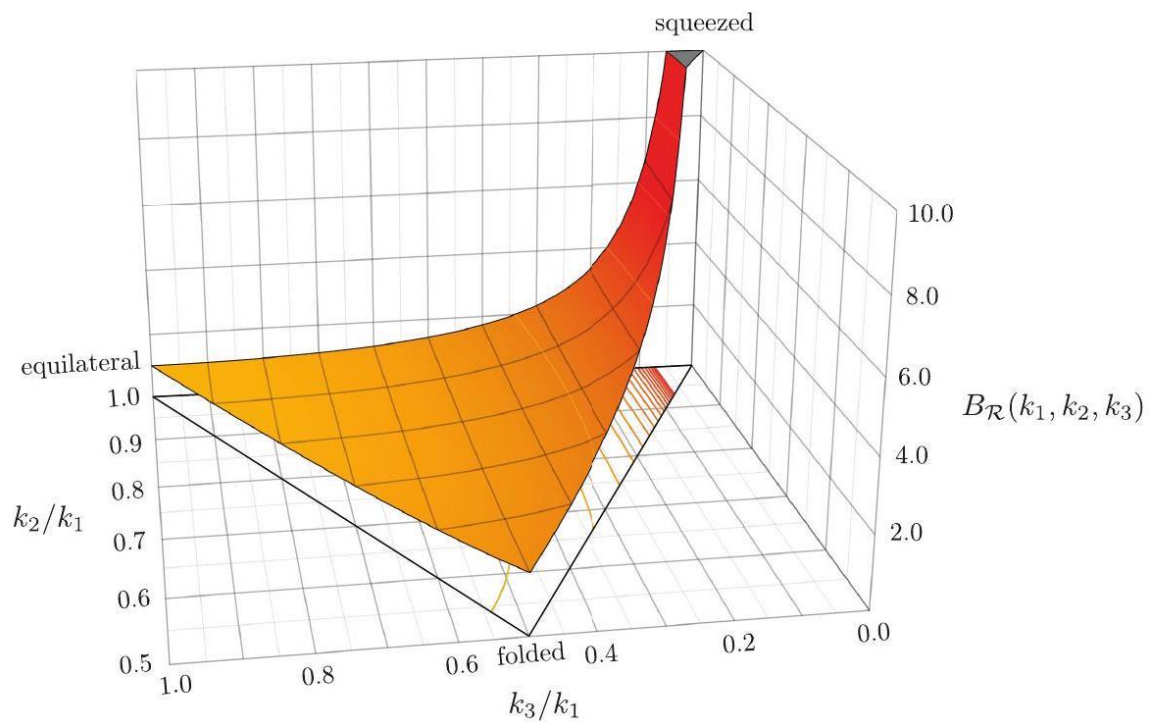
$$B_{\text{local}} \equiv \frac{6}{5}(P_1 P_2 + \text{perms.})$$

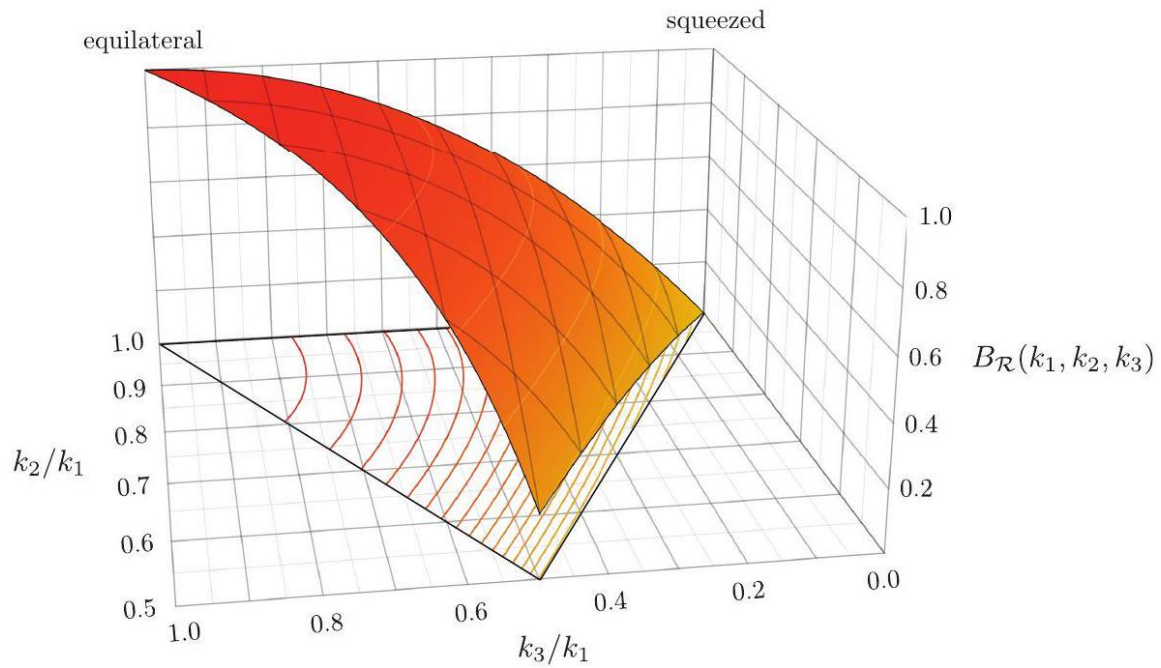
$$B_{\text{equil}} \equiv \frac{3}{5}(6(P_1^3 P_2^2 P_3)^{1/3} - 3P_1 P_2 - 2(P_1 P_2 P_3)^{2/3} + \text{perms.})$$

$$B_{\text{ortho}} \equiv \frac{3}{5}(18(P_1^3 P_2^2 P_3)^{1/3} - 9P_1 P_2 - 8(P_1 P_2 P_3)^{2/3} + \text{perms.})$$

$$\mathcal{R}(\mathbf{x}) \equiv \mathcal{R}_g(\mathbf{x}) + \frac{3}{5} f_{\text{NL}}^{\text{local}} [\mathcal{R}_g^2(\mathbf{x}) - \langle \mathcal{R}_g^2 \rangle]$$

$$\lim_{k_1 \rightarrow 0} \frac{B_{\mathcal{R}}(k_1, k_2, k_3)}{P_{\mathcal{R}}(k_1) P_{\mathcal{R}}(k_2)} = (1 - n_s) \ll 1$$





Figuras 34 y 35. Curvatura local por deformación del espacio – tiempo cuántico, a propósito de una particular supermasiva, esto en dimension $\mathbb{R}^4$.

$$S_{\pi}^{(3)} = \int d^4x \sqrt{-g} \frac{M_{\text{pl}}^2 \dot{H}}{c_s^2} (1 - c_s^2) \left(\frac{\dot{\pi}(\partial_i \pi)^2}{a^2} + \frac{A}{c_s^2} \dot{\pi}^3 \right)$$

$$\begin{aligned} f_{\text{NL}}^{\text{local}} &= 2.7 \pm 5.8, \\ f_{\text{NL}}^{\text{equil}} &= -42 \pm 75, \\ f_{\text{NL}}^{\text{ortho}} &= -25 \pm 39. \end{aligned}$$

Adiabaticity.

$$\delta_I(t, \mathbf{x}) \equiv \frac{\bar{\rho}_I(t + \pi(t, \mathbf{x})) - \bar{\rho}_I(t)}{\bar{\rho}_I(t)} \approx \frac{\dot{\bar{\rho}}_I}{\bar{\rho}_I} \pi(t, \mathbf{x}).$$

$$\mathcal{S} \equiv \delta_c - \frac{3}{4} \delta_\gamma$$

$$\alpha \equiv \frac{P_{\mathcal{S}}}{P_{\mathcal{R}}}$$

$$\alpha_0 < 0.036$$

$$\alpha_{+1} < 0.0025.$$

Curvatura isotrópica.

$$n_t = 2 \frac{\dot{H}}{H^2}$$

$$n_t = -\frac{r}{8}$$

$$\alpha_s = 16\varepsilon^2 - 6\varepsilon\tilde{\eta} + \tilde{\eta}\chi$$

$$\mathcal{N}^{\rm CMB} \sim \left(\frac{\ell_{\rm max}}{\ell_{\rm min}}\right)^2 \sim 10^6$$

$$\mathcal{N}_{\rm linear}^{\rm LSS} \sim \left(\frac{k_{\rm max}}{k_{\rm min}}\right)^3 \sim 10^9$$

$$\mathcal{N}_{\rm linear}^{\rm Euclid} \sim \left(\frac{k_{\rm max}}{k_{\rm min}}\right)^3 \sim 10^6,$$

$$|f_{\rm NL}^{\rm local}| \gtrsim \mathcal{O}(1)$$

$$f_{\rm NL}^{\rm equil} \sim \mathcal{O}(1)$$

$$M_{\rm pl} \equiv \frac{1}{\sqrt{8\pi G_N}} = 2.4 \times 10^{18} \text{GeV}$$

Acción efectiva.

$$\mathcal{L}[\phi,\Psi] = \mathcal{L}_l[\phi] + \mathcal{L}_h[\Psi] + \mathcal{L}_{lh}[\phi,\Psi]$$

$$e^{iS_{\rm eff}[\phi]} = \int [\mathcal{D}\Psi] e^{iS[\phi,\Psi]}$$

$$\phi (-\Box + M^2)^{-1} \phi = \frac{\phi}{M^2} \left(1 + \frac{\Box}{M^2} + \cdots \right) \phi,$$

$$\mathcal{L}_{\rm eff}[\phi] = \mathcal{L}_l[\phi] + \sum_i c_i(g) \frac{\mathcal{O}_i[\phi]}{M^{\delta_i-4}},$$

$$\mathcal{L}[\phi,\Psi] = -\frac{1}{2}(\partial\phi)^2 - \frac{1}{2}m^2\phi^2 - \frac{1}{4!}\lambda\phi^4 - \frac{1}{2}(\partial\Psi)^2 - \frac{1}{2}M^2\Psi^2 - \frac{1}{4}g\phi^2\Psi^2.$$

$$\mathcal{L}_{\rm eff}[\phi] = -\frac{1}{2}(\partial\phi)^2 - \frac{1}{2}m_{\rm R}^2\phi^2 - \frac{1}{4!}\lambda_{\rm R}\phi^4 - \sum_{i=1}^{\infty} \left(\frac{c_i(g)}{M^{2i}}\phi^{4+2i} + \frac{d_i(g)}{M^{2i}}(\partial\phi)^2\phi^{2i} + \cdots \right)$$



$$m_{\text{R}}^2 = \text{---} + \text{---} \bigcirc \text{---} + \dots ,$$

$$\lambda_{\text{R}} = \text{X} + \text{---} \bigcirc \text{---} + \dots ,$$

$$m_{\text{R}}^2 = m^2 + \frac{g}{32\pi^2}(\Lambda^2 - M^2 L) + \dots$$

$$\lambda_{\text{R}} = \lambda - \frac{3g^2}{32\pi^2}L + \dots$$

$$L \rightarrow \frac{1}{\epsilon} + \gamma - \ln(4\pi)$$

$$c_1 = \text{---} \bigcirc \text{---} + \dots \sim \mathcal{O}(g^3) + \dots .$$

Parametrizaciones de simetría.

$$\mathcal{L}_{\text{eff}}[\phi] = \mathcal{L}_l[\phi] + \sum_i c_i \frac{\mathcal{O}_i[\phi]}{\Lambda^{\delta_i-4}},$$

$$\Delta m_e = \alpha \Lambda$$

$$\Delta m_e = \alpha m_e \ln(\Lambda/m_e)$$

$$m_{\pi^+}^2 - m_{\pi^0}^2 = \frac{3\alpha}{4\pi} \Lambda^2$$

$$\frac{m_{K_L^0} - m_{K_S^0}}{m_{K_L^0}} = \frac{G_F^2 f_K^2}{6\pi^2} \sin^2 \theta_c \Lambda^2,$$

$$\Delta m_H^2 \sim \frac{y_t^2}{(4\pi)^2} \Lambda^2$$

$$\langle T_{\mu\nu} \rangle = -\rho_{\text{vac}} g_{\mu\nu}$$

$$\Delta \rho_{\text{vac}} \sim \Lambda^4$$

$$\Delta m^2 \propto \Lambda^2$$

$$\phi \mapsto \phi + \text{const}.$$

$$\Delta m^2 \propto m^2$$

$$\lim_{g \rightarrow 0} c_i(g) = 0$$

$$S_{\text{EH}} = \frac{M_{\text{pl}}^2}{2} \int d^4x \sqrt{-g} R$$

$$S_{\text{EH}} = \int d^4x \left[(\partial h)^2 + \frac{1}{M_{\text{pl}}} h (\partial h)^2 + \frac{1}{M_{\text{pl}}^2} h^2 (\partial h)^2 + \dots \right]$$

$$S_{\text{YM}} = \int d^4x [(\partial A)^2 + g A^2 \partial A + g^2 A^4]$$

$$S_g = \int d^4x \sqrt{-g} \left[M_{\Lambda}^4 + \frac{M_{\text{pl}}^2}{2} R + c_1 R^2 + c_2 R_{\mu\nu} R^{\mu\nu} + \frac{1}{M^2} (d_1 R^3 + \dots) + \dots \right]$$

$$S_g = \int d^{10}X \sqrt{-G} \left[\frac{M_{10}^2}{2} R + \frac{\zeta(3)}{3 \cdot 2^5} \frac{1}{M^6} \mathcal{R}^4 + \dots \right]$$

$$M^2 = \frac{4}{\alpha'}$$

$$S_{\text{eff}}[\phi, g] = S_g + S_{\text{eff}}[\phi] + S_{g, \phi}$$

$$S_{g, \phi} = \int d^4x \sqrt{-g} \left[\sum_i c_i \frac{\mathcal{O}_i[g, \phi]}{\Lambda^{\delta_i - 4}} \right]$$

$$S_{g, \phi}^{(4)} = \int d^4x \sqrt{-g} \xi \phi^2 R$$

$$g_{\mu\nu} \mapsto \bar{g}_{\mu\nu} \equiv e^{2\omega(\phi)} g_{\mu\nu},$$

Dinámicas y perturbaciones inflacionarias.

$$S = \int d^4x \sqrt{-g} \left[\frac{M_{\text{pl}}^2}{2} R - \frac{1}{2} (\partial\phi)^2 - V(\phi) \right]$$

$$3M_{\text{pl}}^2 H^2 = \frac{1}{2} \dot{\phi}^2 + V \text{ and } \ddot{\phi} + 3H\dot{\phi} = -V',$$

$$\varepsilon = -\frac{\dot{H}}{H^2} = \frac{\frac{1}{2}\dot{\phi}^2}{M_{\text{pl}}^2 H^2}$$



$$\epsilon \equiv \frac{M_{\text{pl}}^2}{2} \left(\frac{V'}{V} \right)^2 \ll 1, |\eta| \equiv M_{\text{pl}}^2 \frac{|V''|}{V} \ll 1.$$

$$\pi = \frac{\delta\phi}{\dot{\phi}}$$

$$\mathcal{R}(t,\boldsymbol{x})=-H\pi(t,\boldsymbol{x})=-\frac{H}{\dot{\phi}}\delta\phi(t,\boldsymbol{x})$$

$$\Delta_{\mathcal{R}}^2=\frac{1}{24\pi^2}\frac{1}{\epsilon}\frac{V}{M_{\text{pl}}^4},\Delta_h^2=\frac{2}{3\pi^2}\frac{V}{M_{\text{pl}}^4}$$

$$\begin{array}{l} n_s-1=2\eta-6\epsilon \\ r=16\epsilon \end{array}$$

$$N_{\star}=\int_{\phi_{\rm end}}^{\phi_{\star}}\frac{\mathrm{d}\phi}{M_{\rm pl}}\frac{1}{\sqrt{2\epsilon}}$$

$$V(\phi)=\mu^{4-p}\phi^p$$

$$\epsilon=\frac{p^2}{2}\left(\frac{M_{\rm pl}}{\phi}\right)^2,\eta=p(p-1)\left(\frac{M_{\rm pl}}{\phi}\right)^2.$$

$$N_{\star}\approx\frac{1}{2p}\left(\frac{\phi_{\star}}{M_{\rm pl}}\right)^2,$$

$$n_s-1=-\frac{(2+p)}{2N_{\star}},r=\frac{4p}{N_{\star}}.$$

$$\begin{array}{lll} p=1: & n_s \approx 0.975, & r \approx 0.07, \quad \phi_{\star} \approx 11 M_{\rm pl}, \\ p=2: & n_s \approx 0.967, & r \approx 0.13, \quad \phi_{\star} \approx 15 M_{\rm pl}, \\ p=3: & n_s \approx 0.958, & r \approx 0.20, \quad \phi_{\star} \approx 19 M_{\rm pl}, \\ p=4: & n_s \approx 0.950, & r \approx 0.27, \quad \phi_{\star} \approx 22 M_{\rm pl}. \end{array}$$

$$V(\phi)=\frac{V_0}{2}\Big[1-\cos\left(\frac{\phi}{f}\right)\Big],$$

$$\begin{array}{l} n_s-1=-\alpha\frac{e^{N_{\star}\alpha}+1}{e^{N_{\star}\alpha}-1}\overset{\alpha\ll 1}{\rightarrow}-\frac{2}{N_{\star}}, \\ r=8\alpha\frac{1}{e^{N_{\star}\alpha}-1}\overset{\alpha\ll 1}{\rightarrow}+\frac{8}{N_{\star}}, \end{array}$$

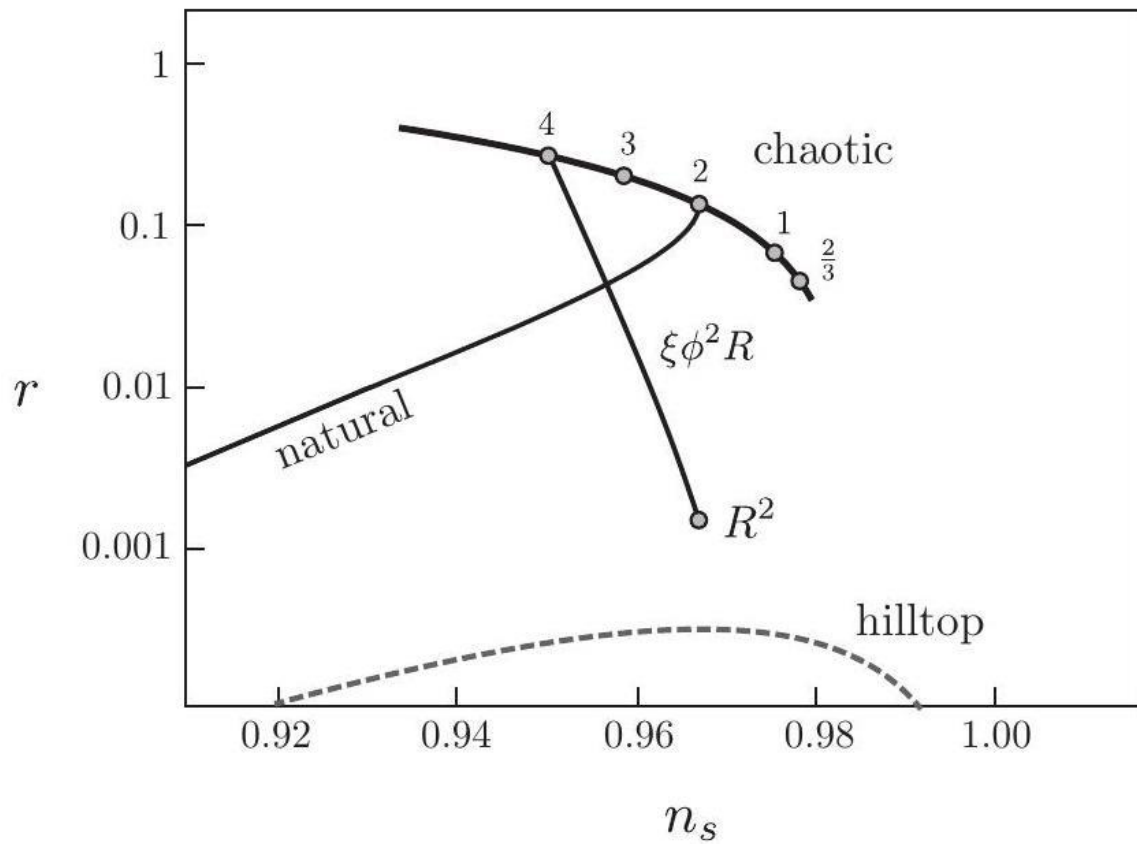
$$V(\phi)=V_0+\frac{1}{2}m^2\phi^2+\cdots$$

$$V(\phi)=V_0\left[1+\frac{1}{2}\eta_0\frac{\phi^2}{M_{\rm pl}^2}+\cdots\right],\text{ where }\eta\approx\eta_0<0$$



$$n_s - 1 = 2\eta_0$$

$$r = 2(1 - n_s)^2 e^{-N_*(1-n_s)} \left(\frac{\phi_{\text{end}}}{M_{\text{pl}}} \right)^2 \approx 10^{-3} \left(\frac{\phi_{\text{end}}}{M_{\text{pl}}} \right)^2.$$



$$V(\phi) = V_0 \left[1 + \lambda_0 \frac{\phi}{M_{\text{pl}}} + \frac{1}{2} \eta_0 \frac{\phi^2}{M_{\text{pl}}^2} + \frac{1}{3!} \mu_0 \frac{\phi^3}{M_{\text{pl}}^3} + \dots \right]$$

$$V(\phi) \approx V_0 \left[1 + \lambda_0 \frac{\phi}{M_{\text{pl}}} + \frac{1}{3!} \mu_0 \frac{\phi^3}{M_{\text{pl}}^3} + \dots \right]$$

$$n_s - 1 = -4 \sqrt{\frac{\lambda_0 \mu_0}{2}} \cot \left(N_* \sqrt{\frac{\lambda_0 \mu_0}{2}} \right),$$

$$r = 16\lambda_0^2$$

$$V(\phi, \Psi) = V(\phi) + V(\Psi) + \frac{1}{2} g \phi^2 \Psi^2,$$

$$V(\Psi) \equiv \frac{1}{4\lambda} (M^2 - \lambda \Psi^2)^2$$

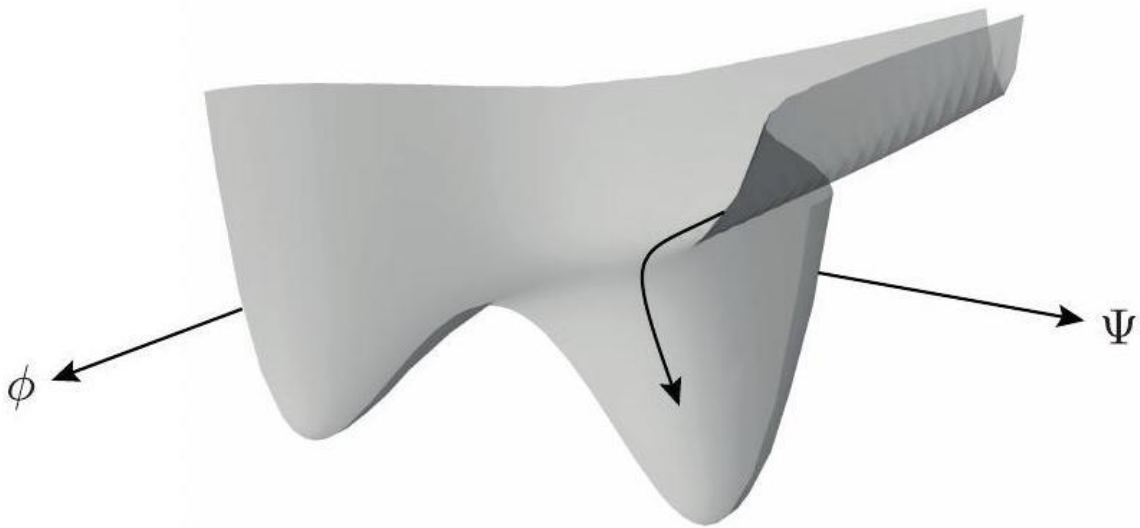


Figura 36. Membranas en dimensión $\mathbb{R}^4$

$$M_{\Psi}^2(\phi) = -M^2 + g\phi^2$$

$$S = \frac{M_{\text{pl}}^2}{2} \int d^4x \sqrt{-g} \left(R + \frac{\alpha}{2M_{\text{pl}}^2} R^2 \right)$$

$$S = \int d^4x \sqrt{-\tilde{g}} \left(\frac{M_{\text{pl}}^2}{2} \tilde{R} - \frac{1}{2} (\partial\phi)^2 - V(\phi) \right)$$

$$V(\phi) = \frac{M_{\text{pl}}^4}{4\alpha} \left(1 - \exp \left[-\sqrt{\frac{2}{3}} \frac{\phi}{M_{\text{pl}}} \right] \right)^2.$$

$$\eta = -\frac{4}{3} e^{-\sqrt{2/3}\phi/M_{\text{pl}}}, \epsilon = \frac{3}{4} \eta^2$$

$$\alpha = 2.2 \times 10^8$$

$$n_s - 1 \approx -\frac{2}{N_*}, r \approx \frac{12}{N_*^2}$$

$$S = \int d^4x \sqrt{-g} \left[\frac{M_{\text{pl}}^2}{2} \left(1 + \xi \frac{\varphi^2}{M_{\text{pl}}^2} \right) R - \frac{1}{2} (\partial\varphi)^2 - \frac{\lambda}{4} \varphi^4 \right]$$

$$S = \int d^4x \sqrt{-\tilde{g}} \left[\frac{M_{\text{pl}}^2}{2} \tilde{R} - \frac{1}{2} k(\varphi) (\partial\varphi)^2 - V(\varphi) \right]$$

$$k(\varphi) = \frac{1 + (6\xi + 1)\psi^2}{(1 + \psi^2)^2}$$

$$V(\varphi) = \frac{\lambda M_{\text{pl}}^4}{4\xi^2} \frac{\psi^4}{(1 + \psi^2)^2}, \psi^2 \equiv \frac{\xi \varphi^2}{M_{\text{pl}}^2}$$

$$\frac{\phi}{M_{\text{pl}}} = \sqrt{\frac{6\xi+1}{\xi}} \sinh^{-1}(\sqrt{6\xi+1}\psi) - \sqrt{6} \sinh^{-1}\left(\sqrt{6\xi} \frac{\psi}{\sqrt{1+\psi^2}}\right).$$

$$\frac{\phi}{M_{\text{pl}}} \approx \sqrt{\frac{3}{2}} \ln(1+\psi^2),$$

$$V(\phi) = \frac{\lambda M_{\text{pl}}^4}{4\xi^2} \left(1 - \exp \left[-\sqrt{\frac{2}{3}} \frac{\phi}{M_{\text{pl}}} \right] \right)^2.$$

$$\xi = 47000\sqrt{\lambda}$$

$$n_s - 1 = -\frac{32\xi}{16\xi N_\star - 1} \stackrel{\xi \gg 1}{\rightarrow} -\frac{2}{N_\star},$$

$$r = +\frac{12}{N_\star^2} \frac{6\xi + 1}{6\xi} \stackrel{\xi \gg 1}{\rightarrow} +\frac{12}{N_\star^2}.$$

Inflación K.

$$S = \int \mathrm{d}^4x \sqrt{-g} \left[\frac{M_{\text{pl}}^2}{2} R + P(X, \phi) \right]$$

$$3M_{\text{pl}}^2H^2 = 2P_{,X}X - P \text{ and } \frac{d}{dt}(a^3P_{,X}\dot{\phi}) = a^3P_{,\phi}$$

$$\varepsilon = -\frac{\dot{H}}{H^2} = \frac{3XP_{,X}}{2XP_{,X} - P}$$

$$c_s^2 = \frac{dP}{d\rho} = \frac{P_{,X}}{P_{,X} + 2XP_{,XX}}.$$

$$\begin{aligned} n_s - 1 &= -2\varepsilon - \tilde{\eta} - \kappa \\ r &= 16\varepsilon c_s \end{aligned}$$

$$S_{\text{eff}}[\phi] = \int \mathrm{d}^4x \sqrt{-g} \left[\frac{M_{\text{pl}}^2}{2} R + \mathcal{L}_l[\phi] + \sum_i c_i \frac{\mathcal{O}_i[\phi]}{\Lambda^{\delta_i-4}} \right]$$

$$\Delta m^2 \sim \Lambda^2.$$

$$\Delta\eta\sim\frac{\Lambda^2}{H^2}\gtrsim1$$

$$\Delta m^2 \sim H^2$$

$$\Delta\eta\sim1.$$

$$\phi \mapsto \phi + \text{const} \, .$$



Operadores en altas dimensiones.

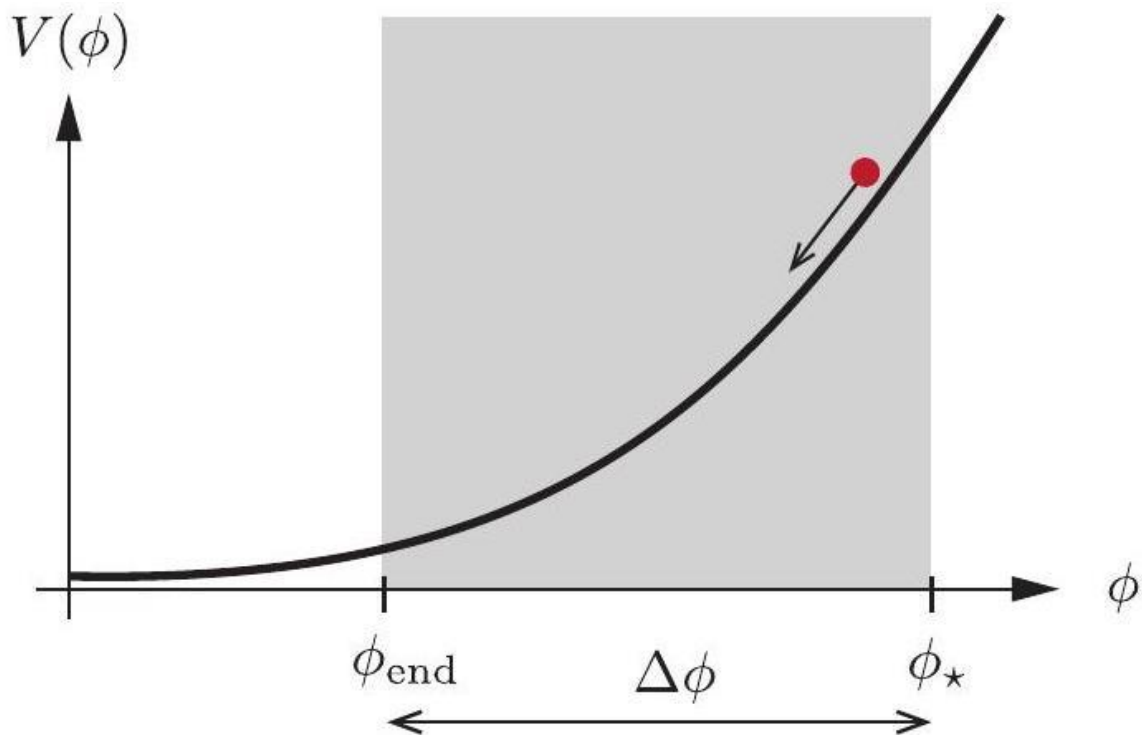
$$\mathcal{O}_6 = c V_I(\phi) \frac{\phi^2}{\Lambda^2}$$

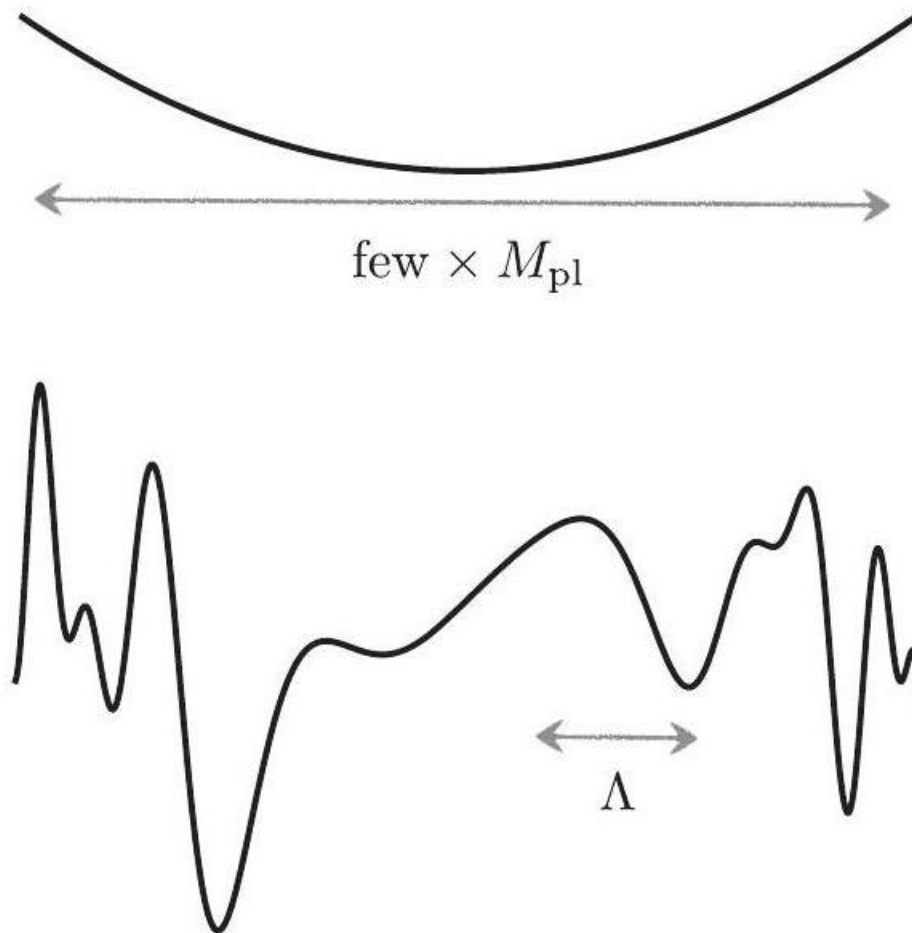
$$\Delta\eta \approx 2c \left(\frac{M_{\text{pl}}}{\Lambda} \right)^2$$

$$\mathcal{O}_\delta = c \langle V \rangle \left(\frac{\phi}{\Lambda} \right)^{\delta-4}$$

$$\Delta\eta \approx c(\delta-4)(\delta-5) \left(\frac{M_{\text{pl}}}{\Lambda} \right)^2 \left(\frac{\phi}{\Lambda} \right)^{\delta-6}.$$

Ondas gravitacionales en campos planckianos y campos inflacionarios.





Figuras 37 y 38. Fluctuaciones gravitacionales en la curvatura local.

$$r = 8 \left(\frac{1}{M_{\text{pl}}} \frac{d\phi}{dN} \right)^2, \text{ where } dN \equiv H dt.$$

$$\frac{\Delta\phi}{M_{\text{pl}}} = \int_0^{N_*} dN \sqrt{\frac{r(N)}{8}}.$$

$$N_{\text{eff}} \equiv \int_0^{N_*} dN \sqrt{\frac{r(N)}{r_*}}$$

$$\frac{\Delta\phi}{M_{\text{pl}}} = N_{\text{eff}} \sqrt{\frac{r_*}{8}}.$$

$$\frac{d \ln r}{dN} = - \left[n_s - 1 + \frac{r}{8} \right],$$

$$\frac{\Delta\phi}{M_{\text{pl}}} \gtrsim \left(\frac{r}{0.01} \right)^{1/2}$$

$$\frac{\Delta\phi}{M_{\rm pl}}\gtrsim 0.25\times\left(\frac{r}{0.01}\right)^{1/2}$$

$$\frac{\Delta V}{V}=c_1\frac{V''}{M_{\rm pl}^2}+c_2\frac{V}{M_{\rm pl}^4},$$

$$\mathcal{L}_{\rm eff}[\phi]=\mathcal{L}_l[\phi]+\sum_{i=1}^{\infty}\left(\frac{c_i}{\Lambda^{2i}}\phi^{4+2i}+\frac{d_i}{\Lambda^{2i}}(\partial\phi)^2\phi^{2i}+\frac{e_i}{\Lambda^{4i}}(\partial\phi)^{2(i+1)}+\cdots\right),$$

$$\phi \mapsto \phi + \text{const} \, .$$

$$\frac{\Delta\phi}{M_{\rm pl}}=(c_sP_{,X})^{-1/2}\sqrt{\frac{r}{8}}\Delta N.$$

$$\Delta\mathcal{L}=P\left(X-V(\phi)\frac{\phi^2}{M_{\rm pl}^2}\right)=P(X)-P_{,X}V(\phi)\frac{\phi^2}{M_{\rm pl}^2}+\cdots$$

$$c_s^2=\frac{P_{,X}}{P_{,X}+2XP_{,XX}}$$

$$P=X+\frac{1}{2}\frac{X^2}{\Lambda^4}+\cdots$$

$$\left|f_{\rm NL}^{\rm equil}\right|\sim \frac{1}{c_s^2}-1\approx \frac{X}{\Lambda^4}+\cdots\ll 1.$$

$$\mathcal{O}_5=\frac{\psi X}{\Lambda}$$

$$\Lambda \gtrsim 10^5 H$$

$$\Lambda \gtrsim \left(\frac{r}{0.01}\right)^{1/2} M_{\rm pl}$$

$$S_{\rm P}=-\frac{1}{4\pi\alpha'}\int\,{\rm d}^2\sigma\sqrt{-h}h^{ab}\partial_aX^M(\sigma)\partial_bX^N(\sigma)\eta_{MN}$$

$$S_{\rm P}=-\frac{1}{4\pi\alpha'}\int\,{\rm d}^2\sigma\partial^aX^M\partial_aX_M$$

$$S_{\sigma}=-\frac{1}{4\pi\alpha'}\int\,{\rm d}^2\sigma\sqrt{-h}([h^{ab}G_{MN}(X)+\epsilon^{ab}B_{MN}(X)]\partial_aX^M\partial_bX^N+\alpha'\Phi(X)R(h))$$

$$S_{\rm B}=\frac{1}{2\kappa_D^2}\int\,{\rm d}^D X\sqrt{-G}e^{-2\Phi}\left(R+4(\partial\Phi)^2-\frac{1}{2}|H_3|^2-\frac{2(D-26)}{3\alpha'}+\mathcal{O}(\alpha')\right)$$



$$e^{-\Phi\chi} = e^{-\Phi(2-2g)} \equiv g_s^{2g-2}$$

$$S = S_{\text{P}} + S_{\text{F}} = -\frac{1}{4\pi\alpha'} \int \text{d}^2\sigma (\partial^a X^M \partial_a X_M - i\bar{\psi}^M \rho^a \partial_a \psi_M)$$

$$\{\rho^a,\rho^b\}=-2\eta^{ab}$$

$$\psi^M\equiv\begin{pmatrix}\psi^M_-\\ \psi^M_+\end{pmatrix}$$

$$S_{\text{F}}=\frac{i}{2\pi\alpha'}\int\text{d}^2\sigma(\psi^M_-\partial_+\psi^N_-+\psi^M_+\partial_-\psi^N_+)\eta_{MN}$$

$$S_{\text{NS}}=\frac{1}{2\kappa^2}\int\text{d}^{10}X\sqrt{-G}e^{-2\Phi}\left(R+4(\partial\Phi)^2-\frac{1}{2}|H_3|^2\right)$$

$$2\kappa^2=(2\pi)^7(\alpha')^4$$

$$S_{\text{IIA}}=S_{\text{NS}}+S_{\text{R}}^{(\text{IIA})}+S_{\text{CS}}^{(\text{IIA})},$$

$$S_{\text{R}}^{(\text{IIA})}=-\frac{1}{4\kappa^2}\int\text{d}^{10}X\sqrt{-G}\left(|F_2|^2+|\tilde{F}_4|^2\right),$$

$$S_{\text{CS}}^{(\text{IIA})}=-\frac{1}{4\kappa^2}\int B_2\wedge F_4\wedge F_4,$$

$$S_{\text{IIB}}=S_{\text{NS}}+S_{\text{R}}^{(\text{IIB})}+S_{\text{CS}}^{(\text{IIB})},$$

$$S_{\text{R}}^{(\text{IIB})}=-\frac{1}{4\kappa^2}\int\text{d}^{10}X\sqrt{-G}\left(|F_1|^2+|\tilde{F}_3|^2+\frac{1}{2}|\tilde{F}_5|^2\right),$$

$$S_{\text{CS}}^{(\text{IIB})}=-\frac{1}{4\kappa^2}\int C_4\wedge H_3\wedge F_3,$$

$$\tilde{F}_5=\star\tilde{F}_5$$

$$G_{E,MN}\equiv e^{-\Phi/2}G_{MN}$$

$$G_3\equiv F_3-\tau H_3$$

$$\tau\equiv C_0+ie^{-\Phi}$$

$$S_{\text{IIB}}=\frac{1}{2\kappa^2}\int\text{d}^{10}X\sqrt{-G_E}\left[R_E-\frac{|\partial\tau|^2}{2(\text{Im}(\tau))^2}-\frac{|G_3|^2}{2\text{Im}(\tau)}-\frac{|\tilde{F}_5|^2}{4}\right]-\frac{i}{8\kappa^2}\int\frac{C_4\wedge G_3\wedge\bar{G}_3}{\text{Im}(\tau)}$$



Membranas en dimensión $\mathbb{R}^4$.

$$S_{\text{CS}} = \mu_p \int_{\Sigma_{p+1}} C_{p+1}$$

$$S_{\text{D}} = -T_p \int d^{p+1} \sigma \sqrt{-\det(G_{ab})}$$

$$G_{ab} \equiv \frac{\partial X^M}{\partial \sigma^a} \frac{\partial X^N}{\partial \sigma^b} G_{MN}$$

$$S_{\text{BI}} = -Q_p \int d^{p+1} \sigma \sqrt{-\det(\eta_{ab} + 2\pi\alpha' F_{ab})} = -Q_p \int d^{p+1} \sigma \left(1 + \frac{(2\pi\alpha')^2}{4} F_{ab} F^{ab} + \dots \right)$$

$$S_{\text{DBI}} = -g_s T_p \int d^{p+1} \sigma e^{-\Phi} \sqrt{-\det(G_{ab} + \mathcal{F}_{ab})}$$

$$\mathcal{F}_{ab} \equiv B_{ab} + 2\pi\alpha' F_{ab}$$

$$T_p \equiv \frac{1}{(2\pi)^p g_s (\alpha')^{(p+1)/2}}$$

$$S_{\text{CS}} = i\mu_p \int_{\Sigma_{p+1}} \sum_n C_n \wedge e^{\mathcal{F}}$$

$$S_{\text{D}p} = S_{\text{DBI}} + S_{\text{CS}}$$

$$r_+^{7-p} = d_p g_s N (\alpha')^{\frac{1}{2}(7-p)}$$

$$e^{\Phi} = g_s \left(1 + \left(\frac{r_+}{r} \right)^{7-p} \right)^{\frac{1}{4}(3-p)}$$

$$1 \ll g_s N \ll N$$

$$\mathcal{M}_{10} = \mathcal{M}_4 \times X_6$$

$$G_{MN} \, dX^M \, dX^N = \eta_{\mu\nu} dx^\mu dx^\nu + g_{mn} dy^m dy^n$$

$$G_{MN} \, dX^M \, dX^N = e^{2A(y)} g_{\mu\nu} dx^\mu dx^\nu + e^{-2A(y)} g_{mn} dy^m dy^n,$$

$$\mathcal{L}_{\Phi} = -K_{i\bar{j}} \partial^{\mu} \phi^i \partial_{\mu} \bar{\phi}^{\bar{j}} - V_F,$$

$$V_F(\phi^i, \bar{\phi}^{\bar{i}}) = e^{K/M_{\text{pl}}^2} \left[K^{i\bar{j}} D_i W \overline{D_j W} - \frac{3}{M_{\text{pl}}^2} |W|^2 \right]$$

$$G_{MN} \, dX^M \, dX^N = e^{-6u(x)} g_{\mu\nu} dx^\mu dx^\nu + e^{2u(x)} \hat{g}_{mn} dy^m dy^n,$$

$$\int_{X_6} d^6 y \sqrt{\hat{g}} \equiv \mathcal{V}$$



$$S_{\text{EH}}^{(10)} = \frac{1}{2\kappa^2} \int d^{10}X \sqrt{-G} e^{-2\Phi} R_{10}$$

$$\bar{g}_{MN} = e^{2\omega(x)} g_{MN}$$

$$e^{2\omega}\bar{R}=R-2(D-1)\nabla^2\omega-(D-2)(D-1)g^{MN}\nabla_M\omega\nabla_N\omega.$$

$$e^{2\omega}\bar{\nabla}^2\!=\!\nabla^2\!+\!(D-2)g^{MN}\nabla_M\omega\nabla_N.$$

$$S_{\text{EH}}^{(10)}=\frac{1}{2\kappa^2}\int d^4x\sqrt{-g}\int_{X_6}d^6y\sqrt{\hat{g}}e^{-2\Phi}\big(R_4+e^{-8u}\hat{R}_6+12\partial_\mu u\partial^\mu u\big)$$

$$S_{\text{EH}}^{(4)}=\frac{M_{\text{pl}}^2}{2}\int d^4x\sqrt{-g}R_4$$

$$M_{\rm pl}^2\equiv\frac{\mathcal{V}}{g_s^2\kappa^2}$$

$$K=-3\ln\left(T+\bar{T}\right)$$

Métrica de Moduli - Kähler.

$$J\equiv ig_{i\overline{j}}\,\mathrm{d}z^i\wedge\,\mathrm{d}\overline{z}^{\overline{j}}$$

$$J=t^I(x)\omega_I$$

$$\delta g_{ij}=\frac{i}{\|\Omega\|^2}\zeta^A(x)(\chi_A)_{i\overline{j}}\Omega^{\overline{i}j}{}_j$$

$$B_2=B_2(x)+b^I(x)\omega_I,$$

$$C_2=C_2(x)+c^I(x)\omega_I,$$

$$C_4=\vartheta^I(x)\tilde{\omega}_I.$$

$$\mathcal{O}=(-1)^{F_L}\Omega_{\mathcal{W}\mathcal{S}}\sigma$$

$$H^{1,1}=H^{1,1}_+\oplus H^{1,1}_-$$

$$J=t^i(x)\omega_i$$

$$B_2=b^\alpha(x)\omega_\alpha,$$

$$C_2=c^\alpha(x)\omega_\alpha,$$

$$C_4=\vartheta^i(x)\tilde{\omega}_i.$$

$$\tau = C_0 + i e^{-\Phi}$$

$$G_\alpha \equiv c_\alpha - \tau b_\alpha$$

$$\mathcal{V}=\frac{1}{6}\int_{X_6}J\wedge J\wedge J=\frac{1}{6}c_{ijk}t^it^jt^k,$$

$$T_i\equiv\frac{1}{2}c_{ijk}t^jt^k+i\vartheta_i+\frac{1}{4}e^\Phi c_{i\alpha\beta}G^\alpha(G-\bar{G})^\beta,$$



$$T_i = \frac{1}{2} c_{ijk} t^j t^k + i \vartheta_i$$

$$\tau_i = \frac{\partial \mathcal{V}}{\partial \bar{t}^i} = \frac{1}{2} c_{ijk} t^j t^k,$$

$$T_i = \tau_i + i \vartheta_i$$

Simetrias de Peccei-Quinn (PQ).

$$a \mapsto a + \text{const} \, .$$

$$b_I=\frac{1}{\alpha'}\int_{\Sigma_2^I}B_2, c_I=\frac{1}{\alpha'}\int_{\Sigma_2^I}C_2, \vartheta_I=\frac{1}{(\alpha')^2}\int_{\Sigma_4^I}C_4$$

$$\int_{\Sigma_2^I}\omega^J=\alpha'\delta_I^J, \int_{\Sigma_4^I}\tilde{\omega}^J=(\alpha')^2\delta_I^J$$

$$S_\sigma \supset -\frac{1}{4\pi\alpha'}\int_{\Sigma_2} \mathrm{d}^2\sigma \epsilon^{ab}\partial_a X^M\partial_b X^N B_{MN}(X),$$

$$S_\sigma \supset -\frac{1}{2\pi\alpha'}\int_{\Sigma_2} B_2 \equiv -\frac{b}{2\pi},$$

$$B_{MN}(X)=B_{MN}(X_{(0)})+X^P\partial_PB_{MN}(X_{(0)})+\cdots$$

$$-\frac{1}{4\pi\alpha'}\int_{\Sigma_2}\mathrm{d}^2\sigma\partial_a\left(\epsilon^{ab}X^M\partial_bX^NB_{MN}(X_{(0)})\right)$$

$$S_{\rm inst} = \exp\left(-\frac{1}{2\pi\alpha'}\int_{\Sigma_2}(J+iB_2)\right) \propto \exp\left(-i\frac{b}{2\pi}\right)$$

$$S_{\mathrm{ED}p} = \exp\left(-T_p \mathrm{Vol}(\Sigma_p)\right),$$

$$\mathcal{L}(a)=-\frac{1}{2}f^2(\partial a)^2-\Lambda^4[1-\cos{(a/2\pi)}]+\cdots,$$

$$B_2=b_\alpha(x)\omega^\alpha.$$

$$\frac{1}{2(2\pi)^7g_s^2(\alpha')^4}\int\mathrm{d}^{10}X|\mathrm{d}B_2|^2\supset\frac{1}{2}\int\mathrm{d}^4x\sqrt{-g}\gamma^{\alpha\beta}(\partial^\mu b_\alpha\partial_\mu b_\beta),$$

$$\gamma^{\alpha\beta}\equiv\frac{1}{6(2\pi)^7g_s^2(\alpha')^4}\int_{X_6}\omega^\alpha\wedge\star_6\omega^\beta.$$

$$\frac{f^2}{M_{\rm pl}^2}\approx\frac{1}{6}\frac{\alpha'^2}{L^4}$$



Compactificaciones de Calabi-Yau.

$$ds^2 = e^{2A(y)} \eta_{\mu\nu} dx^\mu dx^\nu + e^{-2A(y)} g_{mn} dy^m dy^n$$

$$\tilde{F}_5 = (1 + \star_{10}) d\alpha(y) \wedge dx^0 \wedge dx^1 \wedge dx^2 \wedge dx^3$$

$$\nabla^2 e^{4A} = \frac{e^{8A}}{2\text{Im}(\tau)} |G_3|^2 + e^{-4A} (|\partial\alpha|^2 + |\partial e^{4A}|^2) + 2\kappa^2 e^{2A} \mathcal{J}_{\text{loc}}$$

$$\mathcal{J}_{\text{loc}} \equiv \frac{1}{4} \left(\sum_{M=4}^9 T^M{}_M - \sum_{M=0}^3 T^M{}_M \right)_{\text{loc}}$$

$$d\tilde{F}_5 = H_3 \wedge F_3 + 2\kappa^2 T_3 \rho_3^{\text{loc}}$$

$$\frac{1}{2\kappa^2 T_3} \int_{X_6} H_3 \wedge F_3 + Q_3^{\text{loc}} = 0$$

$$\nabla^2 (e^{4A} - \alpha) = \frac{e^{8A}}{24\text{Im}(\tau)} |iG_3 - \star_6 G_3|^2 + e^{-4A} |\partial(e^{4A} - \alpha)|^2 + 2\kappa^2 e^{2A} (\mathcal{J}_{\text{loc}} - \mathcal{Q}_{\text{loc}})$$

$$\mathcal{J}_{\text{loc}} \geq \mathcal{Q}_{\text{loc}}$$

$$\star_6 G_3 = iG_3,$$

$$e^{4A} = \alpha$$

$$V_{\text{flux}} = \frac{1}{2\kappa^2} \int d^{10}X \sqrt{-G_E} \left[-\frac{|G_3|^2}{2\text{Im}(\tau)} \right]$$

$$K_0 = -2\ln(\mathcal{V}) - \ln(-i(\tau - \bar{\tau})) - \ln\left(-i \int \Omega \wedge \bar{\Omega}\right)$$

$$W_0 = \frac{c}{\alpha'} \int G_3 \wedge \Omega$$

$$V_F = e^{K_0} \left[K_0^{I\bar{J}} D_I W_0 \overline{D_J W_0} - 3|W_0|^2 \right]$$

$$D_I W_0 \equiv \partial_I W_0 + (\partial_I K) W_0 = 0$$

$$\sum_{I,J=T_i} K_0^{I\bar{J}} \partial_I K_0 \partial_{\bar{J}} K_0 = 3.$$

$$V_F = e^{K_0} \sum_{I,J \neq T_i} K_0^{I\bar{J}} D_I W_0 \overline{D_J W_0}$$

$$K(T, \bar{T}, z_\alpha, \bar{z}_\alpha) = -3\ln [T + \bar{T} - \gamma k(z_\alpha, \bar{z}_\alpha)],$$



Correcciones perturbativas Kaluza-Klein (KK) y singularidad.

$$K=-2\ln\left[\mathcal{V}+\frac{\xi}{2g_s^{3/2}}\right], \xi\equiv-\frac{\chi(X_6)\zeta(3)}{2(2\pi)^3}$$

$$\delta K_{(g_s)}=\delta K_{(g_s)}^{\rm KK}+\delta K_{(g_s)}^{\rm W}$$

$$\delta K_{(g_s)}^{\rm KK}=-\frac{1}{128\pi^2}\sum_{i=1}^3\frac{\mathcal{E}_i^{\rm KK}(\zeta,\bar{\zeta})}{\mathrm{Re}(\tau)\tau_i},$$

$$\delta K_{(g_s)}^{\rm W}=-\frac{1}{128\pi^2}\sum_{i=1}^3\frac{\mathcal{E}_i^{\rm W}(\zeta,\bar{\zeta})}{\tau_j\tau_k}\Big|_{j\neq k\neq i}$$

$$S=\frac{1}{2g_7^2}\int_{\Sigma_4} \mathrm{d}^4\sigma\sqrt{g_{\mathrm{ind}}} \, e^{-4A(y)} \cdot \int \mathrm{d}^4x\sqrt{-g}\mathrm{Tr}[F_{\mu\nu}F^{\mu\nu}]$$

$$g_7^2=2(2\pi)^5(\alpha')^2$$

$$\frac{1}{g^2}=\frac{T_3\mathcal{V}_4}{8\pi^2}\,,$$

$$\mathcal{V}_4\equiv\int_{\Sigma_4}\mathrm{d}^4\sigma\sqrt{g_{\mathrm{ind}}} \, e^{-4A(y)}$$

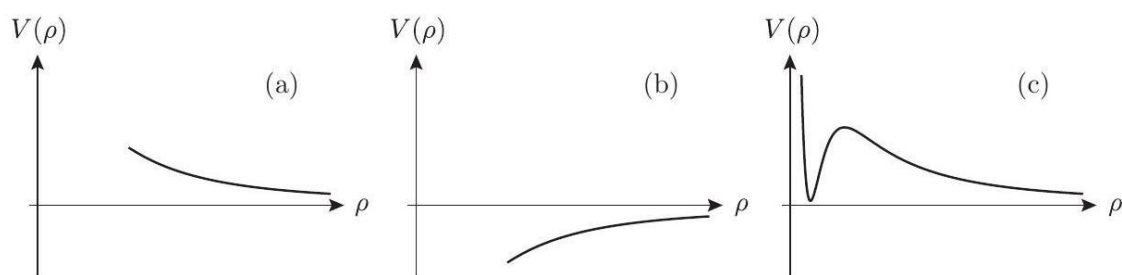
$$|W_{\lambda\lambda}|=16\pi^2M_{\rm UV}^3\mathrm{exp}\left(-\frac{1}{N_c}\frac{8\pi^2}{g^2}\right)\propto\mathrm{exp}\left(-\frac{T_3\mathcal{V}_4}{N_c}\right).$$

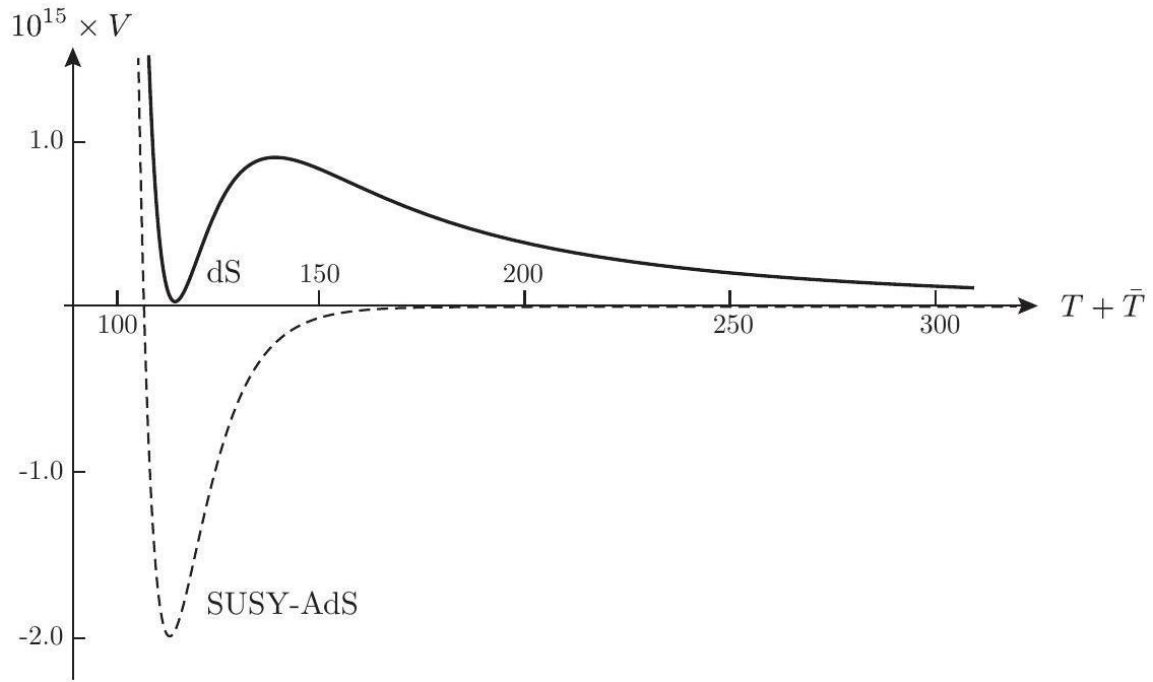
$$W_{\lambda\lambda}=\mathcal{A}e^{-aT}$$

$$W_{\rm ED3}=\mathcal{A}e^{-aT},$$

$$\chi(\mathcal{O}_D)\equiv\sum_{i=0}^3(-1)^ih^{0,i}(D)=1$$

$$h^{0,1}(D)=h^{0,2}(D)=h^{0,3}(D)=0$$





Figuras 39 y 40. Fluctuaciones del centro de masa – energía de la partícula supermasiva.

$$W = W_0 + \sum_{i=1}^{h_+^{1,1}} \mathcal{A}_i e^{-a_i T_i} + \dots$$

$$V_{(np)} = e^K K^{j\bar{i}} \left[a_j \mathcal{A}_j a_{\bar{i}} \overline{\mathcal{A}}_{\bar{i}} e^{-(a_j T_j + a_{\bar{i}} \bar{T}_{\bar{i}})} - (a_j \mathcal{A}_j e^{-a_j T_j} \bar{W} \partial_{\bar{i}} K + a_{\bar{i}} \overline{\mathcal{A}}_{\bar{i}} e^{-a_{\bar{i}} \bar{T}_{\bar{i}}} W \partial_j K) \right]$$

$$V_{(np)} = \frac{a \mathcal{A} e^{-a(T+\bar{T})}}{2(T+\bar{T})^2} \left[\left(1 + \frac{T+\bar{T}}{3} \right) a \mathcal{A} e^{-a(T+\bar{T})} + W_0 \right]$$

$$W_0 = -\mathcal{A} e^{-a(T+\bar{T})_*} \left(1 + \frac{2}{3} a(T+\bar{T})_* \right)$$

$$W_0 = W_0|_{Z=0} + \ell_A Z^A + m_{AB} Z^A Z^B + \dots$$

$$\delta V_{(\alpha')} = 3\hat{\xi} e^K \frac{(\hat{\xi}^2 + 7\hat{\xi}\mathcal{V} + \mathcal{V}^2)}{(\mathcal{V} - \hat{\xi})(2\mathcal{V} + \hat{\xi})^2} W_0^2 \approx \frac{3}{4} \hat{\xi} W_0^2 \frac{1}{\mathcal{V}^3},$$

$$V_{(np)} + \delta V_{(\alpha')} = e^K \left\{ K^{j\bar{i}} \left[a_j \mathcal{A}_j a_{\bar{i}} \overline{\mathcal{A}}_{\bar{i}} e^{-(a_j T_j + a_{\bar{i}} \bar{T}_{\bar{i}})} - (a_j \mathcal{A}_j e^{-a_j T_j} \bar{W} \partial_{\bar{i}} K + a_{\bar{i}} \overline{\mathcal{A}}_{\bar{i}} e^{-a_{\bar{i}} \bar{T}_{\bar{i}}} W \partial_j K) \right] + \frac{3}{4} \hat{\xi} W_0^2 \frac{1}{\mathcal{V}} \right\}$$

$$\mathcal{V} \rightarrow \infty, \text{ with } a_s \tau_s = \ln \mathcal{V}$$

$$\{T_i\} = \{T_b^\rho\} \cup \{T_s^r\},$$

$$\mathcal{V} \rightarrow \infty, \text{ with } a_s^r \tau_s^r = \ln \mathcal{V} \text{ for all } r = 1, \dots, N_s$$

$$\mathcal{V}=\alpha \tau_b^{3/2}-p^{(3/2)}(\tau_s^r),$$

$$\mathcal{V}=\alpha\left(\tau_b^{3/2}-\sum_{r=1}^{N_s}\lambda_r(\tau_s^r)^{3/2}\right)$$

$$\rho_{\mathcal{S}}=\frac{D}{\mathcal{V}^{\alpha}}$$

$$\frac{1}{(2\pi)^2\alpha'}\int_{S^3}F_3\equiv M$$

$$R^4=4\pi g_s N(\alpha')^2,$$

$$\nabla^2\Phi=4\pi e\delta(\boldsymbol{x}),$$

$$\Phi(\boldsymbol{x}')=-\frac{e}{|\boldsymbol{x}-\boldsymbol{x}'|}.$$

$$\mathcal{L}=-\frac{1}{2}(\partial\phi)^2-\frac{1}{4}F_{\mu\nu}F^{\mu\nu}-\frac{\phi}{f}F_{\mu\nu}F_{\rho\sigma}\epsilon^{\mu\nu\rho\sigma}$$

$$\boldsymbol{E}\cdot\boldsymbol{B}=-\frac{eB}{r^2}\cos\theta$$

$$\nabla^2\phi=\frac{eB}{fr^2}\cos\theta$$

$$\mathrm{d}G_{-}=-\mathrm{d}\Big(\frac{\Phi_{-}G_{+}}{\Phi_{+}}\Big)$$

$$V=V_{\mathfrak{F}}(\zeta_1,\ldots,\zeta_{h^{2,1}})$$

$$\mathrm{d}\mathcal{N}_{\min}(\zeta)\equiv\sum_i\delta(\zeta-\zeta_i)$$

$$\mathrm{d}\mathcal{I}_{\min}(\zeta)\equiv\sum_i\delta(\zeta-\zeta_i)(-1)^{F_i}$$

$$\sum_{L\leq L_{\max}}\mathrm{d}\mathcal{I}_{\min}=\frac{(2\pi L_{\max})^{b_3}}{\pi^{b_3/2}b_3!}\det(-\mathcal{R}-\omega)$$

Gravedad cuántica y simetrías de gauge local.

$$\mathcal{N}_{dS} = \mathcal{N}_{\text{c.p.}} \times f_{dS}$$

$$\mathcal{N}_{\text{min}} = \mathcal{N}_{\text{c.p.}} \times f_{\text{min}}$$

$$\mathfrak{C} \equiv \bigcup_{\mathfrak{F}_\star} \left\{ p_i^{(\mathfrak{F}_\star)} \right\},$$

$$\mathcal{N}_{\text{c.p.}} \gtrsim \mathcal{N}_{\mathfrak{F}}$$

$$V=e^K(F_a\bar{F}^a-3|W|^2)$$

$$\mathcal{H} = \begin{pmatrix} \partial_{a\bar{b}}^2 V & \partial_{ab}^2 V \\ \partial_{\bar{a}\bar{b}}^2 V & \partial_{\bar{a}b}^2 V \end{pmatrix}$$

$$f_{\text{min}} = P(\lambda_1 > 0)$$

$$F_a\equiv\mathcal{D}_aW, Z_{ab}\equiv\mathcal{D}_a\mathcal{D}_bW, U_{abc}\equiv\mathcal{D}_a\mathcal{D}_b\mathcal{D}_cW$$

$$\mathcal{H} = \begin{pmatrix} Z_a^{\bar{c}}\bar{Z}_{\bar{b}\bar{c}} - F_a\bar{F}_{\bar{b}} - R_{a\bar{b}c\bar{d}}\bar{F}^cF^{\bar{d}} & U_{abc}\bar{F}^c - Z_{ab}\bar{W} \\ \bar{U}_{\bar{a}\bar{b}c}F^{\bar{c}} - \bar{Z}_{\bar{a}\bar{b}}W & \bar{Z}_{\bar{a}}^cZ_{bc} - F_b\bar{F}_{\bar{a}} - R_{b\bar{a}c\bar{d}}\bar{F}^cF^{\bar{d}} \end{pmatrix} + \mathbb{1}(F^2 - 2|W|^2)$$

$$F \ll m_{\rm susy} \, M_{\rm pl}$$

$$F \gtrsim m_{\rm susy} \, M_{\rm pl}$$

$$F_{\rm rms} \sim Z_{\rm rms} \sim U_{\rm rms}$$

$$M=A+A^\dagger$$

$$\rho(\lambda)=\frac{1}{2\pi N\sigma^2}\sqrt{4N\sigma^2-\lambda^2}$$

$$M=AA^\dagger$$

$$\rho(\lambda)=\frac{1}{2\pi N\sigma^2\lambda}\sqrt{(\eta_+-\lambda)(\lambda-\eta_-)}$$

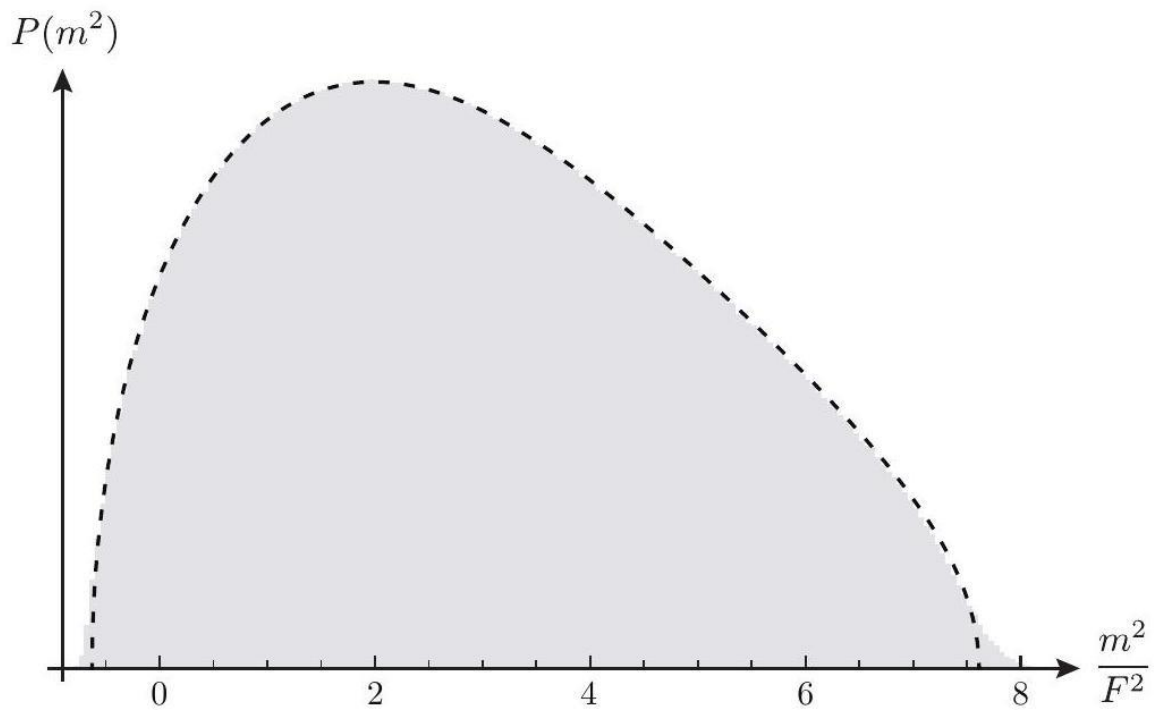
$$\mathcal{H}_Z \equiv \begin{pmatrix} Z_a^{\bar{c}}\bar{Z}_{\bar{b}\bar{c}} & 0 \\ 0 & \bar{Z}_{\bar{a}}^cZ_{bc} \end{pmatrix}$$

$$\mathcal{H}_{WWW} \equiv \mathcal{H}_{\text{Wigner}} + \mathcal{H}_{\text{Wishart}}^{(I)} + \mathcal{H}_{\text{Wishart}}^{(II)},$$

$$\rho(\mathcal{H}_{WWW}) = \rho(\mathcal{H}_{\text{Wigner}}) \boxplus \rho\left(\mathcal{H}_{\text{Wishart}}^{(I)}\right) \boxplus \rho\left(\mathcal{H}_{\text{Wishart}}^{(II)}\right).$$

$$f_{>} \equiv \frac{\int_0^\infty \rho(\lambda)}{\int_{-\infty}^\infty \rho(\lambda)}$$





$$P(\lambda_1 > \zeta) = \exp \left(-N^2 \Psi(\zeta) \right)$$

$$f_{\min} \equiv P(\lambda_1 > 0) = \exp \left(-cN^2 \right),$$

$$f_{\min} \equiv P(\lambda_1 > 0) = \exp \left(-dN \right)$$

$$\begin{aligned} K &= K_l(\phi, \bar{\phi}) + K_h(\Sigma, \bar{\Sigma}) \\ W &= W_l(\phi) + W_h(\Sigma) \end{aligned}$$

$$f_{\min} \equiv P(\lambda_1 > 0) = \exp \left(-cN_l^2 \right),$$

$$S_{10}[\mathcal{C}] \mapsto S_4[\Phi(t)].$$

$$M_{\rm pl} \sim g_s^{-1} (M_{\rm s}/M_{\rm KK})^3 M_{\rm s} \gg M_{\rm s}.$$

$$M_{\rm SUSY} < H < M_{\rm KK} < M_{\rm s} < M_{\rm pl}$$

$$\Delta\mathcal{L} \supset \frac{1}{M_{AB}^{\delta_A+\delta_B-4}} \mathcal{O}_A^{(\delta_A)} \mathcal{O}_B^{(\delta_B)}$$

$$\Delta\mathcal{L} \supset \frac{V_0}{M_{AB}^2} \phi^2$$

$$\frac{M_{AB}}{M_{\rm pl}} \lesssim g_s \left(\frac{\ell_s}{L}\right)^2$$

$$\frac{M_{AB}}{M_{\rm pl}} \lesssim g_{\rm s} \left(\frac{\ell_{\rm s}}{L}\right)^{\frac{1}{2}p-1} \left(\frac{\ell_{\rm s}}{S}\right)^{\frac{1}{2}(6-p)}$$

$$V(r) = 2T_3 \left(1 - \frac{1}{2\pi^3} \frac{T_3 g_s^2 \kappa^2}{r^4} \right),$$

$$\eta \approx -\frac{10}{\pi^3} \frac{\mathcal{V}}{r^6},$$

$$\nabla_y^2(\delta e^{-4A(y_b;y)}) = -\mathcal{C}\left(\frac{\delta(y_b-y)}{\sqrt{g(y)}} - \bar{\rho}(y)\right)$$

$$\delta e^{-4A(y_b;y)} = \mathcal{C}\left(\mathcal{G}(y_b;y) - \int \mathrm{d}^6y' \sqrt{g} \mathcal{G}(y;y') \bar{\rho}(y')\right)$$

$$\nabla_{y'}^2 \mathcal{G}(y;y') = \nabla_y^2 \mathcal{G}(y;y') = -\frac{\delta(y-y')}{\sqrt{g}} + \frac{1}{\mathcal{V}}.$$

$$\nabla_{y_b}^2(\delta e^{-4A(y_b;y)}) = -\mathcal{C}\left(\frac{\delta(y_b-y)}{\sqrt{g(y_b)}} - \frac{1}{\mathcal{V}}\right)$$

$$V(y_b) = 2T_3 e^{4A(y_b)} \approx 2T_3 (1 - \delta e^{-4A(y_b)})$$

$$\mathrm{Tr}(\eta) \approx -\frac{M_{\mathrm{pl}}^2}{T_3} \nabla_{y_b}^2(\delta e^{-4A(y_b;y)}) = -2$$

$$\mathcal{L} = -K_{\varphi\bar{\varphi}}\partial_{\mu}\varphi\partial^{\mu}\bar{\varphi} - e^{K/M_{\mathrm{pl}}^2}\left[K^{\varphi\bar{\varphi}}D_{\varphi}W\overline{D_{\varphi}W} - \frac{3}{M_{\mathrm{pl}}^2}|W|^2\right].$$

$$K=K(0)+K_{\varphi\bar{\varphi}}(0)\varphi\bar{\varphi}+\cdots,$$

$$\mathcal{L} \approx -\partial_{\mu}\phi\partial^{\mu}\bar{\phi} - V(0)\left(1+\frac{\phi\bar{\phi}}{M_{\mathrm{pl}}^2}+\cdots\right),$$

$$m_{\phi}^2 = \frac{V(0)}{M_{\mathrm{pl}}^2} + \cdots = 3H^2 + \cdots \Rightarrow \eta = 1 + \cdots$$

$$K=-3\ln\left[T+\bar{T}-\gamma k(z_{\alpha},\bar{z}_{\alpha})\right]\equiv-2\ln\mathcal{V},$$

$$K=(\phi-\bar{\phi})^2$$

$$\Delta\phi^2<\frac{1}{8\pi}\frac{M_s^2}{g_s}\Big(\frac{L}{\ell_s}\Big)^{p-1}.$$

$$M_{\mathrm{pl}}^2 = \frac{1}{\pi} \frac{M_s^2}{g_s^2} \left(\frac{L}{\ell_s}\right)^6.$$

$$\frac{\Delta\phi^2}{M_{\mathrm{pl}}^2} < \frac{g_s}{8} \Big(\frac{\ell_s}{L}\Big)^{7-p}$$



$$\frac{\Delta\phi^2}{M_{\rm pl}^2} < \frac{g_s}{4\pi} \frac{L}{\ell_s}$$

$$\Delta\Phi=\sqrt{N}\Delta\phi$$

$$\delta M_{\rm pl}^2 \sim \frac{N}{16\pi^2} \Lambda_{\rm UV}^2$$

$$V^{1/4} \ll M_{\rm KK} \ll M_s \ll M_{\rm pl}$$

$$\mathcal{L}_{\rm eff}[\Phi] = \mathcal{L}_l[\Phi] + \sum_i c_i \frac{\mathcal{O}_i[\Phi]}{\Lambda^{\delta_i-4}},$$

Campos multiperturbativos y espacios anti – De Sitter.

$$\mathcal{R}=f(\pi_{\star},\psi_{\star})$$

$$\mu_{p,q}=\frac{1}{2\pi\alpha'}\sqrt{(p-C_0q)^2+e^{-2\Phi}q^2}$$

$$\mathrm{d}s^2=e^{2A(r)}\eta_{\mu\nu}\mathrm{d}x^\mu\mathrm{d}x^\nu+e^{-2A(r)}\big(\mathrm{d}r^2+r^2\,\mathrm{d}\Omega_{S_5}^2\big)$$

$$e^{-4A(r)}=1+\frac{L^4}{r^4},\,\,\text{with}\,\,\frac{L^4}{(\alpha')^2}=4\pi g_s N$$

$$\alpha(r)\equiv (C_4)_{tx^i}=e^{4A(r)}$$

$$\mathrm{d} s_{AdS_5}^2 = \frac{L^2}{r^2} \, \mathrm{d} r^2 + \frac{r^2}{L^2} \eta_{\mu\nu} \mathrm{d} x^\mu \mathrm{d} x^\nu,$$

$$S_{\rm D3}=-T_3\int\,\,\mathrm{d}^4\sigma\sqrt{-\mathrm{det}(G_{ab}^E)}+\mu_3\int\,\,C_4$$

$$\mathcal{L}=-T_3e^{4A(r)}\sqrt{1+e^{-4A(r)}g^{\mu\nu}\partial_\mu r\partial_\nu r}+T_3\alpha(r)$$

$$\mathcal{L}\approx -\frac{1}{2}(\partial\phi)^2-T_3(e^{4A(\phi)}-\alpha(\phi)),$$

$$\sum_{A=1}^4 z_A^2=0$$

$$x\cdot x=\frac{1}{2}\rho^2,y\cdot y=\frac{1}{2}\rho^2,x\cdot y=0$$

$$T^{1,1}=[SU(2)\times SU(2)]/U(1)$$

$$d\Omega_{T^{1,1}}^2\equiv\frac{1}{9}\bigg(\mathrm{d}\psi+\sum_{i=1}^2\cos\theta_i\,\mathrm{d}\phi_i\bigg)^2+\frac{1}{6}\sum_{i=1}^2\big(\,\mathrm{d}\theta_i^2+\sin^2\theta_i\,\mathrm{d}\phi_i^2\big)$$



$$\mathrm{d}s^2 = \mathrm{d}r^2 + r^2 \, \mathrm{d}\Omega_{T^{1,1}}^2$$

$$\mathrm{d}s^2 = k_{\alpha\bar{\beta}} \mathrm{d}z^\alpha \mathrm{d}\bar{z}^\beta$$

$$k(z_\alpha,\bar{z}_\alpha)=\frac{3}{2}\Big(\sum_{A=1}^4|z^A|^2\Big)^{2/3}$$

$$\sum_{A=1}^4 z_A^2 = \varepsilon^2$$

$$\begin{aligned}x\cdot x-y\cdot y&=\varepsilon^2\\x\cdot x+y\cdot y&=\rho^2\end{aligned}$$

$$\mathrm{d}s^2 = e^{2A(r)} \eta_{\mu\nu} \mathrm{d}x^\mu \mathrm{d}x^\nu + e^{-2A(r)} (\mathrm{d}r^2 + r^2 \, \mathrm{d}\Omega_{T^{1,1}}^2),$$

$$e^{-4A(r)}=1+\frac{L^4}{r^4} \text{ with } L^4\equiv \frac{27\pi}{4}g_{\rm s}N(\alpha')^2$$

$$\frac{1}{(2\pi)^2\alpha'}\int_A F_3=M \text{ and } \frac{1}{(2\pi)^2\alpha'}\int_B H_3=K$$

$$\mathrm{d}s^2 = e^{2A(r)} \eta_{\mu\nu} \mathrm{d}x^\mu \mathrm{d}x^\nu + e^{-2A(r)} \mathrm{d}\tilde{s}^2,$$

$$e^{-4A(r)}=\frac{L^4}{r^4}\Big(1+\frac{3g_{\rm s}M}{8\pi K}+\frac{3g_{\rm s}M}{2\pi K}\ln\frac{r}{r_{\rm UV}}\Big)$$

$$L^4\equiv \frac{27\pi}{4}g_{\rm s}N(\alpha')^2, N\equiv MK$$

$$e^{A_{\rm IR}} = \exp\left(-\frac{2\pi K}{3g_{\rm s}M}\right)$$

$$\mathcal{V}_{\mathcal{T}} \equiv \int \, \mathrm{d}\Omega_{T^{1,1}}^2 \int_{r_{\mathrm{IR}}}^{r_{\mathrm{UV}}} r^5 \, \mathrm{d}r e^{-4A(r)} = 2\pi^4 g_{\mathrm{s}} N(\alpha')^2 r_{\mathrm{UV}}^2$$

$$M_{\rm pl}^2 > \frac{N}{4} \frac{r_{\rm UV}^2}{(2\pi^3) g_{\rm s}(\alpha')^2}$$

$$\Delta\phi^2 < T_3 r_{\rm UV}^2 = \frac{r_{\rm UV}^2}{(2\pi^3) g_{\rm s}(\alpha')^2}$$

$$\frac{\Delta\phi}{M_{\rm pl}} \leq \frac{2}{\sqrt{N}}$$

$$V(\phi)=T_3(e^{4A(\phi)}-\alpha(\phi)).$$



$$V_{\mathcal{C}}(\phi) = D_0 \left(1 - \frac{27}{64\pi^2} \frac{D_0}{\phi^4} \right)$$

$$D_0 \equiv 2T_3 e^{4A(r_{\rm IR})}$$

$$V_{\mathcal{R}}(\phi) = \frac{1}{12} R \phi^2.$$

$$V(\phi) = V_{\mathcal{C}}(\phi) + V_{\mathcal{R}}(\phi) + \cdots \approx V_0 + H^2 \phi^2 + \cdots \Rightarrow \eta \approx \frac{2}{3} + \cdots$$

$$V_F = e^K [K^{I\bar{J}} D_I W \overline{D_J W} - 3 |W|^2]$$

$$K=-2\ln\left(\mathcal{V}\right)$$

$$\mathcal{V} = \left(T + \bar{T} - \gamma k(z_\alpha, \bar{z}_\alpha)\right)^{3/2}$$

$$\gamma \equiv \frac{T_3}{6} (T + \bar{T})_{\rm IR}$$

$$K(Z^I,\bar{Z}^I)=-3\ln\left[T+\bar{T}-\gamma k(z_\alpha,\bar{z}_\alpha)\right]\equiv-3\ln\left[U(Z^I,\bar{Z}^I)\right].$$

$$V_F(T,z_\alpha)=\frac{1}{3U^2}\big[(T+\bar{T}+\gamma(k_\gamma k^{\gamma\bar{\delta}}k_{\bar{\delta}}-k))\big]|W_{,T}|^2-3(\bar{W}W_{,T}+c.c.)$$

$$+ \underbrace{\left(k^{\alpha\bar{\delta}}k_{\bar{\delta}}W_{,T}W_{,\alpha}+c.c.\right)+\frac{k^{\alpha\bar{\beta}}}{\gamma}W_{,\alpha}W_{,\beta}}_{\Delta\hat{V}_F}$$

$$V_F(r) \approx \frac{V_0}{\left(1-\frac{1}{6}\phi^2\right)^2} \approx V_0 + \frac{1}{3}\frac{V_0}{M_{\rm pl}^2}\phi^2,$$

$$|\Delta W| \propto \exp\left(-\frac{2\pi}{N_c}\mathcal{V}_4\right)$$

$$f(z_\alpha)=0$$

$$W(T,z_\alpha)=W_0+\mathcal{A}(z_\alpha)e^{-aT}, a\equiv \frac{2\pi}{N_c}$$

$$\mathcal{A}(z_\alpha)=\mathcal{A}_0\left(\frac{f(z_\alpha)}{f(0)}\right)^{1/N_c}$$

$$f(z_1)=\mu-z_1$$

$$V_F(\phi) \approx V_0 + \cdots + \lambda \phi^{3/2} + \frac{V_0}{M_{\rm pl}^2} \phi^2 + \cdots$$



$$\delta\Phi(r)=\delta\Phi(r_{\rm UV})\left(\frac{r}{r_{\rm UV}}\right)^\Delta.$$

$$\Phi_- \equiv e^{4A} - \alpha$$

$${\rm d} s^2 = e^{2A(y)} g_{\mu\nu} {\rm d} x^\mu {\rm d} x^\nu + e^{-2A(y)} g_{mn} \,{\rm d} y^m \,{\rm d} y^n,$$

$$\nabla^2\Phi_- = R_4 + \frac{g_s}{96}|\Lambda|^2 + e^{-4A}|\nabla\Phi_-|^2 + \mathcal{S}_{\rm loc}$$

$$\Lambda \equiv \Phi_+ G_- + \Phi_- G_+$$

$$G_{\pm} \equiv (\star_6 \pm i) G_3 \text{ and } \Phi_{\pm} \equiv e^{4A} \pm \alpha$$

$${\rm d}\Lambda + \frac{i}{2}\frac{{\rm d}\tau}{{\rm Im}\tau}\wedge (\Lambda + \bar{\Lambda}) = 0$$

$$V(x,\Psi)=V_0+V_{\mathcal{C}}(x)+V_{\mathcal{R}}(x)+V_{\mathcal{B}}(x,\Psi),$$

$$V_{\mathcal{C}}(x)=D_0\left(1-\frac{27}{64\pi^2}\frac{D_0}{T_3^2r_{\rm UV}^4}\frac{1}{x^4}\right)$$

$$V_{\mathcal{R}}(x)=\frac{1}{3}\mu^4x^2+\cdots, \text{ where } \mu^4\equiv (V_0+D_0)\frac{T_3r_{\rm UV}^2}{M_{\rm pl}^2}$$

$$V_{\mathcal{B}}(x,\Psi)=\mu^4\sum_{LM}c_{LM}x^{\Delta(L)}f_{LM}(\Psi)$$

$$\nabla^2\Phi_h=0$$

$$\nabla^2\Phi_f=\frac{g_s}{96}|\Lambda|^2$$

$$\Delta_h(L)\equiv -2\sqrt{H(j_1,j_2,R)+4}$$

$$H(j_1,j_2,R)\equiv 6\left[j_1(j_1+1)+j_2(j_2+1)-\frac{1}{8}R^2\right]$$

$$\Delta_h=\frac{3}{2},2,3,\sqrt{28}-2,\cdots$$

$$\Delta_f(L)=\delta_i(L)+\delta_j(L)-4$$

$$\delta_1(L)\equiv -1+\sqrt{H(j_1,j_2,R+2)+4},$$

$$\delta_2(L)\equiv \sqrt{H(j_1,j_2,R)+4},$$

$$\delta_3(L)\equiv 1+\sqrt{H(j_1,j_2,R-2)+4}.$$

$$\Delta_f=1,2,\frac{5}{2},\sqrt{28}-\frac{5}{2},\cdots$$

$$\Delta=\{\Delta_h,\Delta_f\}=1,\frac{3}{2},\sqrt{28}-\frac{5}{2},3,\sqrt{28}-2,\frac{7}{2},\sqrt{28}-\frac{3}{2},\cdots$$



$$\mathcal{M} = \begin{pmatrix} A\bar{A} + B\bar{B} & C \\ \bar{C} & \bar{A}A + \bar{B}B \end{pmatrix}$$

$$\tau_{\rm therm} \ll \tau_{\rm graviton} \ll \tau_{\rm tunnel} \, ,$$

$$V(\phi) \approx V_0 \left[1 + \lambda_0 \frac{\phi}{M_{\rm pl}} + \frac{1}{3!} \mu_0 \frac{\phi^3}{M_{\rm pl}^3} + \cdots \right],$$

$$n_s-1 \approx -\frac{4\pi}{N_{\rm tot}} \cot\left(\pi\frac{N_\star}{N_{\rm tot}}\right) \approx -\frac{4}{N_\star} \left(1+\mathcal{O}\left(\frac{N_\star^2}{N_{\rm tot}^2}\right)\right),$$

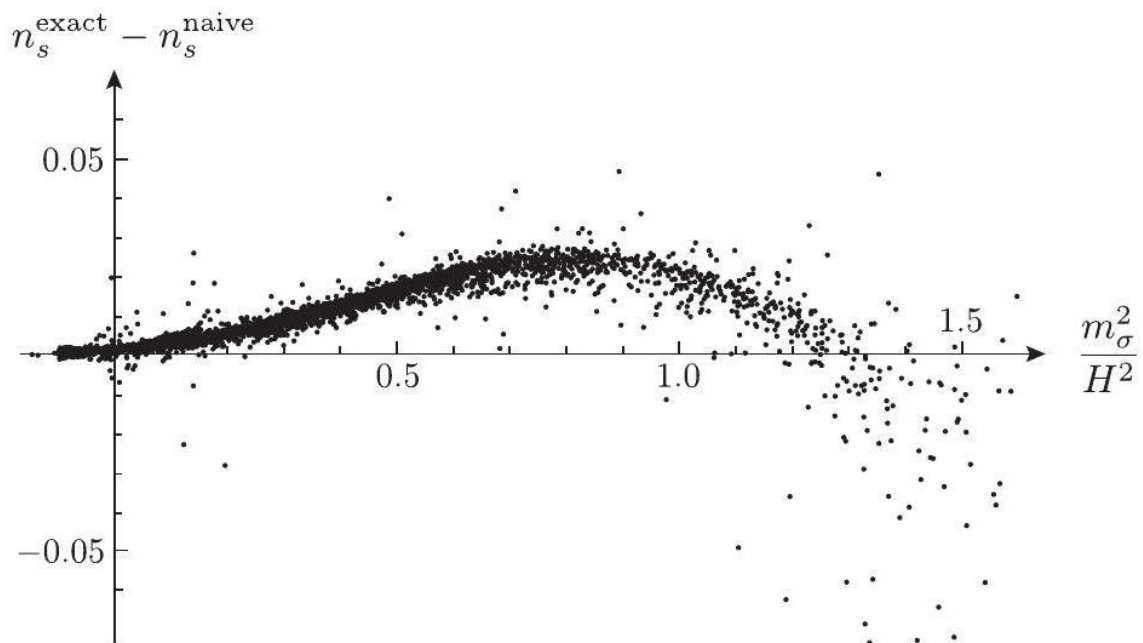
$$N_{\rm tot} = \int_{-\infty}^{\infty} \frac{1}{\sqrt{2\epsilon}} \frac{\mathrm{d}\phi}{M_{\rm pl}} = \pi \sqrt{\frac{2}{\lambda_0 \lambda_1}}.$$

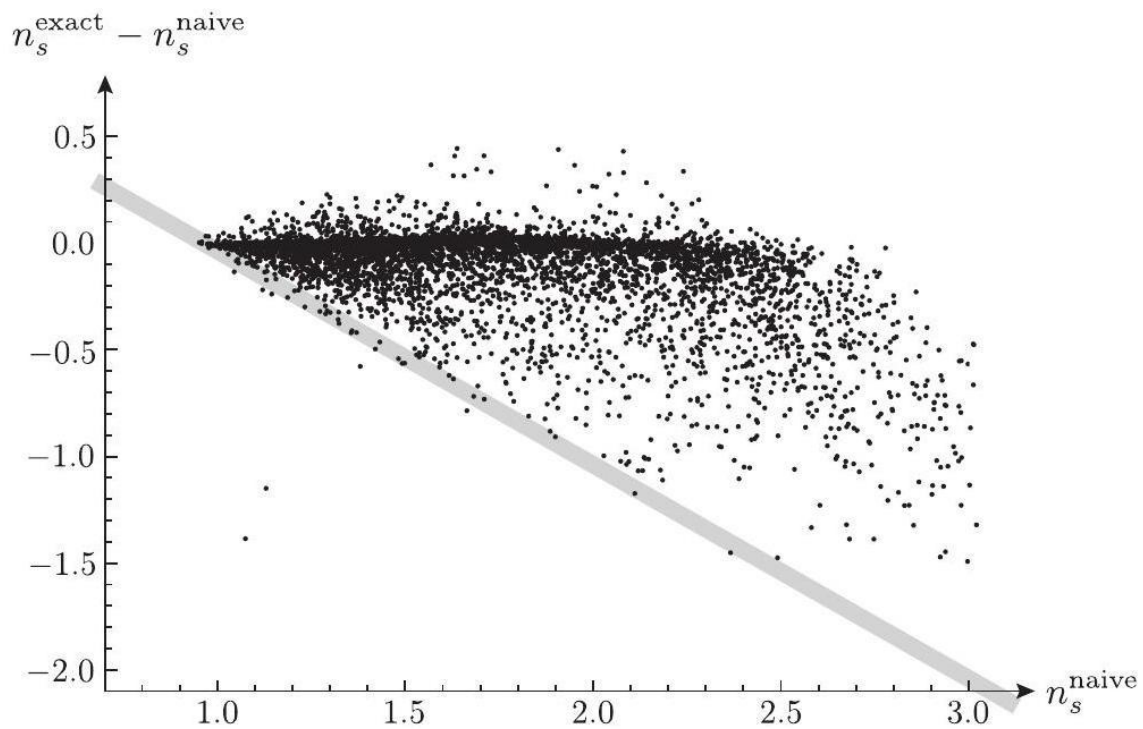
$$N_e(\phi) = \int_{\phi_{\rm end}}^{\phi} \frac{1}{\sqrt{2\epsilon}} \frac{\mathrm{d}\phi}{M_{\rm pl}} = \frac{N_{\rm tot}}{\pi} \arctan\left(\frac{\eta(\phi)N_{\rm tot}}{2\pi}\right) \Big|_{\phi_{\rm end}}^{\phi}$$

$$\alpha_s = -\frac{4\pi^2}{N_{\rm tot}^2} \sin^{-2}\left(\pi\frac{N_\star}{N_{\rm tot}}\right) \approx -\frac{4}{N_\star^2} \left(1+\mathcal{O}\left(\frac{N_\star^2}{N_{\rm tot}^2}\right)\right).$$

$$r < \frac{4}{N} \times 0.01 \ll 0.01$$

$$P(N_e) = P(N_\star) \left(\frac{N_\star}{N_e}\right)^3$$





$$T_{(p,q)} \approx \frac{e^{2A_{\text{IR}}}}{2\pi\alpha'} \sqrt{\frac{q^2}{g_s^2} + \left(\frac{bM}{\pi}\right)^2 \sin^2\left(\frac{\pi(p-qC_0)}{M}\right)}$$

$$V_D(\phi)=\frac{g^2\xi^2}{2}\bigg(1+\frac{g^2}{16\pi^2}U(x)\bigg),$$

$$U(x)\equiv (x^2+1)^2\ln{(x^2+1)}+(x^2-1)^2\ln{(x^2-1)}-4x^4\ln{(x)}-4\ln{(2)},$$

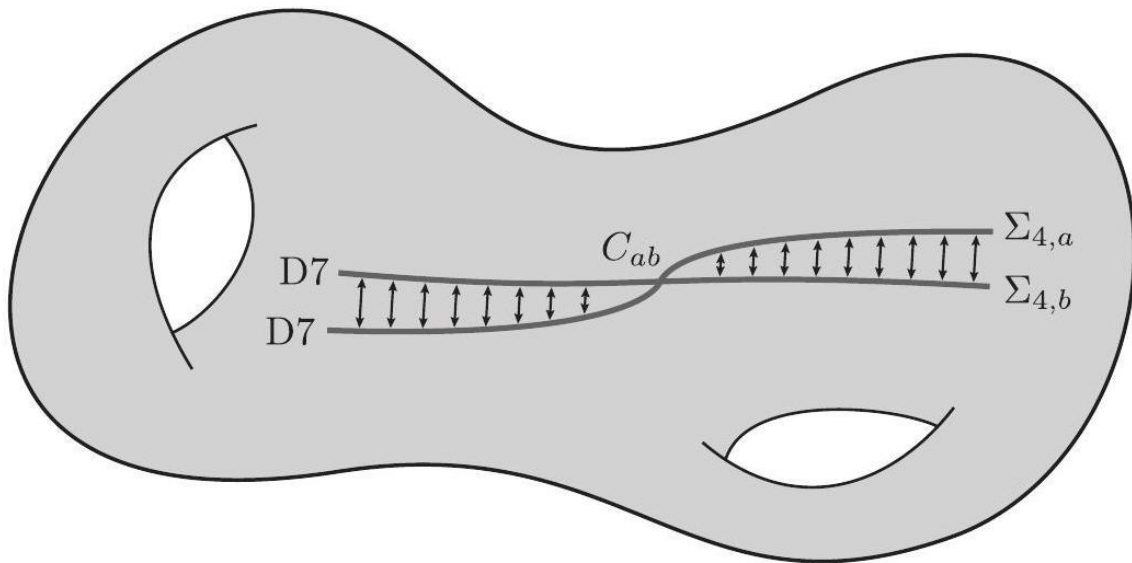
$$K(Z^I,\bar{Z}^I)=-3\ln\left[T+\bar{T}-\gamma k(z_\alpha,\bar{z}_\alpha)\right]$$

$$k=\frac{1}{2}(z_1+\bar{z}_1)^2=x^2,$$

$$W=W_0+\left[\vartheta_1(\sqrt{2\pi}(z_1+\mu),\zeta)\vartheta_1(\sqrt{2\pi}(z_1-\mu),\zeta)\right]^{-1/N_c}e^{-2\pi T/N_c}$$

$$W=\Lambda^{3-1/N_c}m^{1/N_c},$$

$$V(\phi)=V_D(\phi)-\frac{m^2}{2}\phi^2+\frac{\lambda}{4}\phi^4$$



$$V(\phi) = V_D(\phi) + V_F(\phi)$$

$$V_D(\phi) = V_0(1 + \alpha \ln(\phi/\phi_0))$$

$$K \supset -\ln \left(S + \bar{S} - \frac{(Y + \bar{Y})^2}{T + \bar{T}} \right)$$

$$M_{\rm pl} \propto \sqrt{{\rm Vol}(K3) L_1 L_2},$$

$$\Delta\phi_1\propto\sqrt{L_1/L_2}$$

$$\mathrm{d}s^2=\left(\frac{r}{L}\right)^2\eta_{\mu\nu}\mathrm{d}x^\mu\mathrm{d}x^\nu+\left(\frac{L}{r}\right)^2\left(\mathrm{d}r^2+r^2\,\mathrm{d}\Omega_{S^5}^2\right),$$

$$\mathcal{L} = -\frac{\phi^4}{\lambda}\left(\sqrt{1+\frac{\lambda}{\phi^4}(\partial\phi)^2}-1\right)-V(\phi),$$

$$\lambda \equiv T_3 L^4 = \frac{1}{2\pi^2} N$$

$$\mathcal{L} = -T(\phi)\left(\sqrt{1+\frac{(\partial\phi)^2}{T(\phi)}}-1\right)-V(\phi),$$

$$\mathrm{d}s^2=e^{2A(r)}\eta_{\mu\nu}\mathrm{d}x^\mu\mathrm{d}x^\nu+e^{-2A(r)}(\mathrm{d}r^2+r^2\,\mathrm{d}\Omega_{X_5}^2),$$

$$e^{-4A(r)}\approx \frac{L^4}{r^4} \text{ with } L^4\equiv \frac{4\pi^4 g_s}{\mathrm{Vol}(X_5)}N(\alpha')^2$$

$$\frac{\Delta\phi}{M_{\rm pl}} \leq \frac{2}{\sqrt{N}}$$

$$V(\phi) = V_0 - \frac{1}{2} \beta H^2 \phi^2,$$

$$\gamma \equiv \left(1 - \frac{\dot{\phi}^2}{T(\phi)}\right)^{-1/2}$$

$$\dot{\phi}^2 < T(\phi).$$

$$\rho \,=\, (\gamma-1)T + V,$$

$$P \,=\, (\gamma-1)\frac{T}{\gamma} - V.$$

$$3M_{\rm pl}^2H^2=(\gamma-1)T+V(\phi)$$

$$\dot{\phi}=-\frac{2M_{\rm pl}^2H'}{\gamma},$$

$$\gamma=\sqrt{1+\frac{(2M_{\rm pl}^2H')^2}{T}}$$

$$\varepsilon=-\frac{\dot{H}}{H^2}=\frac{2M_{\rm pl}^2}{\gamma}\left(\frac{H'}{H}\right)^2$$

$$\tilde{\eta}=\frac{\dot{\varepsilon}}{H\varepsilon}=\frac{2M_{\rm pl}^2}{\gamma}\left[2\left(\frac{H'}{H}\right)^2-2\frac{H''}{H}+\frac{H'}{H}\frac{\gamma'}{\gamma}\right]$$

$$\frac{V}{\gamma T} \gg 1$$

$$\gamma^2=\frac{2}{3}\epsilon\frac{V}{T}\gg1$$

dilatation : x^μ

$$\phi(x)\mapsto \tilde{x}^\mu\equiv (1+c)x^\mu$$

$$\text{SCTs}:\;\mapsto \phi(\tilde{x})+c$$

$$x^\mu\mapsto \tilde{x}^\mu\equiv x^\mu+(b\cdot x)x^\mu-\frac{1}{2}\Big(x^2+\frac{\lambda}{\phi^2}\Big)b^\mu$$

$$\phi(x)\mapsto \phi(\tilde{x})(1-(b\cdot x))$$

$$S=\int\,\mathrm{d}^4x\phi^4f((\partial\phi)^2/\phi^4)+\cdots$$

$$f(z)=\alpha[\sqrt{1+\lambda z}-\beta]$$



$$V_c(\phi) = D_0 \left(1 - \frac{\pi}{4 \text{Vol}(X_5)} \frac{D_0}{\phi^4} \right).$$

$$\frac{V(\phi)}{\gamma T(\phi)} = \frac{\lambda}{2\gamma} \frac{m^2}{\phi^2}$$

$$M_{\rm KK} \sim \frac{1}{L} e^{A_{\rm IR}} \sim \frac{r_{\rm IR}}{L^2},$$

$$r_{\rm IR} \sim m L^2$$

$$\phi^2 \gtrsim \phi_{\rm IR}^2 \sim T_3 (m L^2)^2 = \lambda m^2.$$

$$\frac{V(\phi)}{\gamma T(\phi)} \lesssim \frac{1}{2\gamma} \ll 1$$

$$\frac{V(\phi)}{\gamma T(\phi)} = \frac{1}{\gamma} \frac{e^{4A} - \alpha}{e^{4A}}$$

$$|\alpha| \gg \gamma e^{4A}$$

$$N_e^{\rm EFT} \approx \frac{N^{1/8}}{\sqrt{\beta}}$$

$$S = \int \, \mathrm{d}^4x \sqrt{-g} \left[\frac{M_{\rm pl}^2}{2} R + P(X,\phi) \right]$$

$$P(X,\phi)=-T(\phi)\left(\sqrt{1-\frac{2X}{T(\phi)}}-1\right)-V(\phi).$$

$$c_s^2 = \frac{P_{,X}}{P_{,X} + 2XP_{,XX}} = \frac{1}{\gamma^2(\phi)}$$

$$c_s^2 = 1 - \frac{2X}{T(\phi)}.$$

$$r=16c_s\varepsilon$$

$$\frac{\Delta\phi}{M_{\rm pl}} = \int_0^{N_\star} \sqrt{\frac{r(N)}{8}} \frac{1}{c_s P_{,X}} \, \mathrm{d}N$$

$$c_s P_{,X} = 1.$$

$$A=c_s^2\left(-1-\frac{2}{3}\frac{XP_{,XXX}}{P_{,XX}}\right).$$

$$A=-1.$$



$$f_{\mathrm{NL}}^{\mathrm{equil}} = (-0.27 + 0.08A) \frac{1 - c_s^2}{c_s^2},$$

$$f_{\mathrm{NL}}^{\mathrm{ortho}} = (+0.02 + 0.02A) \frac{1 - c_s^2}{c_s^2}.$$

$$f_{\mathrm{NL}}^{\mathrm{equil}} = -\frac{35}{108} \left(\frac{1}{c_s^2} - 1 \right) \approx -\frac{35}{108} \gamma^2,$$

$$\gamma \lesssim 24 \text{ (95\% C.L.)}$$

$$2\left(\frac{M_{\mathrm{pl}}}{\phi}\right)^2 = \varepsilon \gamma = \frac{1}{16} r \gamma^2,$$

$$2\left(\frac{M_{\mathrm{pl}}}{\phi}\right)^2 = \frac{27}{140} r \left| f_{\mathrm{NL}}^{\mathrm{equil}} \right|$$

$$N < \frac{27}{70} r \left| f_{\mathrm{NL}}^{\mathrm{equil}} \right| \lesssim 9 \text{ (95\% C.L.)}$$

$$\Delta_{\mathcal{R}}^2(k_{\star}) = \left(\frac{32}{3\pi}\right)^2 \frac{3}{r^4 \left(f_{\mathrm{NL}}^{\mathrm{equil}}\right)^2} \frac{\mathrm{Vol}(X_5)}{N} \gtrsim 15 \frac{\mathrm{Vol}(X_5)}{N},$$

$$N \gtrsim 7 \times 10^9 \mathrm{Vol}(X_5)$$

$$f_{\mathrm{NL}}^{\mathrm{equil}} \approx -\frac{35}{108} \gamma^2 \cos^2 \theta$$

$$\mathcal{L}(\phi) = -\frac{1}{2}(\partial\phi)^2 - \Lambda^4 \left[1 - \cos\left(\frac{\phi}{f}\right) \right] + \cdots$$

$$\mathcal{L} \supset \sum_{i=1}^2 \frac{\phi_i}{f_i} \Big(\frac{c_{ia}}{32\pi^2} \mathrm{Tr}[F^{(a)} \wedge F^{(a)}] + \frac{c_{ib}}{32\pi^2} \mathrm{Tr}[F^{(b)} \wedge F^{(b)}] \Big)$$

$$V = \Lambda_a^4 \left[1 - \cos \left(c_{1a} \frac{\phi_1}{f_1} + c_{2a} \frac{\phi_2}{f_2} \right) \right] + \Lambda_b^4 \left[1 - \cos \left(c_{1b} \frac{\phi_1}{f_1} + c_{2b} \frac{\phi_2}{f_2} \right) \right].$$

$$\frac{c_{1a}}{c_{2a}} = \frac{c_{1b}}{c_{2b}}$$

$$f_{\xi} = \frac{(c_{2a}^2 f_1^2 + f_2^2)^{1/2}}{|c_{2b} - c_{2a}|}, \text{ with } \xi = \frac{\phi_2 f_2 - c_{2a} \phi_1 f_1}{c_{2a}^2 f_1^2 + f_2^2}$$

$$\mathcal{L} = \sum_{i=1}^N \left[-\frac{1}{2} (\partial \phi_i)^2 - V_i(\phi_i) \right],$$

$$\ddot{\phi}_i + 3H\dot{\phi}_i = -\partial_i V_i, \text{ where } 3M_{\mathrm{pl}}^2 H^2 \approx \sum_{i=1}^N V_i.$$

$$V(\Phi)=\frac{1}{2}m^2\Phi^2$$

$$\delta M_{\rm pl}^2 \sim \frac{N}{16\pi^2} \Lambda_{\rm UV}^2$$

$$W=W_0+\sum_{i=1}^N\mathcal{A}_ie^{-a_iT_i}$$

$$K=-2\mathrm{ln}\left[\mathcal{V}(T_i,\bar{T}_i)\right],$$

$$\mathcal{L}=-\frac{1}{2}M_{\rm pl}^2K_{ij}\partial_\mu\vartheta^i\partial^\mu\vartheta^j-\frac{1}{2}M_{ij}\vartheta^i\vartheta^j+\cdots$$

$$\mathcal{L}=\sum_{i=1}^N\left[-\frac{1}{2}(\partial\phi_i)^2-\frac{1}{2}m_i^2\phi_i^2\right]$$

$$\rho(m^2)=\frac{1}{2\pi N\sigma^2m^2}\sqrt{(\eta_+-m^2)(m^2-\eta_-)}.$$

$$\mathcal{L}=\frac{M_{\rm pl}^2}{2}\bigg(1+\frac{\zeta(3)\chi(X_6)}{(2\pi)^3}\frac{(\alpha')^3}{\mathcal{V}}+\cdots\bigg)R_4,$$

$$\chi(X_6)=2h^{1,1}-2h^{2,1},$$

$$S_{\rm D5}=\frac{1}{(2\pi)^5g_s(\alpha')^3}\int_{\mathcal{M}_4\times\Sigma_2}\mathrm{d}^6\sigma\sqrt{-\mathrm{det}(G_{ab}+B_{ab})}$$

$$V(b)=\frac{\varrho}{(2\pi)^6g_s(\alpha')^2}\sqrt{(2\pi)^2\ell^4+b^2}$$

$$V(c)=\frac{\varrho}{(2\pi)^6g_s^2(\alpha')^2}\sqrt{(2\pi)^2\ell^4+g_s^2c^2}$$

$$\mathcal{L}=-\frac{1}{2}(\partial\phi)^2-\mu^3\phi, \text{ with } \mu^3\equiv\frac{1}{f}\frac{\varrho}{(2\pi)^6g_s(\alpha')^2},$$

$$\mathcal{L}=-\frac{1}{2}(\partial\phi)^2-\mu^{4-p}\phi^p$$

$$S\supset\int\,\,\mathrm{d}^{10}X|C_2\wedge H_3|^2$$

$$S_{\rm CS}\supset i\mu_5\int_{\mathcal{M}_4\times\Sigma_2}C_4\wedge\mathcal{F}_2$$

$$S_{\rm CS}^{(\mathrm{D}3)}=i\mu_3\int_{\mathcal{M}_4}C_4,$$

$$\int_{\Sigma_2} \mathcal{F}_2 \neq 0$$

$$\begin{aligned} W &= W_0 + \mathcal{A} e^{-2\pi T} \\ K &= -3\ln\left(T + \bar{T} + \gamma b^2\right) \end{aligned}$$

$$S\supset\int_{\Sigma_4}C_2\wedge\mathcal{F}_2$$

$$W_{\lambda\lambda}=\mathcal{A}e^{-f/N_c},$$

$$f=f_0+f_1+f_{np}$$

$$W_{\lambda\lambda}=\mathcal{A}e^{-(f_0+f_1)/N_c}[1+\mathcal{O}(e^{-S})g(c)],$$

$$V(c)=\mu^3fc+\hat{V}[T_i(c)]$$

$$\mathrm{Re}(T_i)=\int_{\Sigma_4^i}\mathrm{d}^4y\sqrt{g}e^{-4A(y;c)}$$

$$\mathcal{L}=-\frac{1}{2}(\partial\phi)^2-V_0(\phi)-\Lambda^4\cos\left(\frac{\phi}{f}\right)$$

$$b_\star \equiv \frac{\Lambda^4}{V_0'(\phi_\star)f} < 1$$

$$\frac{f^2}{M_{\rm pl}^2} > \frac{\sqrt{g_s}}{(2\pi)^3 \mathcal{V}}$$

$$f\gg (2\Delta_{\mathcal{R}})^{-1/2}\frac{H}{2\pi}$$

$$\frac{f}{M_{\rm pl}} \gg \sqrt{\frac{r}{16}} \Delta_{\mathcal{R}}^{1/2} \approx 4.5 \times 10^{-4} \left(\frac{r}{0.07}\right)^{1/2}$$

$$\Delta_{\mathcal{R}}^2(k)=\Delta_{\mathcal{R}}^2(k_\star)\left(\frac{k}{k_\star}\right)^{n_s-1}\left[1+\mathcal{A}\cos\left(\frac{\phi_k}{f}\right)\right]$$

$$\mathcal{A} \equiv 3b_\star \left(\frac{2\pi f}{\sqrt{2\epsilon_\star}}\right)^{1/2}$$

$$\mathcal{A}_{p=1}=\frac{6b_\star}{\sqrt{1+(3f\phi_\star)^2}}\sqrt{\frac{\pi}{2}\coth\left(\frac{\pi}{2f\phi_\star}\right)f\phi_\star}$$

$$B_{\mathcal{R}}(k_1,k_2,k_3)=f_{\rm NL}^{\rm res}\times\frac{(2\pi\Delta_{\mathcal{R}})^4}{k_1^2k_2^2k_3^2}\Bigg[\sin\left(\frac{\sqrt{2\epsilon_\star}}{f}\ln\frac{K}{k_\star}\right)+\frac{f}{\sqrt{2\epsilon_\star}}\sum_{i\neq j}\cos\left(\frac{\sqrt{2\epsilon_\star}}{f}\ln\frac{K}{k_\star}\right)+\cdots\Bigg]$$



$$f_{\rm NL}^{\rm res} \equiv \frac{3\sqrt{2\pi}}{8} b_\star \left(\frac{\sqrt{2\epsilon_\star}}{f} \right)^{3/2}$$

$$f_{\rm NL}^{\rm res} \ll \frac{3\sqrt{\pi}}{2} (2\Delta_{\mathcal{R}})^{-3/4} \approx 3 \times 10^3$$

$$\mathcal{L} \supset -\frac{\alpha}{4} \frac{\phi}{f} F \tilde{F}$$

$$\left(\frac{\partial^2}{\partial \tau^2} + k^2 \mp 2aHk\xi\right) A_{\pm}(\tau,k) = 0, \text{ where } \xi \equiv \frac{\alpha\dot{\phi}}{2fH}.$$

$$A_+(\tau,k)\simeq \frac{1}{\sqrt{2k}}\Big(\frac{k}{2\xi aH}\Big)^{1/4}e^{\pi\xi-2\sqrt{2\xi k/(aH)}}$$

$$f_{\rm NL}^{\rm equil} \simeq \frac{\Delta_{\mathcal{R},0}^6}{\Delta_{\mathcal{R}}^4} f_3(\xi) e^{6\pi\xi}$$

$$\begin{aligned} f_3(\xi) &= 2.8 \times 10^{-7} \xi^{-9} \text{ for } \xi \gg 1 \\ f_3(\xi) &\approx 7.4 \times 10^{-8} \xi^{-8.1} \text{ for } 2 < \xi < 3 \end{aligned}$$

$$f > \frac{\alpha}{10\pi} \frac{H}{\Delta_{\mathcal{R}}}.$$

$$\frac{f}{M_{\rm pl}} > \frac{\alpha}{10} \sqrt{\frac{r}{2}} \approx 2 \times 10^{-2} \left(\frac{\alpha}{1}\right) \left(\frac{r}{0.07}\right)^{1/2}.$$

$$\Delta_{\mathcal{R}}^2(k) = \Delta_{\mathcal{R},0}^2(k) [1 + \Delta_{\mathcal{R},0}^2(k) f_2(\xi) e^{4\pi\xi}]$$

$$\begin{aligned} f_2(\xi) &= 7.5 \times 10^{-5} \xi^{-6} \text{ for } \xi \gg 1 \\ f_2(\xi) &\approx 3.0 \times 10^{-5} \xi^{-5.4} \text{ for } 2 < \xi < 3 \end{aligned}$$

$$W=W_0+\mathcal{A}e^{-aT}+\mathcal{B}e^{-bT},$$

$$K=-3\ln\left(T+\bar{T}\right),$$

$$\delta V = \frac{\varrho}{(T+\bar{T})^2}$$

$$K=K_0+\delta K_{(\alpha')}=-2\ln\left(\mathcal{V}\right)-\frac{\hat{\xi}}{\mathcal{V}}$$

$$K=K_0+\delta K_{(\alpha')}+\delta K_{(g_s)}$$

$$\delta K_{(g_s)}^{\rm KK} \sim g_s \sum_{i=1}^{h^{1,1}} \frac{c_i^{\rm KK}(\zeta,\bar{\zeta}) M_{\rm KK}^{-2}}{\mathcal{V}} \sim g_s \sum_{i=1}^{h^{1,1}} \frac{c_i^{\rm KK}(\zeta,\bar{\zeta}) (a_{ij} t^j)}{\mathcal{V}},$$

$$\delta K_{(g_s)}^{\rm W} \sim \sum_i \frac{c_i^{\rm W}(\zeta,\bar{\zeta}) M_W^{-2}}{\mathcal{V}} \sim \sum_i \frac{c_i^{\rm W}(\zeta,\bar{\zeta})}{(b_{ij} t^j) \mathcal{V}},$$

$$\delta K_{(g_s)}^{\text{KK}} \sim \sum_i \frac{t_i}{\mathcal{V}} > \delta K_{(\alpha')} \sim \frac{1}{\mathcal{V}}.$$

$$K=-2\ln\left(\mathcal{V}\right)-\frac{\hat{\xi}}{\mathcal{V}}+\frac{\sqrt{\tau}}{\mathcal{V}}$$

$$V=\frac{W_0^2}{\mathcal{V}^3}\Big[0+\hat{\xi}+0\cdot\sqrt{\tau}+\frac{1}{\sqrt{\tau}}+\frac{1}{\tau^{3/2}}\Big].$$

$$\delta V_{\text{CW}} \simeq 0 \cdot \Lambda^4 + \Lambda^2 \, \text{STr}(M^2) + \text{STr} \left[M^4 \ln \left(\frac{M^2}{\Lambda^2} \right) \right]$$

$$\Lambda=M_{\text{KK}}\simeq \frac{M_{\text{pl}}}{\mathcal{V}^{2/3}}\,,\,\,\text{STr}(M^2)\simeq \frac{M_{\text{pl}}^2}{\mathcal{V}^2}\,,$$

$$\delta V_{\text{CW}} \simeq 0 \cdot \frac{1}{\mathcal{V}^{8/3}} + \frac{1}{\mathcal{V}^{10/3}} + \frac{1}{\mathcal{V}^4}$$

$$\delta V_{\text{CW}} \simeq \frac{1}{\mathcal{V}^3} \Big[0 \cdot \sqrt{\tau} + \frac{1}{\sqrt{\tau}} + \frac{1}{\tau^{3/2}} \Big]$$

$$\mathcal{V}=\alpha\left(\tau_b^{3/2}-\lambda_\phi\tau_\phi^{3/2}-\lambda_s\tau_s^{3/2}\right)$$

$$W=W_0+\mathcal{A}_\phi e^{-a_\phi T\phi}+\mathcal{A}_s e^{-a_s T s}$$

$$V=W_0^2\left(\mathfrak{a}\frac{\sqrt{\tau_\phi}e^{-2a_\phi\tau_\phi}}{\mathcal{V}}-\mathfrak{b}\frac{\tau_\phi e^{-a_\phi\tau_\phi}}{\mathcal{V}^2}+\mathfrak{c}\frac{\hat{\xi}}{\mathcal{V}^3}\right),$$

$$V(\phi) \simeq V_0 \big(1 - c_1 \mathcal{V}^{5/3} \phi^{4/3} \exp\big[-c_2 \mathcal{V}^{2/3} \phi^{4/3}\big]\big),$$

$$\phi \equiv \sqrt{4\lambda_\phi/(3\mathcal{V})}\tau_\phi^{3/4},$$

$$a_\phi\langle\tau_\phi\rangle\approx\mathcal{O}(1)\times\ln\left(\mathcal{V}\right)\Leftrightarrow\langle\phi\rangle\approx\mathcal{O}(1)\times\frac{\ln\left(\mathcal{V}\right)^{3/4}}{\mathcal{V}^{1/2}}$$

$$V(\hat{\phi})\simeq V_0\left(1-\kappa_1e^{-\kappa_2\hat{\phi}}\right)$$

$$\kappa_1\equiv c_1\mathcal{V}^{5/3}\langle\phi\rangle^{4/3}\approx\mathcal{O}(\mathcal{V}\ln\left(\mathcal{V}\right))$$

$$\kappa_2\equiv\frac{4}{3}c_2\mathcal{V}^{2/3}\langle\phi\rangle^{1/3}=\mathcal{O}(\mathcal{V}^{1/2}\ln\left(\mathcal{V}\right)^{1/4})$$

$$\mathcal{V}^{2/3}\phi^{4/3}\gg 1 \text{ and } \phi\ll 1$$

$$\delta V_{(g_s)}\sim \frac{1}{\sqrt{\tau_\phi}\mathcal{V}^3}\sim \frac{1}{\phi^{2/3}\mathcal{V}^{10/3}}$$

$$\delta\eta\sim\frac{\delta V_{(g_s)}''}{V_0}\sim\frac{1}{\phi^{8/3}\mathcal{V}^{1/3}}\sim\frac{\mathcal{V}}{\tau_\phi^2}$$



$$\delta\eta\approx a_\phi^2\frac{\mathcal{V}}{\ln(\mathcal{V})^2}\gg1$$

$$\mathcal{V}=\frac{1}{2}\sqrt{\tau_1\tau_2},$$

$$\mathcal{V}=\alpha\left(\sqrt{\tau_1}\tau_2-\lambda_s\tau_s^{3/2}\right)$$

$$W=W_0+\mathcal{A}_se^{-a_sT_s}$$

$$V=a_s^2\mathcal{A}_s^2\frac{\sqrt{\tau_s}}{\mathcal{V}}e^{-2a_s\tau_s}-a_s\mathcal{A}_sW_0\frac{\tau_s}{\mathcal{V}}e^{-a_s\tau_s}+\hat{\xi}W_0^2\frac{1}{\mathcal{V}^3}.$$

$$\delta V_{(g_s)}=\frac{W_0^2}{\mathcal{V}^2}\left(\mathfrak{a}\frac{g_s^2}{\tau_1^2}-\mathfrak{b}\frac{1}{\sqrt{\tau_1}\mathcal{V}}+\mathfrak{c}\frac{g_s^2\tau_1}{\mathcal{V}^2}\right)$$

$$V(\phi)=V_0\Big(1-\frac{4}{3}e^{-\phi/\sqrt{3}}+\frac{1}{3}e^{-4\phi/\sqrt{3}}+\frac{\mathfrak{C}}{3}e^{2\phi/\sqrt{3}}\Big),$$

$$\phi\equiv\frac{\sqrt{3}}{2}\ln\,\tau_1$$

$$V(\phi)\simeq V_0\left(1-\frac{4}{3}e^{-\phi/\sqrt{3}}\right)$$

$$W=\mathcal{A}_a\exp\left(-S_a+\mathcal{A}_be^{-S_b}\right),$$

$$\mathcal{V}=\alpha\left(\sqrt{\tau_1}\tau_2-\lambda_s\tau_s^{3/2}\right)$$

$$\mathcal{V}=\tau_b^{3/2}-\tau_s^{3/2}-(\tau_s+\tau_w)^{3/2}$$

$$W=W_0+\mathcal{A}\mathrm{exp}\left[-aT_s\right]-B\mathrm{exp}\left[-bT_s\right],$$

$$W=W_0+\mathcal{A}\mathrm{exp}\left[-a(T_s+c_1e^{-2\pi T_1})\right]-B\mathrm{exp}\left[-b(T_s+c_2e^{-2\pi T_1})\right],$$

$$V=\frac{F_{\mathrm{poly}}}{\mathcal{V}^{3+p}}\big(1-(1+2\pi\hat{\tau}_1)e^{-2\pi\hat{\tau}_1}\big)$$

$$V(\hat{\phi})\simeq V_0\left(1-\kappa_2\hat{\phi}e^{-\kappa_2\hat{\phi}}\right), \hat{\phi}\approx \frac{\sqrt{3}}{2}\frac{\hat{\tau}_1}{\langle\tau_1\rangle},$$

$$V(\phi)\approx V_0(1-\kappa_1e^{-\kappa_2\phi})$$

$$\eta\simeq -\kappa_1\kappa_2^2e^{-\kappa_2\phi} \text{ and } \epsilon\simeq \frac{1}{2}\frac{\eta^2}{\kappa_2^2}$$

$$n_s\simeq 1+2\eta$$

$$r \simeq \frac{2}{\kappa_2^2} (n_s - 1)^2 \stackrel{n_s=0.96}{\rightarrow} \frac{3 \times 10^{-3}}{\kappa_2^2}$$

$$r \simeq 6(n_s - 1)^2 \stackrel{n_s=0.96}{\rightarrow} 0.01.$$

$$r\sim 10^{-5}$$

$$S=\int\,\mathrm{d}^4x\sqrt{-g}\left[\frac{M_{\mathrm{pl}}^2}{2}R-\frac{1}{2}(\partial\phi)^2-V(\phi)+\mathcal{O}(\phi,\psi)\right].$$

$$\mathcal{L}_{\mathrm{int}}=-\frac{1}{2}g^2(\phi-\phi_0)^2\psi^2$$

$$\phi(t) \approx \phi_0 + \dot{\phi}_0(t-t_0)$$

$$m_\psi^2(t) \equiv g^2(\phi-\phi_0)^2 \approx k_\star^4(t-t_0)^2$$

$$\ddot{\psi}_k+3H\dot{\psi}_k+\underbrace{\left(\frac{k^2}{a^2}+k_\star^4(t-t_0)^2\right)}_{\equiv\omega_k^2(t)}\psi_k=0$$

$$|\dot{\omega}_k|>\omega_k^2$$

$$n_k=e^{-\pi k^2/k_\star^2}$$

$$n_\psi(t_0)=\int\frac{\mathrm{d}^3k}{(2\pi)^3}n_k\approx\frac{k_\star^3}{(2\pi)^3}$$

$$n_\psi(t)=\frac{k_\star^3}{(2\pi)^3}\frac{a^3(t_0)}{a^3(t)}\Theta(t-t_0)$$

$$\ddot{\phi}+3H\dot{\phi}+V'=-g^2(\phi-\phi_0)\langle\psi^2\rangle,$$

$$\langle\psi^2\rangle\approx\frac{n_\psi(t)}{g|\phi-\phi_0|}$$

$$\mathcal{L}_{\mathrm{int}}=-\frac{1}{2}g^2\sum_{i=1}^N(\phi-\phi_i)^2\psi_i^2$$

$$n_{\psi_i}(t_i)\simeq \frac{\left(g\dot{\phi}(t_i)\right)^{3/2}}{(2\pi)^3}\equiv \frac{\dot{\phi}^{3/2}(t_i)}{(2\pi)^3}$$

$$\ddot{\phi}+3H\dot{\phi}+V'+\sum_i\frac{g\dot{\phi}^{3/2}(t_i)}{(2\pi)^3}\frac{a^3(t_i)}{a^3(t)}=0.$$



$$\sum_i \frac{g \dot{\phi}^{3/2}(t_i) a^3(t_i)}{(2\pi)^3 a^3(t)} \approx \int^t \frac{dt'}{\Delta} \frac{dt'}{\Delta} \frac{\dot{\phi}^{5/2}(t') a^3(t')}{(2\pi)^3 a^3(t)} \approx \frac{1}{3H\Delta} \frac{\dot{\phi}^{5/2}(t)}{(2\pi)^3},$$

$$\ddot{\phi} + 3H\dot{\phi} + V' + \frac{1}{24\pi^3} \frac{g^{5/2}}{H\Delta} \dot{\phi}^{5/2} = 0$$

$$\dot{\phi} = g\dot{\phi} \approx -(24\pi^3 H\Delta V')^{2/5}.$$

$$3M_{\rm pl}^2 H^2 = \rho_\phi + \rho_\psi \approx V(\phi)$$

$$2M_{\rm pl}^2 \dot{H} \approx -\rho_\psi,$$

$$\varepsilon = -\frac{\dot{H}}{H^2} \approx \frac{3}{2} \frac{\rho_\psi}{V},$$

$$\rho_\psi(t) = \sum_i g|\phi - \phi_i| n_{\psi_i}(t) \approx \int^t \frac{dt'}{\Delta} \frac{dt'}{\Delta} |\phi(t) - \phi(t')| \frac{\dot{\phi}^{5/2}(t') a^3(t')}{(2\pi)^3 a^3(t)}.$$

$$\rho_\psi(t) \approx \frac{1}{(3H)^2} \frac{1}{g\Delta} \frac{\dot{\phi}^{7/2}(t)}{(2\pi)^3}$$

$$\varepsilon \sim \frac{\epsilon^{7/10}}{g} \left(\frac{H}{M_{\rm pl}} \frac{\Delta^2}{M_{\rm pl}^2} \right)^{1/5}$$

$$\begin{aligned} t_x &: (x, u_1, u_2) \mapsto (x+1, u_1, u_2) \\ t_{u_1} &: (x, u_1, u_2) \mapsto (x-Mu_2, u_1+1, u_2) \\ t_{u_2} &: (x, u_1, u_2) \mapsto (x, u_1, u_2+1), \end{aligned}$$

$$\frac{ds^2}{\alpha'} = L_{u_1}^2 du_1^2 + \underbrace{L_{u_2}^2 du_2^2 + L_x^2 (dx + Mu_1 du_2)^2}_{T^2}$$

$$\frac{ds_{T^2}^2(u_1)}{\alpha'} = L_{u_2}^2 du_2^2 + L_x^2 (dx + Mu_1 du_2)^2$$

$$\frac{ds_{T^2,\perp}^2}{\alpha'} = L_x^2 dy_1^2 + L_{u_2}^2 dy_2^2$$

$$S_{\rm D4} = -\frac{1}{(2\pi)^4 g_s (\alpha')^2} \int d^4x \sqrt{-g} \sqrt{L_{u_2}^2 + L_x^2 M^2 u_1^2} \left(1 - \frac{1}{2} \alpha' L_{u_1}^2 \dot{u}_1^2\right).$$

$$S_{\rm D4} = \int d^4x \sqrt{-g} \left(\frac{1}{2} \dot{\phi}^2 - \mu^{10/3} \phi^{2/3} \right)$$

$$\frac{\phi^2}{M_{\rm pl}^2} = \frac{2}{9} (2\pi)^3 g_s \frac{M}{L^3} \frac{L_{u_1}}{L_{u_2}} u_1^3$$

$$\frac{\mu}{M_{\rm pl}} = \frac{M_s}{M_{\rm pl}} \left(\frac{9}{4} \frac{M^2}{(2\pi)^8 g_s^2} \left(\frac{L_x}{L} \right)^3 \frac{L_{u_2}}{L_{u_1}} \right)^{1/10}$$



$$\Delta u_1^3 \gg \frac{L^3}{M}$$

$$\int_{\Sigma_2} \mathcal{F}_2 = n \in \mathbb{Z}$$

$$T_{\mathrm{D3}/\Sigma_2}=T_3\sqrt{\ell^2+(c\,g_s+n)^2},$$

$$\mathcal{M}_{10}=dS_4\times X_6,$$

$$\int_{dS_4\times\Sigma_{p-2}}F_{p+2}=Q_0\gg1.$$

$$\mathrm{d}s_{\mathrm{E}}^2=H^{-2}(\,\mathrm{d}\xi^2+\sin^2\,\xi\,\mathrm{d}\Omega_3^2)+\mathrm{d}z^2,$$

$$S=\int\,\mathrm{d}z\int\,\mathrm{d}\mathcal{H}_3\,\mathrm{d}t\frac{\sinh^3\left(Ht\right)}{H^3}\Bigg(-2\sigma\delta(z-z_b)\sqrt{1-(\partial z_b)^2}-\frac{F_5^2}{2\cdot5!}\Bigg)$$

$$\frac{F_5^2}{5!}=\mu^5Q^2=\mu^5\left(Q_0+\sum_{j=-\infty}^{\infty}[\Theta(z-z_b+j\ell)-\Theta(z+z_b+j\ell)]\right)^2$$

$$S=\int\,\mathrm{d}t\,\mathrm{d}^3x a^3(t)\left(-2\sigma\sqrt{1-(\partial z_b)^2}-V(z_b)\right)$$

$$V(z_b)\sim \mu^5\left(Q_0-\frac{z_b}{\ell}\right)^2.$$

$$\mathcal{L} \supset - \sum_{i=1}^N \alpha_i \frac{\phi}{f} F_i \tilde{F}_i$$

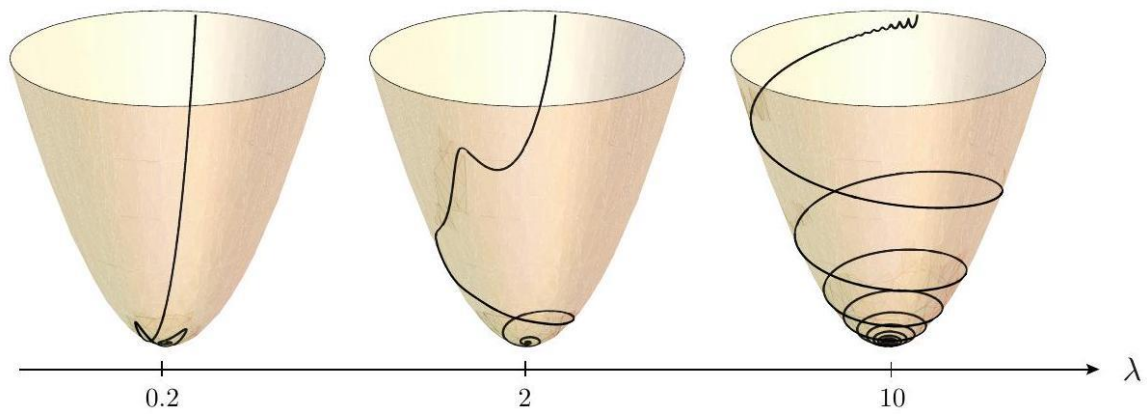
$$S_{\rm CS}=\frac{i}{2\pi}\int_{\mathcal{M}_4}C_0{\rm Tr}[\mathcal{F}_2\wedge\mathcal{F}_2]$$

$$S_{\rm eff}=\int\,\mathrm{d}^4x\left\{a^3\left[\gamma_C\dot{C}_0^2-V(C_0)+\gamma_A\frac{\mathrm{Tr}(\dot{A}^2)}{a^2}+\gamma_A\frac{\mathrm{Tr}([A,A]^2)}{a^4}\right]+\kappa C_0\mathrm{Tr}(\dot{A}[A,A])\right\}$$

$$C_0(t)\equiv \frac{\phi(t)}{\sqrt{\gamma_C}}$$

$$A_0\equiv 0, A_i(t)\equiv \frac{a(t)\psi(t)}{\sqrt{\gamma_A v}}J_i$$

$$S_{\rm eff}=\int\,\mathrm{d}^4x a^3\left[\frac{1}{2}\dot{\phi}^2-V(\phi)+\frac{3}{2}(\dot{\psi}+H\psi)^2-\frac{3}{2}g^2\psi^4-\frac{3g\lambda}{f}\phi\psi^2(\dot{\psi}+H\psi)\right]$$



$$\ddot{\phi} + 3H\dot{\phi} + V_{,\phi} = -3\frac{g\lambda}{f}\psi^2(\dot{\psi} + H\psi)$$

$$(N_e)_{\rm max} \approx \frac{3}{5}\lambda$$

$$S_{\rm CS} = \frac{i}{(2\pi)^3} \int_{\mathcal{M}_4 \times \Sigma_4} \frac{C_0}{24} \Big(\mathrm{Tr}[\mathcal{F}_2 \wedge \mathcal{F}_2 \wedge \mathcal{F}_2 \wedge \mathcal{F}_2] + \frac{1}{2} \mathrm{Tr}[\mathcal{F}_2 \wedge \mathcal{F}_2] \hat{\chi}(R) \Big)$$

$$\lambda=\left[K+\frac{\chi(\Sigma_4)}{24}\right]\times g_{\rm s}\times \frac{\ell_{\rm s}^4}{\mathcal{V}_4}$$

$$\delta\ddot{\phi}+\left(M^2+\frac{k^2}{a^2}\right)\delta\phi+\int^t\mathrm{d}\,\mathrm{d}t' M^2\left(\frac{5}{2}\dot{\delta\phi}(t')-3H\delta\phi(t')\right)\frac{a^3(t')}{a^3(t)}=-g\Delta n_{\psi}(k,t)$$

$$M^2 \equiv \frac{g^{5/2}}{(2\pi)^3} \frac{\dot{\phi}^{3/2}}{\Delta}$$

$$\Delta_{\mathcal{R}}^2 \approx g^{8/3} \frac{H}{M} \left(\frac{M}{\Delta}\right)^{2/3}$$

$$\Delta_{\mathcal{R}}^2 \approx g^{9/4} \left(\frac{H}{\Delta}\right)^{1/2} \left(\frac{H^2}{\dot{\phi}}\right)^{1/4},$$

$$n_s-1=\frac{\dot{H}}{H^2}-\frac{1}{4}\frac{\ddot{\phi}}{H\dot{\phi}}.$$

$$r=g^{-8/2}\frac{HM}{M_{\rm pl}^2}\Big(\frac{\Delta}{M}\Big)^{2/3}.$$

$$f_{\rm NL}^{\rm equil} \simeq \frac{M^2}{H^2}$$

Interacciones gravitónicas y transferencia de gravedad al centro de masa de las partículas portadoras, volviéndolas supermasivas. Formalismo de Batalin-Vilkovisky por permeabilidad de un campo gravitónico para inmersión gravitacional de un espacio – tiempo cuántico específico.

Propagadores y campos fantasma.

$$Q = \oint_0 j_B(z) dz + \oint_0 \bar{j}_B(\bar{z}) d\bar{z}$$

$$j_B = cT_m + bc\partial c, \bar{j}_B = \bar{c}\bar{T}_m + \bar{b}\bar{c}\bar{\partial}\bar{c}$$

$$\oint_0 \frac{dz}{z} = 1, \oint_0 \frac{d\bar{z}}{\bar{z}} = 1$$

$$b(z)c(w) \simeq \frac{1}{z-w}, \bar{b}(\bar{z})\bar{c}(\bar{w}) \simeq \frac{1}{\bar{z}-\bar{w}}$$

$$T_m(z)T_m(w) \simeq \frac{26}{2} \frac{1}{(z-w)^4} + \frac{2}{(z-w)^2} T_m(w) + \frac{1}{z-w} \partial_w T_m(w)$$

$$\bar{T}_m(\bar{z})\bar{T}_m(\bar{w}) \simeq \frac{26}{2} \frac{1}{(\bar{z}-\bar{w})^4} + \frac{2}{(\bar{z}-\bar{w})^2} \bar{T}_m(\bar{w}) + \frac{1}{\bar{z}-\bar{w}} \bar{\partial}_w \bar{T}_m(\bar{w})$$

$$T_{bc}(z) = -\partial bc - 2b\partial c, \bar{T}_{bc}(\bar{z}) = -\bar{\partial}\bar{b}\bar{c} - 2\bar{b}\bar{\partial}\bar{c}$$

$$\partial X^\mu(z)\partial X^\nu(w) = -\frac{\eta^{\mu\nu}}{2(z-w)^2}, \bar{\partial}X^\mu(\bar{z})\bar{\partial}X^\nu(\bar{w}) = -\frac{\eta^{\mu\nu}}{2(\bar{z}-\bar{w})^2}$$

$$T(z) = -\eta_{\mu\nu}\partial X^\mu\partial X^\nu, \bar{T}(\bar{z}) = -\eta_{\mu\nu}\bar{\partial}X^\mu\bar{\partial}X^\nu$$

$$|k\rangle \equiv e^{ik\cdot X}(0)|0\rangle,$$

$$\langle k|c_{-1}\bar{c}_{-1}c_0\bar{c}_0c_1\bar{c}_1|k'\rangle = -(2\pi)^D\delta^{(D)}(k+k')$$

$$\langle k|c_{-1}c_0c_1|k'\rangle' = -(2\pi)^{p+1}\delta^{(p+1)}(k+k')$$

$$\langle k|c_{-1}c_0c_1|k'\rangle = -(2\pi)^{p+1}K\delta^{(p+1)}(k+k').$$

$$|s\rangle \in \mathcal{H}_c \text{ iff } b_0^-|s\rangle = 0, L_0^-|s\rangle = 0$$

$$b_0^\pm = b_0 \pm \bar{b}_0, L_0^\pm = L_0 \pm \bar{L}_0, c_0^\pm = (c_0 \pm \bar{c}_0)/2$$

$$\langle A, B \rangle \equiv \langle A | c_0^- | B \rangle \quad A, B \in \mathcal{H}_c$$

$$\langle A, B \rangle' \equiv \langle A \mid B \rangle', \langle A, B \rangle \equiv \langle A \mid B \rangle.$$

$$\gamma = \eta e^\phi, \beta = \partial \xi e^{-\phi}, \delta(\gamma) = e^{-\phi}, \delta(\beta) = e^\phi$$

$$[\gamma] = -\frac{1}{2}, [\beta] = \frac{3}{2}, [\eta] = 1, [\xi] = 0, [\phi] = 0, [e^{q\phi}] = -\frac{1}{2}q(q+2)$$



$$\xi(z)\eta(w)\simeq\frac{1}{z-w},\partial\phi(z)\partial\phi(w)\simeq-\frac{1}{(z-w)^2},e^{q_1\phi}(z)e^{q_2\phi}(w)\simeq(z-w)^{-q_1q_2}e^{(q_1+q_2)\phi}(w)$$

$$T_{\beta\gamma}=\frac{3}{2}\beta\partial\gamma+\frac{1}{2}\gamma\partial\beta=T_{\xi\eta}+T_{\phi},T_{\xi\eta}=-\eta\partial\xi,T_{\phi}=-\frac{1}{2}\partial\phi\partial\phi-\partial^2\phi.$$

$$T_m(z)T_m(w)\simeq\frac{15}{2}\frac{1}{(z-w)^4}+\frac{2}{(z-w)^2}T_m(w)+\frac{1}{z-w}\partial_wT_m(w)$$

$$T_F(z)T_F(w)\simeq\frac{5}{2}\frac{1}{(z-w)^3}+\frac{1}{2}\frac{1}{z-w}T_m(w)$$

$$T_m(z)T_F(w)\simeq\frac{3}{2}\frac{1}{(z-w)^2}T_F(w)+\frac{1}{z-w}T_F(w)$$

$$\psi^\mu(z)\psi^\nu(w)\simeq-\frac{\eta^{\mu\nu}}{2(z-w)}$$

$$j_B=c(T_m+T_{\beta\gamma})+bc\partial c+\gamma T_F-\frac{1}{4}\gamma^2b$$

$$\mathcal{X}(z)=\{Q,\xi(z)\}=c\partial\xi+e^{\phi}T_F-\frac{1}{4}\partial\eta e^{2\phi}b-\frac{1}{4}\partial(\eta e^{2\phi}b).$$

$$\begin{array}{l} \mathrm{gh}(c)=\mathrm{gh}(\eta)=\mathrm{gh}(\gamma)=1\\ \mathrm{gh}(b)=\mathrm{gh}(\xi)=\mathrm{gh}(\beta)=-1 \end{array}$$

$$\mathrm{pic}(\eta)=-1,\mathrm{pic}(\xi)=1,\mathrm{pic}(e^{q\phi})=q,\mathrm{pic}(\mathcal{X})=1$$

$$\text{GSO odd fields: } \beta, \gamma, \psi^\mu, T_F$$

$$|k;m,n\rangle\equiv e^{ik\cdot X}(0)e^{m\phi}(0)e^{n\bar\phi}(0)|0\rangle,|k;m\rangle\equiv e^{ik\cdot X}(0)e^{m\phi}(0)|0\rangle$$

$$\langle k;0,0|c_{-1}\bar{c}_{-1}c_0\bar{c}_0c_1\bar{c}_1e^{-2\phi}e^{-2\bar{\phi}}|k';0,0\rangle=-(2\pi)^D\delta^{(D)}(k+k'),$$

$$\langle k;0|c_{-1}c_0c_1e^{-2\phi}|k';0\rangle'=(2\pi)^{p+1}\delta^{(p+1)}(k+k'),$$

$$\langle k;0|c_{-1}c_0c_1e^{-2\phi}|k';0\rangle=(2\pi)^D\delta^{(D)}(k+k')$$

$$\mathcal{C}=\left\langle \prod_i f_i\circ V_i(0)\right\rangle_{\Sigma},$$

$$f\circ V(w)=(f'(w))^h(\bar{f}'(\bar{w}))^{\bar{h}}V(f(w))$$

$$V(w)(dw)^h(d\bar{w})^{\bar{h}}=V(z)(dz)^h(d\bar{z})^{\bar{h}}\rightarrow V(w)=V(z(w))\left(\frac{dz}{dw}\right)^h\left(\frac{d\bar{z}}{d\bar{w}}\right)^{\bar{h}}$$



$$\mathcal{C}=\left\langle \prod_iV_i(w_i=0)\right\rangle _{\Sigma}$$

$$\langle \chi^c_s \mid \chi_r \rangle = \delta_{rs} \Leftrightarrow \sum_r |\chi_r \rangle \langle \chi^c_r| = I$$

$$w_1w_2=1$$

$$w_1w_2=-1$$

$$d_{g,n}\equiv \dim_{\mathbb{R}}(\mathcal{M}_{g,n})=6g-6+2n$$

$$\chi_{g,n}=2-2g-n$$

$$\mathcal{A}_g(V_1,\cdots,V_n)=(g_s)^{-\chi_{g,n}}\int_{\mathcal{F}_{g,n}}\Omega^{(g,n)}_{6g-6+2n}(V_1,\cdots,V_n)$$

$$\sigma_s = F_s(\tau_s,u)$$

$$\mathcal{B}\left[\frac{\partial}{\partial u^i}\right]\equiv \sum_s\left[\oint_{c_s}\frac{\partial F_s}{\partial u^i}d\sigma_sb(\sigma_s)+\oint_{c_s}\frac{\partial \bar{F}_s}{\partial u^i}d\bar{\sigma}_s\bar{b}(\bar{\sigma}_s)\right]$$

$$\Omega_p^{(g,n)}(A_1,\cdots,A_n)\Big[\frac{\partial}{\partial u^{j_1}},\cdots,\frac{\partial}{\partial u^{j_p}}\Big]\equiv\Big(-\frac{1}{2\pi i}\Big)^{3g-3+n}\Big\langle\mathcal{B}\Big[\frac{\partial}{\partial u^{j_1}}\Big]\cdots\mathcal{B}\Big[\frac{\partial}{\partial u^{j_p}}\Big]A_1\cdots A_n\Big\rangle_{\Sigma_{g,n}}$$

$$z=F(w,u)$$

$$y(u)=F(0,u)$$

$$\Omega_2(c\bar{c}V)=\Big(-\frac{1}{2\pi i}\Big)du^1\wedge du^2\Big\langle\cdots\mathcal{B}\Big[\frac{\partial}{\partial u^1}\Big]\mathcal{B}\Big[\frac{\partial}{\partial u^2}\Big]c\bar{c}V(w=0)\Big\rangle$$

$$\mathcal{B}\left[\frac{\partial}{\partial u^1}\right]\mathcal{B}\left[\frac{\partial}{\partial u^2}\right]=\left[\oint b(z)dz\frac{\partial F(w,u)}{\partial u^1}+\oint \bar{b}(\bar{z})d\bar{z}\frac{\partial \overline{F(w,u)}}{\partial u^1}\right] \\ \left[\oint b(z)dz\frac{\partial F(w,u)}{\partial u^2}+\oint \bar{b}(\bar{z})d\bar{z}\frac{\partial \overline{F(w,u)}}{\partial u^2}\right]$$

$$\Omega_2(c\bar{c}V)=-\Big(-\frac{1}{2\pi i}\Big)du^1\wedge du^2\Big(\frac{\partial y}{\partial u^1}\frac{\partial \bar{y}}{\partial u^2}-\frac{\partial y}{\partial u^2}\frac{\partial \bar{y}}{\partial u^1}\Big)\langle\cdots V(z=y)\rangle$$

$$du^1\wedge du^2\Big(\frac{\partial y}{\partial u^1}\frac{\partial \bar{y}}{\partial u^2}-\frac{\partial y}{\partial u^2}\frac{\partial \bar{y}}{\partial u^1}\Big)=dy\wedge d\bar{y}$$

$$\Omega_2(c\bar{c}V)=\frac{1}{2\pi i}\langle\cdots(dy\wedge d\bar{y}V(z=y))\rangle$$

$$\Omega_p^{(g,n)}(QA_1,A_2,\cdots,A_n)+\cdots+(-1)^{A_1+\cdots+A_{n-1}}\Omega_p^{(g,n)}(A_1,A_2,\cdots,QA_n)=(-1)^p\,\mathrm{d}\Omega_{p-1}^{(g,n)}(A_1,\cdots,A_n)$$



$$\begin{aligned}
& \int_{\mathcal{F}'_{g,n}} \Omega_{6g-6+2n}^{(g,n)}(V_1, \dots, V_n) - \int_{\mathcal{F}_{g,n}} \Omega_{6g-6+2n}^{(g,n)}(V_1, \dots, V_n) = \int_{\mathcal{R}_{g,n}} d\Omega_{6g-5+2n}^{(g,n)}(V_1, \dots, V_n) \\
& \Omega_{6g-6+2n}^{(g,n)}(QW_1, V_2, \dots, V_n) = d\Omega_{6g-7+2n}^{(g,n)}(W_1, V_2, \dots, V_n) \\
& n_{-1,-1}, n_{-1,-1/2}, n_{-1/2,-1}, n_{-1/2,-1/2} \\
& 2g-2 = N_L - n_{-1,-1} - n_{-1,-1/2} - \frac{1}{2}n_{-1/2,-1} - \frac{1}{2}n_{-1/2,-1/2} \\
& 2g-2 = N_R - n_{-1,-1} - n_{-1/2,-1} - \frac{1}{2}n_{-1,-1/2} - \frac{1}{2}n_{-1/2,-1/2} \\
& p = k + \ell + \bar{\ell} \\
& \Omega_p^{(g,n)}(A_1, \dots, A_n) \left[\frac{\partial}{\partial u^{j_1}}, \dots, \frac{\partial}{\partial u^{j_k}}, \frac{\partial}{\partial y_{\alpha_1}}, \dots, \frac{\partial}{\partial y_{\alpha_\ell}}, \frac{\partial}{\partial \bar{y}_{\beta_1}}, \dots, \frac{\partial}{\partial \bar{y}_{\beta_{\bar{\ell}}}} \right] \\
& \equiv \left(-\frac{1}{2\pi i} \right)^{3g-3+n} \left\langle \mathcal{B} \left[\frac{\partial}{\partial u^{j_1}} \right] \dots \mathcal{B} \left[\frac{\partial}{\partial u^{j_k}} \right] (-\partial \xi(y_{\alpha_1})) \dots (-\partial \xi(y_{\alpha_\ell})) (-\bar{\partial} \bar{\xi}(\bar{y}_{\beta_1})) \dots (-\bar{\partial} \bar{\xi}(\bar{y}_{\beta_{\bar{\ell}}})) \prod_{\alpha=l+1}^{N_R} \mathcal{X}(y_\alpha) \prod_{\beta=\bar{l}+1}^{N_L} \bar{\mathcal{X}}(\bar{y}_\beta) A_1 \dots A_n \right\rangle_{\Sigma_{g,n}} \\
& - \int_{y_\alpha}^{y'_\alpha} \partial \xi(y) dy = \xi(y_\alpha) - \xi(y'_\alpha) \\
& \Omega_p^{(g,n)}(A_1, \dots, A_n) \left[\frac{\partial}{\partial u^{j_1}}, \dots, \frac{\partial}{\partial u^{j_k}}, \frac{\partial}{\partial y_{\alpha_1}}, \dots, \frac{\partial}{\partial y_{\alpha_\ell}} \right] \\
& = \left(-\frac{1}{2\pi i} \right)^{3g-3+n} \left\langle \mathcal{B} \left[\frac{\partial}{\partial u^{j_1}} \right] \dots \mathcal{B} \left[\frac{\partial}{\partial u^{j_k}} \right] (-\partial \xi(y_{\alpha_1})) \dots (-\partial \xi(y_{\alpha_\ell})) \prod_{\alpha=l+1}^{N_R} \mathcal{X}(y_\alpha) A_1 \dots A_n \right\rangle_{\Sigma_{g,n}}
\end{aligned}$$

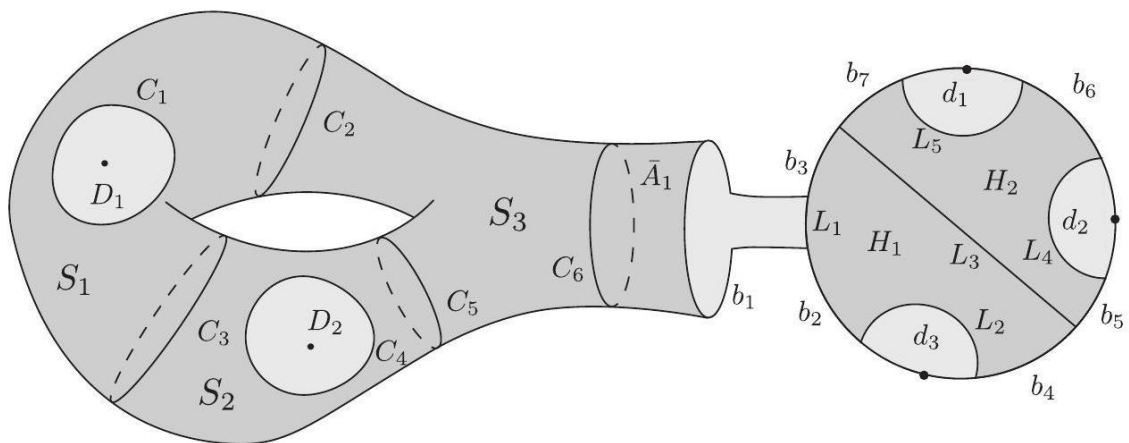


Figura 41. Puente Einstein – Rosen en dimensión $\mathbb{R}^4$ para entrelazamiento cuántico.

$$d_{g,b,n_c,n_o} \equiv \dim_{\mathbb{R}}(\mathcal{M}_{g,b,n_c,n_o}) = 6g - 6 + 3b + 2n_c + n_o$$

$$\chi_{g,b,n_c,n_o} = 2 - 2g - n_c - b - \frac{1}{2}n_o$$

$$S + \frac{1}{2} \#H + \frac{1}{2} \# \bar{A} = -\chi_{g,b,n_c,n_o}$$

$$2\#C = 3S + n_c + A + \bar{A}$$

$$2\#L = 3\#H + n_o + \#\bar{A}$$

$$\mathcal{B}\left[\frac{\partial}{\partial u^i}\right] \equiv \sum_s \left[\oint_{c_s} \frac{\partial F_s}{\partial u^i} d\sigma_s b(\sigma_s) + \oint_{c_s} \frac{\partial \bar{F}_s}{\partial u^i} d\bar{\sigma}_s \bar{b}(\bar{\sigma}_s) \right] + \sum_m \left[\int_{L_m} \frac{\partial G_m}{\partial u^i} d\sigma_m b(\sigma_m) + \int_{L_m} \frac{\partial \bar{G}_m}{\partial u^i} d\bar{\sigma}_m \bar{b}(\bar{\sigma}_m) \right]$$

$$\sigma_m = G_m(\tau_m,u^i)$$

$$\begin{aligned} &\widehat{\Omega}_p^{(g,b,n_c,n_o)}(A_1^c,\cdots,A_{n_c}^c;A_1^o,\cdots,A_{n_o}^o)\left[\frac{\partial}{\partial u^{j_1}},\cdots,\frac{\partial}{\partial u^{j_p}}\right] \\ &\sim \left\langle \mathcal{B}\left[\frac{\partial}{\partial u^{j_1}}\right]\cdots\mathcal{B}\left[\frac{\partial}{\partial u^{j_p}}\right]A_1^c\cdots A_{n_c}^c;A_1^o\cdots A_{n_o}^o\right\rangle_{\Sigma_{g,b,n_c,n_o}} \end{aligned}$$

$$\Omega_p^{(g,b,n_c,n_o)}(A_1^c,\cdots,A_{n_c}^c;A_1^o,\cdots,A_{n_o}^o)\equiv N_{g,b,n_c,n_o}\hat{\Omega}_p^{(g,b,n_c,n_o)}(A_1^c,\cdots,A_{n_c}^c;A_1^o,\cdots,A_{n_o}^o)$$

$$N_{g,b,n_c,n_o}=\eta_c^{3g-3+n_c+\frac{3}{2}b+\frac{3}{4}n_o},\eta_c\equiv-\frac{1}{2\pi i}=\frac{i}{2\pi}$$

$$\mathcal{A}_{g,b}(A_1^c,\cdots,A_{n_c}^c;A_1^o,\cdots,A_{n_o}^o)=(g_s)^{-\chi_{g,b,n_c,n_o}}\int_{\mathcal{F}_{g,b,n_c,n_o}}\Omega_{d_{g,b,n_c,n_o}}^{(g,b,n_c,n_o)}(A_1^c,\cdots,A_{n_c}^c;A_1^o,\cdots,A_{n_o}^o)$$

$$\widehat{\Omega}_0^{(0,1,1,0)}(A^c)=\eta_c\langle c_0^-A^c\rangle,$$

$$\Omega_0^{(0,1,1,0)}(A^c)=\eta_c N_{0,1,1,0}\langle c_0^-A^c\rangle.$$

$$c_o \equiv N_{g,b,n_c,n_o+1}/N_{g,b,n_c,n_o}$$

$$z=G(w,u)$$

$$\Omega_1(cV_o)=c_o du \left\langle \cdots \mathcal{B}\left[\frac{\partial}{\partial u}\right] cV_o(w=0) \right\rangle$$

$$\Omega_1(cV_o)=c_o du \left\langle \cdots \oint b(z) dz \frac{\partial G(w,u)}{\partial u} cV_o(z=y) \right\rangle$$

$$\Omega_1(cV_o)=-c_o du \frac{\partial y}{\partial u} \langle \cdots V_o(z=y) \rangle = -c_o dy \langle \cdots V_o(y) \rangle$$



$$\Omega_p^{(g,b,n_c,n_o)}(A_1^c,\cdots,A_{n_c}^c;A_1^o,\cdots,A_{n_o}^o)\left[\frac{\partial}{\partial u^{j_1}},\cdots,\frac{\partial}{\partial u^{j_k}},\frac{\partial}{\partial y_{\alpha_1}},\cdots,\frac{\partial}{\partial y_{\alpha_\ell}},\frac{\partial}{\partial \bar{y}_{\beta_1}},\cdots,\frac{\partial}{\partial \bar{y}_{\beta_\ell}}\right]$$

$$\sim N_{g,b,n_c,n_o}\left(\mathcal{B}\left[\frac{\partial}{\partial u^{j_1}}\right]\cdots\mathcal{B}\left[\frac{\partial}{\partial u^{j_k}}\right](-\partial\xi(y_{\alpha_1}))\cdots(-\partial\xi(y_{\alpha_\ell}))(-\bar{\partial}\bar{\xi}(\bar{y}_{\beta_1}))\cdots(-\bar{\partial}\bar{\xi}(\bar{y}_{\beta_\ell}))\prod_{\alpha=l+1}^{N_R}\mathcal{X}(y_\alpha)\prod_{\beta=l+1}^{N_L}\overline{\mathcal{X}}(\bar{y}_\beta)A_1^c\cdots A_{n_c}^cA_1^o\cdots A_{n_o}^o\right)_{\Sigma_{g,b,n_c,n_o}}$$

$$2(2g-2+b)=N-2n_{-1,-1}-\frac{3}{2}(n_{-1,-1/2}+n_{-1/2,-1})-n_{-1/2,-1/2},-n_{-1}-\frac{1}{2}n_{-1/2}$$

$$|{\bf b}\rangle=-\frac{dq_r}{q_r}(e^{-\Lambda}q_r)^{L_0+\bar{L}_0}b_0^+|B\rangle, q_r\in[0,1]$$

$$dyV_o(y)$$

$$\widehat{\Omega}^{(0,1,1,1)}(c\bar{c}V_c;cV_o)=-\langle cV_o c\bar{c}V_c\rangle$$

$$N_{g,b,n_c,n_o}=\eta_c^{3g-3+n_c+\frac{3}{2}b+\frac{3}{4}n_o}$$

$$\eta_c\equiv -\frac{1}{2\pi i}=\frac{i}{2\pi}$$

$$\widehat{\Omega}_p^{(g,b,n_c,n_o)}(QA_1^c,\cdots,A_{n_c};A_1^o,\cdots,A_{n_o}^o)+\cdots$$

$$+(-1)^{A_1^c+\cdots+A_n^c+A_1^o+\cdots+A_{n_o-1}^o}\widehat{\Omega}_p^{(g,b,n_c,n_o)}(A_1^c,\cdots,A_{n_c};A_1^o,\cdots,QA_{n_o}^o)$$

$$=(-1)^p\,\mathrm{d}\widehat{\Omega}_{p-1}^{(g,b,n_c,n_o)}(A_1^c,\cdots,A_{n_c};A_1^o,\cdots,A_{n_o}^o)$$

$$\{F,G\}=\frac{\partial_r F}{\partial \psi_m}\frac{\partial_\ell G}{\partial \psi_m^*}-\frac{\partial_r F}{\partial \psi_m^*}\frac{\partial_\ell G}{\partial \psi_m},$$

$$\frac{\partial_r A}{\partial \psi}=(-1)^{\psi(A+1)}\frac{\partial_l A}{\partial \psi}$$

$$\{F,G\}=-(-1)^{(F+1)(G+1)}\{G,F\}$$

$$\Delta F = (-1)^{\psi_m} \frac{\partial_\ell}{\partial \psi_m} \frac{\partial_\ell F}{\partial \psi_m^*} = (-1)^{\psi_m F} \frac{\partial_r}{\partial \psi_m} \left(\frac{\partial_\ell F}{\partial \psi_m^*} \right)$$

$$\{S_{cl},S_{cl}\}=0$$

$$\frac{1}{2}\{S,S\}+\Delta S=0$$

$$\Delta e^S=0$$

$$\psi_m^*=\frac{\partial_\ell\Psi}{\partial\psi_m}$$

$$\tilde{\psi}_m^*=\psi_m^*-\frac{\partial_\ell\Psi}{\partial\psi_m},\tilde{\psi}_m=\psi_m$$



$$\{F, G\} = \frac{\partial_r F}{\partial \psi^i} \omega^{ij} \frac{\partial_l G}{\partial \psi^j}, \omega^{ij} = -\omega^{ji}$$

$$\Delta F = \frac{1}{2} (-1)^{\psi^i} \frac{\partial_\ell}{\partial \psi^i} \left(\omega^{ij} \frac{\partial_\ell F}{\partial \psi^j} \right) = \frac{1}{2} (-1)^{\psi^i F} \frac{\partial_r}{\partial \psi^i} \left(\omega^{ij} \frac{\partial_\ell F}{\partial \psi^j} \right)$$

Campos bosónicos.

$$\Psi = \sum_i \varphi_i \psi^i$$

$$|\Psi\rangle = \sum_i (-1)^{\psi^i \epsilon_0} |\varphi_i\rangle \psi^i$$

$$b_0^- |\Psi\rangle = 0, L_0^- |\Psi\rangle = 0$$

$$\Psi \in \mathcal{H}_c.$$

$$\langle A, B \rangle \equiv \langle A | c_0^- | B \rangle \quad |A\rangle, |B\rangle \in \mathcal{H}_c.$$

$$\langle A, B \rangle = (-1)^{(A+1)(B+1)} \langle B, A \rangle$$

$$\langle \lambda A, B \eta \rangle = \lambda \langle A, B \rangle \eta,$$

$$\partial \mathcal{V}_{g,n} = -\Delta_c \mathcal{V}_{g-1,n+2} - \frac{1}{2} \sum_{\substack{g_1+g_2=g \\ n_1+n_2=n+2}} \{ \mathcal{V}_{g_1,n_1}, \mathcal{V}_{g_2,n_2} \}_c$$

$$\text{General collection of } \mathcal{V}_{g,n} = \begin{cases} \mathcal{V}_{0,n}, & n \geq 3, \\ \mathcal{V}_{1,n}, & n \geq 1, \\ \mathcal{V}_{g,n}, & n \geq 0, g \geq 2. \end{cases}$$

$$w_1w_2=e^{i\theta}, 0\leq \theta\leq 2\pi$$

$$\{A_1,\cdots,A_n\}\equiv\sum_{g=0}^{\infty}(g_s)^{-\chi_{g,n}}\{A_1,\cdots,A_n\}_g\equiv\sum_{g=0}^{\infty}g_s^{2g+n-2}\int_{\mathcal{V}_{g,n}}\Omega_{d_{g,n}}^{(g,n)}(A_1,\cdots,A_n)$$

$$\langle A_0, [A_1, \cdots, A_n] \rangle = \{A_0, \cdots, A_n\}, \forall A_0 \in \mathcal{H}_c$$

$$[A_1,\cdots,A_n]\equiv\sum_{g=0}^{\infty}g_s^{2g+n-1}[A_1,\cdots,A_n]_g$$



$$\begin{aligned}
& \sum_{i=1}^N \{A'_1 \dots A'_{i-1} (QA'_i) A'_{i+1} \dots A'_N\} \\
&= -\frac{1}{2} \sum_{\substack{\ell, k \geq 0 \\ \ell+k=N}} \sum_{\substack{\{i_a; a=1, \dots, \ell\} \\ \{j_b; b=1, \dots, k\} \\ \{i_a\} \cup \{j_b\} = \{1, \dots, N\}}} \{A'_{i_1} \dots A'_{i_\ell} \varphi_s\} \{\varphi_r A'_{j_1} \dots A'_{j_k}\} \langle \varphi_s^c, \varphi_r^c \rangle \\
&\quad - \frac{1}{2} \{A'_1 \dots A'_N \varphi_s \varphi_r\} \langle \varphi_s^c, \varphi_r^c \rangle
\end{aligned}$$

$$\langle \varphi_r^c, \varphi_s \rangle = \langle \varphi_s, \varphi_r^c \rangle = \delta_{rs}$$

$$\langle \varphi_i, \varphi_j \rangle = \langle \varphi_j, \varphi_i \rangle, \text{ and } \langle \varphi_i^c, \varphi_j^c \rangle = \langle \varphi_j^c, \varphi_i^c \rangle.$$

$$\begin{aligned}
Q[A'_1 \dots A'_N] &= - \sum_{i=1}^N [A'_1 \dots A'_{i-1} (QA'_i) A'_{i+1} \dots A'_N] \\
&\quad - \sum_{\substack{\ell, k \geq 0 \\ \ell+k=N}} \sum_{\substack{\{i_a; a=1, \dots, \ell\} \\ \{j_b; b=1, \dots, k\} \\ \{i_a\} \cup \{j_b\} = \{1, \dots, N\}}} [A'_{i_1} \dots A'_{i_\ell} [A'_{j_1} \dots A'_{j_k}]] - \frac{1}{2} [A'_1 \dots A'_N \varphi_s \varphi_r] \langle \varphi_s^c, \varphi_r^c \rangle
\end{aligned}$$

$$S = \frac{1}{2} \langle \Psi, Q\Psi \rangle + \sum_{n=1}^{\infty} \frac{1}{n!} \{\Psi^n\}$$

$$\Psi = \sum_r \varphi_r \psi^r \rightarrow \delta \Psi = \sum_r \varphi_r \delta \psi^r$$

$$\delta F = \sum_k \frac{\partial^r F}{\partial \psi^k} \delta \psi^k = \langle F_R, \delta \Psi \rangle \rightarrow \frac{\partial^r F}{\partial \psi^k} = \langle F_R, \varphi_k \rangle$$

$$\delta G = \sum_k \delta \psi^k \frac{\partial^l G}{\partial \psi^k} = \langle \delta \Psi, G_L \rangle \rightarrow \frac{\partial^l G}{\partial \psi^k} = (-1)^{\varphi_k} \langle \varphi_k, G_L \rangle$$

$$\{F, G\} = -\langle F_R, G_L \rangle$$

$$\omega^{rs} = (-1)^{\varphi_s+1} \langle \varphi_r^c, \varphi_s^c \rangle, \omega_{rs} = (-1)^{\varphi_r+1} \langle \varphi_r, \varphi_s \rangle,$$

$$Q\Psi + \sum_{n=1}^{\infty} \frac{1}{n!} [\Psi^n] = 0$$

$$S_{\text{cl}} = \frac{1}{2} \langle \Psi, Q\Psi \rangle + \sum_{n=3}^{\infty} \frac{1}{n!} \{\Psi^n\}_0$$

$$\delta \Psi = Q\Lambda + \sum_{n=1}^{\infty} \frac{1}{n!} [\Lambda, \Psi^n]_0 = Q\Lambda + [\Lambda, \Psi]_0 + \frac{1}{2!} [\Lambda, \Psi, \Psi]_0 + \frac{1}{3!} [\Lambda, \Psi, \Psi, \Psi]_0 + \dots$$



$$\begin{aligned}\Psi \in \mathcal{H}_c &\equiv \mathcal{H}_{-1,-1} \oplus \mathcal{H}_{-1,-1/2} \oplus \mathcal{H}_{-1/2,-1} \oplus \mathcal{H}_{-1/2,-1/2}, \\ \tilde{\Psi} \in \tilde{\mathcal{H}}_c &\equiv \mathcal{H}_{-1,-1} \oplus \mathcal{H}_{-1,-3/2} \oplus \mathcal{H}_{-3/2,-1} \oplus \mathcal{H}_{-3/2,-3/2}.\end{aligned}$$

$$\begin{aligned}&\sum_{i=1}^N \{A'_1 \dots A'_{i-1} (QA'_i) A'_{i+1} \dots A'_N\} \\ &= -\frac{1}{2} \sum_{\substack{\ell, k \geq 0 \\ \ell+k=N}} \sum_{\substack{\{i_a; a=1, \dots, \ell\} \\ \{j_b; b=1, \dots, k\} \\ \{i_a\} \cup \{j_b\} = \{1, \dots, N\}}} \{A'_{i_1} \dots A'_{i_\ell} \varphi_s\} \{\varphi_r A'_{j_1} \dots A'_{j_k}\} \langle \varphi_s^c, \mathcal{G} \varphi_r^c \rangle \\ &\quad - \frac{1}{2} \{A'_1 \dots A'_N \varphi_s \varphi_r\} \langle \varphi_s^c, \mathcal{G} \varphi_r^c \rangle\end{aligned}$$

$$\mathcal{X}_0 = \oint \frac{dz}{z} \mathcal{X}(z), \overline{\mathcal{X}}_0 = \oint \frac{d\bar{z}}{\bar{z}} \overline{\mathcal{X}}(\bar{z})$$

$$\mathcal{G} = \begin{cases} \mathbf{1} \text{ on } \mathcal{H}_{-1,-1} \\ \mathcal{X}_0 \text{ on } \mathcal{H}_{-1,-3/2} \\ \overline{\mathcal{X}}_0 \text{ on } \mathcal{H}_{-3/2,-1} \\ \mathcal{X}_0 \overline{\mathcal{X}}_0 \text{ on } \mathcal{H}_{-3/2,-3/2} \end{cases}$$

$$S = -\frac{1}{2} \langle \tilde{\Psi}, Q\mathcal{G}\tilde{\Psi} \rangle + \langle \tilde{\Psi}, Q\Psi \rangle + \sum_{n=1}^{\infty} \frac{1}{n!} \{\Psi^n\}$$

$$\delta F = \langle F_R, \delta \tilde{\Psi} \rangle + \langle \tilde{F}_R, \delta \Psi \rangle = \langle \delta \tilde{\Psi}, F_L \rangle + \langle \delta \Psi, \tilde{F}_L \rangle$$

$$\{F,G\}=-\langle F_R,\tilde{G}_L\rangle-\langle \tilde{F}_R,G_L\rangle-\langle \tilde{F}_R,\mathcal{G}\tilde{G}_L\rangle$$

$$\langle A_1,[A_1,\cdots,A_n]\rangle=\{A_1,\cdots,A_n\}$$

$$\delta|\tilde{\Psi}\rangle=Q|\tilde{\Lambda}\rangle+\sum_{n=2}^{\infty}\frac{1}{n!}[\Lambda\Psi^n]_0,\delta|\Psi\rangle=Q|\Lambda\rangle+\sum_{n=2}^{\infty}\frac{1}{n!}\mathcal{G}[\Lambda\Psi^n]_0$$

Campos heteróticos.

$$\begin{aligned}\Psi \in \mathcal{H}_c &\equiv \mathcal{H}_{-1} \oplus \mathcal{H}_{-1/2} \\ \tilde{\Psi} \in \tilde{\mathcal{H}}_c &\equiv \mathcal{H}_{-1} \oplus \mathcal{H}_{-3/2}\end{aligned}$$

$$\mathcal{G} = \left\{ \begin{array}{l} \mathbf{1} \text{ on } \mathcal{H}_{-1} \\ \mathcal{X}_0 \text{ on } \mathcal{H}_{-3/2} \end{array} \right.$$

$$S = -\frac{1}{2} \langle \tilde{\Psi}, Q\mathcal{G}\tilde{\Psi} \rangle + \langle \tilde{\Psi}, Q\Psi \rangle + \sum_{n=1}^{\infty} \frac{1}{n!} \{\Psi^n\}$$

$$\Omega_{n-3}^{o(0,n)}(A_1,\cdots,A_n)\equiv K^{-1}\widehat{\Omega}_{n-3}^{(0,1,0,n)}(A_1,\cdots,A_n)$$



$$\partial \mathcal{V}_{0,n}^o = -\frac{1}{2} \sum_{\substack{n_1, n_2 \\ n_1 + n_2 = n+2}} \{\mathcal{V}_{0,n_1}^o, \mathcal{V}_{0,n_2}^o\}_o$$

$$\{A_1,\cdots,A_n\}\equiv g_o^{n-2}\int_{\mathcal{V}_{0,n}^o}\Omega_{n-3}^{o(0,n)}(A_1,\cdots,A_n)$$

$$S_o=\frac{1}{2}\langle\psi_o,Q\psi_o\rangle'+\sum_{n=3}^{\infty}\frac{1}{n!}\{\psi_o^n\}$$

$$\mathcal{T}=\frac{1}{2\pi^2g_o^2}$$

$$\{A_1\cdots A_iA_{i+1}\cdots A_n\}=(-1)^{A_iA_{i+1}+1}\{A_1\cdots A_{i+1}A_i\cdots A_n\}.$$

$$\begin{aligned} Q[A'_1\cdots A'_N] &= \sum_{i=1}^N (-1)^{i-1} [A'_1\cdots A'_{i-1} (QA'_i) A'_{i+1}\cdots A'_N] \\ &\quad + \sum_{\substack{\ell,k\geqslant 0 \\ \ell+k=N}} \sum_{\substack{\{i_a;a=1,\ldots\ell\} \\ \{j_b;b=1,\ldots k\} \\ \{i_a\}\cup\{j_b\}=\{1,\ldots N\}}} \left[[A'_{j_1}\cdots A'_{j_k}] A'_{i_1}\cdots A'_{i_\ell} \right] \end{aligned}$$

$$\langle A_0,[A_1,\cdots,A_n]\rangle'=\{A_0,A_1,\cdots,A_n\}$$

$$\delta F = \langle F_R^o, \delta \psi_o \rangle' = \langle \delta \psi_o, F_L^o \rangle',$$

$$\{F,G\}=-\langle F_R^o,G_L^o\rangle'$$

$$\delta\Phi=Q\Lambda+[\Lambda\Phi]+\frac{1}{2!}[\Lambda,\Phi,\Phi]+\cdots$$

$$[A_1,A_2]=A_1\star A_2-(-1)^{A_1A_2}A_2\star A_1$$

$$A_1\star(A_2\star A_3)=(A_1\star A_2)\star A_3$$

$$[A_1,A_2,\cdots,A_n]=0\;\;\text{for}\;\;n\geqslant 3$$

$$\langle A_1,A_2\star A_3\rangle'=\langle f_1\circ A_1(0)f_2\circ A_2(0)f_3\circ A_3(0)\rangle'_{\text{UHP}}$$

$$\begin{aligned} f_1(w_1) &= h^{-1}\Big(e^{2\pi i/3}\big(h(w_1)\big)^{2/3}\Big) \\ f_2(w_2) &= h^{-1}\Big(\big(h(w_2)\big)^{2/3}\Big), h(u) \equiv \frac{1+iu}{1-iu} \\ f_3(w_3) &= h^{-1}\Big(e^{-2\pi i/3}\big(h(w_3)\big)^{2/3}\Big) \end{aligned}$$

$$\begin{aligned}\psi_o\in\mathcal{H}_o&\equiv\mathcal{H}_{-1}\oplus\mathcal{H}_{-1/2},\\ \tilde{\psi}_o\in\tilde{\mathcal{H}}_o&\equiv\mathcal{H}_{-1}\oplus\mathcal{H}_{-3/2}.\end{aligned}$$

$$S = -\frac{1}{2}\langle\tilde{\psi}_o, QG\tilde{\psi}_o\rangle' + \langle\tilde{\psi}_o, Q\psi_o\rangle' + \sum_{n=1}^{\infty} \frac{1}{n!}\{\psi_o^n\}$$

$$\begin{array}{l} \Psi_c \in \mathcal{H}_c, \tilde{\Psi}_c \in \tilde{\mathcal{H}}_c, \\ \Psi_o \in \mathcal{H}_o, \tilde{\Psi}_o \in \tilde{\mathcal{H}}_o \end{array}$$

$$\begin{aligned} \partial \mathcal{V}_{g,b,n_c,n_o} = & -\Delta_c \mathcal{V}_{g-1,b,n_c+2,n_o} - \Delta_o' \mathcal{V}_{g,b-1,n_c,n_o+2} - \Delta_o \mathcal{V}_{g-1,b+1,n_c,n_o+2} \\ & -\frac{1}{2} \sum_{\substack{g_1+g_2=g, b_1+b_2=b \\ n_{c1}+n_{c2}=n_c+2, n_{o1}+n_{o2}=n_o}} \{\mathcal{V}_{g_1,b_1,n_{c1},n_{o1}}, \mathcal{V}_{g_2,b_2,n_{c2},n_{o2}}\}_c \\ & -\frac{1}{2} \sum_{\substack{g_1+g_2=g, b_1+b_2=b+1 \\ n_{c1}+n_{c2}=n_c, n_{o1}+n_{o2}=n_o+2}} \{\mathcal{V}_{g_1,b_1,n_{c1},n_{o1}}, \mathcal{V}_{g_2,b_2,n_{c2},n_{o2}}\}_o \end{aligned}$$

$$\begin{aligned} \{A_1^c, \cdots, A_{n_c}^c; A_1^o, \cdots, A_{n_o}^o\} &= \sum_{g=0}^{\infty} \sum_{b=0}^{\infty} \{A_1^c, \cdots, A_{n_c}^c; A_1^o, \cdots, A_{n_o}^o\}_{g,b} \\ \{A_1^c, \cdots, A_{n_c}^c; A_1^o, \cdots, A_{n_o}^o\}_{g,b} &= (g_s)^{-\chi_{g,b,n_c,n_o}} \int_{\mathcal{V}_{g,b,n_c,n_o}} \Omega_{d_{g,b,n_c,n_o}}^{(g,b,n_c,n_o)}(A_1^c, \cdots, A_{n_c}^c; A_1^o, \cdots, A_{n_o}^o) \end{aligned}$$

$$\{\tilde{A}^c\}_D=\Omega_0^{(0,1,1,0)}(\tilde{A}^c)\equiv -\frac{1}{2\pi i}N_{0,1,1,0}\langle\hat{\mathcal{G}}\tilde{A}^c|c_0^-e^{-\Lambda(L_0+\bar{L}_0)}|B\rangle,$$

$$\hat{\mathcal{G}}\equiv\begin{cases}1\text{ on }\mathcal{H}_{-1,-1}\\ \frac{1}{2}(\mathcal{X}_0+\overline{\mathcal{X}_0})\text{ on }\mathcal{H}_{-3/2,-3/2}\end{cases}$$

$$\{A^c\}_D=\Omega_0^{(0,1,1,0)}(A^c)\equiv -\frac{1}{2\pi i}N_{0,1,1,0}\langle A^c|c_0^-e^{-\Lambda(L_0+\bar{L}_0)}|B\rangle.$$

$$\begin{aligned} \langle A_0^c|c_0^-|[A_1^c\cdots A_N^c;A_1^o\cdots A_M^o]_c\rangle &= \{A_0^cA_1^c\cdots A_N^c;A_1^o\cdots A_M^o\}, \forall |A_0^c\rangle\in\mathcal{H}_c, \\ \langle A_0^o|[A_1^c\cdots A_N^c;A_1^o\cdots A_M^o]_o\rangle &= \{A_1^c\cdots A_N^c;A_0^oA_1^o\cdots A_M^o\}, \forall |A_0^o\rangle\in\mathcal{H}_o, \\ \langle \tilde{A}^c|c_0^-|\square_c\rangle &= \{\tilde{A}^c\}_D, (4.65) \end{aligned}$$

$$\{(Q\tilde{A}^c)\}_D=0$$

$$\begin{aligned} &\sum_{i=1}^N\{A_1^c\cdots A_{i-1}^c(QA_i^c)A_{i+1}^c\cdots A_N^c;A_1^o\cdots A_M^o\}+\sum_{j=1}^M\{A_1^c\cdots A_N^c;A_1^o\cdots A_{j-1}^o(QA_j^o)A_{j+1}^o\cdots A_M^o\}(-1)^{j-1}\\ &=-\frac{1}{2}\sum_{k=0}^N\sum_{\ell=0}^M\sum_{\substack{\{i_1,\cdots,i_k\}\subset\{1,\cdots,N\}\\ \{j_1,\cdots,j_\ell\}\subset\{1,\cdots,M\}}}(\{A_{i_1}^c\cdots A_{i_k}^c\mathcal{B}^c;A_{j_1}^o\cdots A_{j_\ell}^o\}+\{A_{i_1}^c\cdots A_{i_k}^c;\mathcal{B}^oA_{j_1}^o\cdots A_{j_\ell}^o\})\\ &\quad -\{[A_1^c\cdots A_N^c;A_1^o\cdots A_M^o]_c\}_D-\frac{1}{2}\{A_1^c\cdots A_N^c\varphi_s\varphi_r;A_1^o\cdots,A_M^o\}\langle\varphi_s^c,\mathcal{G}\varphi_r^c\rangle\\ &\quad -\frac{1}{2}(-1)^{\hat{\varphi}_s}\{A_1^c\cdots A_N^c;\hat{\varphi}_s\hat{\varphi}_rA_1^o\cdots,A_M^o\}\langle\hat{\varphi}_s^c,\mathcal{G}\hat{\varphi}_r^c\rangle \end{aligned}$$

$$\begin{aligned}
\mathcal{B}^c &\equiv \mathcal{G}[A_{i_1}^c \cdots A_{i_{N-k}}^c; A_{j_1}^o \cdots A_{j_{M-\ell}}^o]_c, \mathcal{B}^o \equiv \mathcal{G}[A_{i_1}^c \cdots A_{i_{N-k}}^c; A_{j_1}^o \cdots A_{j_{M-\ell}}^o]_o \\
\{i_1, \dots, i_k\} \cup \{\bar{i}_1, \dots, \bar{i}_{N-k}\} &= \{1, \dots, N\}, \{j_1, \dots, j_\ell\} \cup \{\bar{j}_1, \dots, \bar{j}_{M-\ell}\} = \{1, \dots, M\}, \\
S &= -\frac{1}{2} \langle \tilde{\Psi}_c, Q \mathcal{G} \tilde{\Psi}_c \rangle + \langle \tilde{\Psi}_c, Q \Psi_c \rangle - \frac{1}{2} \langle \tilde{\Psi}_o, Q \mathcal{G} \tilde{\Psi}_o \rangle + \langle \tilde{\Psi}_o, Q \Psi_o \rangle + \{\tilde{\Psi}_c\}_D + \sum_{N=0}^{\infty} \sum_{M=0}^{\infty} \frac{1}{N! M!} \{(\Psi_c)^N; (\Psi_o)^M\} \\
\delta F &= \langle F_R^c, \delta \tilde{\Psi}_c \rangle + \langle \tilde{F}_R^c, \delta \Psi_c \rangle + \langle F_R^o, \delta \tilde{\Psi}_o \rangle + \langle \tilde{F}_R^o, \delta \Psi_o \rangle = \langle \delta \tilde{\Psi}_c, F_L^c \rangle + \langle \delta \Psi_c, \tilde{F}_L^c \rangle + \langle \delta \tilde{\Psi}_o, F_L^o \rangle + \langle \delta \Psi_o, \tilde{F}_L^o \rangle \\
\{F, G\} &= -(\langle F_R^c, \tilde{G}_L^c \rangle + \langle \tilde{F}_R^c, G_L^c \rangle + \langle \tilde{F}_R^c, \mathcal{G} \tilde{G}_L^c \rangle) - (\langle F_R^o, \tilde{G}_L^o \rangle + \langle \tilde{F}_R^o, G_L^o \rangle + \langle \tilde{F}_R^o, \mathcal{G} \tilde{G}_L^o \rangle) \\
S &= \frac{1}{2} \langle \Psi_c, Q \Psi_c \rangle + \frac{1}{2} \langle \Psi_o, Q \Psi_o \rangle + \{\Psi_c\}_D + \sum_{N=0}^{\infty} \sum_{M=0}^{\infty} \frac{1}{N! M!} \{(\Psi_c)^N; (\Psi_o)^M\} \\
S_{oc} &= \frac{1}{2} \langle \Psi_o, Q \Psi_o \rangle + \sum_{n=3}^{\infty} \frac{1}{n!} g_s^{(n-2)/2} \int \Omega_{n-3}^{(0,1,0,n)}(\Psi_o, \dots, \Psi_o) \\
&= \frac{1}{2} K \langle \Psi_o, Q \Psi_o \rangle' + \sum_{n=3}^{\infty} \frac{1}{n!} g_s^{(n-2)/2} N_{0,1,0,n} K \int \Omega_{n-3}^{o(0,n)}(\Psi_o, \dots, \Psi_o) \\
S_o &= \frac{1}{2} \langle \psi_o, Q \psi_o \rangle' + \sum_{n=3}^{\infty} \frac{1}{n!} g_o^{n-2} \int \Omega_{n-3}^{o(0,n)}(\psi_o, \dots, \psi_o) \\
\Psi_o &= K^{-1/2} \psi_o \\
S_{oc} &= \frac{1}{2} \langle \psi_o, Q \psi_o \rangle' + \sum_{n=3}^{\infty} \frac{1}{n!} g_s^{(n-2)/2} K^{(2-n)/2} \eta_c^{3(n-2)/4} \int \Omega_{n-3}^{o(0,n)}(\psi_o, \dots, \psi_o) \\
g_o &= g_s^{1/2} K^{-1/2} \eta_c^{3/4} \\
g_o^2 &= (g_s/K) \eta_c^{3/2}
\end{aligned}$$

Partículas supermasivas – campos de gauge – transformaciones de gauge – campos cuánticos gravitacionales y simetría – osciladores y propagadores.

$$\begin{aligned}
|\psi_o\rangle &= (\alpha')^{(p+1)/2} \int \frac{d^{p+1}k}{(2\pi)^{p+1}} \left(A_\mu(k) c_1 \alpha_{-1}^\mu - i \sqrt{\frac{1}{2}} B(k) c_0 \right) |k\rangle, \\
|\epsilon\rangle &= (\alpha')^{(p+1)/2} \int \frac{d^{p+1}k}{(2\pi)^{p+1}} \frac{i}{\sqrt{2}} \epsilon(k) |k\rangle. \\
\phi(x) &= (\alpha')^{(p+1)/2} \int \frac{d^{p+1}k}{(2\pi)^{p+1}} \tilde{\phi}(k) e^{ikx}, \text{ so that } ik_\mu \leftrightarrow \partial_\mu
\end{aligned}$$



$$\text{bpz}(\phi_n) = (-1)^{n+d} \phi_{-n}$$

$$\langle \psi_o | = (\alpha')^{(p+1)/2} \int \frac{d^{p+1}k}{(2\pi)^{p+1}} \langle k | \left(A_\mu(k) c_{-1} \alpha_1^\mu + i \sqrt{\frac{1}{2}} B(k) c_0 \right),$$

$$Q=\sum_n c_n(L_{-n}^m-\delta_{n,0})+\frac{1}{2}\sum_{m,n} (m-n)\colon c_m c_n b_{-n-m}\colon$$

$$L_n^m=\frac{1}{2}\sum_m\alpha_m\cdot\alpha_{n-m},(n\neq0),L_0^m=\frac{1}{2}\alpha_0^2+\sum_{n\geqslant1}\alpha_{-n}\cdot\alpha_n,\alpha_0^\mu=\sqrt{2\alpha'}k^\mu$$

$$Q=c_0(\alpha'k^2-1)+c_0(\alpha_{-1}\cdot\alpha_1+b_{-1}c_1+c_{-1}b_1)+\sqrt{2\alpha'}k_\mu(\alpha_{-1}^\mu c_1+c_{-1}\alpha_1^\mu)-2b_0c_{-1}c_1+\cdots$$

$$Qc_1\alpha_{-1}^\mu|k\rangle=(\alpha'k^2c_0c_1\alpha_{-1}^\mu+\sqrt{2\alpha'}k^\mu c_{-1}c_1)|k\rangle$$

$$Qc_0|k\rangle=(\sqrt{2\alpha'}k_\mu\alpha_{-1}^\mu c_1c_0-2c_{-1}c_1)|k\rangle$$

$$Q|k\rangle=(\alpha'k^2c_0+\sqrt{2\alpha'}k_\mu\alpha_{-1}^\mu c_1)|k\rangle$$

$$\langle k|c_{-1}c_0c_1|k'\rangle' = -(2\pi)^{p+1}(\alpha')^{-(p+1)/2}\delta^{(p+1)}(k+k').$$

$$S_2=\frac{1}{2}\langle\psi_o|Q|\psi_o\rangle'$$

$$=(\alpha')^{(p+1)/2}\int\frac{d^{p+1}k}{(2\pi)^{p+1}}\Big(-\frac{1}{2}A^\mu(-k)\alpha'p^2A_\mu(k)-\sqrt{\alpha'}A^\mu(-k)ik_\mu B(k)\\-\frac{1}{2}B(-k)B(k)\Big)$$

$$S_2=(\alpha')^{-(p+1)/2}\int d^{p+1}x\left(\frac{\alpha'}{2}A^\mu\Box A_\mu-\sqrt{\alpha'}A^\mu\partial_\mu B-\frac{1}{2}B^2\right)$$

$$\delta A_\mu(k)=i\sqrt{\alpha'}k_\mu\epsilon(k),\delta B(k)=-\alpha'k^2\epsilon(k)$$

$$\delta A_\mu=\sqrt{\alpha'}\partial_\mu\epsilon,\delta B=\alpha'\Box\epsilon$$

$$S=(\alpha')^{-(p-1)/2}\int d^{p+1}x\Big(\frac{1}{2}A_\mu\Box A^\mu+\frac{1}{2}(\partial\cdot A)^2\Big)=(\alpha')^{-(p-1)/2}\int d^{p+1}$$

$$x\left(-\frac{1}{2}\partial_\mu A_\nu\partial^\mu A^\nu+\frac{1}{2}\partial_\mu A_\nu\partial^\nu A^\mu\right)=(\alpha')^{-(p-1)/2}\int d^{p+1}x\left(-\frac{1}{4}F^{\mu\nu}F_{\mu\nu}\right)$$

$$S_{\rm ws}=-\frac{1}{4\pi\alpha'}\int dxdyg_{\mu\nu}(\partial_xX^\mu\partial_xX^\nu+\partial_yX^\mu\partial_yX^\nu)=-\frac{1}{\pi\alpha'}\int dxdyg_{\mu\nu}\partial X^\mu\bar\partial X^\nu,z\equiv x+iy$$



$$g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu}$$

$$S_{\text{ws}}|_h = -\frac{1}{\pi\alpha'} \int dx dy h_{\mu\nu} \partial X^\mu \bar{\partial} X^\nu \\ -g_s \frac{1}{\pi} \int dx dy V$$

$$\mathcal{O}_h = \frac{1}{g_s \alpha'} h_{\mu\nu} c \bar{c} \partial X^\mu \bar{\partial} X^\nu = \frac{1}{g_s} \left(-\frac{1}{2} h_{\mu\nu} \right) c \bar{c} i \sqrt{\frac{2}{\alpha'}} \partial X^\mu i \sqrt{\frac{2}{\alpha'}} \bar{\partial} X^\nu$$

$$|\Psi\rangle = (\alpha')^{D/2} \frac{1}{g_s} \int \frac{d^D k}{(2\pi)^D} \left(-\frac{1}{2} e_{\mu\nu}(k) \alpha_{-1}^\mu \bar{\alpha}_{-1}^\nu c_1 \bar{c}_1 + e(k) c_1 c_{-1} + \bar{e}(k) \bar{c}_1 \bar{c}_{-1} \right. \\ \left. + i \sqrt{\frac{1}{2}} (f_\mu(k) c_0^+ c_1 \alpha_{-1}^\mu + \bar{f}_\mu(k) c_0^+ \bar{c}_1 \bar{\alpha}_{-1}^\mu) \right) |k\rangle$$

$$e_{\mu\nu} = h_{\mu\nu} + b_{\mu\nu}, \text{ with } h_{\mu\nu} = h_{\nu\mu}, b_{\mu\nu} = -b_{\nu\mu}$$

$$\langle k|c_{-1}\bar{c}_{-1}c_0^-c_0^+c_1\bar{c}_1|k'\rangle = -\frac{1}{2}(\alpha')^{-D/2}(2\pi)^D\delta^{(D)}(k+k').$$

$$\text{bpz}(\phi_n) = (-1)^d \phi_{-n},$$

$$\langle \Psi | = (\alpha')^{D/2} \frac{1}{g_s} \int \frac{d^D k}{(2\pi)^D} \langle k | \left(-\frac{1}{2} e_{\mu\nu}(k) \alpha_1^\mu \bar{\alpha}_1^\nu c_{-1} \bar{c}_{-1} + e(k) c_{-1} c_1 + \bar{e}(k) \bar{c}_{-1} \bar{c}_1 \right. \\ \left. - i \sqrt{\frac{1}{2}} (f_\mu(k) c_0^+ c_{-1} \alpha_1^\mu + \bar{f}_\mu(k) c_0^+ \bar{c}_{-1} \bar{\alpha}_1^\mu) \right)$$

$$S^{(2)} = \frac{1}{2} \langle \Psi | c_0^- Q | \Psi \rangle$$

$$Q = c_0^+ \left(\frac{\alpha'}{2} k^2 - 2 \right) + c_0^+ (\alpha_{-1} \cdot \alpha_1 + b_{-1} c_1 + c_{-1} b_1 + \bar{\alpha}_{-1} \cdot \bar{\alpha}_1 + \bar{b}_{-1} \bar{c}_1 + \bar{c}_{-1} \bar{b}_1) + \sqrt{\frac{\alpha'}{2}} k$$

$$\cdot (\alpha_{-1} c_1 + c_{-1} \alpha_1) + \sqrt{\frac{\alpha'}{2}} k \cdot (\bar{\alpha}_{-1} \bar{c}_1 + \bar{c}_{-1} \bar{\alpha}_1) - b_0^+ (c_{-1} c_1 + \bar{c}_{-1} \bar{c}_1) + \cdots$$

$$S^{(2)} = (\alpha')^{-D/2} \frac{1}{8g_s^2} \int d^D x \left[\frac{\alpha'}{4} e^{\mu\nu} \square e_{\mu\nu} + 2\alpha' \bar{e} \square e - f^\mu f_\mu - \bar{f}^\mu \bar{f}_\mu \right. \\ \left. - \sqrt{\alpha'} f^\mu (\partial^\nu e_{\mu\nu} - 2\partial_\mu \bar{e}) + \sqrt{\alpha'} \bar{f}^\nu (\partial^\mu e_{\mu\nu} + 2\partial_\nu e) \right]$$

$$|\Lambda\rangle = (\alpha')^{D/2} \frac{1}{g_s} \int \frac{d^D k}{(2\pi)^D} \left(\frac{i}{\sqrt{2}} \lambda_\mu(k) \alpha_{-1}^\mu c_1 - \frac{i}{\sqrt{2}} \bar{\lambda}_\mu(k) \bar{\alpha}_{-1}^\mu \bar{c}_1 + \mu(k) c_0^+ \right) |k\rangle.$$

$$\delta|\Psi\rangle = Q|\Lambda\rangle$$



$$\delta e_{\mu\nu} = \sqrt{\alpha'} (\partial_\mu \bar{\lambda}_\nu + \partial_\nu \lambda_\mu)$$

$$\delta f_\mu = -\frac{1}{2} \alpha' \square \lambda_\mu + \sqrt{\alpha'} \partial_\mu \mu$$

$$\delta \bar{f}_\nu = \frac{1}{2} \alpha' \square \bar{\lambda}_\nu + \sqrt{\alpha'} \partial_\nu \mu$$

$$\delta e = -\frac{1}{2} \sqrt{\alpha'} \partial \cdot \lambda + \mu$$

$$\delta \bar{e} = \frac{1}{2} \sqrt{\alpha'} \partial \cdot \bar{\lambda} + \mu$$

$$d = \frac{1}{2} (e - \bar{e}), \text{ and } \chi = \frac{1}{2} (e + \bar{e})$$

$$\delta d = -\frac{1}{4} \sqrt{\alpha'} (\partial \cdot \lambda + \partial \cdot \bar{\lambda})$$

$$\delta \chi = -\frac{1}{4} \sqrt{\alpha'} (\partial \cdot \lambda - \partial \cdot \bar{\lambda}) + \mu$$

$$\chi = 0$$

$$f_\mu = -\frac{1}{2} \sqrt{\alpha'} (\partial^\nu e_{\mu\nu} - 2\partial_\mu \bar{e}), \bar{f}_\nu = \frac{1}{2} \sqrt{\alpha'} (\partial^\mu e_{\mu\nu} + 2\partial_\nu e)$$

$$S^{(2)} = (\alpha')^{-(D-2)/2} \frac{1}{8g_s^2} \int d^D x \left[\frac{1}{4} e_{\mu\nu} \square e^{\mu\nu} + \frac{1}{4} (\partial^\nu e_{\mu\nu})^2 + \frac{1}{4} (\partial^\mu e_{\mu\nu})^2 - 2d\partial^\mu \partial^\nu e_{\mu\nu} - 4d \square d \right]$$

$$\delta e_{\mu\nu} = \sqrt{\alpha'} (\partial_\nu \lambda_\mu + \partial_\mu \bar{\lambda}_\nu)$$

$$\delta d = -\frac{1}{4} \sqrt{\alpha'} (\partial \cdot \lambda + \partial \cdot \bar{\lambda})$$

$$S^{(2)} = (\alpha')^{-(D-2)/2} \frac{1}{8g_s^2} \int d^D x L[h, b, d]$$

$$L[h, b, d] = \frac{1}{4} h^{\mu\nu} \partial^2 h_{\mu\nu} + \frac{1}{2} (\partial^\nu h_{\mu\nu})^2 - 2d\partial^\mu \partial^\nu h_{\mu\nu} - 4d\partial^2 d + \frac{1}{4} b^{\mu\nu} \partial^2 b_{\mu\nu} + \frac{1}{2} (\partial^\nu b_{\mu\nu})^2$$

$$S_{\text{st}} = \frac{1}{2\kappa^2} \int d^D x \sqrt{-g} e^{-2\phi} \left[R - \frac{1}{12} H^2 + 4(\partial\phi)^2 \right]$$

$$g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu}, \phi = d + \frac{1}{4} \eta^{\mu\nu} h_{\mu\nu}, H_{\mu\nu\rho} = \partial_\mu b_{\nu\rho} + \dots$$

$$S_{\text{st}}^{(2)} = \frac{1}{2\kappa^2} \int d^D x L[h, b, d]$$

$$\kappa = (\alpha')^{(D-2)/4} (2g_s)$$



$$\begin{aligned}\delta h_{\mu\nu} &= \partial_\mu \epsilon_\nu + \partial_\nu \epsilon_\mu \\ \delta b_{\mu\nu} &= -\partial_\mu \tilde{\epsilon}_\nu + \partial_\nu \tilde{\epsilon}_\mu \\ \delta d &= -\frac{1}{2} \partial \cdot \epsilon\end{aligned}\quad (4.124)$$

$$\epsilon_\mu \equiv \frac{1}{2} \sqrt{\alpha'} (\lambda_\mu + \bar{\lambda}_\mu), \tilde{\epsilon}_\mu \equiv \frac{1}{2} \sqrt{\alpha'} (\lambda_\mu - \bar{\lambda}_\mu)$$

$$b_0^+|\Psi\rangle=0$$

$$S=\frac{1}{2}\langle\Psi|c_0^-c_0^+L_0^+|\Psi\rangle+\sum_{n=1}^{\infty}\frac{1}{n!}\{\Psi^n\}$$

$$\mathcal{P}_b=-b_0^+b_0^-(L_0+\bar{L}_0)^{-1}\delta_{L_0,\bar{L}_0}$$

$$\mathcal{P}_b=-\frac{1}{2\pi}b_0^+b_0^-\int_0^{2\pi}d\theta\int_0^\infty ds e^{-s(L_0+\bar{L}_0)}e^{i\theta(L_0-\bar{L}_0)}=\frac{1}{\pi}b_0\bar{b}_0\int_{|q|\leqslant 1}\frac{d^2q}{|q|^2}q^{L_0}\bar{q}^{\bar{L}_0}$$

$$q\equiv e^{-s+i\theta}, d^2q=d\theta ds|q|^2\equiv\frac{i}{2}dq\wedge d\bar{q}$$

$$w_1w_2=q$$

$$(6g_1-6+2n_1)+(6g_2-6+2n_2)+2$$

$$\mathcal{P}_b=\int_{|q|\leqslant 1}\left[\left(-\frac{1}{2\pi i}\right)b_0\bar{b}_0\frac{dq\wedge d\bar{q}}{|q|^2}\right]q^{L_0}\bar{q}^{\bar{L}_0}$$

$$\Omega_{\mathcal{P}_b}=-\frac{1}{2\pi i}b_0\bar{b}_0\frac{dq\wedge d\bar{q}}{|q|^2}$$

$$\frac{\partial F}{\partial q}=\frac{1}{w_2}=\frac{w_1}{q},\frac{\partial \bar{F}}{\partial q}=0;\frac{\partial F}{\partial \bar{q}}=0,\frac{\partial \bar{F}}{\partial \bar{q}}=\frac{\bar{w}_1}{\bar{q}}.$$

$$\widehat{\Omega}_b=\mathcal{B}\left[\frac{\partial}{\partial q}\right]\mathcal{B}\left[\frac{\partial}{\partial \bar{q}}\right]dq\wedge d\bar{q}=\frac{1}{q}\oint\ w_1b(w_1)dw_1\times\frac{1}{\bar{q}}\oint\ \bar{w}_1\bar{b}(\bar{w}_1)d\bar{w}_1\times dq\wedge d\bar{q}=\frac{b_0}{q}\frac{\bar{b}_0}{\bar{q}}dq\wedge d\bar{q}$$

$$\Omega_{\mathcal{P}_b}=-\frac{1}{2\pi i}\hat{\Omega}_b$$

$$\left(-\frac{1}{2\pi i}\right)^{3g_1-3+n_1}\left(-\frac{1}{2\pi i}\right)^{3g_2-3+n_2}\left(-\frac{1}{2\pi i}\right)^1=\left(-\frac{1}{2\pi i}\right)^{3(g_1+g_2)-3+(n_1+n_2-2)}$$

$$S=-\frac{1}{2}\langle\tilde{\Psi}|c_0^-c_0^+L_0^+\mathcal{G}|\tilde{\Psi}\rangle+\langle\tilde{\Psi}|c_0^-c_0^+L_0^+|\Psi\rangle+\sum_{n=1}^{\infty}\frac{1}{n!}\{\Psi^n\}$$

$$K_0=c_0^-c_0^+L_0^+\left(\begin{smallmatrix}\mathcal{G}&-1\\-1&0\end{smallmatrix}\right)$$

$$\mathcal{P}_s=-b_0^+b_0^-(L_0+\bar{L}_0)^{-1}\delta_{L_0,\bar{L}_0}\left(\begin{smallmatrix}0&1\\1&\mathcal{G}\end{smallmatrix}\right)$$

$$\mathcal{P}_o = -b_o L_0^{-1} = -b_o \int_0^1 \frac{dq_o}{q_o} q_o^{L_0} = \int_0^1 \left[(-b_o) \frac{dq_o}{q_o} \right] q_o^{L_0}$$

$$\Omega_{\mathcal{P}_o} = -b_o \frac{dq_o}{q_o}$$

$$w_1 w_2 = -q_o, q_o \in [0,1]$$

$$\widehat{\Omega}_o = \mathcal{B} \left[\frac{\partial}{\partial q_o} \right] dq_o$$

$$\mathcal{B} \left[\frac{\partial}{\partial q_o} \right] = \int dw_1 b(w_1) \frac{w_1}{q_o} + \int d\bar{w}_1 \bar{b}(\bar{w}_1) \frac{\bar{w}_1}{q_o} = \frac{1}{q_o} \oint dw_1 b(w_1) w_1$$

$$\mathcal{B} \left[\frac{\partial}{\partial q_o} \right] = -\frac{1}{q_o} b_o, \widehat{\Omega}_o = -b_o \frac{dq_o}{q_o}$$

$$\Omega_{\mathcal{P}_o} = \widehat{\Omega}_o$$

$$N_{g_1,b_1,n_{c1},n_{o1}} N_{g_2,b_2,n_{c2},n_{o2}} \sim N_{g_1+g_2,b_1+b_2-1,n_{c1}+n_{c2},n_{o1}+n_{o2}-2}$$

$$N_{g,b,n_c,n_o} \sim N_{g,b+1,n_c,n_o+2}, N_{g,b,n_c,n_o} \sim N_{g+1,b-1,n_c,n_o+2}$$

$$\eta_c N_{g_1,b_1,n_{c1},n_{o1}} N_{g_2,b_2,n_{c2},n_{o2}} \sim N_{g_1+g_2,b_1+b_2,n_{c1}+n_{c2}-2,n_{o1}+n_{o2}} \\ \eta_c N_{g-1,b,n_c+2,n_o} \sim N_{g,b,n_c,n_o}$$

Membranas.

$$\mathcal{O}_h = \frac{1}{g_s} h_{\mu\nu} c \bar{c} \partial X^\mu \bar{\partial} X^\nu$$

$$g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu}, \phi = \frac{1}{4} \eta^{\mu\nu} h_{\mu\nu}$$

$$S_p = -\mathcal{T} \int d^{p+1} x e^{-\phi} \sqrt{-g}$$

$$S_p|_h = -\mathcal{T} \left\{ \frac{1}{2} \eta^{\mu\nu} h_{\mu\nu}^{\parallel}(x_{\perp}=0) - \phi(x_{\perp}=0) \right\} \int d^{p+1} x = -\frac{\mathcal{T}}{4} \eta^{\mu\nu} \left(h_{\mu\nu}^{\parallel} - h_{\mu\nu}^{\perp} \right) (2\pi)^{p+1} \delta^{(p+1)}(0)$$

$$\Omega_0^{(0,1,1,0)}(A) = -\frac{1}{2\pi i} \eta_c^{-1/2} \langle c_0^- A \rangle = \eta_c^{1/2} \langle c_0^- A \rangle$$

$$\Omega_0^{(0,1,1,0)}(\mathcal{O}_h) = \eta_c^{1/2} \frac{1}{g_s} h_{\mu\nu} \langle c_0^- c \bar{c} \partial X^\mu \bar{\partial} X^\nu \rangle$$

$$\Omega_0^{(0,1,1,0)}(\mathcal{O}_h) = \eta_c^{1/2} \frac{1}{g_s} \left(h_{\mu\nu}^{\parallel} - h_{\mu\nu}^{\perp} \right) \frac{K}{2} \eta^{\mu\nu} (2\pi)^{p+1} \delta^{(p+1)}(0)$$

$$K = -g_s \frac{\mathcal{T}}{2\sqrt{\eta}_c}$$



$$g_o^2 = -\frac{2}{\mathcal{T}}\eta_c^2 = \frac{1}{2\pi^2\mathcal{T}} \rightarrow \mathcal{T}g_o^2 = \frac{1}{2\pi^2}$$

$$-\frac{g_s}{\pi}\int dy_R dy_IV_c(y)$$

$$g_o\int dxV_o(x)$$

$$-g_s\frac{\mathcal{T}}{2}\langle c_0^-c\bar{c}V_c\rangle'$$

$$-i\pi\mathcal{T}g_sg_o\langle c\bar{c}V_ccV_o\rangle'$$

$$\frac{i}{2}g_s^2\mathcal{T}\int dy\left\langle c\bar{c}V_c^{(1)}(i)(c+\bar{c})V_c^{(2)}(iy)\right\rangle'$$

$$w_1w_2=q, |q|\leqslant 1$$

$$\mathcal{V}_{0,n}^{1\mathrm{PI}}=\mathcal{V}_{0,n}$$

$$\mathcal{V}_{1,1}^{1\mathrm{PI}}=\mathcal{F}_{1,1}$$

$$\mathcal{V}_{1\mathrm{PI}}\equiv\sum_{g,n}\mathcal{V}_{g,n}^{1\mathrm{PI}}$$

$$\partial \mathcal{V}_{1\mathrm{PI}}+\frac{1}{2}\{\mathcal{V}_{1\mathrm{PI}},\mathcal{V}_{1\mathrm{PI}}\}=0$$

$$\{A_1,\cdots,A_n\}_{1\mathrm{PI}}=\sum_{g=0}^{\infty}g_s^{2g+n-2}\int_{\mathcal{V}_{g,n}^{1\mathrm{PI}}}\Omega_{6g-6+2n}^{(g,n)}(A_1,\cdots,A_n)$$

$$S_{1\mathrm{PI}}=-\frac{1}{2}\langle\widetilde{\Psi},Q\mathcal{G}\widetilde{\Psi}\rangle+\langle\widetilde{\Psi},Q\Psi\rangle+\sum_{n=1}^{\infty}\frac{1}{n!}\{\Psi^n\}_{1\mathrm{PI}}$$

$$\delta|\widetilde{\Psi}\rangle=Q|\widetilde{\Lambda}\rangle+\sum_{n=1}^{\infty}\frac{1}{n!}[\Lambda\Psi^n]_{1\mathrm{PI}},\delta|\Psi\rangle=Q|\Lambda\rangle+\sum_{n=1}^{\infty}\frac{1}{n!}\mathcal{G}[\Lambda\Psi^n]_{1\mathrm{PI}}$$

$$\{S_{1\mathrm{PI}},S_{1\mathrm{PI}}\}=0$$

Acción Wilsoniana.

$$[P,L_0^\pm]=0,[P,b_0^\pm]=0,[P,c_0^\pm]=0,[P,Q]=0,[P,\mathcal{G}]=0$$

$$\{A_1,\cdots,A_n\}_{\mathrm{eff}},A_1,\cdots,A_n\in P\mathcal{H}_c.$$

$$\Psi\in P\mathcal{H}_c,\tilde{\Psi}\in P\tilde{\mathcal{H}}_c$$

$$S_{\mathrm{eff}}=-\frac{1}{2}\langle\widetilde{\Psi},Q\mathcal{G}\widetilde{\Psi}\rangle+\langle\widetilde{\Psi},Q\Psi\rangle+\sum_n\frac{1}{n!}\{\Psi^n\}_{\mathrm{eff}}$$



$$\begin{aligned}\mathcal{V}^{1\text{PI}} &= \sum_{g,n} g_s^{2g-n+2} \mathcal{V}_{g,n}^{1\text{PI}} \\ \partial \mathcal{V}^{1\text{PI}} + \frac{1}{2} \{ \mathcal{V}^{1\text{PI}}, \mathcal{V}^{1\text{PI}} \} &= 0 \\ S'_{1\text{PI}} - S_{1\text{PI}} &= \sum_{g=0}^{\infty} \sum_{n=1}^{\infty} g_s^{2g-n+2} \left(\int_{\mathcal{V}_{g,n}^{1\text{PI}}} - \int_{\mathcal{V}_{g,n}^{1\text{PI}}} \right) \Omega_{6g-6+2n}^{(g,n)}(\Psi^{\otimes n}) \\ \langle \Phi | c_0^- | \delta \Psi \rangle &= - \sum_{g=0}^{\infty} \sum_{n=1}^{\infty} g_s^{2g-n+2} \frac{1}{(n-1)!} \int_{\mathcal{V}_{g,n}^{1\text{PI}}} \Omega_{6g-5+2n}^{(g,n)}(\Phi, \Psi^{\otimes(n-1)}) [\hat{U}_{g,n}]\end{aligned}$$

Teorema del dilatón.

$$\begin{aligned}D(z, \bar{z}) &= \frac{1}{2} (c \partial^2 c - \bar{c} \partial^2 \bar{c}) \\ |D\rangle &= (c_1 c_{-1} - \bar{c}_1 \bar{c}_{-1}) |0\rangle \\ b_0^- |D\rangle &= 0, L_0^- |D\rangle = 0 \\ |D\rangle &= Q c_0^- |0\rangle = -Q |\chi\rangle, |\chi\rangle = -c_0^- |0\rangle \\ L_1 |D\rangle &= c_0 c_1 |0\rangle, \bar{L}_1 |D\rangle = -\bar{c}_0 \bar{c}_1 |0\rangle \\ \Omega_2(D) &= \left(-\frac{1}{2\pi i}\right) du^1 \wedge du^2 \mathcal{B} \left[\frac{\partial}{\partial u^1}\right] \mathcal{B} \left[\frac{\partial}{\partial u^2}\right] D(w=0) \\ \Omega_2(D) &= -d\Omega_1(\chi) \\ \Omega_1(\chi) &= \left(-\frac{1}{2\pi i}\right) \left(du^1 \mathcal{B} \left[\frac{\partial}{\partial u^1}\right] + du^2 \mathcal{B} \left[\frac{\partial}{\partial u^2}\right]\right) \chi(w=0) \\ \int_M \Omega_2(D) + \int_{\partial M} \Omega_1(\chi) \\ \chi(M) &= \frac{1}{2\pi} \int_M K^{(2)} + \frac{1}{2\pi} \int_{\partial M} k^{(1)} \\ K^{(2)} &= -2i\partial\bar{\partial}\rho dz \wedge d\bar{z} \\ k^{(1)} &= d\theta_\gamma - i[dz\partial\log\rho - d\bar{z}\bar{\partial}\log\rho] \\ \Omega_2(D) &= -\frac{1}{2\pi} K^{(2)} \text{ and } \Omega_1(\chi) = -\frac{1}{2\pi} k^{(1)} \\ [Q, \mathcal{P}_b] &= \frac{1}{2\pi} b_0^- \int_0^{2\pi} d\theta \int_0^\infty ds \frac{\partial}{\partial s} e^{-s(L_0 + \bar{L}_0)} e^{i\theta(L_0 - \bar{L}_0)} \\ z &= F(w; u) = F(w; u^1, u^2)\end{aligned}$$



$$y(u) = F(0; u)$$

$$z = F(w; u^1, u^2) = y(u) + a(u)w + \frac{1}{2}b(u)w^2 + \frac{1}{3!}c(u)w^3 + \mathcal{O}(w^4)$$

$$\begin{aligned}\mathcal{B}\left[\frac{\partial}{\partial u^i}\right] &= \oint b(z)dz \frac{\partial F}{\partial u^i} + \oint \bar{b}(\bar{z})d\bar{z} \frac{\partial \bar{F}}{\partial u^i} \\ &= -\oint b(w)dw \left(\frac{\partial F}{\partial w}\right)^{-1} \frac{\partial F}{\partial u^i} - \oint \bar{b}(\bar{w})d\bar{w} \left(\frac{\partial \bar{F}}{\partial \bar{w}}\right)^{-1} \frac{\partial \bar{F}}{\partial u^i}\end{aligned}$$

$$\left(\frac{\partial F}{\partial w}\right)^{-1} \frac{\partial F}{\partial u^i} = \alpha_i + \beta_i w + \gamma_i w^2 + \mathcal{O}(w^3)$$

$$\begin{aligned}\alpha_i &= \frac{1}{a} \frac{\partial y}{\partial u^i} \\ \beta_i &= \frac{1}{a} \frac{\partial a}{\partial u^i} - \frac{b}{a^2} \frac{\partial y}{\partial u^i} \\ \gamma_i &= a \frac{\partial}{\partial u^i} \left(\frac{b}{2a^2} \right) + \left(\frac{b^2}{a^3} - \frac{1}{2} \frac{c}{a^2} \right) \frac{\partial y}{\partial u^i}\end{aligned}$$

$$\mathcal{B}\left[\frac{\partial}{\partial u^i}\right] = -(\alpha_i b_{-1} + \beta_i b_0 + \gamma_i b_1 + \bar{\alpha}_i \bar{b}_{-1} + \bar{\beta}_i \bar{b}_0 + \bar{\gamma}_i \bar{b}_1 + \dots)$$

$$\begin{aligned}\Omega_2(D) &= \left(-\frac{1}{2\pi i}\right) du^1 \wedge du^2 \mathcal{B}\left[\frac{\partial}{\partial u^1}\right] \mathcal{B}\left[\frac{\partial}{\partial u^2}\right] (c_1 c_{-1} - \bar{c}_1 \bar{c}_{-1}) |0\rangle \\ &= \left(\frac{1}{2\pi i}\right) du^1 \wedge du^2 [\alpha_1 \gamma_2 - \alpha_2 \gamma_1 - (\bar{\alpha}_1 \bar{\gamma}_2 - \bar{\alpha}_2 \bar{\gamma}_1)] |0\rangle\end{aligned}$$

$$\Omega_2(D) = \left(\frac{1}{2\pi i}\right) du^1 \wedge du^2 \left[\frac{\partial y}{\partial u^1} \frac{\partial}{\partial u^2} \left(\frac{b}{2a^2} \right) - \frac{\partial y}{\partial u^2} \frac{\partial}{\partial u^1} \left(\frac{b}{2a^2} \right) - (\text{c.c.}) \right] |0\rangle$$

$$\Omega_2(D) = \frac{1}{2\pi i} \left[dy \wedge d\left(\frac{b}{2a^2}\right) - d\bar{y} \wedge d\left(\frac{\bar{b}}{2\bar{a}^2}\right) \right] |0\rangle$$

$$\Omega_2(D) = \frac{1}{2\pi i} dy \wedge d\bar{y} \left[\frac{\partial}{\partial \bar{y}} \left(\frac{b}{2a^2} \right) + \frac{\partial}{\partial y} \left(\frac{\bar{b}}{2\bar{a}^2} \right) \right]$$

$$\begin{aligned}\Omega_1(\chi) &= -\frac{1}{2\pi i} \left(-\frac{1}{2} du^1 \mathcal{B}\left[\frac{\partial}{\partial u^1}\right] - \frac{1}{2} du^2 \mathcal{B}\left[\frac{\partial}{\partial u^2}\right] \right) (c_0 - \bar{c}_0) |0\rangle \\ &= -\frac{1}{2\pi i} \left(\frac{1}{2} du^1 (\beta_1 - \bar{\beta}_1) + \frac{1}{2} du^2 (\beta_2 - \bar{\beta}_2) \right) |0\rangle\end{aligned}$$

$$\Omega_1(\chi) = \frac{1}{2\pi i} \left[-\frac{1}{2} d\left(\ln \frac{a}{\bar{a}}\right) + \frac{b}{2a^2} dy - \frac{\bar{b}}{2\bar{a}^2} d\bar{y} \right]$$

$$\rho^w(w) = \rho \left| \frac{dz}{dw} \right| = \left(\rho(y) + (z-y) \frac{\partial \rho}{\partial z} \Big|_y + (\bar{z}-\bar{y}) \frac{\partial \rho}{\partial \bar{z}} \Big|_y + \dots \right) |a + bw + \dots|$$



$$\rho^w(w) = \rho \left| \frac{dz}{dw} \right| = \left(\rho(y) + aw \frac{\partial \rho}{\partial y} \Big|_y + \bar{a} \bar{w} \frac{\partial \rho}{\partial \bar{y}} \Big|_y + \dots \right) |a| \left(1 + \frac{b}{2a} w + \frac{\bar{b}}{2\bar{a}} \bar{w} + \dots \right)$$

$$= \rho(y) |a| \left[1 + aw \left(\frac{\partial}{\partial y} \log \rho + \frac{b}{2a^2} \right) + \bar{a} \bar{w} \left(\frac{\partial}{\partial \bar{y}} \log \rho + \frac{\bar{b}}{2\bar{a}^2} \right) + \dots \right]$$

$$\frac{b}{2a^2} = -\frac{\partial}{\partial y} \log \rho, \frac{\bar{b}}{2\bar{a}^2} = -\frac{\partial}{\partial \bar{y}} \log \rho$$

$$\Omega_2(D) = \left(\frac{1}{2\pi i} \right) dy \wedge d\bar{y} \left[-2 \frac{\partial}{\partial \bar{y}} \frac{\partial}{\partial y} \log \rho \right]$$

$$\Omega_1(\chi) = \frac{1}{2\pi i} (-i) \left[d\theta_a - idy \frac{\partial}{\partial y} \log \rho + id\bar{y} \frac{\partial}{\partial \bar{y}} \log \rho \right]$$

$$\Omega_2(D) = -\frac{1}{2\pi} K^{(2)} \text{ and } \Omega_1(\chi) = -\frac{1}{2\pi} k^{(1)}$$

$$(2-2g-n)\Omega_{6g-6+2n}^{(g,n)}(A_1,\cdots,A_n)$$

Estructuras algebraicas de campo – álgebra de Lie y álgebra cuántica.

$$\deg(b_n(A_1,\cdots,A_n)) = -1 + \sum_{i=1}^n d_{A_i}$$

$$T(V) \equiv V \oplus (V \otimes V) \oplus (V \otimes V \otimes V) \oplus \cdots$$

$$\mathbf{b} = \sum_{i=1}^n \sum_{j=0}^{n-i} \mathbb{1}^{\otimes j} \otimes b_i \otimes \mathbb{1}^{n-i-j}, \text{ on } V^{\otimes n}$$

$$\mathbf{b} = \sum_{i=1}^{\infty} \mathbf{b}_i = \mathbf{b}_1 + \mathbf{b}_2 + \cdots$$

$$\mathbf{b}_i = \sum_{j=0}^{n-i} \mathbb{1}^{\otimes j} \otimes b_i \otimes \mathbb{1}^{n-i-j}, \text{ on } V^{\otimes n}, \text{ for } i \leq n$$

$$\mathbf{b} = b_1 \otimes \mathbb{1} \otimes \mathbb{1} + \mathbb{1} \otimes b_1 \otimes \mathbb{1} + \mathbb{1} \otimes \mathbb{1} \otimes b_1 + b_2 \otimes \mathbb{1} + \mathbb{1} \otimes b_2 + b_3, \text{ on } V^{\otimes 3},$$

$$\begin{aligned} \mathbf{b}(A \otimes B \otimes C) &= b_1(A) \otimes B \otimes C + (-1)^{d_A A} \otimes b_1(B) \otimes C + (-1)^{d_A + d_B A} \otimes B \otimes b_1(C) \\ &+ b_2(A, B) \otimes C + (-1)^{d_A A} \otimes b_2(B, C) + b_3(A, B, C) \end{aligned}$$

$$\mathbf{b}^2=0$$

$$\sum_{i=1}^n \mathbf{b}_i \mathbf{b}_{n+1-i} = 0, \text{ on } V^{\otimes n}$$



$$\begin{aligned}
V: 0 &= \mathbf{b}_1 \mathbf{b}_1 \\
V^{\otimes 2}: 0 &= \mathbf{b}_1 \mathbf{b}_2 + \mathbf{b}_2 \mathbf{b}_1 \\
V^{\otimes 3}: 0 &= \mathbf{b}_1 \mathbf{b}_3 + \mathbf{b}_2 \mathbf{b}_2 + \mathbf{b}_3 \mathbf{b}_1
\end{aligned}$$

$$V: 0 = b_1 \circ b_1$$

$$V^{\otimes 2}: 0 = b_1 \circ b_2 + b_2 \circ (b_1 \otimes 1 + 1 \otimes b_1)$$

$$\begin{aligned}
V^{\otimes 3}: 0 &= b_1 \circ b_3 + b_2 \circ (b_2 \otimes 1 + 1 \otimes b_2), \\
&+ b_3(b_1 \otimes 1 \otimes 1 + 1 \otimes b_1 \otimes 1 + 1 \otimes 1 \otimes b_1)
\end{aligned}$$

$$0 = Q^2 A$$

$$0 = Q(AB) + (QA)B + (-1)^{d_A A}(QB)$$

$$0 = Q(A, B, C) + (AB)C + (-1)^{d_A A}(BC) + (QA, B, C) + (-1)^{d_A}(A, QB, C) + (-1)^{d_A + d_B}(A, B, QC)$$

$$(A, B) = -(-1)^{d_A d_B}(B, A)$$

$$(QA, B) = -(-1)^{d_A}(A, QB)$$

$$(A_1, b_n(A_2, \dots, A_{n+1})) = (-1)^{\#}(A_2, b_n(A_3, \dots, A_{n+1}, A_1)) = d_{A_2} + d_{A_1}(1 + d_{A_2} + d_{A_3} + \dots + d_{A_{n+1}})$$

$$S(\Phi) = \sum_{n=1}^{\infty} \frac{1}{n+1} (\Phi, b_n(\Phi, \dots, \Phi))$$

$$(A, B) = (-1)^{A+1} \langle A, B \rangle'$$

$$A = d_A + 1(\text{mod}2)$$

$$\langle B, A \rangle' = (-1)^{AB} \langle A, B \rangle'$$

$$(B, A) = (-1)^{B+1} \langle B, A \rangle' = (-1)^{AB+B+1} \langle A, B \rangle' = (-1)^{AB+A+B}(A, B) = (-1)^{d_A d_B + 1}(A, B).$$

$$A \star B = (-1)^{A+1} b_2(A, B) = (-1)^{A+1} (AB)$$

$$Q(A \star B) = (-1)^{A+1} Q(AB) = (-1)^A ((QA)B) - (A(QB)) = (QA) \star B + (-1)^A A \star QB.$$

$$[A_1 \cdots A_n] = (-1)^{n(n+1)/2} \sum_{\sigma} (-1)^{s_{\sigma}} (-1)^{A_{\sigma(1)} + 2A_{\sigma(2)} + \cdots + nA_{\sigma(n)}} b_n(A_{\sigma(1)}, A_{\sigma(2)}, \dots, A_{\sigma(n)})$$

$$[A_1, A_2] = A_1 \star A_2 - (-1)^{A_1 A_2} A_2 \star A_1$$

$$[A_1 \cdots A_n] = \sum_{\sigma} b_n(A_{\sigma(1)}, A_{\sigma(2)}, \dots, A_{\sigma(n)})$$

$$b_1(B_1), b_2(B_1, B_2), b_3(B_1, B_2, B_3), \dots$$

$$b_1(B_1) = QB_1, b_1(B_1, B_2) = [B_1, B_2]_0, b_3(B_1, B_2, B_3) = [B_1, B_2, B_3]_0, \dots,$$



$$b_n(B_1, \dots, B_k, B_{k+1}, \dots, B_n) = (-1)^{d_{B_k} d_{B_{k+1}}} b_n(B_1, \dots, B_{k+1}, B_k, \dots, B_n)$$

$$\deg(b(B_1, \dots, B_n)) = -1 + \sum_{i=1}^n d_{B_i}$$

$$T(W) = \sum_{n=1}^{\infty} SW^{\otimes n}$$

$$B_1 \wedge B_2 \wedge \cdots \wedge B_n$$

$$B_1 \wedge \cdots \wedge B_n = \epsilon(\sigma; B) B_{\sigma(1)} \wedge \cdots \wedge B_{\sigma(n)}$$

$$\mathbf{b}(B_1 \wedge \cdots \wedge B_n) = \sum_{l=1}^n \sum_{\sigma'} \epsilon(\sigma', B) (b_l(B_{i_1}, \dots, B_{i_l}) \wedge B_{j_1} \wedge \cdots \wedge B_{j_{n-l}}),$$

$$\mathbf{b}(B_1) = b_1(B_1)$$

$$\mathbf{b}(B_1 \wedge B_2) = b_1(B_1) \wedge B_2 + (-1)^{d_{B_1} d_{B_2}} b_1(B_2) \wedge B_1 + b_1(B_1, B_2)$$

$$= b_1(B_1) \wedge B_2 + (-1)^{d_{B_1} B_1 \wedge b_1(B_2) + b_1(B_1, B_2)}$$

$$\mathbf{b}^2 = 0$$

$$\mathbf{b}^2(B_1) = \mathbf{b}\mathbf{b}(B_1) = \mathbf{b}(QB_1) = Q(QB_1) = 0$$

$$\mathbf{b}^2(B_1 \wedge B_2) = \mathbf{b}(QB_1 \wedge B_2) + (-1)^{d_{B_1}} \mathbf{b}(B_1 \wedge QB_2) + \mathbf{b}(b_2(B_1, B_2)) = 0$$

$$\begin{aligned} \mathbf{b}^2(B_1 \wedge B_2) &= \left((-1)^{d_{B_1}+1} QB_1 \wedge QB_2 + b_2(QB_1, B_2) \right) \\ &+ (-1)^{d_{B_1}} (QB_1 \wedge QB_2 + b_2(B_1, QB_2)) + Qb_2(B_1, B_2) = 0 \end{aligned}$$

$$Qb_2(B_1, B_2) + b_2(QB_1, B_2) + (-1)^{d_{B_1}} b_2(B_1, QB_2) = 0$$

$$0 = Qb_3(B_1, B_2, B_3) + b_3(QB_1, B_2, B_3) + (-1)^{d_{B_1}} b_3(B_1, QB_2, B_3) + (-1)^{d_{B_1}+d_{B_2}} b_3(B_1, B_2, QB_3)$$

$$+ b_2(b_2(B_1, B_2), B_3) + (-1)^{d_{B_2} d_{B_3}} b_2(b_2(B_1, B_3), B_2)$$

$$+ (-1)^{d_{B_1}(d_{B_2}+d_{B_3})} b_2(b_2(B_2, B_3), B_1)$$

$$0 = \sum_{l=1}^n \sum_{\sigma'} \epsilon(\sigma', B) b_{n-l+1}(b_l(B_{i_1}, \dots, B_{i_l}), B_{j_1}, \dots, B_{j_{n-l}}).$$

$$(B_1, B_2) = (-1)^{(d_{B_1}+1)(d_{B_2}+1)} (B_2, B_1),$$

$$(QB_1, B_2) = (-1)^{d_{B_1}} (B_1, QB_2)$$

$$\{B_1, \dots, B_n\}_{L_\infty} \equiv (B_1, b_{n-1}(B_2, \dots, B_n))$$

$$\{B_1, \dots, B_k, B_{k+1}, \dots, B_n\}_{L_\infty} = (-1)^{d_{B_k} d_{B_{k+1}}} \{B_1, \dots, B_{k+1}, B_k, \dots, B_n\}_{L_\infty}$$



$$S = \frac{1}{2}(\Psi, Q\Psi) + \sum_{n=3}^{\infty} \frac{1}{n!} \{\Psi^n\}_{L_{\infty}}$$

$$\bar{b}_1(A_1) \equiv b_1(A_1)$$

$$\bar{b}_2(A_1,A_2)\equiv b_2(A_1,A_2)+(-1)^{d_{A_1}d_{A_2}}b_2(A_2,A_1)$$

$$\bar{b}_1\bar{b}_2(A_1,A_2)+\bar{b}_2(\bar{b}_1(A_1),A_2)+(-1)^{d_{A_1}}\bar{b}_2\left(A_1,\bar{b}_1(A_2)\right)=0$$

$$b_2(b_2(A_1,A_2),A_3)+(-1)^{d_{A_1}}b_2(A_1,b_2(A_2,A_3))=0$$

$$\bar{b}_2\big(\bar{b}_2(A_1,A_2),A_3\big)+(-1)^{d_{A_2}d_{A_3}}\bar{b}_2\big(\bar{b}_2(A_1,A_3),A_2\big)+(-1)^{d_{A_1}(d_{A_2}+d_{A_3})}\bar{b}_2\big(\bar{b}_2(A_2,A_3),A_1\big)=0$$

$$\bar{b}_n(A_1,\cdots,A_n)\equiv\sum_{\sigma}\epsilon(\sigma,A)b_n(A_{\sigma(1)},\cdots,A_{\sigma(n)})$$

$$[A_1\cdots A_n]=(-1)^{d_{A_1}+2d_{A_2}+\cdots+nd_{A_n}}\bar{b}_n(A_1,\cdots,A_n)$$

$$(A_1,A_2)=\langle A_1,A_2\rangle'=(-1)^{d_{A_1}}(A_1,A_2).$$

$$\{A_0A_1\cdots A_n\}_{L_{\infty}}=\left(A_0,\bar{b}_n(A_1A_2\cdots A_n)+(-1)^{d_{A_0}}\left(A_0,\bar{b}_n(A_1A_2\cdots A_n)\right)\right)$$

$$=(-1)^{d_{A_0}}\sum_{\sigma}\epsilon(\sigma,A)\left(A_0,b_n(A_{\sigma(1)},\cdots,A_{\sigma(n)})\right)$$

$$S(\psi_o)=\sum_{n=1}^{\infty}\frac{1}{n+1}(\psi_o,b_n(\psi_o^n))=\frac{1}{2}f\psi_o,Q\psi_ot+\sum_{n=2}^{\infty}\frac{1}{(n+1)!}\{\psi_o^{n+1}\}_{L_{\infty}}$$

$$[A_1\cdots A_n]=\bar{b}_n(A_1,\cdots,A_n)$$

$$b_n(B_1,\cdots,B_n)=\sum_{g=0}^{\infty}g_s^{2g+n-1}b_n^g(B_1,\cdots,B_n),n=0,1,\cdots$$

$$\begin{aligned} 0 = \sum_{l=1}^n \sum_{\sigma'} \epsilon(\sigma',B) b_{n-l+1} & \big(b_l(B_{i_1},\cdots,B_{i_l}), B_{j_1},\cdots,B_{j_{n-l}} \big) \\ & + \frac{1}{2} b_{n+2}(B_1,\cdots,B_n,\varphi_s,\varphi_r)(\varphi_s^c,\varphi_r^c) \end{aligned}$$

$$\theta(A_1\wedge\cdots\wedge A_n)=\frac{1}{2}A_1\wedge\cdots\wedge A_n\wedge\varphi_r\wedge\varphi_s(\varphi_r^c,\varphi_s^c)$$

$$\pi_1(\mathbf{b}^2+\mathbf{b}\boldsymbol{\theta})=0$$

$$(\mathbf{b}+\boldsymbol{\theta})^2=0$$

$$\{A_1^c,\cdots,A_{n_c}^c;A_1^o,\cdots,A_{n_o}^o\}_{L_{\infty}}\equiv (-1)^{d_{A_2^o}+2d_{A_3^o}+\cdots+(n-1)d_{A_n^o}}\{A_1^c,\cdots,A_{n_c}^c;A_1^o,\cdots,A_{n_o}^o\}$$



Transferencia homotópica.

$$\mathbf{M} = \mathbf{Q} + \mathbf{m}, \mathbf{m} = \sum_{n=2}^{\infty} \mathbf{b}_n$$

$$\mathbf{Q}^2 = 0, \text{ and } \mathbf{Q}\mathbf{m} + \mathbf{m}\mathbf{Q} + \mathbf{m}^2 = 0$$

$$P: V \rightarrow \bar{V}, P^2 = P$$

$$(A_1, PA_2) = (PA_1, A_2)$$

$$QP = PQ$$

$$\mathbf{P}(A_1 \otimes A_2 \otimes \cdots \otimes A_n) \equiv PA_1 \otimes PA_2 \otimes \cdots \otimes PA_n$$

$$\mathbf{P}^2 = \mathbf{P}$$

$$\mathbf{QP} = \mathbf{PQ}$$

$$hP = Ph = 0, h^2 = 0, Qh + hQ = 1 - P.$$

$$h(A_1 \otimes A_2 \otimes \cdots \otimes A_n) \equiv (hA_1) \otimes PA_2 \otimes \cdots \otimes PA_n + (-1)^{d_{A_1}} A_1 \otimes hA_2 \otimes \cdots \otimes PA_n :$$

$$\vdots + (-1)^{d_{A_1} + d_{A_2} + \cdots + d_{A_{n-1}}} A_1 \otimes A_2 \otimes \cdots \otimes hA_n$$

$$\mathbf{hP} = \mathbf{Ph} = 0, \mathbf{h}^2 = 0, \mathbf{Qh} + \mathbf{hQ} = \mathbf{1} - \mathbf{P}$$

$$\overline{\mathbf{M}} = \mathbf{PQP} + \mathbf{Pm} \frac{1}{1 + \mathbf{h} \mathbf{m}} \mathbf{P}$$

$$\frac{1}{1 + \mathbf{h} \mathbf{m}} = 1 - \mathbf{h} \mathbf{m} + \mathbf{h} \mathbf{m} \mathbf{h} \mathbf{m} - \cdots$$

$$\overline{\mathbf{M}}^2 = \mathbf{P} \left\{ \mathbf{Q}, \mathbf{m} \frac{1}{1 + \mathbf{h} \mathbf{m}} \right\} \mathbf{P} + \mathbf{Pm} \frac{1}{1 + \mathbf{h} \mathbf{m}} \mathbf{Pm} \frac{1}{1 + \mathbf{h} \mathbf{m}} \mathbf{P}$$

$$\begin{aligned} \mathbf{P} \left\{ \mathbf{Q}, \mathbf{m} \frac{1}{1 + \mathbf{h} \mathbf{m}} \right\} \mathbf{P} &= \mathbf{P} \{ \mathbf{Q}, \mathbf{m} \} \frac{1}{1 + \mathbf{h} \mathbf{m}} \mathbf{P} + \mathbf{Pm} \frac{1}{1 + \mathbf{h} \mathbf{m}} [\mathbf{Q}, \mathbf{1} + \mathbf{h} \mathbf{m}] \frac{1}{1 + \mathbf{h} \mathbf{m}} \mathbf{P} \\ &= -\mathbf{Pm}^2 \frac{1}{1 + \mathbf{h} \mathbf{m}} \mathbf{P} + \mathbf{Pm} \frac{1}{1 + \mathbf{h} \mathbf{m}} ((1 - \mathbf{P})\mathbf{m} + \mathbf{h}^2) \frac{1}{1 + \mathbf{h} \mathbf{m}} \end{aligned}$$

$$\overline{\mathbf{Q}} = \mathbf{PQP}$$

$$\overline{\mathbf{m}}_2 = \mathbf{Pm}_2\mathbf{P}$$

$$\overline{\mathbf{m}}_3 = \mathbf{P}[\mathbf{m}_3 - \mathbf{m}_2\mathbf{h}\mathbf{m}_2]\mathbf{P},$$

$$\overline{\mathbf{m}}_4 = \mathbf{P}[\mathbf{m}_4 - \mathbf{m}_2\mathbf{h}\mathbf{m}_3 - \mathbf{m}_3\mathbf{h}\mathbf{m}_2 + \mathbf{m}_2\mathbf{h}\mathbf{m}_2\mathbf{h}\mathbf{m}_2]\mathbf{P}.$$

$$\bar{Q}A_1 = PQ\bar{A}_1$$

$$\bar{m}_2(A_1 \otimes A_2) = P(\bar{A}_1 \star \bar{A}_2)$$

$$\bar{m}_3(A_1 \otimes A_2 \otimes A_3) = Pm_3(\bar{A}_1, \bar{A}_2, \bar{A}_3) - P(h(\bar{A}_1 \star \bar{A}_2) \star \bar{A}_3 + \bar{A}_1 \star h(\bar{A}_2 \star \bar{A}_3)).$$



$$\mathbf{L} = \mathbf{Q} + \boldsymbol{\ell}, \boldsymbol{\ell} = \sum_{n=2}^{\infty} \mathbf{b}_n$$

$$\mathbf{Q}^2 = 0, \text{ and } \mathbf{Q}\boldsymbol{\ell} + \boldsymbol{\ell}\mathbf{Q} + \boldsymbol{\ell}^2 = 0$$

$$P^2 = P, PQ = QP, (B_1, PB_2) = (PB_1, B_2)$$

$$\mathbf{P}(B_1 \wedge B_2 \wedge \cdots \wedge B_n) \equiv PB_1 \wedge PB_2 \wedge \cdots \wedge PB_n$$

$$hP = Ph = 0, h^2 = 0, Qh + hQ = 1 - P.$$

$$h(A_1 \wedge \dots \wedge A_n) \equiv \frac{1}{n!} \sum_{\sigma \in \mathcal{S}_n} \epsilon(\sigma; A) (h(A_{\sigma(1)} \wedge \dots \wedge A_{\sigma(n-1)}) \wedge A_{\sigma(n)} \\ + (-1)^{A_{\sigma(1)} + \dots + A_{\sigma(n-1)}} \bar{A}_{\sigma(1)} \wedge \dots \wedge \bar{A}_{\sigma(n-1)} \wedge hA_{\sigma(n)})$$

$$h(A_1 \wedge A_2) = \frac{1}{2} (hA_1 \wedge A_2 + hA_1 \wedge \bar{A}_2 + (-1)^{A_1} (\bar{A}_1 \wedge hA_2 + A_1 \wedge hA_2))$$

$$\mathbf{hP} = \mathbf{Ph} = 0, \mathbf{h}^2 = 0, \mathbf{Qh} + \mathbf{hQ} = \mathbf{1} - \mathbf{P}$$

$$\bar{\mathbf{L}} = \mathbf{PQP} + \mathbf{P}\boldsymbol{\ell} \frac{1}{1 + \mathbf{h}\boldsymbol{\ell}} \mathbf{P}$$

$$\bar{Q}A_1 = Q\bar{A}_1\bar{\ell}_2(A_1 \wedge A_2) = P[\bar{A}_1, \bar{A}_2]\bar{\ell}_3(A_1 \wedge A_2 \wedge A_3) \\ = P\ell_3(\bar{A}_1 \wedge \bar{A}_2 \wedge \bar{A}_3) \\ - P([h[\bar{A}_1, \bar{A}_2], \bar{A}_3] + (-1)^{A_2A_3}[h[\bar{A}_1, \bar{A}_3], \bar{A}_2] + [\bar{A}_1, h[\bar{A}_2, \bar{A}_3]])$$

Vórtices – espacios de Moduli.

$$\Delta(\Delta X) = \Delta^2 X = 0, \forall X \in \mathcal{C}$$

$$\Delta(XYZ) = \Delta(XY)Z + (-1)^X X\Delta(YZ) + (-1)^{XY+Y} Y\Delta(XZ) \\ - \Delta X(YZ) - (-1)^X X(\Delta Y)Z - (-1)^{X+Y} XY\Delta Z$$

$$\{X,Y\} \equiv (-1)^X \Delta(XY) - (-1)^X (\Delta X)Y - X\Delta Y$$

$$\{X,Y\} = -(-1)^{(X+1)(Y+1)}\{Y,X\} \\ 0 = (-1)^{(X+1)(Z+1)}\{\{X,Y\},Z\} + \mathfrak{I}_{\text{cyclic}} \\ \{X,YZ\} = \{X,Y\}Z + (-1)^{XY+Y}Y\{X,Z\}$$

$$\mathcal{A}_{g_1,n_1} \times \dots \times \mathcal{A}_{g_r,n_r}$$

$$\deg(\mathcal{A}_{g,n}) = \dim(\mathcal{A}_{g,n})(\text{mod}2)$$

$$\mathcal{A}_{g_1,n_1} \times \mathcal{A}_{g_2,n_2} = (-1)^{\mathcal{A}_{g_1,n_1} \cdot \mathcal{A}_{g_2,n_2}} \mathcal{A}_{g_2,n_2} \times \mathcal{A}_{g_1,n_1}$$

$$[\mathcal{A}_{g_1,n_1}, \dots, \mathcal{A}_{g_r,n_r}] \equiv \frac{1}{n_1! \dots n_r!} \sum_{\sigma \in S_N} \mathbf{P}_{\sigma}(\mathcal{A}_{g_1,n_1} \times \dots \times \mathcal{A}_{g_r,n_r}) \in \mathcal{C}$$



$$\left[\left[\mathcal{A}_{g_1,n_1},\dots,\mathcal{A}_{g_r,n_r}\right]\right]\cdot\left[\left[\mathcal{B}_{g'_1,n'_1},\dots,\mathcal{B}_{g'_p,n'_p}\right]\right]=\left[\left[\mathcal{A}_{g_1,n_1},\dots,\mathcal{A}_{g_r,n_r},\mathcal{B}_{g'_1,n'_1},\dots,\mathcal{B}_{g'_p,n'_p}\right]\right]$$

$$\Delta X \equiv \Delta_{ij} X \equiv \frac{1}{n_1! \cdots n_r!} \sum_{\sigma \in S_N} \Delta_{ij} \mathbf{P}_{\sigma}(\mathcal{A}_{n_1, g_1} \times \cdots \times \mathcal{A}_{g_r, n_r})$$

$$\Delta_{12}(\Delta_{12}\Sigma_X)\left[\frac{\partial}{\partial\theta_{34}},\frac{\partial}{\partial\theta_{12}},\{X\}\right]$$

$$\Delta_{12}(\Delta_{34}\Sigma_X)\left[\frac{\partial}{\partial\theta_{12}},\frac{\partial}{\partial\theta_{34}},\{X\}\right]$$

$$\{X,Y\}\equiv (-1)^X\Delta(XY)-(-1)^X(\Delta X)Y-X\Delta Y$$

$$(-1)^X\Delta(XY)=S_{XX}+S_{YY}+S_{XY}$$

$$(-1)^X(\Delta X)Y=S_{XX},X\Delta Y=S_{YY}$$

$$\partial(XY)=(\partial X)Y+(-1)^XX\partial Y$$

$$\Delta\partial X=-\partial\Delta X$$

$$\partial\{X,Y\}=\{\partial X,Y\}+(-1)^{X+1}\{X,\partial Y\}$$

Estructuras de campo en CFT.

$$\left\langle \Omega_k^{(g,n)} \right| \in (\mathcal{H}^*)^{\otimes n}$$

$$\left\langle \Omega_k^{(g,n)} \mid A_1 \right\rangle \otimes \cdots \otimes \mid A_n \rangle = \Omega_k^{(g,n)}(A_1,\cdots,A_n),$$

$$\int_{\mathcal{A}_{g,n}} \left\langle \Omega_k^{(g,n)} \mid \sum_{i=1}^n Q^{(i)} = (-1)^k \int_{\partial \mathcal{A}_{g,n}} \left\langle \Omega_{k-1}^{(g,n)} \mid \right.$$

$$f\left(\mathcal{A}_{g,n}^{(k)}\right)\equiv\frac{1}{n!}\int_{\mathcal{A}_{g,n}^{(k)}}\left\langle\Omega_k^{(g,n)}\mid\Psi\right\rangle_1\cdots\mid\Psi\rangle_n$$

$$f\left(\left[\left[\mathcal{A}_{g_1,n_1}^{(k_1)}\cdots\mathcal{A}_{g_r,n_r}^{(k_r)}\right]\right]\right)\equiv\prod_{i=1}^r\frac{1}{n_i!}\int_{\mathcal{A}_{g_i,n_i}^{(k_i)}}\left\langle\Omega_{k_i}^{(g_i,n_i)}\mid\Psi\right\rangle_1\cdots\mid\Psi\rangle_{n_i}$$

$$f(XY)=f(X)f(Y),X,Y\in\mathcal{C}$$

$$_1\langle A\mid _2\langle B\mid R_{12}\rangle=\langle A\mid B\rangle,\langle R_{12}\mid A\rangle_1\mid B\rangle_2=\langle A\mid B\rangle,$$

$$\langle\omega_{12}|=\langle R_{12}|c_0^{-(2)}$$

$$|S_{12}\rangle=b_0^{-(1)}|R_{12}\rangle$$

$$\{F,G\}=(-1)^{G+1}\frac{\partial F}{\partial|\Psi\rangle}\frac{\partial G}{\partial|\Psi\rangle}|S\rangle$$



$$S_{0,2} = \frac{1}{2} \langle \omega_{12} | Q^{(2)} | \Psi \rangle_1 | \Psi \rangle_2$$

$$\Delta F = \frac{1}{2} (-1)^{F+1} \left(\frac{\partial}{\partial |\Psi\rangle} \frac{\partial}{\partial |\Psi\rangle} F \right) |S\rangle$$

$$f(\Delta X)=-\Delta f(X)$$

$$f(\{X,Y\})=-\{f(X),f(Y)\},$$

$$\{S_{0,2},f(X)\}=-f(\partial X)$$

$$\mathcal{V}=\sum_{g,n}g_s^{2g+n-2}\mathcal{V}_{g,n},\text{ with }\begin{cases}n\geqslant 3,\text{ for }g=0\\n\geqslant 1,\text{ for }g=1\\n\geqslant 0,\text{ for }g\geqslant 2\end{cases}$$

$$\partial \mathcal{V} + \Delta \mathcal{V} + \frac{1}{2} \{ \mathcal{V}, \mathcal{V} \} = 0$$

$$S=S_{0,2}+f(\mathcal{V})$$

$$\{S_{0,2},S_{0,2}\}=0,\Delta S_{0,2}=0$$

$$0=\frac{1}{2}\{S,S\}+\Delta S=\{S_{0,2},f(\mathcal{V})\}+\frac{1}{2}\{f(\mathcal{V}),f(\mathcal{V})\}+\Delta f(\mathcal{V})$$

$$=-f(\partial \mathcal{V})-\frac{1}{2}f(\{\mathcal{V},\mathcal{V}\})-f(\Delta \mathcal{V})-f\left(\partial \mathcal{V}+\frac{1}{2}\{\mathcal{V},\mathcal{V}\}+\Delta \mathcal{V}\right)=0$$

$$\partial \bigcirc_{\mathcal{V}_{g,n}} = -\frac{1}{2} \sum_{\substack{g_1+g_2=g \\ n_1+n_2=n+2}} \bigcirc_{\mathcal{V}_{g_1,n_1}} \text{---} \bigcirc_{\mathcal{V}_{g_2,n_2}} - \frac{1}{2} \bigcirc_{\mathcal{V}_{g-1,n+2}}^{\text{bubble}}$$

$$\partial \mathcal{V}_{g,n} = -\Delta \mathcal{V}_{g-1,n+2} - \frac{1}{2} \sum_{\substack{g_1+g_2=g \\ n_1+n_2=n+2}} \{ \mathcal{V}_{g_1,n_1}, \mathcal{V}_{g_2,n_2} \} \equiv \mathcal{O}_{g,n}$$

$$\partial \mathcal{V}_{g,n} = -\partial_p R_1,$$



Transformaciones canónicas.

$$\mathcal{F}_{g,n} = \mathcal{V}_{g,n} \oplus R_1 \oplus \cdots \oplus R_{3g-3+n}$$

$$\partial \mathcal{V}_{0,4} = -\frac{1}{2}\{\mathcal{V}_{0,3}, \mathcal{V}_{0,3}\}, \text{ and } \partial \mathcal{V}_{1,1} = -\Delta \mathcal{V}_{0,3}.$$

$$\partial \mathcal{V}_{g,n} = \mathcal{O}_{g,n}$$

$$[\mathcal{O}_{g,n}] \in H_{6g-6+2n-1}(\hat{\mathcal{P}}_{g,n})^{S_n}$$

$$\delta S = \Delta \epsilon + \{S, \epsilon\}$$

$$\delta_{\mathcal{W}} \mathcal{V} = \{\mathcal{V}, \mathcal{W}\} + \Delta \mathcal{W} + \partial \mathcal{W},$$

$$\mathcal{V} + \delta_{\mathcal{W}} \mathcal{V}$$

$$\begin{aligned} & \frac{1}{2}\{\mathcal{V} + \delta_{\mathcal{W}} \mathcal{V}, \mathcal{V} + \delta_{\mathcal{W}} \mathcal{V}\} + \Delta(\mathcal{V} + \delta_{\mathcal{W}} \mathcal{V}) + \partial(\mathcal{V} + \delta_{\mathcal{W}} \mathcal{V}) \\ &= \left\{ \partial \mathcal{V} + \Delta \mathcal{V} + \frac{1}{2}\{\mathcal{V}, \mathcal{V}\}, \mathcal{W} \right\} + O((\delta_{\mathcal{W}} \mathcal{V})^2) = 0 \end{aligned}$$

$$S_{\mathcal{V} + \delta_{\mathcal{W}} \mathcal{V}} = S_{\mathcal{V}} + f(\delta_{\mathcal{W}} \mathcal{V}) = S_{\mathcal{V}} + f(\{\mathcal{V}, \mathcal{W}\} + \Delta \mathcal{W} + \partial \mathcal{W})$$

$$= S_{\mathcal{V}} - \{f(\mathcal{V}), f(\mathcal{W})\} - \Delta f(\mathcal{W}) - \{S_{0,2}, \mathcal{W}\} = S_{\mathcal{V}} - \{S_{\mathcal{V}}, f(\mathcal{W})\} - \Delta f(\mathcal{W})$$

$$\frac{d}{dt} \mathcal{V}(t) = \delta_{\mathcal{W}} \mathcal{V}(t)$$

$$M_{\mathcal{V}}(t) \equiv \partial \mathcal{V}(t) + \Delta \mathcal{V}(t) + \frac{1}{2}\{\mathcal{V}(t), \mathcal{V}(t)\}$$

$$\frac{dM_{\mathcal{V}}}{dt} = \{M_{\mathcal{V}}(t), \mathcal{W}\}$$

$$\mathcal{V}(t) = \mathcal{V} + t\delta_{\mathcal{W}} \mathcal{V} + \frac{1}{2}t^2\{\delta_{\mathcal{W}} \mathcal{V}, \mathcal{W}\} + \frac{1}{3!}t^3\{\{\delta_{\mathcal{W}} \mathcal{V}, \mathcal{W}\}, \mathcal{W}\} + \cdots$$

$$\exp(\delta_{\mathcal{W}})\mathcal{V} \equiv \mathcal{V} + \delta_{\mathcal{W}} \mathcal{V} + \frac{1}{2}\{\delta_{\mathcal{W}} \mathcal{V}, \mathcal{W}\} + \frac{1}{3!}\{\{\delta_{\mathcal{W}} \mathcal{V}, \mathcal{W}\}, \mathcal{W}\} + \cdots$$

$$\exp(\delta_{\mathcal{W}})\mathcal{V} = \mathcal{V}'$$

$$\partial \begin{array}{c} g, b \\ n_c, n_o \end{array} = - \begin{array}{c} \text{sun} \\ g-1, b \\ n_c+2, n_o \end{array} - \begin{array}{c} \text{loop} \\ g, b-1 \\ n_c, n_o+2 \end{array} - \begin{array}{c} \text{double loop} \\ g-1, b+1 \\ n_c, n_o+2 \end{array}$$

$$-\frac{1}{2} \begin{array}{c} g_1, b_1 \\ n_{c1}, n_{o1} \end{array} \text{---} \begin{array}{c} g_2, b_2 \\ n_{c2}, n_{o2} \end{array}$$

$$-\frac{1}{2} \begin{array}{c} g_1, b_1 \\ n_{c1}, n_{o1} \end{array} \text{---} \begin{array}{c} g_2, b_2 \\ n_{c2}, n_{o2} \end{array}$$

$$(b_1 + b_2 = b + 1)$$

$$\mathcal{V} = \sum_{g,b,n_c,n_o} (g_s)^{-\chi_{g,b,n_c,n_o}} \mathcal{V}_{g,b,n_c,n_o},$$

$$\partial \mathcal{V} + \Delta \mathcal{V} + \frac{1}{2} \{ \mathcal{V}, \mathcal{V} \} = 0$$

$$\partial \mathcal{V}_c + \Delta \mathcal{V}_c + \frac{1}{2} \{ \mathcal{V}_c, \mathcal{V}_c \}_c = 0$$

$$\partial \hat{\mathcal{V}}_c + \frac{1}{2} \{ \hat{\mathcal{V}}_c, \hat{\mathcal{V}}_c \}_c = 0$$

$$\partial \mathcal{V}_o + \frac{1}{2} \{ \mathcal{V}_o, \mathcal{V}_o \}_o = 0$$

$$\partial \mathcal{V}_1 + \{ \mathcal{V}_o, \mathcal{V}_1 \}_o = 0$$

$$\partial \mathcal{V}_2 + \{ \mathcal{V}_o, \mathcal{V}_2 \}_o + \frac{1}{2} \{ \mathcal{V}_1, \mathcal{V}_1 \}_o + \{ \mathcal{V}_1, \mathcal{V}_{0,3} \}_c = 0$$

$$\bar{\mathcal{V}} \equiv \mathcal{V}_o + \sum_{n=1}^{\infty} \mathcal{V}_n, \mathcal{V}_n \equiv \sum_{n_o=0}^{\infty} (g_s)^{n+\frac{1}{2}n_o-1} \mathcal{V}_{0,1,n,n_o}$$



$$\partial \overline{\mathcal{V}} + \frac{1}{2} \{ \overline{\mathcal{V}}, \overline{\mathcal{V}} \}_o + \{ \overline{\mathcal{V}}, \hat{\mathcal{V}}_c \}_c = 0$$

$$\tilde{\mathcal{V}} = \sum_{n_c, b, n_o} (g_s)^{n_c + b + \frac{1}{2}n_o - 2} \mathcal{V}_{0, b, n_c, n_o}$$

$$\partial \tilde{\mathcal{V}} + \frac{1}{2} \{ \tilde{\mathcal{V}}, \tilde{\mathcal{V}} \} + \Delta'_o \tilde{\mathcal{V}} = 0$$

Vórtice de Witten.

$$\begin{aligned} w_1 w_2 &= -1, & \text{for } |w_1| = 1, \operatorname{Re} w_1 \leq 0 \\ w_2 w_3 &= -1, & \text{for } |w_2| = 1, \operatorname{Re} w_2 \leq 0 \\ w_3 w_1 &= -1, & \text{for } |w_3| = 1, \operatorname{Re} w_3 \leq 0 \end{aligned}$$

$$\varphi=\phi(w_i)dw_i^2=-\frac{1}{w_i^2}dw_i^2$$

$$ds^2=|\phi(w_i)||dw_i|^2$$

$$h(u)=\frac{1+iu}{1-iu}$$

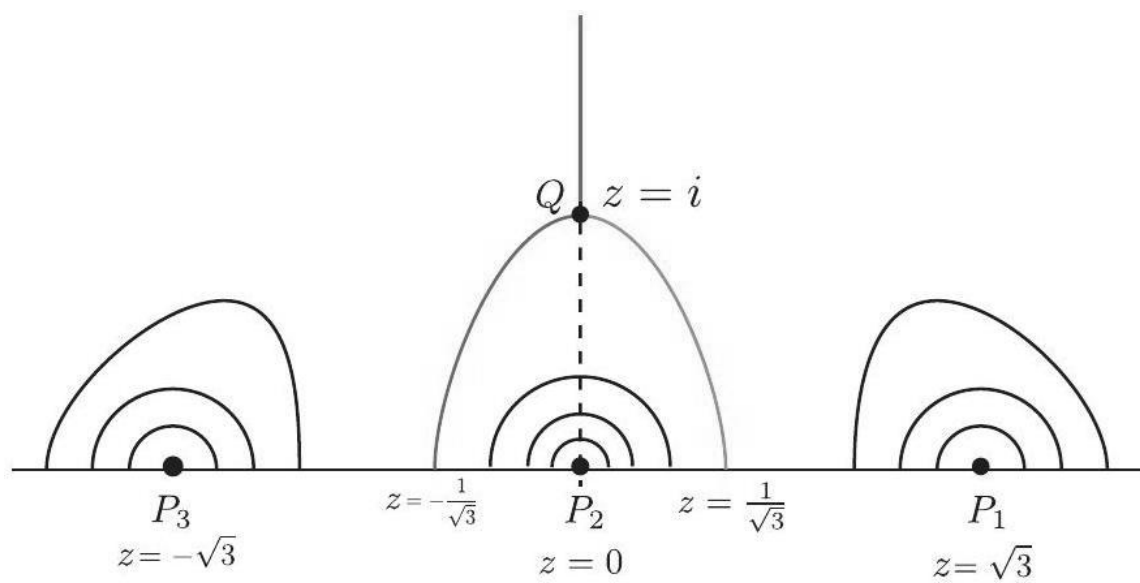
$$w(w_1)=e^{2\pi i/3}\big(h(w_1)\big)^{2/3}, w(w_2)=\big(h(w_2)\big)^{2/3}, w(w_3)=e^{-2\pi i/3}\big(h(w_3)\big)^{2/3},$$

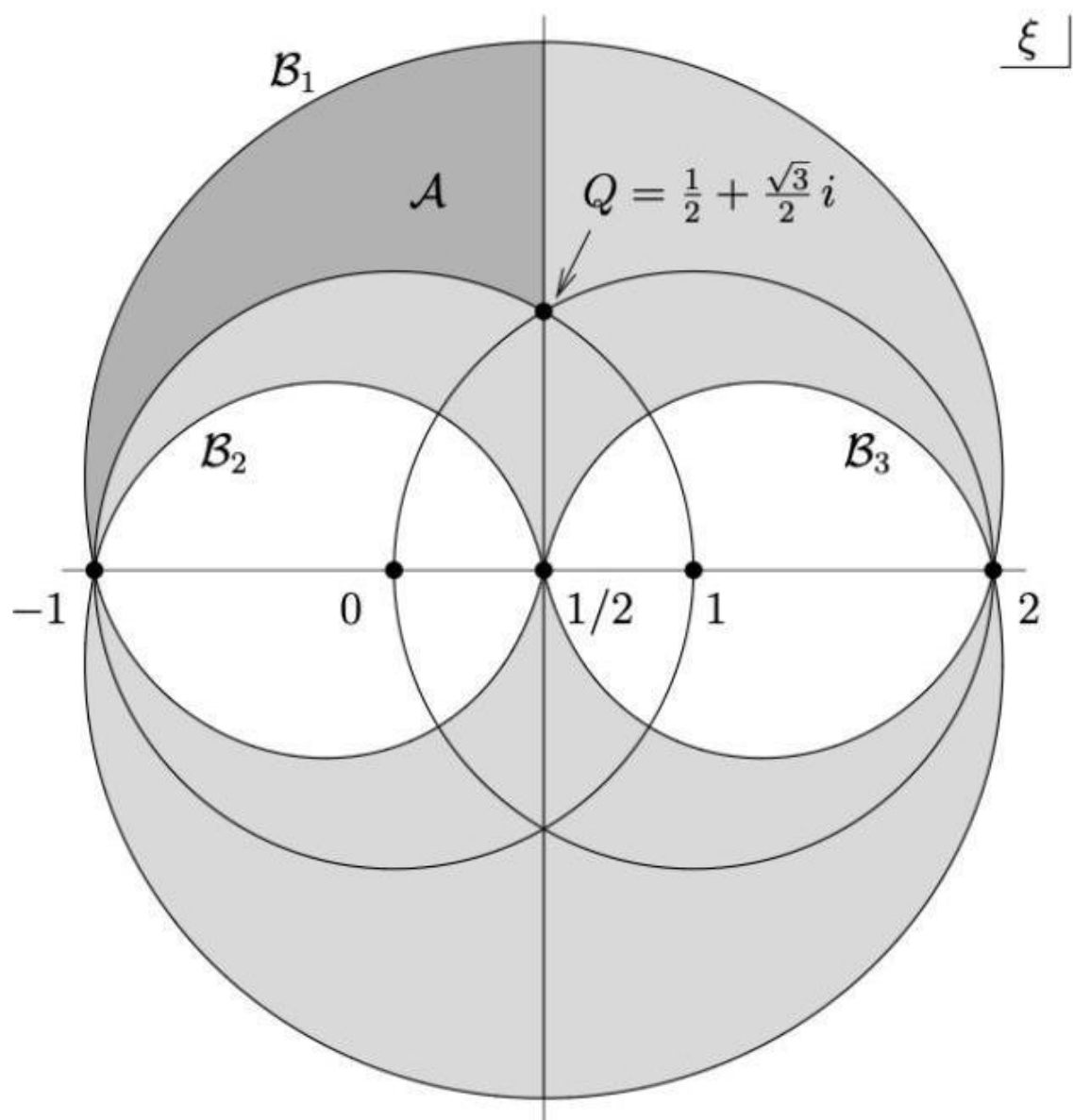
$$z=f_1(w_1)=h^{-1}(w(w_1))=\sqrt{3}+\frac{8}{3}w_1+\frac{16}{9}\sqrt{3}w_1^2+\frac{248}{81}w_1^3+\mathcal{O}(w_1^4)$$

$$z=f_2(w_2)=h^{-1}(w(w_2))=\frac{2}{3}w_2-\frac{10}{81}w_2^3+\mathcal{O}(w_2^5)$$

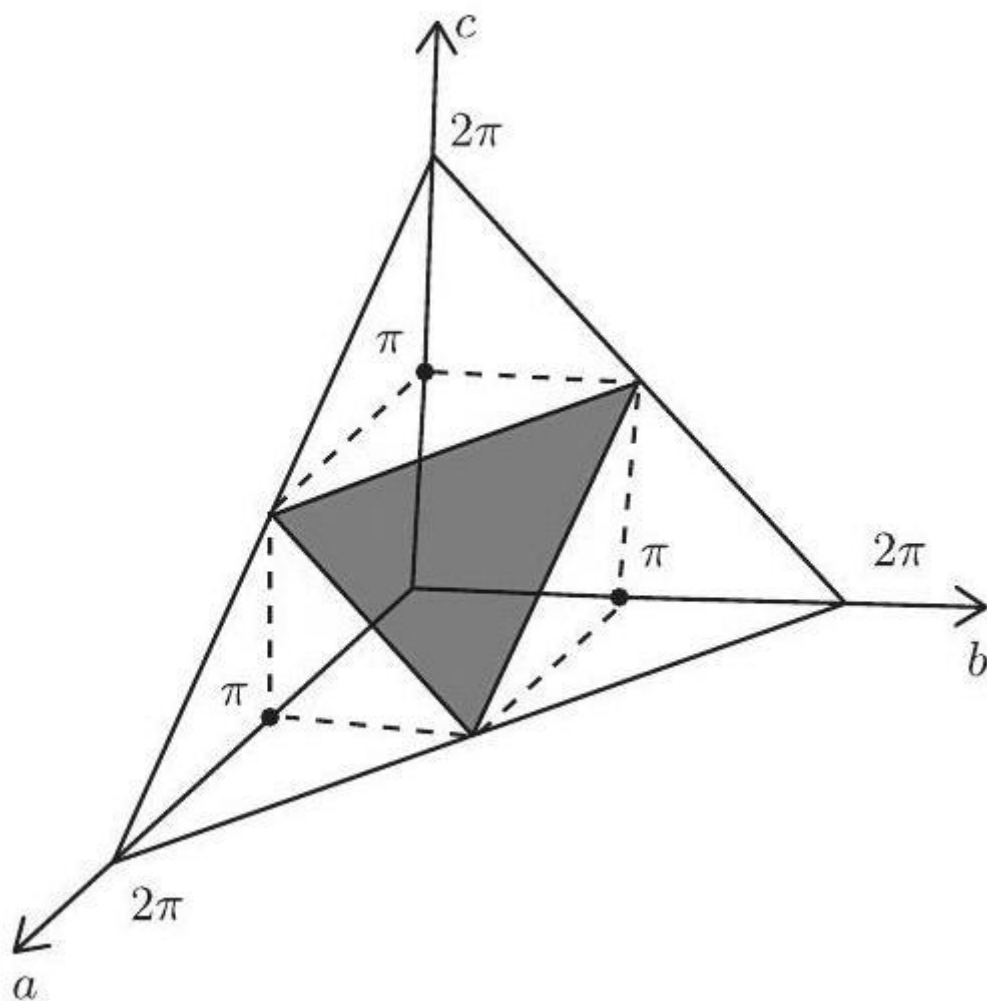
$$z=f_3(w_3)=h^{-1}(w(w_3))=-\sqrt{3}+\frac{8}{3}w_3-\frac{16}{9}\sqrt{3}w_3^2+\frac{248}{81}w_3^3+\mathcal{O}(w_3^4)$$







$$a, b, c \leq \pi, a', b', c' \leq \pi$$



Métricas hiperbólicas y curvatura de Gauss – Riemann – Espacios de Hilbert – Einstein deformados.

$$\mathrm{gr}_{\infty} \colon \Sigma_{g,n} \rightarrow \hat{\mathcal{P}}_{g,n}$$

$$\pi \circ \mathrm{gr}_{\infty} \colon \mathcal{M}_{g,n,L} \rightarrow \mathcal{M}_{g,n}$$

$$\tilde{\mathcal{V}}_{g,n}(L) \equiv \{ \Sigma \in \mathcal{M}_{g,n,L} \mid \text{sys } \Sigma \geqslant L \}$$

$$\mathcal{V}_{g,n}(L) \equiv \mathrm{gr}_{\infty}(\tilde{\mathcal{V}}_{g,n}(L))$$

$$\sinh\left(\frac{1}{2}w_i\right)\sinh\left(\frac{1}{2}\ell_i\right)=1$$

$$L\leqslant L_*\rightarrow w\geqslant L$$

$$L_c\leqslant L'_*,\text{ and }\sinh L_c\sinh L_o\leqslant 1$$

$$\mathbf{B}^{[0]}(t)=\sum_{n=0}^{\infty}t^n\mathbf{b}_{n+1}^{(n)}$$

$$(\mathbf{B}^{[0]}(t))^2 = 0$$

$$\{\boldsymbol{\eta}_0, \mathbf{B}^{[0]}(t)\} = 0$$

$$\mathbf{B}^{[m]}(t) = \sum_{n=0}^{\infty} t^n \mathbf{b}_{m+n+1}^{(n)}, m = 0, 1, \dots$$

$$\mathbf{B}(s, t) = \sum_{m=0}^{\infty} s^m \mathbf{B}^{[m]}(t) = \sum_{m,n=0}^{\infty} s^m t^n \mathbf{b}_{m+n+1}^{(n)}$$

$$\boldsymbol{\mu}(s, t) = \sum_{m=0}^{\infty} s^m \boldsymbol{\mu}^{[m]}(t) = \sum_{m,n=0}^{\infty} s^m t^n \boldsymbol{\mu}_{m+n+2}^{(n+1)}$$

$$\{\mathbf{B}(s, t), \mathbf{B}(s, t)\} = 0, \text{ and } \{\boldsymbol{\eta}_0, \mathbf{B}(s, t)\} = 0$$

$$\frac{\partial}{\partial t} \mathbf{B}(s, t) = [\mathbf{B}(s, t), \boldsymbol{\mu}(s, t)], \frac{\partial}{\partial s} \mathbf{B}(s, t) = [\boldsymbol{\eta}_0, \boldsymbol{\mu}(s, t)].$$

$$\mathbf{B}(s, t) = \mathbf{Q} + t \mathbf{b}_2^{(1)} + s \mathbf{b}_2^{(0)} + \dots, \boldsymbol{\mu}(s, t) = \boldsymbol{\mu}_2^{(1)} + \dots$$

$$\mathbf{b}_2^{(1)} = [\mathbf{Q}, \boldsymbol{\mu}_2^{(1)}], \mathbf{b}_2^{(0)} = [\boldsymbol{\eta}_0, \boldsymbol{\mu}_2^{(1)}].$$

$$\boldsymbol{\mu}_2^{(1)} = \frac{1}{3} (\xi_0 \mathbf{b}_2^{(0)} - \mathbf{b}_2^{(0)} \xi_0).$$

$$\boldsymbol{\mu}_2^{(1)}(A \otimes B) = \frac{1}{3} (\xi_0(A, B) - (\xi_0 A, B) - (-1)^A(A, \xi_0 B)).$$

$$\begin{aligned} [\boldsymbol{\eta}_0, \boldsymbol{\mu}_2^{(1)}](A \otimes B) &= \boldsymbol{\eta}_0 \boldsymbol{\mu}_2^{(1)}(A \otimes B) = \frac{1}{3} \boldsymbol{\eta}_0 (\xi_0(A, B) - (\xi_0 A, B) - (-1)^A(A, \xi_0 B)) \\ &= \frac{1}{3} (\boldsymbol{\eta}_0 \xi_0(A, B) + (\boldsymbol{\eta}_0 \xi_0 A, B) + (A, \boldsymbol{\eta}_0 \xi_0 B)) = \frac{1}{3} ((A, B) + (A, B) + (A, B)) \end{aligned}$$

$$= (A, B) = \mathbf{b}_2^{(0)}(A \otimes B)$$

$$(A, B)^* \equiv \mathbf{b}_2^{(1)}(A \otimes B) = [\mathbf{Q}, \boldsymbol{\mu}_2^{(1)}](A \otimes B)$$

$$= \mathbf{Q} \boldsymbol{\mu}_2^{(1)}(A \otimes B) - \boldsymbol{\mu}_2^{(1)}(QA \otimes B) - (-1)^A \boldsymbol{\mu}_2^{(1)}(A \otimes QB)$$

$$= \frac{1}{3} Q (\xi_0(A, B) - (\xi_0 A, B) - (-1)^A(A, \xi_0 B))$$

$$- \frac{1}{3} (\xi_0(QA, B) - (\xi_0 QA, B) + (-1)^A(QA, \xi_0 B))$$

$$- \frac{1}{3} (-1)^A (\xi_0(A, QB) - (\xi_0 A, QB) - (-1)^A(A, \xi_0 QB))$$



$$(A, B)^* \equiv \frac{1}{3} (\mathcal{X}_0(A, B) + (\mathcal{X}_0 A, B) + (A, \mathcal{X}_0 B))$$

$$S(\Phi) = \frac{1}{2} \langle \Phi, Q\Phi \rangle' + \sum_{n=1}^{\infty} \frac{1}{n+2} \left\langle \Phi, b_{n+1}^{(n)}(\Phi, \dots, \Phi) \right\rangle'$$

$$[A,B]^* \equiv \frac{1}{3} (\mathcal{X}_0[A,B] + [\mathcal{X}_0 A,B] + [A,\mathcal{X}_0 B])$$

Métrica de Berkovits para espacios de Hilbert – Einstein deformados.

$$\{\eta_0,\xi_0\}=1$$

$$\{Q,\eta_0\}=0$$

$$S=\frac{1}{2g^2}\left\langle\left\langle(e^{-\Phi}\eta_0e^{\Phi})(e^{-\Phi}Qe^{\Phi})+\int_0^1dt\Phi\{e^{-t\Phi}Qe^{t\Phi},e^{-t\Phi}\eta_0e^{t\Phi}\}\right\rangle\right\rangle$$

$$\langle\langle A_1\ldots A_n\rangle\rangle=\left\langle h^{-1}\circ f_1^{(n)}\circ A_1(0)\cdots h^{-1}\circ f_n^{(n)}\circ A_n(0)\right\rangle_L',n\geqslant 2,$$

$$f_1^{(n)}(z)=\left(\frac{1+iz}{1-iz}\right)^{2/n}, f_k^{(n)}(z)=e^{2\pi i(k-1)/n}f_1^{(n)}(z), k=2,\ldots n$$

$$\left\langle \xi c \partial c \partial^2 c e^{-2\phi} \right\rangle_L' \neq 0$$

$$\langle\langle A_1\ldots A_n\rangle\rangle=\langle A_1,A_2\star A_3\star\cdots\star A_n\rangle_L'$$

$$\eta_0(e^{-\Phi}Qe^{\Phi})=0$$

$$\delta e^{\Phi}=(Q\Lambda)e^{\Phi}+e^{\Phi}(\eta_0\Omega)$$

$$S=\frac{1}{2g^2}\left\langle\left\langle\frac{1}{2}(Q\Phi)(\eta_0\Phi)+\frac{1}{6}(Q\Phi)(\Phi(\eta_0\Phi)-(\eta_0\Phi)\Phi)\right\rangle\right\rangle+\mathcal{O}(\Phi^3)$$

$$S_{\rm kin} \sim \langle\langle (Q\Phi)(\eta_0\Phi)\rangle\rangle = \langle Q\Phi, \eta_0\Phi\rangle_L',$$

$$Q\eta_0|\Phi\rangle=0$$

$$\delta\Phi=Q|\Lambda\rangle,\delta\Phi=\eta_0|\Omega\rangle$$

$$|\Phi\rangle=|\phi'\rangle+\xi_0|\phi\rangle, \text{ with } \eta_0|\phi'\rangle=\eta_0|\phi\rangle=0$$

$$\delta|\Phi\rangle=\eta_0|\Omega\rangle=\eta_0(|\omega\rangle+\xi_0|\omega'\rangle)=|\omega'\rangle$$

$$0=Q\eta_0|\Phi\rangle=Q\eta_0\xi_0|\phi\rangle=Q\{\eta_0,\xi_0\}|\phi\rangle=Q|\phi\rangle$$



Sector holomórfico.

$$\langle \xi c \partial c \partial^2 c \bar{c} \partial \bar{c} \partial^2 \bar{c} e^{-2\phi} \rangle_L \neq 0$$

$$S = 2 \left[\frac{1}{2} \langle \eta_0 V, QV \rangle_L + \frac{\kappa}{3!} \langle \eta_0 V, [V, QV]_0 \rangle_L \right] + \mathcal{O}(V^3)$$

$$\delta V = Q\Lambda + \eta_0\Omega$$

$$G(\tau V + d\tau V) = G(\tau V) + Q(d\tau V) + \sum_{n=1}^{\infty} \frac{\kappa^n}{n!} [G(\tau V)^n, d\tau V]_0 + \mathcal{O}(d\tau^2)$$

$$\partial_\tau G(\tau V) = QV + \sum_{n=1}^{\infty} \frac{\kappa^n}{n!} [G(\tau V)^n, V]_0$$

$$\partial_\tau G = QV + \kappa[G, V]_0 + \frac{\kappa^2}{2} [G, G, V]_0 + \mathcal{O}(\kappa^3)$$

$$G = G^{(0)} + \kappa G^{(1)} + \kappa^2 G^{(2)} + \mathcal{O}(\kappa^3)$$

$$G(V) = QV + \frac{\kappa}{2} [V, QV]_0 + \frac{\kappa^2}{3!} ([V, QV, QV]_0 + [V, [V, QV]_0]_0) + \mathcal{O}(\kappa^3)$$

$$S = 2 \int_0^1 dt \langle \eta_0 V, G(tV) \rangle_L$$

$$\langle \xi \bar{\xi} c \partial c \partial^2 c \bar{c} \partial \bar{c} \partial^2 \bar{c} e^{-2\phi} e^{-2\bar{\phi}} \rangle_L \neq 0$$

$$S_2 = \frac{1}{2} \langle \eta_0 \Psi, Q \bar{\eta}_0 \Psi \rangle_L$$

$$\delta \Psi = Q\Lambda + \eta_0\Omega + \bar{\eta}_0\bar{\Omega}$$

$$S_3 = \frac{1}{3!} \langle \eta_0 \Psi, [Q \bar{\eta}_0 \Psi, \bar{\eta}_0 \Psi]^* \rangle_L$$

$$[A,B]^* \equiv \frac{1}{3} \left(\overline{\mathcal{X}}_0[A,B]_0 + [\overline{\mathcal{X}}_0 A, B]_0 + [A, \overline{\mathcal{X}}_0 B]_0 \right)$$

$$S = \int_0^1 dt \langle \eta_0 \Psi_t, \mathcal{G}^*(\Psi(t)) \rangle_L$$

$$\Psi_t = \Psi + \frac{\kappa}{2} t [\bar{\eta}_0 \Psi, \Psi]^* + \mathcal{O}(\kappa^2)$$

$$\mathcal{G}^*(\Psi) = Q \bar{\eta}_0 \Psi + \frac{\kappa}{2} t [Q \bar{\eta}_0 \Psi, \bar{\eta}_0 \Psi]^* + \mathcal{O}(\kappa^2)$$



Sector de Ramond.

$$\beta(z) = \sum_{n \in \mathbb{Z}} \beta_n z^{-n-\frac{3}{2}}, \gamma(z) = \sum_{n \in \mathbb{Z}} \gamma_n z^{-n+\frac{1}{2}}$$

$$\beta_n|-1/2\rangle=0 \text{ for } n\geqslant 0, \gamma_n|-1/2\rangle=0 \text{ for } n\geqslant 1$$

$$X=-\delta(\beta_0)G_0^m+\delta'(\beta_0)b_0, Y=-c_0\delta'(\gamma_0)$$

$$\Psi=\phi-(\gamma_0+2c_0b_0\gamma_0+c_0G_0^m)\psi$$

$$b_0\phi=0,\beta_0\phi=0,\eta_0\phi=0,b_0\psi=0,\beta_0\psi=0,\eta_0\psi=0$$

$$XY\Psi=\Psi,\eta_0\Psi=0$$

$$S_R^{(0)}=-\frac{1}{2}\langle\Psi,YQ\Psi\rangle'$$

$$\delta|\Psi\rangle=Q|\lambda\rangle,$$

$$XY\lambda=\lambda,\eta_0\lambda=0$$

$$S=-\frac{1}{2}\langle\Psi,YQ\Psi\rangle'-\int_0^1 dt\langle A_t(t),QA_\eta(t)+(F(t)\Psi)^2\rangle_L'$$

$$A_\eta(t)=\eta e^{\Phi(t)}e^{-\Phi(t)}, A_t(t)=\partial_te^{\Phi(t)}e^{-\Phi(t)}$$

$$F(t)\Psi=\Theta(\beta_0)\big\{A_\eta(t),\Theta(\beta_0)\{A_\eta(t),\cdots\Theta(\beta_0)\{A_\eta(t),\Psi\}\cdots\}\big\},$$

Hiperpartículas – Condensación de campo.

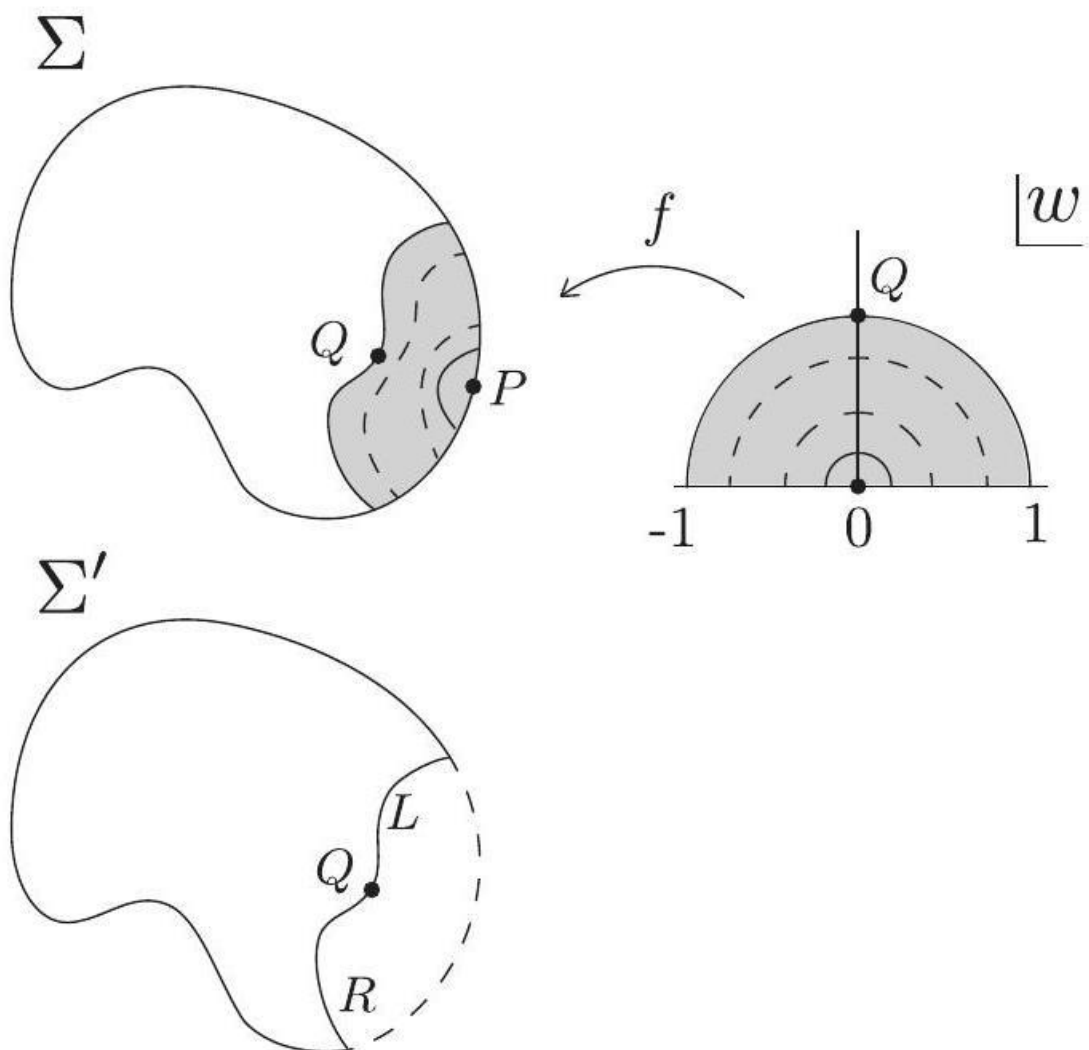
$$-S=Vg_o^{-2}\left[-\frac{1}{2}\phi^2+\frac{1}{3}\left(\frac{3\sqrt{3}}{4}\right)^3\phi^3\right]$$

$$\mathcal{T}=\frac{1}{2\pi^2g_o^2}$$

$$-S=-Vg_o^{-2}\frac{1}{6}\left(\frac{16}{27}\right)^3=-V\mathcal{T}2\pi^2\frac{1}{6}\left(\frac{16}{27}\right)^3\simeq-0.68V\mathcal{T}$$

$$\langle \Sigma \mid \phi \rangle' = \langle \phi(w=0) \rangle_\Sigma' = \langle f \circ \phi(0) \rangle_\Sigma'$$





$$z = \frac{2}{\pi} \arctan w$$

$$|\Omega_\alpha\rangle \star |\Omega_\beta\rangle = |\Omega_{\alpha+\beta}\rangle,$$

$$\langle \phi, \Omega_\alpha \rangle' = \langle f \circ \phi(0) \rangle'_{C_{\alpha+1}}$$

$$\langle \phi, \Omega_\alpha \rangle' = \langle f_\alpha \circ f \circ \phi(0) \rangle'_{C_1}$$

$$f_{\alpha+\delta\alpha} = f_{\delta\alpha/(1+\alpha)} \circ f_\alpha$$

$$\langle \phi, \Omega_{\alpha+\delta\alpha} \rangle' = \langle f_{\delta\alpha/(1+\alpha)} \circ f_\alpha \circ f \circ \phi(0) \rangle'_{C_1}$$

$$\langle \phi, \Omega_{\alpha+\delta\alpha} \rangle' = \langle \phi, \Omega_\alpha \rangle' - \frac{\delta\alpha}{1+\alpha} \left\langle \oint_0 \frac{dz}{2\pi i} z T(z) f_\alpha \circ f \circ \phi(0) \right\rangle'_{C_1}$$

$$= \langle \phi, \Omega_\alpha \rangle' - \frac{\delta\alpha}{1+\alpha} \left\langle \oint_0 \frac{dz}{2\pi i} z T(z) f \circ \phi(0) \right\rangle'_{C_{\alpha+1}}$$

$$\left\langle \phi, \frac{d}{d\alpha} \Omega_\alpha \right\rangle' = -\frac{1}{\alpha+1} \left\langle \phi, \frac{dz}{2\pi i} z T(z) f \circ \phi(0) \right\rangle'_{C_{\alpha+1}}$$

$$\left\langle \phi, \frac{d}{d\alpha} \Omega_\alpha \right\rangle' = - \left\langle \int_R \frac{dz}{2\pi i} T(z) f \circ \phi(0) \right\rangle'_{C_{\alpha+1}}$$

$$\frac{d}{d\alpha}\Omega_\alpha=-\Omega_\alpha\mathcal{K}=-\mathcal{K}\Omega_\alpha$$

$$\mathcal{K}=\left(\int_L T\right)\mathcal{I}$$

$$\Omega_\alpha=e^{-\alpha\mathcal{K}}$$

$$B=\left(\int_L b\right)\mathcal{I}$$

$$\begin{array}{l} c^2=0, B^2=0, \{c,B\}=1 \\ [\mathcal{K},B]=0, [\mathcal{K},c]=\partial c \\ Q\mathcal{K}=0, QB=\mathcal{K}, Qc=c\mathcal{K}c \end{array}$$

$$\Psi = Fc \frac{\mathcal{K}}{1-F^2} BcF, \text{ with } F = F(\mathcal{K})$$

$$(1-FBcF)^{-1}=1+\frac{F}{1-F^2}BcF$$

$$\Psi=(1-FBcF)Q(1-FBcF)^{-1}$$

$$A=\frac{1-F^2}{\mathcal{K}}B, Q_\Psi A=\mathcal{I}$$

$$F(\mathcal{K})=e^{-\mathcal{K}/2}=\Omega_{1/2}$$

$$\Psi=\Omega_{1/2}c\frac{\mathcal{K}}{1-\Omega_1}Bc\Omega_{1/2}=\Omega_{1/2}c(1+\Omega_1+\Omega_2+\cdots)\mathcal{K}Bc\Omega_{1/2}$$

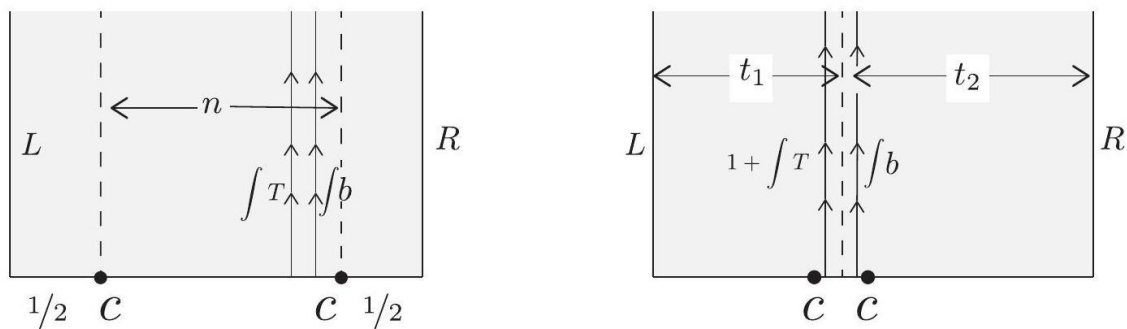
$$=\Omega_{1/2}c\mathcal{K}Bc\Omega_{1/2}+\Omega_{1/2}c\left(\sum_{n=1}^{\infty}\Omega_n\right)\mathcal{K}Bc\Omega_{1/2}$$

$$\Psi=\frac{1}{\sqrt{1+\mathcal{K}}}c(1+\mathcal{K})Bc\frac{1}{\sqrt{1+\mathcal{K}}}$$

$$\frac{1}{\sqrt{1+\mathcal{K}}}=\frac{1}{\sqrt{\pi}}\int_0^{\infty}\frac{dt}{\sqrt{t}}e^{-t}\Omega_t$$

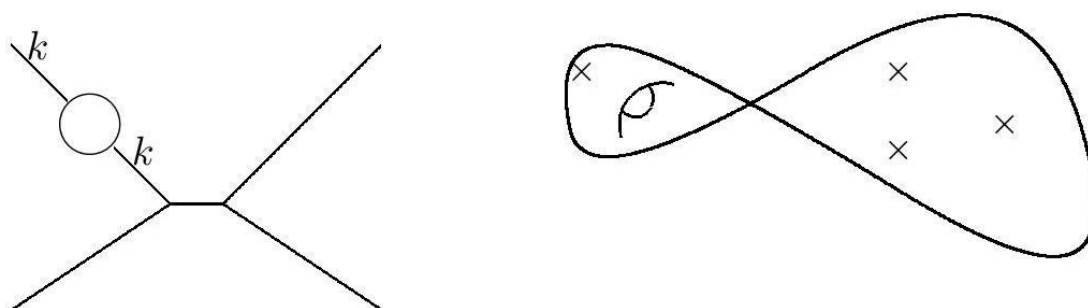
$$\Psi=\frac{1}{\pi}\int_0^{\infty}\int_0^{\infty}\frac{dt_1}{\sqrt{t_1}}\frac{dt_2}{\sqrt{t_2}}e^{-(t_1+t_2)}\Omega_{t_1}c(1+\mathcal{K})Bc\Omega_{t_2}$$





$$-4\pi i \langle E[\mathcal{O}_c] \mid \hat{\Phi} \rangle = \langle \mathcal{O}_c \mid c_0^- \mid B_{\hat{\Phi}} \rangle - \langle \mathcal{O}_c \mid c_0^- \mid B_0 \rangle.$$

Renormalización de masa.



$$\frac{1}{L_0 + \bar{L}_0} \delta_{L_0, \bar{L}_0} = \frac{1}{2\pi} \int_{|q| \leq 1} \frac{d^2 q}{|q|^2} q^{L_0} \bar{q}^{\bar{L}_0}$$



$$S=-\int d^Dx\left[\frac{1}{2}\partial_\mu\phi\partial^\mu\phi+\frac{g_s}{3!}\phi^3\right]$$

$$S=-\int d^Dx\left[\frac{1}{2}\partial_\mu\phi\partial^\mu\phi+\frac{g_s}{3!}\phi^3+c g_s\phi+\mathcal{O}(\phi^2)\right]$$

$$\phi=\pm\sqrt{-2c}+\mathcal{O}(g_s)$$

$$Q(|\Psi\rangle-\mathcal{G}|\tilde{\Psi}\rangle)=0$$

$$Q|\tilde{\Psi}\rangle+\sum_{n=0}^{\infty}\frac{1}{n!}[\Psi^n]_{1\text{PI}}=0$$

$$|\tilde{\Psi}\rangle = |\Psi\rangle,$$

$$Q|\Psi\rangle + \sum_{n=0}^{\infty} \frac{1}{n!} [\Psi^n]_{1\text{PI}} = 0$$

$$|\Psi_{k+1}\rangle = -\frac{b_0^+}{L_0^+} \sum_{n=0}^{k+1} \frac{1}{n!} [\Psi_k^n]_{1\text{PI}} + \mathcal{O}(g_s^{k+2})$$

$$|\Psi_{k+1}\rangle = -\frac{b_0^+}{L_0^+} (1-\mathbf{P}) \sum_{n=0}^{k+1} \frac{1}{n!} [\Psi_k^n]_{1\text{PI}} + |\psi_{k+1}\rangle + \mathcal{O}(g_s^{k+2}), \mathbf{P}|\psi_{k+1}\rangle = |\psi_{k+1}\rangle$$

$$Q|\psi_{k+1}\rangle = -\mathbf{P} \sum_{n=0}^{k+1} \frac{1}{n!} [\Psi_k^n]_{1\text{PI}} + \mathcal{O}(g_s^{k+2})$$

$$|\Psi\rangle = |\Psi_v\rangle + |\Phi\rangle, |\tilde{\Psi}\rangle = |\Psi_v\rangle + |\tilde{\Phi}\rangle$$

$$Q(|\Phi\rangle - \mathcal{G}|\tilde{\Phi}\rangle) = 0, Q|\tilde{\Phi}\rangle + K|\Phi\rangle = 0$$

$$K|A\rangle \equiv \sum_{n=0}^{\infty} \frac{1}{n!} [\Psi_v^n A]$$

$$(Q+\mathcal{G}K)|\Phi\rangle=0$$

$$QK+KQ+KGK=0$$

$$2L_0^+=\left(k^2+\hat{M}^2\right)$$

$$|\Phi_0\rangle=|\phi_n\rangle, |\Phi_{\ell+1}\rangle=-\frac{b_0^+}{L_0^+}(1-P)\mathcal{G}K|\Phi_{\ell}\rangle+|\phi_n\rangle$$

$$P|\phi_n\rangle=|\phi_n\rangle, Q|\phi_n\rangle=-P\mathcal{G}K|\Phi_{n-1}\rangle+\mathcal{O}(g_s^{n+1})$$

$$(Q+\mathcal{G}K)|\Phi_n\rangle=\mathcal{O}(g_s^{n+1})$$

$$N\equiv \exp\left[\int_0^\infty \frac{dt}{t} F(t)\right]$$

$$F(t)=\frac{1}{2}\mathrm{Tr}\{(-1)^fe^{-2\pi tL_0}\},$$

$$N=\exp\left[\int_0^\infty\frac{dt}{2t}\left(\sum_b e^{-2\pi th_b}-\sum_f e^{-2\pi th_f}\right)\right]$$

$$\int_0^\infty \frac{dt}{2t} (e^{-2\pi h_b t} - e^{-2\pi h_f t}) = \ln \sqrt{\frac{h_f}{h_b}}$$

$$N=\prod_{b,f}\sqrt{\frac{h_f}{h_b}}=(\text{sdet}(L_0))^{-1/2}$$

$$\text{sdet}M=1, \text{ with } M_{rs}\equiv \langle \phi_r|c_0|\phi_s\rangle'.$$

$$\int \frac{d\phi_b}{\sqrt{2\pi}}e^{-\frac{1}{2}h_b\phi_b^2}=h_b^{-1/2},\int dp_f dq_f e^{-h_f p_f q_f}=h_f$$

$$|\psi_o\rangle=\sum_r|\phi_r\rangle\psi_r$$

$$N=(\text{sdet}(L_0))^{-1/2}=\int \prod_r D\psi_r \exp\left[\frac{1}{2}\langle \psi_o|c_0L_0|\psi_o\rangle'\right],$$

$$\xi^\mu=\frac{1}{\sqrt{2\pi}g_o}\delta y^\mu$$

$$|\psi_o\rangle=p|0\rangle+qc_1c_{-1}|0\rangle+\cdots$$

$$\frac{1}{2}\langle \psi_o|c_0L_0|\psi_o\rangle'=-hpq+\cdots$$

$$|\psi_o\rangle=i\phi c_0|0\rangle+\cdots$$

$$\delta\phi=h\theta$$

$$\int dpdq\rightarrow \int d\phi e^{S_{\phi,\xi^0}}/\int d\theta$$

$$S_{\phi,\xi^0}=\frac{1}{2}\langle \psi_o|Q|\psi_o\rangle'\Big|_{\phi,\xi^0}=-\left(\phi-\sqrt{\frac{h}{2}}\xi^0\right)^2$$

$$\delta\phi=h,\delta\xi^0=\sqrt{2h}$$

$$\int d\phi e^{-\phi^2}=\sqrt{\pi}$$

$$\theta=\alpha/g_o$$

$$\int d\theta=2\pi/g_o$$

$$\int dpdq\prod_{\mu}\frac{d\xi^{\mu}}{\sqrt{2\pi}}\rightarrow (2\pi\sqrt{\pi}g_o)^{-d}\frac{g_o}{2\pi}\sqrt{\pi}\int\prod_{\mu}dy^{\mu}$$

$$N=(2\pi\sqrt{\pi}g_o)^{-d}\frac{g_o}{2\pi}\sqrt{\pi}\prod_b{}'h_b^{-1/2}\prod_f{}'h_f^{1/2}(2\pi)^d\delta^{(d)}\left(\sum_ip_i\right)$$



$$\prod_b^{' } h_b^{-1/2} \prod_f^{' } h_f^{1/2} = \prod_b \prod_b^{'' } h_b^{-1/2} \prod_f^{'' } h_f^{1/2} \exp \left[- \int_0^\infty \frac{dt}{2t} \left(\sum_b^{'''' } e^{-2\pi t h_b} - \sum_f^{'''' } e^{-2\pi t h_f} \right) \right]$$

$$\mathcal{N}(e^{S_1}-1)$$

$$S=S_0+\mathcal{N}(e^{S_1}-1)+\mathcal{O}(\mathcal{N}^2)$$

$$\frac{1}{2}\{S_0,S_0\}+\Delta S_0=0$$

$$\frac{1}{2}\{S_0+S_1,S_0+S_1\}+\Delta(S_0+S_1)=0$$

$$\{S_0,S_1\}+\frac{1}{2}\{S_1,S_1\}+\Delta S_1=0$$

$$\frac{1}{2}\{S,S\}=\frac{1}{2}\{S_0,S_0\}+\mathcal{N}\{S_0,S_1\}e^{S_1}+\mathcal{O}(\mathcal{N}^2)$$

$$\Delta S=\Delta S_0+\mathcal{N}\Delta e^{S_1}=\Delta S_0+\mathcal{N}\left(\frac{1}{2}\{S_1,S_1\}+\Delta S_1\right)e^{S_1}$$

$$\frac{1}{2}\{S,S\}+\Delta S=\mathcal{O}(\mathcal{N}^2)$$

$$\frac{1}{2}\{S_0+2S_1+S_2,S_0+2S_1+S_2\}+\Delta(S_0+2S_1+S_2)=0$$

$$\{S_0,S_2\}+2\{S_1,S_2\}+\{S_1,S_1\}+\frac{1}{2}\{S_2,S_2\}=0$$

$$S=S_0+\mathcal{N}(e^{S_1}-1)+\frac{1}{2}\mathcal{N}^2(e^{S_2}-1)e^{2S_1}+\mathcal{O}(\mathcal{N}^3)$$

$$\frac{1}{2}\{S,S\}+\Delta S=\mathcal{O}(\mathcal{N}^3)$$

Simetrias.

$$T-T^\dagger=i\sum_nT^\dagger|n\rangle\langle n|T\Leftrightarrow\sum_nS^\dagger|n\rangle\langle n|S=1$$

$$\bar{\Delta}(A_1\otimes\cdots\otimes A_n)=\sum_{k=1}^{n-1}(A_1\otimes\cdots A_k)\otimes'(A_{k+1}\cdots A_n),n\geqslant 2,\bar{\Delta}(A_1)=0$$



Apéndice A.

Gravedad cuántica endógena, es decir, por acción de la partícula supermasiva, a razón de su masa extremadamente densa. Modelo Weyl – Polyakov.

$$d^2z \equiv 2dxdy, \delta^2(z, \bar{z}) \equiv \frac{1}{2} \delta(x) \delta(y)$$

$$M_P = G_N^{-1/2} = 1.22 \times 10^{19} \text{GeV}$$

$$M_P^{-1} = 1.6 \times 10^{-33} \text{ cm.}$$

$$G_N^2 E^2 \int dE' E' = \frac{E^2}{M_P^4} \int dE' E'$$

$$X'^\mu(\tau'(\tau)) = X^\mu(\tau)$$

$$S_{pp} = -m \int d\tau (-\dot{X}^\mu \dot{X}_\mu)^{1/2}$$

$$\delta S_{pp} = -m \int d\tau \dot{u}_\mu \delta X^\mu$$

$$u^\mu = \dot{X}^\mu (-\dot{X}^\nu \dot{X}_\nu)^{-1/2}$$

$$S'_{pp} = \frac{1}{2} \int d\tau (\eta^{-1} \dot{X}^\mu \dot{X}_\mu - \eta m^2)$$

$$\eta'(\tau') d\tau' = \eta(\tau) d\tau$$

$$\eta^2 = -\dot{X}^\mu \dot{X}_\mu / m^2$$

$$h_{ab} = \partial_a X^\mu \partial_b X_\mu$$

$$S_{NG} = \int_M d\tau d\sigma \mathcal{L}_{NG}$$

$$\mathcal{L}_{NG} = -\frac{1}{2\pi\alpha'} (-\text{deth}_{ab})^{1/2}$$

$$T = \frac{1}{2\pi\alpha'}$$

$$X'^\mu(\tau, \sigma) = \Lambda^\mu_\nu X^\nu(\tau, \sigma) + a^\mu$$

$$X'^\mu(\tau', \sigma') = X^\mu(\tau, \sigma)$$

$$S_P[X, \gamma] = -\frac{1}{4\pi\alpha'} \int_M d\tau d\sigma (-\gamma)^{1/2} \gamma^{ab} \partial_a X^\mu \partial_b X_\mu$$



$$\delta_\gamma S_P[X, \gamma] = -\frac{1}{4\pi\alpha'} \int_M d\tau d\sigma (-\gamma)^{1/2} \delta\gamma^{ab} \left(h_{ab} - \frac{1}{2} \gamma_{ab} \gamma^{cd} h_{cd} \right)$$

$$\delta\gamma = \gamma\gamma^{ab}\delta\gamma_{ab} = -\gamma\gamma_{ab}\delta\gamma^{ab}$$

$$h_{ab} = \frac{1}{2} \gamma_{ab} \gamma^{cd} h_{cd}$$

$$h_{ab}(-h)^{-1/2} = \gamma_{ab}(-\gamma)^{-1/2}$$

$$S_P[X, \gamma] \rightarrow -\frac{1}{2\pi\alpha'} \int d\tau d\sigma (-h)^{1/2} = S_{NG}[X]$$

$$X'^\mu(\tau, \sigma) = \Lambda^\mu_\nu X^\nu(\tau, \sigma) + a^\mu$$

$$\gamma'_{ab}(\tau, \sigma) = \gamma_{ab}(\tau, \sigma)$$

$$X'^\mu(\tau', \sigma') = X^\mu(\tau, \sigma)$$

$$\frac{\partial\sigma'^c}{\partial\sigma^a} \frac{\partial\sigma'^d}{\partial\sigma^b} \gamma'_{cd}(\tau', \sigma') = \gamma_{ab}(\tau, \sigma)$$

$$X'^\mu(\tau, \sigma) = X^\mu(\tau, \sigma)$$

$$\gamma'_{ab}(\tau, \sigma) = \exp\left(2\omega(\tau, \sigma)\right)\gamma_{ab}(\tau, \sigma)$$

$$T^{ab}(\tau, \sigma) = -4\pi(-\gamma)^{-1/2} \frac{\delta}{\delta\gamma_{ab}} S_P = -\frac{1}{\alpha'} \left(\partial^a X^\mu \partial^b X_\mu - \frac{1}{2} \gamma^{ab} \partial_c X^\mu \partial^c X_\mu \right)$$

$$\gamma_{ab} \frac{\delta}{\delta\gamma_{ab}} S_P = 0 \Rightarrow T^a_a = 0$$

$$T_{ab} = 0$$

$$\partial_a [(-\gamma)^{1/2} \gamma^{ab} \partial_b X^\mu] = (-\gamma)^{1/2} \nabla^2 X^\mu = 0$$

$$-\infty < \tau < \infty, 0 \leq \sigma \leq \ell.$$

$$\delta S_P = \frac{1}{2\pi\alpha'} \int_{-\infty}^{\infty} d\tau \int_0^{\ell} d\sigma (-\gamma)^{1/2} \delta X^\mu \nabla^2 X_\mu$$

$$- \frac{1}{2\pi\alpha'} \int_{-\infty}^{\infty} d\tau (-\gamma)^{1/2} \delta X^\mu \partial^\sigma X_\mu \Big|_{\sigma=0}^{\sigma=\ell}$$

$$\partial^\sigma X^\mu(\tau, 0) = \partial^\sigma X^\mu(\tau, \ell) = 0$$

$$n^a \partial_a X_\mu = 0 \text{ on } \partial M$$

$$X^\mu(\tau, \ell) = X^\mu(\tau, 0), \partial^\sigma X^\mu(\tau, \ell) = \partial^\sigma X^\mu(\tau, 0)$$

$$\gamma_{ab}(\tau, \ell) = \gamma_{ab}(\tau, 0)$$

$$\chi = \frac{1}{4\pi} \int_M d\tau d\sigma (-\gamma)^{1/2} R$$



$$(-\gamma')^{1/2}R' = (-\gamma)^{1/2}(R - 2\nabla^2\omega).$$

$$S'_{\rm p} = S_{\rm p} - \lambda \chi = - \int_M d\tau d\sigma (-\gamma)^{1/2} \left(\frac{1}{4\pi\alpha'} \gamma^{ab} \partial_a X^\mu \partial_b X_\mu + \frac{\lambda}{4\pi} R \right)$$

$$x^\pm = 2^{-1/2}(x^0 \pm x^1), x^i, i = 2, \ldots, D-1$$

$$\begin{aligned} a^\mu b_\mu &= -a^+b^- - a^-b^+ + a^ib^i \\ a_- &= -a^+, a_+ = -a^-, a_i = a^i \end{aligned}$$

$$X^+(\tau)=\tau$$

$$S'_{\rm pp} = \frac{1}{2} \int \, d\tau (-2\eta^{-1}\dot{X}^- + \eta^{-1}\dot{X}^i\dot{X}^i - \eta m^2)$$

$$p_-=-\eta^{-1}, p_i=\eta^{-1}\dot{X}^i$$

$$H=p_- \dot{X}^- + p_i \dot{X}^i - L = \frac{p^i p^i + m^2}{2p^+}$$

$$[p_i,X^j]=-i\delta_i^j,[p_-,X^-]=-i$$

$$\begin{aligned} X^+ &= \tau \\ \partial_\sigma \gamma_{\sigma\sigma} &= 0 \\ \text{det} \gamma_{ab} &= -1 \end{aligned}$$

$$f'd\sigma'=fd\sigma$$

$$\left[\begin{array}{cc} \gamma^{\tau\tau} & \gamma^{\tau\sigma} \\ \gamma^{\sigma\tau} & \gamma^{\sigma\sigma} \end{array}\right] = \left[\begin{array}{cc} -\gamma_{\sigma\sigma}(\tau) & \gamma_{\tau\sigma}(\tau,\sigma) \\ \gamma_{\tau\sigma}(\tau,\sigma) & \gamma_{\sigma\sigma}^{-1}(\tau)(1-\gamma_{\tau\sigma}^2(\tau,\sigma)) \end{array}\right].$$

$$L=-\frac{1}{4\pi\alpha'}\int_0^\ell d\sigma[\gamma_{\sigma\sigma}(2\partial_\tau x^- - \partial_\tau X^i\partial_\tau X^i)-2\gamma_{\sigma\tau}(\partial_\sigma Y^- - \partial_\tau X^i\partial_\sigma X^i)+\gamma_{\sigma\sigma}^{-1}(1-\gamma_{\tau\sigma}^2)\partial_\sigma X^i\partial_\sigma X^i]$$

$$\begin{aligned} x^-(\tau) &= \frac{1}{\ell} \int_0^\ell d\sigma X^-(\tau,\sigma) \\ Y^-(\tau,\sigma) &= X^-(\tau,\sigma) - x^-(\tau) \end{aligned}$$

$$\gamma_{\tau\sigma}\partial_\tau X^\mu-\gamma_{\tau\tau}\partial_\sigma X^\mu=0\,\,\text{at}\,\,\sigma=0,\ell.$$

$$\gamma_{\tau\sigma}=0\,\,\text{at}\,\,\sigma=0,\ell$$

$$\partial_\sigma X^i=0\,\,\text{at}\,\,\sigma=0,\ell$$

$$L=-\frac{\ell}{2\pi\alpha'}\gamma_{\sigma\sigma}\partial_\tau x^-+\frac{1}{4\pi\alpha'}\int_0^\ell d\sigma(\gamma_{\sigma\sigma}\partial_\tau X^i\partial_\tau X^i-\gamma_{\sigma\sigma}^{-1}\partial_\sigma X^i\partial_\sigma X^i)$$



$$p_- = -p^+ = \frac{\partial L}{\partial(\partial_\tau x^-)} = -\frac{\ell}{2\pi\alpha'}\gamma_{\sigma\sigma}$$

$$\Pi^i = \frac{\delta L}{\delta(\partial_\tau X^i)} = \frac{1}{2\pi\alpha'}\gamma_{\sigma\sigma}\partial_\tau X^i = \frac{p^+}{\ell}\partial_\tau X^i$$

$$H = p_- \partial_\tau x^- - L + \int_0^\ell d\sigma \Pi_i \partial_\tau X^i = \frac{\ell}{4\pi\alpha' p^+} \int_0^\ell d\sigma \left(2\pi\alpha' \Pi^i \Pi_i + \frac{1}{2\pi\alpha'} \partial_\sigma X^i \partial_\sigma X^i \right)$$

$$\partial_\tau x^- = \frac{\partial H}{\partial p_-} = \frac{H}{p^+}, \partial_\tau p^+ = \frac{\partial H}{\partial x^-} = 0$$

$$\partial_\tau X^i = \frac{\delta H}{\delta \Pi^i} = 2\pi\alpha' c \Pi^i, \partial_\tau \Pi^i = -\frac{\delta H}{\delta X^i} = -\frac{c}{2\pi\alpha'} \partial_\sigma^2 X^i$$

$$\partial_\tau^2 X^i = c^2 \partial_\sigma^2 X^i$$

$$X^i(\tau,\sigma) = x^i + \frac{p^i}{p^+}\tau + i(2\alpha')^{1/2} \sum_{\substack{n=-\infty \\ n \neq 0}}^{\infty} \frac{1}{n} \alpha_n^i \exp\left(-\frac{\pi i n c \tau}{\ell}\right) \cos \frac{\pi n \sigma}{\ell}$$

$$x^i(\tau) = \frac{1}{\ell} \int_0^\ell d\sigma X^i(\tau,\sigma)$$

$$p^i(\tau) = \int_0^\ell d\sigma \Pi^i(\tau,\sigma) = \frac{p^+}{\ell} \int_0^\ell d\sigma \partial_\tau X^i(\tau,\sigma)$$

$$[x^-,p^+] = i\eta^{-+} = -i$$

$$[X^i(\sigma),\Pi^j(\sigma')] = i\delta^{ij}\delta(\sigma-\sigma')$$

$$[x^i,p^j] = i\delta^{ij}$$

$$[\alpha_m^i,\alpha_n^j] = m\delta^{ij}\delta_{m,-n}$$

$$\alpha_m^i \sim m^{1/2}a, \alpha_{-m}^i \sim m^{1/2}a^\dagger, m > 0$$

$$p^+|0;k\rangle = k^+|0;k\rangle, p^i|0;k\rangle = k^i|0;k\rangle$$

$$\alpha_m^i|0;k\rangle = 0, m > 0$$

$$|N;k\rangle = \left[\prod_{i=2}^{D-1} \prod_{n=1}^{\infty} \frac{(\alpha_{-n}^i)^{N_{in}}}{(n^{N_{in}} N_{in}!)^{1/2}} \right] |0;k\rangle.$$

$$\mathcal{H} = | \text{ vacuum } \rangle \oplus \mathcal{H}_1 \oplus \mathcal{H}_2 \oplus \ldots$$

$$H = \frac{p^i p^i}{2p^+} + \frac{1}{2p^+ \alpha'} \left(\sum_{n=1}^{\infty} \alpha_{-n}^i \alpha_n^i + A \right)$$

$$A = \frac{D-2}{2} \sum_{n=1}^{\infty} n$$



$$\sum_{n=1}^{\infty}n\rightarrow -\frac{1}{12}$$

$$\exp\left(-\epsilon\gamma_{\sigma\sigma}^{-1/2}|k_{\sigma}|\right)$$

$$A\rightarrow \frac{D-2}{2}\sum_{n=1}^{\infty}n\exp\left[-\epsilon n(\pi/2p^+\alpha'\ell)^{1/2}\right]=\frac{D-2}{2}\left(\frac{2\ell p^+\alpha'}{\epsilon^2\pi}-\frac{1}{12}+O(\epsilon)\right)$$

$$A=\frac{2-D}{24}$$

$$m^2=2p^+H-p^ip^i=\frac{1}{\alpha'}\Big(N+\frac{2-D}{24}\Big),$$

$$N=\sum_{i=2}^{D-1}\sum_{n=1}^{\infty}nN_{in}$$

$$|0;k\rangle,m^2=\frac{2-D}{24\alpha'}$$

$$\alpha_{-1}^i|0;k\rangle,m^2=\frac{26-D}{24\alpha'}$$

$$A=-1,D=26$$

$$S^{ij}=-i\sum_{n=1}^{\infty}\frac{1}{n}(\alpha_{-n}^i\alpha_n^j-\alpha_{-n}^j\alpha_n^i)$$

$$\mathbf{v}=(v^1,0,\ldots,0)+(0,v^2,\ldots,v^{D-1})$$

$$S^{23}\leq 1+\alpha' m^2$$

$$\sigma'=\sigma+s(\tau)\mathrm{mod}\ell$$

$$\gamma_{\tau\sigma}(\tau,0)=0$$

$$\sigma'=\sigma+s\mathrm{mod}\ell.$$

$$X^i(\tau,\sigma)=x^i+\frac{p^i}{p^+}\tau+i\left(\frac{\alpha'}{2}\right)^{1/2}\times\sum_{\substack{n=-\infty\\n\neq0}}^{\infty}\left\{\frac{\alpha_n^i}{n}\exp\left[-\frac{2\pi in(\sigma+c\tau)}{\ell}\right]+\frac{\tilde{\alpha}_n^i}{n}\exp\left[\frac{2\pi in(\sigma-c\tau)}{\ell}\right]\right\}$$

$$\alpha_n^i,\tilde{\alpha}_n^i,x^i,p^i,x^-,p^+$$



$$\begin{aligned}
[x^-, p^+] &= -i \\
[x^i, p^j] &= i\delta^{ij} \\
[\alpha_m^i, \alpha_n^j] &= m\delta^{ij}\delta_{m,-n} \\
[\tilde{\alpha}_m^i, \tilde{\alpha}_n^j] &= m\delta^{ij}\delta_{m,-n}
\end{aligned}$$

$$|N, \tilde{N}; k\rangle = \left[\prod_{i=2}^{D-1} \prod_{n=1}^{\infty} \frac{(\alpha_{-n}^i)^{N_{in}} (\tilde{\alpha}_{-n}^i)^{\tilde{N}_{in}}}{(n^{N_{in}} N_{in}! n^{\tilde{N}_{in}} \tilde{N}_{in}!)^{1/2}} \right] |0, 0; k\rangle.$$

$$m^2=2p^+H-p^ip^i=\frac{2}{\alpha'}\left[\sum_{n=1}^{\infty}(\alpha_{-n}^i\alpha_n^i+\tilde{\alpha}_{-n}^i\tilde{\alpha}_n^i)+A+\tilde{A}\right]=\frac{2}{\alpha'}(N+\tilde{N}+A+\tilde{A})$$

$$A=\tilde{A}=\frac{2-D}{24}$$

$$P=-\int_0^\ell d\sigma \Pi^i \partial_\sigma X^i=-\frac{2\pi}{\ell}\left[\sum_{n=1}^{\infty}(\alpha_{-n}^i\alpha_n^i-\tilde{\alpha}_{-n}^i\tilde{\alpha}_n^i)+A-\tilde{A}\right]=-\frac{2\pi}{\ell}(N-\tilde{N})$$

$$N=\tilde{N}$$

$$|0,0;k\rangle,m^2=\frac{2-D}{6\alpha'}$$

$$\alpha_{-1}^i\tilde{\alpha}_{-1}^j|0,0;k\rangle,m^2=\frac{26-D}{6\alpha'}$$

$$A=\tilde{A}=-1,D=26$$

$$e^{ij}=\frac{1}{2}\Big(e^{ij}+e^{ji}-\frac{2}{D-2}\delta^{ij}e^{kk}\Big)+\frac{1}{2}(e^{ij}-e^{ji})+\frac{1}{D-2}\delta^{ij}e^{kk},$$

$$[\alpha_m^\mu,\alpha_n^\nu]=m\eta^{\mu\nu}\delta_{m,-n}$$

$$\sigma'=\ell-\sigma, \tau'=\tau$$

$$\Omega\alpha_n^i\Omega^{-1}=(-1)^n\alpha_n^i$$

$$\begin{aligned}\Omega\alpha_n^i\Omega^{-1}&=\tilde{\alpha}_n^i\\ \Omega\tilde{\alpha}_n^i\Omega^{-1}&=\alpha_n^i\end{aligned}$$

$$\begin{aligned}\Omega|N;k\rangle &= (-1)^N|N;k\rangle \\ \Omega|N,\tilde{N};k\rangle &= |\tilde{N},N;k\rangle\end{aligned}$$

$$\chi=\frac{1}{4\pi}\int_Md\tau d\sigma(-\gamma)^{1/2}R+\frac{1}{2\pi}\int_{\partial M}dsk$$

$$k=\pm t^an_b\nabla_at^b$$

$$\sum_{n=1}^{\infty}(n-\theta)$$

$$X^{25}(\tau,0)=0,X^{25}(\tau,\ell)=y$$

$$X^{25}(\tau,0)=0,\partial^\sigma X^{25}(\tau,\ell)=0$$

$$X^{25}(\tau,\sigma+\ell)=X^{25}(\tau,\sigma)+2\pi R$$

$$X^{25}(\tau,\sigma+\ell)=-X^{25}(\tau,\sigma)$$

$$S=\frac{1}{4\pi\alpha'}\int\,d^2\sigma(\partial_1X^\mu\partial_1X_\mu+\partial_2X^\mu\partial_2X_\mu)$$

$$z=\sigma^1+i\sigma^2,\bar{z}=\sigma^1-i\sigma^2$$

$$\partial_z=\frac{1}{2}(\partial_1-i\partial_2),\partial_{\bar{z}}=\frac{1}{2}(\partial_1+i\partial_2)$$

$$\partial_zz=1,\partial_z\bar{z}=0,\partial_{\bar{z}}z=0,\partial_{\bar{z}}\bar{z}=1$$

$$v^z=v^1+iv^2,v^{\bar{z}}=v^1-iv^2,v_z=\frac{1}{2}(v^1-iv^2),v_{\bar{z}}=\frac{1}{2}(v^1+iv^2)$$

$$g_{z\bar{z}}=g_{\bar{z}z}=\frac{1}{2},g_{zz}=g_{\bar{z}\bar{z}}=0,g^{z\bar{z}}=g^{\bar{z}z}=2,g^{zz}=g^{\bar{z}\bar{z}}=0$$

$$d^2z=2d\sigma^1d\sigma^2$$

$$\int\,d^2z\delta^2(z,\bar{z})=1$$

$$\int_R d^2z(\partial_zv^z+\partial_{\bar{z}}v^{\bar{z}})=i\oint_{\partial R}(v^zd\bar{z}-v^{\bar{z}}dz)$$

$$S=\frac{1}{2\pi\alpha'}\int\,d^2z\partial X^\mu\bar{\partial}X_\mu$$

$$\partial\bar{\partial}X^\mu(z,\bar{z})=0$$

$$\partial(\bar{\partial}X^\mu)=\bar{\partial}(\partial X^\mu)=0$$

$$\langle \mathcal{F}[X]\rangle = \int [dX] \exp{(-S)} \mathcal{F}[X]$$

$$0=\int\,[dX]\frac{\delta}{\delta X_\mu(z,\bar{z})}\exp{(-S)}=-\int\,[dX]\exp{(-S)}\frac{\delta S}{\delta X_\mu(z,\bar{z})}=-\left\langle\frac{\delta S}{\delta X_\mu(z,\bar{z})}\right\rangle=\frac{1}{\pi\alpha'}\langle\partial\bar{\partial}X^\mu(z,\bar{z})\rangle$$



$$\langle \partial \bar{\partial} X^\mu(z, \bar{z}) \dots \rangle = 0$$

$$\partial \bar{\partial} \hat{X}^\mu(z, \bar{z}) = 0$$

$$\begin{aligned} 0 &= \int [dX] \frac{\delta}{\delta X_\mu(z, \bar{z})} [\exp(-S) X^v(z', \bar{z}')] \\ &= \int [dX] \exp(-S) \left[\eta^{\mu\nu} \delta^2(z - z', \bar{z} - \bar{z}') + \frac{1}{\pi \alpha'} \partial_z \partial_{\bar{z}} X^\mu(z, \bar{z}) X^v(z', \bar{z}') \right] \\ &= \eta^{\mu\nu} \langle \delta^2(z - z', \bar{z} - \bar{z}') \rangle + \frac{1}{\pi \alpha'} \partial_z \partial_{\bar{z}} \langle X^\mu(z, \bar{z}) X^v(z', \bar{z}') \rangle \\ \frac{1}{\pi \alpha'} \partial_z \partial_{\bar{z}} \langle X^\mu(z, \bar{z}) X^v(z', \bar{z}') \dots \rangle &= -\eta^{\mu\nu} \langle \delta^2(z - z', \bar{z} - \bar{z}') \dots \rangle \\ \frac{1}{\pi \alpha'} \partial_z \partial_{\bar{z}} X^\mu(z, \bar{z}) X^v(z', \bar{z}') &= -\eta^{\mu\nu} \delta^2(z - z', \bar{z} - \bar{z}') \end{aligned}$$

$$: X^\mu(z, \bar{z}) := X^\mu(z, \bar{z})$$

$$: X^\mu(z_1, \bar{z}_1) X^v(z_2, \bar{z}_2) := X^\mu(z_1, \bar{z}_1) X^v(z_2, \bar{z}_2) + \frac{\alpha'}{2} \eta^{\mu\nu} \ln |z_{12}|^2$$

$$z_{ij} = z_i - z_j$$

$$\partial_1 \bar{\partial}_1 : X^\mu(z_1, \bar{z}_1) X^v(z_2, \bar{z}_2) : = 0$$

$$\partial \bar{\partial} \ln |z|^2 = 2\pi \delta^2(z, \bar{z}).$$

$$\langle \mathcal{A}_{i_1}(z_1, \bar{z}_1) \mathcal{A}_{i_2}(z_2, \bar{z}_2) \dots \mathcal{A}_{i_n}(z_n, \bar{z}_n) \rangle,$$

$$\mathcal{A}_i(\sigma_1) \mathcal{A}_j(\sigma_2) = \sum_k c_{ij}^k(\sigma_1 - \sigma_2) \mathcal{A}_k(\sigma_2).$$

$$\langle \mathcal{A}_i(\sigma_1) \mathcal{A}_j(\sigma_2) \dots \rangle = \sum_k c_{ij}^k(\sigma_1 - \sigma_2) \langle \mathcal{A}_k(\sigma_2) \dots \rangle$$

$$X^\mu(z_1, \bar{z}_1) X^v(z_2, \bar{z}_2)$$

$$\begin{aligned} &= -\frac{\alpha'}{2} \eta^{\mu\nu} \ln |z_{12}|^2 + : X^v X^\mu(z_2, \bar{z}_2) \\ &+ \sum_{k=1}^{\infty} \frac{1}{k!} \left[(z_{12})^k : X^v \partial^k X^\mu(z_2, \bar{z}_2) : + (\bar{z}_{12})^k : X^v \bar{\partial}^k X^\mu(z_2, \bar{z}_2) : \right] \end{aligned}$$

$$: X^{\mu_1}(z_1, \bar{z}_1) \dots X^{\mu_n}(z_n, \bar{z}_n) : = X^{\mu_1}(z_1, \bar{z}_1) \dots X^{\mu_n}(z_n, \bar{z}_n) + \sum \text{subtractions}$$



$$\begin{aligned}
& :X^{\mu_1}(z_1, \bar{z}_1)X^{\mu_2}(z_2, \bar{z}_2)X^{\mu_3}(z_3, \bar{z}_3): \\
& = X^{\mu_1}(z_1, \bar{z}_1)X^{\mu_2}(z_2, \bar{z}_2)X^{\mu_3}(z_3, \bar{z}_3) \\
& + \left(\frac{\alpha'}{2} \eta^{\mu_1 \mu_2} \ln |z_{12}|^2 X^{\mu_3}(z_3, \bar{z}_3) + 2 \text{ permutations} \right) \\
& : \mathcal{F} := \exp \left(\frac{\alpha'}{4} \int d^2 z_1 d^2 z_2 \ln |z_{12}|^2 \frac{\delta}{\delta X^\mu(z_1, \bar{z}_1)} \frac{\delta}{\delta X_\mu(z_2, \bar{z}_2)} \right) \mathcal{F} \\
& \mathcal{F} = \exp \left(-\frac{\alpha'}{4} \int d^2 z_1 d^2 z_2 \ln |z_{12}|^2 \frac{\delta}{\delta X^\mu(z_1, \bar{z}_1)} \frac{\delta}{\delta X_\mu(z_2, \bar{z}_2)} \right) : \mathcal{F} := : \mathcal{F} : + \sum \text{ contractions} \\
& : \mathcal{F} : : \mathcal{G} : = : \mathcal{F} \mathcal{G} : + \sum \text{ cross-contractions} \\
& : \mathcal{F} : : \mathcal{G} : = \exp \left(-\frac{\alpha'}{2} \int d^2 z_1 d^2 z_2 \ln |z_{12}|^2 \frac{\delta}{\delta X_F^\mu(z_1, \bar{z}_1)} \frac{\delta}{\delta X_{G\mu}(z_2, \bar{z}_2)} \right) : \mathcal{F} \mathcal{G} :, \\
& : \partial X^\mu(z) \partial X_\mu(z) : : \partial' X^\nu(z') \partial' X_\nu(z') : \\
& = : \partial X^\mu(z) \partial X_\mu(z) \partial' X^\nu(z') \partial' X_\nu(z') : - 4 \cdot \frac{\alpha'}{2} (\partial \partial' \ln |z - z'|^2) : \partial X^\mu(z) \partial' X_\mu(z') : + 2 \\
& \cdot \eta_\mu^\mu \left(-\frac{\alpha'}{2} \partial \partial' \ln |z - z'|^2 \right)^2 \\
& \sim \frac{D \alpha'^2}{2(z - z')^4} - \frac{2 \alpha'}{(z - z')^2} : \partial' X^\mu(z') \partial' X_\mu(z') : - \frac{2 \alpha'}{z - z'} : \partial'^2 X^\mu(z') \partial' X_\mu(z') : \\
& \mathcal{F} = e^{ik_1 \cdot X(z, \bar{z})}, \mathcal{G} = e^{ik_2 \cdot X(0,0)}. \\
& : e^{ik_1 \cdot X(z, \bar{z})} : : e^{ik_2 \cdot X(0,0)} : = \exp \left(\frac{\alpha'}{2} k_1 \cdot k_2 \ln |z|^2 \right) : e^{ik_1 \cdot X(z, \bar{z})} e^{ik_2 \cdot X(0,0)} = |z|^{\alpha' k_1 \cdot k_2} : e^{ik_1 \cdot X(z, \bar{z})} e^{ik_2 \cdot X(0,0)} : \\
& : e^{ik_1 \cdot X(z, \bar{z})} : : e^{ik_2 \cdot X(0,0)} : = |z|^{\alpha' k_1 \cdot k_2} : e^{i(k_1 + k_2) \cdot X(0,0)} [1 + O(z, \bar{z})] : . \\
& c_{ij}^k (\sigma_1 - \sigma_2)_{\text{sym}} = \pm c_{ji}^k (\sigma_2 - \sigma_1)_{\text{sym}} \\
& \phi'_\alpha(\sigma) = \phi_\alpha(\sigma) + \delta \phi_\alpha(\sigma), \\
& [d\phi'] \exp(-S[\phi']) = [d\phi] \exp(-S[\phi]). \\
& \phi'_\alpha(\sigma) = \phi_\alpha(\sigma) + \rho(\sigma) \delta \phi_\alpha(\sigma) \\
& [d\phi'] \exp(-S[\phi']) = [d\phi] \exp(-S[\phi]) \left[1 + \frac{i\epsilon}{2\pi} \int d^d \sigma g^{1/2} j^a(\sigma) \partial_a \rho(\sigma) + O(\epsilon^2) \right] \\
& .
\end{aligned}$$

$$0 = \int [d\phi'] \exp(-S[\phi']) \dots - \int [d\phi] \exp(-S[\phi]) \dots = \frac{\epsilon}{2\pi i} \int d^d \sigma g^{1/2} \rho(\sigma) \langle \nabla_a j^a(\sigma) \dots \rangle$$

$$\nabla_a j^a = 0$$

$$\delta \mathcal{A}(\sigma_0) + \frac{\epsilon}{2\pi i} \int_R d^d \sigma g^{1/2} \nabla_a j^a(\sigma) \mathcal{A}(\sigma_0) = 0$$

$$\nabla_a j^a(\sigma) \mathcal{A}(\sigma_0) = g^{-1/2} \delta^d(\sigma - \sigma_0) \frac{2\pi}{i\epsilon} \delta \mathcal{A}(\sigma_0) + \text{total } \sigma\text{-derivative}$$

$$\int_{\partial R} dA n_a j^a \mathcal{A}(\sigma_0) = \frac{2\pi}{i\epsilon} \delta \mathcal{A}(\sigma_0)$$

$$\oint_{\partial R} (j dz - \bar{j} d\bar{z}) \mathcal{A}(z_0, \bar{z}_0) = \frac{2\pi}{\epsilon} \delta \mathcal{A}(z_0, \bar{z}_0)$$

$$\text{Res}_{z \rightarrow z_0} j(z) \mathcal{A}(z_0, \bar{z}_0) + \overline{\text{Res}_{\bar{z} \rightarrow \bar{z}_0} \bar{j}(\bar{z}) \mathcal{A}(z_0, \bar{z}_0)} = \frac{1}{i\epsilon} \delta \mathcal{A}(z_0, \bar{z}_0)$$

$$\delta S = \frac{\epsilon a_\mu}{2\pi \alpha'} \int d^2 \sigma \partial^a X^\mu \partial_a \rho$$

$$j_a^\mu = \frac{i}{\alpha'} \partial_a X^\mu$$

$$j^\mu(z):e^{ik\cdot X(0,0)}:\sim\frac{k^\mu}{2z}:e^{ik\cdot X(0,0)}:$$

$$\bar{j}^\mu(\bar{z}):e^{ik\cdot X(0,0)}:\sim\frac{k^\mu}{2\bar{z}}:e^{ik\cdot X(0,0)}:$$

$$j_a = i v^b T_{ab}$$

$$T_{ab} = -\frac{1}{\alpha'}: \left(\partial_a X^\mu \partial_b X_\mu - \frac{1}{2} \delta_{ab} \partial_c X^\mu \partial^c X_\mu \right):.$$

$$T_{z\bar{z}}=0$$

$$\bar{\partial} T_{zz} = \partial T_{\bar{z}\bar{z}} = 0$$

$$T(z) \equiv T_{zz}(z), \tilde{T}(\bar{z}) \equiv T_{\bar{z}\bar{z}}(\bar{z})$$

$$T(z) = -\frac{1}{\alpha'}: \partial X^\mu \partial X_\mu:, \tilde{T}(\bar{z}) = -\frac{1}{\alpha'}: \bar{\partial} X^\mu \bar{\partial} X_\mu:,$$

$$j(z) = i v(z) T(z), \bar{j}(\bar{z}) = i v(z)^* \tilde{T}(\bar{z})$$

$$T(z) X^\mu(0) \sim \frac{1}{z} \partial X^\mu(0), \tilde{T}(\bar{z}) X^\mu(0) \sim \frac{1}{\bar{z}} \bar{\partial} X^\mu(0)$$

$$\delta X^\mu = -\epsilon v(z) \partial X^\mu - \epsilon v(z)^* \bar{\partial} X^\mu$$



$$X'^{\mu}(z',\bar{z}')=X^{\mu}(z,\bar{z}),z'=f(z)$$

$$z'=\zeta z$$

$$ds'^2=dz'd\bar{z}'=\frac{\partial z'}{\partial z}\frac{\partial \bar{z}'}{\partial \bar{z}}dzd\bar{z}$$

$$T(z)\mathcal{A}(0,0)\sim\sum_{n=0}^{\infty}\frac{1}{z^{n+1}}\mathcal{A}^{(n)}(0,0)$$

$$\delta \mathcal{A}(z,\bar{z})=-\epsilon \sum_{n=0}^{\infty}\frac{1}{n!}[\partial^n v(z)\mathcal{A}^{(n)}(z,\bar{z})+\bar{\partial}^n v(z)^*\tilde{\mathcal{A}}^{(n)}(z,\bar{z})]$$

$$\mathcal{A}'(z',\bar{z}')=\zeta^{-h\bar{\zeta}-\tilde{h}}\mathcal{A}(z,\bar{z})$$

$$T(z)\mathcal{A}(0,0)=\cdots+\frac{h}{z^2}\mathcal{A}(0,0)+\frac{1}{z}\partial\mathcal{A}(0,0)+\cdots,$$

$$\mathcal{O}'(z',\bar{z}')=(\partial_z z')^{-h}(\partial_{\bar{z}}\bar{z}')^{-\tilde{h}}\mathcal{O}(z,\bar{z})$$

$$T(z)\mathcal{O}(0,0)=\frac{h}{z^2}\mathcal{O}(0,0)+\frac{1}{z}\partial\mathcal{O}(0,0)+\cdots$$

$$\left. \begin{array}{ll} X^{\mu} & (0,0), \quad \partial X^{\mu} \\ \bar{\partial} X^{\mu} & (0,1), \quad \partial^2 X^{\mu} \\ :e^{ik\cdot X}: & \left(\frac{\alpha'k^2}{4},\frac{\alpha'k^2}{4}\right) \end{array} \right\}$$

$$:\Big(\prod_i\partial^{m_i}X^{\mu_i}\Big)\Big(\prod_j\bar{\partial}^{n_j}X^{v_j}\Big)e^{ik\cdot X}:,$$

$$\left(\frac{\alpha'k^2}{4}+\sum_im_i,\frac{\alpha'k^2}{4}+\sum_jn_j\right)$$

$$\mathcal{A}_i(z_1,\bar{z}_1)\mathcal{A}_j(z_2,\bar{z}_2)=\sum_kz_{12}^{h_k-h_i-h_j}\bar{z}_{12}^{\tilde{h}_k-\tilde{h}_i-\tilde{h}_j}c_{ij}^k\mathcal{A}_k(z_2,\bar{z}_2)$$

$$\delta e^{ik\cdot X}=-\epsilon v(z)\partial e^{ik\cdot X}-\epsilon v(z)^*\bar{\partial} e^{ik\cdot X}$$

$$T(z)T(0)=\frac{\eta_{\mu}^{\mu}}{2z^4}-\frac{2}{\alpha'z^2}:\partial X^{\mu}(z)\partial X_{\mu}(0):+ :T(z)T(0): \sim \frac{D}{2z^4}+\frac{2}{z^2}T(0)+\frac{1}{z}\partial T(0)$$

$$\epsilon^{-1}\delta T(z)=-\frac{D}{12}\partial_z^3v(z)-2\partial_zv(z)T(z)-v(z)\partial_zT(z).$$



$$\epsilon^{-1}\delta T(z)=-\frac{c}{12}\partial_z^3v(z)-2\partial_zv(z)T(z)-v(z)\partial_zT(z)$$

$$T(z)T(0)\sim \frac{c}{2z^4}+\frac{2}{z^2}T(0)+\frac{1}{z}\partial T(0)$$

$$(\partial_z z')^2 T'(z')=T(z)-\frac{c}{12}\{z',z\}$$

$$\{f,z\}=\frac{2\partial_z^3f\partial_zf-3\partial_z^2f\partial_z^2f}{2\partial_zf\partial_zf}$$

$$T(z)=-\frac{1}{\alpha'}\colon\partial X^\mu\partial X_\mu\colon+V_\mu\partial^2X^\mu,$$

$$\tilde{T}(\bar{z})=-\frac{1}{\alpha'}\colon\bar{\partial}X^\mu\bar{\partial}X_\mu\colon+V_\mu\bar{\partial}^2X^\mu,$$

$$c=\tilde{c}=D+6\alpha'V_\mu V^\mu$$

$$\delta X^\mu=-\epsilon v\partial X^\mu-\epsilon v^*\bar{\partial}X^\mu-\frac{\epsilon}{2}\alpha'V^\mu[\partial v+(\partial v)^*]$$

$$S=\frac{1}{2\pi}\int\,d^2zb\bar{\partial}c$$

$$h_b=\lambda,h_c=1-\lambda$$

$$\bar{\partial}c(z)=\bar{\partial}b(z)=0\\ \bar{\partial}b(z)c(0)=2\pi\delta^2(z,\bar{z})$$

$$\colon b(z_1)c(z_2)\colon = b(z_1)c(z_2) - \frac{1}{z_{12}}$$

$$\bar{\partial}\frac{1}{z}=\partial\frac{1}{\bar{z}}=2\pi\delta^2(z,\bar{z})$$

$$b(z_1)c(z_2)\sim \frac{1}{z_{12}},c(z_1)b(z_2)\sim \frac{1}{z_{12}}$$

$$b(z_1)b(z_2)=O(z_{12}),c(z_1)c(z_2)=O(z_{12})$$

$$T(z)=:(\partial b)c:-\lambda\partial(:bc:)\nonumber\\ \tilde{T}(\bar{z})=0$$

$$c=-3(2\lambda-1)^2+1,\tilde{c}=0$$

$$S=\frac{1}{2\pi}\int\,d^2z\tilde{b}\partial\tilde{c}$$

$$j=-\colon bc\colon$$

$$T(z)j(0)\sim \frac{1-2\lambda}{z^3}+\frac{1}{z^2}j(0)+\frac{1}{z}\partial j(0).$$



$$\epsilon^{-1}\delta j=-v\partial j-j\partial v+\frac{2\lambda-1}{2}\partial^2v,$$

$$(\partial_z z')j_{z'}(z')=j_z(z)+\frac{2\lambda-1}{2}\frac{\partial_z^2 z'}{\partial_z z'}.$$

$$\psi\,=\,2^{-1/2}(\psi_1+i\psi_2),\bar{\psi}\,=\,2^{-1/2}(\psi_1-i\psi_2)$$

$$S\,=\,\frac{1}{4\pi}\int\,d^2z(\psi_1\bar{\partial}\psi_1+\psi_2\bar{\partial}\psi_2)$$

$$T\,=\,-\frac{1}{2}\psi_1\partial\psi_1-\frac{1}{2}\psi_2\partial\psi_2$$

$$\beta\gamma_{CFT}$$

$$h_\beta=\lambda, h_\gamma=1-\lambda.$$

$$S=\frac{1}{2\pi}\int\,d^2z\beta\bar{\partial}\gamma$$

$$\bar{\partial}\gamma(z)=\bar{\partial}\beta(z)=0.$$

$$\beta(z_1)\gamma(z_2)\sim-\frac{1}{z_{12}},\gamma(z_1)\beta(z_2)\sim\frac{1}{z_{12}}.$$

$$T =: (\partial \beta) \gamma : - \lambda \partial (: \beta \gamma :) \\ \tilde{T} = 0$$

$$c=3(2\lambda-1)^2-1, \tilde{c}=0$$

$$\sigma^1\sim\sigma^1+2\pi$$

$$-\infty<\sigma^2<\infty$$

$$w=\sigma^1+i\sigma^2$$

$$z=\exp{(-iw)}=\exp{(-i\sigma^1+\sigma^2)}$$

$$T_{zz}(z)=\sum_{m=-\infty}^{\infty}\frac{L_m}{z^{m+2}},\tilde{T}_{\bar{z}\bar{z}}(\bar{z})=\sum_{m=-\infty}^{\infty}\frac{\tilde{L}_m}{\bar{z}^{m+2}}.$$

$$L_m=\oint_c\frac{dz}{2\pi iz}z^{m+2}T_{zz}(z)$$

$$T_{ww}(w)=-\sum_{m=-\infty}^{\infty}\exp{(im\sigma^1-m\sigma^2)}T_m$$

$$T_{\bar{w}\bar{w}}(\bar{w})=-\sum_{m=-\infty}^{\infty}\exp{(-im\sigma^1-m\sigma^2)}\tilde{T}_m$$

$$T_m=L_m-\delta_{m,0}\frac{c}{24},\tilde{T}_m=\tilde{L}_m-\delta_{m,0}\frac{\tilde{c}}{24}$$

$$T_{ww} = (\partial_w z)^2 T_{zz} + \frac{c}{24}$$

$$H = \int_0^{2\pi} \frac{d\sigma^1}{2\pi} T_{22} = L_0 + \tilde{L}_0 - \frac{c + \tilde{c}}{24}$$

$$Q_i\{C\}=\oint_c\frac{dz}{2\pi i}j_i$$

$$Q_1\{C_1\}Q_2\{C_2\}-Q_1\{C_3\}Q_2\{C_2\}$$

$$\hat{Q}_1\hat{Q}_2-\hat{Q}_2\hat{Q}_1\equiv[\hat{Q}_1,\hat{Q}_2]$$

$$[Q_1,Q_2]\{C_2\}=\oint_{c_2}\frac{dz_2}{2\pi i}\mathrm{Res}_{z_1\rightarrow z_2}j_1(z_1)j_2(z_2)$$

$$[Q,\mathcal{A}(z_2,\bar{z}_2)]=\mathrm{Res}_{z_1\rightarrow z_2}j(z_1)\mathcal{A}(z_2,\bar{z}_2)=\frac{1}{i\epsilon}\delta\mathcal{A}(z_2,\bar{z}_2)$$

$$\tilde{Q}\{C\}=-\oint_c\frac{d\bar{z}}{2\pi i}\tilde{j}$$

$$[\tilde{Q},\mathcal{A}(z_2,\bar{z}_2)]=\overline{\mathrm{Res}}_{\bar{z}_1\rightarrow\bar{z}_2}\tilde{j}(\bar{z}_1)\mathcal{A}(z_2,\bar{z}_2)=\frac{1}{i\epsilon}\delta\mathcal{A}(z_2,\bar{z}_2)$$

$$\begin{aligned}\mathrm{Res}_{z_1\rightarrow z_2}z_1^{m+1}T(z_1)z_2^{n+1}T(z_2)&=\mathrm{Res}_{z_1\rightarrow z_2}z_1^{m+1}z_2^{n+1}\left(\frac{c}{2z_{12}^4}+\frac{2}{z_{12}^2}T(z_2)+\frac{1}{z_{12}}\partial T(z_2)\right)\\&=\frac{c}{12}(\partial^3z_2^{m+1})z_2^{n+1}+2(\partial z_2^{m+1})z_2^{n+1}T(z_2)+z_2^{m+n+2}\partial T(z_2)=\frac{c}{12}(m^3-m)z_2^{m+n-1}+(m-n)z_2^{m+n+1}T(z_2)\end{aligned}$$

$$[L_m,L_n]=(m-n)L_{m+n}+\frac{c}{12}(m^3-m)\delta_{m,-n}$$

$$[L_0,L_n]=-nL_n$$

$$L_0L_n|\psi\rangle=L_n(L_0-n)|\psi\rangle=(h-n)L_n|\psi\rangle$$

$$[L_0,L_1]=-L_1,[L_0,L_{-1}]=L_{-1},[L_1,L_{-1}]=2L_0$$

$$\mathcal{O}(z)=\sum_{m=-\infty}^{\infty}\frac{\mathcal{O}_m}{z^{m+h}}$$

$$[L_m,\mathcal{O}_n]=[(h-1)m-n]\mathcal{O}_{m+n}$$

$$0\leq \mathrm{Re} w\leq \pi\iff \mathrm{Im} z\geq 0$$

$$T_{ab}n^at^b=0$$

$$T_{ww}=T_{\bar{w}\bar{w}}, \mathrm{Re} w=0, \pi \Leftrightarrow T_{zz}=T_{\bar{z}\bar{z}}, \mathrm{Im} z=0$$

$$T_{zz}(z)\equiv T_{\bar{z}\bar{z}}(\bar{z}'), \mathrm{Im} z<0$$



$$L_m = \frac{1}{2\pi i} \int_c (dz z^{m+1} T_{zz} - d\bar{z} \bar{z}^{m+1} T_{\bar{z}\bar{z}}) = \frac{1}{2\pi i} \oint dz z^{m+1} T_{zz}(z)$$

$$[L_m, L_n] = (m-n)L_{m+n} + \frac{c}{12}(m^3-m)\delta_{m,-n}$$

$$\partial X^\mu(z) = -i \left(\frac{\alpha'}{2}\right)^{1/2} \sum_{m=-\infty}^{\infty} \frac{\alpha_m^\mu}{z^{m+1}}, \bar{\partial} X^\mu(\bar{z}) = -i \left(\frac{\alpha'}{2}\right)^{1/2} \sum_{m=-\infty}^{\infty} \frac{\tilde{\alpha}_m^\mu}{\bar{z}^{m+1}}.$$

$$\alpha_m^\mu = \left(\frac{2}{\alpha'}\right)^{1/2} \oint \frac{dz}{2\pi} z^m \partial X^\mu(z)$$

$$\tilde{\alpha}_m^\mu = -\left(\frac{2}{\alpha'}\right)^{1/2} \oint \frac{d\bar{z}}{2\pi} \bar{z}^m \bar{\partial} X^\mu(\bar{z})$$

$$p^\mu = \frac{1}{2\pi i} \oint_c (dz j^\mu - d\bar{z} \bar{j}^\mu) = \left(\frac{2}{\alpha'}\right)^{1/2} \alpha_0^\mu = \left(\frac{2}{\alpha'}\right)^{1/2} \tilde{\alpha}_0^\mu$$

$$X^\mu(z,\bar{z})=x^\mu-i\frac{\alpha'}{2}p^\mu\ln|z|^2+i\left(\frac{\alpha'}{2}\right)^{1/2}\sum_{\substack{m=-\infty\\m\neq 0}}^{\infty}\frac{1}{m}\left(\frac{\alpha_m^\mu}{z^m}+\frac{\tilde{\alpha}_m^\mu}{\bar{z}^m}\right)$$

$$[\alpha_m^\mu,\alpha_n^\nu]=[\tilde{\alpha}_m^\mu,\tilde{\alpha}_n^\nu]=m\delta_{m,-n}\eta^{\mu\nu}$$

$$[x^\mu,p^\nu]=i\eta^{\mu\nu}$$

$$L_m\sim\frac{1}{2}\sum_{n=-\infty}^{\infty}\alpha_{m-n}^\mu\alpha_{\mu n}$$

$$L_0=\frac{\alpha'p^2}{4}+\sum_{n=1}^{\infty}(\alpha_{-n}^\mu\alpha_{\mu n})+a^X$$

$$2L_0|0;0\rangle=(L_1L_{-1}-L_{-1}L_1)|0;0\rangle=0$$

$$a^X=0$$

$$L_m=\frac{1}{2}\sum_{n=-\infty}^{\infty}\alpha_{m-n}^\mu\alpha_{\mu n}^\circ$$

$$X^\mu(z,\bar{z})X^\nu(z',\bar{z}')=\mathop{\circ}\nolimits X^\mu(z,\bar{z})X^\nu(z',\bar{z}')\mathop{\circ}\nolimits_2+\frac{\alpha'}{2}\eta^{\mu\nu}\left[-\ln|z|^2+\sum_{m=1}^{\infty}\frac{1}{m}\left(\frac{z'^m}{z^m}+\frac{\bar{z}'^m}{\bar{z}^m}\right)\right]$$

$$=\mathop{\circ}\nolimits X^\mu(z,\bar{z})X^\nu(z',\bar{z}')\mathop{\circ}\nolimits-\frac{\alpha'}{2}\eta^{\mu\nu}\ln|z-z'|^2$$

$$\mathop{\circ}\nolimits X^\mu(z,\bar{z})X^\nu(z',\bar{z}')\mathop{\circ}\nolimits=:X^\mu(z,\bar{z})X^\nu(z',\bar{z}'):$$

$$[X^\mu(z,\bar{z})X^\nu(z',\bar{z}')]_1=[X^\mu(z,\bar{z})X^\nu(z',\bar{z}')]_2+\eta^{\mu\nu}\Delta(z,\bar{z},z',\bar{z}'),$$



$$[\mathcal{F}]_1 = \exp \left(\frac{1}{2} \int d^2 z d^2 z' \Delta(z, \bar{z}, z', \bar{z}') \frac{\delta}{\delta X^\mu(z, \bar{z})} \frac{\delta}{\delta X_\mu(z', \bar{z}')} \right) [\mathcal{F}]_2$$

$$L_m = \frac{1}{2} \sum_{n=-\infty}^{\infty} \circ \alpha_{m-n}^\mu \alpha_{\mu n}^\circ + i \left(\frac{\alpha'}{2} \right)^{1/2} (m+1) V^\mu \alpha_{\mu m}$$

$$b(z) = \sum_{m=-\infty}^{\infty} \frac{b_m}{z^{m+\lambda}}, c(z) = \sum_{m=-\infty}^{\infty} \frac{c_m}{z^{m+1-\lambda}}$$

$$\{b_m, c_n\} = \delta_{m,-n}$$

$$\begin{aligned} b_0|\downarrow\rangle &= 0, b_0|\uparrow\rangle = |\downarrow\rangle, \\ c_0|\downarrow\rangle &= |\uparrow\rangle, c_0|\uparrow\rangle = 0, \\ b_n|\downarrow\rangle &= b_n|\uparrow\rangle = c_n|\downarrow\rangle = c_n|\uparrow\rangle = 0, n > 0. \end{aligned}$$

$$L_m = \sum_{n=-\infty}^{\infty} (m\lambda - n)^\circ b_n c_{m-n}^\circ + \delta_{m,0} a^{\mathfrak{g}}$$

$$2L_0|\downarrow\rangle = (L_1L_{-1} - L_{-1}L_1)|\downarrow\rangle = (\lambda b_0c_1)[(1-\lambda)b_{-1}c_0]|\downarrow\rangle = \lambda(1-\lambda)|\downarrow\rangle$$

$$L_m = \sum_{n=-\infty}^{\infty} (m\lambda - n)^\circ b_n c_{m-n}^\circ + \frac{\lambda(1-\lambda)}{2} \delta_{m,0}$$

$$N^{\mathfrak{g}} = -\frac{1}{2\pi i} \int_0^{2\pi} dw j_w = \sum_{n=1}^{\infty} (c_{-n} b_n - b_{-n} c_n) + c_0 b_0 - \frac{1}{2}$$

$$[N^{\mathfrak{g}},b_m]=-b_m, [N^{\mathfrak{g}},c_m]=c_m$$

$$N^{\mathfrak{g}}|\downarrow\rangle = -\frac{1}{2}|\downarrow\rangle, N^{\mathfrak{g}}|\uparrow\rangle = \frac{1}{2}|\uparrow\rangle.$$

$$\alpha_0^\mu = (2\alpha')^{1/2} p^\mu$$

$$X^\mu(z,\bar{z})=x^\mu-i\alpha'p^\mu\ln|z|^2+i\left(\frac{\alpha'}{2}\right)^{1/2}\sum_{\substack{m=-\infty\\m\neq 0}}^{\infty}\frac{\alpha_m^\mu}{m}(z^{-m}+\bar{z}^{-m})$$

$$L_0=\alpha'p^2+\sum_{n=1}^{\infty}\alpha_{-n}^\mu\alpha_{\mu n}$$

$$[\alpha_m^\mu,\alpha_n^\nu]=m\delta_{m,-n}\eta^{\mu\nu}, [x^\mu,p^\nu]=i\eta^{\mu\nu}$$

$$c(z)=\tilde{c}(\bar{z}), b(z)=\tilde{b}(\bar{z}), \mathrm{Im} z=0$$

$$c(z)\equiv \tilde{c}(\bar{z}'), b(z)\equiv \tilde{b}(\bar{z}'), \mathrm{Im}(z)\leq 0, z'= \bar{z}$$

$$0 \leq \text{Re} w \leq 2\pi, w \sim w + 2\pi, \text{Im} w \leq 0$$

$$|1\rangle = |0;0\rangle$$

$$\alpha_{-m}^{\mu}=\left(\frac{2}{\alpha'}\right)^{1/2}\oint\frac{dz}{2\pi}z^{-m}\partial X^{\mu}(z)\rightarrow\left(\frac{2}{\alpha'}\right)^{1/2}\frac{i}{(m-1)!}\partial^mX^{\mu}(0)$$

$$\alpha_{-m}^{\mu}|1\rangle\cong\left(\frac{2}{\alpha'}\right)^{1/2}\frac{i}{(m-1)!}\partial^mX^{\mu}(0),m\geq1$$

$$\tilde{\alpha}_{-m}^{\mu}|1\rangle\cong\left(\frac{2}{\alpha'}\right)^{1/2}\frac{i}{(m-1)!}\bar{\partial}^mX^{\mu}(0),m\geq1.$$

$$\alpha_{-m}^{\mu}:\mathcal{A}(0,0):=\alpha_{-m}^{\mu}\mathcal{A}(0,0):$$

$$\alpha_{-m}^{\mu}\rightarrow i\left(\frac{2}{\alpha'}\right)^{1/2}\frac{1}{(m-1)!}\partial^mX^{\mu}(0),\quad m\geq1$$

$$\tilde{\alpha}_{-m}^{\mu}\rightarrow i\left(\frac{2}{\alpha'}\right)^{1/2}\frac{1}{(m-1)!}\bar{\partial}^mX^{\mu}(0),\quad m\geq1$$

$$x_0^{\mu}\rightarrow X^{\mu}(0,0)$$

$$|0;k\rangle\cong:e^{ik\cdot X(0,0)}.$$

$$b_m|1\rangle=0, m\geq-1, c_m|1\rangle=0, m\geq2$$

$$|1\rangle=b_{-1}|\downarrow\rangle$$

$$b_{-m}\rightarrow \frac{1}{(m-2)!}\partial^{m-2}b(0), m\geq 2$$

$$c_{-m}\rightarrow \frac{1}{(m+1)!}\partial^{m+1}c(0), m\geq -1$$

$$(\partial_z w)j_w(w)=j_z(z)+q_0\frac{\partial_z^2 w}{\partial_z w}=j_z(z)-\frac{q_0}{z}$$

$$Q^{\rm g}\equiv\frac{1}{2\pi i}\oint\,dzj_z=N^{\rm g}+q_0$$

$$0\leq \text{Re} w\leq \pi, \text{Im} w\leq 0$$

$$\Psi_{\mathcal{A}}[\phi_{\rm b}]=\int\,[d\phi_{\rm i}]_{\phi_{\rm b}}\exp\left(-S[\phi_{\rm i}]\right)\mathcal{A}(0)$$

$$\int\,[d\phi'_{\rm b}][d\phi_{\rm i}]_{\phi_{\rm b},\phi'_{\rm b}}\exp\left(-S[\phi_{\rm i}]\right)r^{-L_0-\tilde{L}_0}\Psi[\phi'_{\rm b}]$$

$$X_{\rm b}(\theta)=\sum_{n=-\infty}^{\infty}X_ne^{in\theta}, X_n^*=X_{-n}$$



$$\Psi_1[X_b] = \int [dX_i]_{X_b} \exp \left(-\frac{1}{2\pi\alpha'} \int d^2z \partial X \bar{\partial} X \right)$$

$$X_i = X_{cl} + X'_i$$

$$X_{cl}(z,\bar{z}) = X_0 + \sum_{n=1}^{\infty} (z^n X_n + \bar{z}^n X_{-n})$$

$$\Psi_1[X_b] = \exp(-S_{cl}) \int [dX'_i]_{X_b=0} \exp \left(-\frac{1}{2\pi\alpha'} \int d^2z \partial X' \bar{\partial} X' \right)$$

$$S_{cl} = \frac{1}{\alpha'} \sum_{m=1}^{\infty} m X_m X_{-m} = \frac{1}{\alpha'} \sum_{m=1}^{\infty} m X_m X_{-m}$$

$$\Psi_1[X_b] \propto \exp \left(-\frac{1}{\alpha'} \sum_{m=1}^{\infty} m X_m X_{-m} \right)$$

$$\alpha_n = -\frac{in}{(2\alpha')^{1/2}} X_{-n} - i \left(\frac{\alpha'}{2} \right)^{1/2} \frac{\partial}{\partial X_n}$$

$$\tilde{\alpha}_n = -\frac{in}{(2\alpha')^{1/2}} X_n - i \left(\frac{\alpha'}{2} \right)^{1/2} \frac{\partial}{\partial X_{-n}}$$

$$\alpha_n \Psi_1[X_b] = \tilde{\alpha}_n \Psi_1[X_b] = 0, n \geq 0,$$

$$|1\rangle \propto |0;0\rangle.$$

$$|\partial^k X\rangle = k! X_k \Psi_1 = -i \left(\frac{\alpha'}{2} \right)^{1/2} (k-1)! \alpha_{-k} |0;0\rangle,$$

$$\mathcal{A}_i(z,\bar{z})\mathcal{A}_j(0,0),$$

$$\mathcal{A}_{ij,z,\bar{z}} = \sum_k z^{h_k-h_i-h_j\bar{z}} \bar{h}_k-\bar{h}_i-\bar{h}_j c_{ij}^k \mathcal{A}_k,$$

$$\sum_l \begin{array}{c} i \\ \diagdown \\ \bullet \\ \diagup \\ k \end{array} \begin{array}{c} \text{---} l \text{---} \end{array} \begin{array}{c} m \\ \diagup \\ \bullet \\ \diagdown \\ j \end{array} = \sum_l \begin{array}{c} i \\ \diagdown \\ \bullet \\ \diagup \\ k \end{array} \begin{array}{c} \text{---} l \text{---} \\ \text{---} l \text{---} \\ \bullet \\ \diagup \\ j \end{array}$$

$$\mathcal{A}_i(0,0)\mathcal{A}_j(1,1)\mathcal{A}_k(z,\bar{z}),$$

$$L_m|\mathcal{A}\rangle\cong\oint\frac{dz}{2\pi i}z^{m+1}T(z)\mathcal{A}(0,0)\cong L_m\cdot\mathcal{A}(0,0)$$

$$T(z)\mathcal{A}(0,0)=\sum_{m=-\infty}^{\infty}z^{-m-2}L_m\cdot\mathcal{A}(0,0).$$

$$\delta\mathcal{A}(z,\bar{z})=-\epsilon\sum_{n=0}^{\infty}\frac{1}{n!}\big[\partial^n v(z)L_{n-1}+(\partial^n v(z))^*\tilde{L}_{n-1}\big]\cdot\mathcal{A}(z,\bar{z})$$

$$\begin{array}{l} L_{-1}\cdot\mathcal{A}=\partial\mathcal{A},\tilde{L}_{-1}\cdot\mathcal{A}=\bar{\partial}\mathcal{A},\\ L_0\cdot\mathcal{A}=h\mathcal{A},\tilde{L}_0\cdot\mathcal{A}=\tilde{h}\mathcal{A}. \end{array}$$

$$\begin{array}{l} L_0|\mathcal{O}\rangle=h|\mathcal{O}\rangle,\tilde{L}_0|\mathcal{O}\rangle=\tilde{h}|\mathcal{O}\rangle\\ L_m|\mathcal{O}\rangle=\tilde{L}_m|\mathcal{O}\rangle=0,m>0 \end{array}$$

$$L_m|1\rangle=\tilde{L}_m|1\rangle=0,m\geq-1$$

$$SL(2,\mathbf{C}).$$

$$L_m^\dagger=L_{-m},\tilde{L}_m^\dagger=\tilde{L}_{-m}$$

$$\langle 0;k\mid 0;k'\rangle=2\pi\delta(k-k')$$

$$\alpha_m^\dagger=\alpha_{-m},\tilde{\alpha}_m^\dagger=\tilde{\alpha}_{-m}$$

$$2h_{\mathcal{O}}\langle \mathcal{O} \mid \mathcal{O} \rangle = 2\langle \mathcal{O} | L_0 | \mathcal{O} \rangle = \langle \mathcal{O} | [L_1, L_{-1}] | \mathcal{O} \rangle = \| L_{-1} | \mathcal{O} \rangle \|^2 \geq 0$$

$$L_{-1}\cdot\mathcal{O}=\tilde{L}_{-1}\cdot\mathcal{O}=0$$

$$\partial\mathcal{A}=0\Leftrightarrow h=0,\bar{\partial}\mathcal{A}=0\Leftrightarrow\tilde{h}=0.$$

$$E=-\frac{c+\tilde{c}}{24}$$

$$E=-\frac{\pi(c+\tilde{c})}{12\ell}$$

$$\sum_{n=1}^{\infty}(n-\theta)=\frac{1}{24}-\frac{1}{8}(2\theta-1)^2$$

$$\partial\bar{\partial}\ln\,|z|^2=\partial\frac{1}{\bar{z}}=\bar{\partial}\frac{1}{z}=2\pi\delta^2(z,\bar{z})$$

$$\left\langle \prod_{i=1}^n:e^{ik_i\cdot X(z_i,\bar{z}_i)}\right\rangle=iC^X(2\pi)^D\delta^D\left(\sum_{i=1}^nk_i\right)\prod_{\substack{i,j=1\\i<j}}^n|z_{ij}|^{\alpha'k_i\cdot k_j}$$

$$S[\phi]=\int\,d^d\sigma\mathcal{L}(\phi(\sigma),\partial_a\phi(\sigma))$$

$$\delta \mathcal{L} = \epsilon \partial_a \mathcal{K}^a$$

$$j^a=2\pi i\bigg(\frac{\partial\mathcal{L}}{\partial\phi_{\alpha,a}}\epsilon^{-1}\delta\phi_{\alpha}-\mathcal{K}^a\bigg).$$

$$\begin{array}{l} \star e^{ik_1\cdot X(y_1)\star\star}e^{ik_2\cdot X(y_2)\star},y_1\rightarrow y_2(y)\\ :e^{ik\cdot X(z,\bar{z})}:,\mathrm{Im}(z)\rightarrow 0 \end{array}$$

$$L_m(L_{-m}|0;0))-L_{-m}(L_m|0;0))$$

$$:b(z)c(z'):-\circ b(z)c(z')\circ=\frac{(z/z')^{1-\lambda}-1}{z-z'}$$

$$\exp\left(iS_{\mathrm{cl}}/\hbar\right)$$

$$X_{\mathrm{upper}}^{\mu}\left(\sigma\right)=X_{\mathrm{lower}}^{\mu}\left(\pi-\sigma\right)$$

$$\int\,[dXdg]\mathrm{exp}\left(-S\right)$$

$$S=S_X+\lambda\chi$$

$$S_X\,=\,\frac{1}{4\pi\alpha'}\int_M\,d^2\sigma g^{1/2}g^{ab}\partial_aX^\mu\partial_bX_\mu$$

$$\chi=\frac{1}{4\pi}\int_Md^2\sigma g^{1/2}R+\frac{1}{2\pi}\int_{\partial M}dsk$$

$$\int\,[d\eta dX]\mathrm{exp}\left[\frac{i}{2}\int\,d\tau\big(\eta^{-1}\dot{X}^\mu\dot{X}_\mu-\eta m^2\big)\right]$$

$$\eta(\tau)\rightarrow e^{-i\theta}\eta(\tau),X^0(\tau)\rightarrow e^{-i\theta}X^0(\tau)$$

$$\int\,[dedX]\mathrm{exp}\left[-\frac{1}{2}\int\,d\tau\left(e^{-1}\sum_{\mu=1}^D\,\dot{X}^\mu\dot{X}^\mu+em^2\right)\right]$$

$$g_0^2\sim g_c\sim e^\lambda$$

$$\chi=\tilde{\chi}+\frac{1}{4}n_{\rm c},$$

$$\exp\left(-\lambda\tilde{\chi}\right)=\exp\left(-\lambda\chi+\lambda n_{\rm c}/4\right)$$

$$\int\,\frac{[dXdg]}{V_{\mathrm{diff}\times\mathrm{Weyl}}}\mathrm{exp}\left(-S\right)\equiv Z$$

$$g_{ab}(\sigma)\rightarrow \hat{g}_{ab}(\sigma)$$

$$\hat{g}_{ab}(\sigma)=\delta_{ab}$$



$$\hat{g}_{ab}(\sigma) = \exp [2\omega(\sigma)]\delta_{ab}$$

$$g'^{1/2}R' = g^{1/2}(R - 2\nabla^2\omega).$$

$$R_{abcd}=\frac{1}{2}(g_{ac}g_{bd}-g_{ad}g_{bc})R$$

$$z'\equiv\sigma'^1+i\sigma'^2=f(z)$$

$$ds'^2=\exp{(2\omega)}|\partial_zf|^{-2}dz'd\bar{z}'$$

$$\omega=\ln|\partial_zf|$$

$$\zeta\colon g\rightarrow g^\zeta, g^\zeta_{ab}(\sigma')=\exp{[2\omega(\sigma)]}\frac{\partial\sigma^c}{\partial\sigma'^a}\frac{\partial\sigma^d}{\partial\sigma'^b}g_{cd}(\sigma)$$

$$1=\Delta_{\rm FP}(g)\int\,[d\zeta]\delta(g-\hat{g}^\zeta)$$

$$Z[\hat{g}]=\int\frac{[d\zeta dXdg]}{V_{\rm diff\times Weyl}}\Delta_{\rm FP}(g)\delta(g-\hat{g}^\zeta)\exp{(-S[X,g])}.$$

$$Z[\hat{g}]=\int\frac{[d\zeta dX^\zeta]}{V_{\rm diff\times Weyl}}\Delta_{\rm FP}(\hat{g}^\zeta)\exp{(-S[X^\zeta,\hat{g}^\zeta])}$$

$$Z[\hat{g}]=\int\frac{[d\zeta dX]}{V_{\rm diff\times Weyl}}\Delta_{\rm FP}(\hat{g})\exp{(-S[X,\hat{g}])}$$

$$\begin{aligned}\Delta_{\rm FP}(g^\zeta)^{-1}&=\int\,[d\zeta']\delta(g^\zeta-\hat{g}^{\zeta'})=\int\,[d\zeta']\delta(g-\hat{g}^{\zeta^{-1},\zeta'})\\&=\int\,[d\zeta'']\delta(g-\hat{g}^{\zeta''})=\Delta_{\rm FP}(g)^{-1},\end{aligned}$$

$$Z[\hat{g}]=\int\,[dX]\Delta_{\rm FP}(\hat{g})\exp{(-S[X,\hat{g}])}$$

$$\delta g_{ab}=2\delta\omega g_{ab}-\nabla_a\delta\sigma_b-\nabla_b\delta\sigma_a=(2\delta\omega-\nabla_c\delta\sigma^c)g_{ab}-2(P_1\delta\sigma)_{ab}$$

$$(P_1\delta\sigma)_{ab}=\frac{1}{2}(\nabla_a\delta\sigma_b+\nabla_b\delta\sigma_a-g_{ab}\nabla_c\delta\sigma^c)$$

$$\begin{aligned}\Delta_{\rm FP}(\hat{g})^{-1}&=\int\,[d\delta\omega d\delta\sigma]\delta\bigl[-(2\delta\omega-\hat{\nabla}\cdot\delta\sigma)\hat{g}+2\hat{P}_1\delta\sigma\bigr] \\&=\int\,[d\delta\omega d\beta d\delta\sigma]\exp\left\{2\pi i\int\,d^2\sigma\hat{g}^{1/2}\beta^{ab}\bigl[-(2\delta\omega-\hat{\nabla}\cdot\delta\sigma)\hat{g}+2\hat{P}_1\delta\sigma\bigr]_{ab}\right\} \\&=\int\,[d\beta'd\delta\sigma]\exp\left\{4\pi i\int\,d^2\sigma\hat{g}^{1/2}\beta'^{ab}(\hat{P}_1\delta\sigma)_{ab}\right\}\end{aligned}$$

$$\begin{array}{l} \delta\sigma^a \rightarrow c^a \\ \beta'_{ab} \rightarrow b_{ab} \end{array}$$

$$\Delta_{\text{FP}}(\hat{g}) = \int [dbdc] \exp(-S_g)$$

$$S_{\rm g} = \frac{1}{2\pi} \int d^2\sigma \hat{g}^{1/2} b_{ab} \hat{\nabla}^a c^b = \frac{1}{2\pi} \int d^2\sigma \hat{g}^{1/2} b_{ab} (\hat{P}_1 c)^{ab}$$

$$Z[\hat{g}] = \int [dXd bdc] \exp \left(-S_X - S_g \right)$$

$$Z[\hat{g}] = (\det \hat{\nabla}^2)^{-D/2} \det \hat{P}_1,$$

$$S_{\rm g} = \frac{1}{2\pi} \int d^2z (b_{zz} \nabla_{\bar{z}} c^z + b_{\bar{z}\bar{z}} \nabla_z c^{\bar{z}}) = \frac{1}{2\pi} \int d^2z (b_{zz} \partial_{\bar{z}} c^z + b_{\bar{z}\bar{z}} \partial_z c^{\bar{z}})$$

$$n_a\delta\sigma^a=0$$

$$n_a c^a = 0$$

$$\int_{\partial M} ds n^a b_{ab} \delta c^b = 0$$

$$n_a t_b b^{ab} = 0$$

$$Z[g^\zeta]=Z[g]$$

$$\langle \ldots \rangle_g \equiv \int [dXd bdc] \exp \left(-S[X,b,c,g] \right) \ldots$$

$$\langle \ldots \rangle_{g^\zeta} = \langle \ldots \rangle_g$$

$$\delta \langle \ldots \rangle_g = -\frac{1}{4\pi} \int d^2\sigma g(\sigma)^{1/2} \delta g_{ab}(\sigma) \langle T^{ab}(\sigma) \ldots \rangle_g$$

$$T^{ab}(\sigma) \stackrel{\text{classical}}{=} \frac{4\pi}{g(\sigma)^{1/2}} \frac{\delta}{\delta g_{ab}(\sigma)} S$$

$$\delta_{\rm W} \langle \ldots \rangle_g = -\frac{1}{2\pi} \int d^2\sigma g(\sigma)^{1/2} \delta \omega(\sigma) \langle T^a_a(\sigma) \ldots \rangle_g$$

$$T^a_a\stackrel{?}{=}0$$

$$T^a_a=a_1R$$

$$T_{z\bar{z}}=\frac{a_1}{2}g_{z\bar{z}}R$$

$$\nabla^{\bar{z}}T_{\bar{z}z}=\frac{a_1}{2}\nabla^{\bar{z}}(g_{z\bar{z}}R)=\frac{a_1}{2}\partial_zR,$$



$$\nabla^z T_{zz} = -\nabla^{\bar{z}} T_{\bar{z}\bar{z}} = -\frac{a_1}{2} \partial_z R$$

$$a_1 \partial_z \nabla^2 \delta \omega \approx 4 a_1 \partial_z^2 \partial_{\bar{z}} \delta \omega,$$

$$\epsilon^{-1} \delta T_{zz}(z) = -\frac{c}{12} \partial_z^3 v^z(z) - 2 \partial_z v^z(z) T_{zz}(z) - v^z(z) \partial_z T_{zz}(z)$$

$$\delta_{\rm W} T_{zz} = -\frac{c}{6} \partial_z^2 \delta \omega$$

$$c=-12a_1, T_a^a=-\frac{c}{12}R$$

$$R \, = \, -2 \exp \left(-2 \omega \right) \partial_a \partial_a \omega \\ \nabla^2 \, = \, \exp \left(-2 \omega \right) \partial_a \partial_a$$

$$\delta_{\rm W} Z[\exp{(2\omega)}\delta] = \frac{a_1}{\pi} Z[\exp{(2\omega)}\delta] \int d^2\sigma \delta\omega \partial_a \partial_a \omega$$

$$Z[\exp{(2\omega)}\delta] = Z[\delta] \exp\left(-\frac{a_1}{2\pi} \int d^2\sigma \partial_a \omega \partial_a \omega\right)$$

$$Z[g] = Z[\delta] \exp \left[\frac{a_1}{8\pi} \int d^2\sigma \int d^2\sigma' g^{1/2} R(\sigma) G(\sigma, \sigma') g^{1/2} R(\sigma') \right]$$

$$g(\sigma)^{1/2}\nabla^2G(\sigma,\sigma')=\delta^2(\sigma-\sigma')$$

$$\ln \frac{Z[\delta+h]}{Z[\delta]} \approx \frac{a_1}{8\pi^2} \int d^2z \int d^2z' (\partial_{\bar{z}}^2 \ln |z-z'|^2) h_{\bar{z}\bar{z}}(z,\bar{z}) \partial_{\bar{z}}^2 h_{\bar{z}\bar{z}}(z',\bar{z}')$$

$$=-\frac{3a_1}{4\pi^2} \int d^2z \int d^2z' \frac{h_{\bar{z}\bar{z}}(z,\bar{z}) h_{\bar{z}\bar{z}}(z',\bar{z}')}{(z-z')^4}$$

$$\ln \frac{Z[\delta+h]}{Z[\delta]} \approx \frac{1}{8\pi^2} \int d^2z \int d^2z' h_{\bar{z}\bar{z}}(z,\bar{z}) h_{\bar{z}\bar{z}}(z',\bar{z}') \langle T_{zz}(z) T_{zz}(z') \rangle_{\delta}$$

$$c=c^X+c^{\mathfrak{g}}=D-26$$

$$T^a_a=a_1R+a_2$$

$$S_{\rm ct} = b \int d^2\sigma g^{1/2}$$

$$T^a_a=a_1R+(a_2+4\pi b)$$

$$\delta_{\rm W} \ln \langle \ldots \rangle_g = -\frac{1}{2\pi} \int_M d^2\sigma g^{1/2} (a_1 R + a_2) \delta \omega - \frac{1}{2\pi} \int_{\partial M} ds (a_3 + a_4 k + a_5 n^a \partial_a) \delta \omega$$

$$S_{\rm ct} = \int_M d^2\sigma g^{1/2} b_1 + \int_{\partial M} ds (b_2 + b_3 k)$$



$$\delta_{\mathrm{W}}S_{\mathrm{ct}}=2\int_Md^2\sigma g^{1/2}b_1\delta\omega+\int_{\partial M}ds(b_2+b_3n^a\partial_a)\delta\omega$$

$$\begin{aligned}\delta_{\mathrm{W}_1}(\delta_{\mathrm{W}_2}\ln\langle...\rangle_g)&=\frac{a_1}{\pi}\int_Md^2\sigma g^{1/2}\delta\omega_2\nabla^2\delta\omega_1-\frac{a_4}{2\pi}\int_{\partial M}ds\delta\omega_2n^a\partial_a\delta\omega_1\\&=-\frac{a_1}{\pi}\int_Md^2\sigma g^{1/2}\partial_a\delta\omega_2\partial^a\delta\omega_1+\frac{2a_1-a_4}{2\pi}\int_{\partial M}dsn^a\delta\omega_2\partial_a\delta\omega_1\end{aligned}$$

$$T_a^a(\sigma)=-\frac{\mathcal{C}(\sigma)}{12}R(\sigma)$$

$$-2\pi t\leq \mathrm{Im} w\leq 0, w\cong w+2\pi$$

$$z=\exp\left(-i w\right), \exp\left(-2\pi t\right)\leq |z|\leq 1$$

$$-2\pi t\leq \mathrm{Im} w\leq 0, 0\leq \mathrm{Re} w\leq \pi$$

$$\exp\left(-2\pi t\right)\leq |z|\leq 1, \mathrm{Im} z\geq 0$$

$$S_{j_1\dots j_n}(k_1,\dots,k_n)=\sum_{\substack{\text{compact} \\ \text{topologies}}}\int\frac{[dXdg]}{V_{\text{diff}\times\text{Weyl}}}\exp\left(-S_X-\lambda\chi\right)\prod_{i=1}^n\int d^2\sigma_i g(\sigma_i)^{1/2}\mathcal{V}_{j_i}(k_i,\sigma_i)$$

$$V_0=2g_{\rm c}\int d^2\sigma g^{1/2}e^{ik\cdot X}\rightarrow g_{\rm c}\int d^2z\colon e^{ik\cdot X}\colon$$

$$m^2=-k^2=-\frac{4}{\alpha'}$$

$$\frac{2g_c}{\alpha'}\int d^2z\colon\partial X^\mu\bar\partial X^{\nu}e^{ik\cdot X}\colon.$$

$$h=\tilde{h}=1+\frac{\alpha'k^2}{4}$$

$$[\mathcal{F}]_{\mathrm{r}}=\exp\left(\frac{1}{2}\int d^2\sigma d^2\sigma'\Delta(\sigma,\sigma')\frac{\delta}{\delta X^\mu(\sigma)}\frac{\delta}{\delta X_\mu(\sigma')}\right)\mathcal{F}$$

$$\Delta(\sigma,\sigma')=\frac{\alpha'}{2}\ln\,d^2(\sigma,\sigma')$$

$$\delta_{\mathrm{W}}[\mathcal{F}]_{\mathrm{r}}=[\delta_{\mathrm{W}}\mathcal{F}]_{\mathrm{r}}+\frac{1}{2}\int d^2\sigma d^2\sigma'\delta_{\mathrm{W}}\Delta(\sigma,\sigma')\frac{\delta}{\delta X^\mu(\sigma)}\frac{\delta}{\delta X_\mu(\sigma')}[\mathcal{F}]_{\mathrm{r}}$$

$$\delta_{\mathrm{W}}V_0=2g_{\rm c}\int d^2\sigma g^{1/2}\left(2\delta\omega(\sigma)-\frac{k^2}{2}\delta_{\mathrm{W}}\Delta(\sigma,\sigma)\right)[e^{ik\cdot X(\sigma)}]_{\mathrm{r}}$$

$$d^2(\sigma,\sigma')\approx (\sigma-\sigma')^2\exp\left(2\omega(\sigma)\right)$$



$$\Delta(\sigma, \sigma') \approx \alpha' \omega(\sigma) + \frac{\alpha'}{2} \ln(\sigma - \sigma')^2$$

$$\delta_W \Delta(\sigma, \sigma) = \alpha' \delta \omega(\sigma)$$

$$\delta^D(X(\sigma) - x_0) = \int \frac{d^D k}{(2\pi)^D} \exp[ik \cdot (X(\sigma) - x_0)]$$

$$\int_M d^2\sigma g(\sigma)^{1/2} \int_M d^2\sigma' g(\sigma')^{1/2} \delta^D(X(\sigma) - X(\sigma'))$$

$$V_1 = \frac{g_c}{\alpha'} \int d^2\sigma g^{1/2} \left\{ (g^{ab} s_{\mu\nu} + i\epsilon^{ab} a_{\mu\nu}) [\partial_a X^\mu \partial_b X^\nu e^{ik \cdot X}]_r + \alpha' \phi R[e^{ik \cdot X}]_r \right\}$$

$$\partial_a \delta_W \Delta(\sigma, \sigma')|_{\sigma'=\sigma} = \frac{1}{2} \alpha' \partial_a \delta \omega(\sigma)$$

$$\partial_a \partial'_b \delta_W \Delta(\sigma, \sigma')|_{\sigma'=\sigma} = \frac{1+\gamma}{2} \alpha' \nabla_a \partial_b \delta \omega(\sigma)$$

$$\nabla_a \partial_b \delta_W \Delta(\sigma, \sigma')|_{\sigma'=\sigma} = -\frac{\gamma}{2} \alpha' \nabla_a \partial_b \delta \omega(\sigma)$$

$$\delta_W V_1 = \frac{g_c}{2} \int d^2\sigma g^{1/2} \delta \omega \left\{ (g^{ab} S_{\mu\nu} + i\epsilon^{ab} A_{\mu\nu}) [\partial_a X^\mu \partial_b X^\nu e^{ik \cdot X}]_r + \alpha' F R[e^{ik \cdot X}]_r \right\}$$

$$S_{\mu\nu} = -k^2 s_{\mu\nu} + k_\nu k^\omega s_{\mu\omega} + k_\mu k^\omega s_{\nu\omega} - (1+\gamma) k_\mu k_\nu s^\omega_\omega + 4k_\mu k_\nu \phi$$

$$A_{\mu\nu} = -k^2 a_{\mu\nu} + k_\nu k^\omega a_{\mu\omega} - k_\mu k^\omega a_{\nu\omega}$$

$$F = (\gamma-1)k^2\phi + \frac{1}{2}\gamma k^\mu k^\nu s_{\mu\nu} - \frac{1}{4}\gamma(1+\gamma)k^2 s^v_v$$

$$[\nabla^2 X^\mu e^{ik \cdot X}]_r = i \frac{\alpha' \gamma}{4} k^\mu R[e^{ik \cdot X}]_r$$

$$S_{\mu\nu} = A_{\mu\nu} = F = 0$$

$$s_{\mu\nu} \rightarrow s_{\mu\nu} + \xi_\mu k_\nu + k_\mu \xi_\nu$$

$$a_{\mu\nu} \rightarrow a_{\mu\nu} + \zeta_\mu k_\nu - k_\mu \zeta_\nu$$

$$\phi \rightarrow \phi + \frac{\gamma}{2} k \cdot \xi$$

$$n^\mu s_{\mu\nu} = n^\mu a_{\mu\nu} = 0$$

$$k^2=0$$

$$k^\nu s_{\mu\nu} = k^\mu a_{\mu\nu} = 0$$

$$\phi = \frac{1+\gamma}{4} s^\mu_\mu$$

$$k^\mu = (1,1,0,0,\ldots,0), n^\mu = \frac{1}{2}(-1,1,0,0,\ldots,0)$$



$$\begin{aligned}
[e^{ik \cdot X}]_{\text{DR}} &= [e^{ik \cdot X}]_{\text{r}} \\
[\partial_a X^\mu e^{ik \cdot X}]_{\text{DR}} &= [\partial_a X^\mu e^{ik \cdot X}]_{\text{r}} \\
[\partial_a X^\mu \partial_b X^\nu e^{ik \cdot X}]_{\text{DR}} &= [\partial_a X^\mu \partial_b X^\nu e^{ik \cdot X}]_{\text{r}} - \frac{\alpha'}{12} g_{ab} \eta^{\mu\nu} R [e^{ik \cdot X}]_{\text{r}} \\
[\nabla_a \partial_b X^\mu e^{ik \cdot X}]_{\text{DR}} &= [\nabla_a \partial_b X^\mu e^{ik \cdot X}]_{\text{r}} + i \frac{\alpha'}{12} g_{ab} k^\mu R [e^{ik \cdot X}]_{\text{r}}
\end{aligned}$$

$$g_o \int_{\partial M} ds [e^{ik \cdot X}]_{\text{r}}$$

$$-i \frac{g_o}{(2\alpha')^{1/2}} e_\mu \int_{\partial M} ds [\dot{X}^\mu e^{ik \cdot X}]_{\text{r}}$$

$$S_{\text{pp}} = \frac{1}{2} \int d\tau (\eta^{-1} G_{\mu\nu}(X) \dot{X}^\mu \dot{X}^\nu - \eta m^2)$$

$$S_\sigma = \frac{1}{4\pi\alpha'} \int_M d^2\sigma g^{1/2} g^{ab} G_{\mu\nu}(X) \partial_a X^\mu \partial_b X^\nu$$

$$G_{\mu\nu}(X) = \eta_{\mu\nu} + \chi_{\mu\nu}(X)$$

$$\exp(-S_\sigma) = \exp(-S_{\text{P}}) \left(1 - \frac{1}{4\pi\alpha'} \int_M d^2\sigma g^{1/2} g^{ab} \chi_{\mu\nu}(X) \partial_a X^\mu \partial_b X^\nu + \dots \right)$$

$$\chi_{\mu\nu}(X) = -4\pi g_c e^{ik \cdot X} S_{\mu\nu}.$$

$$S_\sigma = \frac{1}{4\pi\alpha'} \int_M d^2\sigma g^{1/2} [(g^{ab} G_{\mu\nu}(X) + i\epsilon^{ab} B_{\mu\nu}(X)) \partial_a X^\mu \partial_b X^\nu + \alpha' R \Phi(X)]$$

$$\delta B_{\mu\nu}(X) = \partial_\mu \zeta_\nu(X) - \partial_\nu \zeta_\mu(X)$$

$$H_{\omega\mu\nu} = \partial_\omega B_{\mu\nu} + \partial_\mu B_{\nu\omega} + \partial_\nu B_{\omega\mu}$$

$$G_{\mu\nu}(X) \partial_a X^\mu \partial_b X^\nu = \left[G_{\mu\nu}(x_0) + G_{\mu\nu,\omega}(x_0) Y^\omega + \frac{1}{2} G_{\mu\nu,\omega\rho}(x_0) Y^\omega Y^\rho + \dots \right] \partial_a Y^\mu \partial_b Y^\nu$$

$$\begin{aligned}
G_{\mu\nu}(X) &= \eta_{\mu\nu} - 4\pi g_c S_{\mu\nu} e^{ik \cdot X} \\
B_{\mu\nu}(X) &= -4\pi g_c a_{\mu\nu} e^{ik \cdot X} \\
\Phi(X) &= -4\pi g_c \phi e^{ik \cdot X}
\end{aligned}$$

$$T^a_\alpha = -\frac{1}{2\alpha'} \beta^G_{\mu\nu} g^{ab} \partial_a X^\mu \partial_b X^\nu - \frac{i}{2\alpha'} \beta^B_{\mu\nu} \epsilon^{ab} \partial_a X^\mu \partial_b X^\nu - \frac{1}{2} \beta^\Phi R$$

$$\beta_{\mu\nu}^G \approx -\frac{\alpha'}{2}(\partial^2\chi_{\mu\nu} - \partial_\nu\partial^\omega\chi_{\mu\omega} - \partial_\mu\partial^\omega\chi_{\omega\nu} + \partial_\mu\partial_\nu\chi_\omega^\omega) + 2\alpha'\partial_\mu\partial_\nu\Phi$$

$$\begin{aligned}\beta_{\mu\nu}^B &\approx -\frac{\alpha'}{2}\partial^\omega H_{\omega\mu\nu} \\ \beta^\Phi &\approx \frac{D-26}{6} - \frac{\alpha'}{2}\partial^2\Phi\end{aligned}$$

$$\begin{aligned}\beta_{\mu\nu}^G &= \alpha' R_{\mu\nu} + 2\alpha'\nabla_\mu\nabla_\nu\Phi - \frac{\alpha'}{4}H_{\mu\lambda\omega}H_{\nu}{}^{\lambda\omega} + O(\alpha'^2) \\ \beta_{\mu\nu}^B &= -\frac{\alpha'}{2}\nabla^\omega H_{\omega\mu\nu} + \alpha'\nabla^\omega\Phi H_{\omega\mu\nu} + O(\alpha'^2) \\ \beta^\Phi &= \frac{D-26}{6} - \frac{\alpha'}{2}\nabla^2\Phi + \alpha'\nabla_\omega\Phi\nabla^\omega\Phi - \frac{\alpha'}{24}H_{\mu\nu\lambda}H^{\mu\nu\lambda} + O(\alpha'^2)\end{aligned}$$

$$\beta_{\mu\nu}^G = \beta_{\mu\nu}^B = \beta^\Phi = 0$$

$$G_{\mu\nu}(X) = \eta_{\mu\nu}, B_{\mu\nu}(X) = 0, \Phi(X) = \Phi_0$$

$$\lambda = \Phi_0$$

$$G_{\mu\nu}(X) = \eta_{\mu\nu}, B_{\mu\nu}(X) = 0, \Phi(X) = V_\mu X^\mu$$

$$V_\mu V^\mu = \frac{26-D}{6\alpha'}.$$

$$S = \frac{1}{2\kappa_0^2} \int d^D x (-G)^{1/2} e^{-2\Phi} \left[-\frac{2(D-26)}{3\alpha'} + R - \frac{1}{12} H_{\mu\nu\lambda} H^{\mu\nu\lambda} + 4\partial_\mu\Phi\partial^\mu\Phi + O(\alpha') \right]$$

$$\begin{aligned}\delta S = & -\frac{1}{2\kappa_0^2\alpha'} \int d^D x (-G)^{1/2} e^{-2\Phi} \left[\delta G_{\mu\nu} \beta^{G\mu\nu} \right. \\ & \left. + \delta B_{\mu\nu} \beta^{B\mu\nu} + \left(2\delta\Phi - \frac{1}{2} G^{\mu\nu} \delta G_{\mu\nu} \right) (\beta_\omega^{G\omega} - 4\beta^\Phi) \right]\end{aligned}$$

$$\tilde{G}_{\mu\nu}(x) = \exp\left(2\omega(x)\right)G_{\mu\nu}(x)$$

$$\tilde{R} = \exp\left(-2\omega\right)\left[R - 2(D-1)\nabla^2\omega - (D-2)(D-1)\partial_\mu\omega\partial^\mu\omega\right]$$

$$\tilde{\Phi} = \Phi - \Phi_0$$

$$S = \frac{1}{2\kappa^2} \int d^D X (-\tilde{G})^{1/2} \left[-\frac{2(D-26)}{3\alpha'} e^{4\tilde{\Phi}/(D-2)} + \tilde{R} - \frac{1}{12} e^{-8\tilde{\Phi}/(D-2)} H_{\mu\nu\lambda} \tilde{H}^{\mu\nu\lambda} - \frac{4}{D-2} \partial_\mu \tilde{\Phi} \partial^\mu \tilde{\Phi} + O(\alpha') \right]$$

$$\kappa = (8\pi G_{\rm N})^{1/2} = \frac{(8\pi)^{1/2}}{M_{\rm P}} = (2.43 \times 10^{18} {\rm GeV})^{-1}$$

$$\begin{aligned}S &= -\frac{1}{4}\int d^Dx {\rm Tr}(F_{\mu\nu}F^{\mu\nu}) \\ F_{\mu\nu} &= \partial_\mu A_\nu - \partial_\nu A_\mu - ig[A_\mu,A_\nu]\end{aligned}$$

$$gA_\mu = A'_\mu, gF_{\mu\nu} = F'_{\mu\nu}$$

$$g_{MN}=\left[\begin{array}{cc}\eta_{\mu\nu}&0\\0&g_{mn}(x^p)\end{array}\right]$$

$$-\frac{1}{4\pi\alpha'}\int_M d^2\sigma g^{1/2}g^{ab}\partial_aX^0\partial_bX^0$$

$$(t,t')=\int\,d^2\sigma g^{1/2}t\cdot t'$$

$$\nabla_a j^a = a R$$

$$\delta\langle f\mid i\rangle=-\frac{1}{4\pi}\int\,d^2\sigma g(\sigma)^{1/2}\delta g_{ab}(\sigma)\langle f\mid T^{ab}(\sigma)\mid i\rangle.$$

$$\langle \psi | T^{ab}(\sigma) | \psi' \rangle = 0$$

$$T_{ab}=T_{ab}^X+T_{ab}^{\mathcal{G}}$$

$$T_{ab}=T_{ab}^{\mathfrak{m}}+T_{ab}^{\mathfrak{g}}$$

$$(L_n^{\mathfrak{m}}+A\delta_{n,0})|\psi\rangle=0\,\,\text{for}\,\,n\geq 0$$

$$\langle \psi | L_n^{\mathfrak{m}} | \psi' \rangle = \langle L_{-n}^{\mathfrak{m}} \psi \mid \psi' \rangle = 0$$

$$L_n^{\mathfrak{m}\dagger}=L_{-n}^{\mathfrak{m}}$$

$$|\chi\rangle=\sum_{n=1}^\infty L_{-n}^{\mathfrak{m}}|\chi_n\rangle$$

$$|\psi\rangle\cong|\psi\rangle+|\chi\rangle.$$

$$\mathcal{H}_{\text{OCQ}}=\frac{\mathcal{H}_{\text{phys}}}{\mathcal{H}_{\text{null}}}.$$

$$\begin{aligned} L_0^{\mathfrak{m}} &= \alpha' p^2 + \alpha_{-1} \cdot \alpha_1 + \cdots \\ L_{\pm 1}^{\mathfrak{m}} &= (2\alpha')^{1/2} p \cdot \alpha_{\pm 1} + \cdots \end{aligned}$$

$$|e;k\rangle=e\cdot\alpha_{-1}|0;k\rangle.$$

$$\langle e;k\mid e;k'\rangle=\langle 0;k|e^*\cdot\alpha_1e\cdot\alpha_{-1}|0;k'\rangle=\langle 0;k|(e^*\cdot e+e^*\cdot\alpha_{-1}e\cdot\alpha_1)|0;k'\rangle$$

$$=e^{\mu*}e_\mu(2\pi)^D\delta^D(k-k')$$

$$\alpha_n^{\mu\dagger}=\alpha_{-n}^\mu,$$

$$\langle 0;k\mid 0;k'\rangle=(2\pi)^D\delta^D(k-k'),$$

$$m^2=\frac{1+A}{\alpha'}$$



$$L_1^m|e;k\rangle \propto p \cdot \alpha_1 e \cdot \alpha_{-1}|0;k\rangle = e \cdot k|0;k\rangle = 0$$

$$L_{-1}^m|0;k\rangle = (2\alpha')^{1/2}k \cdot \alpha_{-1}|0;k\rangle.$$

$$k \cdot e = 0, e_\mu \cong e_\mu + \gamma k_\mu$$

$$\mathcal{H}_{\text{OCQ}} = \mathcal{H}_{\text{light-cone}} ,$$

$$|0;k\rangle, m^2 = -\frac{4}{\alpha'}$$

$$e_{\mu\nu}\alpha_{-1}^\mu\tilde{\alpha}_{-1}^\nu|0;k\rangle, m^2=0, k^\mu e_{\mu\nu}=k^\nu e_{\mu\nu}=0$$

$$e_{\mu\nu}\cong e_{\mu\nu}+a_\mu k_\nu+k_\mu b_\nu, a\cdot k=b\cdot k=0$$

$$(L_0^{\rm m}+L_0^{\rm g})|\psi,\downarrow\rangle=0$$

$$(H^{\rm m}+H^{\rm g})|\psi,\downarrow\rangle=0$$

Cuantización BRST:

$$[\delta_\alpha,\delta_\beta]=f^\gamma{}_{\alpha\beta}\delta_\gamma$$

$$F^A(\phi)=0$$

$$\int \frac{[d\phi_i]}{V_{\rm gauge}} \exp{(-S_1)} \rightarrow \int [d\phi_i dB_A db_A dc^\alpha] \exp{(-S_1-S_2-S_3)}$$

$$S_2=-iB_AF^A(\phi)$$

$$S_3=b_Ac^\alpha\delta_\alpha F^A(\phi)$$

$$\delta_{\rm B}\phi_i=-i\epsilon c^\alpha\delta_\alpha\phi_i$$

$$\delta_{\rm B}B_A=0$$

$$\delta_{\rm B}b_A=\epsilon B_A$$

$$\delta_{\rm B}c^\alpha=\frac{i}{2}\epsilon f^\alpha_{\beta\gamma}c^\beta c^\gamma$$

$$\delta_{\rm B}(b_AF^A)=i\epsilon(S_2+S_3).$$

$$\epsilon\delta\langle f\mid i\rangle=i\langle f\mid\delta_{\rm B}(b_A\delta F^A)\mid i\rangle=-\epsilon\langle f\mid\{Q_{\rm B},b_A\delta F^A\}\mid i\rangle$$

$$\langle\psi|\{Q_{\rm B},b_A\delta F^A\}|\psi'\rangle=0$$

$$Q_{\rm B}|\psi\rangle=Q_{\rm B}|\psi'\rangle=0$$

$$\epsilon^{-1}\delta_{\rm B}(b_AB_B M^{AB})=-B_AB_B M^{AB}$$

$$0=[Q_{\rm B},\{Q_{\rm B},b_A\delta F^A\}]=Q_{\rm B}^2b_A\delta F^A-Q_{\rm B}b_A\delta F^AQ_{\rm B}+Q_{\rm B}b_A\delta F^AQ_{\rm B}-b_A\delta F^AQ_{\rm B}^2=[Q_{\rm B}^2,b_A\delta F^A]$$

$$Q_{\rm B}^2=0$$



$$\delta_{\mathrm{B}}(\delta'_{\mathrm{B}}c^{\alpha})=-\frac{1}{2}\epsilon\epsilon'f_{\beta\gamma}^{\alpha}f_{\delta\epsilon}^{\gamma}c^{\beta}c^{\delta}c^{\epsilon}=0$$

$$Q_{\mathrm{B}}|\chi\rangle$$

$$\langle\psi|(Q_{\mathrm{B}}|\chi\rangle)=(\langle\psi|Q_{\mathrm{B}})|\chi\rangle=0$$

$$|\psi'\rangle=|\psi\rangle+Q_{\mathrm{B}}|\chi\rangle$$

$$\mathcal{H}_{\mathrm{BRST}}=\frac{\mathcal{H}_{\mathrm{closed}}}{\mathcal{H}_{\mathrm{exact}}}$$

$$\delta_{\tau_1}X^\mu(\tau)=-\delta(\tau-\tau_1)\partial_\tau X^\mu(\tau),\delta_{\tau_1}e(\tau)=-\partial_\tau[\delta(\tau-\tau_1)e(\tau)]$$

$$[\delta_{\tau_1},\delta_{\tau_2}]X^\mu(\tau)=-[\delta(\tau-\tau_1)\partial_\tau\delta(\tau-\tau_2)-\delta(\tau-\tau_2)\partial_\tau\delta(\tau-\tau_1)]\partial_\tau X^\mu(\tau)$$

$$\equiv \int \, d\tau_3 f^{\tau_3}{}_{\tau_1\tau_2} \delta_{\tau_3} X^\mu(\tau)$$

$$f^{\tau_3}{}_{\tau_1\tau_2}=\delta(\tau_3-\tau_1)\partial_{\tau_3}\delta(\tau_3-\tau_2)-\delta(\tau_3-\tau_2)\partial_{\tau_3}\delta(\tau_3-\tau_1)$$

$$\delta_{\mathrm{B}}X^\mu=i\epsilon c\dot{X}^\mu$$

$$\delta_{\mathrm{B}}e=i\epsilon(\dot{c}e)$$

$$\delta_{\mathrm{B}}B=0$$

$$\delta_{\mathrm{B}}b=\epsilon B$$

$$\delta_{\mathrm{B}}c=i\epsilon c\dot{c}$$

$$S=\int\,d\tau\left(\frac{1}{2}e^{-1}\dot{X}^\mu\dot{X}_\mu+\frac{1}{2}em^2+iB(e-1)-e\dot{b}c\right)$$

$$S=\int\,d\tau\left(\frac{1}{2}\dot{X}^\mu\dot{X}_\mu+\frac{1}{2}m^2-\dot{b}c\right)$$

$$\delta_{\mathrm{B}}X^\mu=i\epsilon c\dot{X}^\mu$$

$$\delta_{\mathrm{B}}b=i\epsilon\left(-\frac{1}{2}\dot{X}^\mu\dot{X}_\mu+\frac{1}{2}m^2-\dot{b}c\right)$$

$$\delta_{\mathrm{B}}c=i\epsilon c\dot{c}$$

$$[p^\mu,X^\nu]=-i\eta^{\mu\nu},\{b,c\}=1$$

$$Q_{\mathrm{B}}=cH$$

$$p^\mu|k,\downarrow\rangle=k^\mu|k,\downarrow\rangle,p^\mu|k,\uparrow\rangle=k^\mu|k,\uparrow\rangle$$

$$b|k,\downarrow\rangle=0,b|k,\uparrow\rangle=|k,\downarrow\rangle$$

$$c|k,\downarrow\rangle=|k,\uparrow\rangle,c|k,\uparrow\rangle=0$$

$$Q_{\mathrm{B}}|k,\downarrow\rangle=\frac{1}{2}(k^2+m^2)|k,\uparrow\rangle,Q_{\mathrm{B}}|k,\uparrow\rangle=0$$

$$|k,\downarrow\rangle,\quad k^2+m^2=0$$

$$|k,\uparrow\rangle,\quad \text{all } k^\mu,$$

$$|k,\uparrow\rangle,k^2+m^2\neq 0$$



$$|k,\downarrow\rangle,k^2+m^2=0;\;|k,\uparrow\rangle,k^2+m^2=0$$

$$b|\psi\rangle=0$$

$$\begin{aligned}\delta_{\rm B}X^\mu &= i\epsilon(c\partial+\tilde{c}\bar{\partial})X^\mu \\ \delta_{\rm B}b &= i\epsilon(T^X+T^{\mathfrak{g}}), \delta_{\rm B}\tilde{b} = i\epsilon(\tilde{T}^X+\tilde{T}^{\mathfrak{g}}) \\ \delta_{\rm B}c &= i\epsilon c\partial c, \delta_{\rm B}\tilde{c} = i\epsilon\tilde{c}\bar{\partial}\tilde{c}\end{aligned}$$

$$\frac{i}{4\pi}\int\,d^2\sigma g^{1/2}B^{ab}(\delta_{ab}-g_{ab})$$

$$j_{\rm B}=cT^{\rm m}+\frac{1}{2}\colon cT^{\mathfrak{g}}\colon+\frac{3}{2}\partial^2c=cT^{\rm m}+\colon bc\partial c\colon+\frac{3}{2}\partial^2c$$

$$\begin{aligned}j_{\rm B}(z)b(0) &\sim \frac{3}{z^3}+\frac{1}{z^2}j^{\mathfrak{g}}(0)+\frac{1}{z}T^{\rm m+\mathfrak{g}}(0),\\ j_{\rm B}(z)c(0) &\sim \frac{1}{z}c\partial c(0),\\ j_{\rm B}(z)\mathcal{O}^{\rm m}(0,0) &\sim \frac{h}{z^2}c\theta^{\rm m}(0,0)+\frac{1}{z}[h(\partial c)\mathcal{O}^{\rm m}(0,0)+c\partial\theta^{\rm m}(0,0)].\end{aligned}$$

$$Q_{\rm B}=\frac{1}{2\pi i}\oint\, (dzj_{\rm B}-d\bar{z}\tilde{j}_{\rm B})$$

$$\{Q_{\rm B},b_m\}=L_m^{\rm m}+L_m^{\mathfrak{g}}$$

$$Q_{\rm B}=\sum_{n=-\infty}^{\infty}\left(c_nL_{-n}^{\rm m}+\tilde{c}_n\tilde{L}_{-n}^{\rm m}\right)+\sum_{m,n=-\infty}^{\infty}\frac{(m-n)}{2}\frac{(c_m c_n b_{-m-n}+\tilde{c}_m\tilde{c}_n\tilde{b}_{-m-n})^{\circ}}{\circ}+a^{\rm B}(c_0+\tilde{c}_0)$$

$$\{Q_{\rm B},b_0\}=L_0^{\rm m}+L_0^{\mathfrak{g}}$$

$$\{Q_{\rm B},Q_{\rm B}\}=0 \text{ only if } c^{\rm m}=26$$

$$j_{\rm B}(z)j_{\rm B}(0)\sim -\frac{c^{\rm m}-18}{2z^3}c\partial c(0)-\frac{c^{\rm m}-18}{4z^2}c\partial^2c(0)-\frac{c^{\rm m}-26}{12z}c\partial^3c(0)$$

$$T(z)j_{\rm B}(0)\sim \frac{c^{\rm m}-26}{2z^4}c(0)+\frac{1}{z^2}j_{\rm B}(0)+\frac{1}{z}\partial j_{\rm B}(0),$$

$$\left[G_I,G_J\right]=ig^K{}_{IJ}G_K$$

$$\{c^I,b_J\}=\delta_J^I,\{c^I,c^J\}=\{b_I,b_J\}=0$$

$$Q_{\rm B}=c^I G_I^{\rm m}-\frac{i}{2}g_{IJ}^K c^I c^J b_K=c^I\left(G_I^{\rm m}+\frac{1}{2}G_I^{\mathfrak{g}}\right)$$

$$G^{\mathfrak{g}}_I=-ig^K{}_{IJ}c^Jb_K$$

$$Q_{\text{B}}^2 = \frac{1}{2} \{Q_{\text{B}}, Q_{\text{B}}\} = -\frac{1}{2} g_{IJ}^K g_{KL}^M c^I c^J c^L b_M = 0$$

$$\begin{aligned} (\alpha_m^\mu)^\dagger &= \alpha_{-m}^\mu, & (\tilde{\alpha}_m^\mu)^\dagger &= \tilde{\alpha}_{-m}^\mu \\ (b_m)^\dagger &= b_{-m}, & (\tilde{b}_m)^\dagger &= \tilde{b}_{-m} \\ (c_m)^\dagger &= c_{-m}, & (\tilde{c}_m)^\dagger &= \tilde{c}_{-m}. \end{aligned}$$

$$\begin{aligned} \langle 0; k | c_0 | 0; k' \rangle &= (2\pi)^{26} \delta^{26}(k - k') \\ \langle 0; k | \tilde{c}_0 c_0 | 0; k' \rangle &= i(2\pi)^{26} \delta^{26}(k - k') \end{aligned}$$

$$b_0|\psi\rangle=0$$

$$L_0|\psi\rangle=\{Q_{\text{B}},b_0\}|\psi\rangle=0$$

$$L_0=\alpha'(p^\mu p_\mu+m^2)$$

$$\alpha' m^2 = \sum_{n=1}^{\infty} n \left(N_{bn} + N_{cn} + \sum_{\mu=0}^{25} N_{\mu n} \right) - 1.$$

$$|0;{\bf k}\rangle,-k^2=-\frac{1}{\alpha'}.$$

$$Q_{\text{B}}|0;{\bf k}\rangle=0$$

$$|\psi_1\rangle=(e\cdot\alpha_{-1}+\beta b_{-1}+\gamma c_{-1})|0;{\bf k}\rangle,-k^2=0$$

$$\begin{aligned} \langle \psi_1 || \psi_1 \rangle &= \langle 0; \mathbf{k} | (e^* \cdot \alpha_1 + \beta^* b_1 + \gamma^* c_1) (e \cdot \alpha_{-1} + \beta b_{-1} + \gamma c_{-1}) | 0; \mathbf{k}' \rangle \\ &= (e^* \cdot e + \beta^* \gamma + \gamma^* \beta) \langle 0; \mathbf{k} || 0; \mathbf{k}' \rangle \end{aligned}$$

$$0=Q_{\text{B}}|\psi_1\rangle=(2\alpha')^{1/2}(c_{-1}k\cdot\alpha_1+c_1k\cdot\alpha_{-1})|\psi_1\rangle=(2\alpha')^{1/2}(k\cdot ec_{-1}+\beta k\cdot\alpha_{-1})|0;\mathbf{k}\rangle$$

$$Q_{\text{B}}|\chi\rangle=(2\alpha')^{1/2}(k\cdot e'c_{-1}+\beta'k\cdot\alpha_{-1})|0;\mathbf{k}\rangle.$$

$$b_0|\psi\rangle=\tilde{b}_0|\psi\rangle=0$$

$$L_0|\psi\rangle=\tilde{L}_0|\psi\rangle=0$$

$$L_0=\frac{\alpha'}{4}(p^2+m^2), \tilde{L}_0=\frac{\alpha'}{4}(p^2+\tilde{m}^2)$$

$$\begin{aligned} \frac{\alpha'}{4}m^2 &= \sum_{n=1}^{\infty} n \left(N_{bn} + N_{cn} + \sum_{\mu=0}^{25} N_{\mu n} \right) - 1, \\ \frac{\alpha'}{4}\tilde{m}^2 &= \sum_{n=1}^{\infty} n \left(\tilde{N}_{bn} + \tilde{N}_{cn} + \sum_{\mu=0}^{25} \tilde{N}_{\mu n} \right) - 1. \end{aligned}$$

$$L_m=L_m^X+L_m^K+L_m^g$$



$$|N, I; k\rangle, |N, \tilde{N}, I; k\rangle,$$

$$-\sum_{\mu=0}^{d-1} k_{\mu} k^{\mu} = m^2$$

$$\alpha' m^2 = \sum_{n=1}^{\infty} n \left(N_{bn} + N_{cn} + \sum_{\mu=0}^{d-1} N_{\mu n} \right) + L_0^K - 1$$

$$-\sum_{\mu=0}^{d-1} k_{\mu} k^{\mu} = m^2 = \tilde{m}^2$$

$$\frac{\alpha'}{4} m^2 = \sum_{n=1}^{\infty} n \left(N_{bn} + N_{cn} + \sum_{\mu=0}^{d-1} N_{\mu n} \right) + L_0^K - 1$$

$$\frac{\alpha'}{4} \tilde{m}^2 = \sum_{n=1}^{\infty} n \left(\tilde{N}_{bn} + \tilde{N}_{cn} + \sum_{\mu=0}^{d-1} \tilde{N}_{\mu n} \right) + \tilde{L}_0^K - 1$$

$$\langle 0, I; \mathbf{k} | 0, I'; \mathbf{k}' \rangle = \langle 0, 0, I; \mathbf{k} | 0, 0, I'; \mathbf{k}' \rangle = 2k^0 (2\pi)^{d-1} \delta^{d-1}(\mathbf{k} - \mathbf{k}') \delta_{I, I'}$$

$$\alpha_m^{\pm}=2^{-1/2}(\alpha_m^0\pm\alpha_m^1)$$

$$[\alpha_m^+, \alpha_n^-] = -m \delta_{m, -n}, [\alpha_m^+, \alpha_n^+] = [\alpha_m^-, \alpha_n^-] = 0$$

$$N^{\rm lc}=\sum_{\substack{m=-\infty\\m\neq 0}}^{\infty}\frac{1}{m}:\alpha_{-m}^+\alpha_m^-:$$

$$Q_{\rm B}=Q_1+Q_0+Q_{-1}$$

$$[N^{\rm lc},Q_j]=jQ_j$$

$$[N^{\rm g},Q_j]=Q_j$$

$$(Q_1^2) + (\{Q_1, Q_0\}) + (\{Q_1, Q_{-1}\} + Q_0^2) + (\{Q_0, Q_{-1}\}) + (Q_{-1}^2) = 0$$

$$Q_1=-(2\alpha')^{1/2}k^+\sum_{\substack{m=-\infty\\m\neq 0}}^{\infty}\alpha_{-m}^-c_m$$

$$R=\frac{1}{(2\alpha')^{1/2}k^+}\sum_{\substack{m=-\infty\\m\neq 0}}^{\infty}\alpha_{-m}^+b_m$$

$$S\equiv\{Q_1,R\}=\sum_{m=1}^{\infty}(mb_{-m}c_m+mc_{-m}b_m-\alpha_{-m}^+\alpha_m^- -\alpha_{-m}^-\alpha_m^+)=\sum_{m=1}^{\infty}m(N_{bm}+N_{cm}+N_m^++N_m^-)$$



$$|\psi\rangle = \frac{1}{S}\{Q_1, R\}|\psi\rangle = \frac{1}{S}Q_1R|\psi\rangle,$$

$$0 = Q_1S|\psi\rangle = SQ_1|\psi\rangle$$

$$S+U\equiv\{Q_{\rm B},R\}$$

$$|\psi\rangle=(1-S^{-1}U+S^{-1}US^{-1}U-\cdots)|\psi_0\rangle$$

$$\langle\psi||\psi'\rangle=\langle\psi_0||\psi'_0\rangle$$

Cuantización BRST-OCQ.

$$\mathcal{H}_{\text{OCQ}}=\mathcal{H}_{\text{BRST}}=\mathcal{H}_{\text{light-cone}}$$

$$|\psi,\downarrow\rangle$$

$$Q_{\rm B}|\psi,\downarrow\rangle=\sum_{n=0}^{\infty}c_{-n}(L_n^{\rm m}-\delta_{n,0})|\psi,\downarrow\rangle=0$$

$$|\psi,\downarrow\rangle-|\psi',\downarrow\rangle$$

$$|\psi,\downarrow\rangle-|\psi',\downarrow\rangle=Q_{\rm B}|\chi\rangle,$$

$$|\chi\rangle=\sum_{n=1}^{\infty}b_{-n}|\chi_n,\downarrow\rangle+\cdots;$$

$$|\psi,\downarrow\rangle-|\psi',\downarrow\rangle=\sum_{m,n=1}^{\infty}c_mL_{-m}^{\rm m}b_{-n}|\chi_n,\downarrow\rangle=\sum_{n=1}^{\infty}L_{-n}^{\rm m}|\chi_n,\downarrow\rangle$$

$$\{[Q_{\rm B},L_m],b_n\}-\{[L_m,b_n],Q_{\rm B}\}-[\{b_n,Q_{\rm B}\},L_m]=0.$$

$$G_I G_J - (-1)^{F_I F_J} G_J G_I = i g^K{}_{IJ} G_K$$

$$(b_{-1})^{N_b}(c_{-1})^{N_c}(\alpha_{-1}^+)^{N^+}(\alpha_{-1}^-)^{N^-}|0\rangle.$$

Simetrías de gauge.

$$\int\,[dedX]\mathrm{exp}\left[-\frac{1}{2}\int\,d\tau(e^{-1}\dot{X}^{\mu}\dot{X}_{\mu}+em^2)\right]$$

$$\frac{\partial \tau'}{\partial \tau}=e(\tau).$$

$$\tau'(\tau)=\int_0^{\tau}d\tau''e(\tau'')$$

$$\tau'(1)=\int_0^1d\tau e(\tau)=l$$



$$e'=l, 0\leq \tau\leq 1$$

$$e'=1, 0\leq \tau\leq l$$

$$\tau\rightarrow\tau+v\bmod 1.$$

$$0\leq \sigma^1\leq 2\pi, 0\leq \sigma^2\leq 2\pi$$

$$(\sigma^1,\sigma^2)\cong (\sigma^1,\sigma^2)+2\pi(m,n)$$

$$ds^2=|d\sigma^1+\tau d\sigma^2|^2$$

$$\tilde{\sigma}^a\cong\tilde{\sigma}^a+2\pi(mu^a+nv^a)$$

$$w\cong w+2\pi(m+n\tau)$$

$$T\colon \tau'=\tau+1, S\colon \tau'=-1/\tau$$

$$\tau'=\frac{a\tau+b}{c\tau+d}$$

$$\begin{bmatrix}\sigma^1\\ \sigma^2\end{bmatrix}=\begin{bmatrix}d&b\\ c&a\end{bmatrix}\begin{bmatrix}\sigma'^1\\ \sigma'^2\end{bmatrix}$$

$$-\frac{1}{2}\leq \operatorname{Re}\tau\leq \frac{1}{2}, |\tau|\geq 1$$

$$\sigma^a\rightarrow\sigma^a+v^a$$

Superficies de Moduli y Riemann.

$$\mathcal{M}_r=\frac{\mathcal{G}_r}{(\text{diff}\times\text{Weyl})_r}$$

$$\mathcal{M}_{r,n}=\frac{\mathcal{G}_r\times\mathcal{M}^n}{(\text{diff}\times\text{Weyl})_r}$$

$$\frac{\text{diff}_r}{\text{diff}_{r0}}$$

$$\delta g_{ab}=-2(P_1\delta\sigma)_{ab}+(2\delta\omega-\nabla\cdot\delta\sigma)g_{ab}$$

$$0=\int\,d^2\sigma g^{1/2}\delta'g_{ab}[-2(P_1\delta\sigma)^{ab}+(2\delta\omega-\nabla\cdot\delta\sigma)g^{ab}]$$

$$=\int\,d^2\sigma g^{1/2}[-2(P_1^T\delta'g)_a\delta\sigma^a+\delta'g_{ab}g^{ab}(2\delta\omega-\nabla\cdot\delta\sigma)]$$

$$g^{ab}\delta'g_{ab}=0$$

$$(P_1^T\delta'g)_a=0$$

$$(P_1\delta\sigma)_{ab}=0$$



$$\begin{aligned}\partial_{\bar{z}}\delta'g_{zz}&=\partial_z\delta'g_{\bar{z}\bar{z}}=0\\ \partial_{\bar{z}}\delta z&=\partial_z\delta\bar{z}=0\end{aligned}$$

$$\mu-\kappa=-3\chi$$

$$\int \, d^2\sigma g^{1/2} (P_1\delta\sigma)_{ab} (P_1\delta\sigma)^{ab} = \int \, d^2\sigma g^{1/2} \delta\sigma_a (P_1^TP_1\delta\sigma)^a = \int \, d^2\sigma g^{1/2} \left(\frac{1}{2}\nabla_a\delta\sigma_b\nabla^a\delta\sigma^b - \frac{R}{4}\delta\sigma_a\delta\sigma^a\right)$$

$$\begin{array}{l} \chi>0: \quad \kappa=3\chi, \mu=0 \\ \chi<0: \quad \kappa=0, \mu=-3\chi \end{array}$$

$$\sigma_m^a=f_{mn}^a(\sigma_n)$$

$$z_m^a=f_{mn}^a(z_n)$$

$$ds^2\propto dz_md\bar{z}_m$$

$$w\cong w+2\pi\cong w+2\pi\tau$$

$$S_{j_1\dots j_n}(k_1,\dots,k_n)=\sum_{\substack{\text{compact} \\ \text{topologies}}}\int\frac{[d\phi dg]}{V_{\text{diff}\times\text{Weyl}}}\exp\left(-S_{\text{m}}-\lambda\chi\right)\prod_{i=1}^n\int\,d^2\sigma_i g(\sigma_i)^{1/2}\mathcal{V}_{j_i}(k_i,\sigma_i)$$

$$[dg]d^{2n}\sigma\rightarrow[d\zeta]d^{\mu}td^{2n-\kappa}\sigma$$

$$1=\Delta_{\text{FP}}(g,\sigma)\int_Fd^{\mu}t\int_{\text{diff}\times\text{Weyl}}[d\zeta]\delta(g-\hat{g}(t)^{\zeta})\prod_{(a,i)\in f}\delta\left(\sigma_i^a-\hat{\sigma}_i^{\zeta a}\right)$$

$$\begin{aligned} S_{j_1\dots j_n}(k_1,\dots,k_n) &= \sum_{\substack{\text{compact} \\ \text{topologies}}} \int_F d^{\mu}t \Delta_{\text{FP}}(\hat{g}(t),\hat{\sigma}) \int [d\phi] \int \prod_{(a,i)\notin f} d\sigma_i^a \\ &\times \exp\left(-S_{\text{m}}[\phi,\hat{g}(t)]-\lambda\chi\right) \prod_{i=1}^n \left[\hat{g}(\sigma_i)^{1/2}\mathcal{V}_{j_i}(k_i;\sigma_i)\right] \end{aligned}$$

$$\delta g_{ab}=\sum_{k=1}^{\mu}\delta t^k\partial_{t^k}\hat{g}_{ab}-2(\hat{p}_1\delta\sigma)_{ab}+(2\delta\omega-\hat{\nabla}\cdot\delta\sigma)\hat{g}_{ab}$$

$$\begin{aligned}\Delta_{\text{FP}}(\hat{g},\hat{\sigma})^{-1}&=n_R\int\,d^{\mu}\delta t[d\delta\omega d\delta\sigma]\delta(\delta g_{ab})\prod_{(a,i)\in f}\delta(\delta\sigma^a(\hat{\sigma}_i))\\ &=n_R\int\,d^{\mu}\delta t d^{\kappa}x[d\beta'd\delta\sigma]\times\exp\left[2\pi i(\beta',2\hat{p}_1\delta\sigma-\delta t^k\partial_k\hat{g})+2\pi i\sum_{(a,i)\in f}x_{ai}\delta\sigma^a(\hat{\sigma}_i)\right]\end{aligned}$$

$$\begin{array}{l} \delta\sigma^a\rightarrow c^a\\ \beta'_{ab}\rightarrow b_{ab}\\ x_{ai}\rightarrow\eta_{ai}\\ \delta t^k\rightarrow\xi^k\end{array}$$

$$\Delta_{\text{FP}}(\hat{g}, \hat{\sigma}) = \frac{1}{n_R} \int [dbdc] d^\mu \xi d^\kappa \eta \times \exp \left[-\frac{1}{4\pi} (b, 2\hat{P}_1 c - \xi^k \partial_k \hat{g}) + \sum_{(a,i) \in f} \eta_{ai} c^a(\hat{\sigma}_i) \right]$$

$$= \frac{1}{n_R} \int [dbdc] \exp(-S_g) \prod_{k=1}^{\mu} \frac{1}{4\pi} (b, \partial_k \hat{g}) \prod_{(a,i) \in f} c^a(\hat{\sigma}_i)$$

$$S_{j_1 \dots j_n}(k_1, \dots, k_n)$$

$$= \sum_{\substack{\text{compact} \\ \text{topologies}}} \int_F \frac{d^\mu t}{n_R} \int [d\phi dbdc] \exp(-S_m - S_g - \lambda \chi)$$

$$\times \prod_{(a,i) \notin f} \int d\sigma_i^a \prod_{k=1}^{\mu} \frac{1}{4\pi} (b, \partial_k \hat{g}) \prod_{(a,i) \in f} c^a(\hat{\sigma}_i) \prod_{i=1}^n \hat{g}(\sigma_i)^{1/2} \mathcal{V}_{j_i}(k_i, \sigma_i)$$

$$c^a(\sigma) = \sum_J c_J C_J^a(\sigma), b_{ab}(\sigma) = \sum_K b_K B_{Kab}(\sigma).$$

$$S_g = \frac{1}{2\pi} (b, P_1 c) = \frac{1}{2\pi} (P_1^T b, c)$$

$$P_1^T P_1 C_J^a = v_J'^2 C_J^a, P_1 P_1^T B_{Kab} = v_K^2 B_{Kab}$$

$$(C_J, C_{J'}) = \int d^2 \sigma g^{1/2} C_J^a C_{J'a} = \delta_{JJ'}$$

$$(B_K, B_{K'}) = \int d^2 \sigma g^{1/2} B_{Kab} B_{K'}^{ab} = \delta_{KK'}$$

$$(P_1 P_1^T) P_1 C_J = P_1 (P_1^T P_1) C_J = v_J'^2 P_1 C_J$$

$$B_{Jab} = \frac{1}{v_J} (P_1 C_J)_{ab}, v_J = v_J' \neq 0$$

$$\int \prod_{k=1}^{\mu} db_{0k} \prod_{j=1}^{\kappa} dc_{0j} \prod_J db_J dc_J \exp \left(-\frac{v_J b_J c_J}{2\pi} \right) \prod_{k'=1}^{\mu} \frac{1}{4\pi} (b, \partial_{k'} \hat{g}) \prod_{(a,i) \in f} c^a(\sigma_i)$$

$$\Delta_{\text{FP}} = \int \prod_{k=1}^{\mu} db_{0k} \prod_{k'=1}^{\mu} \left[\sum_{k''=1}^{\mu} \frac{b_{0k''}}{4\pi} (B_{0k''}, \partial_{k'} \hat{g}) \right] \times \int \prod_{j=1}^{\kappa} dc_{0j} \prod_{(a,i) \in f} \left[\sum_{j'=1}^{\kappa} c_{0j'} C_{0j'}^a(\sigma_i) \right]$$

$$\times \int \prod_J db_J dc_J \exp \left(-\frac{v_J b_J c_J}{2\pi} \right)$$

$$\Delta_{\text{FP}} = \det \frac{(B_{0k}, \partial_{k'} \hat{g})}{4\pi} \det C_{0j}^a(\sigma_i) \det' \left(\frac{P_1^T P_1}{4\pi^2} \right)^{1/2}$$



Teorema de Riemann-Roch.

$$\nabla_a j^a = \frac{1-2\lambda}{4} R$$

$$\frac{\delta([d\phi]\exp(-S))}{[d\phi]\exp(-S)} = \frac{i\epsilon}{2\pi} \int d^2\sigma g^{1/2} \nabla_a j^a \rightarrow -i\epsilon \frac{2\lambda-1}{2} \chi$$

$$\frac{1}{2}(\kappa-\mu)=\frac{1}{2}(\dim\ker P_1-\dim\ker P_1^T).$$

$$\dim\ker P_n-\dim\ker P_n^T=(2n+1)\chi,$$

$$d^\mu t'=\left|\det\frac{\partial t'}{\partial t}\right|d^\mu t\prod_{k=1}^\mu\frac{1}{4\pi}(b,\partial'_k\hat{g})=\det\left(\frac{\partial t}{\partial t'}\right)\prod_{k=1}^\mu\frac{1}{4\pi}(b,\partial_k\hat{g})$$

$$\delta \hat{g}'_{ab}(t;\sigma)=2\delta\omega(t;\sigma)\hat{g}_{ab}(t;\sigma),$$

$$(b,\partial_k\hat{g}')=\int d^2\sigma\hat{g}'^{1/2}b_{ab}\hat{g}'^{ac}\hat{g}'^{bd}\partial_k\hat{g}'_{cd}=\int d^2\sigma\hat{g}^{1/2}b_{ab}(\hat{g}^{ac}\hat{g}^{bd}\partial_k\hat{g}_{cd}+2\hat{g}^{ab}\partial_k\omega)$$

$$=(b,\partial_k\hat{g})$$

$$\delta(b,\partial_k\hat{g})=-2(b,P_1\partial_k\xi)=-2(P_1^Tb,\partial_k\xi)=0$$

Invariancia BRST.

$$\delta_{\rm B}\mathcal{V}_{\rm m}=i\epsilon\partial_a(c^a\mathcal{V}_{\rm m})$$

$$\delta_{\rm B}(b,\partial_k\hat{g})=i\epsilon(T,\partial_k\hat{g})$$

$$\mu_{ka}^b=\frac{1}{2}\hat{g}^{bc}\partial_k\hat{g}_{ac}$$

$$\frac{1}{2\pi}(b,\mu_k)=\frac{1}{2\pi}\int d^2z(b_{zz}\mu_{k\bar{z}}^z+b_{\bar{z}\bar{z}}\mu_{kz}^{\bar{z}})$$

$$z'_m=z_m+\delta t^k v_{km}^{z_m}(z_m,\bar{z}_m)$$

$$dz'_m d\bar{z}'_m \propto dz_md\bar{z}_m + \delta t^k \left(\mu_{kz_m}^{\bar{z}_m} dz_mdz_m + \mu_{k\bar{z}_m}^{z_m} d\bar{z}_md\bar{z}_m \right)$$

$$\mu_{kz_m}^{\bar{z}_m}=\partial_{z_m}v_{km}^{\bar{z}_m},\mu_{k\bar{z}_m}^{z_m}=\partial_{\bar{z}_m}v_{km}^{z_m}$$

$$\frac{1}{2\pi}(b,\mu_k)=\frac{1}{2\pi i}\sum_m\oint_{c_m}(dz_mv_{km}^{z_m}b_{z_mz_m}-d\bar{z}_mv_{km}^{\bar{z}_m}b_{\bar{z}_m\bar{z}_m})$$

$$\frac{dz_m}{dt^k}=v_{km}^{z_m}$$



$$\left.\frac{\partial z_m}{\partial t^k}\right|_{z_n}=v_{km}^{z_m}-\left.\frac{\partial z_m}{\partial z_n}\right|_tv_{kn}^{z_n}=v_{km}^{z_m}-v_{kn}^{z_m}$$

$$\frac{1}{2\pi}(b,\mu_k)=\frac{1}{2\pi i}\sum_{(mn)}\int_{C_{mn}}\left(dz_m\left.\frac{\partial z_m}{\partial t^k}\right|_{z_n}b_{z_mz_m}-d\bar{z}_m\left.\frac{\partial \bar{z}_m}{\partial t^k}\right|_{z_n}b_{\bar{z}_m\bar{z}_m}\right)$$

$$z=z'+z_v$$

$$\left.\frac{\partial z'}{\partial z_v}\right|_z=-1.$$

$$\int_c\frac{dz'}{2\pi i}b_{z'z'}\int_c\frac{d\bar{z}'_m}{-2\pi i}b_{\bar{z}'\bar{z}'}=b_{-1}\tilde{b}_{-1}$$

$$S(1;\ldots;n)=\sum_{\substack{\text{compact} \\ \text{topologies}}}e^{-\lambda\chi}\int_F\frac{d^mt}{n_R}\Big\langle\prod_{k=1}^mB_k\prod_{i=1}^n\hat{\nu}_i\Big\rangle$$

$$m=\mu+2n_{\rm c}+n_{\rm o}-\kappa=-3\chi+2n_{\rm c}+n_{\rm o}$$

$$b_{-1}\tilde{b}_{-1}\cdot\tilde{c}c\mathcal{V}_{\mathfrak{m}}=\mathcal{V}_{\mathfrak{m}}$$

$$u=1/z.$$

$$ds^2=\exp\left(2\omega(z,\bar{z})\right)dzd\bar{z}$$

$$ds^2=\frac{4r^2dzd\bar{z}}{(1+z\bar{z})^2}=\frac{4r^2dud\bar{u}}{(1+u\bar{u})^2}$$

$$\delta u=\frac{\partial u}{\partial z}\delta z=-z^{-2}\delta z$$

$$\delta g_{uu}=\left(\frac{\partial u}{\partial z}\right)^{-2}\delta g_{zz}=z^4\delta g_{zz}$$

$$\begin{array}{l} \delta z = a_0 + a_1 z + a_2 z^2 \\ \delta \bar{z} = a_0^* + a_1^* \bar{z} + a_2^* \bar{z}^2 \end{array}$$

$$z'=\frac{\alpha z+\beta}{\gamma z+\delta}$$

$$z'=1/\bar{z}$$

$$z'=\bar{z}$$

$$z'=-1/\bar{z}$$

$$Z[J]=\left\langle \exp\left(i\int\,d^2\sigma J(\sigma)\cdot X(\sigma)\right)\right\rangle$$



$$X^{\mu}(\sigma) = \sum_I x_I^{\mu} X_I(\sigma)$$

$$\nabla^2 X_I = -\omega_I^2 X_I$$

$$\int_M d^2\sigma g^{1/2} X_I X_{I'} = \delta_{II'}$$

$$Z[J] = \prod_{I,\mu} \int dx_I^{\mu} \exp \left(-\frac{\omega_I^2 x_I^{\mu} x_{I\mu}}{4\pi\alpha'} + i x_I^{\mu} J_{I\mu} \right)$$

$$J_I^{\mu} = \int d^2\sigma J^{\mu}(\sigma) X_I(\sigma)$$

$$X_0 = \left(\int d^2\sigma g^{1/2} \right)^{-1/2}$$

$$Z[J] = i(2\pi)^d \delta^d(J_0) \prod_{I \neq 0} \left(\frac{4\pi^2 \alpha'}{\omega_I^2} \right)^{\frac{d}{2}} \exp \left(-\frac{\pi \alpha' J_I \cdot J_I}{\omega_I^2} \right)$$

$$= i(2\pi)^d \delta^d(J_0) \left(\det' \frac{-\nabla^2}{4\pi^2 \alpha'} \right)^{-d/2} \times \exp \left(-\frac{1}{2} \int d^2\sigma d^2\sigma' J(\sigma) \cdot J(\sigma') G'(\sigma, \sigma') \right)$$

$$G'(\sigma_1, \sigma_2) = \sum_{I \neq 0} \frac{2\pi \alpha'}{\omega_I^2} X_I(\sigma_1) X_I(\sigma_2)$$

$$-\frac{1}{2\pi\alpha'} \nabla^2 G'(\sigma_1, \sigma_2) = \sum_{I \neq 0} X_I(\sigma_1) X_I(\sigma_2) = g^{-1/2} \delta^2(\sigma_1 - \sigma_2) - X_0^2$$

$$G'(\sigma_1, \sigma_2) = -\frac{\alpha'}{2} \ln |z_{12}|^2 + f(z_1, \bar{z}_1) + f(z_2, \bar{z}_2)$$

$$f(z, \bar{z}) = \frac{\alpha' X_0^2}{4} \int d^2 z' \exp [2\omega(z', \bar{z}')] \ln |z - z'|^2 + k$$

$$A_{S_2}^n(k,\sigma)=\left\langle [e^{ik_1\cdot X(\sigma_1)}]_{\rm r}[e^{ik_2\cdot X(\sigma_2)}]_{\rm r}\ldots [e^{ik_n\cdot X(\sigma_n)}]_{\rm r}\right\rangle_{S_2}$$

$$J(\sigma) = \sum_{i=1}^n k_i \delta^2(\sigma - \sigma_i)$$

$$A_{S_2}^n(k,\sigma) = iC_{S_2}^X(2\pi)^d\delta^d\left(\sum_i k_i\right)\times \exp\left(-\sum_{\substack{i,j=1\\i<j}}^n k_i\cdot k_j G'(\sigma_i,\sigma_j)-\frac{1}{2}\sum_{i=1}^n k_i^2 G'_{\rm r}(\sigma_i,\sigma_i)\right)$$

$$C_{S_2}^X = X_0^{-d} \left(\det' \frac{-\nabla^2}{4\pi^2 \alpha'} \right)_{S_2}^{-d/2}$$



$$G'_r(\sigma, \sigma') = G'(\sigma, \sigma') + \frac{\alpha'}{2} \ln d^2(\sigma, \sigma')$$

$$G_r'(\sigma, \sigma) = 2f(z, \bar{z}) + \alpha' \omega(z, \bar{z})$$

$$A_{S_2}^n(k, \sigma) = iC_{S_2}^X (2\pi)^d \delta^d \left(\sum_i k_i \right) \exp \left(-\frac{\alpha'}{2} \sum_i k_i^2 \omega(\sigma_i) \right) \prod_{\substack{i,j=1 \\ i < j}}^n |z_{ij}|^{\alpha' k_i \cdot k_j}$$

$$\left\langle \prod_{i=1}^n [e^{ik_i \cdot X(z_i, \bar{z}_i)}]_{\mathbf{r}} \prod_{j=1}^p \partial X^{\mu_j}(z'_j) \prod_{k=1}^q \bar{\partial} X^{\nu_k}(\bar{z}''_k) \right\rangle_{S_2}$$

$$iC_{S_2^X}^X(2\pi)^d\delta^d\left(\sum_i k_i\right)\prod_{\substack{i,j=1 \\ i < j}}^n |z_{ij}|^{\alpha'k_i\cdot k_j}\times\left\langle\prod_{j=1}^p [v^{\mu_j}(z'_j)+q^{\mu_j}(z'_j)]\prod_{k=1}^q [\tilde{v}^{v_k}(\bar{z}''_k)+\tilde{q}^{v_k}(\bar{z}''_k)]\right\rangle$$

$$v^\mu(z) = -i \frac{\alpha'}{2} \sum_{i=1}^n \frac{k_i^\mu}{z - z_i}, \tilde{v}^\mu(\bar{z}) = -i \frac{\alpha'}{2} \sum_{i=1}^n \frac{k_i^\mu}{\bar{z} - \bar{z}_i}$$

$$\langle \partial X^\mu(z_1) \partial X^\nu(z_2) \rangle_{S_2}$$

$$-\frac{\alpha'\eta^{\mu\nu}}{2z_{12}^2}\langle 1\rangle_{S_2}+g(z_1,z_2)$$

$$\partial_u X^\mu = -z^2 \partial_z X^\mu$$

$$A_{S_2}^n(k, \sigma) = \left\langle \prod_{i=1}^n : e^{ik_i \cdot X(z_i, \bar{z}_i)} : \right\rangle_{S_2}$$

$$\left\langle \partial X^\mu(z) \prod_{i=1}^n : e^{ik_i \cdot X(z_i, \bar{z}_i)} : \right\rangle_{S_2} = -\frac{i\alpha'}{2} A_{S_2}^n(k, \sigma) \sum_{i=1}^n \frac{k_i^\mu}{z - z_i} + \lambda_{\text{terms holomorphic in } z}$$

$$A_{S_2}^n(k, \sigma) \sum_{i=1}^n k_i^\mu = 0$$

$$p^\mu = \frac{1}{2\pi i} \oint_c (dz j_z^\mu - d\bar{z} j_{\bar{z}}^\mu)$$

$$-\frac{i\alpha'}{2}A_{S_2}^n(k,\sigma)\left(\frac{k_1^\mu}{z-z_1}+\sum_{i=2}^n\frac{k_i^\mu}{z_1-z_i}+O(z-z_1)\right).$$

$$ik_1 \cdot \partial X(z) : e^{ik_1 \cdot X(z_1, \bar{z}_1)} : = \frac{\alpha' k_1^2}{2(z - z_1)} : e^{ik_1 \cdot X(z_1, \bar{z}_1)} : + \partial_{z_1} : e^{ik_1 \cdot X(z_1, \bar{z}_1)} : + O(z - z_1)$$



$$\partial_{z_1} A_{S_2}^n(k, \sigma) = \frac{\alpha'}{2} A_{S_2}^n(k, \sigma) \sum_{i=2}^n \frac{k_1 \cdot k_i}{z_{1i}}.$$

$$A_{S_2}^n(k, \sigma) \propto \delta^d \left(\sum_i k_i \right) \prod_{\substack{i,j=1 \\ i < j}}^n |z_{ij}|^{\alpha' k_i \cdot k_j},$$

$$G'(\sigma_1, \sigma_2) = -\frac{\alpha'}{2} \ln |z_1 - z_2|^2 - \frac{\alpha'}{2} \ln |z_1 - \bar{z}_2|^2,$$

$$\left\langle \prod_{i=1}^n : e^{ik_i \cdot X(z_i, \bar{z}_i)} : \right\rangle_{D_2}$$

$$= iC_{D_2}^X (2\pi)^d \delta^d \left(\sum_i k_i \right) \prod_{i=1}^n |z_i - \bar{z}_i|^{\alpha' k_i^2/2} \times \prod_{\substack{i,j=1 \\ i < j}}^n |z_i - z_j|^{\alpha' k_i \cdot k_j} |z_i - \bar{z}_j|^{\alpha' k_i \cdot k_j}$$

$${}^*X^\mu(y_1)X^v(y_2){}_* = X^\mu(y_1)X^v(y_2) + 2\alpha'\eta^{\mu v}\ln|y_1-y_2|,$$

$$\left\langle \prod_{i=1}^n {}^*e^{ik_i \cdot X(y_i)} {}^* \right\rangle_{D_2} = iC_{D_2}^X (2\pi)^d \delta^d \left(\sum_i k_i \right) \prod_{\substack{i,j=1 \\ i < j}}^n |y_i - y_j|^{2\alpha' k_i \cdot k_j}$$

$$\left\langle \prod_{i=1}^n \star \star e^{ik_i \cdot X(y_i)} \prod_{j=1}^p \partial_y X^{\mu_j}(y'_j) \right\rangle_{D_2}$$

$$= iC_{D_2}^X (2\pi)^d \delta^d \left(\sum_i k_i \right) \times \prod_{\substack{i,j=1 \\ i < j}}^n |y_{ij}|^{2\alpha' k_i \cdot k_j} \left\langle \prod_{j=1}^p [v^{\mu_j}(y'_j) + q^{\mu_j}(y'_j)] \right\rangle_{D_2}$$

$$v^\mu(y)=-2i\alpha'\sum_{i=1}^n\frac{k_i^\mu}{y-y_i}$$

$$G'(\sigma_1, \sigma_2) = -\frac{\alpha'}{2} \ln |z_1 - z_2|^2 - \frac{\alpha'}{2} \ln |1 + z_1 \bar{z}_2|^2.$$



$$\left\langle \prod_{i=1}^n : e^{ik_i \cdot X(z_i, \bar{z}_i)} : \right\rangle_{RP_2}$$

$$= iC_{RP_2}^X (2\pi)^d \delta^d \left(\sum_i k_i \right) \prod_{i=1}^n |1 + z_i \bar{z}_i|^{\alpha' k_i^2/2}$$

$$\times \prod_{\substack{i,j=1 \\ i < j}}^n |z_i - z_j|^{\alpha' k_i \cdot k_j} |1 + z_i \bar{z}_j|^{\alpha' k_i \cdot k_j}$$

Métrica CFT.

$$\langle c(z_1)c(z_2)c(z_3)\tilde{c}(\bar{z}_4)\tilde{c}(\bar{z}_5)\tilde{c}(\bar{z}_6)\rangle_{S_2}$$

$$\det C_{0j}^a(\sigma_i)$$

$$C^Z=1, z, z^2, C^{\bar{Z}}=0$$

$$C^Z=0, C^{\bar{Z}}=1, \bar{z}, \bar{z}^2$$

$$C_{S_2}^{\rm g} \det \begin{vmatrix} 1 & 1 & 1 \\ z_1 & z_2 & z_3 \\ z_1^2 & z_2^2 & z_3^2 \end{vmatrix} \det \begin{vmatrix} 1 & 1 & 1 \\ \bar{z}_4 & \bar{z}_5 & \bar{z}_6 \\ \bar{z}_4^2 & \bar{z}_5^2 & \bar{z}_6^2 \end{vmatrix} = C_{S_2}^{\rm g} z_{12} z_{13} z_{23} \bar{z}_{45} \bar{z}_{46} \bar{z}_{56}$$

$$\left\langle \prod_{i=1}^{p+3} c(z_i) \prod_{j=1}^p b(z'_j) \cdot (\text{anti}) \right\rangle_{S_2} = C_{S_2}^{\rm g} \frac{z_{p+1,p+2} z_{p+1,p+3} z_{p+2,p+3}}{(z_1-z'_1) \dots (z_p-z'_p)} \cdot (\mathfrak{I}_{\text{anti}}) \pm \boxtimes_{\text{permutations}}$$

$$\oint_c \frac{dz}{2\pi i} j_z = -\oint_c \frac{du}{2\pi i} j_u + 3 \rightarrow 3$$

$$z_{12}z_{13}z_{23}\bar{z}_{45}\bar{z}_{46}\bar{z}_{56}F(z_1,z_2,z_3)\tilde{F}(\bar{z}_4,\bar{z}_5,\bar{z}_6)$$

$$C_{S_2}^{\rm g} \prod_{\substack{i,i'=1 \\ i < i'}}^{p+3} z_{ii'} \prod_{\substack{j,j'=1 \\ j < j'}}^p z'_{jj'} \prod_{i=1}^{p+3} \prod_{j=1}^p (z_i - z'_j)^{-1} \cdot (\text{anti}) ,$$

$$\tilde{b}(\bar{z})=b(z'), \tilde{c}(\bar{z})=c(z'), z'= \bar{z}, \text{Im} z > 0$$

$$\langle c(z_1)c(z_2)c(z_3)\rangle_{D_2}=C_{D_2}^{\rm g} z_{12}z_{13}z_{23}$$

$$\langle c(z_1)c(z_2)\tilde{c}(\bar{z}_3)\rangle_{D_2}=\langle c(z_1)c(z_2)c(z'_3)\rangle_{D_2}=C_{D_2}^{\rm g} z_{12}(z_1-\bar{z}_3)(z_2-\bar{z}_3)$$

$$\tilde{b}(\bar{z})=\left(\frac{\partial z'}{\partial \bar{z}}\right)^2 b(z')=z'^4 b(z'), \tilde{c}(\bar{z})=\left(\frac{\partial z'}{\partial \bar{z}}\right)^{-1} c(z')=z'^{-2} c(z')$$

$$\langle c(z_1)c(z_2)c(z_3)\rangle_{RP_2}=C_{RP_2}^{\rm g} z_{12}z_{13}z_{23}$$



$$\langle c(z_1)c(z_2)\tilde{c}(\bar{z}_3)\rangle_{RP_2} = z_3'^{-2}\langle c(z_1)c(z_2)c(z_3')\rangle_{RP_2} = C_{RP_2}^{\mathbf{g}} z_{12}(1+z_1\bar{z}_3)(1+z_2\bar{z}_3)$$

$$S_{D_2}(k_1;k_2;k_3)=g_0^3e^{-\lambda}\big\langle \star c^1e^{ik_1\cdot X}(y_1)_{\star\star}^{\star}c^1e^{ik_2\cdot X}(y_2)_{\star\star}^{\star}c^1e^{ik_3\cdot X}(y_3)_{\star}^{\star}\big\rangle_{D_2}+(k_2\leftrightarrow k_3)$$

$$S_{D_2}(k_1;k_2;k_3)=ig_0^3C_{D_2}(2\pi)^{26}\delta^{26}\left(\sum_i k_i\right)\times |y_{12}|^{1+2\alpha'k_1\cdot k_2}|y_{13}|^{1+2\alpha'k_1\cdot k_3}|y_{23}|^{1+2\alpha'k_2\cdot k_3}\\ + (k_2\leftrightarrow k_3)$$

$$2\alpha'k_1\cdot k_2=\alpha'(k_3^2-k_1^2-k_2^2)=-1$$

$$S_{D_2}(k_1;k_2;k_3)=2ig_0^3C_{D_2}(2\pi)^{26}\delta^{26}\left(\sum_i k_i\right)$$

$$S_{D_2}(k_1;k_2;k_3;k_4)=g_0^4e^{-\lambda}\int_{-\infty}^{\infty}dy_4\left\langle\prod_{i=1}^3\star c^1(y_i)e^{ik_i\cdot X(y_i)}\star\star e^{ik_4\cdot X(y_4)}\star\right\rangle+(k_2\leftrightarrow k_3)\\ =ig_0^4C_{D_2}(2\pi)^{26}\delta^{26}\left(\sum_i k_i\right)|y_{12}y_{13}y_{23}|\int_{-\infty}^{\infty}dy_4\prod_{i<j}|y_{ij}|^{2\alpha'k_i\cdot k_j}+(k_2\leftrightarrow k_3)$$

$$s=-(k_1+k_2)^2,t=-(k_1+k_3)^2,u=-(k_1+k_4)^2.$$

$$s+t+u=\sum_i m_i^2=-\frac{4}{\alpha'}$$

$$S_{D_2}(k_1;k_2;k_3;k_4)=ig_0^4C_{D_2}(2\pi)^{26}\delta^{26}\left(\sum_i k_i\right)\times\left[\int_{-\infty}^{\infty}dy_4|y_4|^{-\alpha'u-2}|1-y_4|^{-\alpha't-2}+(t\rightarrow s)\right]$$

$$S_{D_2}(k_1;k_2;k_3;k_4)=2ig_0^4C_{D_2}(2\pi)^{26}\delta^{26}\left(\sum_i k_i\right)[I(s,t)+I(t,u)+I(u,s)]$$

$$I(s,t)=\int_0^1 dy y^{-\alpha's-2}(1-y)^{-\alpha't-2}$$

$$I(s,t)=\int_0^r dy y^{-\alpha's-2}+\,\,\, \varsigma_{\text{terms analytic at }\alpha's=-1}=-\frac{r^{-\alpha's-1}}{\alpha's+1}+\,\,\, \varsigma_{\text{terms analytic at }\alpha's=-1}$$

$$=-\frac{1}{\alpha's+1}+\,\,\, \varsigma_{\text{terms analytic at }\alpha's=-1}$$

$$\frac{1}{\alpha's+1}\equiv\frac{1}{\alpha's+1+i\epsilon}\equiv\text{P}\frac{1}{\alpha's+1}-i\pi\delta(\alpha's+1)$$

$$S_{D_2}(k_1; k_2; k_3; k_4) = i \int \frac{d^{26}k}{(2\pi)^{26}} \frac{S_{D_2}(k_1; k_2; k) S_{D_2}(-k; k_3; k_4)}{-k^2 + \alpha'^{-1} + i\epsilon} + \text{q terms analytic at } k^2=1/\alpha'$$

$$C_{D_2} = e^{-\lambda} C_{D_2}^X C_{D_2}^g = \frac{1}{\alpha' g_0^2}$$

$$S_{D_2}(k_1;k_2;k_3)=\frac{2ig_o}{\alpha'}(2\pi)^{26}\delta^{26}\left(\sum_i k_i\right)$$

$$I(s,t)=\int_0^r dy[y^{-\alpha's-2}+(\alpha't+2)y^{-\alpha's-1}+\cdots]$$

$$I(s,t)=\frac{u-t}{2s}+\text{ terms analytic at } \alpha's=0$$

$$\alpha's=-1,0,1,2,\ldots$$

$$B(a,b)=\int_0^1 dy y^{a-1}(1-y)^{b-1}$$

$$I(s,t)=B(-\alpha_0(s),-\alpha_0(t)),\alpha_0(x)=1+\alpha'x$$

$$w^{a+b-1}B(a,b)=\int_0^w dv v^{a-1}(w-v)^{b-1}$$

$$\Gamma(a+b)B(a,b)=\int_0^\infty dv v^{a-1}e^{-v}\int_0^\infty d(w-v)(w-v)^{b-1}e^{-(w-v)}=\Gamma(a)\Gamma(b)$$

$$S_{D_2}(k_1;k_2;k_3;k_4)$$

$$=\frac{2ig_0^2}{\alpha'}(2\pi)^{26}\delta^{26}\left(\sum_i k_i\right)$$

$$\times [B(-\alpha_o(s),-\alpha_o(t))+B(-\alpha_o(s),-\alpha_o(u))+B(-\alpha_o(t),-\alpha_o(u))]$$

$$B(-\alpha_o(x),-\alpha_o(y))=\frac{\Gamma(-\alpha'x-1)\Gamma(-\alpha'y-1)}{\Gamma(-\alpha'x-\alpha'y-2)}$$

$$s\rightarrow\infty,t\,\mathrm{fixed}$$

$$s\rightarrow\infty,t/s\,\mathrm{fixed}$$

$$s=E^2,t=(4m^2-E^2)\sin^2\frac{\theta}{2},u=(4m^2-E^2)\cos^2\frac{\theta}{2}$$

$$S_{D_2}(k_1;k_2;k_3;k_4)\propto s^{\alpha_o(t)}\Gamma(-\alpha_o(t))$$

$$S_{D_2}(k_1;k_2;k_3;k_4)\approx \exp\left[-\alpha'(\mathrm{sln}\,\,s\alpha'+t\mathrm{ln}\,\,t\alpha'+u\mathrm{ln}\,\,u\alpha')\right]=\exp\left[-\alpha'sf(\theta)\right]$$



$$f(\theta) \approx -\sin^2 \frac{\theta}{2} \ln \sin^2 \frac{\theta}{2} - \cos^2 \frac{\theta}{2} \ln \cos^2 \frac{\theta}{2}$$

Factores Chan-Paton e interacciones de gauge.

$$|N; k; ij\rangle,$$

$$\mathrm{Tr}(\lambda^a \lambda^b) = \delta^{ab},$$

$$|N; k; a\rangle = \sum_{i,j=1}^n |N; k; ij\rangle \lambda_{ij}^a$$

$$\mathrm{Tr}(\lambda^{a_1} \lambda^{a_2} \lambda^{a_3} \lambda^{a_4})$$

$$S_{D_2}(k_1, a_1; k_2, a_2; k_3, a_3) = \frac{ig_o}{\alpha'} (2\pi)^{26} \delta^{26} \left(\sum_i k_i \right) \mathrm{Tr}(\lambda^{a_1} \lambda^{a_2} \lambda^{a_3} + \lambda^{a_1} \lambda^{a_3} \lambda^{a_2})$$

$$S_{D_2}(k_1, a_1; k_2, a_2; k_3, a_3; k_4, a_4)$$

$$= \frac{ig_o^2}{\alpha'} (2\pi)^{26} \delta^{26} \left(\sum_i k_i \right)$$

$$\times [\mathrm{Tr}(\lambda^{a_1} \lambda^{a_2} \lambda^{a_4} \lambda^{a_3} + \lambda^{a_1} \lambda^{a_3} \lambda^{a_4} \lambda^{a_2}) B(-\alpha_o(s), -\alpha_o(t))$$

$$+ \mathrm{Tr}(\lambda^{a_1} \lambda^{a_3} \lambda^{a_2} \lambda^{a_4}$$

$$+ \lambda^{a_1} \lambda^{a_4} \lambda^{a_2} \lambda^{a_3}) B(-\alpha_o(t), -\alpha_o(u)) + \mathrm{Tr}(\lambda^{a_1} \lambda^{a_2} \lambda^{a_3} \lambda^{a_4}$$

$$+ \lambda^{a_1} \lambda^{a_4} \lambda^{a_3} \lambda^{a_2}) B(-\alpha_o(s), -\alpha_o(u))].$$

$$\frac{1}{4} \mathrm{Tr}(\{\lambda^{a_1}, \lambda^{a_2}\} \{\lambda^{a_3}, \lambda^{a_4}\})$$

$$\frac{1}{4} \sum_a \mathrm{Tr}(\{\lambda^{a_1}, \lambda^{a_2}\} \lambda^a) \mathrm{Tr}(\{\lambda^{a_3}, \lambda^{a_4}\} \lambda^a)$$

$$\mathrm{Tr}(A \lambda^a) \mathrm{Tr}(B \lambda^a) = \mathrm{Tr}(AB)$$

$$S_{D_2}(k_1, a_1, e_1; k_2, a_2; k_3, a_3)$$

$$= -ig'_o g_o^2 e^{-\lambda} e_{1\mu}$$

$$\times \left\langle {}^*c^1 \dot{X}^\mu e^{ik_1 \cdot X} (y_1)_* {}^*c^1 e^{ik_2 \cdot X} (y_2)_* {}^*c^1 e^{ik_3 \cdot X} (y_3)_* \right\rangle_{D_2} \mathrm{Tr}(\lambda^{a_1} \lambda^{a_2} \lambda^{a_3}) + (k_2, a_2)$$

$$\leftrightarrow (k_3, a_3)$$



$$\begin{aligned}
& \langle {}^* \dot{X}^\mu e^{ik_1 \cdot X} (y_1)_{**}^* e^{ik_2 \cdot X} (y_2)_{**}^* e^{ik_3 \cdot X} (y_3)_{**}^* \rangle_{D_2} \\
&= -2i\alpha' \left(\frac{k_2^\mu}{y_{12}} + \frac{k_3^\mu}{y_{13}} \right) \\
&\quad \times iC_{D_2}^X (2\pi)^{26} \delta^{26} \left(\sum_i k_i \right) |y_{12}|^{2\alpha' k_1 \cdot k_2} |y_{13}|^{2\alpha' k_1 \cdot k_3} |y_{23}|^{2\alpha' k_2 \cdot k_3} \\
S_{D_2}(k_1, a_1, e_1; k_2, a_2; k_3, a_3) &= -ig'_o e_1 \cdot k_{23} (2\pi)^{26} \delta^{26} \left(\sum_i k_i \right) \text{Tr}(\lambda^{a_1} [\lambda^{a_2}, \lambda^{a_3}]) \\
&\quad \text{Tr}([\lambda^{a_1}, \lambda^{a_2}] [\lambda^{a_3}, \lambda^{a_4}]) \\
g'_o &= (2\alpha')^{-1/2} g_o
\end{aligned}$$

$$\begin{aligned}
S_{D_2}(k_1, a_1, e_1; k_2, a_2, e_2; k_3, a_3, e_3) \\
&= ig'_o (2\pi)^{26} \delta^{26} \left(\sum_i k_i \right) (e_1 \cdot k_{23} e_2 \cdot e_3 + e_2 \cdot k_{31} e_3 \cdot e_1 + e_3 \cdot k_{12} e_1 \cdot e_2 + \frac{\alpha'}{2} e_1 \\
&\quad \cdot k_{23} e_2 \cdot k_{31} e_3 \cdot k_{12}) \text{Tr}(\lambda^{a_1} [\lambda^{a_2}, \lambda^{a_3}])
\end{aligned}$$

$$S = \frac{1}{g_o'^2} \int d^{26}x \left[-\frac{1}{2} \text{Tr}(D_\mu \varphi D^\mu \varphi) + \frac{1}{2\alpha'} \text{Tr}(\varphi^2) + \frac{2^{1/2}}{3\alpha'^{1/2}} \text{Tr}(\varphi^3) - \frac{1}{4} \text{Tr}(F_{\mu\nu} F^{\mu\nu}) \right]$$

$$-\frac{2i\alpha'}{3g_o'^2} \text{Tr}(F_\mu{}^\nu F_\nu{}^\omega F_\omega{}^\mu)$$

$$\lambda^a \rightarrow U \lambda^a U^\dagger,$$

$$\Omega \alpha_n^\mu \Omega^{-1} = (-1)^n \alpha_n^\mu$$

$$\Omega \alpha_n^\mu \Omega^{-1} = \tilde{\alpha}_n^\mu$$

$$\Omega|N;k\rangle=\omega_N|N;k\rangle, \omega_N=(-1)^{1+\alpha' m^2}$$

$$\Omega|N;k;ij\rangle=\omega_N|N;k;ji\rangle.$$

$$\Omega|N;k;a\rangle=\omega_N s^a|N;k;a\rangle$$

$$\alpha' m^2 \text{ even: } \lambda^a \text{ antisymmetric}$$

$$\alpha' m^2 \text{ odd: } \lambda^a \text{ symmetric}$$

$$\Omega_\gamma|N;k;ij\rangle=\omega_N\gamma_{jj'}|N;k;j'i'\rangle\gamma_{i'i}^{-1}$$

$$\Omega_\gamma^2|N;k;ij\rangle=[(\gamma^T)^{-1}\gamma]_{ii'}|N;k;i'j'\rangle(\gamma^{-1}\gamma^T)_{j'j}.$$



$$\gamma^T = \pm \gamma$$

$$|N;k;ij\rangle' = U_{ii'}^{-1}|N;k;i'j'\rangle U_{j'j},$$

$$\gamma'=U^T\gamma U.$$

$$\gamma=M\equiv i\left[\begin{array}{cc}0&I\\-I&0\end{array}\right]$$

$$\begin{array}{l} \alpha' m^2 \text{ even: } M(\lambda^a)^T M = -\lambda^a \\ \alpha' m^2 \text{ odd: } M(\lambda^a)^T M = +\lambda^a \end{array}$$

$$\lambda \rightarrow (\gamma^T)^{-1} \gamma \lambda \gamma^{-1} \gamma^T = \lambda,$$

$$S_{S_2}(k_1;k_2;k_3)=g_{\rm c}^3e^{-2\lambda}\left\langle\prod_{i=1}^3:\tilde{c}ce^{ik_i\cdot X}(z_i,\bar{z}_i):\right\rangle_{S_2}$$

$$S_{S_2}(k_1;k_2;k_3)=ig_{\rm c}^3C_{S_2}(2\pi)^{26}\delta^{26}\left(\sum_i k_i\right)$$

$$S_{S_2}(k_1;k_2;k_3;k_4)=g_{\rm c}^4e^{-2\lambda}\int_{\rm c}d^2z_4\left\langle\prod_{i=1}^3:\tilde{c}ce^{ik_i\cdot X}(z_i,\bar{z}_i)::e^{ik_4\cdot X}(z_4,\bar{z}_4):\right\rangle_{S_2}$$

$$S_{S_2}(k_1;k_2;k_3;k_4)=ig_{\rm c}^4C_{S_2}(2\pi)^{26}\delta^{26}\left(\sum_i k_i\right)J(s,t,u)$$

$$J(s,t,u)=\int_{\rm c}d^2z_4|z_4|^{-\alpha'u/2-4}|1-z_4|^{-\alpha't/2-4}$$

$$\alpha's,\alpha't,\alpha'u=-4,0,4,8,\ldots,$$

$$ig_{\rm c}^4C_{S_2}\int_{|z_4|>1/\epsilon}d^2z_4|z_4|^{\alpha's/2}\sim-\frac{8\pi ig_{\rm c}^4C_{S_2}}{\alpha's+4}$$

$$C_{S_2}=\frac{8\pi}{\alpha'g_{\rm c}^2}$$

$$S_{S_2}(k_1;k_2;k_3)=\frac{8\pi ig_c}{\alpha'}(2\pi)^{26}\delta^{26}\left(\sum_i k_i\right).$$

$$S_{S_2}(k_1;k_2;k_3;k_4)=\frac{8\pi ig_{\rm c}^2}{\alpha'}(2\pi)^{26}\delta^{26}\left(\sum_i k_i\right)C(-\alpha_c(t),-\alpha_c(u)),$$

$$C(a,b)=\int_{\rm c}d^2z|z|^{2a-2}|1-z|^{2b-2}=2\pi\frac{\Gamma(a)\Gamma(b)\Gamma(c)}{\Gamma(a+b)\Gamma(a+c)\Gamma(b+c)},a+b+c=1$$



$$S_{S_2}(k_1;k_2;k_3;k_4)\propto s^{2\alpha_c(t)}\frac{\Gamma(-\alpha_c(t))}{\Gamma(1+\alpha_c(t))},$$

$$S_{S_2}(k_1;k_2;k_3;k_4)\propto \exp\left[-\frac{\alpha'}{2}(\mathrm{sln}\;s\alpha'+\mathrm{tln}\;t\alpha'+\mathrm{uln}\;u\alpha')\right]$$

$$S_{S_2}(k_1,e_1;k_2;k_3)=g_c^2g'_ce^{-2\lambda}e_{1\mu\nu}\langle:\tilde{c}c\partial X^\mu\bar{\partial}X^\nu e^{ik_1\cdot X}(z_1,\bar{z}_1)::\tilde{c}ce^{ik_2\cdot X}(z_2,\bar{z}_2)::\tilde{c}ce^{ik_3\cdot X}(z_3,\bar{z}_3):\rangle_{S_2}$$

$$=-\frac{\pi i\alpha'}{2}g'_ce_{1\mu\nu}k_{23}^\mu k_{23}^\nu (2\pi)^{26}\delta^{26}\left(\sum_i k_i\right)$$

$$g'_c=\frac{2}{\alpha'}g_c$$

$$\mathcal{S}_T=-\frac{1}{2}\int\,d^{26}x(-G)^{1/2}e^{-2\tilde{\Phi}}\left(G^{\mu\nu}\partial_\mu T\partial_\nu T-\frac{4}{\alpha'}T^2\right)$$

$$\tilde{G}_{\mu\nu}=\eta_{\mu\nu}-2\kappa e_{\mu\nu}e^{ik\cdot x}$$

$$\kappa=\pi\alpha'g'_c=2\pi g_c$$

$$S_{S_2}(k_1,e_1;k_2,e_2;k_3,e_3)=\frac{i\kappa}{2}(2\pi)^{26}\delta^{26}\left(\sum_i k_i\right)e_{1\mu\nu}e_{2\alpha\beta}e_{3\gamma\delta}T^{\mu\alpha\gamma}T^{\nu\beta\delta}$$

$$T^{\mu\alpha\gamma}=k_{23}^\mu\eta^{\alpha\gamma}+k_{31}^\alpha\eta^{\gamma\mu}+k_{12}^\gamma\eta^{\mu\alpha}+\frac{\alpha'}{8}k_{23}^\mu k_{31}^\alpha k_{12}^\gamma$$

$$J(s,t,u,\alpha')=-2\sin\pi\alpha_c(t)I(s,t,4\alpha')I(t,u,4\alpha')$$

$$\int_{\mathbb{C}} d^2z z^{a-1+m_1} \bar{z}^{a-1+n_1} (1-z)^{b-1+m_2} (1-\bar{z})^{b-1+n_2}$$

$$=2\sin\left[\pi(b+n_2)\right]B(a+m_1,b+m_2)B(b+n_2,1-a-b-n_1-n_2)m_1-n_1$$

$$\in \mathbf{Z}, m_2-n_2\in \mathbf{Z}$$

$$A_c(s,t,u,\alpha',g_c)=\frac{\pi i g_c^2 \alpha'}{g_o^4} \sin\left[\pi \alpha_c(t)\right] A_o\left(s,t,\frac{1}{4}\alpha',g_o\right) A_o\left(t,u,\frac{1}{4}\alpha',g_o\right)^*$$

$$:e^{ik_1\cdot X(z_1,\bar{z}_1)}::e^{ik_4\cdot X(z_4,\bar{z}_4)}:$$

$$=|z_{14}|^{\alpha'k_1\cdot k_4}:(1+iz_{14}k_1\cdot\partial X+i\bar{z}_{14}k_1\cdot\bar{\partial}X-z_{14}\bar{z}_{14}k_1\cdot\partial Xk_1\cdot\bar{\partial}X$$

$$+\ldots)e^{i(k_1+k_4)\cdot X}(z_4,\bar{z}_4):$$

$$\alpha'k_1\cdot k_4=\frac{\alpha'}{2}(k_1+k_4)^2-4>-2$$

$$-\frac{1}{4g_o'^2}\int\,d^{26}x\mathrm{Tr}(F_{\mu\nu}F^{\mu\nu})$$



$$-\frac{1}{4g_o'^2}\int d^{26}x(-G)^{1/2}e^{-\tilde{\Phi}}\text{Tr}(F_{\mu\nu}F^{\mu\nu})$$

$$-\Lambda\int d^{26}x(-G)^{1/2}e^{-\tilde{\Phi}}$$

Invariancia Möbius.

$$z'=\frac{\alpha z+\beta}{\gamma z+\delta}$$

$$\langle \mathcal{A}_i(z_1,\bar{z}_1)\ldots \mathcal{A}_k(z_n,\bar{z}_n)\rangle_{S_2}=\langle \mathcal{A}'_i(z_1,\bar{z}_1)\ldots \mathcal{A}'_k(z_n,\bar{z}_n)\rangle_{S_2}.$$

$$\langle \mathcal{A}_i(0,0)\rangle_{S_2}=\langle \mathcal{A}'_i(0,0)\rangle_{S_2}=\gamma^{-h_i}\bar{\gamma}^{-\tilde{h}_i}\langle \mathcal{A}_i(0,0)\rangle_{S_2}.$$

$$\langle \mathcal{A}_i(z_1,\bar{z}_1)\mathcal{A}_j(z_2,\bar{z}_2)\rangle_{S_2}=z_{12}^{-h_i-h_j}\bar{z}_{12}^{-\tilde{h}_i-\tilde{h}_j}\langle \mathcal{A}_i(1,1)\mathcal{A}_j(0,0)\rangle_{S_2},$$

$$\langle \mathcal{O}_p(z_1,\bar{z}_1)\mathcal{O}_q(z_2,\bar{z}_2)\rangle_{S_2}=0\text{ unless }h_p=h_q,\tilde{h}_p=\tilde{h}_q.$$

$$\left\langle \prod_{i=1}^3 \mathcal{O}_{p_i}(z_i,\bar{z}_i) \right\rangle_{S_2} = C_{p_1 p_2 p_3} \prod_{\substack{i,j=1 \\ i < j}}^3 z_{ij}^{h-2(h_i+h_j)} \bar{z}_{ij}^{\tilde{h}-2(\tilde{h}_i+\tilde{h}_j)}$$

$$\left\langle \prod_{i=1}^4 \mathcal{O}_{p_i}(z_i,\bar{z}_i) \right\rangle_{S_2} = C_{p_1 p_2 p_3 p_4}(z_c,\bar{z}_c)(z_{12}z_{34})^h(\bar{z}_{12}\bar{z}_{34})^{\tilde{h}} \times \prod_{\substack{i,j=1 \\ i < j}}^4 z_{ij}^{-h_i-h_j}\bar{z}_{ij}^{-\tilde{h}_i-\tilde{h}_j}$$

$$\langle \mathcal{A}'_i(\infty,\infty)\mathcal{A}_j(0,0)\rangle_{S_2}.$$

$$\int [d\phi_{\rm b}]\Psi_{\mathcal{A}_i}[\phi_{\rm b}^{\Omega}]\Psi_{\mathcal{A}_j}[\phi_{\rm b}]$$

$$\langle i\mid j\rangle=\langle \mathcal{A}'_i(\infty,\infty)\mathcal{A}_j(0,0)\rangle_{S_2}.$$

$$\mathcal{O}_i(z,\bar{z})\mathcal{O}_j(0,0)=\frac{\langle i\mid j\rangle}{z^{h_i+h_j}\tilde{z}_i^{\tilde{h}_i}+\tilde{h}_j\langle 1\rangle_{S_2}}+\cdots$$

$$\langle\langle i\mid j\rangle=\pm\langle j\mid i\rangle,$$

$$\langle \mathcal{A}'_i(\infty,\infty)\mathcal{A}_k(1,1)\mathcal{A}_j(0,0)\rangle_{S_2}=\langle i|\hat{\mathcal{A}}_k(1,1)\mid j\rangle,$$

$$\sum_l c^l_{kj}\langle \mathcal{A}'_i(\infty,\infty)\mathcal{A}_l(0,0)\rangle_{S_2}=c_{ikj}$$

$$\langle \mathcal{A}'_i(\infty,\infty)\mathcal{A}_k(z_1,\bar{z}_1)\mathcal{A}_j(0,0)\rangle_{S_2}=z_1^{h_i-h_k-h_j}\bar{z}_1^{\tilde{h}_i-\tilde{h}_k-\tilde{h}_j}c_{ikj}$$

$$\langle \mathcal{A}'_i(\infty,\infty)\mathcal{A}_k(z_1,\bar{z}_1)\mathcal{A}_l(z_2,\bar{z}_2)\mathcal{A}_j(0,0)\rangle_{S_2}=\langle i|\text{T}[\hat{\mathcal{A}}_k(z_1,\bar{z}_1)\hat{\mathcal{A}}_l(z_2,\bar{z}_2)]\mid j\rangle$$



$$1 = |m\rangle \mathcal{G}^{mn} \langle n|.$$

$$\sum_m z_1^{h_i-h_k-h_m} \bar{z}_1^{\tilde{h}_i-\tilde{h}_k-\tilde{h}_m} z_2^{h_m-h_l-h_j} \bar{z}_2^{\tilde{h}_m-\tilde{h}_l-\tilde{h}_j} c_{ikm} c^m_{lj}$$

Cálculos de operador.

$$\left\langle :e^{ik_4\cdot X(\infty,\infty)}:':e^{ik_1\cdot X(z_1,\bar{z}_1)}::e^{ik_2\cdot X(z_2,\bar{z}_2)}::e^{ik_3\cdot X(0,0)}:\right\rangle_{S_2}=\left\langle 0;k_4|T\left[\circ e^{ik_1\cdot X_1}\circ \circ e^{ik_2\cdot X_2}\circ \right]|0;k_3\right\rangle$$

$$e^{ik\cdot X_{\circ}}=e^{ik\cdot X_C}e^{ik\cdot X_A},$$

$$X_C^\mu(z,\bar{z})=x^\mu-i\left(\frac{\alpha'}{2}\right)^{1/2}\sum_{m=1}^\infty\frac{1}{m}(\alpha_{-m}^\mu z^m+\tilde{\alpha}_{-m}^\mu\bar{z}^m)$$

$$X_A^\mu(z,\bar{z})=-i\frac{\alpha'}{2}p^\mu\ln\,|z|^2+i\left(\frac{\alpha'}{2}\right)^{1/2}\sum_{m=1}^\infty\frac{1}{m}\left(\frac{\alpha_m^\mu}{z^m}+\frac{\tilde{\alpha}_m^\mu}{\bar{z}^m}\right).$$

$$\langle 0;k_4|e^{ik_1\cdot X_{1C}}e^{ik_1\cdot X_{1A}}e^{ik_2\cdot X_{2C}}e^{ik_2\cdot X_{2A}}|0;k_3\rangle.$$

$$e^{ik_1\cdot X_{1A}}e^{ik_2\cdot X_{2C}}=e^{ik_2\cdot X_{2C}}e^{ik_1\cdot X_{1A}}e^{-[k_1\cdot X_{1A},k_2\cdot X_{2C}]}=e^{ik_2\cdot X_{2C}}e^{ik_1\cdot X_{1A}}|_{z_{12}}|\alpha'k_1\cdot k_2$$

$$|_{z_{12}}|\alpha'k_1\cdot k_2\langle 0;k_4|e^{ik_1\cdot X_{1C}+ik_2\cdot X_{2C}}e^{ik_1\cdot X_{1A}+ik_2\cdot X_{2A}}|0;k_3\rangle$$

$$=|_{z_{12}}|\alpha'k_1\cdot k_2\left\langle 0;k_4|e^{i(k_1+k_2)\cdot x}e^{\alpha'(k_1\ln|z_1|+k_2\ln|z_2|)\cdot p}\,|\,0;k_3\right\rangle$$

$$=|_{z_{12}}|\alpha'k_1\cdot k_2|_{z_1}|\alpha'k_1\cdot k_3|_{z_2}|\alpha'k_2\cdot k_3\langle 0;k_1+k_2+k_4\,|\,0;k_3\rangle$$

$$=iC_{S_2}^X(2\pi)^d\delta^d\left(\sum_i k_i\right)|_{z_{12}}|\alpha'k_1\cdot k_2|_{z_1}|\alpha'k_1\cdot k_3|_{z_2}|\alpha'k_2\cdot k_3$$

$$\langle 0;k\,|\,0;l\rangle=iC_{S_2}^X(2\pi)^d\delta^d(k+l)$$

$$\langle 0;k\,|\,0;l\rangle=(2\pi)^d\delta^d(l-k)$$

$$\langle\langle 0;k| = iC_{S_2}^X\langle 0;-k|.$$

$$\overline{\mathcal{A}(p)}=\mathcal{A}'(p')^\dagger.$$

$$\mathcal{O}(z)=i^h\sum_{n=-\infty}^{\infty}\frac{\mathcal{O}_n}{z^{n+h}}$$

$$\mathcal{O}(z)^\dagger=i^{-h}\sum_{n=-\infty}^{\infty}\frac{\mathcal{O}_n^\dagger}{\bar{z}^{n+h}}.$$

$$\overline{\mathcal{O}(z)}=i^{-h}(-z^{-2})^h\sum_{n=-\infty}^{\infty}\frac{\mathcal{O}_n^\dagger}{z^{-n-h}}=i^h\sum_{n=-\infty}^{\infty}\frac{\mathcal{O}_{-n}^\dagger}{z^{n+h}}.$$



$$\left\langle \overline{\mathcal{A}_i} \right| = K \langle \mathcal{A}_i |$$

$$\langle \bar{i} \mid j \rangle \equiv \left\langle 1 \mid \overline{\mathcal{A}'_i(\infty, \infty)} \mid j \right\rangle = K \langle 1 \mid \overline{\mathcal{A}'_i(\infty, \infty)} \mid j \rangle = K \langle 1 \mid \mathcal{A}_i(0, 0)^\dagger \mid j \rangle = K \langle j \mid \mathcal{A}_i(0, 0) \mid 1 \rangle^* = K \langle j \mid i \rangle^* = K \langle i \mid j \rangle$$

$$k_1^ik_1^jP_{ij,kl}^Jk_3^k k_3^l$$

$$|z|^{-2a}=\frac{1}{\Gamma(a)}\int_0^\infty dt t^{a-1}\exp\left(-tz\bar{z}\right)$$

$$w\cong w+2\pi\cong w+2\pi\tau$$

$$(\sigma^1,\sigma^2)\cong(\sigma^1+2\pi,\sigma^2)\cong(\sigma^1+2\pi\tau_1,\sigma^2+2\pi\tau_2)$$

$$z\cong z\exp\left(-2\pi i\tau\right)$$

$$1\leq |z|\leq \exp\left(2\pi\tau_2\right)$$

$$0\leq \mathrm{Re} w\leq \pi, w\cong w+2\pi it$$

$$w'=-\bar{w}$$

$$(\sigma^1,0)\cong(\sigma^1,2\pi t)$$

$$w\cong w+2\pi\cong-\bar{w}+2\pi it$$

$$(\sigma^1,\sigma^2)\cong(\sigma^1+2\pi,\sigma^2)\cong(-\sigma^1,\sigma^2+2\pi t)$$

$$w'=-\bar{w}+2\pi it$$

$$0\leq \mathrm{Re} w\leq \pi, w\cong-\bar{w}+\pi+2\pi it$$

$$w'=-\bar{w} \text{ and } w'=w+\pi(2it+1)$$

Correlaciones escalares.

$$\frac{2}{\alpha'}\bar{\partial}\partial G'(w,\bar{w};w',\bar{w}')=-2\pi\delta^2(w-w')+\frac{1}{4\pi\tau_2}.$$

$$G'(w,\bar{w};w',\bar{w}')\sim -\frac{\alpha'}{2}\ln\left|\vartheta_1\left(\frac{w-w'}{2\pi}\middle|\tau\right)\right|^2.$$

$$G'(w,\bar{w};w',\bar{w}')=-\frac{\alpha'}{2}\ln\left|\vartheta_1\left(\frac{w-w'}{2\pi}\middle|\tau\right)\right|^2+\alpha'\frac{[\mathrm{Im}(w-w')]^2}{4\pi\tau_2}+k(\tau,\bar{\tau}).$$



$$\left\langle \prod_{i=1}^n : e^{ik_i \cdot X(z_i, \bar{z}_i)} : \right\rangle_{T^2}$$

$$= iC_{T^2}^X(\tau)(2\pi)^d \delta^d \left(\sum_i k_i \right)$$

$$\times \prod_{i < j} \left| \frac{2\pi}{\partial_v \vartheta_1(0 \mid \tau)} \vartheta_1 \left(\frac{w_{ij}}{2\pi} \mid \tau \right) \exp \left[-\frac{(\text{Im} w_{ij})^2}{4\pi\tau_2} \right] \right|^{\alpha' k_i \cdot k_j}$$

$$Z(\tau) = \text{Tr}[\exp{(2\pi i \tau_1 P - 2\pi \tau_2 H)}] = (q\bar{q})^{-d/24} \text{Tr}(q^{L_0} \bar{q}^{\bar{L}_0})$$

$$V_d(q\bar{q})^{-d/24}\int\frac{d^dk}{(2\pi)^d}\exp{(-\pi\tau_2\alpha'k^2)}\prod_{\mu,n}\sum_{N_{\mu n},\tilde{N}_{\mu n}=0}^{\infty}q^{nN_{\mu n}}\bar{q}^{n\tilde{N}_{\mu n}}$$

$$\sum_{N=0}^{\infty}q^{nN}=(1-q^n)^{-1}$$

$$Z(\tau)=iV_dZ_X(\tau)^d$$

$$Z_X(\tau)=(4\pi^2\alpha'\tau_2)^{-1/2}|\eta(\tau)|^{-2}$$

$$\eta(\tau)=q^{1/24}\prod_{n=1}^{\infty}(1-q^n)$$

$$ds^2 = dw d\bar{w} + \epsilon^* dw^2 + \epsilon d\bar{w}^2 = (1 + \epsilon^* + \epsilon) d[w + \epsilon(\bar{w} - w)] d[\bar{w} + \epsilon^*(w - \bar{w})] + O(\epsilon^2)$$

$$w' \cong w' + 2\pi \cong w' + 2\pi(\tau - 2i\tau_2\epsilon)$$

$$\delta\tau=-2i\tau_2\epsilon$$

$$\delta Z(\tau) \, = -\frac{1}{2\pi} \int \, d^2w [\delta g_{\bar{w}\bar{w}} \langle T_{ww}(w) \rangle + \delta g_{ww} \langle T_{\bar{w}\bar{w}}(\bar{w}) \rangle] = -2\pi i [\delta \tau \langle T_{ww}(0) \rangle - \delta \bar{\tau} \langle T_{\bar{w}\bar{w}}(0) \rangle]$$

$$\partial_w X^\mu(w) \partial_w X_\mu(0) = -\frac{\alpha' d}{2w^2} - \alpha' T_{ww}(0) + O(w)$$

$$Z(\tau)^{-1}\langle\partial_wX^\mu(w)\partial_wX_\mu(0)\rangle=d\partial_w\partial_{w'}G'(w,\bar{w};w',\bar{w}')|_{w'=0}=\frac{\alpha'd}{2}\frac{\vartheta_1\partial_w^2\vartheta_1-\partial_w\vartheta_1\partial_w\vartheta_1}{\vartheta_1^2}+\frac{\alpha'd}{8\pi\tau_2}$$

$$\frac{\alpha'd}{6}\frac{\partial_w^3\vartheta_1}{\partial_w\vartheta_1}+\frac{\alpha'd}{8\pi\tau_2}$$

$$\langle T_{ww}(0) \rangle = \left(-\frac{d}{6} \frac{\partial_w^3 \vartheta_1(0|\tau)}{\partial_w \vartheta_1(0|\tau)} - \frac{d}{8\pi\tau_2} \right) Z(\tau)$$

$$\partial_\tau \ln Z(\tau) = \frac{\pi i d}{3} \frac{\partial_w^3 \vartheta_1(0|\tau)}{\partial_w \vartheta_1(0|\tau)} + \frac{id}{4\tau_2}$$

$$\partial_w^2 \vartheta_1\left(\frac{w}{2\pi} \middle| \tau\right) = \frac{i}{\pi} \partial_\tau \vartheta_1\left(\frac{w}{2\pi} \middle| \tau\right)$$

$$\partial_\tau \ln Z(\tau) = -\frac{d}{3} \partial_\tau \ln \partial_w \vartheta_1(0|\tau) + \frac{id}{4\tau_2}$$

$$Z(\tau) = |\partial_w \vartheta_1(0|\tau)|^{-2d/3} \tau_2^{-d/2}$$

Propagador CFT.

$$\mathrm{Tr}[\exp{(2\pi i \tau_1 P - 2\pi \tau_2 H)}] = (q\bar{q})^{13/12} \mathrm{Tr}(q^{L_0} \bar{q}^{\bar{L}_0}) = 4(q\bar{q})^{1/12} \prod_{n=1}^{\infty} |1 + q^n|^4$$

$$Z(\tau)=\mathrm{Tr}[(-1)^F\exp{(2\pi i \tau_1 P-2\pi \tau_2 H)}]=0$$

$$\langle c(w_1)b(w_2)\tilde{c}(\bar{w}_3)\tilde{b}(\bar{w}_4)\rangle.$$

$$\mathrm{Tr}[(-1)^F c_0 b_0 \tilde{c}_0 \tilde{b}_0 \exp{(2\pi i \tau_1 P - 2\pi \tau_2 H)}]$$

$$(q\bar{q})^{1/12}\prod_{n=1}^{\infty}|1-q^n|^4=|\eta(\tau)|^4$$

$$Z(\tau)=\sum_i q^{h_i-c/24}\bar{q}^{\tilde{h}_i-\tilde{c}/24}(-1)^{F_i}$$

$$h_i-\tilde{h}_i\in\mathbf{Z}$$

$$Z(i\ell) \stackrel{\ell \rightarrow 0}{\approx} \exp\left[\frac{\pi(c+\tilde{c})}{12\ell}\right]$$

$$\vartheta(v,\tau)=\sum_{n=-\infty}^{\infty}\exp{(\pi i n^2\tau+2\pi i n v)}$$

$$\vartheta(v+1,\tau)=\vartheta(v,\tau)\\ \vartheta(v+\tau,\tau)=\exp{(-\pi i \tau-2\pi i v)}\vartheta(v,\tau)$$

$$\vartheta(v,\tau+1) \, = \, \vartheta(v+1/2,\tau) \\ \vartheta(v/\tau,-1/\tau) \, = \, (-i\tau)^{1/2} \exp{(\pi i v^2/\tau)} \vartheta(v,\tau)$$

$$\vartheta(v,\tau)=\prod_{m=1}^{\infty}(1-q^m)(1+zq^{m-1/2})(1+z^{-1}q^{m-1/2})$$



$$q = \exp(2\pi i\tau), z = \exp(2\pi iv)$$

$$\vartheta \begin{bmatrix} a \\ b \end{bmatrix} (v, \tau) = \exp[\pi i a^2 \tau + 2\pi i a(v + b)] \vartheta(v + a\tau + b, \tau) = \sum_{n=-\infty}^{\infty} \exp[\pi i(n + a)^2 \tau + 2\pi i(n + a)(v + b)]$$

$$\vartheta_{00}(v, \tau) = \vartheta_3(v | \tau) = \vartheta \begin{bmatrix} 0 \\ 0 \end{bmatrix} (v, \tau) = \sum_{n=-\infty}^{\infty} q^{n^2/2} z^n$$

$$\vartheta_{01}(v, \tau) = \vartheta_4(v | \tau) = \vartheta \begin{bmatrix} 0 \\ 1/2 \end{bmatrix} (v, \tau) = \sum_{n=-\infty}^{\infty} (-1)^n q^{n^2/2} z^n$$

$$\vartheta_{10}(v, \tau) = \vartheta_2(v | \tau) = \vartheta \begin{bmatrix} 1/2 \\ 0 \end{bmatrix} (v, \tau) = \sum_{n=-\infty}^{\infty} q^{(n-1/2)^2/2} z^{n-1/2}$$

$$\vartheta_{11}(v, \tau) = -\vartheta_1(v | \tau) = \vartheta \begin{bmatrix} 1/2 \\ 1/2 \end{bmatrix} (v, \tau)$$

$$= -i \sum_{n=-\infty}^{\infty} (-1)^n q^{(n-1/2)^2/2} z^{n-1/2}$$

$$\vartheta_{00}(v, \tau) = \prod_{m=1}^{\infty} (1 - q^m)(1 + zq^{m-1/2})(1 + z^{-1}q^{m-1/2})$$

$$\vartheta_{01}(v, \tau) = \prod_{m=1}^{\infty} (1 - q^m)(1 - zq^{m-1/2})(1 - z^{-1}q^{m-1/2})$$

$$\vartheta_{10}(v, \tau) = 2 \exp(\pi i \tau / 4) \cos \pi v \prod_{m=1}^{\infty} (1 - q^m)(1 + zq^m)(1 + z^{-1}q^m)$$

$$\vartheta_{11}(v, \tau) = -2 \exp(\pi i \tau / 4) \sin \pi v \prod_{m=1}^{\infty} (1 - q^m)(1 - zq^m)(1 - z^{-1}q^m)$$

$$\vartheta_{00}(v, \tau + 1) = \vartheta_{01}(v, \tau)$$

$$\vartheta_{01}(v, \tau + 1) = \vartheta_{00}(v, \tau)$$

$$\vartheta_{10}(v, \tau + 1) = \exp(\pi i / 4) \vartheta_{10}(v, \tau)$$

$$\vartheta_{11}(v, \tau + 1) = \exp(\pi i / 4) \vartheta_{11}(v, \tau)$$

$$\vartheta_{00}(v/\tau, -1/\tau) = (-i\tau)^{1/2} \exp(\pi i v^2 / \tau) \vartheta_{00}(v, \tau)$$

$$\vartheta_{01}(v/\tau, -1/\tau) = (-i\tau)^{1/2} \exp(\pi i v^2 / \tau) \vartheta_{10}(v, \tau)$$

$$\vartheta_{10}(v/\tau, -1/\tau) = (-i\tau)^{1/2} \exp(\pi i v^2 / \tau) \vartheta_{01}(v, \tau)$$

$$\vartheta_{11}(v/\tau, -1/\tau) = -i(-i\tau)^{1/2} \exp(\pi i v^2 / \tau) \vartheta_{11}(v, \tau)$$

$$\vartheta_{00}^4(0, \tau) - \vartheta_{01}^4(0, \tau) - \vartheta_{10}^4(0, \tau) = 0$$

$$\vartheta_{11}(0, \tau) = 0$$

$$\eta(\tau) = q^{1/24} \prod_{m=1}^{\infty} (1 - q^m) = \left[\frac{\partial_v \vartheta_{11}(0, \tau)}{-2\pi} \right]^{1/3}$$



$$\begin{aligned}\eta(\tau+1) &= \exp(i\pi/12)\eta(\tau) \\ \eta(-1/\tau) &= (-i\tau)^{1/2}\eta(\tau)\end{aligned}$$

$$S_{T^2}(1;2;\ldots;n)=\frac{1}{2}\int_{F_0}d\tau d\bar\tau\left\langle B\tilde B\tilde c c\mathcal{V}_1(w_1,\bar w_1)\prod_{i=2}^n\int dw_id\bar w_i\mathcal{V}_i(w_i,\bar w_i)\right\rangle_{T^2}$$

$$B=\frac{1}{4\pi}(b,\partial_\tau g)=\frac{1}{2\pi}\int d^2wb_{ww}(w)\partial_\tau g_{\bar w\bar w}=\frac{i}{4\pi\tau_2}\int d^2wb_{ww}(w)\rightarrow 2\pi ib_{ww}$$

$$\int \frac{dw d\bar w}{2(2\pi)^2\tau_2}$$

$$S_{T^2}(1;2;\ldots;n)=\int_{F_0}\frac{d\tau d\bar\tau}{4\tau_2}\left\langle b(0)\tilde b(0)\tilde c(0)c(0)\prod_{i=1}^n\int dw_id\bar w_i\mathcal{V}_i(w_i,\bar w_i)\right\rangle_{T^2}$$

$$Z_{T^2}=\int_{F_0}\frac{d\tau d\bar\tau}{4\tau_2}\langle b(0)\tilde b(0)\tilde c(0)c(0)\rangle_{T^2}$$

$$Z_{T^2}=iV_{26}\int_{F_0}\frac{d\tau d\bar\tau}{4\tau_2}(4\pi^2\alpha'\tau_2)^{-13}|\eta(\tau)|^{-48}$$

$$\frac{d\tau d\bar\tau}{\tau_2^2}$$

$$\begin{aligned}Z_{T^2}&=V_d\int_{F_0}\frac{d\tau d\bar\tau}{4\tau_2}\int\frac{d^dk}{(2\pi)^d}\exp\left(-\pi\tau_2\alpha'k^2\right)\sum_{i\in\mathcal{H}^\perp}q^{h_i-1}\bar q^{\tilde h_i-1}\\&=iV_d\int_{F_0}\frac{d\tau d\bar\tau}{4\tau_2}(4\pi^2\alpha'\tau_2)^{-d/2}\sum_{i\in\mathcal{H}^\perp}q^{h_i-1}\bar q^{\tilde h_i-1}\end{aligned}$$

$$Z_{S_1}(m^2)=V_d\int\frac{d^dk}{(2\pi)^d}\int_0^\infty\frac{dl}{2l}\exp\left[-(k^2+m^2)l/2\right]=iV_d\int_0^\infty\frac{dl}{2l}(2\pi l)^{-d/2}\exp\left(-m^2l/2\right)$$

$$m^2=\frac{2}{\alpha'}(h+\tilde h-2)$$

$$\delta_{h,\tilde h}=\int_{-\pi}^{\pi}\frac{d\theta}{2\pi}\exp\left[i(h-\tilde h)\theta\right]$$

$$\begin{aligned}\sum_{i\in\mathcal{H}^\perp}Z_{S_1}(m_i^2)&=iV_d\int_0^\infty\frac{dl}{2l}\int_{-\pi}^{\pi}\frac{d\theta}{2\pi}(2\pi l)^{-\frac{d}{2}}\times\sum_{i\in\mathcal{H}^\perp}\exp\left[-\frac{(h_i+\tilde h_i-2)l}{\alpha'}+i(h_i-\tilde h_i)\theta\right] \\&=iV_d\int_R\frac{d\tau d\bar\tau}{4\tau_2}(4\pi^2\alpha'\tau_2)^{-d/2}\sum_{i\in\mathcal{H}^\perp}q^{h_i-1}\bar q^{\tilde h_i-1}\end{aligned}$$

$$R\colon \tau_2>0, |\tau_1|<\frac{1}{2}$$

$$F_0\colon |\tau|>1, |\tau_1|<\frac{1}{2}, \tau_2>0$$

$$iV_{26}\int\limits^{\infty}\frac{d\tau_2}{2\tau_2}\frac{d\tau_2}{2\tau_2}(4\pi^2\alpha'\tau_2)^{-13}[\exp{(4\pi\tau_2)}+24^2+\cdots]$$

$$iV_d\int\limits^{\infty}\frac{d\tau_2}{2\tau_2}\frac{d\tau_2}{2\tau_2}(4\pi^2\alpha'\tau_2)^{-d/2}\sum_i\exp{(-\pi\alpha'm_i^2\tau_2)}$$

$$\mathbf{Z}_{\rm vac}(m^2)=\exp\left[Z_{S_1}(m^2)\right]$$

$$\mathbf{Z}_{\rm vac}(m^2)=\langle 0|\exp{(-iHT)}|0\rangle=\exp{(-i\rho_0V_d)}$$

$$\rho_0=\frac{i}{V_d}Z_{S_1}(m^2).$$

$$\int_0^\infty \frac{dl}{2l} \exp{[-(k^2+m^2)l/2]} \rightarrow -\frac{1}{2} \ln{(k^2+m^2)}$$

$$i\int_0^\infty \frac{dl}{2l} \int_{-\infty}^\infty \frac{dk^0}{2\pi} \exp{[-(k^2+m^2)l/2]} \rightarrow \frac{\omega_{\mathbf{k}}}{2}$$

$$\rho_0=\int\frac{d^{d-1}\mathbf{k}}{(2\pi)^{d-1}}\frac{\omega_{\mathbf{k}}}{2}$$

$$\ln \mathbf{Z}_{\rm vac}(m^2)=-\frac{1}{2}{\rm Tr} \ln{(-\partial^2+m^2)}=-\frac{V_d}{2}\int\frac{d^dk}{(2\pi)^d}\ln{(k^2+m^2)}$$

$$\rho_0=\frac{i}{V_d}\sum_i(-1)^{\mathbf{F}_i}Z_{S_1}(m_i^2).$$

$$|\rho_0|\lesssim 10^{-44}\mathrm{GeV}^4$$

$$m_{\rm ew}^4\approx 10^8\mathrm{GeV}^4$$

$$Z_{C_2}=\int_0^\infty\frac{dt}{2t}\mathrm{Tr}'_0[\exp{(-2\pi tL_0)}]=iV_d\int_0^\infty\frac{dt}{2t}(8\pi^2\alpha't)^{-d/2}\sum_{i\in\mathcal{H}^1_0}\exp{[-2\pi t(h_i-1)]}$$

$$\rightarrow iV_{26}n^2\int_0^\infty\frac{dt}{2t}(8\pi^2\alpha't)^{-13}\eta(it)^{-24}$$

$$\eta(it)=t^{-1/2}\eta(i/t)$$

$$Z_{C_2}=i\frac{V_{26}n^2}{2\pi(8\pi^2\alpha')^{13}}\int_0^\infty ds\eta(is/\pi)^{-24}$$



$$\eta(is/\pi)^{-24}=\exp{(2s)}\prod_{n=1}^{\infty}[1-\exp{(-2ns)}]^{-24}=\exp{(2s)}+24+O(\exp{(-2s)})$$

$$\int_0^{\infty}ds\exp{(\beta s)}\equiv-\frac{1}{\beta}$$

$$\left.\frac{1}{k^2}\right|_{k^\mu=0}$$

$$-\frac{1}{g^2}\int\,d^dx\left(\frac{1}{2}\partial_\mu\phi\partial^\mu\phi+g\Lambda\phi\right)$$

$$\partial^2\phi=g\Lambda$$

$$-\Lambda\int\,d^{26}x(-G)^{1/2}e^{-\tilde{\Phi}}$$

$$\exp{(-\alpha'k^2s/2)}$$

$$\langle B|c_0b_0\exp\left[-s(L_0+\tilde{L}_0)\right]|B\rangle.$$

$$(\alpha_n^\mu+\tilde{\alpha}_{-n}^\mu)|B\rangle=(c_n+\tilde{c}_{-n})|B\rangle=(b_n-\tilde{b}_{-n})|B\rangle=0,\,\,\text{all}\,\,n$$

$$|B\rangle\propto (c_0+\tilde{c}_0)\exp\left[-\sum_{n=1}^{\infty}\left(n^{-1}\alpha_{-n}\cdot\tilde{\alpha}_{-n}+b_{-n}\tilde{c}_{-n}+\tilde{b}_{-n}c_{-n}\right)\right]|0;0\rangle.$$

$$Z_{K_2}=\int_0^{\infty}\frac{dt}{4t}\mathrm{Tr}'_{\mathrm{c}}\{\Omega\exp\left[-2\pi t(L_0+\tilde{L}_0)\right]\}=iV_d\int_0^{\infty}\frac{dt}{4t}(4\pi^2\alpha't)^{-d/2}\sum_{i\in\mathcal{H}_c^{\perp}}\Omega_i\exp\left[-2\pi t(h_i+\tilde{h}_i-2)\right]$$

$$Z_{K_2}\rightarrow iV_{26}\int_0^{\infty}\frac{dt}{4t}(4\pi^2\alpha't)^{-13}\eta(2it)^{-24}$$

$$\begin{array}{l} 0\leq\sigma^1\leq2\pi,0\leq\sigma^2\leq2\pi t\\ 0\leq\sigma^1\leq\pi,0\leq\sigma^2\leq4\pi t \end{array}$$

$$w\cong w+2\pi\cong-\bar{w}+2\pi it$$

$$w\cong w+4\pi it, w+\pi\cong-(\bar{w}+\pi)+2\pi it$$

$$Z_{K_2}=i\frac{2^{26}V_{26}}{4\pi(8\pi^2\alpha')^{13}}\int_0^{\infty}ds\eta(is/\pi)^{-24}$$

$$Z_{M_2}=iV_d\int_0^{\infty}\frac{dt}{4t}(8\pi^2\alpha't)^{-d/2}\sum_{i\in\mathcal{H}_0^{\perp}}\Omega_i\exp\left[-2\pi t(h_i-1)\right]$$

$$\exp{(2\pi t)}\prod_{n=1}^{\infty}\left[1-(-1)^n\exp{(-2\pi nt)}\right]^{-24}=\vartheta_{00}(0,2it)^{-12}\eta(2it)^{-12}$$



$$Z_{M_2}=\pm inV_{26}\int_0^\infty\frac{dt}{4t}(8\pi^2\alpha't)^{-13}\vartheta_{00}(0,2it)^{-12}\eta(2it)^{-12}$$

$$Z_{M_2}=\pm 2in\frac{2^{13}V_{26}}{4\pi(8\pi^2\alpha')^{13}}\int_0^\infty ds\vartheta_{00}(0,2is/\pi)^{-12}\eta(2is/\pi)^{-12}$$

$$i\frac{24V_{26}}{4\pi(8\pi^2\alpha')^{13}}(2^{13}\mp n)^2\int_0^\infty ds$$

$$\mathrm{Tr}(\lambda^a\lambda^b)\mathrm{Tr}(\lambda^c\lambda^d)=\mathrm{Tr}(\lambda^a\lambda^b\lambda^e\lambda^c\lambda^d\lambda^e).$$

$$x^4\cong x^4+2\pi R$$

$$ds^2=G^D_{MN}dx^Mdx^N=G_{\mu\nu}dx^\mu dx^\nu+G_{dd}\big(dx^d+A_\mu dx^\mu\big)^2$$

$$x'^d=x^d+\lambda(x^\mu).$$

$$A'_\mu=A_\mu-\partial_\mu\lambda$$

$$\phi(x^M)=\sum_{n=-\infty}^{\infty}\phi_n(x^\mu)\mathrm{exp}\left(inx^d/R\right)$$

$$\partial_\mu\partial^\mu\phi_n(x^\mu)=\frac{n^2}{R^2}\phi_n(x^\mu)$$

$$-p^\mu p_\mu=\frac{n^2}{R^2}$$

$$\boldsymbol{R}=\boldsymbol{R}_d-2e^{-\sigma}\nabla^2e^{\sigma}-\frac{1}{4}e^{2\sigma}F_{\mu\nu}F^{\mu\nu}$$

$$S_1=\frac{1}{2\kappa_0^2}\int\,d^Dx(-G_D)^{1/2}e^{-2\Phi}(\boldsymbol{R}+4\nabla_\mu\Phi\nabla^\mu\Phi)$$

$$=\frac{\pi R}{\kappa_0^2}\int\,d^dx(-G_d)^{1/2}e^{-2\Phi+\sigma}\times\left(\boldsymbol{R}_d-4\partial_\mu\Phi\partial^\mu\sigma+4\partial_\mu\Phi\partial^\mu\Phi-\frac{1}{4}e^{2\sigma}F_{\mu\nu}F^{\mu\nu}\right)$$

$$=\frac{\pi R}{\kappa_0^2}\int\,d^dx(-G_d)^{1/2}e^{-2\Phi_d}\times\left(\boldsymbol{R}_d-\partial_\mu\sigma\partial^\mu\sigma+4\partial_\mu\Phi_d\partial^\mu\Phi_d-\frac{1}{4}e^{2\sigma}F_{\mu\nu}F^{\mu\nu}\right)$$

$$\partial_\mu + ip_d A_\mu = \partial_\mu + in \tilde{A}_\mu,$$

$$g_d^2=\frac{\kappa_0^2e^{2\Phi_d}}{\pi R^3e^{2\sigma}}=\frac{2\kappa_d^2}{\rho^2}$$

$$\frac{1}{\kappa_d^2}=\frac{2\pi\rho}{\kappa^2}$$

$$\begin{aligned}
S_2 &= -\frac{1}{24\kappa_0^2} \int d^D x (-G_D)^{1/2} e^{-2\Phi} H_{MNL} H^{MNL} \\
&= -\frac{\pi R}{12\kappa_0^2} \int d^d x (-G_d)^{1/2} e^{-2\Phi_d} (\tilde{H}_{\mu\nu\lambda} \tilde{H}^{\mu\nu\lambda} + 3e^{-2\sigma} H_{d\mu\nu} H_d^{\mu\nu}) \\
\tilde{H}_{\mu\nu\lambda} &= (\partial_\mu B_{\nu\lambda} - A_\mu H_{d\nu\lambda}) + \mathfrak{C} p_{\text{cyclic permutations}} \\
B'_{v\lambda} &= B_{v\lambda} - \lambda H_{dv\lambda} \\
X &\cong X + 2\pi R \\
k &= \frac{n}{R}, n \in \mathbf{Z} \\
X(\sigma + 2\pi) &= X(\sigma) + 2\pi R w, w \in \mathbf{Z}
\end{aligned}$$

$$\partial X(z) = -i \left(\frac{\alpha'}{2} \right)^{1/2} \sum_{m=-\infty}^{\infty} \frac{\alpha_m}{z^{m+1}}, \bar{\partial} X(\bar{z}) = -i \left(\frac{\alpha'}{2} \right)^{1/2} \sum_{m=-\infty}^{\infty} \frac{\tilde{\alpha}_m}{\bar{z}^{m+1}}.$$

$$2\pi R w = \oint (dz \partial X + d\bar{z} \bar{\partial} X) = 2\pi (\alpha'/2)^{1/2} (\alpha_0 - \tilde{\alpha}_0)$$

$$p = \frac{1}{2\pi\alpha'} \oint (dz \partial X - d\bar{z} \bar{\partial} X) = (2\alpha')^{-1/2} (\alpha_0 + \tilde{\alpha}_0)$$

$$\begin{aligned}
p_L &\equiv (2/\alpha')^{1/2} \alpha_0 = \frac{n}{R} + \frac{wR}{\alpha'} \\
p_R &\equiv (2/\alpha')^{1/2} \tilde{\alpha}_0 = \frac{n}{R} - \frac{wR}{\alpha'}
\end{aligned}$$

$$\begin{aligned}
L_0 &= \frac{\alpha' p_L^2}{4} + \sum_{n=1}^{\infty} \alpha_{-n} \alpha_n, \\
\tilde{L}_0 &= \frac{\alpha' p_R^2}{4} + \sum_{n=1}^{\infty} \tilde{\alpha}_{-n} \tilde{\alpha}_n.
\end{aligned}$$

Función de Partición.

$$\begin{aligned}
(q\bar{q})^{-1/24} \text{Tr}(q^{L_0} \bar{q}^{\tilde{L}_0}) &= |\eta(\tau)|^{-2} \sum_{n,w=-\infty}^{\infty} q^{\alpha' p_L^2/4} \bar{q}^{\alpha' p_R^2/4} \\
&= |\eta(\tau)|^{-2} \sum_{n,w=-\infty}^{\infty} \exp \left[-\pi \tau_2 \left(\frac{\alpha' n^2}{R^2} + \frac{w^2 R^2}{\alpha'} \right) + 2\pi i \tau_1 n w \right] \\
\sum_{n=-\infty}^{\infty} \exp(-\pi a n^2 + 2\pi i b n) &= a^{-1/2} \sum_{m=-\infty}^{\infty} \exp \left[-\frac{\pi(m-b)^2}{a} \right].
\end{aligned}$$



$$2\pi R Z_X(\tau) \sum_{m,w=-\infty}^{\infty} \exp\left(-\frac{\pi R^2 |m - w\tau|^2}{\alpha' \tau_2}\right)$$

$$\begin{aligned} X(\sigma^1 + 2\pi, \sigma^2) &= X(\sigma^1, \sigma^2) + 2\pi w R \\ X(\sigma^1 + 2\pi\tau_1, \sigma^2 + 2\pi\tau_2) &= X(\sigma^1, \sigma^2) + 2\pi m R \end{aligned}$$

$$X_{\text{cl}} = \sigma^1 w R + \sigma^2 (m - w\tau_1) R / \tau_2$$

Operadores de vórtice.

$$[x_L, p_L] = [x_R, p_R] = i$$

$$X(z, \bar{z}) = X_L(z) + X_R(\bar{z})$$

$$\begin{aligned} X_L(z) &= x_L - i \frac{\alpha'}{2} p_L \ln z + i \left(\frac{\alpha'}{2} \right)^{1/2} \sum_{\substack{m=-\infty \\ m \neq 0}}^{\infty} \frac{\alpha_m}{m z^m}, \\ X_R(\bar{z}) &= x_R - i \frac{\alpha'}{2} p_R \ln \bar{z} + i \left(\frac{\alpha'}{2} \right)^{1/2} \sum_{\substack{m=-\infty \\ m \neq 0}}^{\infty} \frac{\tilde{\alpha}_m}{m \bar{z}^m}. \end{aligned}$$

$$\begin{aligned} X_L(z_1) X_L(z_2) &\sim -\frac{\alpha'}{2} \ln z_{12}, \quad X_R(\bar{z}_1) X_R(\bar{z}_2) \sim -\frac{\alpha'}{2} \ln \bar{z}_{12} \\ X_L(z_1) X_R(\bar{z}_2) &\sim 0 \end{aligned}$$

$$\mathcal{V}_{k_L k_R}(z, \bar{z}) =: e^{ik_L X_L(z) + ik_R X_R(\bar{z})};$$

$$\mathcal{V}_{k_L k_R}(z_1, \bar{z}_1) \mathcal{V}_{k'_L k'_R}(z_2, \bar{z}_2) \sim z_{12}^{\alpha' k'_L / 2} \bar{z}_{12}^{\alpha' k_R k'_R / 2} \mathcal{V}_{(k+k')_L (k+k')_R}(z_2, \bar{z}_2).$$

$$\exp [\pi i \alpha' (k_L k'_L - k_R k'_R)] = \exp [2\pi i (n w' + w n')] = 1.$$

$$[X_L(z_1), X_L(z_2)] = \frac{\pi i \alpha'}{2} \text{sign}(\sigma_1^1 - \sigma_2^1).$$

$$\mathcal{V}_{k_L k_R}(z, \bar{z}) = \exp [\pi i (k_L - k_R)(p_L + p_R) \alpha' / 4] \circ e^{ik_L X_L(z) + ik_R X_R(\bar{z})} \circ$$

$$\begin{aligned} \exp \{ \pi i [(k_L - k_R)(k'_L + k'_R) - (k'_L - k'_R)(k_L + k_R)] \alpha' / 4 \} \\ = \exp [\pi i (n w' - w n')] \end{aligned}$$

$$\prod_{\substack{i,j=1 \\ i < j}}^n |z_{ij}|^{\alpha' k_i k_j} \rightarrow \prod_{\substack{i,j=1 \\ i < j}}^n z_{ij}^{\alpha' k_L i k_L j / 2} \bar{z}_{ij}^{\alpha' k_R i k_R j / 2}$$

$$2\pi R \delta_{\Sigma_i n_i, 0} \delta_{\Sigma_i w_i, 0}$$



Operadores DDF.

$$X^+(z)X^-(0) \sim \frac{\alpha'}{2} \ln z, X^+(z)X^+(0) \sim X^-(z)X^-(0) \sim 0.$$

$$V^i(nk_0, z) = \partial X^i(z) e^{ink_0 X^+(z)} (2/\alpha')^{1/2}$$

$$V^i(nk_0, z) V^j(mk_0, 0) \sim -\frac{\delta^{ij}}{z^2} e^{i(n+m)k_0 X^+(0)} - \frac{ink_0 \delta^{ij}}{z} \partial X^+(0) e^{i(n+m)k_0 X^+(0)}$$

$$A_n^i = \oint \frac{dz}{2\pi} V^i(nk_0, z)$$

$$[A_m^i, A_n^j] = m \delta^{ij} \delta_{m,-n} \frac{\alpha' k_0 p^+}{2}$$

$$m^2 = -k^\mu k_\mu = (k_L^{25})^2 + \frac{4}{\alpha'} (N-1) = (k_R^{25})^2 + \frac{4}{\alpha'} (\tilde{N}-1)$$

$$m^2 = \frac{n^2}{R^2} + \frac{w^2 R^2}{\alpha'^2} + \frac{2}{\alpha'} (N + \tilde{N} - 2)$$

$$0 = nw + N - \tilde{N}$$

$$\alpha_{-1}^\mu \tilde{\alpha}_{-1}^\nu |0; k\rangle, (\alpha_{-1}^\mu \tilde{\alpha}_{-1}^{25} + \alpha_{-1}^{25} \tilde{\alpha}_{-1}^\mu) |0; k\rangle (\alpha_{-1}^\mu \tilde{\alpha}_{-1}^{25} - \alpha_{-1}^{25} \tilde{\alpha}_{-1}^\mu) |0; k\rangle, \alpha_{-1}^{25} \tilde{\alpha}_{-1}^{25} |0; k\rangle$$

$$\partial X^\mu \bar{\partial} X^{25} - \partial X^{25} \bar{\partial} X^\mu = \bar{\partial} (X^{25} \partial X^\mu) - \partial (X^{25} \bar{\partial} X^\mu)$$

$$\frac{2^{1/2} g_{c,25}}{\alpha'} : (\partial X^\mu \bar{\partial} X^{25} \pm \partial X^{25} \bar{\partial} X^\mu) e^{ik \cdot X} :$$

$$g_{c,25} : e^{ik_L \cdot X_L(z) + ik_R \cdot X_R(\bar{z})} :$$

$$-2^{-1/2} \pi i g_{c,25} (2\pi)^{25} \delta^{25} \left(\sum_i k_i \right) k_{23}^\mu (k_{L23}^{25} \pm k_{R23}^{25})$$

$$\rightarrow -2^{3/2} \pi i g_{c,25} (2\pi)^{25} \delta^{25} \left(\sum_i k_i \right) k_2^\mu (k_{L2}^{25} \pm k_{R2}^{25})$$

Simetrías de gauge.

$$(n+w)^2 + 4N = (n-w)^2 + 4\tilde{N} = 4$$

$$n = w = \pm 1, N = 0, \tilde{N} = 1, n = -w = \pm 1, N = 1, \tilde{N} = 0$$

$$n = \pm 2, w = N = \tilde{N} = 0, w = \pm 2, n = N = \tilde{N} = 0$$

$$: \bar{\partial} X^\mu e^{ik \cdot X} \exp [\pm 2i \alpha'^{-1/2} X_L^{25}] : , : \partial X^\mu e^{ik \cdot X} \exp [\pm 2i \alpha'^{-1/2} X_R^{25}] : .$$



$$\begin{aligned}j^1(z) &=: \cos \left[2\alpha'^{-1/2} X_L^{25}(z) \right]: \\j^2(z) &=: \sin \left[2\alpha'^{-1/2} X_L^{25}(z) \right]: \\j^3(z) &= i\partial X_L^{25}(z)/\alpha'^{1/2}\end{aligned}$$

$$j^i(z)j^j(0)\sim \frac{\delta^{ij}}{2z^2}+i\frac{\epsilon^{ijk}}{z}j^k(0)$$

$$\begin{aligned}j^i(z) &= \sum_{m=-\infty}^{\infty} \frac{j_m^i}{z^{m+1}} \\[j_m^i,j_n^j] &= \frac{m}{2} \delta_{m,-n} \delta^{ij} + i \epsilon^{ijk} j_{m+n}^k\end{aligned}$$

Escalares y acoplamientos.

$$g_{25}^2=2\kappa_{25}^2/\alpha'$$

$$g_4^2=2\kappa_4^2/\alpha'$$

$$g_{\rm G,4}^2(E)=\kappa_4^2E^2$$

$$g_5^2=2\pi\rho_5g_4^2$$

$$\hat{g}_5^2=g_5^2E=2\pi\rho_5Eg_4^2$$

Mecanismo de Higgs.

$$m=\frac{|R^2-\alpha'|}{R\alpha'}\approx\frac{2}{\alpha'}|R-\alpha'^{1/2}|,$$

$$:j^i(z)\tilde{j}^j(\bar{z})e^{ik\cdot X(z,\bar{z})}:\,.$$

$$U(M) \propto \epsilon^{ijk} \epsilon^{i'j'k'} M_{ii'} M_{jj'} M_{kk'} = \det M$$

$$U(M)=\frac{\partial U(M)}{\partial M_{ij}}=0$$

$$M_{11}M_{22}M_{33}=M_{11}M_{22}=M_{11}M_{33}=M_{22}M_{33}=0$$

$$m^2=\frac{n^2}{R^2}+\frac{w^2R^2}{\alpha'^2}+\frac{2}{\alpha'}(N+\tilde{N}-2),$$

$$R\rightarrow R'=\frac{\alpha'}{R}, n\leftrightarrow w.$$

$$p_L^{25}\rightarrow p_L^{25}, p_R^{25}\rightarrow -p_R^{25}$$

$$X'^{25}(z,\bar{z})=X_L^{25}(z)-X_R^{25}(\bar{z})$$

$$R_{\rm self\text{-}dual}=R_{SU(2)\times SU(2)}=\alpha'^{1/2}$$



$$\rho' = \frac{\alpha'}{\rho}, \kappa' = \frac{\alpha'^{1/2}}{\rho} \kappa$$

$$e^{\Phi'} = \frac{\alpha'^{1/2}}{\rho} e^{\Phi}$$

Dimensiones múltiples.

$$X^m \cong X^m + 2\pi R, 26 - k \leq m \leq 25$$

$$\begin{aligned} S = & \frac{(2\pi R)^k}{2\kappa_0^2} \int d^d x (-G_d)^{1/2} e^{-2\Phi_d} [\mathbf{R}_d + 4\partial_\mu \Phi_d \partial^\mu \Phi_d \\ & - \frac{1}{4} G^{mn} G^{pq} (\partial_\mu G_{mp} \partial^\mu G_{nq} + \partial_\mu B_{mp} \partial^\mu B_{nq}) - \frac{1}{4} G_{mn} F_{\mu\nu}^m F^{n\mu\nu} - \frac{1}{4} G^{mn} H_{m\mu\nu} H_n^{\mu\nu} \\ & - \frac{1}{12} H_{\mu\nu\lambda} H^{\mu\nu\lambda}] \end{aligned}$$

$$B_{mn}\partial_a(g^{1/2}\epsilon^{ab}X^m\partial_bX^n)$$

$$X^m(\sigma^1,\sigma^2)=x^m(\sigma^2)+w^mR\sigma^1$$

$$L=\frac{1}{2\alpha'}G_{mn}(\dot{x}^m\dot{x}^n+w^mw^nR^2)-\frac{i}{\alpha'}B_{mn}\dot{x}^mw^nR$$

$$p_m=-\frac{\partial L}{\partial v^m}=\frac{1}{\alpha'}(G_{mn}v^n+B_{mn}w^nR),$$

$$v_m=\alpha'\frac{n_m}{R}-B_{mn}w^nR$$

$$\frac{1}{2\alpha'}G_{mn}(v^mv^n+w^mw^nR^2)$$

$$\begin{aligned} m^2 = & \frac{1}{2\alpha'^2}G_{mn}(v_L^mv_L^n+v_R^mv_R^n)+\frac{2}{\alpha'}(N+\tilde{N}-2) \\ & v_{L,R}^m=v^m\pm w^mR \end{aligned}$$

$$\begin{aligned} 0 & = G_{mn}(v_L^mv_L^n-v_R^mv_R^n)+4\alpha'(N-\tilde{N}) \\ & = 4\alpha'(n_mw^m+N-\tilde{N})(8.4.10) \end{aligned}$$

$$X^m=(w_1^m\sigma^1+w_2^m\sigma^2)R$$

$$2\pi i b_{mn}w_1^mw_2^n$$

$$G_{mn}=e^r_me^r_n$$

$$k_{rL}=e_r^m\frac{v_{mL}}{\alpha'}, k_{rR}=e_r^m\frac{v_{mR}}{\alpha'}$$



$$m^2 = \frac{1}{2}(k_{rL}k_{rL} + k_{rR}k_{rR}) + \frac{2}{\alpha'}(N + \tilde{N} - 2)$$

$$0 = \alpha'(k_{rL}k_{rL} - k_{rR}k_{rR}) + 4(N - \tilde{N})$$

Compactaci3n de Narain.

$$:e^{ik_L\cdot X_L(z)+ik_R\cdot X_R(\bar{z})}: :e^{ik'_L\cdot X_L(0)+ik'_R\cdot X_R(0)}: \sim z^{l_L\cdot l'_L}\bar{z}^{l_R\cdot l'_R}\cdot l'_R:e^{i(k_L+k'_L)\cdot X_L(0)+i(k_R+k'_R)\cdot X_R(0)}:$$

$$l_L\cdot l'_L-l_R\cdot l'_R\equiv l\circ l'\in\mathbf{Z}$$

$$\Gamma\subset\Gamma^*$$

$$l\circ l\in 2\mathbf{Z} \text{ for all } l\in \Gamma.$$

$$2l\circ l'=(l+l')\circ(l+l')-l\circ l-l'\circ l'\in 2\mathbf{Z}$$

$$Z_\Gamma(\tau)=|\eta(\tau)|^{-2k}\sum_{l\in\Gamma}\exp\left(\pi i\tau l_L^2-\pi i\bar\tau l_R^2\right)$$

$$\sum_{l'\in\Gamma}\delta(l-l')=V_\Gamma^{-1}\sum_{l''\in\Gamma^*}\exp\left(2\pi il''\circ l\right).$$

$$Z_\Gamma(\tau)=V_\Gamma^{-1}|\eta(\tau)|^{-2k}\sum_{l''\in\Gamma^*}\int d^{2k}l\exp\left(2\pi il''\circ l+\pi i\tau l_L^2-\pi i\bar\tau l_R^2\right)$$

$$=V_\Gamma^{-1}(\tau\bar\tau)^{-k/2}|\eta(\tau)|^{-2k}\sum_{l''\in\Gamma^*}\exp\left(-\pi il_L'^{''2}/\tau+\pi il_R'^{''2}/\bar\tau\right)=V_\Gamma^{-1}Z_{\Gamma^*}(-1/\tau)$$

$$\Gamma=\Gamma^*$$

$$\Gamma'=\Lambda\Gamma$$

$$l_{L,R}=\frac{n}{r}\pm\frac{mr}{2}$$

$$l'_L=l_L\cosh\lambda+l_R\sinh\lambda,l'_R=l_L\sinh\lambda+l_R\cosh\lambda$$

$$\frac{O(k,k,\mathbf{R})}{O(k,\mathbf{R})\times O(k,\mathbf{R})}$$

$$\frac{2k(2k-1)}{2}-k(k-1)=k^2$$

$$\Lambda\Gamma_0\cong\Lambda'\Lambda\Lambda''\Gamma_0$$

$$\Lambda'\in O(k,\mathbf{R})\times O(k,\mathbf{R}),\Lambda''\in O(k,k,\mathbf{Z})$$

$$\frac{O(k,k,\mathbf{R})}{O(k,\mathbf{R})\times O(k,\mathbf{R})\times O(k,k,\mathbf{Z})}$$

$$x'^m=L^m{}_n x^n$$

$$b_{mn}\rightarrow b_{mn}+N_{mn}$$



$$-\frac{1}{2}g_{ij}(\phi)\partial_{\mu}\phi^i\partial^{\mu}\phi^j$$

$$\frac{16}{d-2}\partial_{\mu}\Phi\partial^{\mu}\Phi+G^{mn}G^{pq}(\partial_{\mu}G_{mp}\partial^{\mu}G_{nq}+\partial_{\mu}B_{mp}\partial^{\mu}B_{nq})$$

$$\rho=\frac{R^2}{\alpha'}\big(B_{24,25}+i\det^{1/2}G_{mn}\big)=b_{24,25}+i\frac{V}{4\pi^2\alpha'}$$

$$ds^2=\frac{\alpha'\rho_2}{R^2\tau_2}\big|dX^{24}+\tau dX^{25}\big|^2$$

$$\mathrm{PSL}(2,\mathbf{Z})\times\mathrm{PSL}(2,\mathbf{Z})\times\mathbf{Z}_2^2$$

$$\frac{\partial_{\mu}\tau\partial^{\mu}\bar{\tau}}{\tau_2^2}+\frac{\partial_{\mu}\rho\partial^{\mu}\bar{\rho}}{\rho_2^2}$$

$$\frac{PSL(2,\mathbf{R})\times PSL(2,\mathbf{R})}{U(1)\times U(1)\times PSL(2,\mathbf{Z})\times PSL(2,\mathbf{Z})\times \mathbf{Z}_2^2}.$$

Orbifolds.

$$X^{25}\cong -X^{25}$$

$$X^m\rightarrow -X^m, 26-k\leq m\leq 25.$$

$$\begin{array}{l} t^m\colon X^{25}\cong X^{25}+2\pi Rm,\\ t^mr\colon X^{25}\cong 2\pi Rm-X^{25}, \end{array}$$

$$X^{25}(\sigma^1+2\pi)=-X^{25}(\sigma^1)$$

$$|N,\tilde{N};k^\mu,n,w\rangle\rightarrow (-1)^{\sum_{\tilde{m}=1}^\infty (N_m^{25}+\tilde{N}_{\tilde{m}}^{25})}|N,\tilde{N};k^\mu,-n,-w\rangle,$$

$$X^{25}(z,\bar{z})=i\left(\frac{\alpha'}{2}\right)^{1/2}\sum_{m=-\infty}^{\infty}\frac{1}{m+1/2}\left(\frac{\alpha_{m+1/2}^{25}}{z^{m+1/2}}+\frac{\tilde{\alpha}_{m+1/2}^{25}}{\bar{z}^{m+1/2}}\right).$$

$$X^{25}(\sigma^1+2\pi)=2\pi R-X^{25}(\sigma^1).$$

$$m^2=\frac{4}{\alpha'}\Big(N-\frac{15}{16}\Big), N=\tilde{N}$$

$$(q\bar{q})^{-1/24}\mathrm{Tr}_U\left(\frac{1+r}{2}q^{L_0}\bar{q}^{\bar{L}_0}\right)$$

$$\frac{1}{2}Z_{\mathrm{tor}}(R,\tau)+\frac{1}{2}(q\bar{q})^{-1/24}\prod_{m=1}^{\infty}|1+q^m|^{-2}$$

$$(q\bar{q})^{1/48}\mathrm{Tr}_T\left(\frac{1+r}{2}q^{L_0}\bar{q}^{\bar{L}_0}\right)=(q\bar{q})^{1/48}\left[\prod_{l=1}^{\infty}|1-q^{m-1/2}|^{-2}+\prod_{m=1}^{\infty}|1+q^{m-1/2}|^{-2}\right]$$



$$Z_{\text{orb}}(R,\tau)=\frac{1}{2}Z_{\text{tor}}(R,\tau)+\left|\frac{\eta(\tau)}{\vartheta_{10}(0,\tau)}\right|+\left|\frac{\eta(\tau)}{\vartheta_{01}(0,\tau)}\right|+\left|\frac{\eta(\tau)}{\vartheta_{00}(0,\tau)}\right|$$

$$\begin{aligned} X^{25}(\sigma^1+2\pi,\sigma^2)&=(-1)^{a+1}X^{25}(\sigma^1,\sigma^2)\\ X^{25}(\sigma^1+2\pi\tau_1,\sigma^2+2\pi\tau_2)&=(-1)^{b+1}X^{25}(\sigma^1,\sigma^2) \end{aligned}$$

Twisting.

$$\phi(\sigma^1+2\pi)=h\cdot\phi(\sigma^1)$$

$$P_H=\frac{1}{\text{order}(H)}\sum_{h_2\in H}\hat{h}_2$$

$$Z=\frac{1}{\text{order}(H)}\sum_{h_1,h_2\in H}Z_{h_1,h_2}$$

$$\phi'(\sigma^1)=h_2\cdot\phi(\sigma^1)$$

$$\phi'(\sigma^1+2\pi)=h_1'\cdot\phi'(\sigma^1)$$

$$h_1'=h_2h_1h_2^{-1}$$

$$c=1CFTs$$

$$r'\colon X^{25}\rightarrow X^{25}+\pi\alpha^{1/2}$$

$$Z_{\text{orb}}\left(\alpha'^{1/2},\tau\right)=Z_{\text{tor}}\left(2\alpha'^{1/2},\tau\right)$$

$$X^{25}\rightarrow X^{25}+\frac{2\pi\alpha'^{1/2}}{k}$$

$$j^3\tilde{j}^3,j^1\tilde{j}^1+j^2\tilde{j}^2,j^1\tilde{j}^2-j^2\tilde{j}^1$$

$$A_{25}(x^M)=-\frac{\theta}{2\pi R}=-i\Lambda^{-1}\frac{\partial\Lambda}{\partial x^{25}},\Lambda(x^{25})=\exp\left(-\frac{i\theta x^{25}}{2\pi R}\right)$$

$$W_q=\exp\left(iq\oint\,dx^{25}A_{25}\right)=\exp\left(-iq\theta\right)$$

$$S=\int\,d\tau\left(\frac{1}{2}\dot{X}^M\dot{X}_M+\frac{m^2}{2}-iqA_M\dot{X}^M\right)$$

$$p_{25}=-\frac{\partial L}{\partial v^{25}}=v^{25}-\frac{q\theta}{2\pi R}$$

$$v_{25}=\frac{2\pi l+q\theta}{2\pi R}$$

$$H=\frac{1}{2}(p_\mu p^\mu+v_{25}^2+m^2)$$



$$v_{25} = p_{25} = \frac{2\pi l + q\theta}{2\pi R}$$

$$A_{25} = -\frac{1}{2\pi R} \text{diag}(\theta_1, \theta_2, \dots, \theta_n)$$

$$v_{25} = \frac{2\pi l - \theta_j + \theta_i}{2\pi R}$$

$$m^2 = \frac{(2\pi l - \theta_j + \theta_i)^2}{4\pi^2 R^2} + \frac{1}{\alpha'} (N - 1)$$

$$m^2 = \frac{(\theta_j - \theta_i)^2}{4\pi^2 R^2}$$

$$U(r_1) \times \dots \times U(r_s), \sum_{i=1}^s r_i = n$$

$$\text{Tr}([A_m, A_n]^2)$$

Dualidad T.

$$X'^{25}(z, \bar{z}) = X_L^{25}(z) - X_R^{25}(\bar{z}).$$

$$\partial_n X^{25} = -i\partial_t X'^{25}$$

$$X'^{25}(\pi) - X'^{25}(0) = \int_0^\pi d\sigma^1 \partial_1 X'^{25} = -i \int_0^\pi d\sigma^1 \partial_2 X^{25} = -2\pi\alpha' v^{25} = -\frac{2\pi\alpha' l}{R} = -2\pi l R'$$

$$\Delta X'^{25} = X'^{25}(\pi) - X'^{25}(0) = -(2\pi l - \theta_j + \theta_i)R'.$$

$$X'^{25} = \theta_i R' = -2\pi\alpha' A_{25,ii}.$$

$$X'^{25}(z, \bar{z}) = \theta_i R' - \frac{iR'}{2\pi} (2\pi l - \theta_j + \theta_i) \ln(z/\bar{z}) + i \left(\frac{\alpha'}{2} \right)^{1/2} \sum_{\substack{m=-\infty \\ m \neq 0}}^{\infty} \frac{\alpha_m^{25}}{m} (z^{-m} - \bar{z}^{-m})$$

$$= \theta_i R' + \frac{\sigma^1}{\pi} \Delta X'^{25} - (2\alpha')^{1/2} \sum_{\substack{m=-\infty \\ m \neq 0}}^{\infty} \frac{\alpha_m^{25}}{m} \exp(-m\sigma^2) \sin m\sigma^1$$

$$m^2 = \left(\frac{\Delta X'^{25}}{2\pi\alpha'} \right)^2 + \frac{1}{\alpha'} (N - 1)$$

Membranas.

$$\begin{aligned} \alpha_{-1}^\mu |k; ii\rangle, \quad \mathcal{V} &= i\partial_t X^\mu \\ \alpha_{-1}^{25} |k; ii\rangle, \quad \mathcal{V} &= i\partial_t X^{25} = \partial_n X'^{25} \end{aligned}$$



$$S_p = -T_p \int d^{p+1} \xi e^{-\Phi} [-\det(G_{ab} + B_{ab} + 2\pi\alpha' F_{ab})]^{1/2}$$

$$G_{ab}(\xi) = \frac{\partial X^\mu}{\partial \xi^a} \frac{\partial X^\nu}{\partial \xi^b} G_{\mu\nu}(X(\xi)), B_{ab}(\xi) = \frac{\partial X^\mu}{\partial \xi^a} \frac{\partial X^\nu}{\partial \xi^b} B_{\mu\nu}(X(\xi))$$

$$X'^2 = -2\pi\alpha' X^1 F_{12}$$

$$\int dX^1[1+(\partial_1X'^2)^2]^{1/2}=\int dX^1[1+(2\pi\alpha'F_{12})^2]^{1/2}$$

$$\frac{i}{4\pi\alpha'}\int_M d^2\sigma g^{1/2}\epsilon^{ab}\partial_aX^\mu\partial_bX^\nu B_{\mu\nu}+i\int_{\partial M}dX^\mu A_\mu$$

$$\delta A_\mu = \partial_\mu \lambda$$

$$\delta B_{\mu\nu} = \partial_\mu \zeta_\nu - \partial_\nu \zeta_\mu$$

$$\delta A_\mu = -\zeta_\mu/2\pi\alpha'$$

$$B_{\mu\nu} + 2\pi\alpha' F_{\mu\nu} \equiv 2\pi\alpha' \mathcal{F}_{\mu\nu}$$

$$V \propto \text{Tr}([X_m,X_n][X^m,X^n]).$$

$$S_p = -T_p \int d^{p+1} \xi \text{Tr}\{e^{-\Phi}[-\det(G_{ab} + B_{ab} + 2\pi\alpha' F_{ab})]^{1/2} + O([X^m,X^n]^2)\}$$

$$T_p e^{-\Phi} \prod_{i=1}^p (2\pi R_i)$$

$$2\pi\alpha'^{1/2}T_p e^{-\Phi'}\prod_{i=1}^{p-1}(2\pi R_i)$$

$$T_{p-1}e^{-\Phi'}\prod_{i=1}^{p-1}(2\pi R_i)$$

$$T_p = T_{p-1}/2\pi\alpha'^{1/2},$$

$$\mathcal{A}=iV_{p+1}\int_0^\infty\frac{dt}{t}(8\pi^2\alpha't)^{-(p+1)/2}\exp\left(-ty^2/2\pi\alpha'\right)\eta(it)^{-24}$$

$$=\frac{iV_{p+1}}{(8\pi^2\alpha')^{(p+1)/2}}\int_0^\infty dt t^{(21-p)/2}\exp\left(-ty^2/2\pi\alpha'\right)\times\left[\exp\left(2\pi/t\right)+24+\cdots\right]$$

$$\mathcal{A}=iV_{p+1}\frac{24}{2^{12}}(4\pi^2\alpha')^{11-p}\pi^{(p-23)/2}\Gamma\left(\frac{23-p}{2}\right)|y|^{p-23}=iV_{p+1}\frac{24\pi}{2^{10}}(4\pi^2\alpha')^{11-p}G_{25-p}(y)$$

$$S=\frac{1}{2\kappa^2}\int d^{26}X(-\tilde{G})^{1/2}\left(\tilde{R}-\frac{1}{6}\nabla_\mu\tilde{\Phi}\tilde{\nabla}^\mu\tilde{\Phi}\right)$$



$$S_p = -\tau_p \int d^{p+1} \xi \exp \left(\frac{p-11}{12} \tilde{\Phi} \right) (-\det \tilde{G}_{ab})^{1/2}$$

$$F_v \equiv \partial^{\hat{\mu}} h_{\mu\nu} - \frac{1}{2} \partial_\nu h_{\hat{\mu}}^{\hat{\mu}} = 0$$

$$S = -\frac{1}{8\kappa^2} \int d^{26}X \left(\partial_\mu h_{\nu\lambda} \partial^{\hat{\mu}} h^{\hat{\nu}\hat{\lambda}} - \frac{1}{2} \partial_\mu h_{\hat{\nu}}^{\hat{\nu}} \partial^{\hat{\mu}} h_{\hat{\lambda}}^{\hat{\lambda}} + \frac{2}{3} \partial_\mu \tilde{\Phi} \partial^{\hat{\mu}} \tilde{\Phi} \right)$$

$$\begin{aligned}\langle \tilde{\Phi} \tilde{\Phi} \rangle &= -\frac{(D-2) i \kappa^2}{4 k^2}, \\ \langle h_{\mu\nu} h_{\sigma\rho} \rangle &= -\frac{2 i \kappa^2}{k^2} \left(\eta_{\mu\sigma} \eta_{\nu\rho} + \eta_{\mu\rho} \eta_{\nu\sigma} - \frac{2}{D-2} \eta_{\mu\nu} \eta_{\sigma\rho} \right).\end{aligned}$$

$$S_p = -\tau_p \int d^{p+1} \xi \left(\frac{p-11}{12} \tilde{\Phi} - \frac{1}{2} h_{aa} \right)$$

$$\mathcal{A} = \frac{i\kappa^2\tau_p^2}{k_\perp^2}V_{p+1}\left\{6\left[\frac{p-11}{12}\right]^2 + \frac{1}{2}\left[2(p+1) - \frac{1}{12}(p+1)^2\right]\right\} = \frac{6i\kappa^2\tau_p^2}{k_\perp^2}V_{p+1}$$

$$\tau_p^2 = \frac{\pi}{256\kappa^2} \left(4\pi^2\alpha'\right)^{11-p}.$$

$$\frac{\tau_{25}}{4}(2\pi\alpha')^2\mathrm{Tr}(F_{\mu\nu}F^{\mu\nu}).$$

$$\frac{g_o^2}{g_c} = \frac{4\pi\alpha' g_o'^2}{\kappa} = 2^{18}\pi^{25/2}\alpha'^6$$

$$\Omega\colon X^M_L(z)\leftrightarrow X^M_R(z)$$

$$\begin{array}{l} \Omega: \quad X'^m(z,\bar{z}) \leftrightarrow -X'^m(\bar{z},z) \\ \quad \quad X^\mu(z,\bar{z}) \leftrightarrow X^\mu(\bar{z},z) \end{array}$$

$$\begin{array}{l} G_{\mu\nu}(x') \,=\, G_{\mu\nu}(x), B_{\mu\nu}(x') \,=\, -B_{\mu\nu}(x) \\ G_{\mu n}(x') \,=\, -G_{\mu n}(x), B_{\mu n}(x') \,=\, B_{\mu n}(x) \\ G_{mn}(x') \,=\, G_{mn}(x), B_{mn}(x') \,=\, -B_{mn}(x) \end{array}$$

$$W = \text{diag}(e^{i\theta_1}, e^{-i\theta_1}, e^{i\theta_2}, e^{-i\theta_2}, \ldots, e^{i\theta_{n/2}}, e^{-i\theta_{n/2}})$$

$$2^{12-k}T_p\int d^{p+1}\xi e^{-\Phi}(-\det G_{ab})^{1/2}$$

$$g_{\mathrm{c},d}\colon e^{ik\cdot X}\colon$$

$$\mathcal{V}_j(k;z,\bar{z})\overline{\mathcal{V}_{j'}(k;0,0)}=\frac{g_{\mathrm{c},d}^2}{z^2\bar{z}^2}+\cdots$$

$$e^{-2\lambda}\langle \tilde{c}(\bar{z}_1)c(z_1)\tilde{c}(\bar{z}_2)c(z_2)\tilde{c}(\bar{z}_3)c(z_3)\colon e^{ik\cdot X}(z,\bar{z})\colon\rangle=\frac{8\pi i}{\alpha'g_{\mathrm{c},d}^2}|z_{12}z_{13}z_{23}|^2(2\pi)^d\delta^d(k)$$



$$S_{S_2}(1; 2; 3) = e^{-2\lambda} \langle \hat{\mathcal{V}}'_1(\infty, \infty) \hat{\mathcal{V}}_2(1, 1) \hat{\mathcal{V}}_3(0, 0) \rangle_{S_2}$$

$$\hat{\mathcal{V}}_j = \tilde{c} \mathcal{V}_j$$

$$S_{S_2}(1; 2; 3) = e^{-2\lambda} \left\langle \left\langle \hat{\mathcal{V}}_1 \middle| \hat{\mathcal{V}}_2(1, 1) \middle| \hat{\mathcal{V}}_3 \right\rangle \right\rangle = e^{-2\lambda} c_{123}.$$

$$|\hat{\mathcal{V}}_3\rangle = Q_{\rm B}|\chi\rangle.$$

$$e^{2\lambda}S_{S_2}(1; 2; 3) = \left\langle \left\langle \hat{\mathcal{V}}_1 \middle| [\hat{\mathcal{V}}_2(1, 1), Q_{\rm B}] \middle| \chi \right\rangle \right\rangle + \left\langle \left\langle \hat{\mathcal{V}}_1 \middle| Q_{\rm B} \hat{\mathcal{V}}_2(1, 1) \middle| \chi \right\rangle \right\rangle = 0$$

$$\hat{\mathcal{V}}_3 = Q_{\rm B} \cdot \mathcal{X}.$$

$$\sum_n S_{mn} S_{pn}^* \stackrel{?}{=} \delta_{mp}$$

$$S_{mn}=\delta_{mn}+iT_{mn}$$

$$T_{mp}-T_{pm}^*=i\sum_nT_{mn}T_{pn}^*$$

$$T_{S_2}(1; 2; 3) = T_{S_2}(-1; -2; -3)^*$$

$$\hat{\mathcal{V}}_{-i}=-\overline{\hat{\mathcal{V}}_i}.$$

$$\langle \hat{\mathcal{V}}'_1(\infty, \infty) \hat{\mathcal{V}}_2(1, 1) \hat{\mathcal{V}}_3(0, 0) \rangle_{S_2} = - \langle \hat{\mathcal{V}}'_{-1}(\infty, \infty) \hat{\mathcal{V}}_{-2}(1, 1) \hat{\mathcal{V}}_{-3}(0, 0) \rangle_{S_2}^*,$$

$$S_{S_2}(1; 2; 3; 4) = e^{-2\lambda} \int_{\mathbf{C}} d^2z \langle \hat{\mathcal{V}}'_1(\infty, \infty) \hat{\mathcal{V}}_2(1, 1) \hat{\mathcal{V}}_3(0, 0) \mathcal{V}_4(z, \bar{z}) \rangle_{S_2}.$$

$$\left\langle \left\langle \hat{\mathcal{V}}_1 \middle| \mathsf{T}[\hat{\mathcal{V}}_2(1, 1) \mathcal{V}_4(z, \bar{z})] \middle| \hat{\mathcal{V}}_3 \right\rangle \right\rangle.$$

$$\mathcal{V}_4(z,\bar{z})=(z\bar{z})^{-1}\{b_0,[\tilde{b}_0,\hat{\mathcal{V}}_4(z,\bar{z})]\}.$$

$$\theta(1-|z|)\left\langle \left\langle \hat{\mathcal{V}}_1 \middle| \hat{\mathcal{V}}_2(1, 1) b_0 \tilde{b}_0 z^{L_0-1} \bar{z}^{\tilde{L}_0-1} \hat{\mathcal{V}}_4(1, 1) \middle| \hat{\mathcal{V}}_3 \right\rangle \right\rangle + \theta(|z|$$

$$-1)\left\langle \left\langle \hat{\mathcal{V}}_1 \middle| \hat{\mathcal{V}}_4(1, 1) b_0 \tilde{b}_0 z^{-L_0-1} \bar{z}^{-\tilde{L}_0-1} \hat{\mathcal{V}}_2(1, 1) \middle| \hat{\mathcal{V}}_3 \right\rangle \right\rangle$$

$$\theta(1-|z|)\sum_{i,i'}z^{\alpha'(k_i^2+m_i^2)/4-1}\bar{z}^{\alpha'(k_i^2+\tilde{m}_i^2)/4-1}\times\left\langle \left\langle \hat{\mathcal{V}}_1 \middle| \hat{\mathcal{V}}_2(1, 1) b_0 \tilde{b}_0 \middle| i \right\rangle \mathcal{G}^{ii'}\left\langle \left\langle i' \middle| \hat{\mathcal{V}}_4(1, 1) \middle| \hat{\mathcal{V}}_3 \right\rangle \right\rangle + \theta(|z|$$

$$-1)\sum_{i,i'}z^{-\alpha'(k_i^2+m_i^2)/4-1}\bar{z}^{-\alpha'(k_i^2+\tilde{m}_i^2)/4-1}$$

$$\times\left\langle \left\langle \hat{\mathcal{V}}_1 \middle| \hat{\mathcal{V}}_4(1, 1) b_0 \tilde{b}_0 \middle| i \right\rangle \mathcal{G}^{ii'}\left\langle \left\langle i' \middle| \hat{\mathcal{V}}_2(1, 1) \middle| \hat{\mathcal{V}}_3 \right\rangle \right\rangle$$

$$k_i^2 = (k_1 + k_2)^2 > \frac{4}{\alpha'},$$

$$\int_{|z|<1} d^2z z^{\alpha'(k^2+m^2)/4-1} \bar{z}^{\alpha'(k^2+\tilde{m}^2)/4-1} = \frac{8\pi}{\alpha'} \frac{\delta_{m^2,\tilde{m}^2}}{k^2+m^2-i\epsilon}$$

$$\left[b_0\tilde{b}_0z^{\pm L_0-1}\bar{z}^{\pm \tilde{L}_0-1},Q_{\rm B}\right]=\pm b_0\partial_{\bar{z}}\big(z^{\pm L_0-1}\bar{z}^{\pm \tilde{L}_0}\big)\mp \tilde{b}_0\partial_z\big(z^{\pm L_0}\bar{z}^{\pm \tilde{L}_0-1}\big)$$

$$T_{S_2}(1;2;3;4)-T_{S_2}(-1;-2;-3;-4)^*$$

$$=i\sum_{j_5}\int\frac{d^{d-1}\mathbf{k}_5}{2E_5(2\pi)^{d-1}}T_{S_2}(1;2;5)T_{S_2}(-5;3;4)+2\text{ permutations}$$

$$=-ie^{-4\lambda}\sum_{j_5}\int\frac{d^{d-1}\mathbf{k}_5}{2E_5(2\pi)^{d-1}}\Big\langle\hat{\mathcal{V}}_1|\hat{\mathcal{V}}_2(1,1)|\hat{\mathcal{V}}_5\Big\rangle\Big\langle\hat{\mathcal{V}}_{-5}|\hat{\mathcal{V}}_3(1,1)|\hat{\mathcal{V}}_4\Big\rangle$$

$$+2\text{ permutations}$$

$$\frac{1}{k^2+m^2-i\epsilon}-\frac{1}{k^2+m^2+i\epsilon}=2\pi i\delta(k^2+m^2)$$

$$T_{S_2}(1;2;3;4)-T_{S_2}(-1;-2;-3;-4)^*$$

$$=\frac{16\pi^2i}{\alpha'e^{2\lambda}}\sum_{i,i'}\delta_{m_i^2,\tilde{m}_i^2}\delta(k_i^2+m_i^2)\langle\hat{\mathcal{V}}_1|\hat{\mathcal{V}}_2(1,1)b_0\tilde{b}_0|i\rangle\mathcal{G}^{ii'}\langle i'|\hat{\mathcal{V}}_3(1,1)|\hat{\mathcal{V}}_4\rangle$$

$$+2\text{ permutations}$$

$$i\leftrightarrow (j,k)$$

$$\sum_{i,i'}\delta_{m_i^2,\tilde{m}_i^2}\delta(k_i^2+m_i^2)b_0\tilde{b}_0|i\rangle\mathcal{G}^{ii'}\langle i'|=\delta_{L_0,\tilde{L}_0}\delta(4L_0/\alpha')b_0\tilde{b}_0$$

$$=-\frac{i\alpha'}{16\pi^2e^{2\lambda}}\int\frac{d^{d-1}\mathbf{k}}{2E(2\pi)^{d-1}}\sum_{\substack{j\in\mathcal{H}\\k^0=\pm\omega_k}}|\hat{\mathcal{V}}_j(k)\rangle\langle\hat{\mathcal{V}}_j(k)|$$

$$\left|\overline{\hat{\mathcal{V}}_j(k)}\right|\tilde{c}_0c_0|\hat{\mathcal{V}}_{j'}(k')\rangle=\frac{8\pi ie^{2\lambda}}{\alpha'}\delta_{jj'}(2\pi)^d\delta^d(k-k')$$

$$\mathcal{P}|b\rangle_{\rm P}=|b\rangle_{\rm P},\mathcal{P}|a\rangle_{\rm U}=0,\mathcal{P}|a\rangle_{\rm N}=0;\\ \mathcal{U}|b\rangle_{\rm P}=0,\mathcal{U}|a\rangle_{\rm U}=0,\mathcal{U}|a\rangle_{\rm N}=|a\rangle_{\rm U}.$$

$$1=\mathcal{P}+Q_{\rm B}\mathcal{U}+\mathcal{U}Q_{\rm B}.$$

$$-i\frac{\eta^{\mu\nu}}{k^2}$$

$$\eta^{\mu\nu}=(\eta^{\mu\nu}-k^\mu n^\nu-n^\mu k^\nu)+k^\mu n^\nu+n^\mu k^\nu.$$



$$\left.\frac{\partial z_1}{\partial z}\right|_{z_2}=\frac{z_1}{z}$$

$$\gamma\colon \frac{z'-U}{z'-V}=K\frac{z-U}{z-V},$$

$$|K|^{-1/2}<\left|\frac{z-U}{z-V}\right|<|K|^{1/2}.$$

$$ds^2=\frac{dzd\bar{z}}{(\mathrm{Im}z)^2}$$

$$(A_1B_1^{-1}A_1^{-1}B_1)(A_2B_2^{-1}A_2^{-1}B_2)\ldots(A_gB_g^{-1}A_g^{-1}B_g)$$

$$\partial_{\bar{z}}\omega_z=0,\omega_{\bar{z}}=0$$

$$\dim\ker P_0^T-\dim\ker P_0=-\chi=2g-2.$$

$$\oint_{A_i} dz \omega_{zj} = \delta_{ij}$$

$$\tau_{ij}=\oint_{B_i} dz \omega_{zj}$$

$$y^2=\prod_{i=1}^{2k}(z-Z_i)$$

$$z_1z_2=q$$

$$(1-\epsilon)^{-1}|q|^{1/2}>|z_1|>(1-\epsilon)|q|^{1/2}$$

$$M=M_1\!\!\infty\!\!M_2(z_1,z_2,q)$$

$$M=M_08(z_1,z_2,q).$$

$$\ln\ |q|<\mathrm{Im}w<0, w\cong w+2\pi,$$

$$z^{(0)}=\frac{az}{z-2}, z^{(1)}=\frac{a(z-1)}{z+1}, z^{(\infty)}=\frac{a}{2z-1}$$

$$z_4=(q-a^2)^2/4a^2q$$

$$r=3g+n-2$$

$$\langle \cdots 1 \cdots 2 \rangle_M = \sum_{ij} \left\langle \cdots 1 \mathcal{A}_i^{(z_1)} \right\rangle_{M_1} G^{ij} \left\langle \mathcal{A}_j^{(z_2)} \cdots 2 \right\rangle_{M_2}.$$



$$\left\langle \cdots \mathcal{A}_k^{(z_1)} \right\rangle_{M_1} = \sum_{ij} \left\langle \cdots \mathcal{A}_i^{(z_1)} \right\rangle_{M_1} G^{ij} \left\langle \mathcal{A}_j^{(u)} \mathcal{A}_k^{(z)} \right\rangle_{S_2}.$$

$$G^{ij}\mathcal{G}_{jk}=\delta_k^i$$

$$G^{ij}=\mathcal{G}^{ij}$$

$$\mathcal{A}_j^{(z_2')} = q^{h_j} \bar{q}^{\tilde{h}_j} \mathcal{A}_j^{(z_2)}$$

$$\langle \cdots 1 \cdots 2 \rangle_M = \sum_{ij} q^{h_j} \bar{q}^{\tilde{h}_j} \left\langle \cdots 1 \mathcal{A}_i^{(z_1)} \right\rangle_{M_1} \mathcal{G}^{ij} \left\langle \mathcal{A}_j^{(z_2)} \cdots 2 \right\rangle_{M_2}$$

$$\langle \cdots \rangle_M = \sum_{ij} q^{h_j} \bar{q}^{\tilde{h}_j} \mathcal{G}^{ij} \left\langle \cdots \mathcal{A}_i^{(z_1)} \mathcal{A}_j^{(z_2)} \right\rangle_{M_0}$$

$$\begin{aligned} z_1' &= z_1 + \sum_{n=-\infty}^{\infty} \epsilon_n z_1^{n+1} \\ z_2' &= z_2 - \sum_{n=-\infty}^{\infty} \epsilon_{-n} q^{-n} z_2^{n+1} \end{aligned}$$

$$\mathcal{A}_i^{(z_1')} \mathcal{G}^{ij} \mathcal{A}_j^{(z_2')} - \mathcal{A}_i^{(z_1)} \mathcal{G}^{ij} \mathcal{A}_j^{(z_2)} = - \sum_{n=-\infty}^{\infty} \epsilon_n \left[L_n \cdot \mathcal{A}_i^{(z_1)} \mathcal{G}^{ij} \mathcal{A}_j^{(z_2)} - \mathcal{A}_i^{(z_1)} \mathcal{G}^{ij} L_{-n} \cdot \mathcal{A}_j^{(z_2)} \right]$$

$$2\pi(m,n).$$

$$(\sigma^1,\sigma^2)=(\sigma'^1,\sigma'^2)\begin{bmatrix}m&n\\q&p\end{bmatrix}$$

$$(\sigma'^1,\sigma'^2)=2\pi(1,0)$$

$$\phi\,\,\left.\frac{dz}{2\pi i}b_{z_1z_1}\frac{\partial z_1}{\partial q}\right|_{z_2}=\frac{b_0}{q}$$

$$\int_{|q|<1} \frac{d^2q}{q\bar{q}} q^{\alpha'(k_i^2+m_i^2)/4} \bar{q}^{\alpha'(k_i^2+\tilde{m}_i^2)/4} \mathcal{G}^{ij} b_0 \tilde{b}_0 = \frac{8\pi \delta_{m_i^2,\tilde{m}_i^2} \mathcal{G}^{ij} b_0 \tilde{b}_0}{\alpha'(k_i^2+m_i^2-i\epsilon)}.$$

$$g(i_1,\ldots,i_n)=\sum_g\int_{\mathcal{V}_{g,n}}d^mt\left\langle\prod_{k=1}^mB_k\prod_{l=1}^n\hat{\mathcal{V}}_{i_l}\right\rangle_g$$

$$\partial X^\mu \bar{\partial} X_\mu \frac{dq d\bar{q}}{q \bar{q}}$$

$$-4\pi\partial X^\mu\bar{\partial}X_\mu\ln\left(ae^{-\omega}\right)$$

$$\delta \mathcal{S} = - \sum_{\chi} \Lambda_{\chi} \int d^d x (-G)^{1/2} e^{-\chi \Phi}$$



$$\int [d\Psi]\exp\left(iS[\Psi]\right)$$

$$\delta\Psi=Q_{\rm B}\Lambda$$

$$S_0=\frac{1}{2}\langle\langle\Psi|Q_{\rm B}|\Psi\rangle.$$

$$Q_{\rm B}\Psi=0$$

$$\Psi[X,c,\tilde{c}]=\sum_i\Phi_i(x)\Psi_i[X',c,\tilde{c}]$$

$$\int [d\Psi] \rightarrow \prod_i \int [d\Phi_i]$$

$$\Psi=[\varphi(x)+A_{\mu}(x)\alpha_{-1}^{\mu}+B(x)b_{-1}+C(x)c_{-1}$$

$$+\varphi'(x)c_0+A_{\mu}'(x)\alpha_{-1}^{\mu}c_0+B'(x)b_{-1}c_0+C'(x)c_{-1}c_0+\cdots]\Psi_0$$

$$\delta\Psi=Q_{\rm B}\Lambda+g\Psi*\Lambda-g\Lambda*\Psi.$$

$$Q_{\rm B}(\Psi_1*\Psi_2)=(Q_{\rm B}\Psi_1)*\Psi_2+\Psi_1*(Q_{\rm B}\Psi_2)$$

$$\int \Psi_1*\Psi_2=\int \Psi_2*\Psi_1$$

$$S=\frac{1}{2}\int \Psi*Q_{\rm B}\Psi+\frac{2g}{3}\int \Psi*\Psi*\Psi$$

$$S=\frac{1}{2}\langle\mathcal{V}_{\Psi}Q_{\rm B}\cdot\mathcal{V}_{\Psi}\rangle_{D_2}+\frac{2g}{3}\langle\mathcal{V}_{\Psi}\mathcal{V}_{\Psi}\mathcal{V}_{\Psi}\rangle_{D_2}$$

$$\int_0^\infty dt \exp\left(-tL_0\right)=L_0^{-1}$$

$$S_0'=\frac{1}{2}\langle\langle\Psi|(c_0-\tilde{c}_0)Q_{\rm B}|\Psi\rangle$$

$$\int_{-\infty}^{\infty} dy \exp\left[-\frac{y^2}{2!}-\lambda\frac{y^3}{3!}\right]=\lambda^{-1}\int_{-\infty}^{\infty} dz \exp\left[-\frac{1}{\lambda^2}\left(\frac{z^2}{2!}+\frac{z^3}{3!}\right)\right]$$

$$\sum_{n=0}^{\infty}\frac{(-\lambda)^n}{6^n n!}\int_{-\infty}^{\infty} dy y^{3n}\exp\left(-y^2/2\right)=(2\pi)^{1/2}\sum_{k=0}^{\infty}\lambda^{2k}C_{2k}$$

$$C_{2k}=\frac{2^k\Gamma(3k+1/2)}{\pi^{1/2}3^{2k}(2k)!}$$

$$C_{2k}\approx k^k\approx k!,$$



$$\sum_{k=0}^{\infty} \lambda^{2k} k! f_{2k}$$

$$\exp\left[-O(1/\lambda^2)\right]$$

$$\int_0^\infty dt \exp(-t) \sum_{k=0}^\infty (t\lambda^2)^k f_{2k}$$

$$\sum_{k=0}^\infty g_o^{2k} O(k!) = \sum_{k=0}^\infty g_c^k O(k!)$$

$$\exp\left[-O(1/g_c)\right]$$

$$\frac{8\pi i g_c^2}{\alpha'} (2\pi)^{26} \delta^{26} \left(\sum_i k_i \right) \int_{\mathbb{C}} d^2 z_4 |z_4|^{-\alpha' u/2-4} |1-z_4|^{-\alpha' t/2-4}$$

$$\frac{\partial}{\partial z_4} (u \ln |z_4|^2 + t \ln |1-z_4|^2) = \frac{u}{z_4} + \frac{t}{z_4-1} = 0$$

$$S \approx \exp\left[-\alpha'(\mathrm{sln}\,s + \mathrm{tln}\,t + \mathrm{uln}\,u)/2\right]$$

$$X_{\mathrm{cl}}^\mu(\sigma)=i\sum_i k_i^\mu G'(\sigma,\sigma_i),$$

$$\exp\left[-\sum_{i<j}k_i\cdot k_jG'(\sigma_i,\sigma_j)\right]$$

$$z_2^N=\frac{(z_1-a_1)(z_1-a_2)}{(z_1-a_3)(z_1-a_4)}$$

$$S_g \approx \exp\left[-\alpha'(\mathrm{sln}\,s + \mathrm{tln}\,t + \mathrm{uln}\,u)/2(g+1)\right]$$

$$\exp\left[\alpha' \mathrm{tln}\,s/2(g+1)\right]$$

$$\langle 0|[X^1(\sigma)-x^1]^2|0\rangle=\sum_{n=1}^\infty\frac{\alpha'}{2n^2}\langle 0|(\alpha_n\alpha_{-n}+\tilde\alpha_n\tilde\alpha_{-n})|0\rangle=\alpha'\sum_{n=1}^\infty\frac{1}{n}$$

$$s^{2+\alpha' t/2} \frac{\Gamma(-\alpha' t/4-1)}{\Gamma(\alpha' t/4+2)} \sim \frac{s^2}{t} \exp\left[-\frac{q^2 \alpha'}{2} \ln(\alpha' s)\right]$$

$$\int_0^\infty dm n(m) \exp\left(-\alpha'\pi m^2\ell\right) \approx \exp\left(4\pi/\ell\right)$$

$$\int_0^\infty dm \exp\left(4\pi m \alpha'^{1/2}\right) \exp\left(-m/T\right)$$



$$T_{\rm H}=\frac{1}{4\pi\alpha'^{1/2}}$$

$$F(T,m^2)=T\int\frac{d^{d-1}\mathbf{k}}{(2\pi)^{d-1}}\ln\left[1-\exp\left(-\omega_k/T\right)\right]=-\int_0^\infty\frac{dt}{t}(2\pi t)^{-d/2}\sum_{r=1}^\infty\exp\left(-\frac{m^2t}{2}-\frac{r^2}{2T^2t}\right)$$

$$F(T)=-\int_R\frac{d\tau d\bar\tau}{2\tau_2}(4\pi^2\alpha'\tau_2)^{-13}|\eta(\tau)|^{-48}\sum_{r=1}^\infty\exp\left(-\frac{r^2}{4\pi T^2\alpha'\tau_2}\right)$$

$$R\colon \tau_1\leq \frac{1}{2}, |\tau_2|>0$$

$$\eta(i\tau_2)=\eta(i/\tau_2)\tau_2^{-1/2}\approx\exp\left(-\pi/12\tau_2\right)\text{ as }\tau_2\rightarrow 0$$

$$\tilde{F}(T)=F(T)+\rho_0$$

$$\frac{1}{T}\tilde{F}(T)=\frac{T}{4T_{\rm H}^2}\tilde{F}(4T_{\rm H}^2/T)$$

$$\tilde{F}(T)\rightarrow \frac{T^2}{4T_{\rm H}^2}\rho_0$$

$$G_{\mu\nu}(x)\,=\,\eta_{\mu\nu}, B_{\mu\nu}(x)=0, \Phi(x)=V_\mu x^\mu$$

$$V_\mu\,=\,\delta_\mu\,^1\left(\frac{26-D}{6\alpha'}\right)^{1/2}$$

$$-\partial^\mu\partial_\mu T(x)+2V^\mu\partial_\mu T(x)-\frac{4}{\alpha'}T(x)=0$$

$$T(x)=\exp{(q\cdot x)},(q-V)^2=\frac{2-D}{6\alpha'}$$

$$q_1=\alpha_{\pm}=\left(\frac{26-D}{6\alpha'}\right)^{1/2}\pm\left(\frac{2-D}{6\alpha'}\right)^{1/2}$$

$$S_\sigma=\frac{1}{4\pi\alpha'}\int_Md^2\sigma g^{1/2}[g^{ab}\eta_{\mu\nu}\partial_aX^\mu\partial_bX^\nu+\alpha'RV_1X^1+T_0\exp{(\alpha_-X^1)}]$$

$$S=\frac{1}{4\pi\alpha'}\int_Md^2\sigma g^{1/2}(g^{ab}\partial_aX^\mu\partial_bX_\mu+\mu)$$

$$g_{ab}(\sigma)=\exp{(2\phi(\sigma))}\hat{g}_{ab}(\sigma)$$

$$g^{1/2}R=\hat{g}^{1/2}(\hat{R}-2\hat{\nabla}^2\phi)$$

$$S = \frac{1}{4\pi\alpha'} \int_M d^2\sigma \hat{g}^{1/2} \left[\hat{g}^{ab} \partial_a X^\mu \partial_b X_\mu + \mu \exp 2\phi + \frac{13\alpha'}{3} (\hat{g}^{ab} \partial_a \phi \partial_b \phi + \hat{R}\phi) \right]$$

$$\begin{aligned}\hat{g}_{ab}(\sigma) &\rightarrow \exp(2\omega(\sigma))\hat{g}_{ab}(\sigma) \\ \phi(\sigma) &\rightarrow \phi(\sigma) - \omega(\sigma)\end{aligned}$$

Campos bosónicos.

$$[\hat{q},\hat{p}]=i\hbar$$

$$\langle q_f|\exp{(-i\hat{H}T/\hbar)}|q_i\rangle$$

$$\hat{A}(t)=\exp{(i\hat{H}t/\hbar)}\hat{A}\exp{(-i\hat{H}t/\hbar)}$$

$$\hat{q}(t)|q,t\rangle=|q,t\rangle q$$

$$|q,t\rangle=\exp{(i\hat{H}t/\hbar)}|q\rangle$$

$$\langle q_f,T\mid q_i,0\rangle=\int dq\langle q_f,T\mid q,t\rangle\langle q,t\mid q_i,0\rangle$$

$$t_m=m\epsilon,\epsilon=T/N$$

$$\langle q_f,T\mid q_i,0\rangle=\int dq_{N-1}\ldots dq_1\prod_{m=0}^{N-1}\langle q_{m+1},t_{m+1}\mid q_m,t_m\rangle$$

$$\langle q_{m+1},t_{m+1}\mid q_m,t_m\rangle=\langle q_{m+1}|\exp{(-i\hat{H}\epsilon/\hbar)}|q_m\rangle=\int dp_m\langle q_{m+1}\mid p_m\rangle\langle p_m|\exp{(-i\hat{H}\epsilon/\hbar)}|q_m\rangle$$

$$\langle p_m|\hat{H}(\hat{p},\hat{q})|q_m\rangle=H(p_m,q_m)\langle p_m\mid q_m\rangle.$$

$$\int dp_m\exp\left[-\frac{iH(p_m,q_m)\epsilon}{\hbar}\right]\langle q_{m+1}\mid p_m\rangle\langle p_m\mid q_m\rangle$$

$$=\int \frac{dp_m}{2\pi\hbar}\exp\left\{-\frac{i}{\hbar}[H(p_m,q_m)\epsilon-p_m(q_{m+1}-q_m)]+O(\epsilon^2)\right\}$$

$$\langle q_f,T\mid q_i,0\rangle\approx\int \frac{dp_{N-1}}{2\pi\hbar}dq_{N-1}\ldots\frac{dp_1}{2\pi\hbar}dq_1\frac{dp_0}{2\pi\hbar}\times\exp\left\{-\frac{i}{\hbar}\sum_{m=0}^{N-1}[H(p_m,q_m)\epsilon-p_m(q_{m+1}-q_m)]\right\}$$

$$\rightarrow\int [dpdq]\exp\left\{\frac{i}{\hbar}\int_0^T dt[p\dot{q}-H(p,q)]\right\}$$

$$0=\frac{\partial}{\partial p}[p\dot{q}-H(p,q)]=\dot{q}-\frac{\partial}{\partial p}H(p,q).$$

$$\langle q_f,T\mid q_i,0\rangle=\int [dq]\exp\left[\frac{i}{\hbar}\int_0^T dtL(q,\dot{q})\right]$$



$$\int [dq]_{q_i,0}^{q_f,T} \exp \left(i \int_0^T dt L \right) = \int dq \int [dq]_{q,t}^{q_f,T} \exp \left(i \int_t^T dt L \right) \int [dq]_{q_i,0}^{q,t} \exp \left(i \int_0^t dt L \right)$$

$$\int [dq]_{q_i,0}^{q_f,T} \exp (iS) q(t) = \int dq \langle q_f, T | q, t \rangle q \langle q, t | q_i, 0 \rangle = \int dq \langle q_f, T | \hat{q}(t) | q, t \rangle \langle q, t | q_i, 0 \rangle$$

$$= \langle q_f, T | \hat{q}(t) | q_i, 0 \rangle$$

$$\int [dq]_{q_i,0}^{q_f,T} \exp (iS) q(t) q(t') = \langle q_f, T | T[\hat{q}(t) \hat{q}(t')] | q_i, 0 \rangle.$$

$$T[\hat{A}(t) \hat{B}(t')] = \theta(t - t') \hat{A}(t) \hat{B}(t') + \theta(t' - t) \hat{B}(t') \hat{A}(t)$$

$$\int [dq]_{q_i,0}^{q_f,T} \exp (iS) \frac{\delta S}{\delta q(t)} \mathcal{F} = -i \int [dq]_{q_i,0}^{q_f,T} \left[\frac{\delta}{\delta q(t)} \exp (iS) \right] \mathcal{F} = i \int [dq]_{q_i,0}^{q_f,T} \exp (iS) \frac{\delta \mathcal{F}}{\delta q(t)}$$

$$\frac{\delta}{\delta q(t)} q(t') = \delta(t - t')$$

$$\frac{1}{\epsilon} \frac{\partial}{\partial q_m}, m = t/\epsilon$$

$$\left\langle q(t') \frac{\delta S}{\delta q(t)} \dots \right\rangle = i \langle \delta(t - t') \dots \rangle$$

$$q(t') \frac{\delta S}{\delta q(t)} = i \delta(t - t')$$

$$S = \frac{1}{2} \int dt (\dot{q}^2 - \omega^2 q^2)$$

$$\frac{\delta S}{\delta q(t)} = -\ddot{q}(t) - \omega^2 q(t)$$

$$\left(\frac{\partial^2}{\partial t^2} + \omega^2 \right) q(t) = 0$$

$$\left(\frac{\partial^2}{\partial t^2} + \omega^2 \right) q(t) q(t') = -i \delta(t - t')$$

$$\left(\frac{\partial^2}{\partial t^2} + \omega^2 \right) \hat{q}(t) = 0$$

$$\left(\frac{\partial^2}{\partial t^2} + \omega^2 \right) T[\hat{q}(t) \hat{q}(t')] = -i \delta(t - t')$$

$$\left(\frac{\partial^2}{\partial t^2} + \omega^2 \right) T[\hat{q}(t) \hat{q}(t')] = T\{[\ddot{\hat{q}}(t) + \omega^2 \hat{q}(t)] \hat{q}(t')\} + \delta(t - t') [\dot{\hat{q}}(t), \hat{q}(t)] = 0 - i \delta(t - t')$$

$$\langle q_f | \exp [-i\hat{H}(T - t_>)] \hat{q} \exp [-i\hat{H}(t_> - t_<)] \hat{q} \exp [-i\hat{H}t_<] | q_i \rangle,$$



$$\langle q_f, u | T \left[\prod_a \hat{q}(u_a) \right] | q_i, 0 \rangle_E$$

$$\langle q_f, t | T \left[\prod_a \hat{q}(t_a) \right] | q_i, 0 \rangle$$

$$\langle q_f, U | T \left[\prod_a \hat{q}(u_a) \right] | q_i, 0 \rangle_E = \int [dq]_{q_i,0}^{q_f,U} \exp \left[\int_0^U du L(q, i \partial_u q) \right] \prod_a q(u_a).$$

$$S_E = \int du L_E(q, \partial_u q) = - \int du L(q, i \partial_u q)$$

$$\text{Tr} \exp (-\hat{H}U) = \sum_i \exp (-E_i U)$$

$$\int dq \langle q | \exp (-\hat{H}U) | q \rangle = \int dq \langle q, U | q, 0 \rangle_E = \int dq \int [dq]_{q,0}^{q,U} \exp (-S_E) = \int [dq]_P \exp (-S_E)$$

$$\int [dq]_A \exp (-S_E) = \int dq \int [dq]_{-q,0}^{q,U} \exp (-S_E) = \int dq \langle q, U | -q, 0 \rangle_E = \int dq \langle q, U | \hat{R} | q, 0 \rangle_E$$

$$= \text{Tr} [\exp (-\hat{H}U) \hat{R}]$$

$$\frac{1}{2} (\text{periodic} + \text{antiperiodic}) = \sum_{R_i = \pm 1} \exp (-E_i U)$$

$$\hat{A}(u)^{\dagger}=\hat{A}(-u)$$

$$\overline{\hat{A}(u)}=\hat{A}(-u)^{\dagger}.$$

$$\frac{1}{2}\int d^dx\phi(x)\Delta\phi(x)$$

$$Z[J]=\left\langle \exp \left[i\int d^dxJ(x)\phi(x)\right] \right\rangle =\int [d\phi] \exp \left[-\frac{1}{2}\int d^dx\phi(x)\Delta\phi(x)+i\int d^dxJ(x)\phi(x)\right]$$

$$Z[J]=Z[0]\exp\left[-\frac{1}{2}\int d^dx d^dy J(x)\Delta^{-1}(x,y)J(y)\right]$$

$$=Z[0]\exp\left[\frac{1}{2}\int d^dx d^dy \Delta^{-1}(x,y)\frac{\delta^2}{\delta\phi(x)\delta\phi(y)}\right]\times\exp\left[i\int d^dxJ(x)\phi(x)\right]\Big|_{\phi=0}$$

$$\mathcal{F}[\phi]=\int [dJ]\exp\left[i\int d^dxJ(x)\phi(x)\right]f[J]$$

$$\langle \mathcal{F} \rangle = Z[0] \exp \left[\frac{1}{2} \int d^d x d^d y \Delta^{-1}(x,y) \frac{\delta^2}{\delta \phi(x) \delta \phi(y)} \right] \mathcal{F}[\phi] \Big|_{\phi=0}.$$



$$\frac{\delta}{\delta\phi(x)} = \frac{\delta}{\delta\phi_1(x)} + \frac{\delta}{\delta\phi_2(x)} + \cdots$$

$$\frac{\delta^2}{\delta\phi_i(x)\delta\phi_i(y)} \text{ (no sum on } i\text{)}$$

$$\phi(x) = \sum_i \phi_i \Phi_i(x)$$

$$\Delta\Phi_i(x) = \lambda_i\Phi_i(x)$$

$$\int d^d x \Phi_i(x)\Phi_j(x) = \delta_{ij}$$

$$[d\phi] = \prod_i d\phi_i$$

$$\langle 1 \rangle = \prod_i \left[\int d\phi_i \exp \left(-\frac{1}{2} \phi_i^2 \lambda_i \right) \right] = \prod_i \left(\frac{2\pi}{\lambda_i} \right)^{1/2} = \left(\det \frac{\Delta}{2\pi} \right)^{-1/2}$$

$$\int [d\phi d\phi^*] \exp \left(- \int d^d x \phi^* \Delta \phi \right) = (\det \Delta)^{-1}$$

$$\int [d\phi^1 d\phi^2] \exp \left(i \int d^d x \phi^1 \Delta \phi^2 \right) = (\det \Delta)^{-1}$$

$$\int_{-\infty}^{\infty} dx \int_{-\infty}^{\infty} dy \exp (i \lambda x y) = 2 \pi \int_{-\infty}^{\infty} dx \delta (\lambda x) = \frac{2 \pi}{\lambda}$$

$$\langle q_f, U \mid q_i, 0 \rangle_{\rm E} = \int [dq]_{q_i, 0}^{q_f, U} \exp (-S_{\rm E})$$

$$S_{\rm E} = \frac{1}{2} \int_0^U du [(\partial_u q)^2 + \omega^2 q^2] + S_{\rm ct}$$

$$q(u)=q_{\rm cl}(u)+q'(u)$$

$$-\ddot{q}_{\rm cl}(u)+\omega^2q_{\rm cl}(u)=0$$

$$q_{\rm cl}(0)=q_i, q_{\rm cl}(U)=q_f$$

$$q'(0)=q'(U)=0$$

$$S_{\rm E}=S_{\rm cl}(q_i,q_f)+S'+S_{\rm ct}$$

$$S_{\rm cl}(q_i,q_f)=\frac{1}{2}\int_0^U du [(\partial_u q_{\rm cl})^2+\omega^2 q_{\rm cl}^2]=\omega\frac{(q_i^2+q_f^2)\cosh\omega U-2q_iq_f}{2\sinh\omega U}$$

$$S'=\frac{1}{2}\int_0^U du [(\partial_u q')^2+\omega^2 q'^2]$$



$$\langle q_f, U | q_i, 0 \rangle = \exp [-S_{\text{cl}}(q_i, q_f) - S_{\text{ct}}] \int [dq']_{0,0}^{0,U} \exp (-S')$$

$$f_j(u) = \left(\frac{2}{U}\right)^{1/2} \sin \frac{j\pi u}{U}, \lambda_j = \frac{j^2\pi^2}{U^2} + \omega^2$$

$$\det \frac{\Delta}{2\pi} = \prod_{j=1}^{\infty} \frac{j^2\pi^2 + \omega^2 U^2}{2\pi U^2}$$

$$\det \frac{\Delta}{2\pi} \stackrel{\text{regulate}}{\rightarrow} \prod_{j=1}^{\infty} \frac{j^2\pi^2 + \omega^2 U^2}{j^2\pi^2 + \Omega^2 U^2}$$

$$\frac{\Omega \sinh \omega U}{\omega \sinh \Omega U}$$

$$\langle q_f, U | q_i, 0 \rangle \rightarrow \left(\frac{\omega}{2\sinh \omega U}\right)^{1/2} \exp \left[-S_{\text{cl}}(q_i, q_f) + \frac{1}{2}(\Omega U - \ln \Omega) - S_{\text{ct}}\right]$$

$$\langle q_f, U | q_i, 0 \rangle \rightarrow \langle q_f | 0 \rangle \langle 0 | q_i \rangle \exp (-E_0 U) = \left(\frac{\omega}{\pi}\right)^{1/2} \exp \left[-\frac{\omega}{2}(q_f^2 + q_i^2 + U)\right]$$

$$\langle q_f, U | q_i, 0 \rangle \rightarrow \omega^{1/2} \exp \left[-\frac{\omega}{2}(q_f^2 + q_i^2 + U) + \frac{1}{2}(\Omega U - \ln \Omega) - S_{\text{ct}}\right]$$

$$S_{\text{ct}} = \frac{1}{2} \left[\int_0^U du \Omega \right] - \frac{1}{2} \ln \frac{\Omega}{\pi}$$

$$\langle q_f, U | q_i, 0 \rangle = \left(\frac{\omega}{2\pi \sinh \omega U}\right)^{1/2} \exp (-S_{\text{cl}}),$$

Campos Fermiónicos.

$$\hat{\psi} | \downarrow \rangle = 0, \hat{\chi} | \downarrow \rangle = | \uparrow \rangle,$$

$$\hat{\psi} | \uparrow \rangle = | \downarrow \rangle, \hat{\chi} | \uparrow \rangle = 0,$$

$$\{\hat{\psi}, \hat{\chi}\} = 1, \hat{\psi}^2 = \hat{\chi}^2 = 0.$$

$$\{\chi_m, \psi_n\} = \{\chi_m, \chi_n\} = \{\psi_m, \psi_n\} = 0.$$

$$\int d\theta = 0, \int d\theta\theta = -\int \theta d\theta = 1$$

$$\int d\theta \frac{d}{d\theta}(a+\theta b) = \int d\theta b = 0$$

$$\int d\theta_1 \ldots d\theta_n f(\theta) = c$$

$$f(\theta) = \cdots + \theta_n \theta_{n-1} \ldots \theta_1 c.$$

$$|\psi\rangle = |\downarrow\rangle + |\uparrow\rangle \psi,$$

$$\hat{\psi} |\psi\rangle = |\psi\rangle \psi.$$



$$\langle \psi | \psi' \rangle = \psi - \psi'.$$

$$(\psi - \psi')f(\psi) = (\psi - \psi')f(\psi')$$

$$\int d\psi (\psi - \psi')f(\psi) = f(\psi')$$

$$\langle \psi | \hat{\psi} = -\psi \langle \psi |,$$

$$\int |\psi\rangle d\psi \langle \psi| = 1$$

$$-\int d\chi \exp [\chi(\psi' - \psi)] = \psi - \psi' = \langle \psi | \psi' \rangle$$

$$\langle \psi | \downarrow \rangle = \psi, \langle \psi | \uparrow \rangle = -1$$

$$\langle \psi | \hat{\chi} | \psi' \rangle = \langle \psi | \uparrow \rangle = - \int d\chi \chi \exp [\chi(\psi' - \psi)]$$

$$-\langle \psi | f(\hat{\chi}, \hat{\psi}) | \psi' \rangle = \int d\chi f(\chi, \psi') \exp [\chi(\psi' - \psi)]$$

$$\langle \psi_f, T | \psi_i, 0 \rangle \approx - \int d\chi_0 d\psi_1 d\chi_1 \dots d\psi_{N-1} d\chi_{N-1}$$

$$\times \exp \sum_{m=0}^{N-1} [-iH(\chi_m, \psi_m)\epsilon - \chi_m(\psi_{m+1} - \psi_m)]$$

$$\rightarrow \int [d\chi d\psi] \exp \left\{ \int_0^T dt [-\chi \dot{\psi} - iH(\chi, \psi)] \right\}$$

$$(A): \quad \langle \uparrow \uparrow \uparrow \rangle = \langle \downarrow \downarrow \downarrow \rangle = 1, \langle \uparrow \uparrow \downarrow \rangle = \langle \downarrow \downarrow \uparrow \rangle = 0,$$

$$(B): \quad \langle \uparrow \uparrow \downarrow \rangle = \langle \downarrow \downarrow \uparrow \rangle = 1, \langle \uparrow \uparrow \uparrow \rangle = \langle \downarrow \downarrow \downarrow \rangle = 0.$$

$$\psi(t)\psi(t')=-\psi(t')\psi(t)$$

$$T[\hat{\psi}(t)\hat{\psi}(t')]\equiv \theta(t-t')\hat{\psi}(t)\hat{\psi}(t')-\theta(t'-t)\hat{\psi}(t')\hat{\psi}(t)$$

$$\hat{A}|i\rangle=|j\rangle A_{ji}$$

$$\int d\psi \langle \psi | \hat{A} | \psi \rangle = A_{\downarrow\downarrow} - A_{\uparrow\uparrow} = \text{Tr}[(-1)^{\hat{F}} \hat{A}]$$

$$\int d\psi \langle \psi, U | \psi, 0 \rangle = \text{Tr}[(-1)^{\hat{F}} \exp(-\hat{H}U)]$$

$$\langle 1 \rangle = \int [d\psi d\chi] \exp \left(\int d^d x \chi \Delta \psi \right)$$



$$\psi(x) = \sum_i \psi_i \Psi_i(x), \chi(x) = \sum_i \chi_i \Upsilon_i(x)$$

$$\Delta \Psi_i(x) = \lambda_i \Psi_i(x), \Delta^T \Upsilon_i(x) = \lambda_i \Upsilon_i(x)$$

$$\int d^d x \Upsilon_i(x) \Psi_j(x) = \delta_{ij}$$

$$\langle 1 \rangle = \prod_i \left[\int d\psi_i d\chi_i \exp(\lambda_i \chi_i \psi_i) \right] = \prod_i \lambda_i = \det \Delta$$

$$\int [d\psi] \exp \left(\int d^d x \psi \Delta \psi \right) = (\det \Delta)^{1/2}$$

$$\int_{-\infty}^{\infty} dx \int_{-\infty}^{\infty} dy \exp(2\pi i \lambda xy) = \frac{1}{\lambda} = \left[\int d\psi \int d\chi \exp(\lambda \chi \psi) \right]^{-1}$$

$$\delta(d^m x d^n \theta) = d^m x d^n \theta \left(\sum_i \frac{\partial}{\partial x_i} \delta x_i - \sum_j \frac{\partial}{\partial \theta_j} \delta \theta_j \right)$$

$$\int [dq dq^*]_{\theta} \exp(-S_E)$$

$$S=\frac{1}{4\pi\alpha'}\int d^2\sigma(\partial_1X\partial_1X+\partial_2X\partial_2X+m^2X^2)$$

$$0\leq\sigma_1\leq2\pi,0\leq\sigma_2\leq T$$

$$H=-\frac{1}{2}\frac{\partial^2}{\partial\phi^2}$$

$$\frac{1}{2\kappa^2}\int d^d x (-G)^{1/2} R$$

$$\chi=\sum_{p=0}^d (-1)^p B_p \frac{i}{2\pi\alpha'} \int_M B + i \int_{\partial M} \mathcal{A}$$

CONCLUSIONES.

En mérito a lo antes expuesto, concluyese que, las ecuaciones de campo de Einstein engranan, en tratándose de un espacio – tiempo cuántico, en tanto se esté ante un efecto cuántico – gravitacional ordinario, es decir, cuando el tejido del espacio – tiempo, en dimensión $\mathbb{R}^4$, se deforma, sin que por esto, estemos ante pluridimensiones, como sí supone la supergravedad cuántica, lo que no ocupa este estudio.

Sin embargo, esta deformación, supone la existencia de otros puntos de realidad en una misma



dimensión, en los que impera la simetría de gauge fija o corregida. Por tanto, la curvatura, es puramente local, por lo que, las transformaciones de gauge son locales, sin que existan contrapuntos pluridimensionales, sino pluriespaciales, es decir, interacciones de partículas focalizadas en puntos arbitrarios del espacio de Hilbert – Einstein deformado. Adicionalmente y como ha quedado anotado, concebida así, la gravedad cuántica, ésta puede ser endógena o exógena, según la naturaleza fenomenológica de la partícula supermasiva en un campo cuántico específico y según las magnitudes escalares y simetrías tensoriales.

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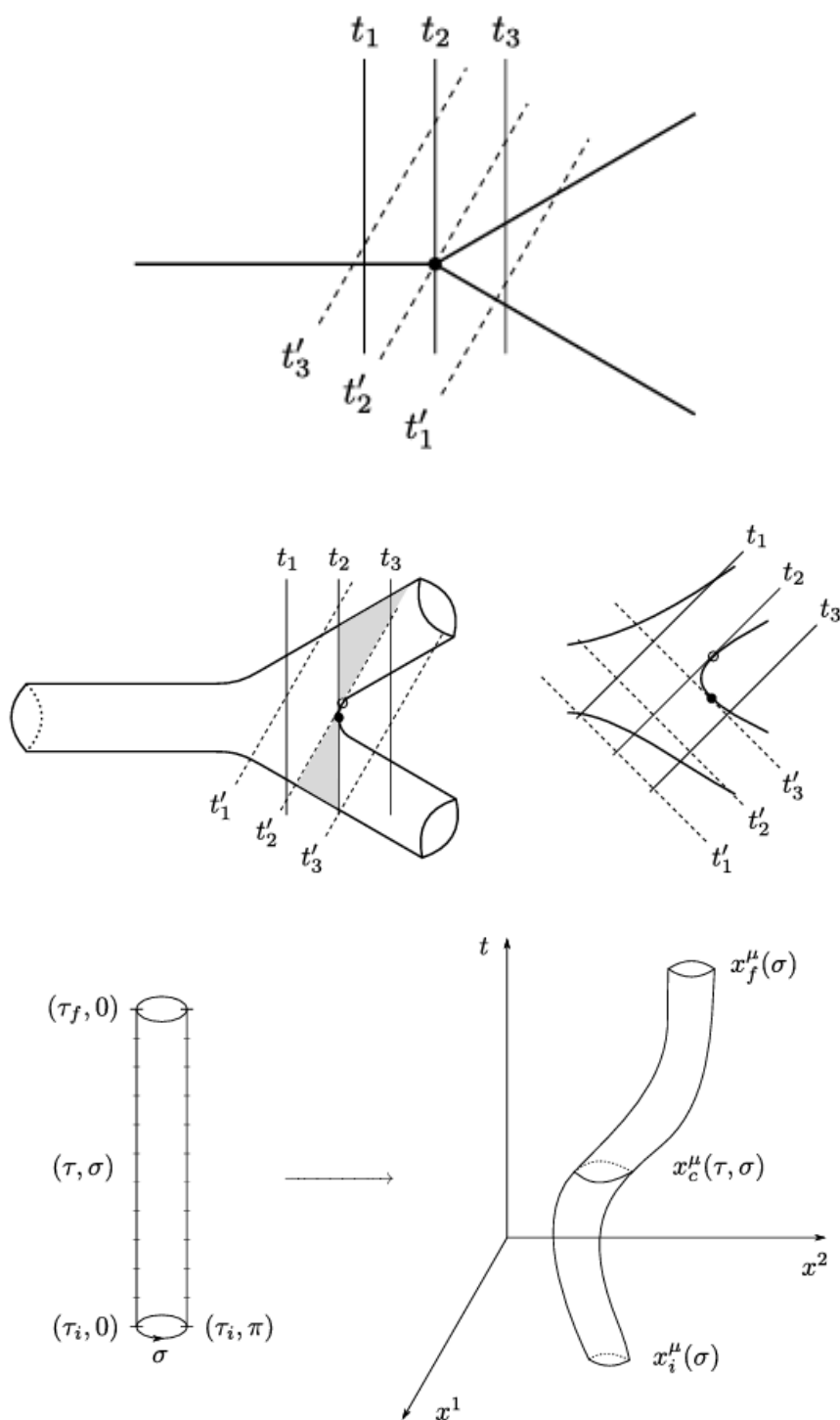
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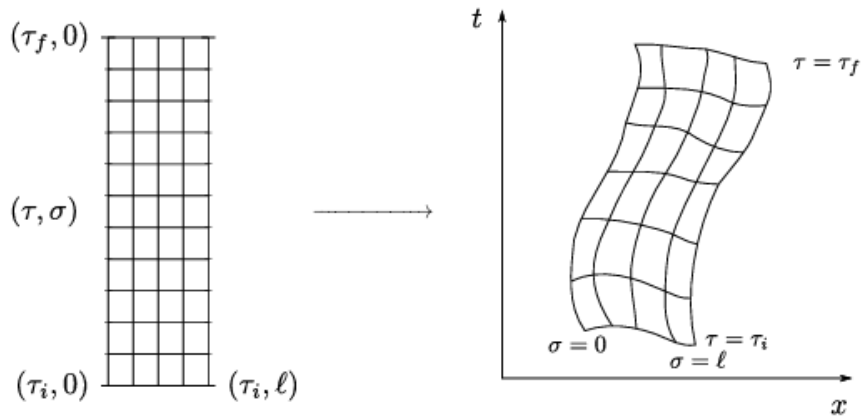
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Anexo A.

Acoplamiento de las ecuaciones de campo de Einstein – Jordan - Weyl a un espacio – tiempo cuántico relativista.





$$S = \int_M d^4x \sqrt{-g} \left(\phi^4 R - \frac{\omega}{\phi} g^{\mu\nu} \partial_\mu \phi \partial_\nu \phi - U(\phi) \right) + 2 \int_{\partial M} d^3x \sqrt{h} \phi K$$

$$R_{\mu\nu} - \frac{1}{2} g_{\mu\nu} R = \frac{\omega}{\phi^2} \left[\partial_\mu \phi \partial_\nu \phi - \frac{1}{2} g_{\mu\nu} g^{\alpha\beta} \partial_\alpha \phi \partial_\beta \phi \right] + \frac{1}{\phi} \left[\nabla_\mu \nabla_\nu \phi - g_{\mu\nu} \square \phi - \frac{1}{2} g_{\mu\nu} U(\phi) \right]$$

$$(3+2\omega)\square\phi=\phi\frac{dU}{d\phi}-2U(\phi).$$

$$\tilde{g}_{\mu\nu}=(16\pi G\phi)g_{\mu\nu},\tilde{\phi}=\phi,$$

$$\tilde{S}=\int_M d^4x\sqrt{-\tilde{g}}\left[\frac{1}{16\pi G}\tilde{R}-\frac{2\omega+3}{(16\pi G)2\phi^2}\tilde{g}^{\mu\nu}\partial_\mu\phi\partial_\nu\phi-V(\phi)\right]+\frac{1}{8\pi G}\int_{\partial M}d^3\sqrt{\tilde{h}}\tilde{K}$$

$$V(\phi)\equiv \frac{U(\phi)}{(16\pi G\phi)^2}$$

$$(g_{\mu\nu}(x),\phi(x))$$

$$(\tilde{g}_{\mu\nu}(x),\phi(x))$$

$$\{\tilde{N}(x),\tilde{\pi}_{\phi}(x')\}=\frac{8\pi GN(x)\delta^{(3)}(x-x')}{\sqrt{16\pi G\phi(x)}}\neq 0.$$

$$g=-N^2(t)dt\otimes dt+a^2(t)dx^i\otimes dx^i$$

$$\mathcal{L}_{FLRW}=-\frac{6a\dot{a}^2}{N}\phi-\frac{6a^2\dot{a}}{N}\dot{\phi}+\frac{\omega a^3}{N\phi}(\dot{\phi})^2-Na^3U(\phi).$$

$$H_T = N \left[-\frac{\omega\pi_a^2}{12a\phi(2\omega+3)} - \frac{\pi_a\pi_\phi}{2a^2(2\omega+3)} + \frac{\phi\pi_\phi^2}{2a^3(2\omega+3)} + a^3 U(\phi) \right] + \lambda_N \pi_N \equiv NH + \lambda_N \pi_N$$

$$\begin{aligned}\dot{N} &\approx \{N, H_T\} \approx \lambda_N, \\ \dot{\pi}_N &\approx \{H, H_T\} \approx -H, \\ \dot{a} &\approx \{a, H_T\} \approx -\frac{N}{2a(2\omega+3)} \left(\frac{\omega\pi_a}{3\phi} + \frac{\pi_\phi}{a} \right), \\ \dot{\pi}_a &\approx \{\pi_a, H_T\} \approx -\frac{N}{2a^2(2\omega+3)} \left(\frac{\omega\pi_a^2}{6\phi} + \frac{2\pi_a\pi_\phi}{a} - \frac{3\phi\pi_\phi^2}{a^2} \right) \\ &\quad - 3Na^2 U(\phi), \\ \dot{\phi} &\approx \{\phi, H_T\} \approx \frac{N}{2a^2(2\omega+3)} \left(-\pi_a + \frac{2\phi\pi_\phi}{a} \right), \\ \dot{\pi}_\phi &\approx \{\pi_\phi, H_T\} \approx -\frac{N}{2a(2\omega+3)} \left(\frac{\omega\pi_a^2}{6\phi^2} + \frac{\pi_\phi^2}{a^2} \right) - Na^3 \frac{dU}{d\phi}.\end{aligned}$$

$$\tilde{g} = -\tilde{N}^2(t)dt \otimes dt + \tilde{a}^2(t)dx^i \otimes dx^i$$

$$\tilde{N} = N(16\pi G\phi)^{\frac{1}{2}}, \tilde{a} = (16\pi G\phi)^{\frac{1}{2}}a, \tilde{\phi} = \phi.$$

$$\tilde{\mathcal{L}}_{FLRW} = -\frac{1}{\tilde{N}(16\pi G\phi)} \left[6\tilde{a}\dot{\tilde{a}}^2\phi - \frac{(2\omega+3)\tilde{a}^3\dot{\phi}^2}{2\phi} \right] - \tilde{N}\tilde{a}^3 V(\phi)$$

$$\begin{aligned}\tilde{H}_T &= \tilde{N}\tilde{a}^3 \left[-\frac{2\pi G\tilde{\pi}_a^2}{3\tilde{a}^4} + \frac{8\pi G\tilde{\pi}_\phi^2\phi^2}{(2\omega+3)\tilde{a}^6} + V(\phi) \right] \\ &+ \tilde{\lambda}_N \tilde{\pi}_N \equiv \tilde{N}\tilde{H} + \tilde{\lambda}_N \tilde{\pi}_N\end{aligned}$$

$$\begin{aligned}\dot{\tilde{N}} &\approx \{\tilde{N}, \tilde{H}_T\} \approx \tilde{\lambda}_N \\ \dot{\tilde{\pi}}_N &\approx \{\tilde{\pi}_N, \tilde{H}_T\} \approx -\tilde{H} \\ \dot{\tilde{a}} &\approx \{\tilde{a}, \tilde{H}_T\} \approx -\tilde{N} \frac{4\pi G\tilde{\pi}_a}{3\tilde{a}} \\ \dot{\tilde{\pi}}_a &\approx \{\tilde{\pi}_a, \tilde{H}_T\} \\ &\approx \tilde{N} \left[-\frac{(2\pi G)\tilde{\pi}_a^2}{3\tilde{a}^2} + \frac{3(8\pi G)\tilde{\pi}_\phi^2\tilde{\phi}^2}{(2\omega+3)\tilde{a}^4} - 3\tilde{a}^2 V(\tilde{\phi}) \right] \\ \dot{\tilde{\phi}} &\approx \{\tilde{\phi}, \tilde{H}_T\} \approx \tilde{N} \frac{(16\pi G)\tilde{\pi}_\phi\tilde{\phi}^2}{(2\omega+3)\tilde{a}^3} \\ \dot{\tilde{\pi}}_\phi &\approx \{\tilde{\pi}_\phi, \tilde{H}_T\} \approx -\tilde{N} \left[\frac{16\pi G\tilde{\pi}_\phi^2\tilde{\phi}}{(2\omega+3)\tilde{a}^3} + \tilde{a}^3 \frac{dV(\tilde{\phi})}{d\tilde{\phi}} \right]\end{aligned}$$

$$\begin{aligned}\tilde{\pi}_a &= \frac{\pi_a}{(16\pi G\phi)^{\frac{1}{2}}}, \tilde{\pi}_\phi = \frac{1}{\phi} \left(\phi\pi_\phi - \frac{1}{2}a\pi_a \right) \\ \tilde{\pi}_N &= \frac{\pi_N}{(16\pi G\phi)^{\frac{1}{2}}}\end{aligned}$$

$$\{\tilde{N}, \tilde{\pi}_N\} = 1, \{\tilde{a}, \tilde{\pi}_a\} = 1, \{\phi, \tilde{\pi}_\phi\} = 1$$



$$\begin{aligned}\{\tilde{N}, \tilde{a}\} &= 0, \{\tilde{N}, \tilde{\pi}_a\} = 0, \{\tilde{N}, \phi\} = 0 \\ \{\tilde{a}, \tilde{\pi}_N\} &= 0, \{\tilde{a}, \phi\} = 0, \{\tilde{a}, \tilde{\pi}_\phi\} = 0 \\ \{\phi, \tilde{\pi}_N\} &= 0, \{\phi, \tilde{\pi}_a\} = 0, \{\tilde{\pi}_N, \tilde{\pi}_a\} = 0\end{aligned}$$

$$\begin{aligned}\{\tilde{N}, \tilde{\pi}_\phi\} &= \left\{ N(16\pi G\phi)^{\frac{1}{2}}, \frac{1}{\phi} \left(\phi\pi_\phi - \frac{1}{2}a\pi_a \right) \right\} \\ &= \frac{8\pi GN}{(16\pi G\phi)^{\frac{1}{2}}} \neq 0\end{aligned}$$

$$\begin{aligned}\{\tilde{\pi}_N, \tilde{\pi}_\phi\} &= \left\{ \frac{\pi_N}{(16\pi G\phi)^{\frac{1}{2}}}, \frac{1}{\phi} \left(\phi\pi_\phi - \frac{1}{2}a\pi_a \right) \right\} \\ &= -\frac{8\pi G\pi_N}{(16\pi G\phi)^{\frac{3}{2}}} \neq 0\end{aligned}$$

$$\begin{aligned}\dot{N} &\approx \frac{\tilde{\lambda}_N}{(16\pi G\phi)^{\frac{1}{2}}} - \frac{N^2}{2a^2(2\omega+3)} \left(\frac{\pi_\phi}{a} - \frac{\pi_a}{2\phi} \right), \\ \dot{\pi}_N &\approx -H + \frac{N\pi_N}{2a^2(2\omega+3)} \left(\frac{\pi_\phi}{a} - \frac{\pi_a}{2\phi} \right), \\ \dot{a} &\approx -\frac{N}{2a(2\omega+3)} \left(\frac{\omega\pi_a}{3\phi} + \frac{\pi_\phi}{a} \right), \\ \dot{\pi}_a &\approx -\frac{N}{2a^2(2\omega+3)} \left(\frac{\omega\pi_a^2}{6\phi} + \frac{2\pi_a\pi_\phi}{a} - \frac{3\phi\pi_\phi^2}{a^2} \right) \\ &\quad - 3Na^2U(\phi), \\ \dot{\phi} &\approx \frac{N}{2a^2(2\omega+3)} \left(-\pi_a + \frac{2\phi\pi_\phi}{a} \right), \\ \dot{\pi}_\phi &\approx -\frac{N}{2a(2\omega+3)} \left(\frac{\omega\pi_a^2}{4\phi^2} - \frac{7\pi_\phi^2}{2a^2} + \frac{\pi_a\pi_\phi}{2a\phi} \right) \\ &\quad - Na^3 \frac{dU}{d\phi} + \frac{Na^3}{2\phi} U(\phi).\end{aligned}$$

$$\dot{\pi}_\phi \approx -\frac{N}{2a(2\omega+3)} \left(\frac{\omega\pi_a^2}{6\phi^2} + \frac{\pi_\phi^2}{a^2} \right) - Na^3 \frac{dU}{d\phi} + \frac{H}{2\phi},$$

$$\tilde{\lambda}_N = (16\pi G\phi)^{\frac{1}{2}} \left[\lambda_N + \frac{N^2}{2a^2(2\omega+3)} \left(\frac{\pi_\phi}{a} - \frac{\pi_a}{2\phi} \right) \right].$$

$$\tilde{\pi}_N = \frac{\pi_N}{(16\pi G\phi)^{\frac{1}{2}}}$$

$$N = c_0$$

$$N - c_0 \approx 0$$

$$\{N - c_0, H\} \approx 0$$



$$\{N - c_0, \pi_N\} \approx 1$$

$$\chi_0 \equiv N - c_0, \text{ and } \chi_1 \equiv \pi_N$$

$$\dot{\chi}_0 \approx \{N - c_0, H_T\} \approx 0,$$

$$\dot{\chi}_1 \approx \{\pi_N, H_T\} \approx 0,$$

$$C_{\alpha\beta} \equiv \{\chi_\alpha, \chi_\beta\}$$

$$C_{\alpha\beta} = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix}, \text{ and } C_{\alpha\beta}^{-1} = \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix}.$$

$$\{\cdot, \cdot\}_{DB} \equiv \{\cdot, \cdot\} - \{\cdot, \chi_\alpha\} C_{\alpha\beta}^{-1} \{\chi_\beta, \cdot\}$$

$$\begin{aligned} \dot{a} &\approx \{a, H_T\}_{DB} \\ &\approx -\frac{N}{2a(2\omega + 3)} \left(\frac{\omega\pi_a}{3\phi} + \frac{\pi_\phi}{a} \right), \\ \dot{\pi}_a &\approx \{\pi_a, H_T\}_{DB} \\ &\approx -\frac{N}{2a^2(2\omega + 3)} \left(\frac{\omega\pi_a^2}{6\phi} + \frac{2\pi_a\pi_\phi}{a} - \frac{3\phi\pi_\phi^2}{a^2} \right) \\ &\quad - 3Na^2U(\phi), \\ \dot{\phi} &\approx \{\phi, H_T\}_{DB} \\ &\approx \frac{N}{2a^2(2\omega + 3)} \left(-\pi_a + \frac{2\phi\pi_\phi}{a} \right), \\ \dot{\pi}_\phi &\approx \{\pi_\phi, H_T\}_{DB} \\ &\approx -\frac{N}{2a(2\omega + 3)} \left(\frac{\omega\pi_a^2}{6\phi^2} + \frac{\pi_\phi^2}{a^2} \right) - Na^3 \frac{dU}{d\phi}. \end{aligned}$$

$$\begin{aligned} \dot{a} &\approx -\frac{c_0}{2a(2\omega + 3)} \left(\frac{\omega\pi_a}{3\phi} + \frac{\pi_\phi}{a} \right), \\ \dot{\pi}_a &\approx -\frac{c_0}{2a^2(2\omega + 3)} \left(\frac{\omega\pi_a^2}{6\phi} + \frac{2\pi_a\pi_\phi}{a} - \frac{3\phi\pi_\phi^2}{a^2} \right) \\ &\quad - 3c_0a^2U(\phi), \\ \dot{\phi} &\approx \frac{c_0}{2a^2(2\omega + 3)} \left(-\pi_a + \frac{2\phi\pi_\phi}{a} \right), \\ \dot{\pi}_\phi &\approx -\frac{c_0}{2a(2\omega + 3)} \left(\frac{\omega\pi_a^2}{6\phi^2} + \frac{\pi_\phi^2}{a^2} \right) - c_0a^3 \frac{dU}{d\phi}. \end{aligned}$$

$$\tilde{N} = c_0(16\pi G\phi)^{\frac{1}{2}},$$

$$\tilde{N} - c_0(16\pi G\phi)^{\frac{1}{2}} \approx 0.$$

$$\left\{ \tilde{N} - c_0(16\pi G\phi)^{\frac{1}{2}}, \tilde{\pi}_N \right\} \approx 1.$$

$$\left\{ \tilde{N} - c_0(16\pi G\phi)^{\frac{1}{2}}, \tilde{H} \right\} = -\frac{c_0(16\pi G\phi)^{3/2}\tilde{\pi}_\phi}{2(2\omega + 3)\tilde{a}^3} \equiv -\eta \neq 0$$



$$\tilde{H}' \equiv \tilde{H} + \eta \tilde{\pi}_N,$$

$$\left\{ \tilde{N} - c_0(16\pi G\phi)^{\frac{1}{2}}, \tilde{H}' \right\} = -\frac{c_0^2(16\pi G)^2\phi\tilde{\pi}_N}{4(2\omega+3)\tilde{a}^3} \approx 0$$

$$\tilde{H}'_T = \tilde{N}\tilde{H}' + \tilde{\lambda}_N\tilde{\pi}_N.$$

$$\tilde{\chi}_0 \equiv \tilde{N} - c_0(16\pi G\phi)^{\frac{1}{2}}, \text{ and } \tilde{\chi}_1 \equiv \tilde{\pi}_N.$$

$$\dot{\tilde{\chi}}_0 \approx \left\{ \tilde{N} - c_0(16\pi G\phi)^{\frac{1}{2}}, \tilde{H}'_T \right\} \approx 0,$$

$$\dot{\tilde{\chi}}_1 \approx \{ \tilde{\pi}_N, \tilde{H}'_T \} \approx 0,$$

$$C_{\alpha\beta} \equiv \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix}, \text{ and } C_{\alpha\beta}^{-1} \equiv \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix}.$$

$$\begin{aligned} \dot{\tilde{a}} &\approx \{ \tilde{a}, \tilde{H}'_T \}_{DB} \approx -\tilde{N} \frac{4\pi G \tilde{\pi}_a}{3\tilde{a}}, \\ \dot{\tilde{\pi}}_a &\approx \{ \tilde{\pi}_a, \tilde{H}'_T \}_{DB} \\ &\approx \tilde{N} \left[-\frac{(2\pi G)\tilde{\pi}_a^2}{3\tilde{a}^2} + \frac{3(8\pi G)\tilde{\pi}_\phi^2\tilde{\phi}^2}{(2\omega+3)\tilde{a}^4} - 3\tilde{a}^2 V(\phi) \right] \\ &\quad - \frac{3c_0(16\pi G\phi)^{3/2}\tilde{\pi}_\phi\tilde{\pi}_N}{(2\omega+3)\tilde{a}^4}, \\ \dot{\phi} &\approx \{ \phi, \tilde{H}'_T \}_{DB} \\ &\approx \tilde{N} \frac{(16\pi G)\tilde{\pi}_\phi\phi^2}{(2\omega+3)\tilde{a}^3} - \frac{c_0(16\pi G\phi)^{3/2}\tilde{\pi}_N}{(2\omega+3)\tilde{a}^3} \\ \dot{\tilde{\pi}}_\phi &\approx \{ \tilde{\pi}_\phi, \tilde{H}'_T \}_{DB} \\ &\approx \{ \tilde{\pi}_\phi, \tilde{H}'_T \} \\ &\quad - \left\{ \tilde{\pi}_\phi, \tilde{N} - c_0(16\pi G\phi)^{\frac{1}{2}} \right\} C_{01}^{-1} \{ \tilde{\pi}_N, \tilde{H}'_T \} \\ &\approx -\tilde{N} \left[\frac{16\pi G\tilde{\pi}_\phi^2\tilde{\phi}}{(2\omega+3)\tilde{a}^3} + \tilde{a}^3 \frac{dV(\phi)}{d\phi} \right] \\ &\quad + \frac{c_0\tilde{\pi}_\phi\tilde{\pi}_N}{(2\omega+3)\tilde{a}^3} \frac{3}{2} (16\pi G\phi)^{1/2} 16\pi G \\ &\quad - \frac{16\pi Gc_0}{2(16\pi G\phi)^{1/2}} \tilde{H}'. \end{aligned}$$

$$\begin{aligned} \dot{\tilde{a}} &\approx -c_0(16\pi G\phi)^{\frac{1}{2}} \frac{4\pi G \tilde{\pi}_a}{3\tilde{a}}, \\ \dot{\tilde{\pi}}_a &\approx c_0(16\pi G\phi)^{\frac{1}{2}} \left[-\frac{(2\pi G)\tilde{\pi}_a^2}{3\tilde{a}^2} + \frac{3(8\pi G)\tilde{\pi}_\phi^2\tilde{\phi}^2}{(2\omega+3)\tilde{a}^4} \right. \\ &\quad \left. - 3\tilde{a}^2 V(\phi) \right] \\ \dot{\phi} &\approx c_0(16\pi G\phi)^{\frac{1}{2}} \frac{(16\pi G)\tilde{\pi}_\phi\phi^2}{(2\omega+3)\tilde{a}^3} \\ \dot{\tilde{\pi}}_\phi &\approx -c_0(16\pi G\phi)^{\frac{1}{2}} \left[\frac{16\pi G\tilde{\pi}_\phi^2\tilde{\phi}}{(2\omega+3)\tilde{a}^3} + \tilde{a}^3 \frac{dV(\phi)}{d\phi} \right]. \end{aligned}$$



$$g = -(N^2 - N_i N^i) dt \otimes dt + N_i (dx^i \otimes dt + dt \otimes dx^i) + h_{ij} dx^i \otimes dx^j,$$

$$\begin{aligned} \mathcal{L}_{ADM} = & \sqrt{h} [N\phi ({}^{(3)}R + K_{ij} K^{ij} - K^2) \\ & - \frac{\omega}{N\phi} (N^2 h^{ij} D_i \phi D_j \phi - (\dot{\phi} - N^i D_i \phi)^2) \\ & + 2K(\dot{\phi} - N^i D_i \phi) - NU(\phi) + 2h^{ij} D_i N D_j \phi] \end{aligned}$$

$$K_{ij} = \frac{1}{2N} \left(-\frac{\partial h_{ij}}{\partial t} + D_i N_j + D_j N_i \right)$$

$$H_T = \int d^3x (\lambda^N \pi_N + \lambda_i \pi^i + N\mathcal{H} + N_i \mathcal{H}^i)$$

$$\begin{aligned} \mathcal{H} = & \sqrt{h} \left\{ \left[-\phi^3 R + \frac{1}{\phi h} \left(\pi^{ij} \pi_{ij} - \frac{\pi_h^2}{2} \right) \right] \right. \\ & + \frac{\omega}{\phi} D_i \phi D^i \phi + 2D^i D_i \phi + U(\phi) \\ & \left. + \frac{1}{2h\phi} \left(\frac{1}{3+2\omega} \right) (\pi_h - \phi \pi_\phi)^2 \right\} \end{aligned}$$

$$\mathcal{H}^i = -2D_j \pi^{ji} + (D^i \phi) \pi_\phi.$$

$$\tilde{g} = -(\tilde{N}^2 - \tilde{N}_i \tilde{N}^i) dt \otimes dt + \tilde{N}_i (dx^i \otimes dt + dt \otimes dx^i) + \tilde{h}_{ij} dx^i \otimes dx^j$$

$$\begin{aligned} \tilde{N} &= N(16\pi G\phi)^{\frac{1}{2}}, \tilde{N}_i = N_i(16\pi G\phi), \\ \tilde{h}_{ij} &= (16\pi G\phi)h_{ij}, \tilde{\phi} = \phi \end{aligned}$$

$$\tilde{\mathcal{L}}_{ADM} = \frac{\sqrt{\tilde{h}}}{16\pi G} [\tilde{N} ({}^{(3)}\tilde{R} + \tilde{K}_{ij} \tilde{K}^{ij} - \tilde{K}^2) - \frac{2\omega + 3}{2\tilde{N}\phi^2} (\tilde{N}^2 \tilde{h}^{ij} \tilde{D}_i \phi \tilde{D}_j \phi - (\dot{\phi} - \tilde{N}^i \tilde{D}_i \phi)^2) - 16\pi G \tilde{N} V(\phi)]$$

$$\tilde{H}_T = \int d^3x (\tilde{\lambda}^N \tilde{\pi}_N + \tilde{\lambda}_i \tilde{\pi}^i + \tilde{N} \tilde{\mathcal{H}} + \tilde{N}_i \tilde{\mathcal{H}}^i)$$

$$\begin{aligned} \tilde{\mathcal{H}} = & \frac{\sqrt{\tilde{h}}}{16\pi G} \left[-{}^{(3)}\tilde{R} + \frac{(16\pi G)^2}{\tilde{h}} \left(\tilde{\pi}^{ij} \tilde{\pi}_{ij} - \frac{\tilde{\pi}_h^2}{2} \right) \right. \\ & \left. + \frac{\left(\omega + \frac{3}{2} \right)}{\phi^2} \tilde{D}_i \phi \tilde{D}^i \phi + \frac{64(\pi G)^2 \phi^2}{\tilde{h} \left(\omega + \frac{3}{2} \right)} \tilde{\pi}_\phi^2 \right] + \sqrt{\tilde{h}} V(\phi) \end{aligned}$$

$$\tilde{\mathcal{H}}^i = -2\tilde{D}_j \tilde{\pi}^{ji} + \tilde{D}^i \phi \tilde{\pi}_\phi,$$



$$\tilde{\pi}_N = \frac{\pi_N}{(16\pi G\phi)^{\frac{1}{2}}}, \tilde{\pi}^i = \frac{\pi^i}{(16\pi G\phi)},$$

$$\tilde{\pi}^{ij} = \frac{\pi^{ij}}{(16\pi G\phi)^{\frac{1}{2}}}, \tilde{\pi}_\phi = \frac{1}{\phi}(\phi\pi_\phi - \pi_h),$$

$$\{\tilde{N}(x), \tilde{\pi}_N(x')\} = \delta^{(3)}(x - x')$$

$$\{\tilde{N}_i(x), \tilde{\pi}^j(x')\} = \delta_i^j \delta^{(3)}(x - x')$$

$$\{\tilde{h}_{ij}(x), \tilde{\pi}^{kl}(x')\} = \delta_i^k \delta_j^l \delta^{(3)}(x - x')$$

$$\{\phi(x), \tilde{\pi}_\phi(x')\} = \delta^{(3)}(x - x')$$

$$\begin{aligned} \{\tilde{N}(x), \tilde{\pi}^j(x')\} &= \{\tilde{N}_i(x), \tilde{\pi}_N(x')\} = \{\tilde{\pi}^i(x), \tilde{\pi}_\phi(x')\} \\ &= \{\tilde{N}(x), \tilde{N}_j(x')\} = \{\tilde{\pi}_N(x), \tilde{\pi}^j(x')\} \\ &= \{\tilde{N}(x), \tilde{h}^{ij}(x')\} = \{\tilde{N}(x), \tilde{\pi}^{ij}(x')\} = \{\tilde{\pi}^i(x), \phi(x')\} \\ &= \{\tilde{N}(x), \phi(x')\} = \{\tilde{N}_i(x), \tilde{h}^{ij}(x')\} = \{\tilde{N}_h(x), \tilde{\pi}^{ij}(x')\} \\ &= \{\tilde{N}_i(x), \phi(x')\} = \{\tilde{\pi}_N(x), \tilde{h}_{ij}(x')\} = \{\tilde{\pi}_N(x), \tilde{\pi}^{ij}(x')\} \\ &= \{\tilde{\pi}^i(x), \tilde{\pi}^{ij}(x')\} = \{\tilde{\pi}^i(x), \tilde{h}_{ij}(x')\} = \{\tilde{\pi}_N(x), \phi(x')\} \\ &= \{\tilde{h}_{ij}(x), \phi(x')\} = \{\tilde{h}_{ij}(x), \tilde{\pi}_\phi(x')\} = \{\tilde{\pi}^{ij}(x), \phi(x')\} \\ &= \{\tilde{\pi}^{ij}(x), \tilde{\pi}_\phi(x')\} = 0, \end{aligned}$$

$$\{\tilde{N}(x), \tilde{\pi}_\phi(x')\} = \frac{8\pi GN(x)\delta^{(3)}(x - x')}{\sqrt{16\pi G\phi(x)}} \neq 0$$

$$\{\tilde{N}_i(x), \tilde{\pi}_\phi(x')\} = 16\pi GN_i(x)\delta^{(3)}(x - x') \neq 0$$

$$\{\tilde{\pi}_N(x), \tilde{\pi}_\phi(x')\} = -\frac{8\pi G\pi_N(x)\delta^{(3)}(x - x')}{\sqrt{(16\pi G\phi(x))^3}} \neq 0$$

$$\{\tilde{\pi}^i(x), \tilde{\pi}_\phi(x')\} = -\frac{\pi^i(x)}{(16\pi G\phi^2)}\delta^{(3)}(x - x') \neq 0$$

$$N = c_0, N_i = c_i$$

$$N - c_0 \approx 0, N_i - c_i \approx 0$$

$$\pi_N \approx 0, \pi^i \approx 0$$

$$\{N(x) - c_0, \pi_N(x')\} \approx \delta^{(3)}(x - x')$$

$$\{N_i(x) - c_i, \pi^j(x')\} = \delta_i^j \delta^{(3)}(x - x')$$

$$\{N - c_0, \mathcal{H}\} \approx 0, \quad \{N_0 - c_0, \mathcal{H}^i\} = 0$$

$$\{N_i - c_i, \mathcal{H}\} \approx 0, \quad \{N_i - c_i, \mathcal{H}^i\} = 0$$

$$\chi_0 \equiv N - c_0, \chi_i \equiv N_i - c_i$$

$$\chi_4 \equiv \pi_N, \chi_{i+4} \equiv \pi_i$$

$$\dot{\chi}_0 \approx \{N - c_0, H_T\} \approx 0,$$

$$\dot{\chi}_i \approx \{N_i - c_i, H_T\} \approx 0,$$



$$\begin{aligned}\dot{\chi}_4 &\approx \{\pi_N, H_T\} \approx 0, \\ \dot{\chi}_{i+4} &\approx \{\pi^i, H_T\} \approx 0,\end{aligned}$$

$$C_{\alpha\beta}^{-1} \equiv \begin{pmatrix} 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & -\mathbb{I} \\ 1 & 0 & 0 & 0 \\ 0 & \mathbb{I} & 0 & 0 \end{pmatrix}$$

$$\begin{aligned}\{\mathcal{H}_i(x), \mathcal{H}_j(x')\}_{DB} &= \{\mathcal{H}_i(x), \mathcal{H}_j(x')\} \\ \{\mathcal{H}_i(x), \mathcal{H}(x')\}_{DB} &= \{\mathcal{H}_i(x), \mathcal{H}(x')\} \\ \{\mathcal{H}(x), \mathcal{H}(x')\}_{DB} &= \{\mathcal{H}(x), \mathcal{H}(x')\}\end{aligned}$$

$$\begin{aligned}h_{ij} &\approx \{h_{ij}, H_T\}_{DB} \\ &\approx D_i N_j + D_j N_i + \frac{2N}{\phi\sqrt{h}}\left(\pi_{ij} - \frac{\pi_h}{2}h_{ij}\right) \\ &\quad + \frac{N}{\phi\sqrt{h}}\frac{(\pi_h - \phi\pi_\phi)}{(2\omega + 3)}h_{ij}.\end{aligned}$$

$$\begin{aligned}\dot{\pi}^{ij} &\approx \{\dot{\pi}^{ij}, H_T\}_{DB} \\ &\approx -N\sqrt{h}\phi\left({}^{(3)}R^{ij} - \frac{h^{ij}}{2}{}^{(3)}R\right) \\ &\quad + \sqrt{h}(D^iD^j - h^{ij}D^kD_k)(N\phi) \\ &\quad + \frac{Nh^{ij}}{2\phi\sqrt{h}}\left(\pi^{ij}\pi_{ij} - \frac{\pi_h^2}{2}\right) + \frac{N}{\phi\sqrt{h}}\pi_h\pi^{ij} \\ &\quad - \sqrt{h}\frac{N\omega}{\phi}D^i\phi D^j\phi - \sqrt{h}\frac{N\omega}{2\phi}h^{ij}D_k\phi D^k\phi \\ &\quad + \sqrt{h}h^{ij}D_kND^k\phi \\ &\quad + \sqrt{h}(D^iND^j\phi + D^jND^i\phi) - \frac{\sqrt{h}}{2}h^{ij}NU(\phi) \\ &\quad + \frac{h^{ij}}{4\sqrt{h}\phi}\left(\frac{N}{2\omega + 3}\right)(\pi_h - \phi\pi_\phi)^2 \\ &\quad - \frac{\pi^{ij}}{\sqrt{h}\phi}\left(\frac{N}{2\omega + 3}\right)(\pi_h - \phi\pi_\phi) \\ &\quad + D_k(N^k\pi^{ij}) - \frac{2N}{\sqrt{h}\phi}\pi^{iq}\pi_q^j \\ &\quad - (D_kN^i)\pi^{kj} - (D_kN^j)\pi^{ki}\end{aligned}$$

$$\begin{aligned}\dot{\phi} &\approx \{\phi, H_T\}_{DB} \\ &\approx N^iD_i\phi - \frac{N}{\sqrt{h}(2\omega + 3)}(\pi_h - \phi\pi_\phi)\end{aligned}$$



$$\begin{aligned}
\dot{\pi}_\phi &\approx \{\pi_\phi, H_T\}_{DB} \\
&\approx \sqrt{h} N^{(3)} R + \frac{N}{\phi^2 \sqrt{h}} \left(\pi^{ij} \pi_{ij} - \frac{\pi_h^2}{2} \right) \\
&\quad + \frac{\sqrt{h} N \omega}{\phi^2} D_i \phi D^i \phi + 2 \sqrt{h} D_i \left(\frac{N \omega}{\phi} D^i \phi \right) \\
&\quad - 2 \sqrt{h} (D^i D_i)(N) - N \sqrt{h} \frac{dU}{d\phi} + D_i (N^i \pi_\phi) \\
&\quad + \frac{N(\pi_h - \phi \pi_\phi)^2}{2 \sqrt{h} \phi^2 (2\omega + 3)} + \frac{N(\pi_h - \phi \pi_\phi) \pi_\phi}{\sqrt{h} \phi (2\omega + 3)} \\
\dot{h}_{ij} &\approx \frac{2c_0}{\phi \sqrt{h}} \left(\pi_{ij} - \frac{\pi_h}{2} h_{ij} \right) + \frac{c_0}{\phi \sqrt{h}} \frac{(\pi_h - \phi \pi_\phi)}{(2\omega + 3)} h_{ij} \\
\dot{\pi}^{ij} &\approx -c_0 \sqrt{h} \phi \left({}^{(3)}R^{ij} - \frac{h^{ij}}{2} {}^{(3)}R \right) \\
&\quad + c_0 \sqrt{h} (D^i D^j - h^{ij} D^k D_k)(\phi) \\
&\quad + \frac{c_0 h^{ij}}{2 \phi \sqrt{h}} \left(\pi^{ij} \pi_{ij} - \frac{\pi_h^2}{2} \right) + \frac{c_0}{\phi \sqrt{h}} \pi_h \pi^{ij} \\
&\quad - \sqrt{h} \frac{c_0 \omega}{\phi} D^i \phi D^j \phi - \sqrt{h} \frac{c_0 \omega}{2 \phi} h^{ij} D_k \phi D^k \phi \\
&\quad - \frac{\sqrt{h}}{2} h^{ij} c_0 U(\phi) \\
&\quad + \frac{h^{ij}}{4 \sqrt{h} \phi} \left(\frac{c_0}{2\omega + 3} \right) (\pi_h - \phi \pi_\phi)^2 \\
&\quad - \frac{\pi^{ij}}{\sqrt{h} \phi} \left(\frac{c_0}{2\omega + 3} \right) (\pi_h - \phi \pi_\phi) \\
&\quad + c^k D_k (\pi^{ij}) - \frac{2c_0}{\sqrt{h} \phi} \pi^{iq} \pi_q^j \\
\dot{\phi} &\approx c^i D_i \phi - \frac{c_0}{\sqrt{h} (2\omega + 3)} (\pi_h - \phi \pi_\phi) \\
\dot{\pi}_\phi &\approx c_0 \sqrt{h} {}^{(3)}R + \frac{c_0}{\phi^2 \sqrt{h}} \left(\pi^{ij} \pi_{ij} - \frac{\pi_h^2}{2} \right) \\
&\quad + c_0 \frac{\sqrt{h} \omega}{\phi^2} D_i \phi D^i \phi + 2 c_0 \sqrt{h} D_i \left(\frac{\omega}{\phi} D^i \phi \right) \\
&\quad - c_0 \sqrt{h} \frac{dU}{d\phi} + c^i D_i (\pi_\phi) \\
&\quad + \frac{c_0 (\pi_h - \phi \pi_\phi)^2}{2 \sqrt{h} \phi^2 (2\omega + 3)} + \frac{c_0 (\pi_h - \phi \pi_\phi) \pi_\phi}{\sqrt{h} \phi (2\omega + 3)} \\
\tilde{N} &= c_0 (16\pi G \phi)^{\frac{1}{2}}, \tilde{N}_i = c_i (16\pi G \phi). \\
\tilde{N} - c_0 (16\pi G \phi)^{\frac{1}{2}} &\approx 0, \tilde{N}_i - c_i (16\pi G \phi) \approx 0
\end{aligned}$$

$$\begin{aligned}\left\{\tilde{N}(x) - c_0(16\pi G\phi)^{\frac{1}{2}}, \tilde{\pi}_N(x')\right\} &\approx \delta^{(3)}(x - x') \\ \left\{\tilde{N}_i(x) - c_i(16\pi G\phi(x)), \tilde{\pi}^j(x')\right\} &\approx \delta_i^j \delta^{(3)}(x - x')\end{aligned}$$

$$\begin{aligned}&\left\{\tilde{N}(x) - c_0(16\pi G\phi(x))^{\frac{1}{2}}, \tilde{\mathcal{H}}(x')\right\} \\ &\approx -\frac{(16\pi G\phi(x))^{\frac{3}{2}}c_0\tilde{\pi}_\phi(x)\delta^{(3)}(x - x')}{4\sqrt{\tilde{h}(x)}\left(\omega + \frac{3}{2}\right)} \\ &\left\{\tilde{N}(x) - c_0(16\pi G\phi(x))^{\frac{1}{2}}, \tilde{\mathcal{H}}^i(x')\right\} \\ &\approx -\frac{8\pi G\phi(x)c_0D^i\phi(x)\delta^{(3)}(x - x')}{(16\pi G\phi(x))^{\frac{1}{2}}} \\ &\left\{\tilde{N}_i(x) - c_i(16\pi G\phi(x)), \tilde{\mathcal{H}}(x')\right\} \\ &\approx -\frac{128c_i\tilde{\pi}_\phi(x)\delta^{(3)}(x - x')(\pi G\phi(x))^2}{\sqrt{\tilde{h}(x)}\left(\omega + \frac{3}{2}\right)} \\ &\left\{\tilde{N}_i(x) - c_i(16\pi G\phi(x)), \tilde{\mathcal{H}}^j(x')\right\} \\ &\approx -(16\pi G)c_iD^j\phi(x)\delta^{(3)}(x - x')\end{aligned}$$

$$\begin{aligned}\tilde{\mathcal{H}}' &\equiv \tilde{\mathcal{H}} + \eta_N\tilde{\pi}_N + \gamma_i\tilde{\pi}^i \\ \tilde{\mathcal{H}}^{i'} &\equiv \tilde{\mathcal{H}}^i + \eta^i\tilde{\pi}_N + \rho_k^i\tilde{\pi}^k\end{aligned}$$

$$\begin{aligned}\eta_N &\equiv \frac{(16\pi G\phi(x))^{\frac{3}{2}}c_0\tilde{\pi}_\phi(x)}{4\sqrt{\tilde{h}(x)}\left(\omega + \frac{3}{2}\right)} \\ \eta^i &\equiv \frac{8\pi G\phi(x)c_0D^i\phi(x)}{(16\pi G\phi(x))^{\frac{1}{2}}} \\ \gamma_i &\equiv \frac{128c_i\tilde{\pi}_\phi(x)(\pi G\phi(x))^2}{\sqrt{\tilde{h}(x)}\left(\omega + \frac{3}{2}\right)} \\ \rho_k^i &\equiv (16\pi G)c_kD^i\phi(x)\end{aligned}$$

$$\begin{aligned}\tilde{\chi}_0 &\equiv \tilde{N} - c_0(16\pi G\phi)^{\frac{1}{2}}, \tilde{\chi}_i \equiv \tilde{N}_i - c_i(16\pi G\phi) \\ \tilde{\chi}_4 &\equiv \tilde{\pi}_N, \tilde{\chi}_{i+4} \equiv \tilde{\pi}^i\end{aligned}$$

$$\tilde{C}_{\alpha\beta}^{-1} \equiv \begin{pmatrix} 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & -\mathbb{I} \\ 1 & 0 & 0 & 0 \\ 0 & \mathbb{I} & 0 & 0 \end{pmatrix}$$

$$\tilde{H}'_T = \int d^3x (\tilde{\lambda}^N\tilde{\pi}_N + \tilde{\lambda}_i\tilde{\pi}^i + \tilde{N}\tilde{\mathcal{H}}' + \tilde{N}_i\tilde{\mathcal{H}}^{i'})$$

$$\dot{\tilde{\chi}}_0(x) \approx \left\{\tilde{N}(x) - c_0(16\pi G\phi(x))^{\frac{1}{2}}, \tilde{H}'_T\right\} \approx 0$$

$$\tilde{\lambda}^N(x) \approx 0$$

$$\dot{\tilde{\chi}}_i(x) \approx \{\tilde{N}_i(x) - c_i(16\pi G\phi(x)), \tilde{H}'_T\} \approx 0$$

$$\tilde{\lambda}_i(x) \approx 0.$$

$$\begin{aligned}\{\tilde{\mathcal{H}}'_i(x), \tilde{\mathcal{H}}'_j(x')\}_{DB} &= \{\tilde{\mathcal{H}}'_i(x), \tilde{\mathcal{H}}'_j(x')\}, \\ \{\tilde{\mathcal{H}}'_i(x), \tilde{\mathcal{H}}'(x')\}_{DB} &= \{\tilde{\mathcal{H}}'_i(x), \tilde{\mathcal{H}}'(x')\}, \\ \{\tilde{\mathcal{H}}'(x), \tilde{\mathcal{H}}'(x')\}_{DB} &= \{\tilde{\mathcal{H}}'(x), \tilde{\mathcal{H}}'(x')\}.\end{aligned}$$

$$\begin{aligned}\dot{\tilde{h}}_{ij} &\approx \{\tilde{h}_{ij}, \tilde{H}'_{ADM}\}_{DB} \\ &\approx \tilde{D}_i \tilde{N}_j + \tilde{D}_j \tilde{N}_i + \frac{(32\pi G)\tilde{N}}{\sqrt{\tilde{h}}} \left(\tilde{\pi}_{ij} - \frac{\tilde{\pi}_h}{2} \tilde{h}_{ij} \right)\end{aligned}$$

$$\begin{aligned}\dot{\tilde{\pi}}^{ij} &\approx \{\tilde{\pi}^{ij}, \tilde{H}'_{ADM}\}_{DB} \\ &\approx -\frac{\tilde{N}\sqrt{\tilde{h}}}{16\pi G} \left({}^{(3)}\tilde{R}^{ij} - \frac{1}{2}(3)\tilde{R}^{ij} \right) \\ &\quad + \frac{(16\pi G)\tilde{N}}{2\sqrt{\tilde{h}}} \tilde{h}^{ij} \left(\tilde{\pi}^{pk} \tilde{\pi}_{pk} - \frac{\tilde{\pi}_h^2}{2} \right) \\ &\quad - 2 \frac{(16\pi G)\tilde{N}}{\sqrt{\tilde{h}}} \left(\tilde{\pi}^{ik} \tilde{\pi}_k^j - \frac{1}{2} \tilde{\pi}_h \tilde{\pi}^{ij} \right) - \tilde{h}^{ij} \frac{\sqrt{\tilde{h}}\tilde{N}}{2} V(\phi) \\ &\quad + \frac{\sqrt{\tilde{h}}}{16\pi G} (\tilde{D}^i \tilde{D}^j \tilde{N} - \tilde{h}^{ij} \tilde{D}^k \tilde{D}_k \tilde{N}) \\ &\quad + \sqrt{\tilde{h}} \tilde{D}_k \left(\frac{1}{\sqrt{\tilde{h}}} \tilde{N}^k \tilde{\pi}^{ij} \right) - \tilde{\pi}^{ki} \tilde{D}_k \tilde{N}^j - \tilde{\pi}^{kj} \tilde{D}_k \tilde{N}^i\end{aligned}$$

$$\begin{aligned}\dot{\phi} &\approx \{\phi, \tilde{H}'_{ADM}\}_{DB} \\ &\approx \frac{8(\pi G)\phi^2}{\sqrt{\tilde{h}} \left(\omega + \frac{3}{2} \right)} \tilde{\pi}_\phi \tilde{N} + \tilde{N}_i \tilde{D}^i \phi\end{aligned}$$

$$\begin{aligned}\dot{\tilde{\pi}}_\phi &\approx \{\tilde{\pi}_\phi, \tilde{H}'_{ADM}\}_{DB} \\ &\approx -\frac{\sqrt{\tilde{h}}}{8\pi G} \left(\omega + \frac{3}{2} \right) \frac{\tilde{N}}{\phi^3} \tilde{D}^i \phi \tilde{D}_i \phi \\ &\quad + \frac{\sqrt{\tilde{h}}}{8\pi G} \frac{\left(\omega + \frac{3}{2} \right)}{\phi^2} \tilde{D}^i \tilde{N} \tilde{D}_i \phi \\ &\quad + \frac{\sqrt{\tilde{h}}}{8\pi G} \left(\omega + \frac{3}{2} \right) \frac{\tilde{N}}{\phi^2} \tilde{D}^i \tilde{D}_i \phi \\ &\quad - \frac{8\tilde{N}(\pi G)\phi \tilde{\pi}_\phi^2}{\sqrt{\tilde{h}} \left(\omega + \frac{3}{2} \right)} + \tilde{D}^i (\tilde{N}_i \tilde{\pi}_\phi) \\ &\quad - \sqrt{\tilde{h}} \tilde{N} \frac{dV(\phi)}{d\phi}\end{aligned}$$

$$\dot{\tilde{h}}_{ij} \approx 16\pi G(c_j \tilde{D}_i \phi + c_i \tilde{D}_j \phi),$$

$$\begin{aligned} \dot{\tilde{\pi}}^{ij} \approx & -\frac{c_0(16\pi G\phi)^{\frac{1}{2}}\sqrt{\tilde{h}}}{16\pi G} \left({}^{(3)}\tilde{R}^{ij} - \frac{1}{2}(3)\tilde{R}^{ij} \right) \\ & + \frac{(16\pi G)c_0(16\pi G\phi)^{\frac{1}{2}}}{2\sqrt{\tilde{h}}} \tilde{h}^{ij} \left(\tilde{\pi}^{pk} \tilde{\pi}_{pk} - \frac{\tilde{\pi}_h^2}{2} \right) \\ & - 2 \frac{(16\pi G)c_0(16\pi G\phi)^{\frac{1}{2}}}{\sqrt{\tilde{h}}} \left(\tilde{\pi}^{ik} \tilde{\pi}_k^j - \frac{1}{2} \tilde{\pi}_h \tilde{\pi}^{ij} \right) \\ & - \tilde{h}^{ij} \frac{\sqrt{\tilde{h}}c_0(16\pi G\phi)^{\frac{1}{2}}}{2} V(\phi) \\ & + \frac{c_0\sqrt{\tilde{h}}}{2(16\pi G\phi)^{\frac{1}{2}}} \left[-\frac{\tilde{D}^i \phi \tilde{D}^j \phi}{2\phi} + \tilde{D}^i \tilde{D}^j \phi \right. \\ & \left. - \tilde{h}^{ij} \left(-\frac{\tilde{D}^k \phi \tilde{D}_k \phi}{2\phi} + \tilde{D}^k \tilde{D}_k \phi \right) \right] \\ & + (16\pi G)\sqrt{\tilde{h}}c^k \tilde{D}_k \left(\frac{1}{\sqrt{\tilde{h}}} \phi \tilde{\pi}^{ij} \right) \\ & - (16\pi G)\tilde{\pi}^{ki} c^j \tilde{D}_k \phi - (16\pi G)\tilde{\pi}^{kj} c^i \tilde{D}_k \phi \end{aligned}$$

$$\begin{aligned} \dot{\phi} \approx & \frac{c_0(16\pi G\phi)^{\frac{3}{2}}}{2\sqrt{\tilde{h}}\left(\omega + \frac{3}{2}\right)} \tilde{\pi}_\phi \phi + (16\pi G\phi)c_i \tilde{D}^i \phi \\ \dot{\tilde{\pi}}_\phi \approx & -\frac{c_0\sqrt{\tilde{h}}}{(16\pi G\phi)^{\frac{1}{2}}\phi^2} \left(\omega + \frac{3}{2} \right) \tilde{D}^i \phi \tilde{D}_i \phi \\ & + \frac{2c_0\sqrt{\tilde{h}}}{(16\pi G\phi)^{\frac{1}{2}}\phi} \left(\omega + \frac{3}{2} \right) \tilde{D}^i \tilde{D}_i \phi \\ & - \frac{c_0(16\pi G\phi)^{\frac{3}{2}}\tilde{\pi}_\phi^2}{2\sqrt{\tilde{h}}\left(\omega + \frac{3}{2}\right)} + 16\pi Gc_i \tilde{D}^i (\phi \tilde{\pi}_\phi) \\ & - c_0\sqrt{\tilde{h}}(16\pi G\phi)^{\frac{1}{2}} \frac{dV(\phi)}{d\phi} \end{aligned}$$

$$H = \frac{p^2}{2m} + \frac{m\omega^2}{2} q^2$$

$$\begin{aligned} q &= \sqrt{\frac{2P}{m\omega}} \sin Q \\ p &= \sqrt{2m\omega P} \cos Q \end{aligned}$$

$$H = \omega P$$

$$P = \frac{E}{\omega}$$



$$\dot{Q} = \frac{\partial H}{\partial P} = \omega$$

$$Q = \omega t + \alpha$$

$$q(t) = \sqrt{\frac{2E}{m\omega^2}} \sin(\omega t + \alpha)$$

$$\tilde{N}^* = N, \tilde{\pi}_{N^*} = \pi_N$$

$$\tilde{N}_i^* = N_i, \tilde{\pi}^{*i} = \pi^i$$

$$\tilde{h}_{ij}^* = (16\pi G\phi)h_{ij}, \tilde{\pi}^{*ij} = \frac{\pi^{ij}}{(16\pi G\phi)^{\frac{1}{2}}}$$

$$\tilde{\phi}^* = \phi, \tilde{\pi}_{\phi}^* = \frac{1}{\phi}(\phi\pi_{\phi} - \pi_h)$$

$$ds^2 = -dt^2 + \lambda^2 dx^2$$

$$C_{\phi} \equiv \frac{1}{2} a \pi_a - \phi \pi_{\phi}$$

$$\begin{aligned} H_{\text{T}}^{(-3/2)} &= N \left[-\frac{\pi_a^2}{24a\phi} + a^3 U(\phi) \right] + \lambda_N \pi_N + \lambda_{\phi} C_{\phi} \\ &= N H^{(-3/2)} + \lambda_N \pi_N + \lambda_{\phi} C_{\phi} \end{aligned}$$

$$\dot{\chi}_1 \approx \left\{ N - c_0, H_{\text{T}}^{(-3/2)} \right\} \approx 0,$$

$$\dot{\chi}_2 \approx \left\{ \pi_N, H_{\text{T}}^{(-3/2)} \right\} \approx 0$$

$$\dot{a} \approx \left\{ a, H_{\text{T}}^{(-3/2)} \right\}_{DB} \approx -\frac{N\pi_a}{12a\phi} + \frac{\lambda_{\phi}a}{2},$$

$$\begin{aligned} \dot{\pi}_a &\approx \left\{ \pi_a, H_{\text{T}}^{(-3/2)} \right\}_{DB} \\ &\approx -\frac{N\pi_a^2}{24\phi a^2} - 3Na^2 U(\phi) - \frac{\lambda_{\phi}\pi_a}{2}, \end{aligned}$$

$$\dot{\phi} \approx \left\{ \phi, H_{\text{T}}^{(-3/2)} \right\}_{DB} \approx -\lambda_{\phi}\phi,$$

$$\begin{aligned} \dot{\pi}_{\phi} &\approx \left\{ \pi_{\phi}, H_{\text{T}}^{(-3/2)} \right\}_{DB} \\ &\approx -N \frac{\pi_a^2}{24a\phi^2} - Na^3 \frac{dU(\phi)}{d\phi} + \lambda_{\phi}\pi_{\phi} \end{aligned}$$

$$\approx -N \frac{\pi_a^2}{24a\phi^2} - \frac{2Na^3 U(\phi)}{\phi} + \lambda_{\phi}\pi_{\phi},$$

$$\phi \frac{dU(\phi)}{d\phi} = 2U(\phi),$$

$$\begin{aligned}\dot{a} &\approx -\frac{c_0\pi_a}{12a\phi} + \frac{\lambda_\phi a}{2} \\ \dot{\pi}_a &\approx -\frac{c_0\pi_a^2}{24\phi a^2} - 3c_0a^2U(\phi) - \frac{\lambda_\phi\pi_a}{2} \\ \dot{\phi} &\approx -\lambda_\phi\phi \\ \dot{\pi}_\phi &\approx -c_0\frac{\pi_a^2}{24a\phi^2} - \frac{2c_0a^3U(\phi)}{\phi} + \lambda_\phi\pi_\phi\end{aligned}$$

$$\begin{aligned}\tilde{H}_T^{(-3/2)} &= \tilde{N}\tilde{a}^3\left[-\frac{2\pi G\tilde{\pi}_a^2}{3\tilde{a}^4} + V(\tilde{\phi})\right] + \tilde{\lambda}_N\tilde{\pi}_N + \lambda_\phi\tilde{C}_\phi \\ &= \tilde{N}\tilde{H}^{(-3/2)} + \tilde{\lambda}_N\tilde{\pi}_N + \tilde{\lambda}_\phi\tilde{C}_\phi\end{aligned}$$

$$\tilde{C}_\phi = -\phi\tilde{\pi}_\phi.$$

$$\tilde{N} - c_0(16\pi G\phi)^{\frac{1}{2}} \approx 0,$$

$$\left\{\tilde{H}^{(-3/2)}, \tilde{N} - c_0(16\pi G\phi)^{\frac{1}{2}}\right\} = 0$$

$$\left\{\tilde{H}^{(-3/2)}, \tilde{C}_\phi\right\} = -\tilde{a}^3\phi\frac{dV(\phi)}{d\phi} = 0$$

$$\left\{\tilde{N} - c_0(16\pi G\phi)^{\frac{1}{2}}, \tilde{C}_\phi\right\} = \frac{c_0(16\pi G\phi)^{\frac{1}{2}}}{2}$$

$$\tilde{C}_\phi - \frac{c_0(16\pi G\phi)^{\frac{1}{2}}}{2}\tilde{\pi}_N,$$

$$\left\{\tilde{N} - c_0(16\pi G\phi)^{\frac{1}{2}}, \tilde{C}_\phi - \frac{c_0(16\pi G\phi)^{\frac{1}{2}}}{2}\tilde{\pi}_N\right\} \approx 0$$

$$\begin{aligned}\tilde{H}_T'^{(-3/2)} &= \tilde{N}\tilde{H}^{(-3/2)} + \tilde{\lambda}_N\tilde{\pi}_N \\ &\quad + \tilde{\lambda}_\phi\left[-\phi\tilde{\pi}_\phi - \frac{c_0(16\pi G\phi)^{\frac{1}{2}}}{2}\tilde{\pi}_N\right].\end{aligned}$$

$$\dot{\chi}_0 \approx \left\{\tilde{N} - c_0(16\pi G\phi)^{\frac{1}{2}}, \tilde{H}_T'^{(-3/2)}\right\} \approx 0$$

$$\dot{\chi}_1 \approx \left\{\tilde{\pi}_N, \tilde{H}_T'^{(-3/2)}\right\} \approx 0$$

$$\begin{aligned}
\dot{\tilde{a}} &\approx \left\{ \tilde{a}, \tilde{H}_T'^{(-3/2)} \right\}_{DB} \approx -\tilde{N} \frac{4\pi G \tilde{\pi}_a}{3\tilde{a}} \\
\dot{\pi}_a &\approx \left\{ \tilde{\pi}_a, \tilde{H}_T'^{(-3/2)} \right\}_{DB} \\
&\approx \tilde{N} \left[-\frac{(2\pi G) \tilde{\pi}_a^2}{3\tilde{a}^2} - 3\tilde{a}^2 V(\phi) \right] \\
\dot{\tilde{\phi}} &\approx \left\{ \tilde{\phi}, \tilde{H}_T'^{(-3/2)} \right\}_{DB} \approx -\tilde{\lambda}_\phi \tilde{\phi} \\
\dot{\tilde{\pi}}_\phi &\approx \left\{ \tilde{\pi}_\phi, \tilde{H}_T'^{(-3/2)} \right\}_{DB} \\
&\approx \left\{ \tilde{\pi}_\phi, \tilde{H}_T'^{(-3/2)} \right\} \\
&- \left\{ \tilde{\pi}_\phi, \tilde{N} - c_0 (16\pi G \phi)^{\frac{1}{2}} \right\} C_{01}^{-1} \left\{ \tilde{\pi}_N, \tilde{H}_T'^{(-3/2)} \right\} \\
&\approx -\tilde{N} \tilde{a}^3 \frac{dV(\phi)}{d\phi} + \tilde{\lambda}_\phi \tilde{\pi}_\phi - \frac{\tilde{\lambda}_\phi c_0}{2} \frac{16\pi G}{(16\pi G \phi)^{\frac{1}{2}}} \tilde{\pi}_N \\
&\quad + \frac{c_0}{2} \frac{16\pi G}{(16\pi G \phi)^{\frac{1}{2}}} \tilde{H}^{(-3/2)} \\
\dot{\tilde{a}} &\approx -c_0 (16\pi G \phi)^{\frac{1}{2}} \frac{4\pi G \tilde{\pi}_a}{3\tilde{a}}, \\
\dot{\tilde{\pi}}_a &\approx c_0 (16\pi G \phi)^{\frac{1}{2}} \left[-\frac{(2\pi G) \tilde{\pi}_a^2}{3\tilde{a}^2} - 3\tilde{a}^2 V(\phi) \right], \\
\dot{\tilde{\phi}} &\approx -\tilde{\lambda}_\phi \tilde{\phi}, \\
\dot{\tilde{\pi}}_\phi &\approx \tilde{\lambda}_\phi \tilde{\pi}_\phi.
\end{aligned}$$

$$\begin{aligned}
H_T^{(-3/2)} &= \int d^3x (\lambda^N \pi_N + \lambda_i \pi^i + \lambda_\phi C_\phi \\
&\quad + N \mathcal{H}^{(-3/2)} + N_i \mathcal{H}^{i(-3/2)})
\end{aligned}$$

$$C_\phi = \pi^{ij} h_{ij} - \phi \pi_\phi$$

$$\begin{aligned}
\mathcal{H}^{(-3/2)} &= \sqrt{h} \left\{ \left[-\phi^3 R + \frac{1}{\phi h} \left(\pi^{ij} \pi_{ij} - \frac{\pi_h^2}{2} \right) \right] \right. \\
&\quad \left. - \frac{3}{2\phi} D_i \phi D^i \phi + 2D^i D_i \phi + U(\phi) \right\}
\end{aligned}$$

$$\mathcal{H}^{i(-3/2)} = -2D_j \pi^{ji} + D^i \phi \pi_\phi.$$

$$N = c_0 \quad N_i = c_i$$

$$\chi_0 \equiv N - c_0 \approx 0, \chi_i \equiv N_i - c_i \approx 0$$

$$\chi_4 \equiv \pi_N \approx 0, \chi_{i+4} \equiv \pi^i \approx 0$$

$$\begin{aligned}
\{N(x) - c_0, \pi_N(x')\} &\approx \delta^{(3)}(x - x') \\
\{N_i(x) - c_i, \pi^j(x')\} &= \delta_i^j \delta^{(3)}(x - x')
\end{aligned}$$

$$\dot{\chi}_0 \approx \{N - c_0, H_T\} \approx 0,$$



$$\dot{\chi}_i \approx \{N_i - c_i, H_T\} \approx 0,$$

$$\begin{aligned}\dot{\chi}_4 &\approx \{\pi_N, H_T\} \approx 0, \\ \dot{\chi}_i + 4 &\approx \{\pi^i, H_T\} \approx 0,\end{aligned}$$

$$\begin{aligned}\dot{h}_{ij} &\approx \{h_{ij}, H_T\}_{DB} \\ &\approx \lambda_\phi h_{ij} + \frac{2c_0}{\phi\sqrt{h}} \left(\pi_{ij} - \frac{\pi_h}{2} h_{ij} \right),\end{aligned}$$

$$\begin{aligned}\dot{\pi}^{ij} &\approx \{\pi^{ij}, H_T\}_{DB} \\ &\approx -\lambda_\phi \pi^{ij} - c_0 \sqrt{h} \phi \left({}^{(3)}R^{ij} - \frac{h^{ij}}{2} {}^{(3)}R \right) \\ &\quad + c_0 \sqrt{h} (D^i D^j - h^{ij} D^k D_k) (\phi) \\ &\quad + \frac{c_0 h^{ij}}{2\phi\sqrt{h}} \left(\pi^{ij} \pi_{ij} - \frac{\pi_h^2}{2} \right) + \frac{c_0}{\phi\sqrt{h}} \pi_h \pi^{ij} \\ &\quad - \sqrt{h} \frac{c_0 \omega}{\phi} D^i \phi D^j \phi - \sqrt{h} \frac{c_0 \omega}{2\phi} h^{ij} D_k \phi D^k \phi \\ &\quad - \frac{\sqrt{h}}{2} h^{ij} c_0 U(\phi) \\ &\quad + c^k D_k (\pi^{ij}) - \frac{2c_0}{\sqrt{h}\phi} \pi^{iq} \pi_q^j \\ \dot{\phi} &\approx \{\phi, H_T\}_{DB} \approx -\lambda_\phi \phi + c^i D_i \phi\end{aligned}$$

$$\begin{aligned}\dot{\pi}_\phi &\approx \{\pi_\phi, H_T\}_{DB} \\ &\approx \lambda_\phi \pi_\phi + c_0 \sqrt{h} {}^{(3)}R + \frac{c_0}{\phi^2 \sqrt{h}} \left(\pi^{ij} \pi_{ij} - \frac{\pi_h^2}{2} \right) \\ &\quad + c_0 \frac{\sqrt{h} \omega}{\phi^2} D_i \phi D^i \phi + 2c_0 \sqrt{h} D_i \left(\frac{\omega}{\phi} D^i \phi \right) \\ &\quad - c_0 \sqrt{h} \frac{dU}{d\phi} + c^i D_i (\pi_\phi)\end{aligned}$$

$$\begin{aligned}\tilde{H}_T^{(-3/2)} &= \int d^3x (\tilde{\lambda}^N \tilde{\pi}_N + \tilde{\lambda}^i \tilde{\pi}_i + \tilde{\lambda}_\phi \tilde{C}_\phi \\ &\quad + \tilde{N} \tilde{\mathcal{H}}^{(-3/2)} + \tilde{N}^i \tilde{\mathcal{H}}_i^{(-3/2)})\end{aligned}$$

$$\tilde{C}_\phi = -\phi \tilde{\pi}_\phi$$

$$\begin{aligned}\tilde{\mathcal{H}}^{(-3/2)} &= \frac{\sqrt{h}}{16\pi G} \left[-{}^3\tilde{R} + \frac{(16\pi G)^2}{h} \left(\tilde{\pi}^{ij} \tilde{\pi}_{ij} - \frac{\tilde{\pi}_h^2}{2} \right) \right] \\ &\quad + \sqrt{h} V(\phi)\end{aligned}$$

$$\tilde{\mathcal{H}}_i^{(-3/2)} = -2\tilde{D}_j \tilde{\pi}_i^j$$

$$\tilde{N} = c_0 (16\pi G \phi)^{\frac{1}{2}}, \tilde{N}_i = c_i (16\pi G \phi),$$

$$\tilde{N} - c_0 (16\pi G \phi)^{\frac{1}{2}} \approx 0, \tilde{N}_i - c_i (16\pi G \phi) \approx 0.$$



$$\left\{\tilde{N}(x) - c_0(16\pi G\phi)^{\frac{1}{2}}, \tilde{\pi}_N(x')\right\} \approx \delta^{(3)}(x - x')$$

$$\left\{\tilde{N}_i(x) - c_i(16\pi G\phi(x)), \tilde{\pi}^j(x')\right\} \approx \delta_i^j \delta^{(3)}(x - x')$$

$$\left\{\tilde{N}(x) - c_0(16\pi G\phi)^{\frac{1}{2}}, \tilde{\mathcal{H}}^{(-3/2)}(x')\right\} \approx 0,$$

$$\left\{\tilde{N}_i(x) - c_i(16\pi G\phi(x)), \tilde{\mathcal{H}}^{(-3/2)}(x')\right\} \approx 0$$

$$\left\{\tilde{N}(x) - c_0(16\pi G\phi)^{\frac{1}{2}}, \tilde{\mathcal{H}}_i^{(-3/2)}(x')\right\} \approx 0,$$

$$\left\{\tilde{N}_i(x) - c_i(16\pi G\phi(x)), \tilde{\mathcal{H}}_i^{(-3/2)}(x')\right\} \approx 0.$$

$$\left\{\tilde{N}(x) - c_0(16\pi G\phi)^{\frac{1}{2}}, \tilde{C}_\phi\right\} \approx \frac{c_0(16\pi G\phi)^{\frac{1}{2}}}{2},$$

$$\left\{\tilde{N}_i(x) - c_i(16\pi G\phi(x)), \tilde{C}_\phi\right\} \approx c_i(16\pi G\phi),$$

$$\tilde{C}'_\phi \equiv -\phi \tilde{\pi}_\phi - \frac{c_0(16\pi G\phi)^{\frac{1}{2}}}{2} \tilde{\pi}_N - c_i(16\pi G\phi) \tilde{\pi}^i$$

$$\begin{aligned} \tilde{H}'^{(-3/2)}_T = \int d^3x & (\tilde{\lambda}^N \tilde{\pi}_N + \tilde{\lambda}^i \tilde{\pi}_i + \tilde{\lambda}_\phi \tilde{C}'_\phi \\ & + \tilde{N} \tilde{\mathcal{H}}^{(-3/2)} + \tilde{N}^i \tilde{\mathcal{H}}_i^{(-3/2)}) \end{aligned}$$

$$\tilde{\chi}_0 \equiv \tilde{N} - c_0(16\pi G\phi)^{\frac{1}{2}}, \tilde{\chi}_i \equiv \tilde{N}_i - c_i(16\pi G\phi)$$

$$\tilde{\chi}_4 \equiv \tilde{\pi}_N, \tilde{\chi}_{i+4} \equiv \tilde{\pi}^i$$

$$\dot{\tilde{\chi}}_0 \approx \left\{\tilde{N}(x) - c_0(16\pi G\phi(x))^{\frac{1}{2}}, \tilde{H}'_T(-3/2)\right\} \approx 0$$

$$\dot{\tilde{\chi}}_i \approx \left\{\tilde{N}_i(x) - c_i(16\pi G\phi(x)), \tilde{H}'_T(-3/2)\right\} \approx 0$$

$$\tilde{C}^{-1}_{\alpha\beta} \equiv \begin{pmatrix} 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & -\mathbb{I} \\ 1 & 0 & 0 & 0 \\ 0 & \mathbb{I} & 0 & 0 \end{pmatrix}$$

$$\begin{aligned} \dot{\tilde{h}}_{ij} & \approx \{\tilde{h}_{ij}, \tilde{H}'_T(-3/2)\}_{DB} \\ & \approx 16\pi G(c_j \tilde{D}_i \phi + c_i \tilde{D}_j \phi) \\ & + \frac{(32\pi G)c_0(16\pi G\phi)^{\frac{1}{2}}}{\sqrt{\tilde{h}}} \left(\tilde{\pi}_{ij} - \frac{\tilde{\pi}_h}{2} \tilde{h}_{ij}\right) \end{aligned}$$

$$\begin{aligned}
\dot{\tilde{\pi}}^{ij} &\approx \{\tilde{\pi}^{ij}, \tilde{H}'_T(-3/2)\}_{DB} \\
&\approx -\frac{c_0(16\pi G\phi)^{\frac{1}{2}}\sqrt{\tilde{h}}}{16\pi G} \left({}^{(3)}\tilde{R}^{ij} - \frac{1}{2}(3)\tilde{R}^{ij} \right) \\
&\quad + \frac{(16\pi G)c_0(16\pi G\phi)^{\frac{1}{2}}}{2\sqrt{\tilde{h}}} \tilde{h}^{ij} \left(\tilde{\pi}^{pk}\tilde{\pi}_{pk} - \frac{\tilde{\pi}_h^2}{2} \right) \\
&\quad - 2\frac{(16\pi G)c_0(16\pi G\phi)^{\frac{1}{2}}}{\sqrt{\tilde{h}}} \left(\tilde{\pi}^{ik}\tilde{\pi}_k^j - \frac{1}{2}\tilde{\pi}_h\tilde{\pi}^{ij} \right) \\
&\quad - \tilde{h}^{ij} \frac{\sqrt{\tilde{h}}c_0(16\pi G\phi)^{\frac{1}{2}}}{2} V(\phi) \\
&\quad + \frac{c_0\sqrt{\tilde{h}}}{2(16\pi G\phi)^{\frac{1}{2}}} \left[-\frac{\tilde{D}^i\phi\tilde{D}^j\phi}{2\phi} + \tilde{D}^i\tilde{D}^j\phi \right. \\
&\quad \left. - \tilde{h}^{ij} \left(-\frac{\tilde{D}^k\phi\tilde{D}_k\phi}{2\phi} + \tilde{D}^k\tilde{D}_k\phi \right) \right] \\
&\quad + (16\pi G)\sqrt{\tilde{h}}c^k\tilde{D}_k \left(\frac{1}{\sqrt{\tilde{h}}} \phi \tilde{\pi}^{ij} \right) \\
&\quad - (16\pi G)\tilde{\pi}^{ki}c^j\tilde{D}_k\phi - (16\pi G)\tilde{\pi}^{kj}c^i\tilde{D}_k\phi \\
\dot{\phi} &\approx \{\phi, \tilde{H}'_T(-3/2)\}_{DB} \approx -\tilde{\lambda}_\phi\phi \\
\dot{\tilde{\pi}}_\phi &\approx \{\tilde{\pi}_\phi, \tilde{H}'_T(-3/2)\}_{DB} \approx \tilde{\lambda}_\phi\tilde{\pi}_\phi
\end{aligned}$$

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